

Studio della produzione di materia (iper-)nucleare a LHC con l'esperimento ALICE

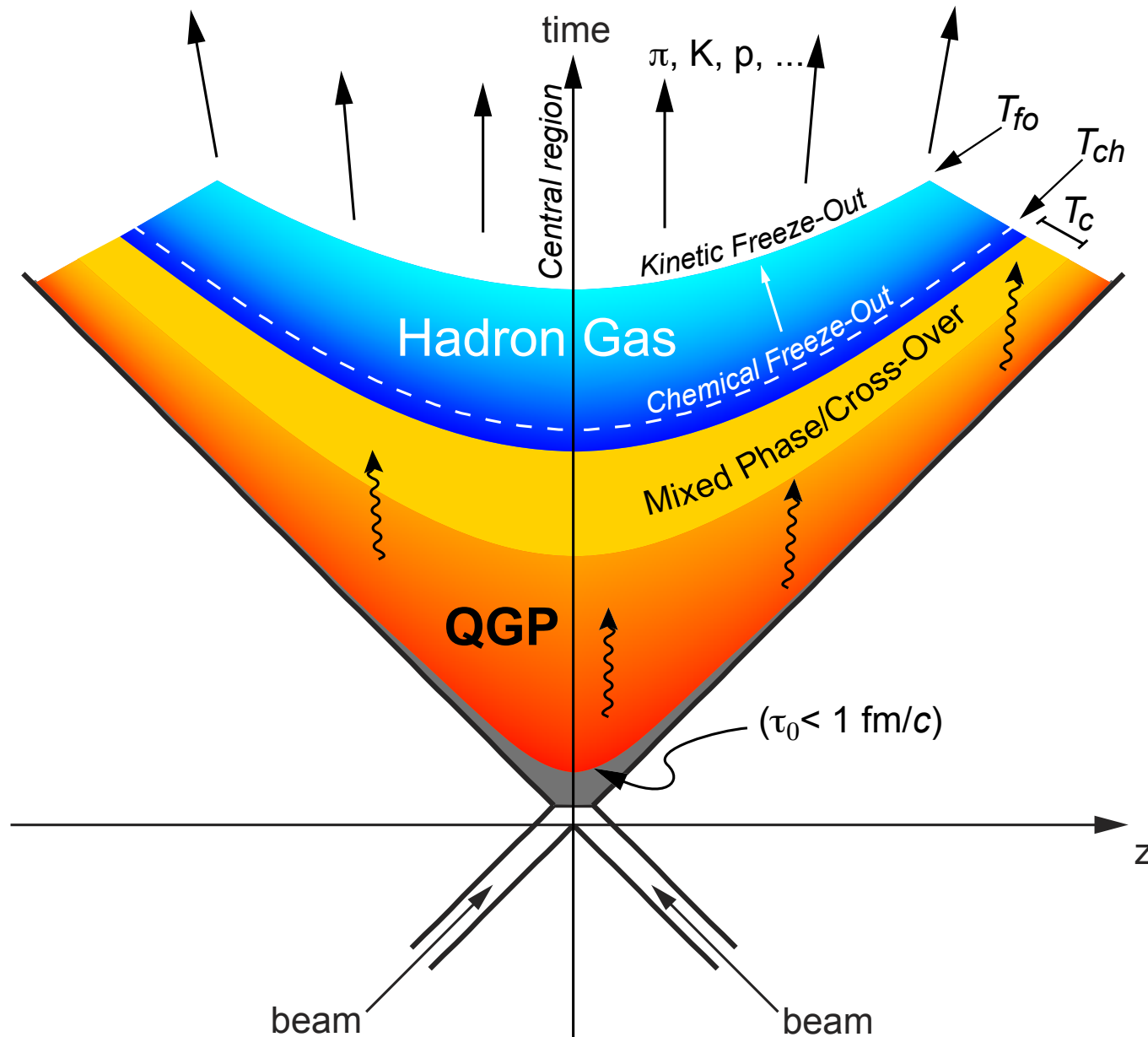
Maximiliano Puccio

Università & INFN Torino

Why studying nuclei in Heavy Ion Collisions?

In high energy Heavy Ions collisions

- a dense and hot partonic phase is created and undergoes a rapid expansion
- hydrodynamical models are used to describe this rapid expansion
- as a result of this evolution it is possible to observe collective phenomena



SOFT PROBES

- Low p_T ($p_T < 2 \text{ GeV}/c$) light flavoured particles coming from the interaction region
- Their constituents are produced in the late stages of the collisions
- They are useful to study the freeze-out conditions

Light nuclei and hyper-nuclei are soft probes and as they are loosely bound composite objects it is interesting to study their production in HI collisions.



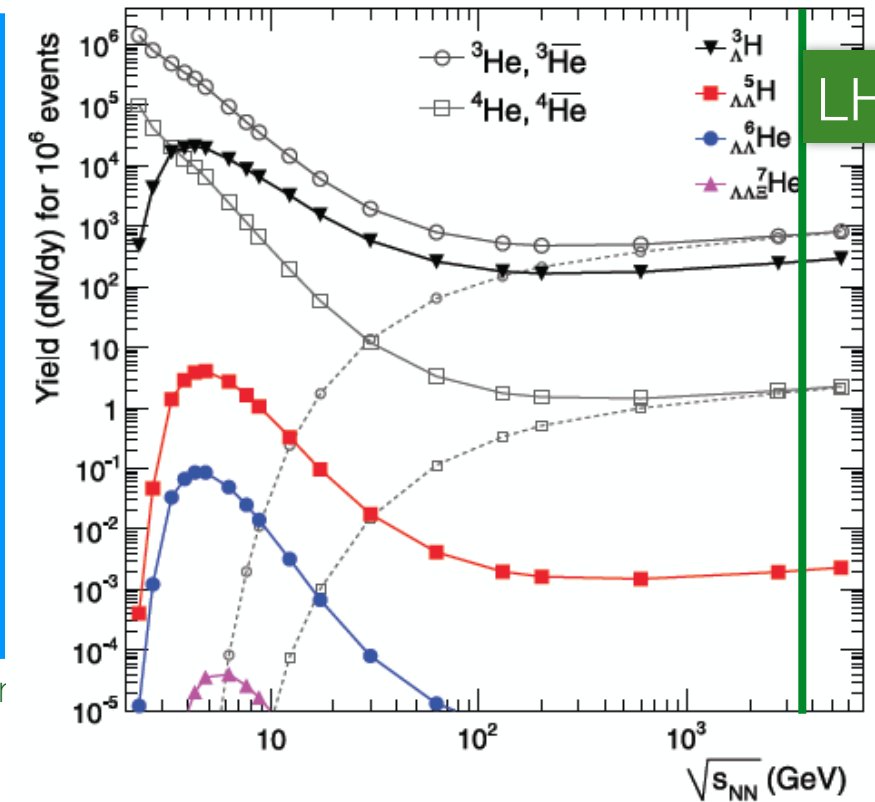
(Hyper-)Nuclei production models

(Hyper-)Nuclei production models

THERMAL MODELS

- Hadrons emitted from the interaction region in statistical equilibrium when the fireball reaches limiting temperature
- Abundances fixed at chemical freeze-out
- Freeze-out temperature T_{chem} is a key parameter
- Abundance of a species $\propto \exp(-m/T_{\text{chem}})$:
 - For nuclei (large m) strong dependence on T_{chem}

A. Andronic, P. Braun-Munzinger, J. Stachel and H. Stoecker
Phys. Lett. B607, 203 (2011), 1010.2995

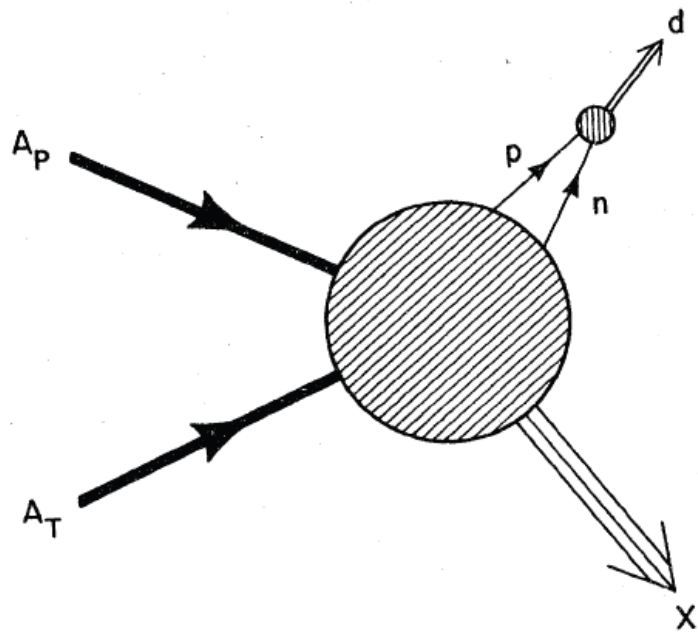
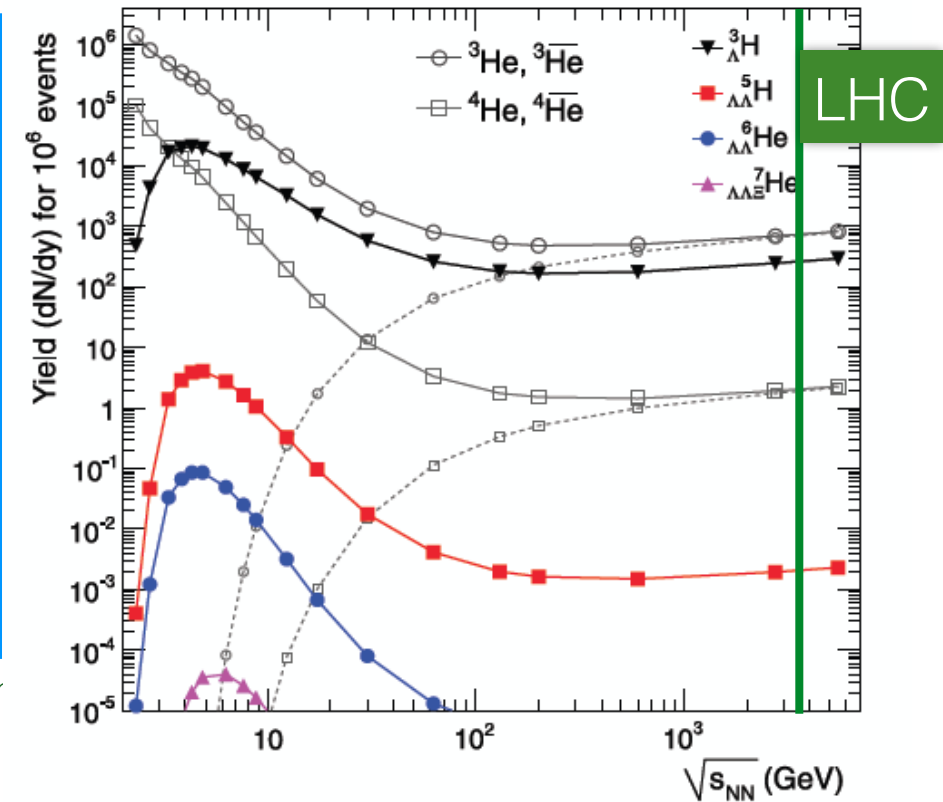


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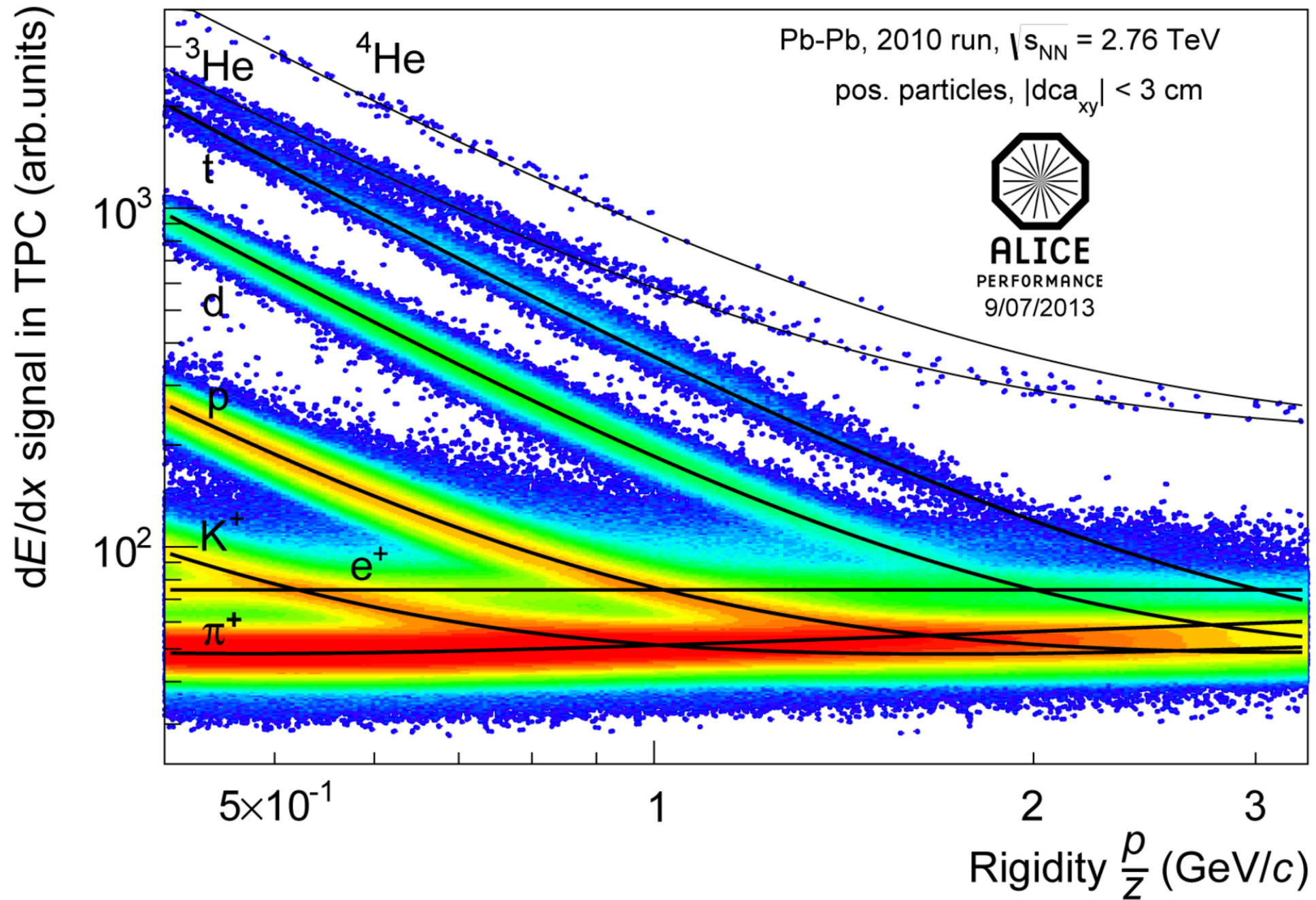


COALESCENCE

- If (anti-)baryons are close in phase space after the kinetic freeze-out they can form a (anti-)nucleus
- (Anti-)nuclei produced at the chemical freeze-out might break and re-form during the time span between the chemical freeze-out and the kinetic freeze-out.

J. I. Kapusta, Phys.Rev. C21, 1301 (1980)

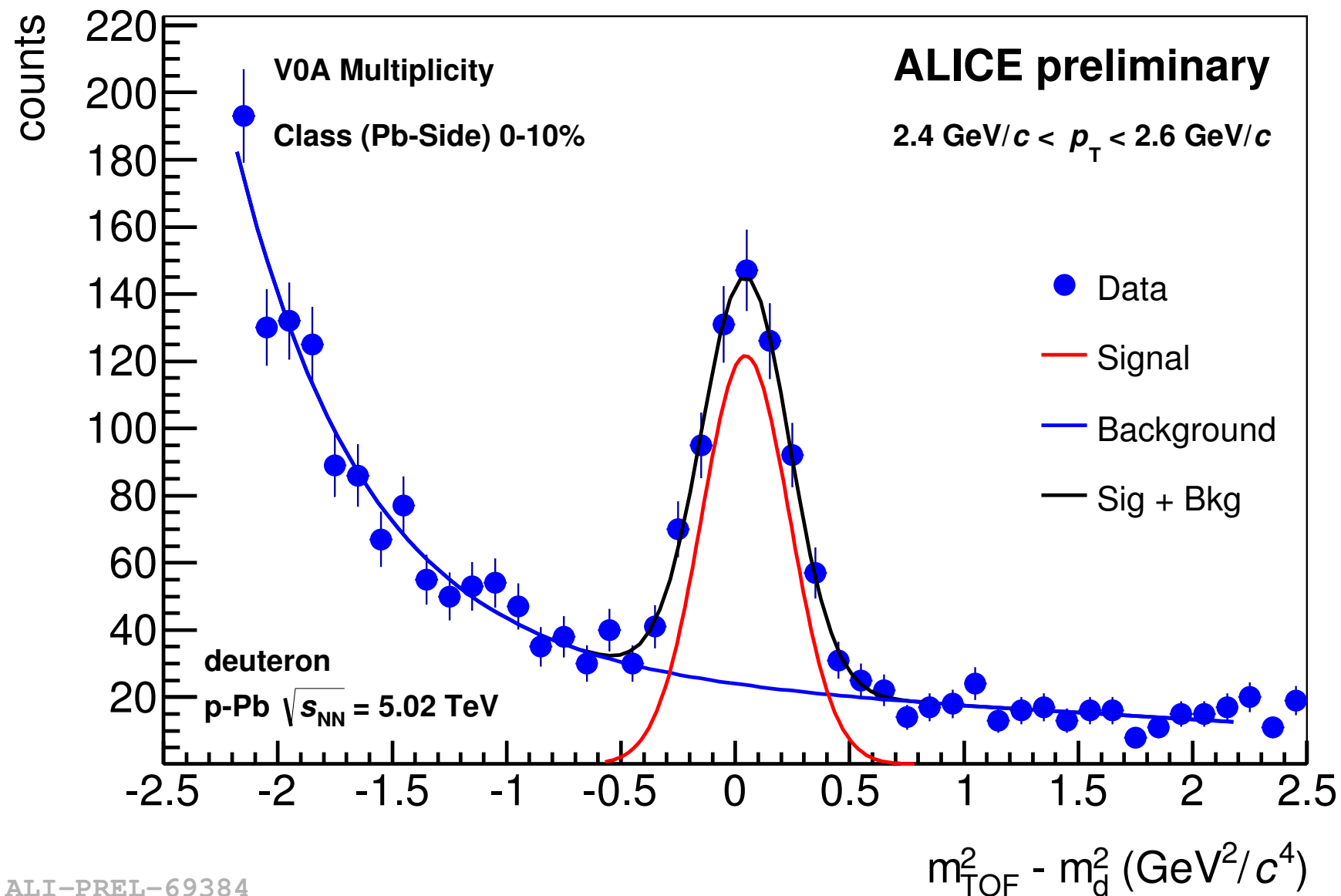
Detecting nuclei: TPC specific energy loss



Specific energy loss in the **T**ime **P**rojection **C**hamber volume provides an excellent PID for light nuclei.

