

Pairing and symmetry energy effects in low density nuclear matter: nuclear structure and neutron stars

Joint LIA COLL - AGAIN
COPIGAL and POLITA Workshop



Laboratori Nazionali del Sud - INFN

Catania, April 26 - 29, 2016

Authors: Burrello S.¹, Zheng H.², Colonna M.², Baran V.³

¹ INFN - LNS and Department of Physics and Astronomy, Catania

² INFN - LNS, Catania

³ Department of Physics, Bucharest - Magurele

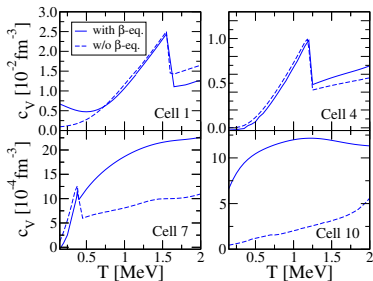


Collective patterns and effective interaction terms

- **Collective phenomena** and **correlations** in nuclear many-body systems
 - Superfluidity in neutron stars: **pairing** effect on heat capacity (cooling)
Collaboration with theoretical group of LPC of Caen, France
 - Dipole excitations in nuclei: **Giant/Pygmy Dipole Resonances (GDR/PDR)**
Collaboration with theoretical group of IPN of Orsay, France
- Isovector term of effective interaction: **symmetry energy** in EoS

Collective patterns and effective interaction terms

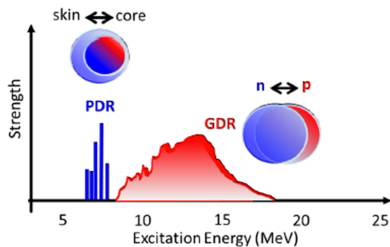
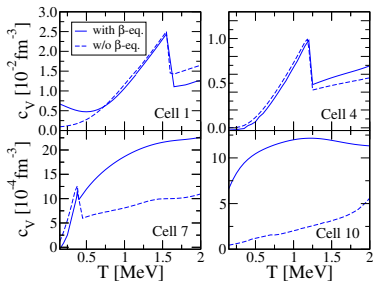
- **Collective phenomena** and **correlations** in nuclear many-body systems
 - **Superfluidity** in neutron stars: **pairing** effect on heat capacity (**cooling**)
Collaboration with theoretical group of LPC of Caen, France
 - **Dipole excitations** in nuclei: **Giant/Pygmy Dipole Resonances (GDR/PDR)**
Collaboration with theoretical group of IPN of Orsay, France
- **Isvector term** of effective interaction: **symmetry energy** in EoS



[Burrello S. et al., PRC 92, 055804, (2015).]

Collective patterns and effective interaction terms

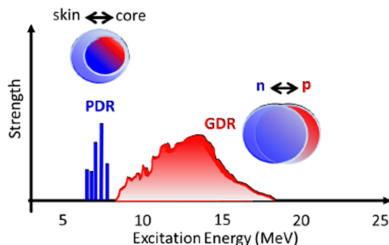
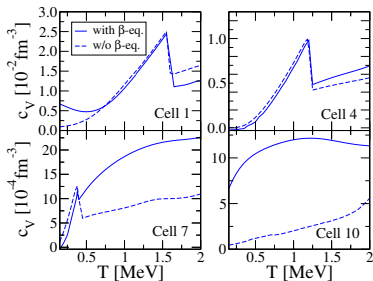
- **Collective phenomena** and **correlations** in nuclear many-body systems
 - **Superfluidity** in neutron stars: **pairing** effect on heat capacity (**cooling**)
Collaboration with theoretical group of LPC of Caen, France
 - **Dipole excitations** in nuclei: **Giant/Pygmy Dipole Resonances (GDR/PDR)**
Collaboration with theoretical group of IPN of Orsay, France
- **Isvector term of effective interaction**: **symmetry energy** in EoS



[Burrello S. et al., PRC 92, 055804, (2015).]

Collective patterns and effective interaction terms

- **Collective phenomena** and **correlations** in nuclear many-body systems
 - **Superfluidity** in neutron stars: **pairing** effect on heat capacity (**cooling**)
Collaboration with theoretical group of LPC of Caen, France
 - **Dipole excitations** in nuclei: **Giant/Pygmy Dipole Resonances (GDR/PDR)**
Collaboration with theoretical group of IPN of Orsay, France
- **Isvector** term of **effective interaction**: **symmetry energy** in EoS



[Burrello S. et al., PRC 92, 055804, (2015).]

Dynamics in many-body systems: the BNV model

- **Small amplitude dynamics** of nuclei

- Quantal approaches: **HF+RPA**, **TDHF**, ...
- Semi-classical approaches: **BNV** model [see V. Baran et al., PRC88, (2013)]

- **Transport equation** for the 1-body distributions $f_q(\mathbf{r}, \mathbf{p}, t) \Rightarrow$ **Vlasov** equations:

$$\frac{\partial f_q}{\partial t} + \frac{\partial \epsilon_q}{\partial \mathbf{p}} \frac{\partial f_q}{\partial \mathbf{r}} - \frac{\partial \epsilon_q}{\partial \mathbf{r}} \frac{\partial f_q}{\partial \mathbf{p}} = 0 \quad q = p, n$$

- **Mean-field with Skyrme interactions: SAMi-J** [X. Roca-Maza et al., PRC87, (2013)]

$$\mathcal{E} = \frac{\hbar^2}{2m} \tau + C_0 \rho^2 + D_0 \rho_3^2 + C_3 \rho^{\sigma+2} + D_3 \rho^\sigma \rho_3^2 + C_{\text{eff}} \rho \tau + D_{\text{eff}} \rho_3 \tau_3 + C_\nabla (\nabla \rho)^2 + D_\nabla (\nabla \rho_3)^2$$

$$\rho = \rho_n + \rho_p \quad \rho_3 = \rho_n - \rho_p$$

$$\tau = \tau_n + \tau_p \quad \tau_3 = \tau_n - \tau_p$$

- **Test-particle method** (**finite width** wave packets)

\Rightarrow no explicitly surface terms ($C_\nabla = D_\nabla = 0$) to reproduce experimental values ($\sqrt{\langle r_p^2 \rangle}$, B/A)

