

Exclusive $b \rightarrow s \ell^+ \ell^-$

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Outline

- 1 $B \rightarrow K^* \ell^+ \ell^-$
 - amplitudes and rates
 - factorization
 - numerical predictions
 - warnings
 - brief comment on new physics
- 2 other modes (briefly)
- 3 discussion
- 4 backup

$B \rightarrow K^* \ell^+ \ell^-$: Amplitudes and rates

doubly differential decay rate:

e.g. [Lee/Ligeti/Stewart/Tackmann 06]

$$\frac{d^2\Gamma}{dq^2 d\cos\theta} = \frac{3}{8} \left[(1 + \cos^2\theta) H_T(q^2) + 2 \cos\theta H_A(q^2) + 2(1 - \cos^2\theta) H_L(q^2) \right].$$

- $H_T(q^2)$ determines rate for transversely polarized K^*
- $H_L(q^2)$ determines rate for longitudinally polarized K^*
- $H_A(q^2)$ determines lepton forward-backward asymmetry

helicity and transversity amplitudes:

[Krüger/Matias 05]

$$H_T(q^2) = |A_{\perp L}|^2 + |A_{\perp R}|^2 + |A_{\parallel L}|^2 + |A_{\parallel R}|^2,$$

$$H_L(q^2) = |A_{0L}|^2 + |A_{0R}|^2,$$

$$H_A(q^2) = 2 \operatorname{Re} [A_{\parallel R} A_{\perp R}^* - A_{\parallel L} A_{\perp L}^*].$$

Following [Krüger/Matias 05] and [Beneke/Feldmann/Seidel 01,04], the transversity amplitudes (relevant for an angular analysis of $B \rightarrow K^*(K\pi)\ell\ell$) can be written as

$$\begin{aligned}
 A_{\perp L,R} &\propto \left[(C_9 \mp C_{10}) \frac{V(q^2)}{m_B + m_{K^*}} + \frac{2m_b}{q^2} T_1(q^2) \right], \\
 A_{\parallel L,R} &\propto \left[(C_9 \mp C_{10}) \frac{A_1(q^2)}{m_B - m_{K^*}} + \frac{2m_b}{q^2} T_2(q^2) \right], \\
 A_{0 L,R} &\propto \left[(C_9 \mp C_{10}) \left\{ \frac{A_1(q^2)}{m_B - m_{K^*}} - \frac{m_B^2 - q^2}{m_B^2} \frac{A_2(q^2)}{m_B + m_{K^*}} \right\} \right. \\
 &\quad \left. + \frac{2m_b}{m_B^2} \left\{ T_2(q^2) - \frac{m_B^2 - q^2}{m_B^2} T_3(q^2) \right\} \right],
 \end{aligned}$$

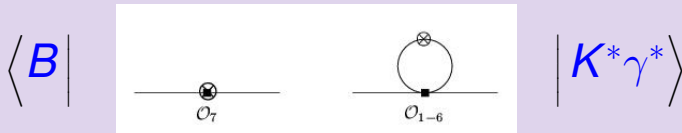
(some terms $m_{K^*}^2/m_B^2$ neglected, kinematic normalization factors can be found in [Krüger/Matias])

- $C_{9,10}$: short-distance Wilson coefficients for $b \rightarrow s\ell\ell$ (to be tested against NP)
- $V, A_{1,2}$: vector/axial-vector $B \rightarrow K^*$ transition form factors (sum rules/lattice)
- $T_i(q^2)$: factorizable and non-factorizable effects from virtual photons, via \mathcal{O}_{1-8} (QCDF/SCET)

factorization

- “naive” factorization:

⊗: photon insertion



$$T_1(q^2) \simeq C_7^{\text{eff}} T_1(q^2) + Y(q^2; C_{1-6}) \frac{q^2}{2m_b} \frac{V(q^2)}{m_B + m_{K^*}},$$

$$T_2(q^2) \simeq C_7^{\text{eff}} T_2(q^2) + Y(q^2; C_{1-6}) \frac{q^2}{2m_b} \frac{A_1(q^2)}{m_B - m_{K^*}},$$