→ $\sigma_{dE/dx} \sim 7\%$ (in Pb-Pb collisions)

Detecting nuclei: Time of flight

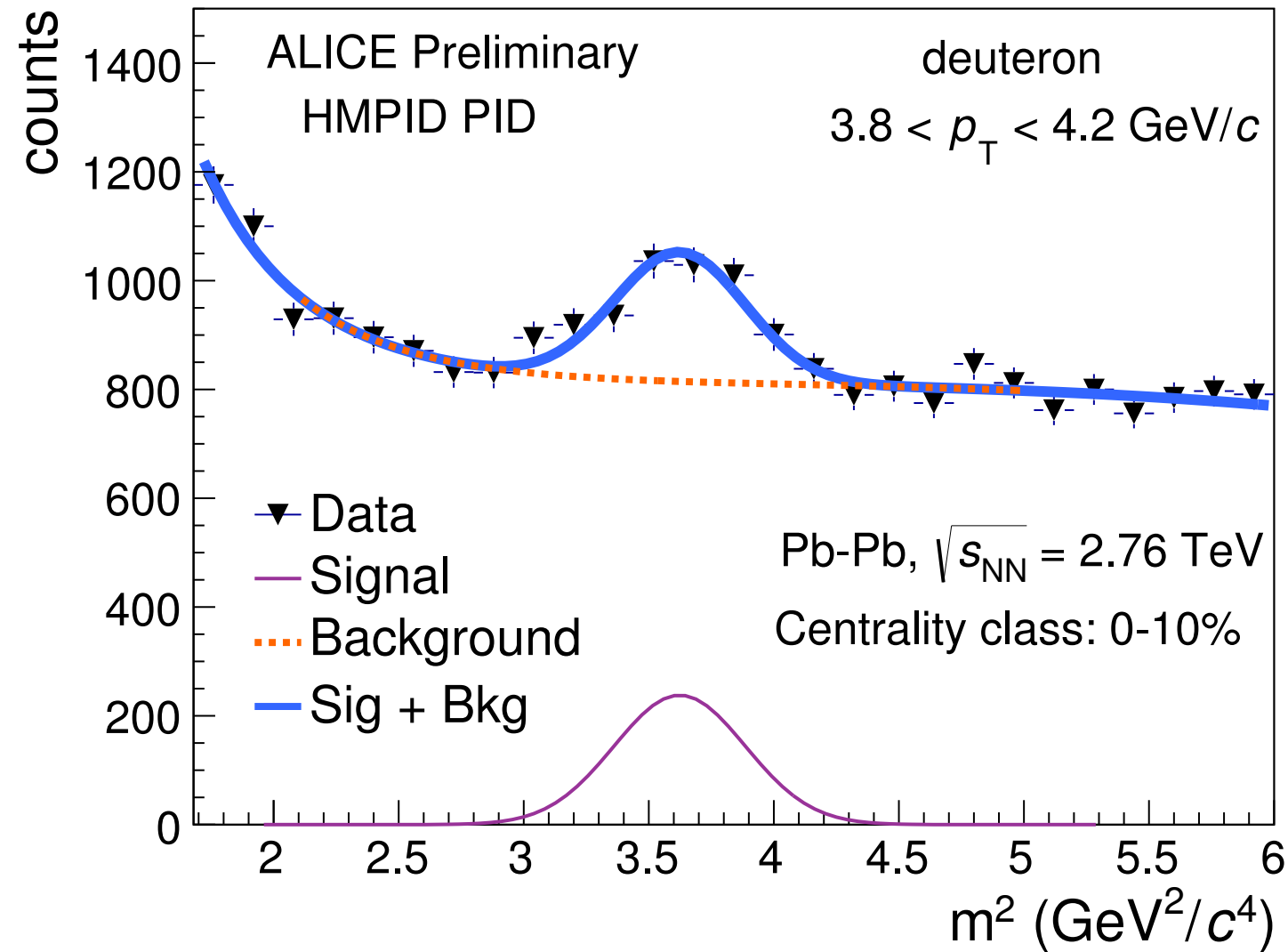


ALI-PREL-69384

Tracking information + time of flight measured by the **T**ime **O**f **F**light detector $\rightarrow m$ of the particle

$\rightarrow \sigma_{\text{time-of-flight}} \sim 80 \text{ ps}$ (in Pb-Pb)

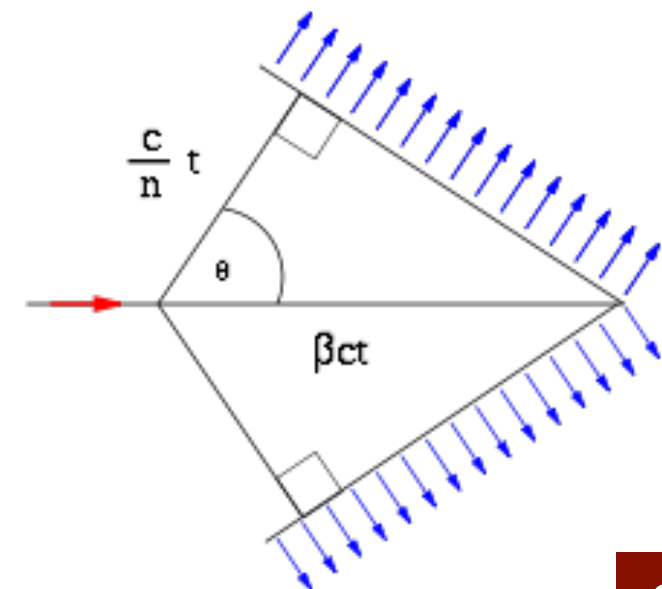
Detecting nuclei: Cherenkov light



ALI-PREL-86759

$$m_{\text{HMPID}} = \frac{p}{\beta\gamma} \quad \text{where} \quad \beta = \frac{1}{n \cdot \cos\theta_c}$$

Tracking information + **H**igh **M**omentum **P**article **I**dentification detector (Cherenkov light detector) signal
 → m of the particle



Detecting hyper-nuclei

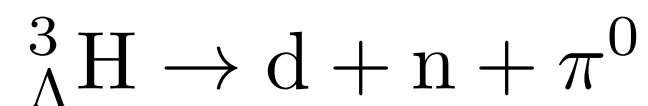
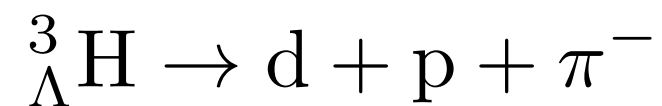
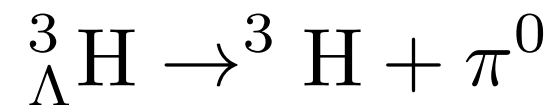
Hyper-triton is the lightest hyper-nucleus. ALICE collaboration measured its production in the charged 2 body decay channel.

Mass = 2.991 GeV/c²
Lifetime ~ 215 ps

Signal Extraction:

- Identify ³He and π
- Evaluate (³He,π) invariant mass
- Apply topological cuts in order to:
 - identify secondary decay vertex
 - reduce combinatorial background

DECAY MODES



+ anti-hypertriton

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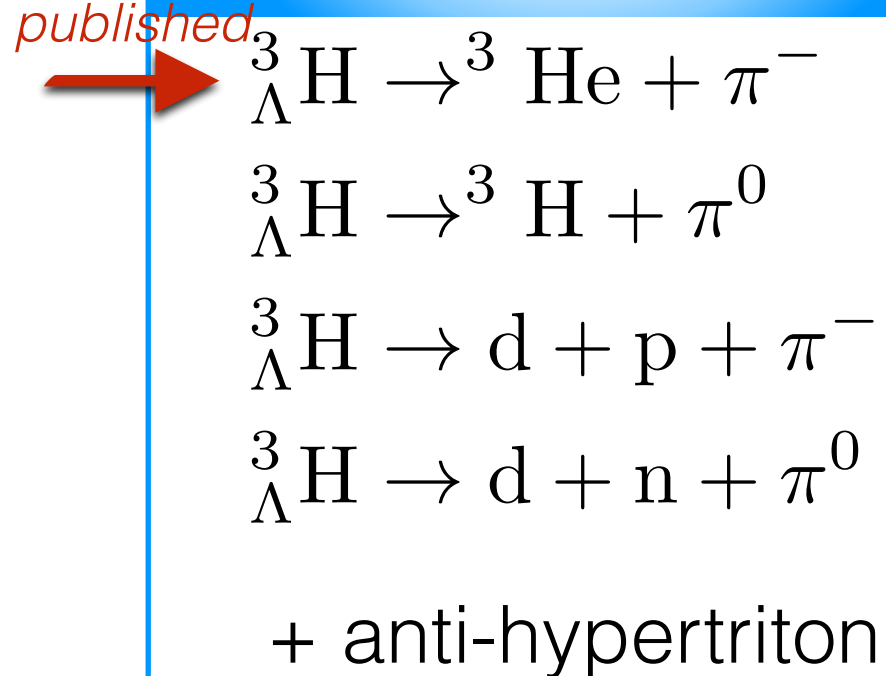
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published



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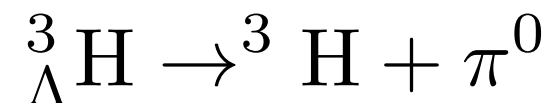
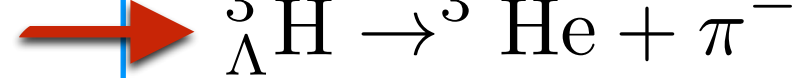
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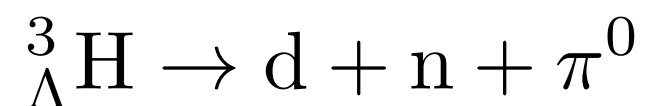
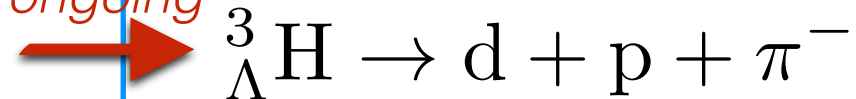
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DECAY MODES