- **Semi-classical model** \Rightarrow **no shell effects**

Dynamics in many-body systems: the BNV model

- **Small amplitude dynamics** of nuclei
 - **Quantal** approaches: **HF+RPA**, **TDHF**, ...
 - **Semi-classical** approaches: **BNV** model [see V. Baran et al., PRC88, (2013)]
- **Transport equation** for the 1-body distributions $f_q(\mathbf{r}, \mathbf{p}, t) \Rightarrow$ **Vlasov** equations:

$$\frac{\partial f_q}{\partial t} + \frac{\partial \epsilon_q}{\partial \mathbf{p}} \frac{\partial f_q}{\partial \mathbf{r}} - \frac{\partial \epsilon_q}{\partial \mathbf{r}} \frac{\partial f_q}{\partial \mathbf{p}} = 0 \quad \Rightarrow \quad \rho_q(\mathbf{r}, t) = \frac{2}{(2\pi\hbar)^3} \int d\mathbf{p} f_q(\mathbf{r}, \mathbf{p}, t) \quad q = p, n$$

- **Mean-field with Skyrme interactions**: **SAMi-J** [X. Roca-Maza et al., PRC87, (2013)]

$$\mathcal{E} = \frac{\hbar^2}{2m} \tau + C_0 \rho^2 + D_0 \rho_3^2 + C_3 \rho^{\sigma+2} + D_3 \rho^\sigma \rho_3^2 + C_{\text{eff}} \rho \tau + D_{\text{eff}} \rho_3 \tau_3 + C_\nabla (\nabla \rho)^2 + D_\nabla (\nabla \rho_3)^2$$

$$\rho = \rho_n + \rho_p \quad \rho_3 = \rho_n - \rho_p$$

$$\tau = \tau_n + \tau_p \quad \tau_3 = \tau_n - \tau_p$$

- **Test-particle method** (**finite width** wave packets)
 \Rightarrow no explicitly surface terms ($C_\nabla = D_\nabla = 0$) to reproduce experimental values ($\sqrt{\langle r_p^2 \rangle}$, B/A)
- **Semi-classical model** \Rightarrow **no shell effects**

Dynamics in many-body systems: the BNV model

- **Small amplitude dynamics** of nuclei
 - **Quantal** approaches: **HF+RPA**, **TDHF**, ...
 - **Semi-classical** approaches: **BNV** model [see V. Baran et al., PRC88, (2013)]
- **Transport equation** for the 1-body distributions $f_q(\mathbf{r}, \mathbf{p}, t) \Rightarrow$ **Vlasov** equations:

$$\frac{\partial f_q}{\partial t} + \frac{\partial \epsilon_q}{\partial \mathbf{p}} \frac{\partial f_q}{\partial \mathbf{r}} - \frac{\partial \epsilon_q}{\partial \mathbf{r}} \frac{\partial f_q}{\partial \mathbf{p}} = 0 \quad \Rightarrow \quad \rho_q(\mathbf{r}, t) = \frac{2}{(2\pi\hbar)^3} \int d\mathbf{p} f_q(\mathbf{r}, \mathbf{p}, t) \quad q = p, n$$

- **Mean-field** with Skyrme interactions: **SAMi-J** [X. Roca-Maza et al., PRC87, (2013)]

$$\mathcal{E} = \frac{\hbar^2}{2m} \tau + C_0 \rho^2 + D_0 \rho_3^2 + C_3 \rho^{\sigma+2} + D_3 \rho^\sigma \rho_3^2 + C_{\text{eff}} \rho \tau + D_{\text{eff}} \rho_3 \tau_3 + C_\nabla (\nabla \rho)^2 + D_\nabla (\nabla \rho_3)^2$$

$$\rho = \rho_n + \rho_p \quad \rho_3 = \rho_n - \rho_p$$

$$\tau = \tau_n + \tau_p \quad \tau_3 = \tau_n - \tau_p$$

- **Test-particle method** (**finite width** wave packets)
 \Rightarrow no explicitly surface terms ($C_\nabla = D_\nabla = 0$) to reproduce experimental values ($\sqrt{\langle r_p^2 \rangle}$, B/A)
- **Semi-classical model** \Rightarrow **no shell effects**

Dynamics in many-body systems: the BNV model

- **Small amplitude dynamics** of nuclei
 - **Quantal** approaches: **HF+RPA**, **TDHF**, ...
 - **Semi-classical** approaches: **BNV** model [see V. Baran et al., PRC88, (2013)]
- **Transport equation** for the 1-body distributions $f_q(\mathbf{r}, \mathbf{p}, t) \Rightarrow$ **Vlasov** equations:

$$\frac{\partial f_q}{\partial t} + \frac{\partial \epsilon_q}{\partial \mathbf{p}} \frac{\partial f_q}{\partial \mathbf{r}} - \frac{\partial \epsilon_q}{\partial \mathbf{r}} \frac{\partial f_q}{\partial \mathbf{p}} = 0 \quad \Rightarrow \quad \rho_q(\mathbf{r}, t) = \frac{2}{(2\pi\hbar)^3} \int d\mathbf{p} f_q(\mathbf{r}, \mathbf{p}, t) \quad q = p, n$$

- **Mean-field** with Skyrme interactions: **SAMi-J** [X. Roca-Maza et al., PRC87, (2013)]

$$\mathcal{E} = \frac{\hbar^2}{2m} \tau + C_0 \rho^2 + D_0 \rho_3^2 + C_3 \rho^{\sigma+2} + D_3 \rho^\sigma \rho_3^2 + C_{\text{eff}} \rho \tau + D_{\text{eff}} \rho_3 \tau_3 + C_\nabla (\nabla \rho)^2 + D_\nabla (\nabla \rho_3)^2$$

$$\rho = \rho_n + \rho_p \quad \rho_3 = \rho_n - \rho_p$$

$$\tau = \tau_n + \tau_p \quad \tau_3 = \tau_n - \tau_p$$

- **Test-particle** method (**finite width** wave packets)
 - \Rightarrow no explicitly surface terms ($C_\nabla = D_\nabla = 0$) to reproduce experimental values ($\sqrt{\langle r_p^2 \rangle}$, B/A)
- **Semi-classical** model \Rightarrow **no shell effects**