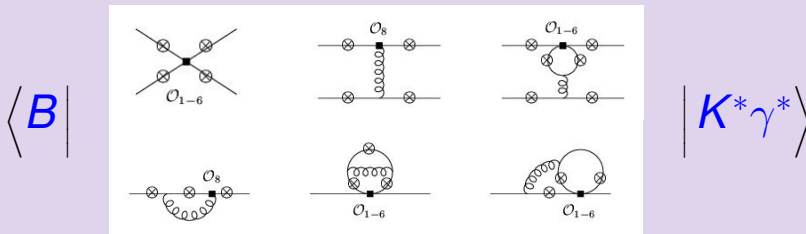
$$T_3(q^2) \simeq C_7^{\text{eff}} T_3(q^2) + Y(q^2; C_{1-6}) \left[\frac{m_B - m_{K^*}}{2m_b} A_2(q^2) - \frac{m_B + m_{K^*}}{2m_b} A_1(q^2) \right],$$

- only additional non-perturbative input: tensor form factors $T_i(q^2)$
- quark-loop function $Y(q^2)$ perturbative for $m_p^2 \ll q^2 \ll 4m_c^2$
sometimes: $C_9 \rightarrow C_9^{\text{eff}}(q^2) = C_9 + Y(q^2)$

factorization

- include annihilation and QCD corrections:

[Beneke/Feldmann/Seidel 01]



- also involve chromomagnetic penguin operator \mathcal{O}_8^g
- leading power in $1/m_b$ can be calculated using **QCDF/SCET** (except for annihilation in $H_L(q^2)$)
- new non-perturbative input in spectator diagrams: **light-cone wave functions**

$$\begin{aligned}
 \mathcal{T}_1(q^2) &\simeq \frac{m_B^2}{m_B^2 - q^2} \mathcal{T}_2(q^2) \\
 \mathcal{T}_3(q^2) &- \frac{m_B^2}{m_B^2 - q^2} \mathcal{T}_2(q^2) \simeq \begin{cases} \xi_{\perp}(q^2, \mu) C_{\perp}(q^2, \mu) + \phi_B^{\pm}(\omega, \mu) \otimes \phi_{K^*}^{\perp}(u, \mu) \otimes T_{\perp}(\omega, u, \mu) \\ \xi_{\parallel}(q^2, \mu) C_{\parallel}(q^2, \mu) + \phi_B^{\pm}(\omega, \mu) \otimes \phi_{K^*}^{\parallel}(u, \mu) \otimes T_{\parallel}(\omega, u, \mu) \end{cases}
 \end{aligned}$$

simplifications

- In the limit $m_b \rightarrow \infty$ and $q^2 \ll m_b^2$,
number of independent $B \rightarrow K^*$ form factors reduces to 2: [Charles et al. 98]

$$\xi_{\perp}(q^2, \mu), \quad \xi_{\parallel}(q^2, \mu) \quad [\text{scale-/scheme-dependent (SCET)}]$$

[→ forward-backward asymmetry zero, ...]

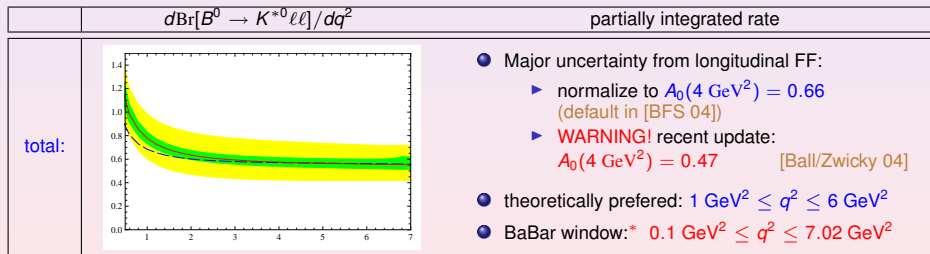
complications

- $1/m_b$ corrections induce new (partly unknown) non-perturbative parameters/functions. [→ isospin asymmetry, ...]
- theoretical working window : $1 \text{ GeV}^2 \lesssim q^2 \lesssim 6 \text{ GeV}^2$
 - ▶ dominance of C_7 at low q^2
 - ▶ presence of light and heavy vector meson resonances
 - ▶ non-factorizable effects in longitudinal rate

[in principle, also the region $4m_c^2 \ll q^2 \leq M_B^2$ can be addressed in HQET ⊗ ChPT]

rates

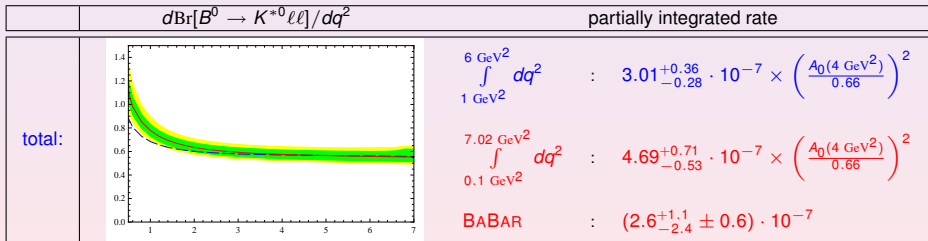
- using input values as in [Beneke/TF/Seidel 04] except for new life times from HFAG.
(tensor form factor $\rightarrow \xi_{\perp}(0)$ from $B \rightarrow K^* \gamma$, all numbers preliminary)
- dashed line: LO approximation (incl. some tree-level $1/m_b$ corrections)
- yellow band: “all” uncertainties
- green band: w/o error on form factors and CKM input.