published



ongoing



+ anti-hypertriton

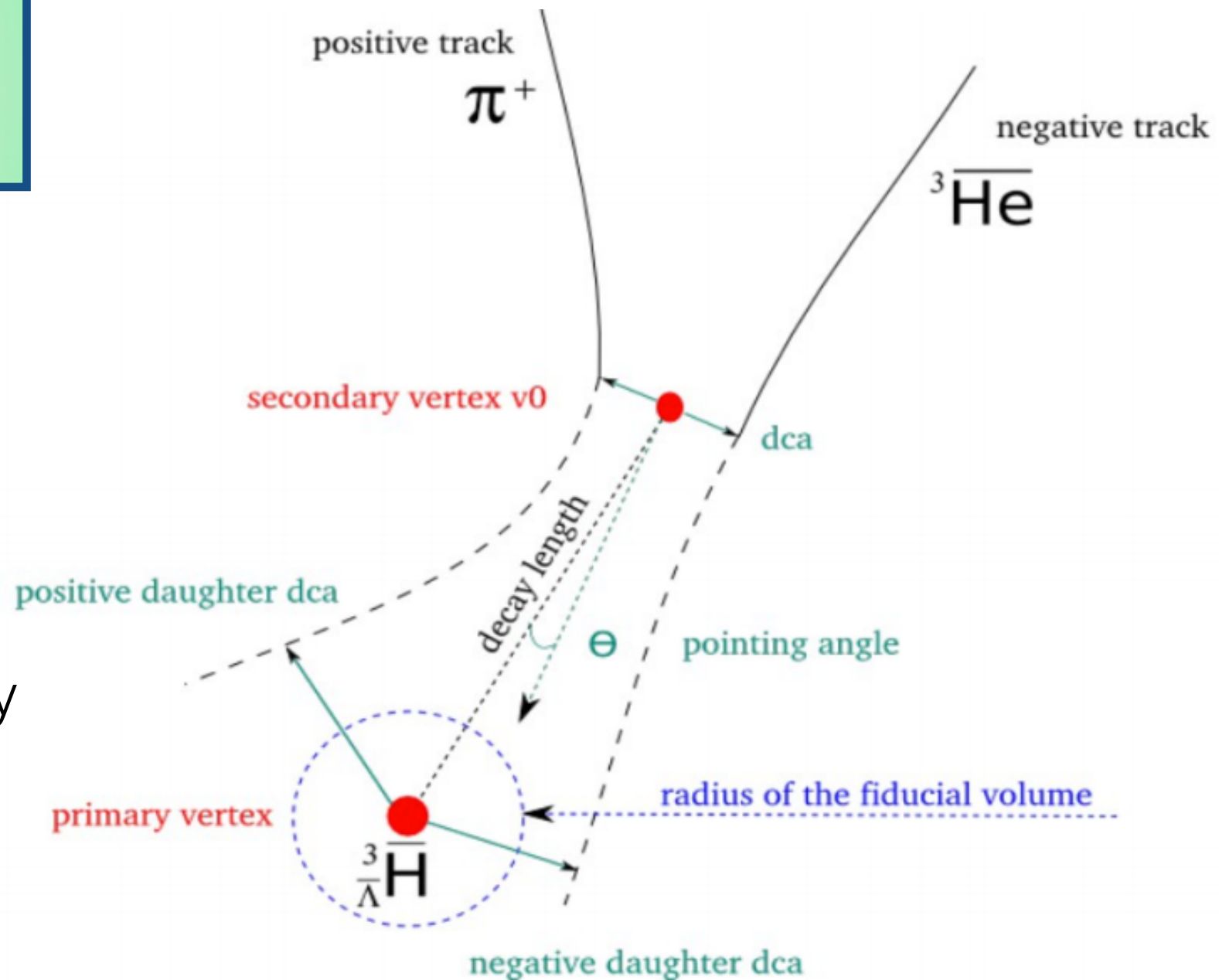
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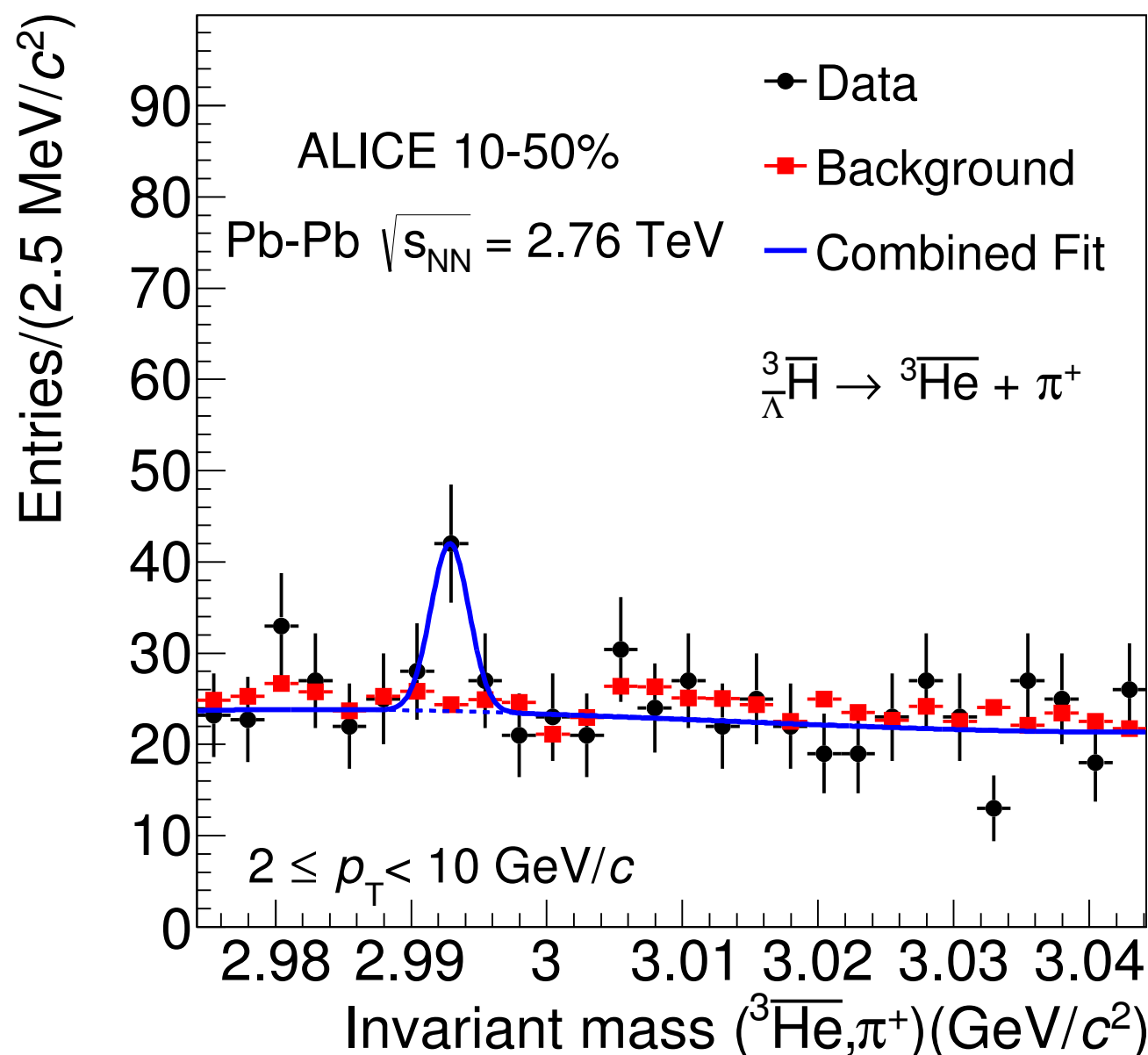
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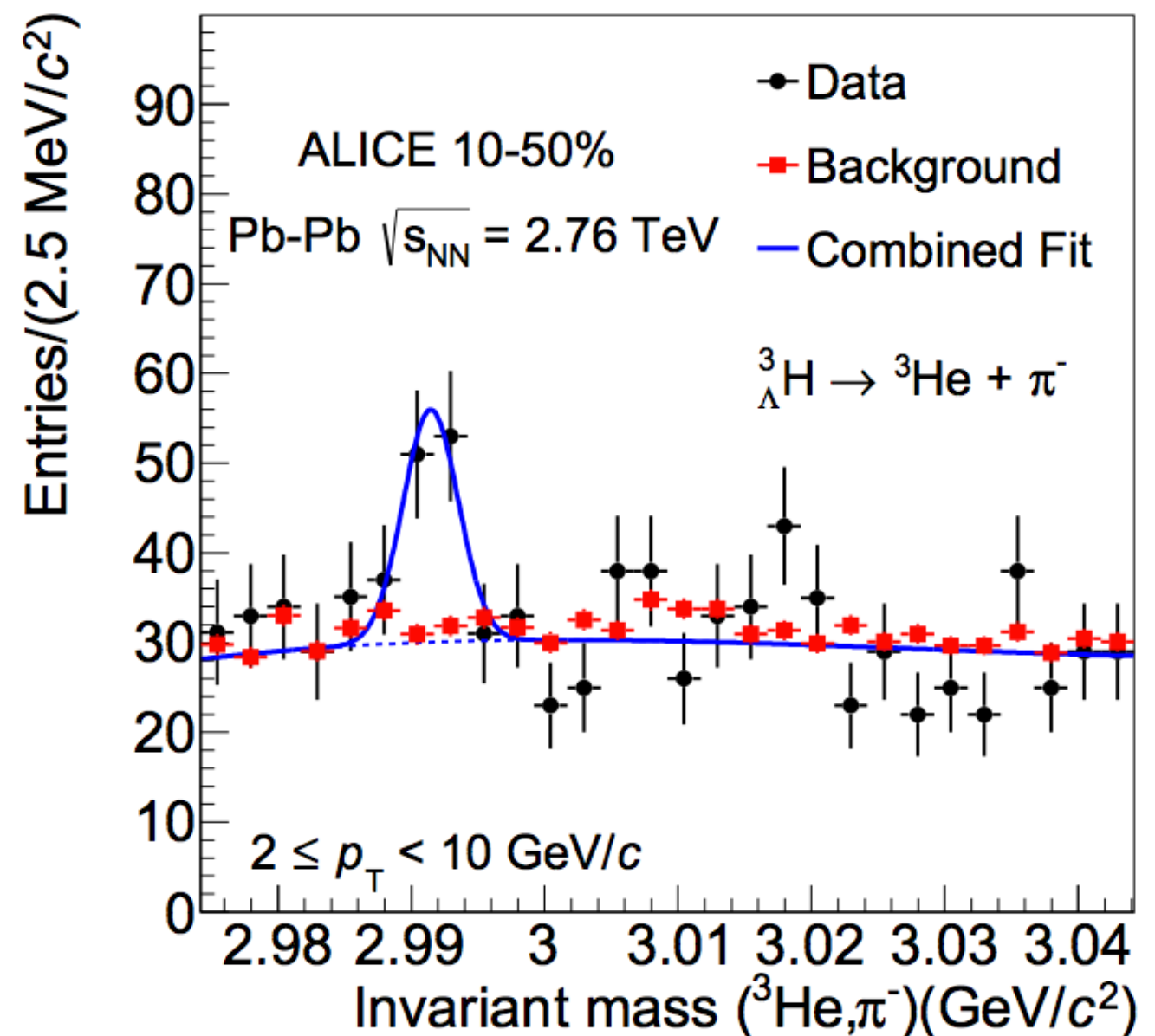
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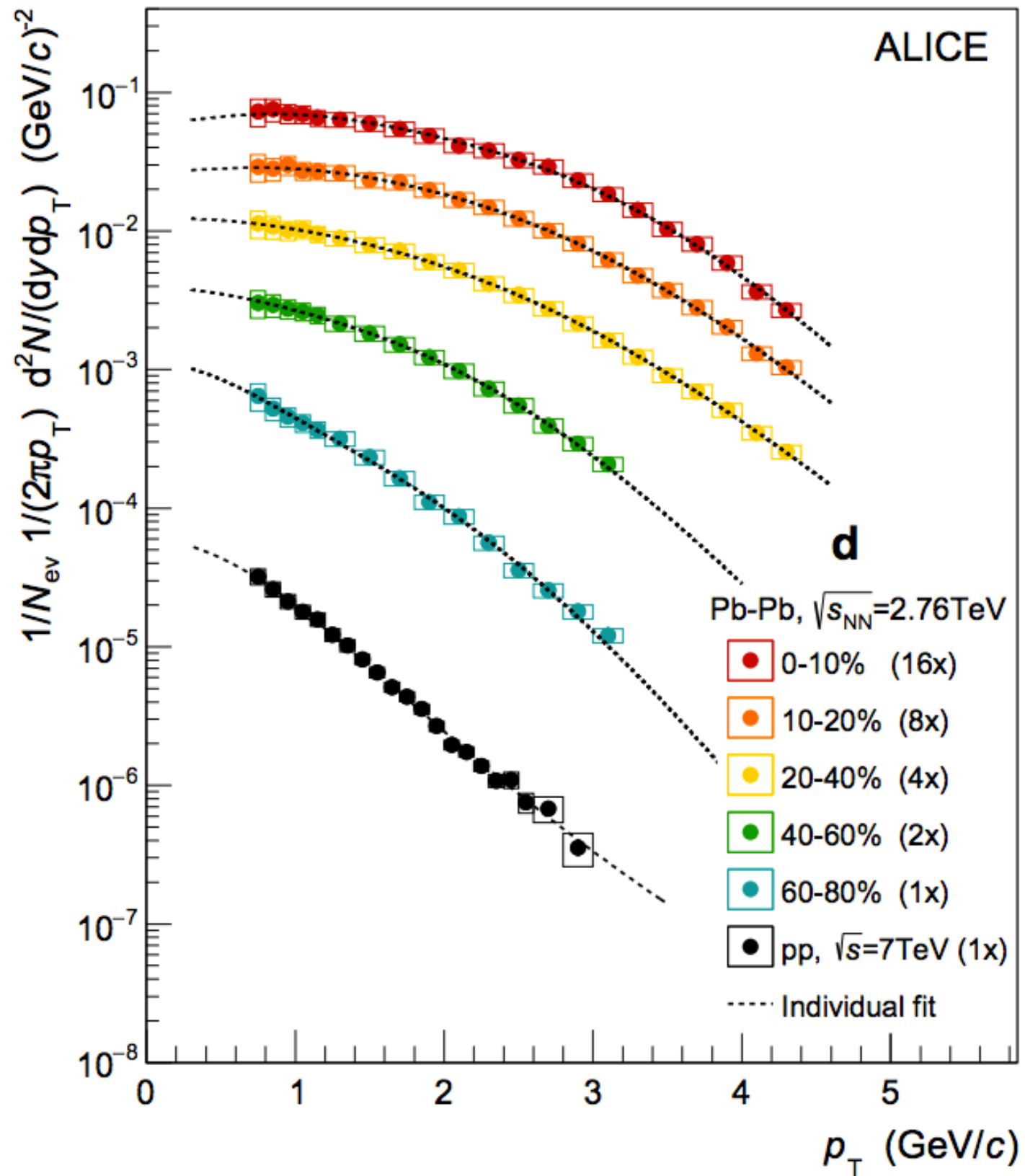
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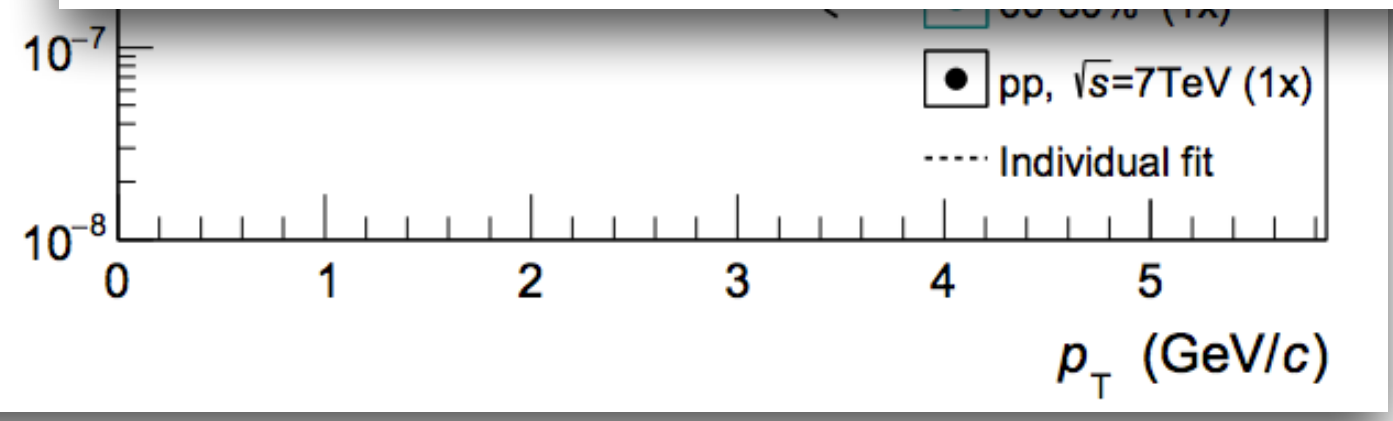
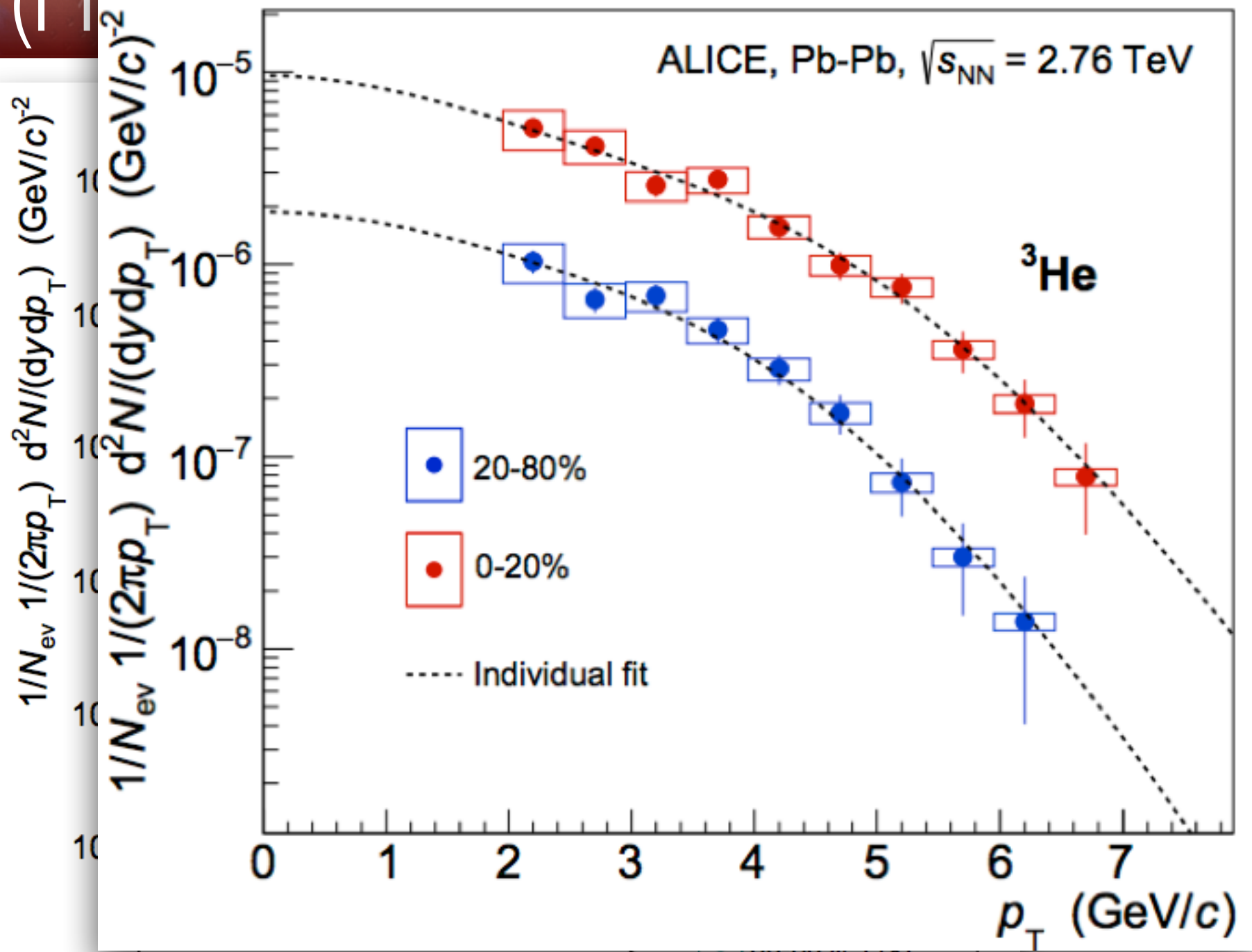
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(Hyper-)(Anti-)nuclei production

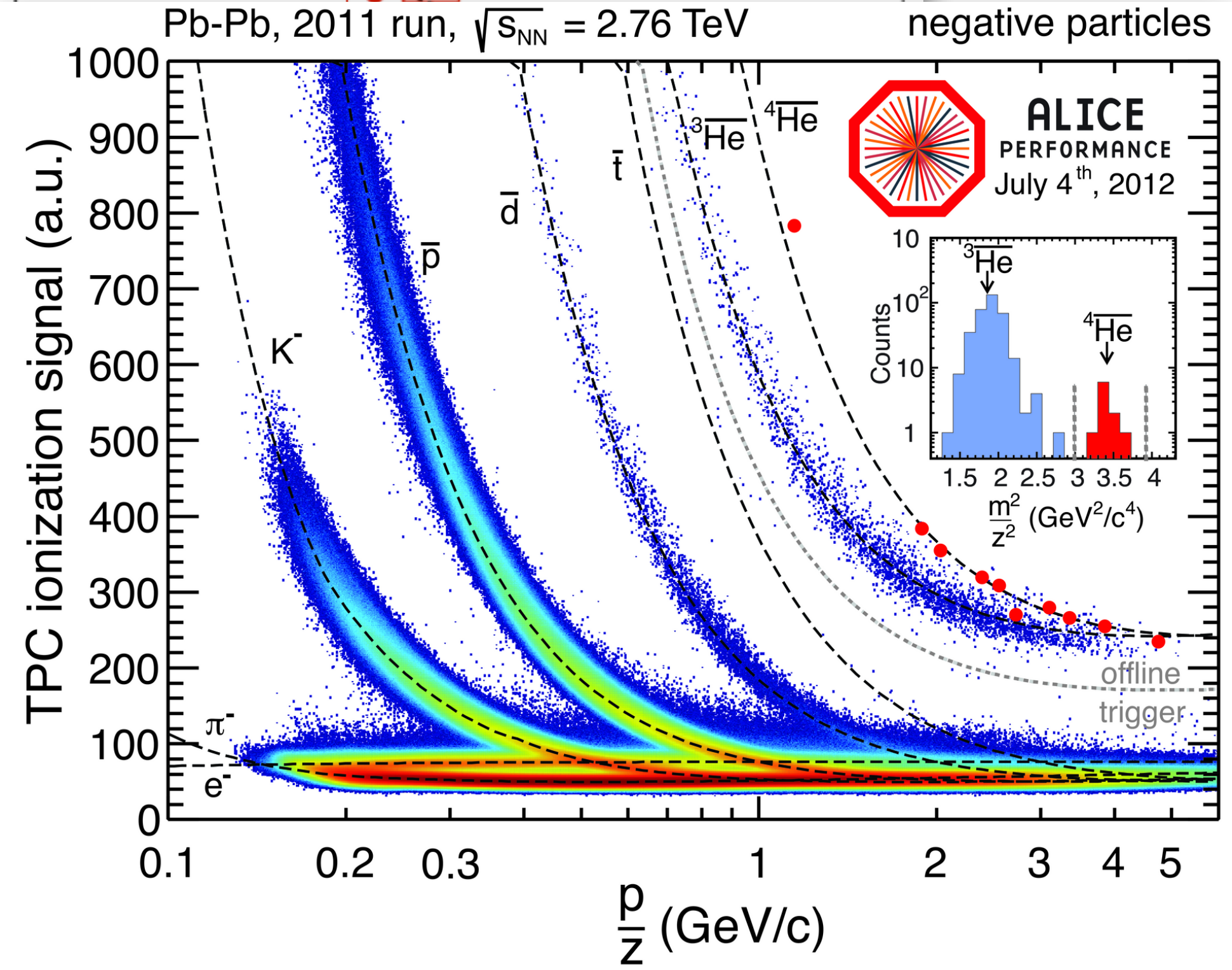
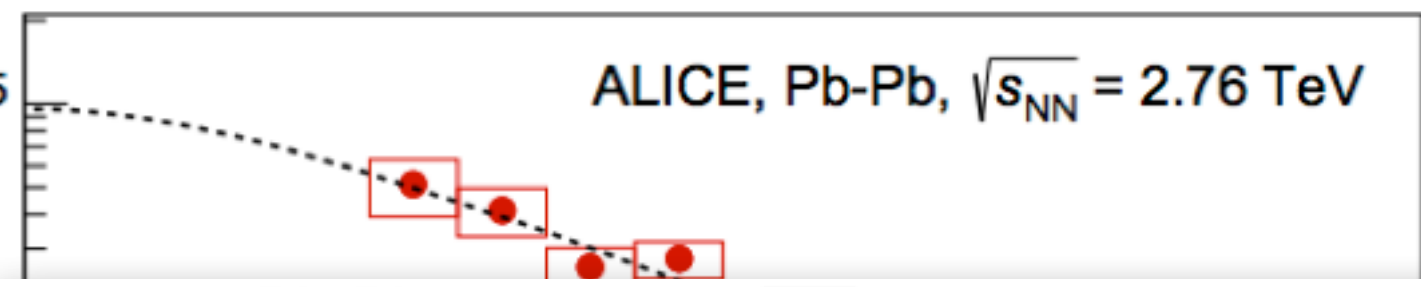


(H)

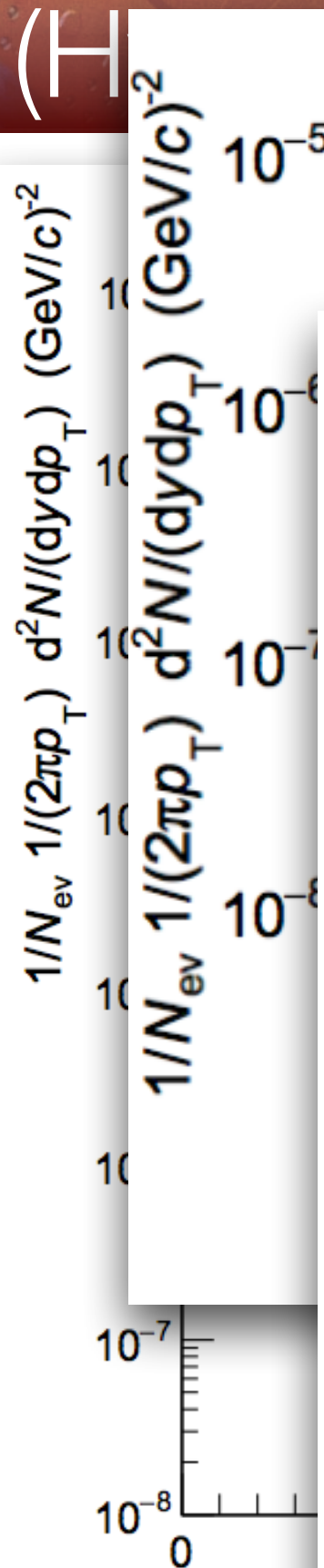


(H)