Dynamics in many-body systems: the BNV model

- **Small amplitude dynamics** of nuclei
 - **Quantal** approaches: **HF+RPA**, **TDHF**, ...
 - **Semi-classical** approaches: **BNV** model [see V. Baran et al., PRC88, (2013)]
- **Transport equation** for the 1-body distributions $f_q(\mathbf{r}, \mathbf{p}, t) \Rightarrow$ **Vlasov** equations:

$$\frac{\partial f_q}{\partial t} + \frac{\partial \epsilon_q}{\partial \mathbf{p}} \frac{\partial f_q}{\partial \mathbf{r}} - \frac{\partial \epsilon_q}{\partial \mathbf{r}} \frac{\partial f_q}{\partial \mathbf{p}} = 0 \quad \Rightarrow \quad \rho_q(\mathbf{r}, t) = \frac{2}{(2\pi\hbar)^3} \int d\mathbf{p} f_q(\mathbf{r}, \mathbf{p}, t) \quad q = p, n$$

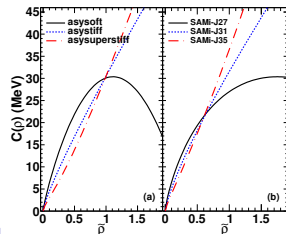
- **Mean-field** with **Skyrme** interactions: **SAMi-J** [X. Roca-Maza et al., PRC87, (2013)]

$$\mathcal{E} = \frac{\hbar^2}{2m} \tau + C_0 \rho^2 + D_0 \rho_3^2 + C_3 \rho^{\sigma+2} + D_3 \rho^\sigma \rho_3^2 + C_{\text{eff}} \rho \tau + D_{\text{eff}} \rho_3 \tau_3 + C_\nabla (\nabla \rho)^2 + D_\nabla (\nabla \rho_3)^2$$

$$\rho = \rho_n + \rho_p \quad \rho_3 = \rho_n - \rho_p$$

$$\tau = \tau_n + \tau_p \quad \tau_3 = \tau_n - \tau_p$$

- **Test-particle** method (**finite width** wave packets)
 \Rightarrow no explicitly surface terms ($C_\nabla = D_\nabla = 0$) to reproduce experimental values ($\sqrt{\langle r_p^2 \rangle}$, B/A)
- **Semi-classical** model \Rightarrow **no shell effects**



Dynamics in many-body systems: the BNV model

- **Small amplitude dynamics** of nuclei
 - **Quantal** approaches: **HF+RPA**, **TDHF**, ...
 - **Semi-classical** approaches: **BNV** model [see V. Baran et al., PRC88, (2013)]
- **Transport equation** for the 1-body distributions $f_q(\mathbf{r}, \mathbf{p}, t) \Rightarrow$ **Vlasov** equations:

$$\frac{\partial f_q}{\partial t} + \frac{\partial \epsilon_q}{\partial \mathbf{p}} \frac{\partial f_q}{\partial \mathbf{r}} - \frac{\partial \epsilon_q}{\partial \mathbf{r}} \frac{\partial f_q}{\partial \mathbf{p}} = 0 \quad \Rightarrow \quad \rho_q(\mathbf{r}, t) = \frac{2}{(2\pi\hbar)^3} \int d\mathbf{p} f_q(\mathbf{r}, \mathbf{p}, t) \quad q = p, n$$

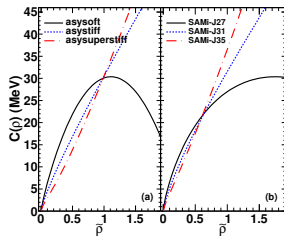
- **Mean-field** with **Skyrme** interactions: **SAMi-J** [X. Roca-Maza et al., PRC87, (2013)]

$$\mathcal{E} = \frac{\hbar^2}{2m} \tau + C_0 \rho^2 + D_0 \rho_3^2 + C_3 \rho^{\sigma+2} + D_3 \rho^\sigma \rho_3^2 + C_{\text{eff}} \rho \tau + D_{\text{eff}} \rho_3 \tau_3 + C_\nabla (\nabla \rho)^2 + D_\nabla (\nabla \rho_3)^2$$

$$\rho = \rho_n + \rho_p \quad \rho_3 = \rho_n - \rho_p$$

$$\tau = \tau_n + \tau_p \quad \tau_3 = \tau_n - \tau_p$$

- **Test-particle** method (**finite width** wave packets)
 - \Rightarrow no **explicitly** surface terms ($C_\nabla = D_\nabla = 0$) to reproduce **experimental** values ($\sqrt{\langle r_p^2 \rangle}$, B/A)
- **Semi-classical** model \Rightarrow **no shell effects**



Dynamics in many-body systems: the BNV model

- **Small amplitude dynamics** of nuclei
 - **Quantal** approaches: **HF+RPA**, **TDHF**, ...
 - **Semi-classical** approaches: **BNV** model [see V. Baran et al., PRC88, (2013)]
- **Transport equation** for the 1-body distributions $f_q(\mathbf{r}, \mathbf{p}, t) \Rightarrow$ **Vlasov** equations:

$$\frac{\partial f_q}{\partial t} + \frac{\partial \epsilon_q}{\partial \mathbf{p}} \frac{\partial f_q}{\partial \mathbf{r}} - \frac{\partial \epsilon_q}{\partial \mathbf{r}} \frac{\partial f_q}{\partial \mathbf{p}} = 0 \quad \Rightarrow \quad \rho_q(\mathbf{r}, t) = \frac{2}{(2\pi\hbar)^3} \int d\mathbf{p} f_q(\mathbf{r}, \mathbf{p}, t) \quad q = p, n$$