* (just for comparison)

rates

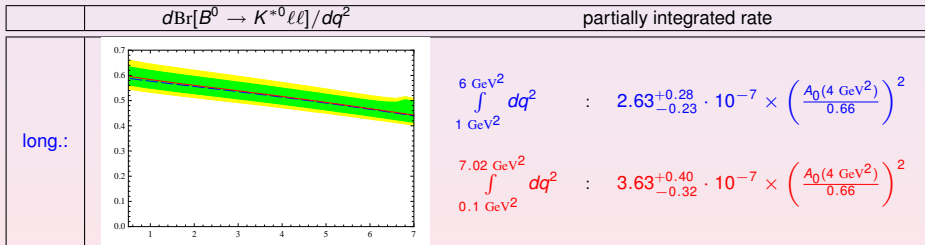
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- **dashed line:** LO approximation (incl. some tree-level $1/m_b$ corrections)
- **yellow band:** “all” uncertainties, but $A_0(q^2 = 4 \text{ GeV}^2)$ factored out.
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rates

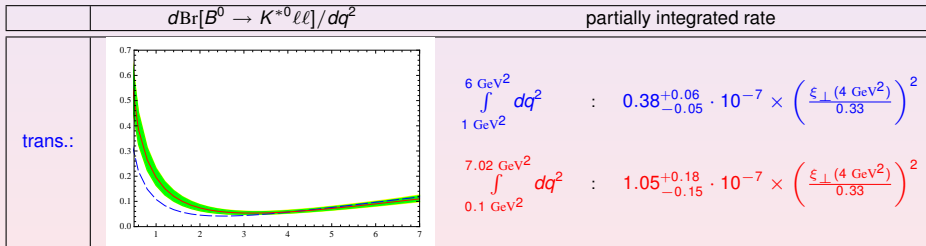
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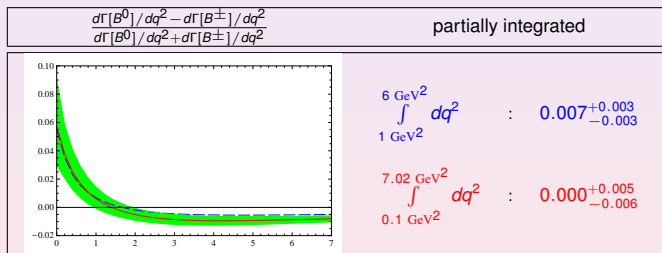
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isospin asymmetry

- using input values as in [Beneke/TF/Seidel 04] except for new life times from HFAG.
(tensor form factor $\rightarrow \xi_{\perp}(0)$ from $B \rightarrow K^* \gamma$, **all numbers preliminary**)
- **dashed line**: LO approximation (incl. some tree-level $1/m_b$ corrections)
- **green band**: all uncertainties.



[see also TF/Matias 03]

forward-backward asymmetry

- using input values as in [Beneke/TF/Seidel 04]
(tensor form factor $\rightarrow \xi_{\perp}(0)$ from $B \rightarrow K^* \gamma$, all numbers preliminary)
- **dashed line:** LO approximation (incl. some tree-level $1/m_b$ corrections)
- **green band:** all parametric uncertainties. (form factor dependence decreases around q_0^2)

| | $B^0 \rightarrow K^{*0} \ell \ell$ | $B^{\pm} \rightarrow K^{*\pm} \ell \ell$ |
|---|--|--|
| $\frac{(d\Gamma^F - d\Gamma^B)/dq^2}{(d\Gamma^F + d\Gamma^B)/dq^2}$ | | |
| FBA zero q_0^2 : | $4.36^{+0.33}_{-0.31} \text{ GeV}^2$ | $4.15^{+0.27}_{-0.27} \text{ GeV}^2$ |
| $\frac{\Gamma^F - \Gamma^B}{\Gamma^F + \Gamma^B}$ | $\int_{1 \text{ GeV}^2}^{q_0^2} dq^2 : -0.062^{+0.018}_{-0.023}$ $\int_{6 \text{ GeV}^2}^{q_0^2} dq^2 : +0.033^{+0.016}_{-0.015}$ | $\int_{1 \text{ GeV}^2}^{q_0^2} dq^2 : -0.058^{+0.016}_{-0.021}$ $\int_{6 \text{ GeV}^2}^{q_0^2} dq^2 : +0.036^{+0.016}_{-0.013}$ |