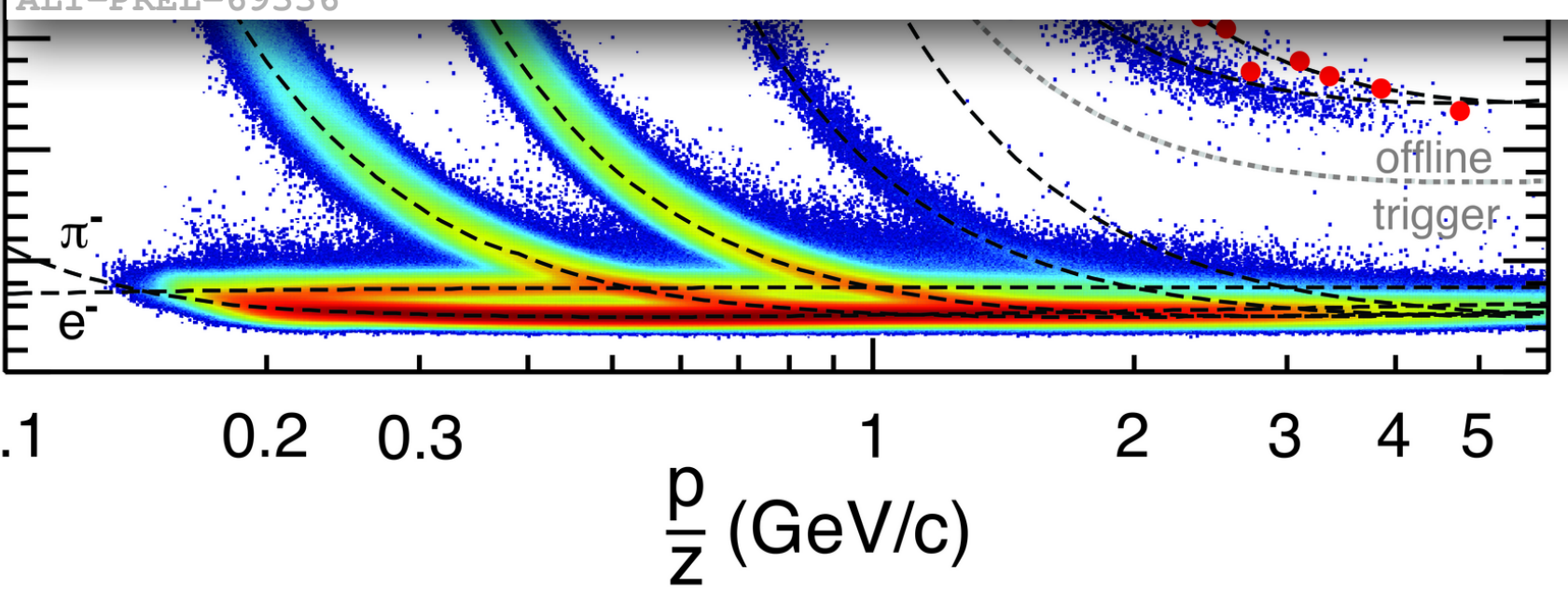
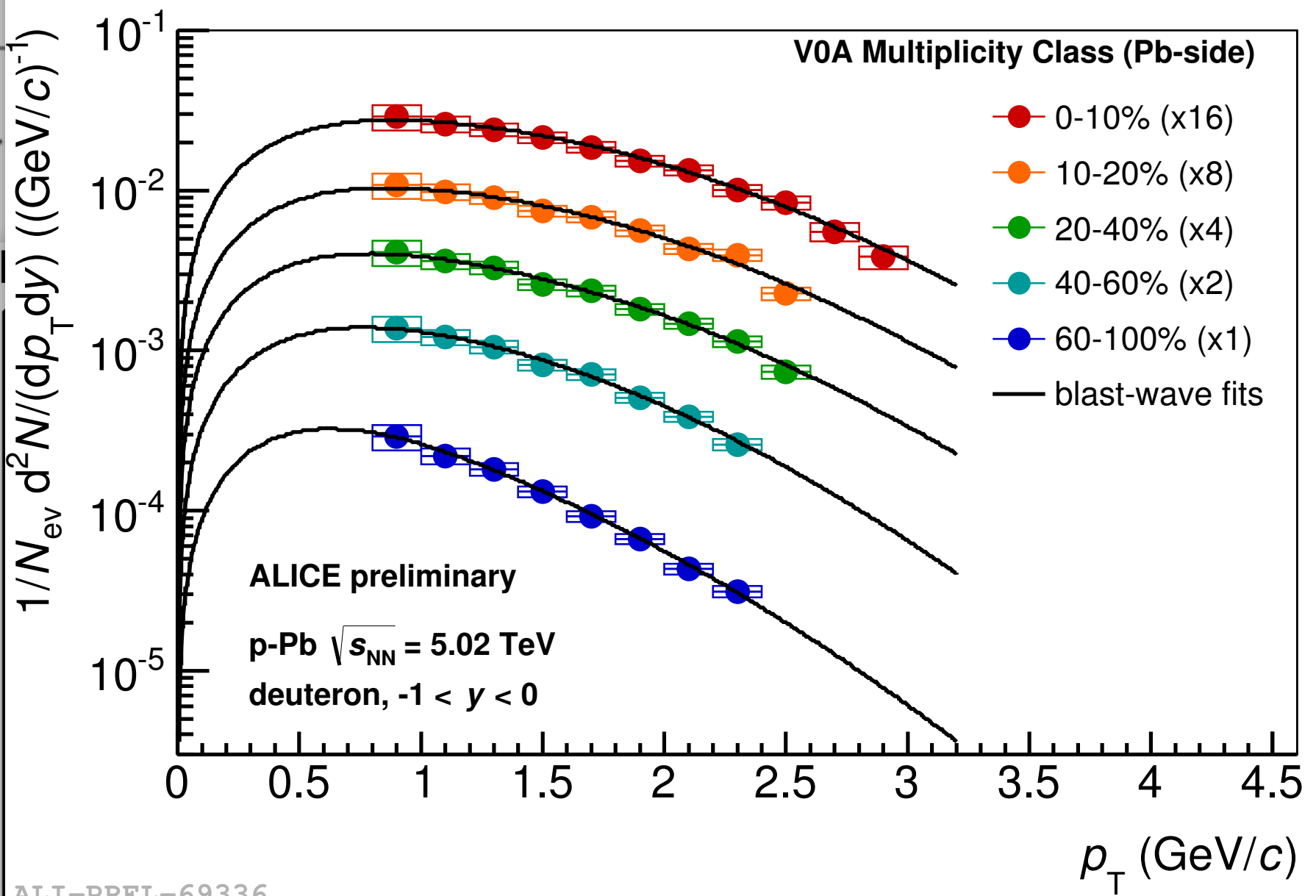
$1/N_{ev} 1/(2\pi p_T) d^2N/(dy dp_T) (GeV/c)^{-2}$
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 10^{-5}
 10^{-6}
 10^{-7}
 10^{-8}



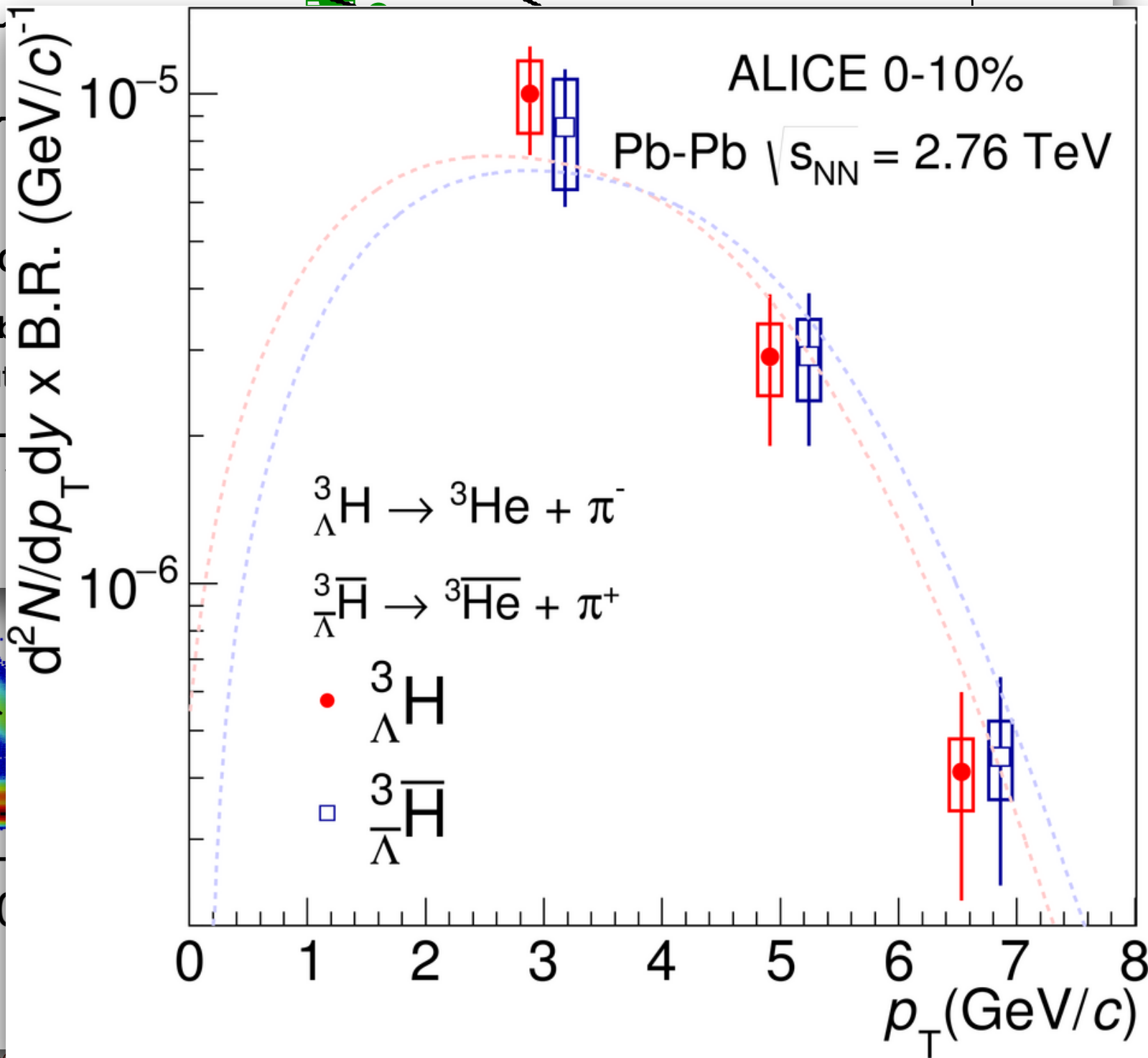
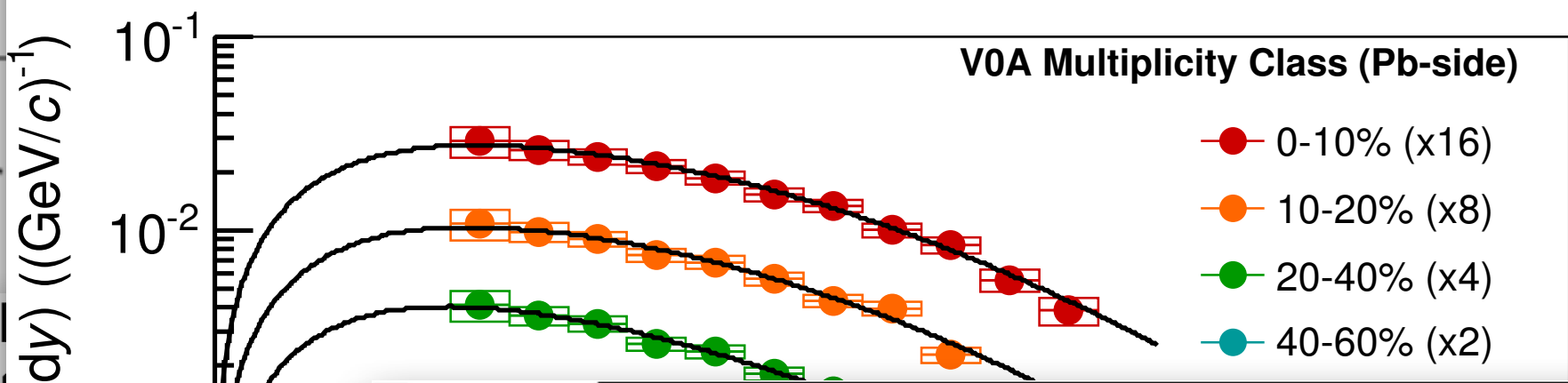
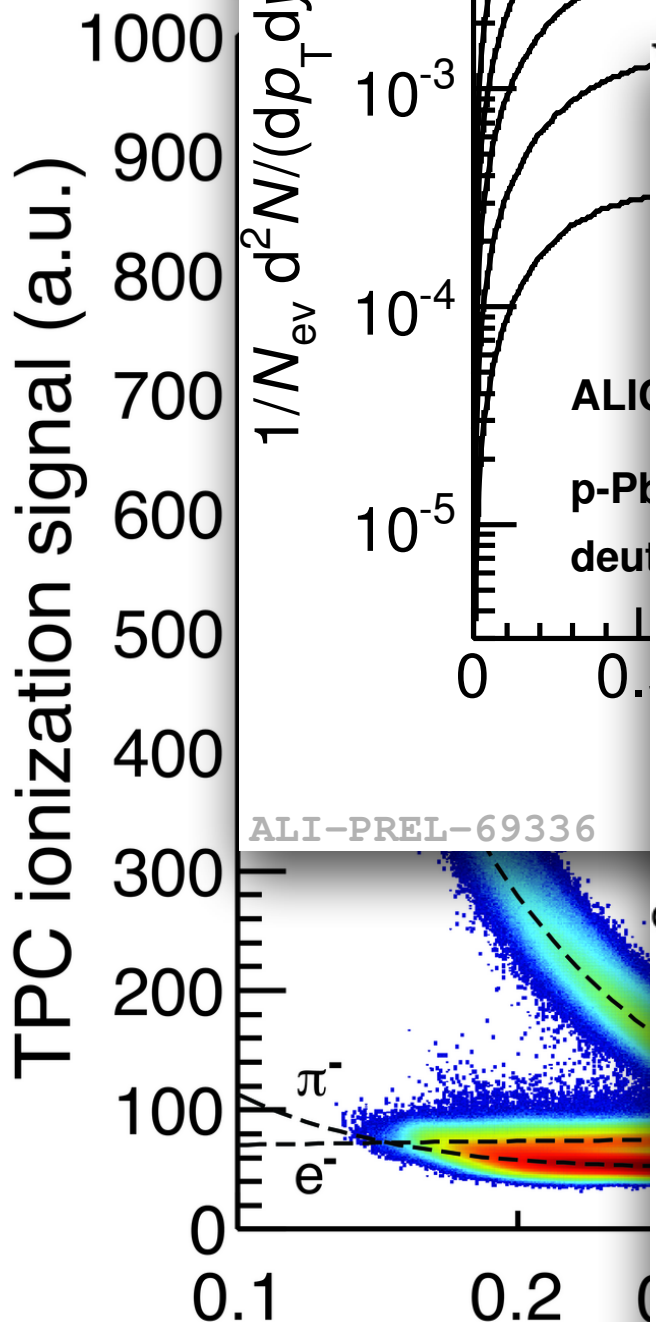
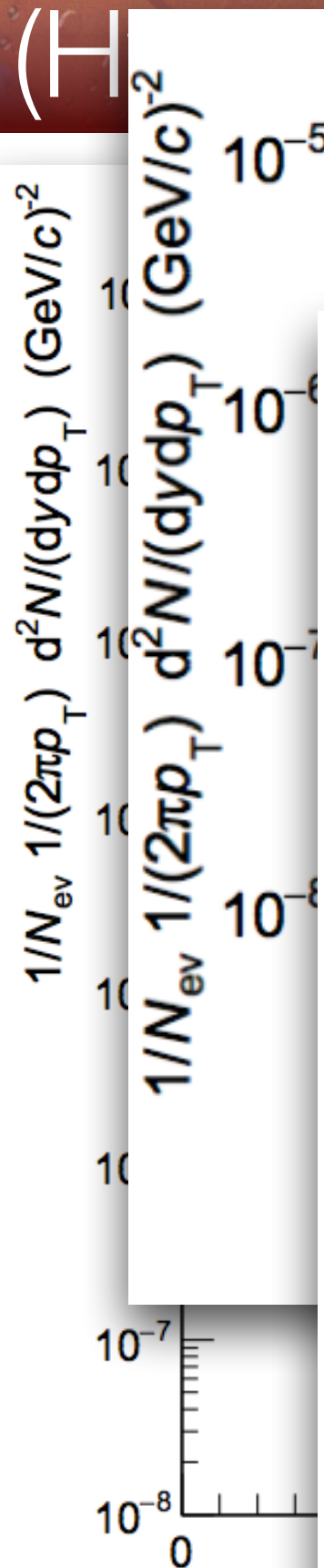
ALI-PERF-36713