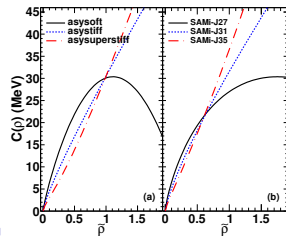
- **Mean-field** with **Skyrme** interactions: **SAMi-J** [X. Roca-Maza et al., PRC87, (2013)]

$$\mathcal{E} = \frac{\hbar^2}{2m} \tau + C_0 \rho^2 + D_0 \rho_3^2 + C_3 \rho^{\sigma+2} + D_3 \rho^\sigma \rho_3^2 + C_{\text{eff}} \rho \tau + D_{\text{eff}} \rho_3 \tau_3 + C_\nabla (\nabla \rho)^2 + D_\nabla (\nabla \rho_3)^2$$

$$\rho = \rho_n + \rho_p \quad \rho_3 = \rho_n - \rho_p$$

$$\tau = \tau_n + \tau_p \quad \tau_3 = \tau_n - \tau_p$$

- **Test-particle** method (**finite width** wave packets)
 - \Rightarrow no **explicitly** surface terms ($C_\nabla = D_\nabla = 0$) to reproduce **experimental** values ($\sqrt{\langle r_p^2 \rangle}$, B/A)
- **Semi-classical** model \Rightarrow **no shell effects**



Dipole oscillations and response functions

- Instantaneous **ground-state perturbation**:

$$\hat{V}_K^{\text{ext}}(\mathbf{r}, t) = \eta_K \delta(t - t_0) \hat{D}_K(\mathbf{r}) \quad K = S, V$$

$$\Rightarrow |\Phi_0\rangle \rightarrow |\Phi_K(t_0)\rangle = e^{i\eta_K \hat{D}_K} |\Phi_0\rangle$$

- Isoscalar (IS) or isovector (IV) dipole operator:

$$\hat{D}_S = \sum_i \left(r_i^2 - \frac{5}{3} \langle r^2 \rangle \right) z_i, \quad \hat{D}_V = \sum_i \tau_i \frac{N}{A} z_i - (1 - \tau_i) \frac{Z}{A} z_i, \quad \tau_i = 0(1) \text{ for } n(p)$$

- Dynamical evolution of the excitation: $D_K(t) = \langle \Phi_K(t) | \hat{D}_K | \Phi_K(t) \rangle$
- Strength function: $S_K(E) = \sum_n | \langle n | \hat{D}_K | 0 \rangle |^2 \delta(E - (E_n - E_0))$

$$S_K(\omega) = \frac{\text{Im } D_K(\omega)}{\pi \eta_K} \quad D_K(\omega) \text{ Fourier Transform of } D_K(t)$$

Dipole oscillations and response functions

- Instantaneous **ground-state perturbation**:

$$\hat{V}_K^{\text{ext}}(\mathbf{r}, t) = \eta_K \delta(t - t_0) \hat{D}_K(\mathbf{r}) \quad K = S, V$$

$$\Rightarrow |\Phi_0\rangle \rightarrow |\Phi_K(t_0)\rangle = e^{i\eta_K \hat{D}_K} |\Phi_0\rangle$$

- Isoscalar (IS)** or **isovector (IV)** dipole operator:

$$\hat{D}_S = \sum_i \left(r_i^2 - \frac{5}{3} \langle r^2 \rangle \right) z_i, \quad \hat{D}_V = \sum_i \tau_i \frac{N}{A} z_i - (1 - \tau_i) \frac{Z}{A} z_i, \quad \tau_i = 0(1) \text{ for } n(p)$$

- Dynamical evolution** of the excitation: $D_K(t) = \langle \Phi_K(t) | \hat{D}_K | \Phi_K(t) \rangle$
- Strength function**: $S_K(E) = \sum_n |\langle n | \hat{D}_K | 0 \rangle|^2 \delta(E - (E_n - E_0))$

$$S_K(\omega) = \frac{\text{Im } D_k(\omega)}{\pi \eta_k} \quad D_k(\omega) \text{ Fourier Transform of } D_k(t)$$

Dipole oscillations and response functions

- Instantaneous **ground-state perturbation**:

$$\hat{V}_K^{\text{ext}}(\mathbf{r}, t) = \eta_K \delta(t - t_0) \hat{D}_K(\mathbf{r}) \quad K = S, V$$

$$\Rightarrow |\Phi_0\rangle \rightarrow |\Phi_K(t_0)\rangle = e^{i\eta_K \hat{D}_K} |\Phi_0\rangle$$

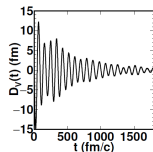
- Isoscalar (IS)** or **isovector (IV)** dipole operator:

$$\hat{D}_S = \sum_i \left(r_i^2 - \frac{5}{3} \langle r^2 \rangle \right) z_i, \quad \hat{D}_V = \sum_i \tau_i \frac{N}{A} z_i - (1 - \tau_i) \frac{Z}{A} z_i, \quad \tau_i = 0(1) \text{ for } n(p)$$

- Dynamical evolution** of the excitation: $D_K(t) = \langle \Phi_K(t) | \hat{D}_K | \Phi_K(t) \rangle$

- Strength function**: $S_K(E) = \sum_n |\langle n | \hat{D}_K | 0 \rangle|^2 \delta(E - (E_n - E_0))$

$$S_K(\omega) = \frac{\text{Im } D_k(\omega)}{\pi \eta_k} \quad D_k(\omega) \text{ Fourier Transform of } D_k(t)$$



Dipole oscillations and response functions

- Instantaneous **ground-state perturbation**:

$$\hat{V}_K^{\text{ext}}(\mathbf{r}, t) = \eta_K \delta(t - t_0) \hat{D}_K(\mathbf{r}) \quad K = S, V$$

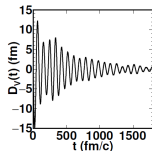
$$\Rightarrow |\Phi_0\rangle \rightarrow |\Phi_K(t_0)\rangle = e^{i\eta_K \hat{D}_K} |\Phi_0\rangle$$

- Isoscalar (IS)** or **isovector (IV)** dipole operator:

$$\hat{D}_S = \sum_i \left(r_i^2 - \frac{5}{3} \langle r^2 \rangle \right) z_i, \quad \hat{D}_V = \sum_i \tau_i \frac{N}{A} z_i - (1 - \tau_i) \frac{Z}{A} z_i, \quad \tau_i = 0(1) \text{ for } n(p)$$