warnings

- systematic uncertainties from (partly) neglected $1/m_b$ corrections!
- how precise are the form factor estimates from sum rules/lattice?
is it better to extract tensor form factors from experimental data on $B \rightarrow K^* \gamma$
or from sum rules/lattice ?
- how reliable are the phenomenological estimates for light-cone wave functions?
- how much do vector meson poles influence the intermediate q^2 region?
- part(!) of perturbative uncertainty could be reduced by resumming logs in SCET.
[Ali/Kramer/Zhu, hep-ph/0601034]
- ...

new physics

- minimal flavour violation:

SM operators, but different values for Wilson coefficients. (✓)

- generic flavour violation:

new operators with different chirality structure and/or new CP phases.

(→ formulas to be generalized → see backup)

⇒ difficult to establish deviations from SM via (partially integrated) rates!

(form factor uncertainties \oplus non-factorizable effects)

⇒ concentrate on asymmetries

(forward-backward, isospin, CP, angular)

other modes (briefly)

- $B \rightarrow K\ell^+\ell^-$:
 - ▶ $R_K = \Gamma[B \rightarrow K\mu\mu]/\Gamma[B \rightarrow Kee]$ [Bobeth/Hiller/Piranishvili, 0709.4174]
- $B_s \rightarrow \mu^+\mu^-$:
 - ▶ enhanced by large powers of $\tan\beta$ in type-II 2HDM [see previous talks; also Nierste arXiv:0807.3733]
- $B_s \rightarrow \phi\ell^+\ell^-$:
 - ▶ NP (model-independent) \oplus naive factorization [Yilmaz, 0806.0269]
 - ▶ numerical analysis incl. non-factorizable effects [Beneke/TF/Seidel, in preparation]
- ...
- $B \rightarrow \rho\ell^+\ell^-$: [Beneke/TF/Seidel, 04]
 - ▶ different CKM hierarchy \rightarrow potentially larger isospin and CP effects
 - ▶ useful cross-check to test factorization approach

- What else should appear in the write-up?
→ open for discussion ...

$$\begin{aligned}
 A_{\perp L,R} &\propto \left[(C_9 \mp C_{10}) \frac{V(q^2)}{m_B + m_{K^*}} + \frac{2m_b}{q^2} T_1(q^2) \right], & C_{7,9,10} &\rightarrow C_{7,9,10} + C'_{7,9,10} \\
 A_{\parallel L,R} &\propto \left[(C_9 \mp C_{10}) \frac{A_1(q^2)}{m_B - m_{K^*}} + \frac{2m_b}{q^2} T_2(q^2) \right], & C_{7,9,10} &\rightarrow C_{7,9,10} - C'_{7,9,10} \\
 A_{0 L,R} &\propto (C_9 \mp C_{10}) \left\{ \frac{A_1(q^2)}{m_B - m_{K^*}} - \frac{m_B^2 - q^2}{m_B^2} \frac{A_2(q^2)}{m_B + m_{K^*}} \right\} \\
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 \end{aligned}$$

- $C_{9,10}$: short-distance Wilson coefficients for $b \rightarrow s\ell\ell$
- $V, A_{1,2}$: vector/axial-vector $B \rightarrow K^*$ transition form factors
- $T_i(q^2)$: factorizable and non-factorizable effects from virtual photons, via $\mathcal{O}_{1-8}, \mathcal{O}'_7$

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 &\quad + \frac{2m_b}{m_B^2} \left\{ T_2(q^2) - \frac{m_B^2 - q^2}{m_B^2} T_3(q^2) \right\}, & C_{7,9,10} &\rightarrow C_{7,9,10} - C'_{7,9,10}
 \end{aligned}$$

- $C_{9,10}'$: short-distance Wilson coefficients for $b \rightarrow s_{L/R} \ell \ell$
- $V, A_{1,2}$: vector/axial-vector $B \rightarrow K^*$ transition form factors
- $T_i(q^2)$: factorizable and non-factorizable effects from virtual photons, via $\mathcal{O}_{1-8}, \mathcal{O}'_7$