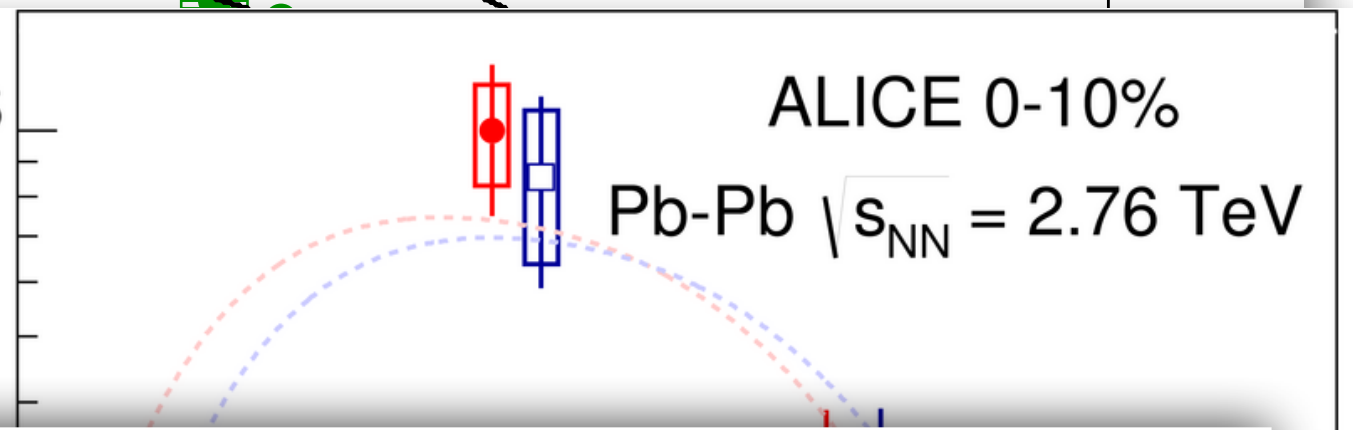
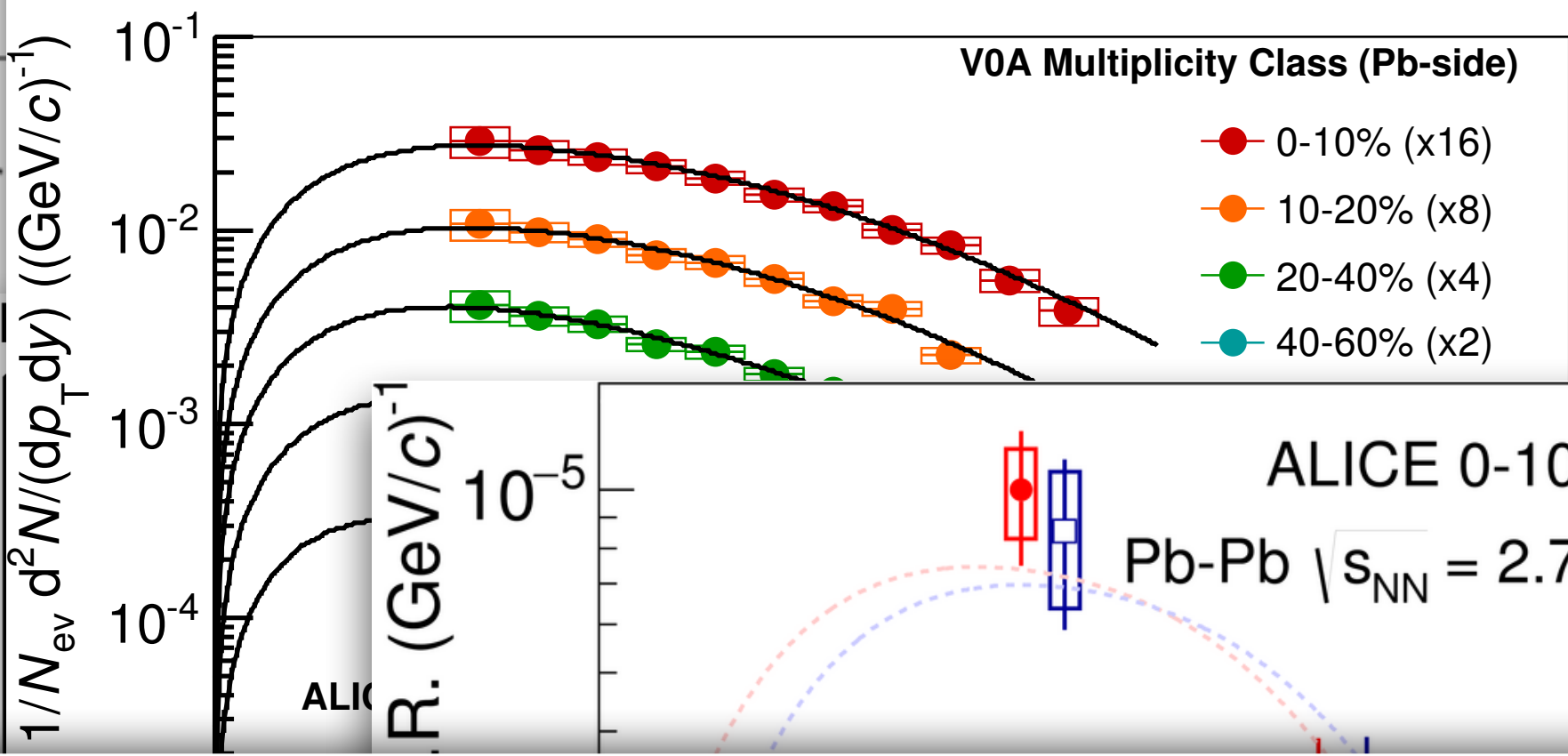
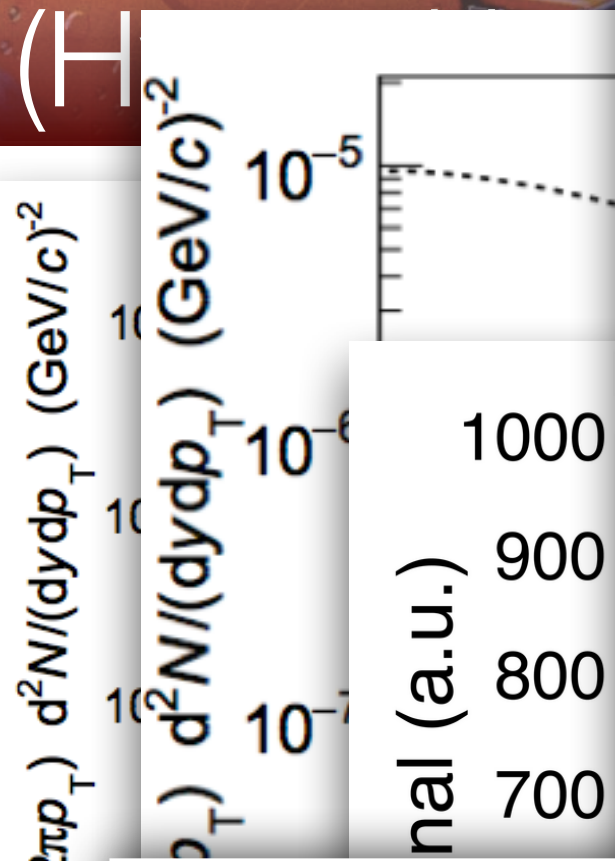
TPC ionization signal (a.u.)



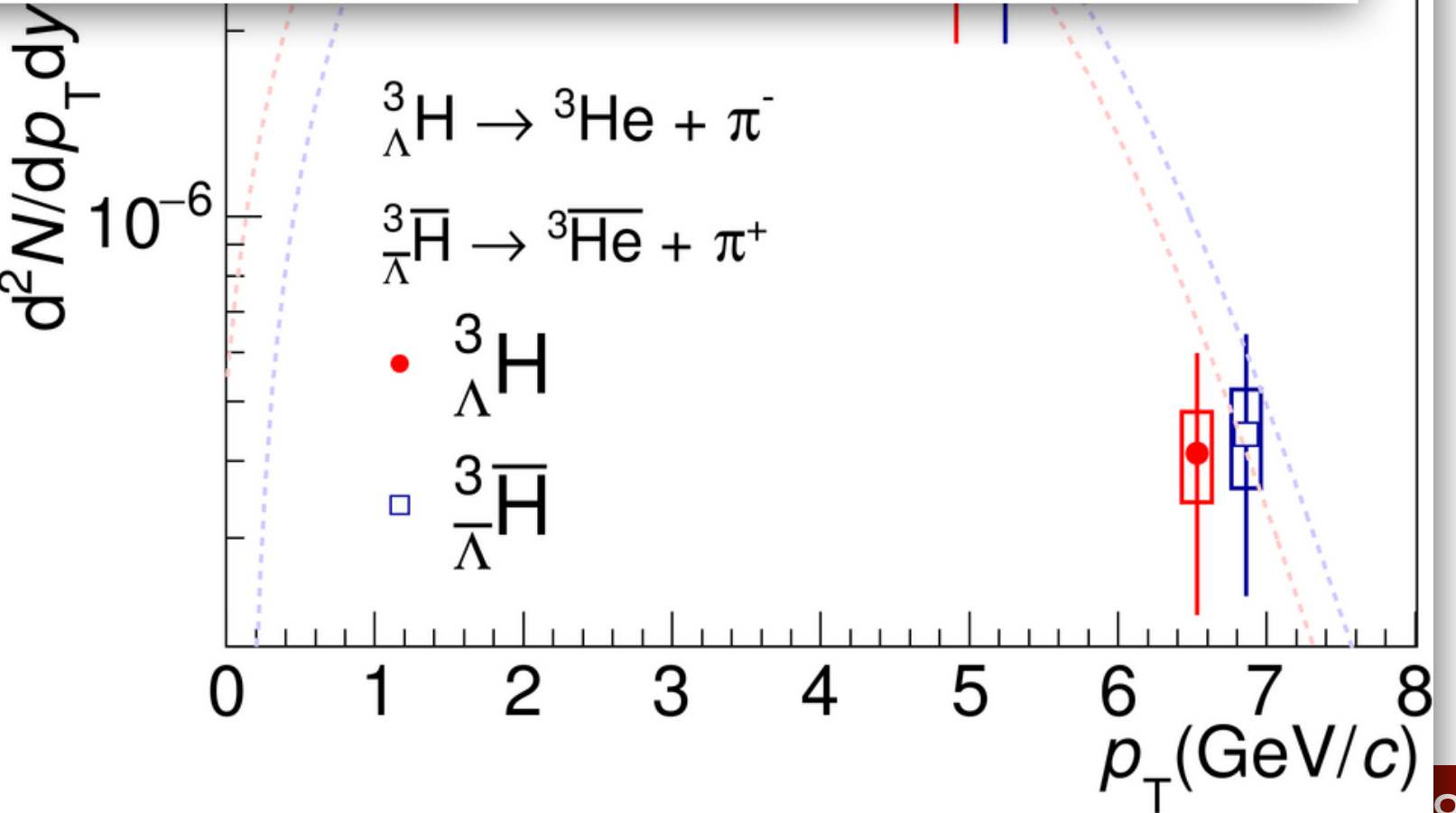
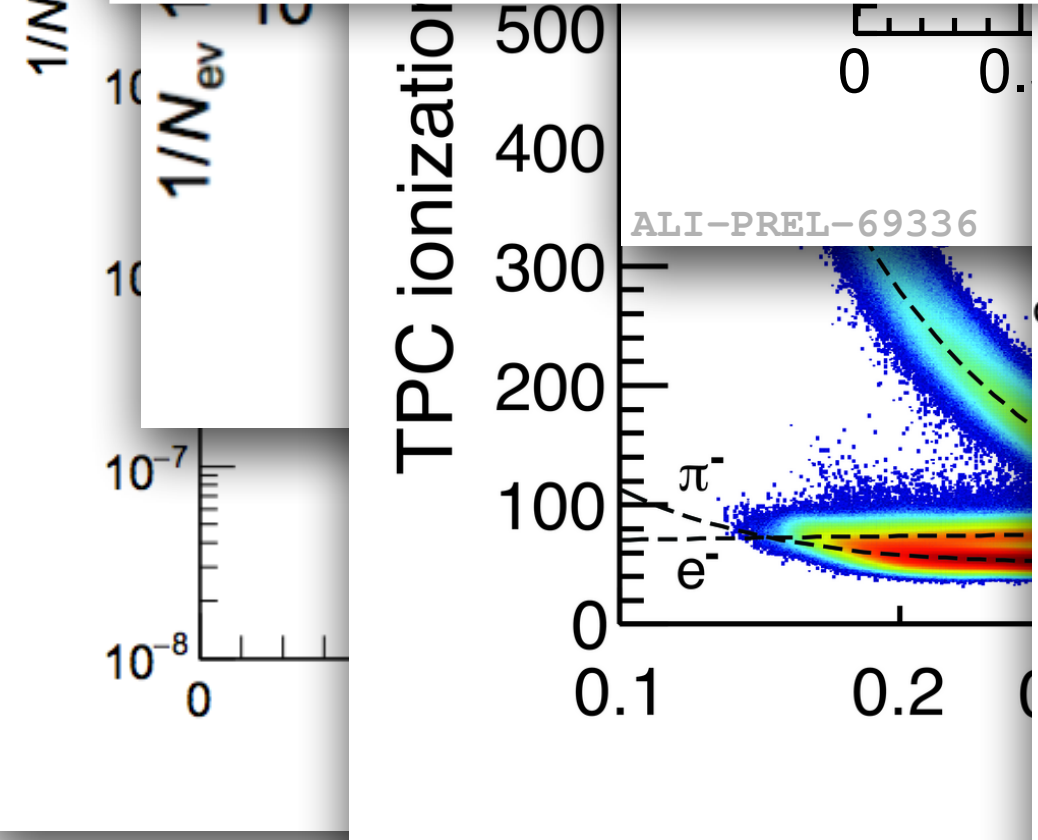
ALI-PERF-36713



ALI-PERF-36713



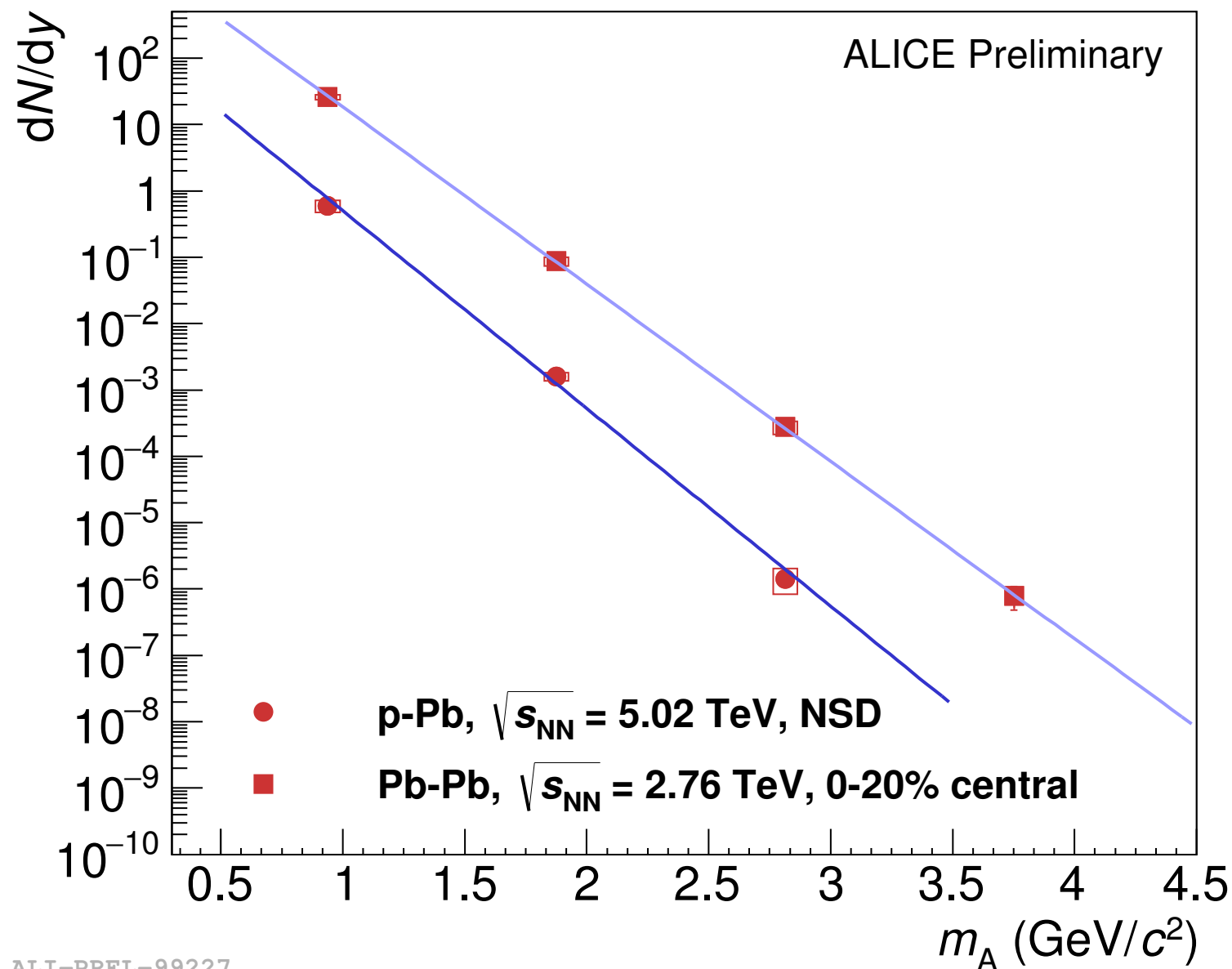
Huge stack of data! But what does it tell about the theory?



ALI-PERF-36713

Mass scaling in nuclei production

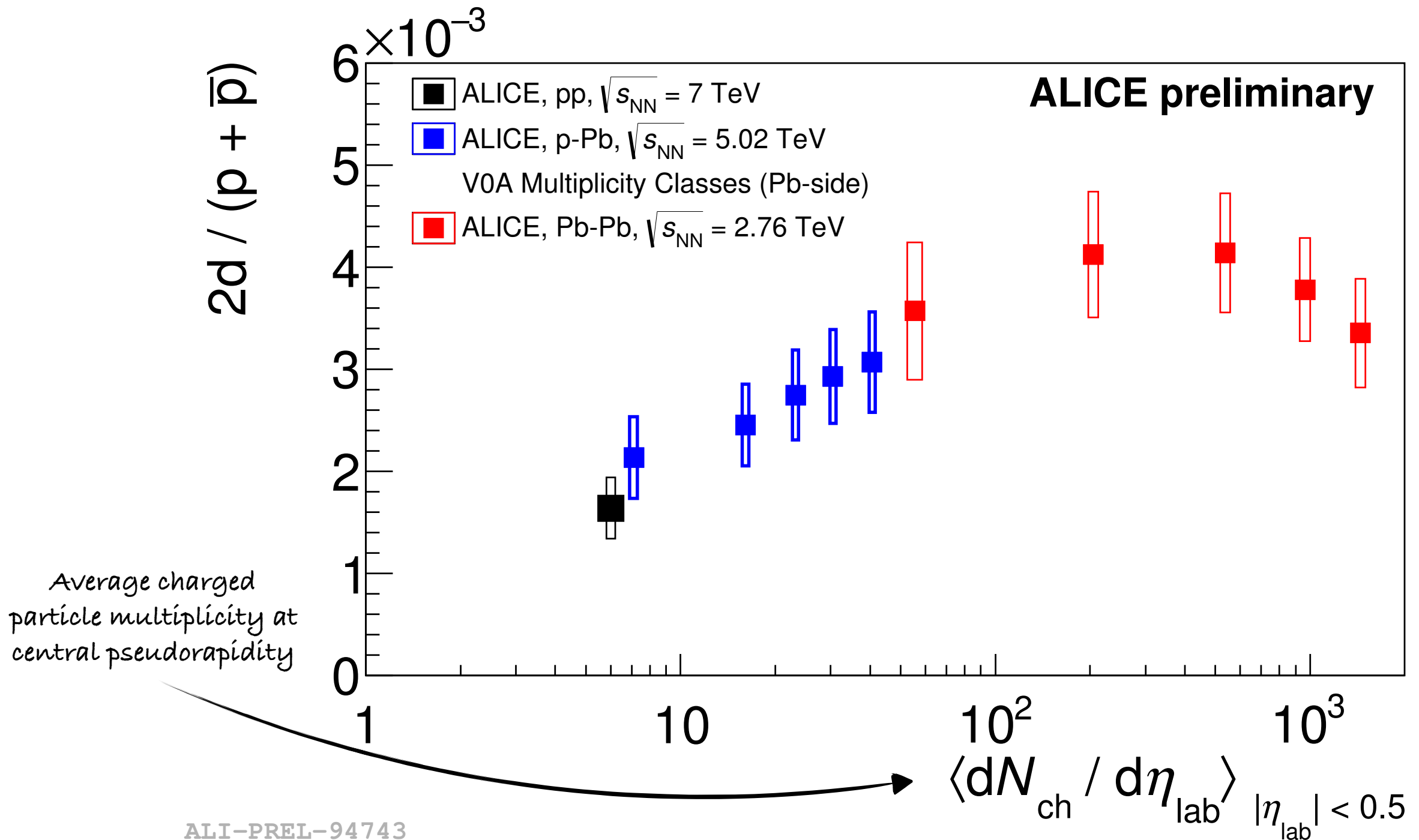
Nuclei production yields follow an exponential decrease with mass as predicted by the thermal model