- Dynamical evolution** of the excitation: $D_K(t) = \langle \Phi_K(t) | \hat{D}_K | \Phi_K(t) \rangle$
- Strength function**: $S_K(E) = \sum_n |\langle n | \hat{D}_K | 0 \rangle|^2 \delta(E - (E_n - E_0))$

$$S_K(\omega) = \frac{\text{Im } D_k(\omega)}{\pi \eta_k} \quad D_k(\omega) \text{ Fourier Transform of } D_k(t)$$



Results (1) : coupling between IS and IV modes

- **Symmetric** nuclear matter: **IS** and **IV** modes are **decoupled**
- Neutron-rich systems: n and p oscillate with different amplitudes \Rightarrow coupling

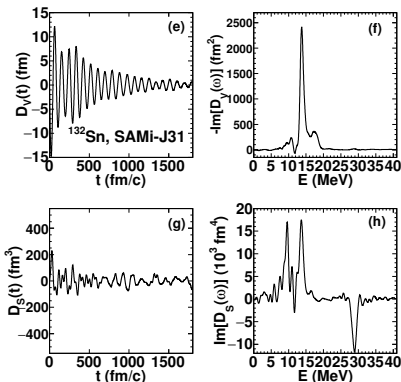
Results (1) : coupling between IS and IV modes

- **Symmetric** nuclear matter: **IS** and **IV** modes are **decoupled**
- **Neutron-rich** systems: n and p oscillate with **different amplitudes** \Rightarrow **coupling**

Results (1) : coupling between IS and IV modes

- **Symmetric** nuclear matter: **IS** and **IV** modes are **decoupled**
- **Neutron-rich** systems: n and p oscillate with **different amplitudes** \Rightarrow **coupling**

[H. Zheng et al., arXiv:1603.03594, (2016)]



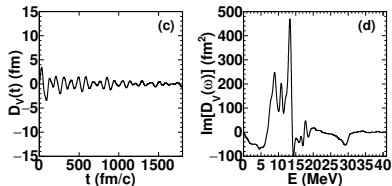
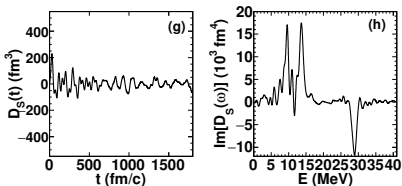
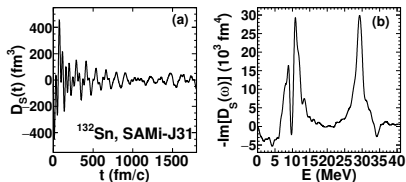
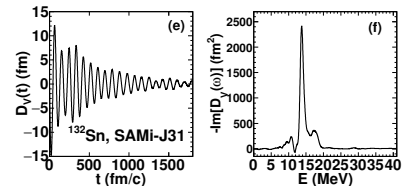
- **IV** response:
 - main **IV GDR** peak
 - smaller IV peak at higher E
 - some strength at lower E (**PDR**)
- **IS** response:
 - larger PDR strength \Rightarrow **isoscalar-like**
 - peak at higher E (**IS GDR**)

IV perturbation \rightarrow **IV + IS** response

Results (1) : coupling between IS and IV modes

- **Symmetric** nuclear matter: **IS** and **IV** modes are **decoupled**
- **Neutron-rich** systems: n and p oscillate with **different amplitudes** \Rightarrow **coupling**

[H. Zheng et al., arXiv:1603.03594, (2016)]



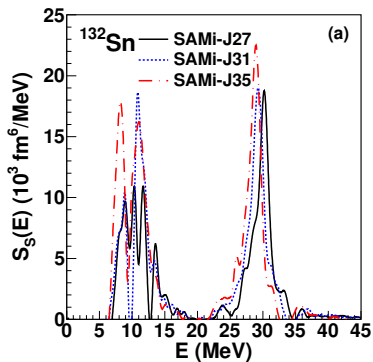
IV perturbation \rightarrow **IV + IS response**

IS perturbation \rightarrow **IS + IV response**

Results (2): influence of IV term of the interaction

- **Three regions of A:** ^{68}Ni ($N/Z = 1.43$), ^{132}Sn ($N/Z = 1.64$), ^{208}Pb ($N/Z = 1.54$)

[H. Zheng et al., arXiv:1603.03594, (2016)]

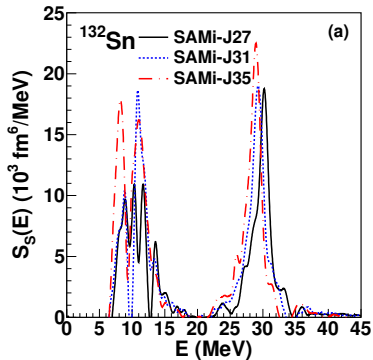


- **Two peaks** in PDR region
[see M. Urban, PRC85, (2012)]
- Semi-classical treatment for **surface**
⇒ **discrepancy** with RPA
- Larger $L \Rightarrow$ higher IV PDR peak
(**mixing**)
- Larger $L \Rightarrow$ larger **neutron-skin**

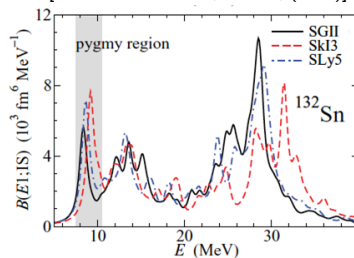
Results (2): influence of IV term of the interaction

- **Three regions of A:** ^{68}Ni ($N/Z = 1.43$), ^{132}Sn ($N/Z = 1.64$), ^{208}Pb ($N/Z = 1.54$)

[H. Zheng et al., arXiv:1603.03594, (2016)]



[X. Roca-Maza et al., PRC85, (2012)]



- **Two peaks** in PDR region

[see M. Urban, PRC85, (2012)]

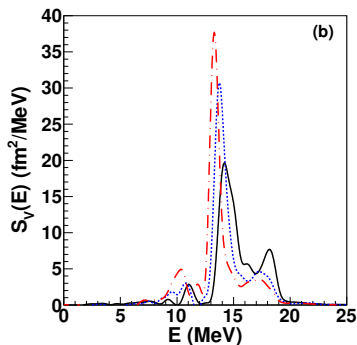
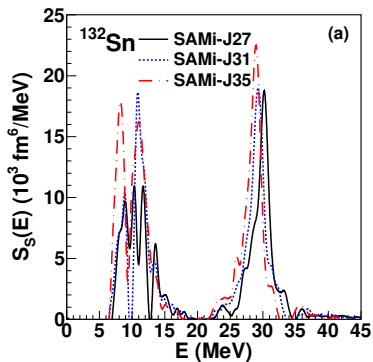
- Semi-classical treatment for **surface**
⇒ **discrepancy** with **RPA**

- Larger L ⇒ higher IV PDR peak
(**mixing**)
- Larger L ⇒ larger **neutron-skin**

Results (2): influence of IV term of the interaction

- **Three regions of A:** ^{68}Ni ($N/Z = 1.43$), ^{132}Sn ($N/Z = 1.64$), ^{208}Pb ($N/Z = 1.54$)

[H. Zheng et al., arXiv:1603.03594, (2016)]



- **Two peaks** in PDR region

[see M. Urban, PRC85, (2012)]

- Semi-classical treatment for **surface**
⇒ **discrepancy** with **RPA**

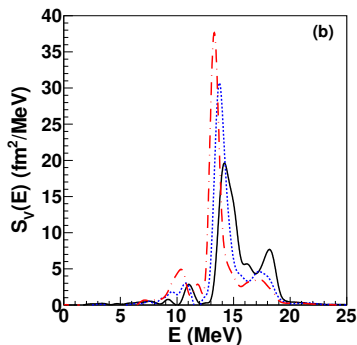
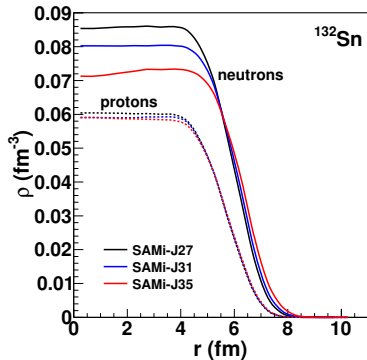
- Larger **L** ⇒ higher IV PDR peak
(**mixing**)

- Larger **L** ⇒ larger **neutron-skin**

Results (2): influence of IV term of the interaction

- **Three regions of A:** ^{68}Ni ($N/Z = 1.43$), ^{132}Sn ($N/Z = 1.64$), ^{208}Pb ($N/Z = 1.54$)

[H. Zheng et al., arXiv:1603.03594, (2016)]



- **Two peaks** in PDR region

[see M. Urban, PRC85, (2012)]

- Semi-classical treatment for **surface**
⇒ **discrepancy** with **RPA**

- Larger **L** ⇒ higher IV PDR peak
(**mixing**)
- Larger **L** ⇒ larger **neutron-skin**

Results (3): transition densities analysis

- **Transition densities:** information on **spatial structure** of excitation dynamics

$$\delta\rho_q(r, E) \propto \int_{t_0}^{\infty} dt \delta\rho_q(r, t) \sin \frac{Et}{\hbar} \quad E \equiv \text{energy with a peak in the strength}$$

- Oscillations: in phase \Rightarrow **isoscalar-like** vs. out of phase \Rightarrow **isovector-like**
- **Simultaneous** agitation of all modes \Rightarrow interference

Results (3): transition densities analysis

- **Transition densities:** information on **spatial structure** of excitation dynamics

$$\delta\rho_q(r, E) \propto \int_{t_0}^{\infty} dt \delta\rho_q(r, t) \sin \frac{Et}{\hbar} \quad E \equiv \text{energy with a peak in the strength}$$

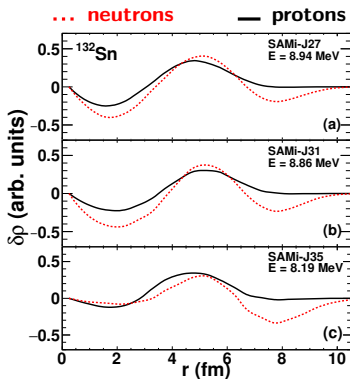
- Oscillations: **in phase** \Rightarrow **isoscalar-like** vs. **out of phase** \Rightarrow **isovector-like**
- **Simultaneous** agitation of all modes \Rightarrow **interference**

Results (3): transition densities analysis

- **Transition densities:** information on **spatial structure** of excitation dynamics

$$\delta\rho_q(r, E) \propto \int_{t_0}^{\infty} dt \delta\rho_q(r, t) \sin \frac{Et}{\hbar} \quad E \equiv \text{energy with a peak in the strength}$$

- Oscillations: **in phase** \Rightarrow **isoscalar-like** vs. **out of phase** \Rightarrow **isovector-like**
- **Simultaneous** agitation of all modes \Rightarrow **interference**



PDR peak in IS response

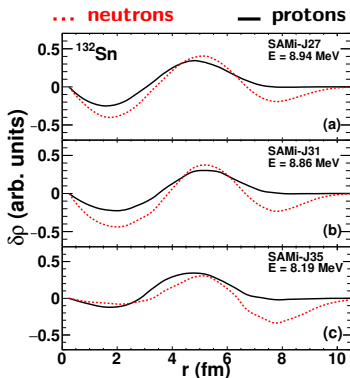
- Larger oscillation for neutrons
- Amplitude in the **surface increases** while increasing L
 - Thicker neutron skin
 - More symmetric central region

Results (3): transition densities analysis

- **Transition densities:** information on **spatial structure** of excitation dynamics

$$\delta\rho_q(r, E) \propto \int_{t_0}^{\infty} dt \delta\rho_q(r, t) \sin \frac{Et}{\hbar} \quad E \equiv \text{energy with a peak in the strength}$$

- Oscillations: **in phase** \Rightarrow **isoscalar-like** vs. **out of phase** \Rightarrow **isovector-like**
- **Simultaneous** agitation of all modes \Rightarrow **interference**



PDR peak in IS response

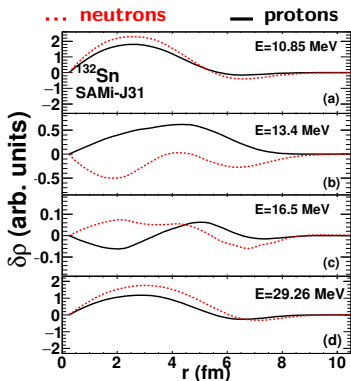
- Larger oscillation for **neutrons**
- Amplitude in the **surface increases** while increasing **L**
 - **Thicker neutron-skin**
 - More **symmetric central** region