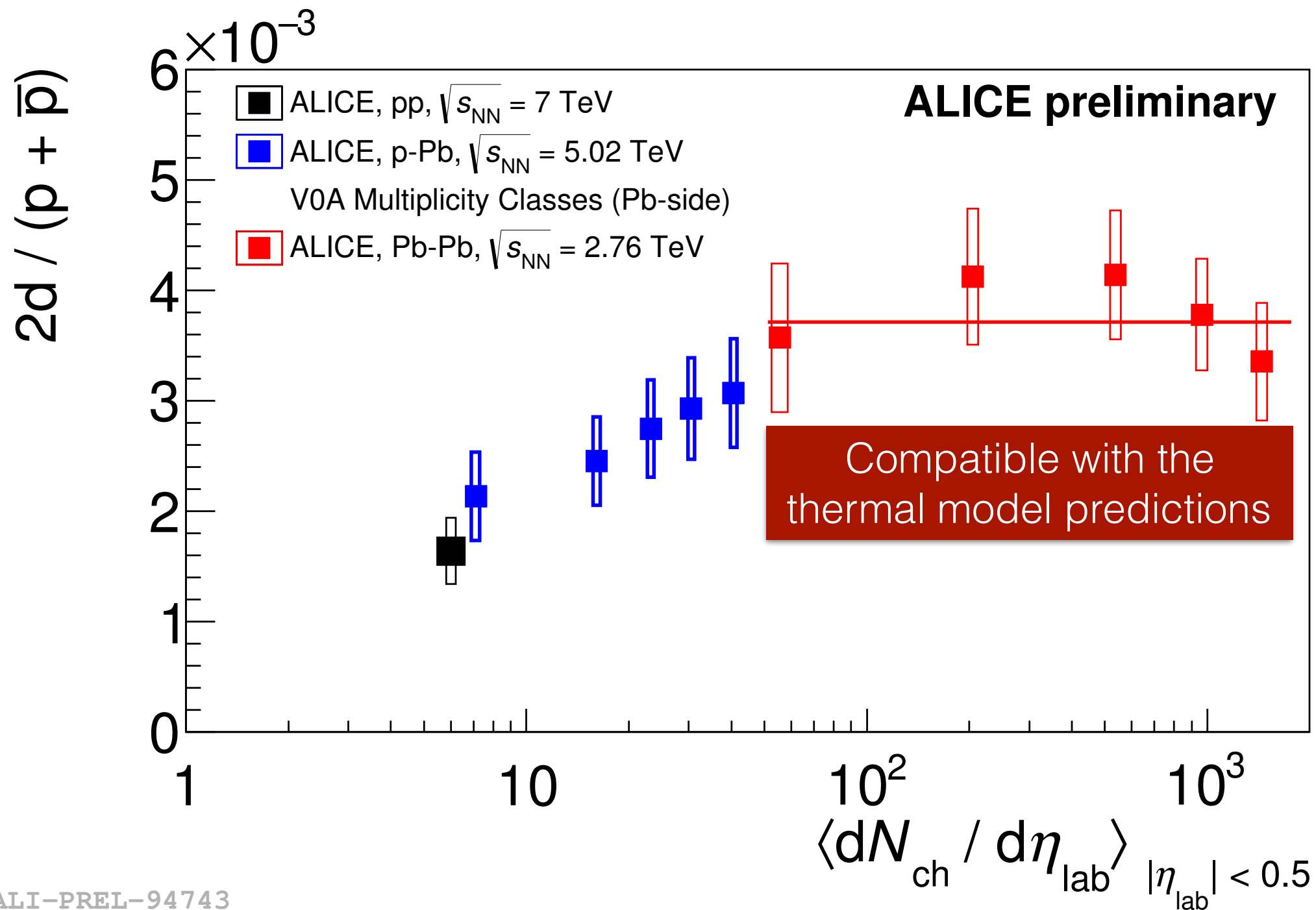
ALI-PREL-99227

The penalty factor for adding one baryon is in Pb-Pb ~ 300 and for p-Pb ~ 600

Deuteron over proton ratio

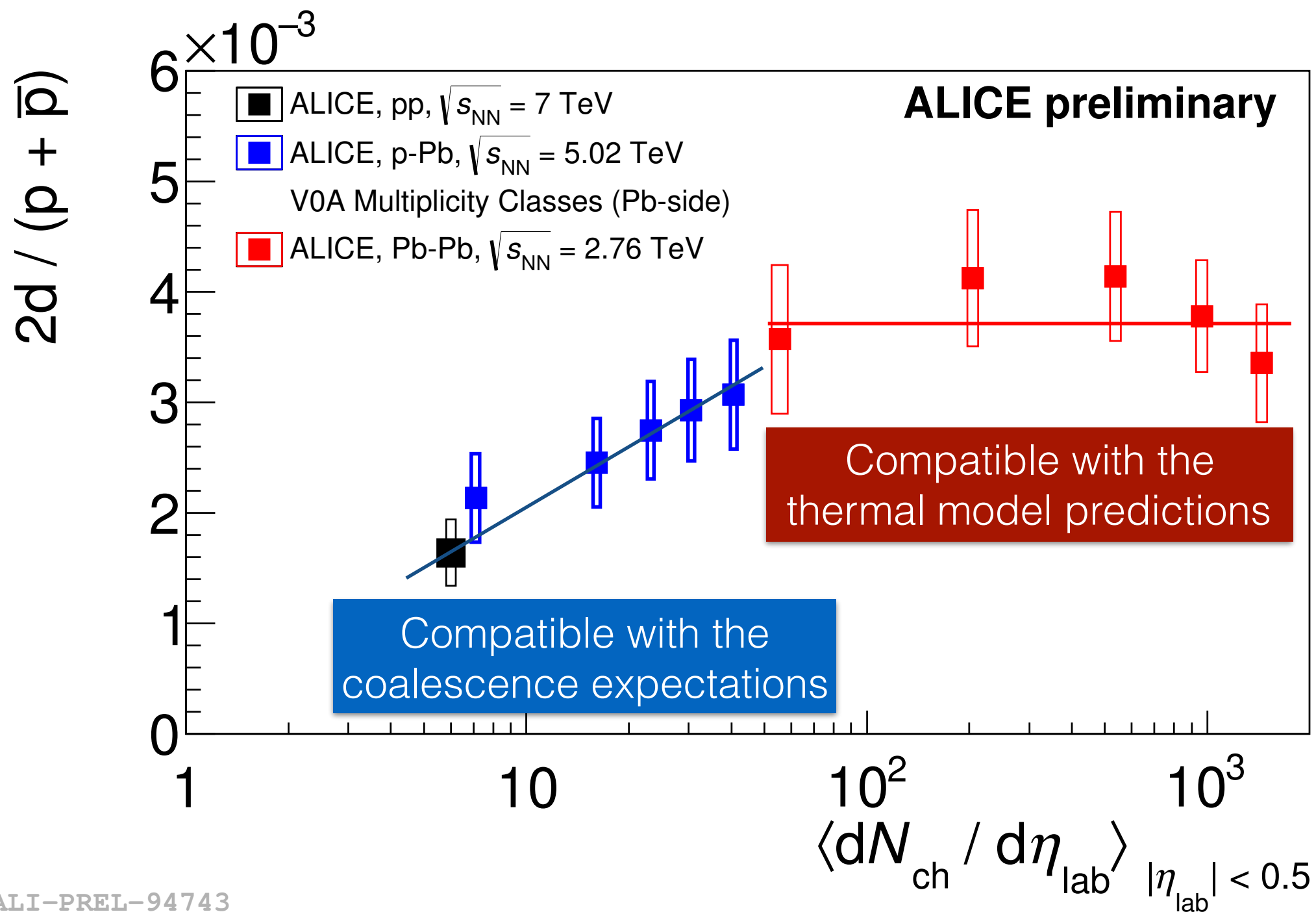


Deuteron over proton ratio



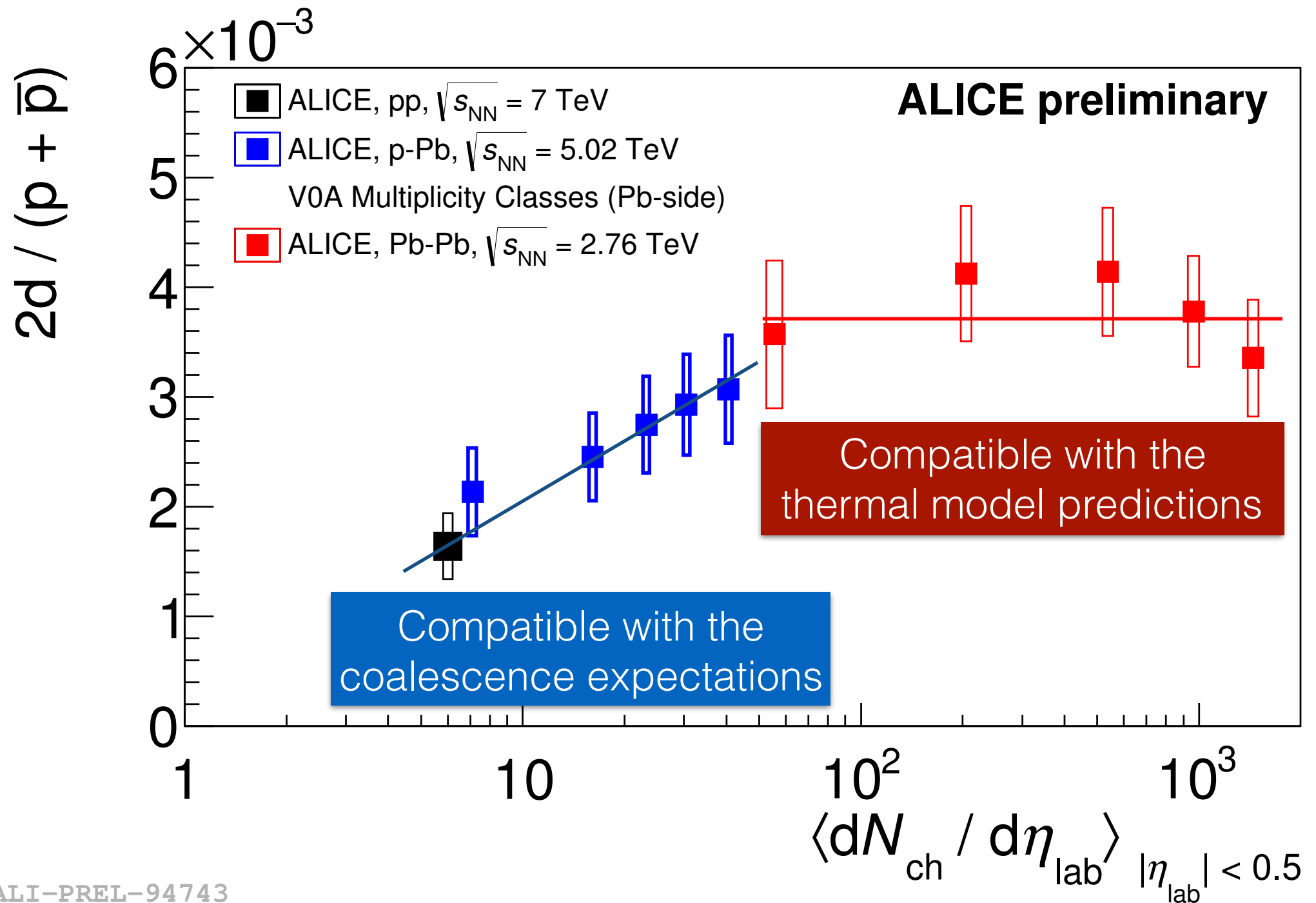
ALI-PREL-94743

Deuteron over proton ratio



ALI-PREL-94743

Deuteron over proton ratio

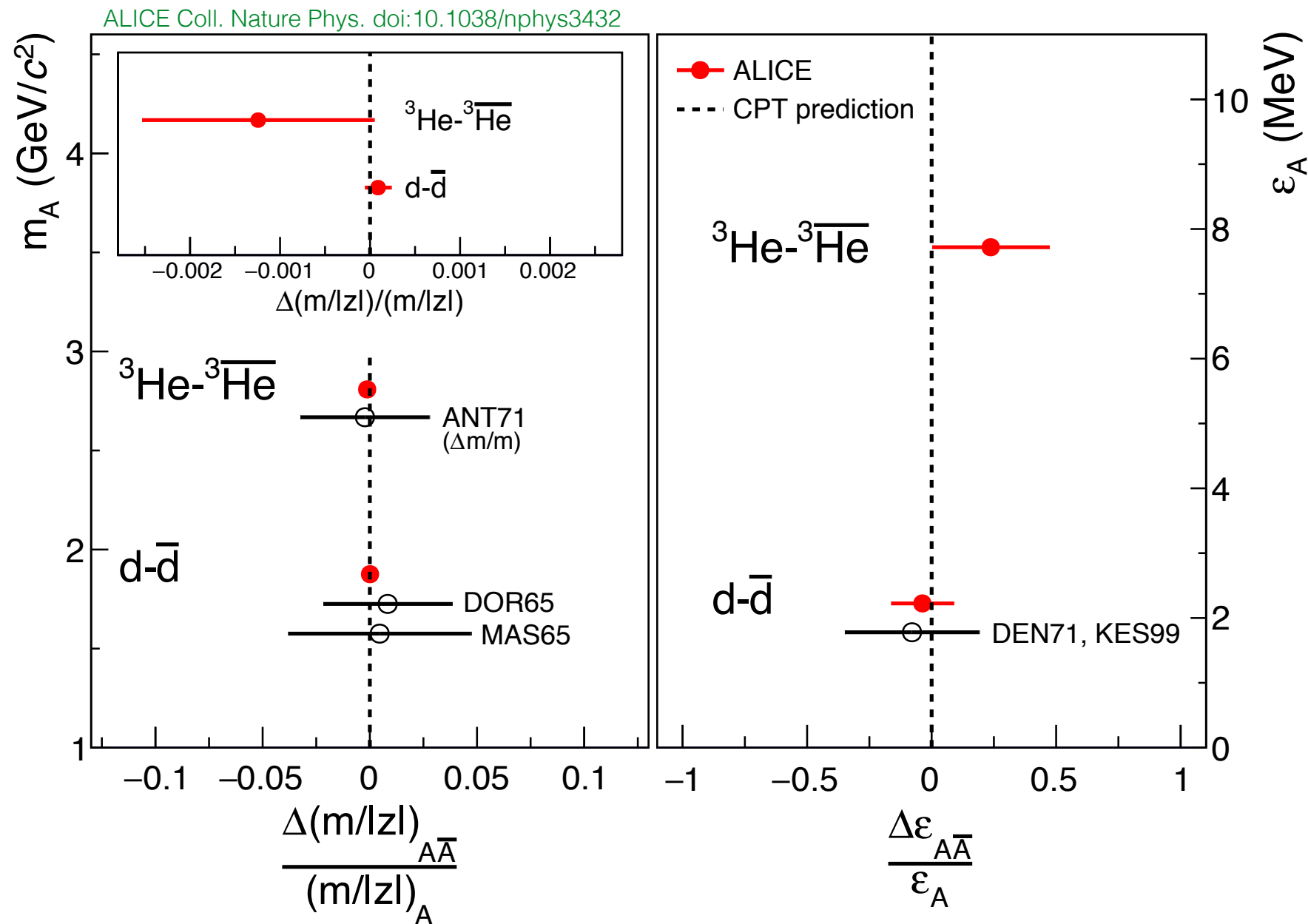


ALI-PREL-94743

Simple coalescence seems to work in the small systems while thermal models describe better the data in Pb-Pb

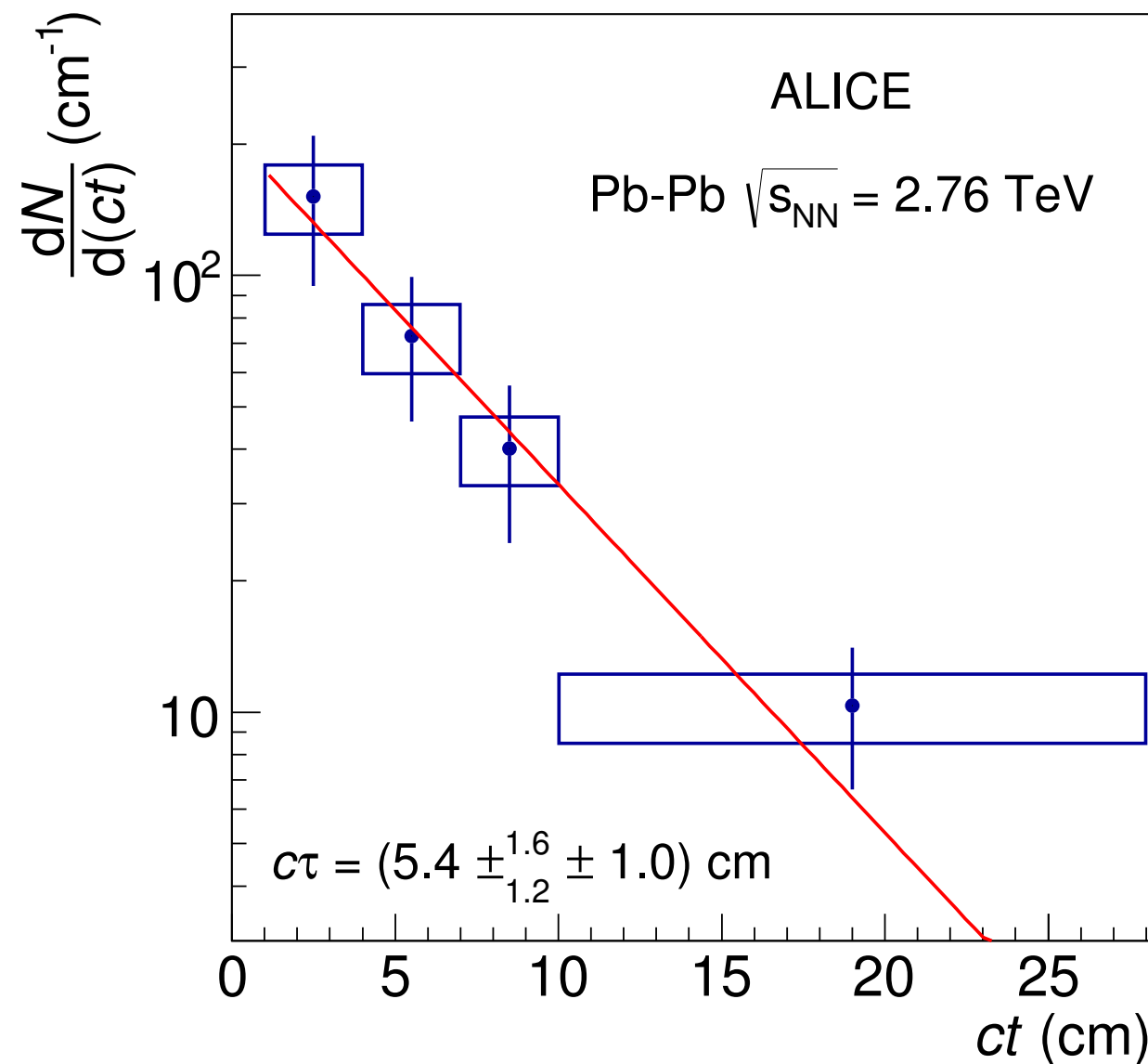
CPT tests in the nuclei sector!

CPT invariance: no mass difference between nuclei and anti-nuclei



This test shows that the **mass of nuclei and anti-nuclei are compatible** within the uncertainties. The binding energies are compatible in nuclei and anti-nuclei as well.

Open puzzle: hyper-triton lifetime



From the exponential fit to the differential yield in different ct bins it is possible to extract the lifetime of the hyper-triton

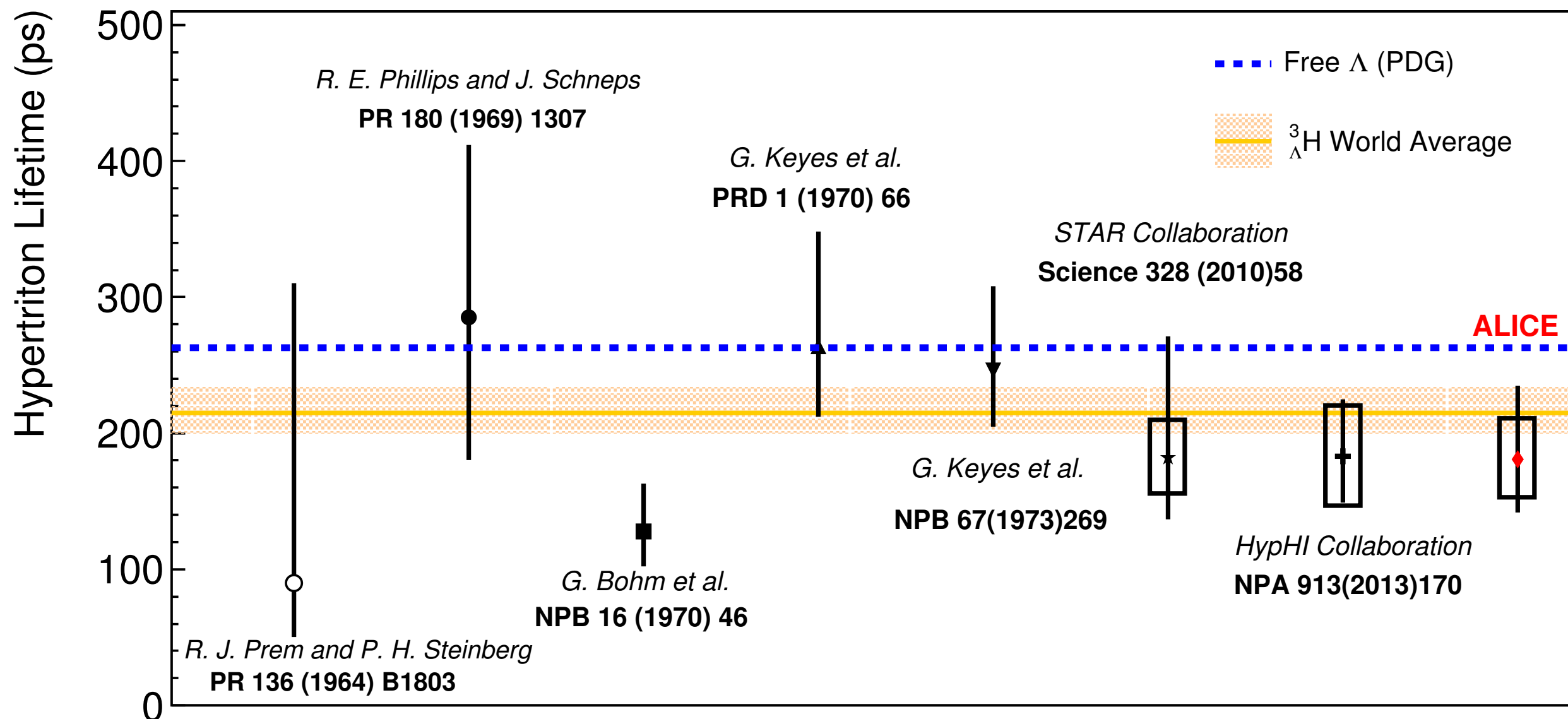
ALICE

$$\tau (181_{-39}^{+54} (\text{stat.}) \pm 33 (\text{syst})) \text{ ps}$$

World average

$$\tau = 215_{-16}^{+18} \text{ ps}$$

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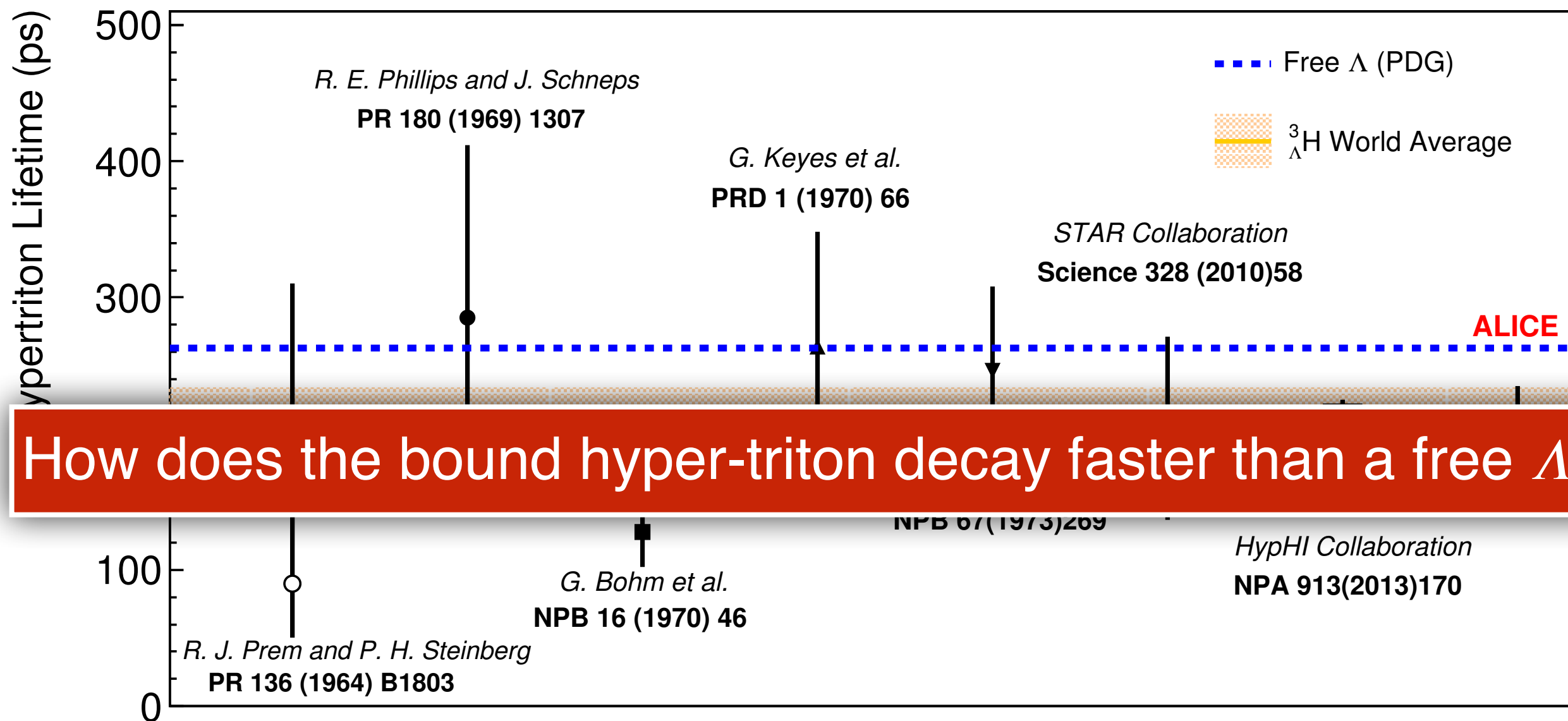
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Open puzzle: hyper-triton lifetime



How does the bound hyper-triton decay faster than a free Λ ?

From the exponential fit to the differential yield in different ct bins it is possible to extract the lifetime of the hyper-triton

ALICE

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World average

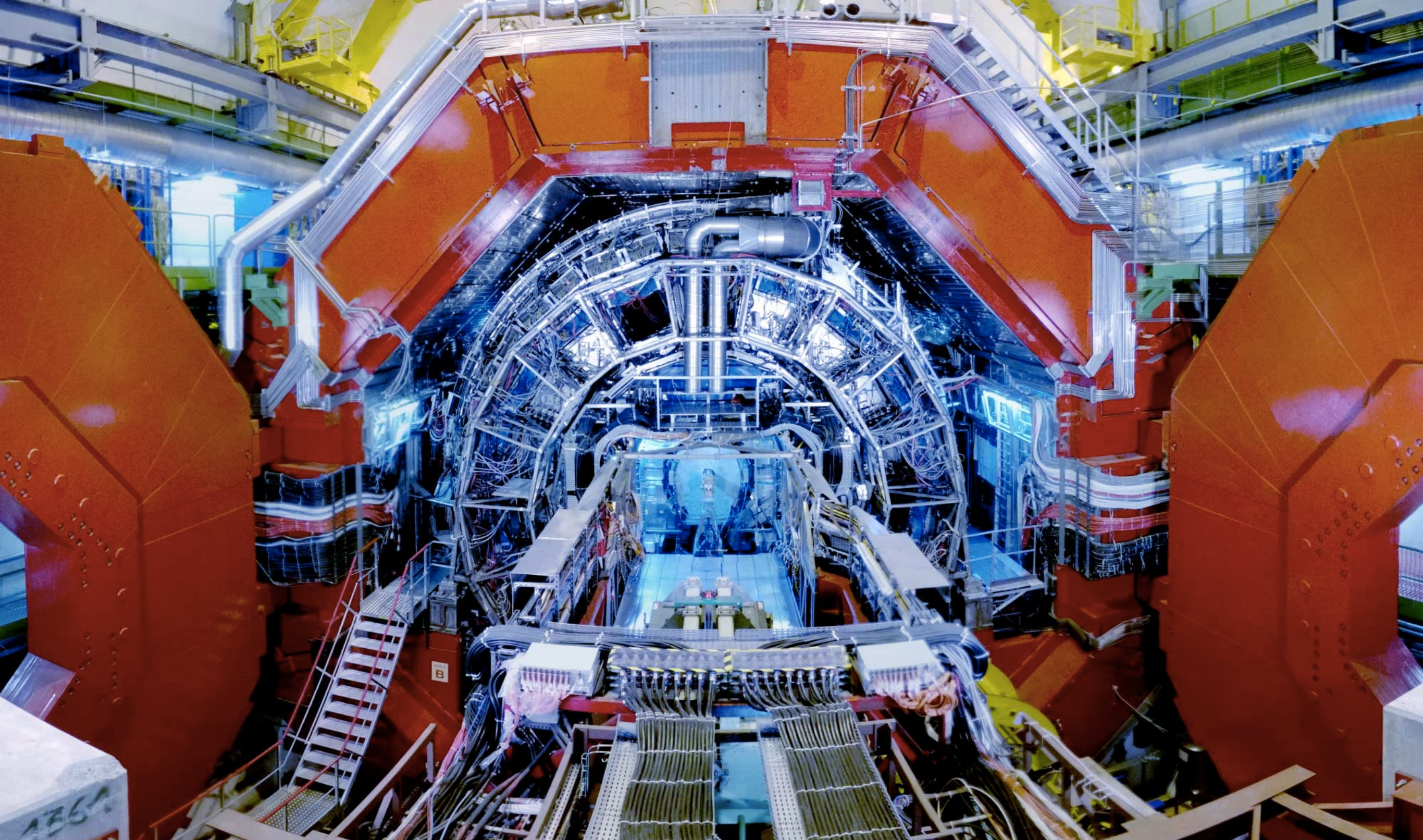
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Conclusions

The study of (hyper-)(anti-)nuclei at the LHC allows one to investigate:

- Coalescence model and thermal models as production mechanisms
- CPT invariance
- Hadronic phase of the QGP
 - How can loosely bound objects survive such conditions?
- Outstanding puzzle of the hyper triton lifetime

Data and results from Run2 of the LHC will help to answer the open questions.



Backup slides



Collision systems at the LHC

	<p>pp</p>	\sqrt{s} 0.9 TeV 2.76 TeV 5.02 TeV 7 TeV 8 TeV 13 TeV	<p>Reference for measurement in other systems</p>
	<p>p-Pb Pb-p</p>	$\sqrt{s_{NN}}$ 5.02 TeV 5.02 TeV	<p>Study of nuclear matter effects</p>
	<p>Pb-Pb</p>	$\sqrt{s_{NN}}$ 2.76 TeV 5.02 TeV	<p>Study of hot and dense QCD matter</p>

Three collision systems: unique opportunity to further study hadronisation and the strong interaction at extreme regimes of energy density and temperature

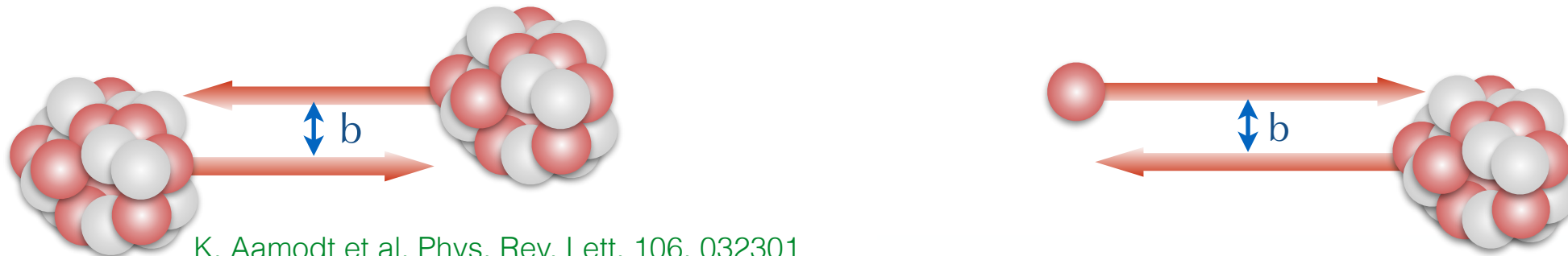
Centrality of a collision

The centrality of a collision is defined by the impact parameter b :

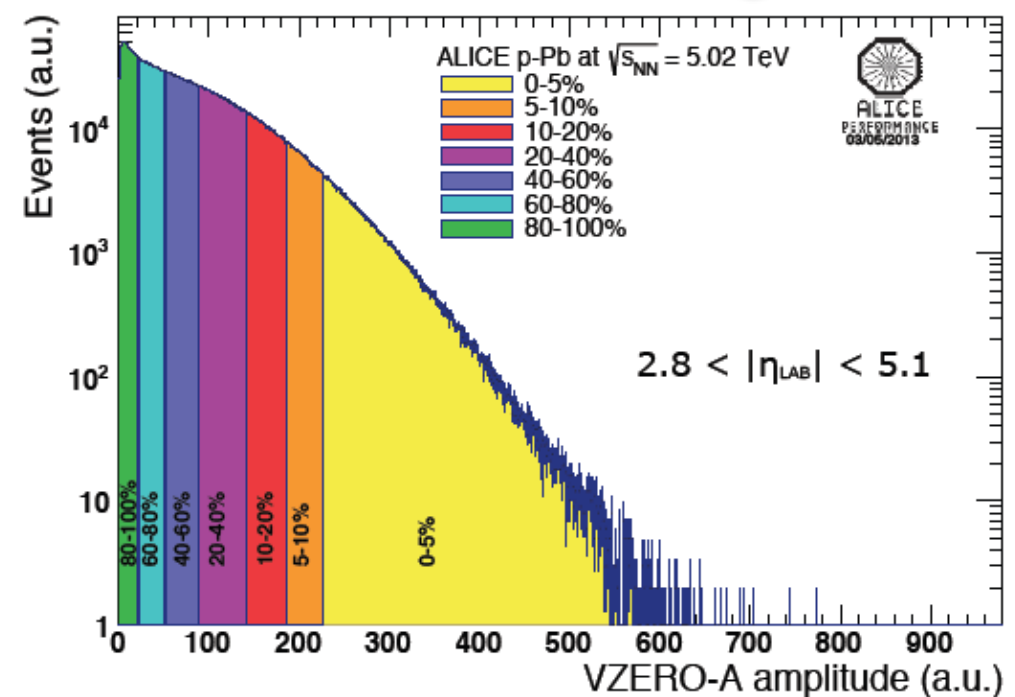
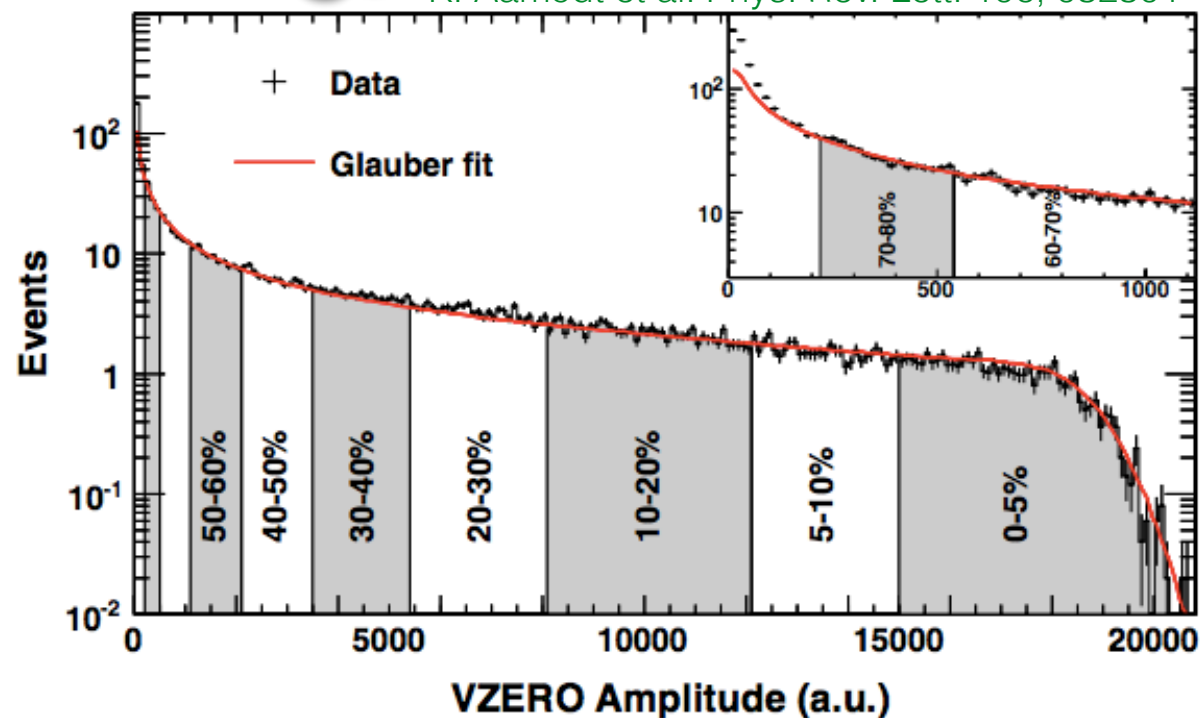
Most central collision \Leftrightarrow *Smallest b*

Experimentally it is possible to correlate the charged particle multiplicity to b by fitting data with the function shape predicted by the Glauber model.

The correlation between charged particle mult. and impact parameter in p-Pb is broader.



K. Aamodt et al. Phys. Rev. Lett. 106, 032301



V-ZERO is a scintillator hodoscope used for centrality estimation and for the trigger.