Results (3): transition densities analysis

- **Transition densities:** information on **spatial structure** of excitation dynamics

$$\delta\rho_q(r, E) \propto \int_{t_0}^{\infty} dt \delta\rho_q(r, t) \sin \frac{Et}{\hbar} \quad E \equiv \text{energy with a peak in the strength}$$

- Oscillations: **in phase** \Rightarrow **isoscalar-like** vs. **out of phase** \Rightarrow **isovector-like**
- **Simultaneous** agitation of all modes \Rightarrow **interference**

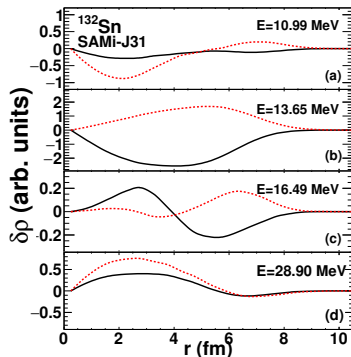


IS - like

IV GDR

IV - like

IS GDR

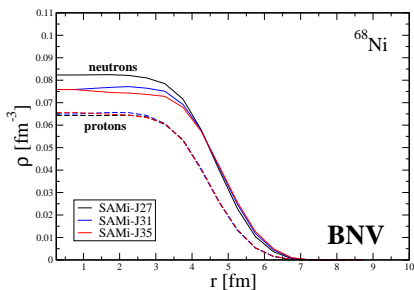
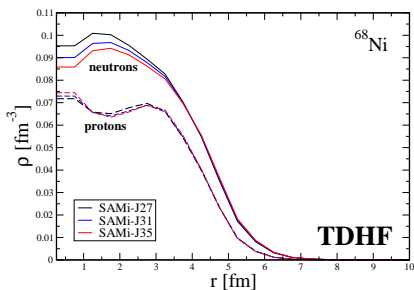


Further developments: comparison with TDHF

- **TDHF** [D. Lacroix, G. Scamps] calculations: **PRELIMINARY RESULTS**
- Difference in the **initial density profile**: **sharper** in TDHF case
- **Overestimation** of PDR energy and **merging** of IS-L peak with GDR
- TDHF + BCS: **negligible pairing effects** for ^{68}Ni (closed-shell) \Rightarrow CHECK!

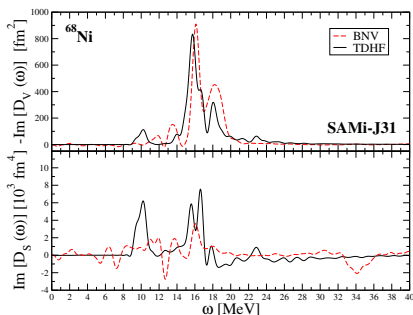
Further developments: comparison with TDHF

- **TDHF** [D. Lacroix, G. Scamps] calculations: **PRELIMINARY RESULTS**
- Difference in the **initial density profile**: **sharper** in TDHF case
- Overestimation of PDR energy and **merging** of IS-L peak with GDR
- **TDHF + BCS**: **negligible pairing effects** for ^{68}Ni (closed-shell) \Rightarrow CHECK!



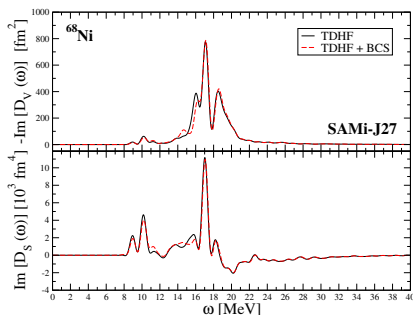
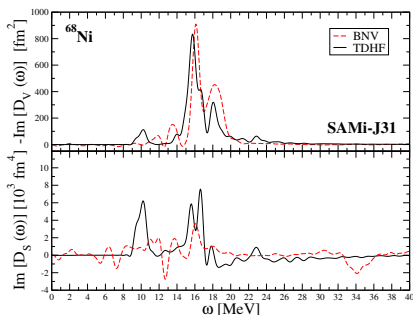
Further developments: comparison with TDHF

- **TDHF** [D. Lacroix, G. Scamps] calculations: **PRELIMINARY RESULTS**
- Difference in the **initial density profile**: **sharper** in TDHF case
- **Overestimation** of PDR energy and **merging** of IS-L peak with GDR
- TDHF + BCS: **negligible pairing effects** for ^{68}Ni (closed-shell) \Rightarrow CHECK!



Further developments: comparison with TDHF

- **TDHF** [D. Lacroix, G. Scamps] calculations: **PRELIMINARY RESULTS**
- Difference in the **initial density profile**: **sharper** in TDHF case
- **Overestimation** of PDR energy and **merging** of IS-L peak with GDR
- **TDHF + BCS**: **negligible pairing effects** for ^{68}Ni (closed-shell) \Rightarrow **CHECK!**



Final remarks and conclusions

Summary

- Small amplitude **dynamics** in nuclei: **semi-classical BNV** transport model
- Examination of **both IS and IV** nuclear E1 **response** in **neutron-rich** systems

Final remarks and conclusions

Summary

- Small amplitude **dynamics** in nuclei: **semi-classical BNV** transport model
- Examination of **both IS and IV** nuclear E1 **response** in **neutron-rich** systems

Main results

- Isoscalar/isovector **mixing** of dipole excitations
- Characterization of the **nature** of low-lying energy IV response (**PDR**)
- Essential role of **L** in shaping **neutron-skin** and E1 response in PDR region
- General good **agreement** with previous semi-classical and **RPA** studies

Final remarks and conclusions

Summary

- Small amplitude **dynamics** in nuclei: **semi-classical BNV** transport model
- Examination of **both IS and IV** nuclear E1 **response** in **neutron-rich** systems

Main results

- Isoscalar/isovector **mixing** of dipole excitations
- Characterization of the **nature** of low-lying energy IV response (**PDR**)
- Essential role of **L** in shaping **neutron-skin** and E1 response in PDR region
- General good **agreement** with previous semi-classical and **RPA** studies