Coalescence parameter

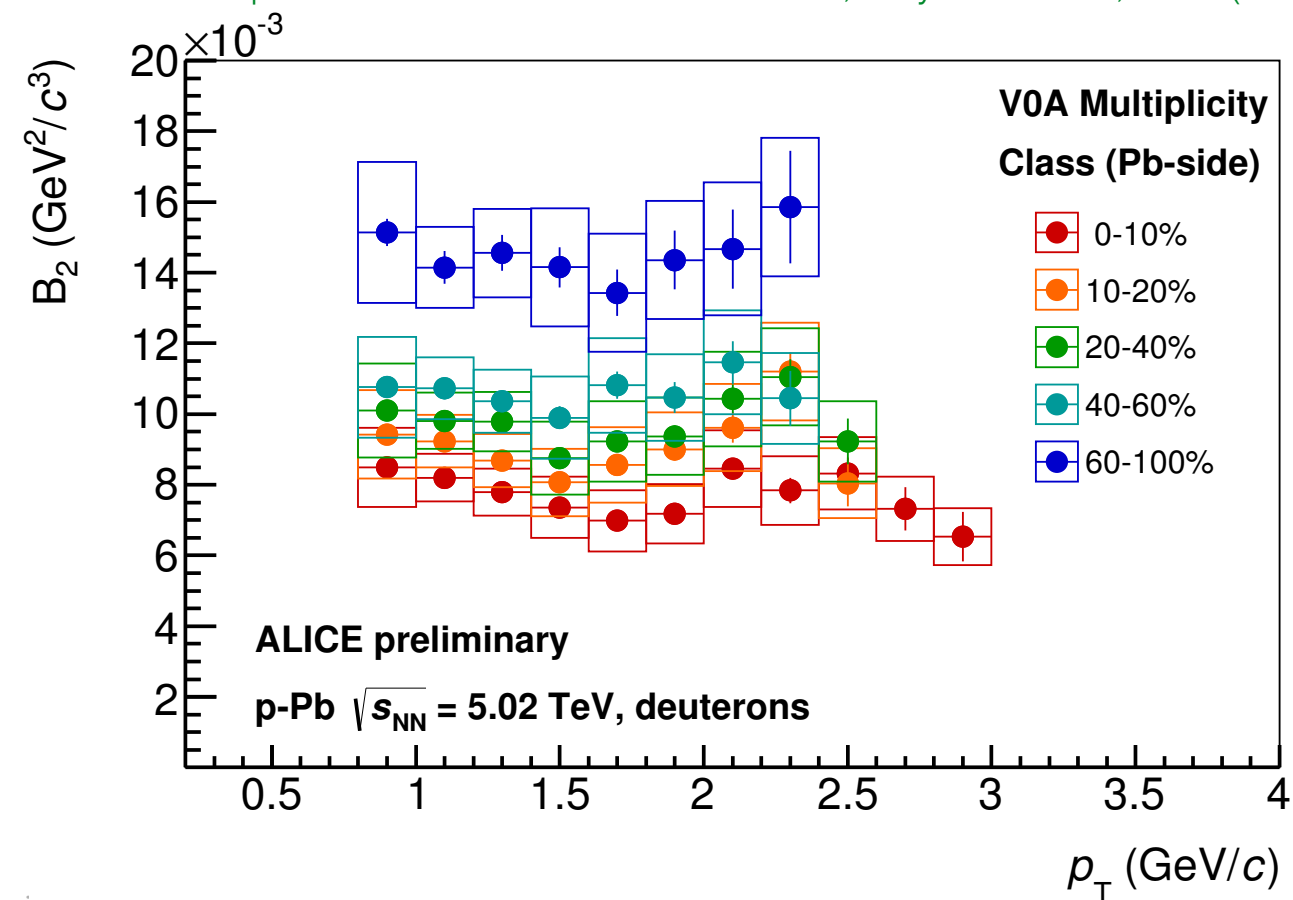
The **coalescence parameter**, defined as:

$$B_2 = \frac{E_d \frac{d^3 N_d}{dp_p^3}}{\left(E_p \frac{d^3 N_p}{dp_p^3} \right)^2}$$

is predicted to be p_T independent by the simplest formulation of coalescence model.

This is observed in p-Pb collisions.

B_2 parameter: R. Scheibl and U. Heinz, Phys.Rev. C59, 1585 (1999)



Coalescence parameter

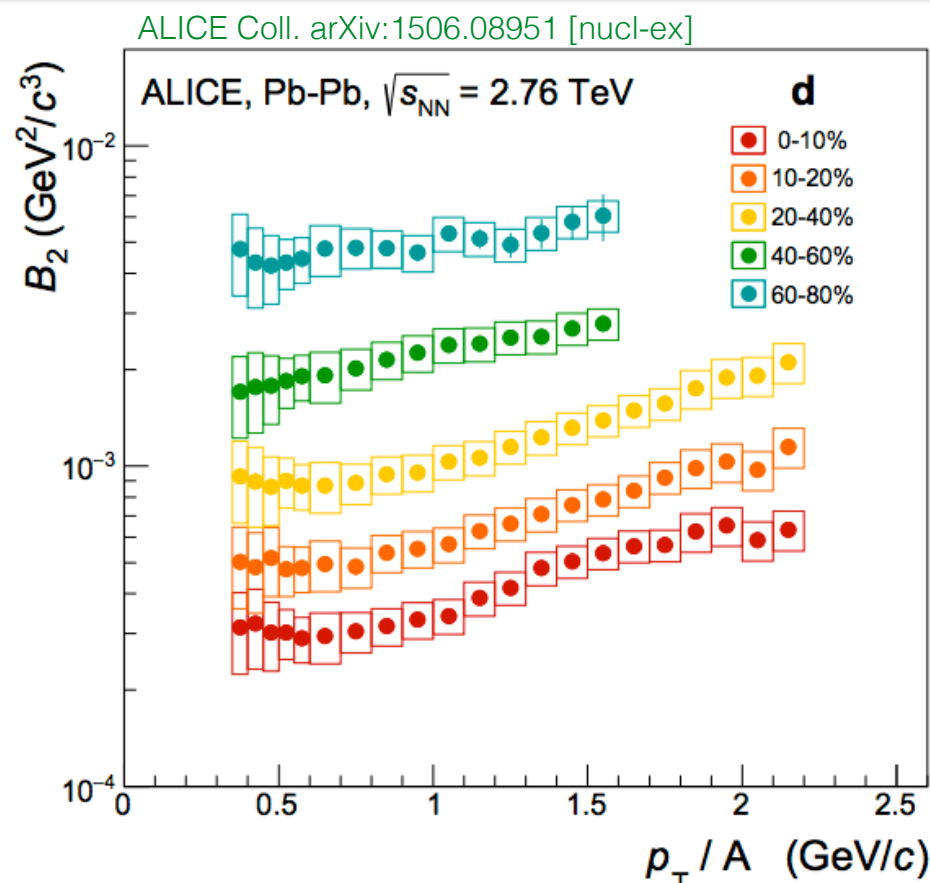
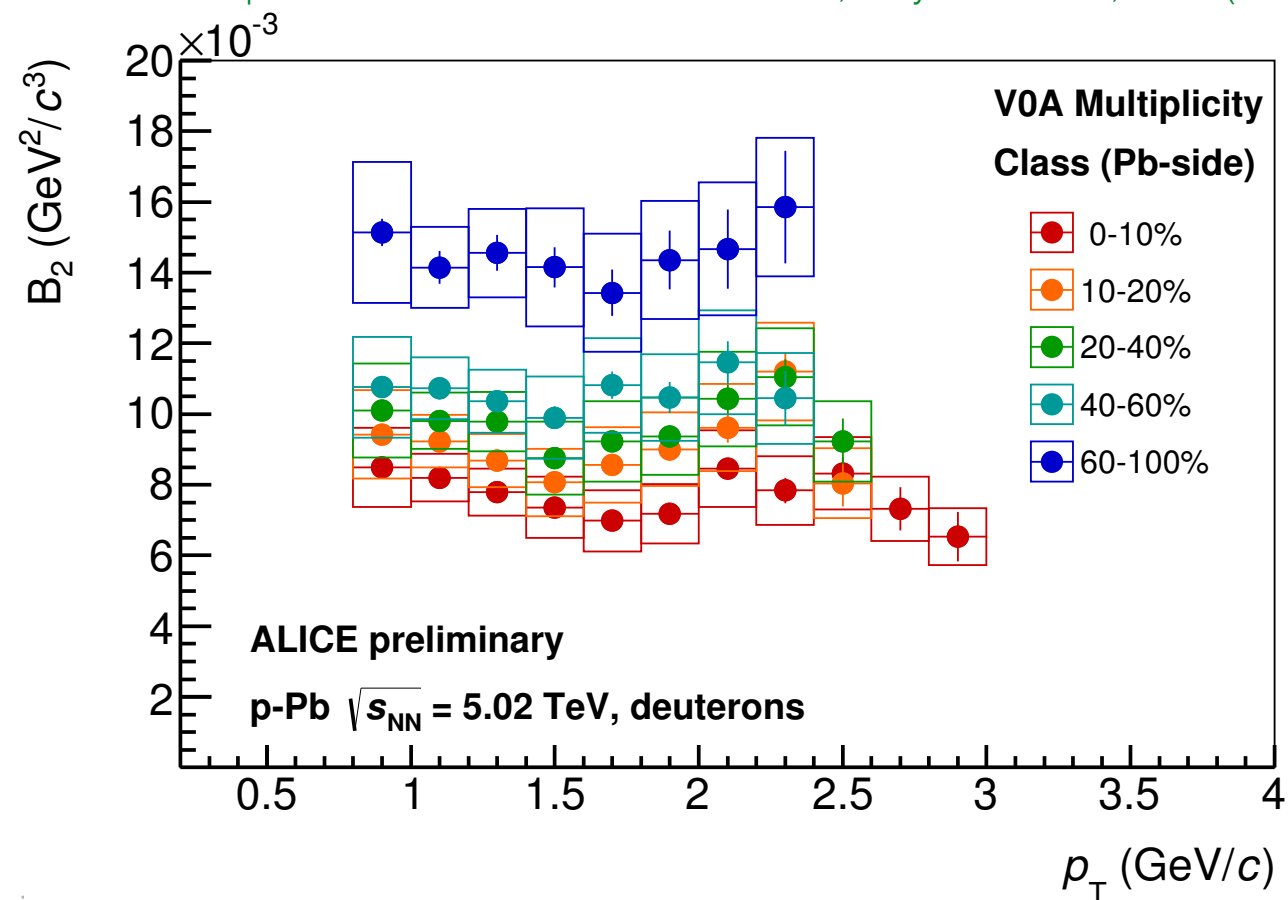
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The coalescence parameter gets smaller as the events are more central. This is due to the increasing size of the emitting source.

The B_2 depends on p_T in central events, which could be explained by looking at the Hanbury Brown and Twiss (HBT) radii dependence of the B_2 .

Coalescence and HBT radii

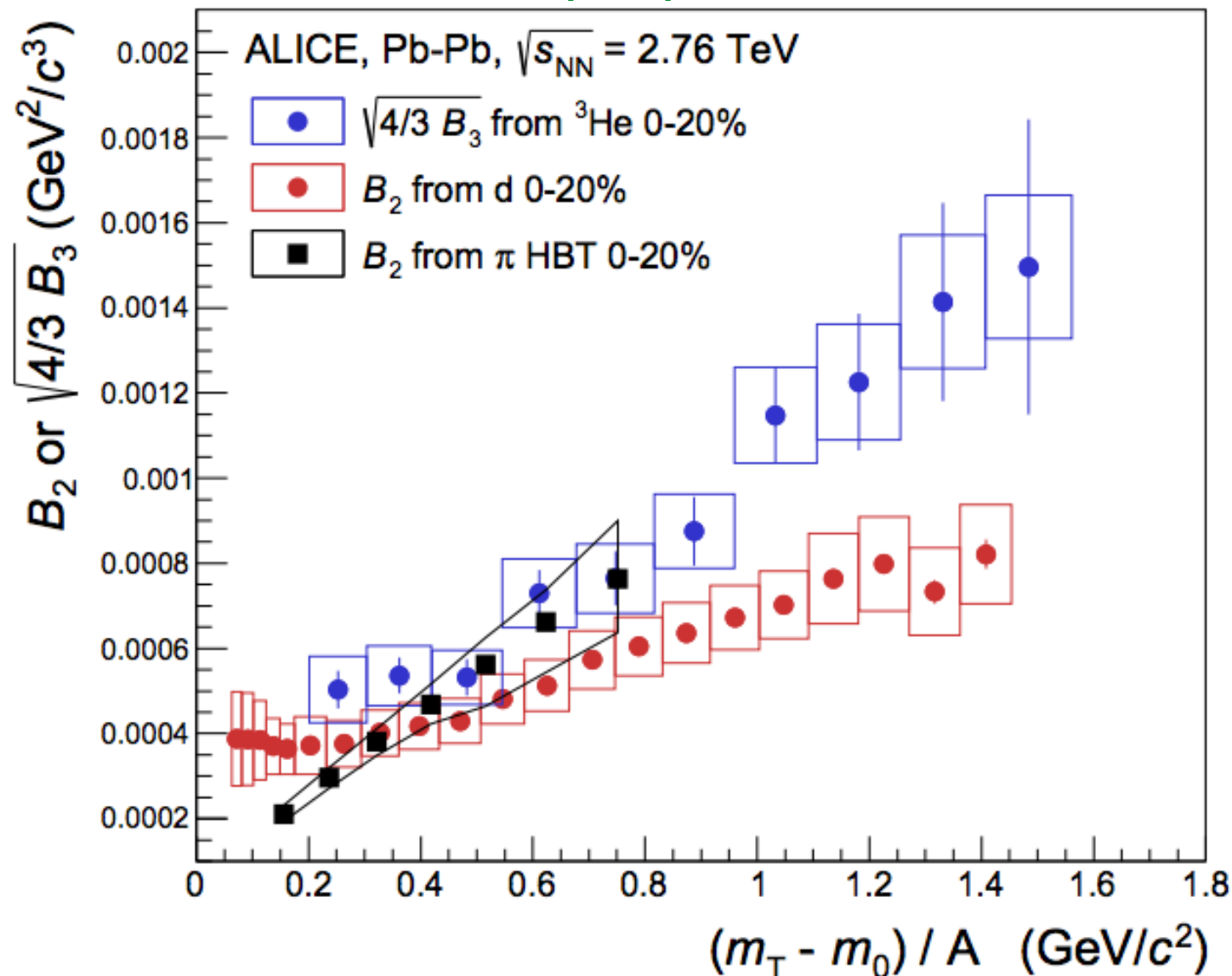
The coalescence parameter can be expressed as a function of the HBT radii:

$$B_2 = \frac{3\pi^{3/2} \langle C_d \rangle}{2m_T R_{\perp}^2(m_T) R_{\parallel}(m_T)}$$

R. Scheibl and U. Heinz, Phys.Rev. C59, 1585 (1999)

A rough agreement is found in terms of magnitude and the dependence on p_T .

ALICE Coll. arXiv:1506.08951 [nucl-ex]



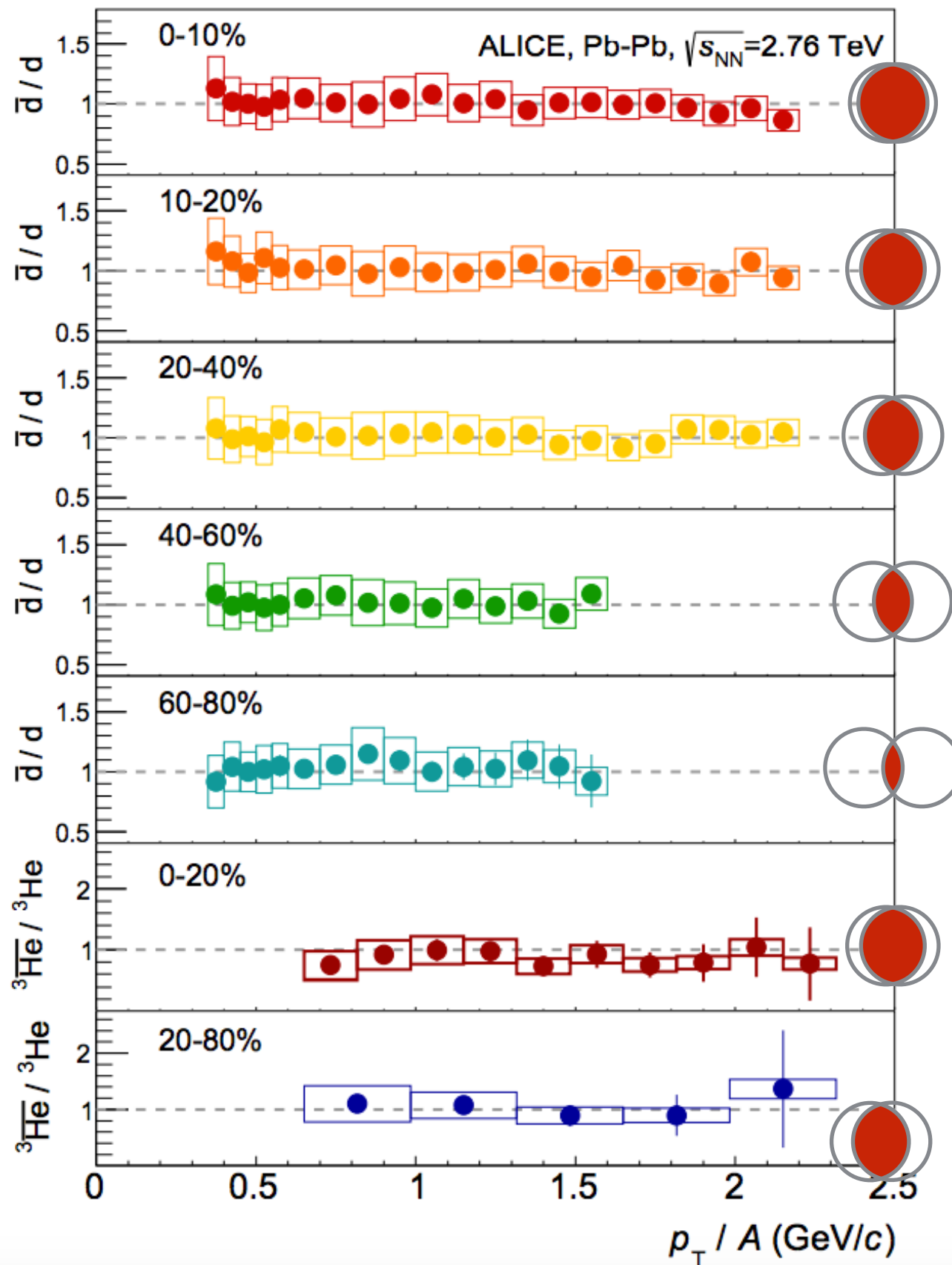
The coalescence parameter for a nucleus i with A nucleons is defined as:

$$E_i \frac{d^3 N_i}{dp_i^3} = B_A \left(E_p \frac{d^3 N_p}{dp_p^3} \right)^A$$

From this it is possible to derive a recursive formula. For instance for the ${}^3\text{He}$:

$$B_3 = B_2^2 \left(\frac{M_{3He} \cdot m}{M_d^2} \right) \approx \frac{3}{4} B_2^2$$

Ratio matter/ antimatter



The ratio nuclei / anti-nuclei is compatible with one at mid-rapidity



The same ratio is seen for other particle species measured at the LHC

A large fraction of the systematic uncertainties on the determination of the ratios is due to the limited knowledge of the cross sections of anti-nuclei interacting with the material of the detector.

Combined Blast Wave fit

BW model fit:

- gives insight into the kinetic freeze-out conditions
- does not describe hard processes that contribute to particle production at high p_T

Fit parameters

$$\langle\beta\rangle = 0.632 \pm 0.01$$

$$T_{\text{kin}} = 113 \pm 12 \text{ MeV}$$

$$n = 0.72 \pm 0.03$$

With respect to the fit performed without the nuclei the $\langle\beta\rangle$ decreased while the T_{kin} increased but they are compatible within the uncertainties.

Solid symbols denote the spectra points used for the fit.