Further developments and outlooks

- Tuning of finite size terms to improve the **treatment of surface** effects
- Comparison with **TDHF calculations** (coll. with D. Lacroix group, Orsay)
- Study some **isotopes chains** to understand **isovector** and **pairing** terms role

French collaborations of our group

Nuclear structure (TDHF)

D. Lacroix (IPN, IN2P3-CNRS, Orsay, France),

G. Scamps (Department of Physics, Tohoku University, Japan)

Neutron stars modelization

F. Gulminelli, F. Aymard (LPC, CNRS and ENSICAEN, Caen, France)

A. Raduta (NIPNE, Bucharest-Magurele, Romania)

Multifragmentation reactions

P. Napolitani (IPN, IN2P3-CNRS, Orsay, France),

M. Di Prima (LNS - INFN, Catania, Italy)

French collaborations of our group

Nuclear structure (TDHF)

D. Lacroix (IPN, IN2P3-CNRS, Orsay, France),

G. Scamps (Department of Physics, Tohoku University, Japan)

Neutron stars modelization

F. Gulminelli, F. Aymard (LPC, CNRS and ENSICAEN, Caen, France)

A. Raduta (NIPNE, Bucharest-Magurele, Romania)

Multifragmentation reactions

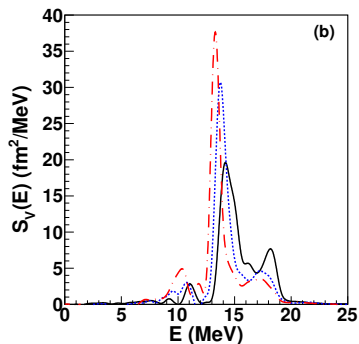
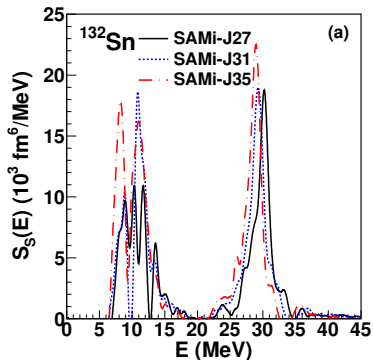
P. Napolitani (IPN, IN2P3-CNRS, Orsay, France),

M. Di Prima (LNS - INFN, Catania, Italy)

THANK YOU FOR YOUR KIND ATTENTION!

Further insight: influence of IV term of the interaction

- Three regions of A: ^{68}Ni ($N/Z = 1.43$), ^{132}Sn ($N/Z = 1.64$), ^{208}Pb ($N/Z = 1.54$)



- Two peaks in PDR region

see M. Urban, PRC85, (2012)

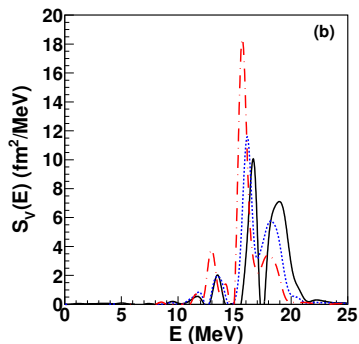
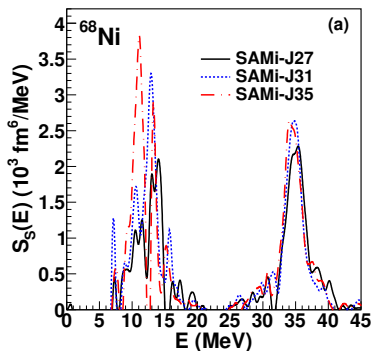
- Semi-classical treatment for **surface**

- Larger $L \Rightarrow$ higher IV PDR peak
(**mixing**)

- Larger $L \Rightarrow$ larger **neutron-skin**

Further insight: influence of IV term of the interaction

- Three regions of A: ^{68}Ni ($N/Z = 1.43$), ^{132}Sn ($N/Z = 1.64$), ^{208}Pb ($N/Z = 1.54$)



- Two peaks in PDR region

see M. Urban, PRC85, (2012)

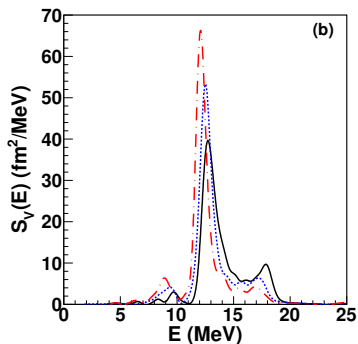
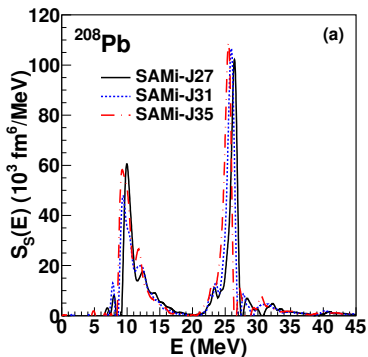
- Semi-classical treatment for surface

- Larger $L \Rightarrow$ higher IV PDR peak
(mixing)

- Larger $L \Rightarrow$ larger neutron-skin

Further insight: influence of IV term of the interaction

- Three regions of A: ^{68}Ni ($N/Z = 1.43$), ^{132}Sn ($N/Z = 1.64$), ^{208}Pb ($N/Z = 1.54$)



- Two peaks in PDR region

see M. Urban, PRC85, (2012)

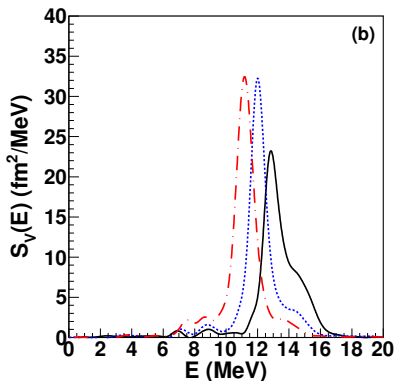
- Semi-classical treatment for **surface**

- Larger $L \Rightarrow$ higher IV PDR peak
(**mixing**)

- Larger $L \Rightarrow$ larger **neutron-skin**

Further insight: influence of momentum dependence

Momentum independent interaction



Momentum dependent interaction

