

# **Sensitivity limits of atom interferometry gravity gradiometers and strainmeters**

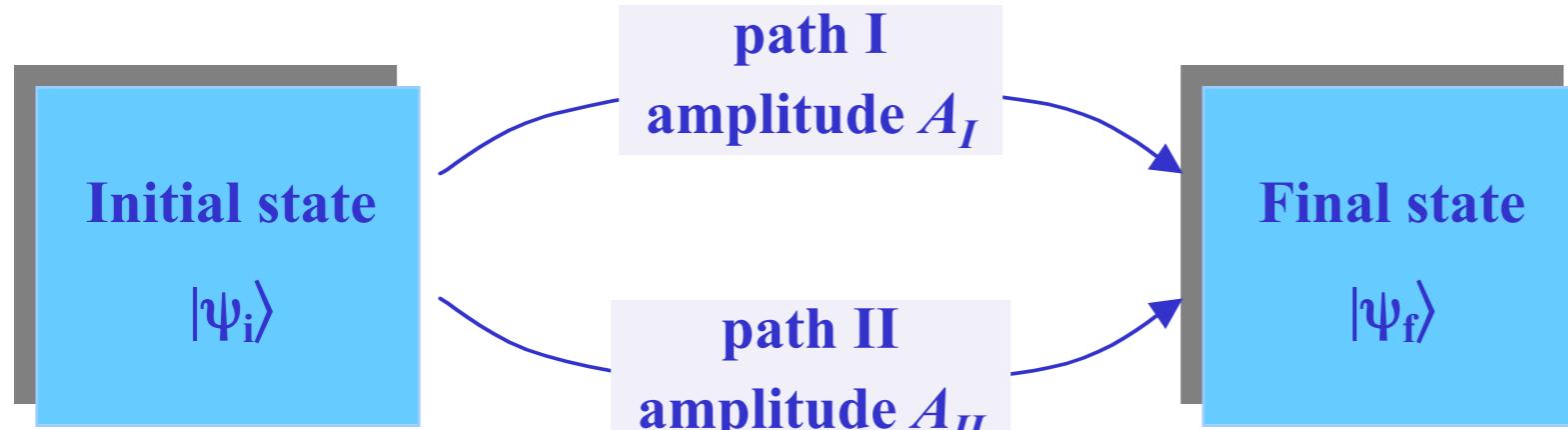
Fiodor Sorrentino  
INFN Genova



# Outline

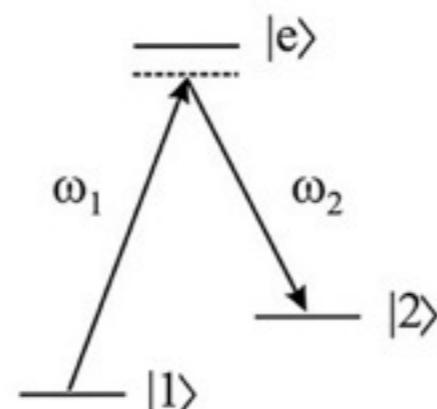
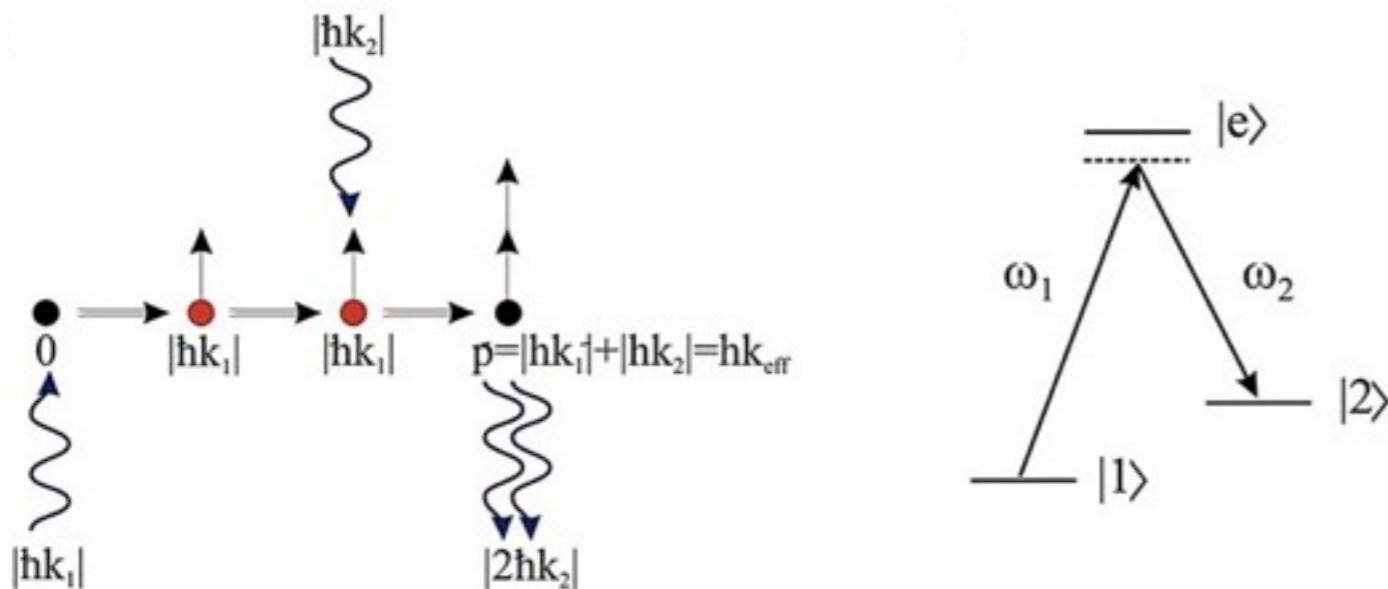
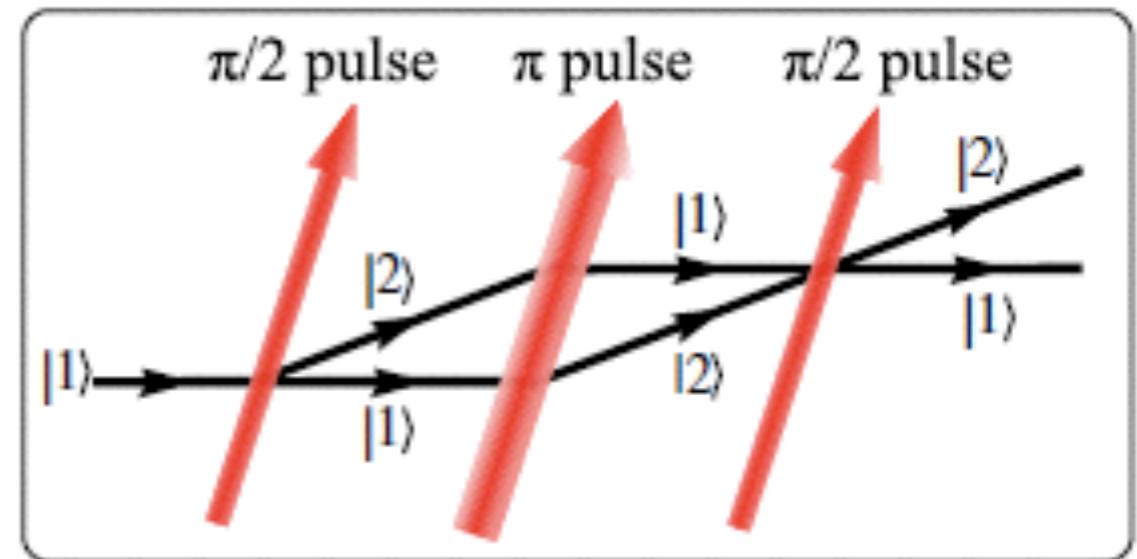
- AI sensors, state of the art performance
- Main noise sources
- Potential improvements
- Technology readiness of AI sensors for field use

# Atom interferometry



Interference of de Broglie amplitudes

Light-pulse beam splitters + fluorescence detection



Output phase selectively sensitive to different effects (inertial, gravitational, external fields, laser phase/frequency, etc) via choice of quantum states

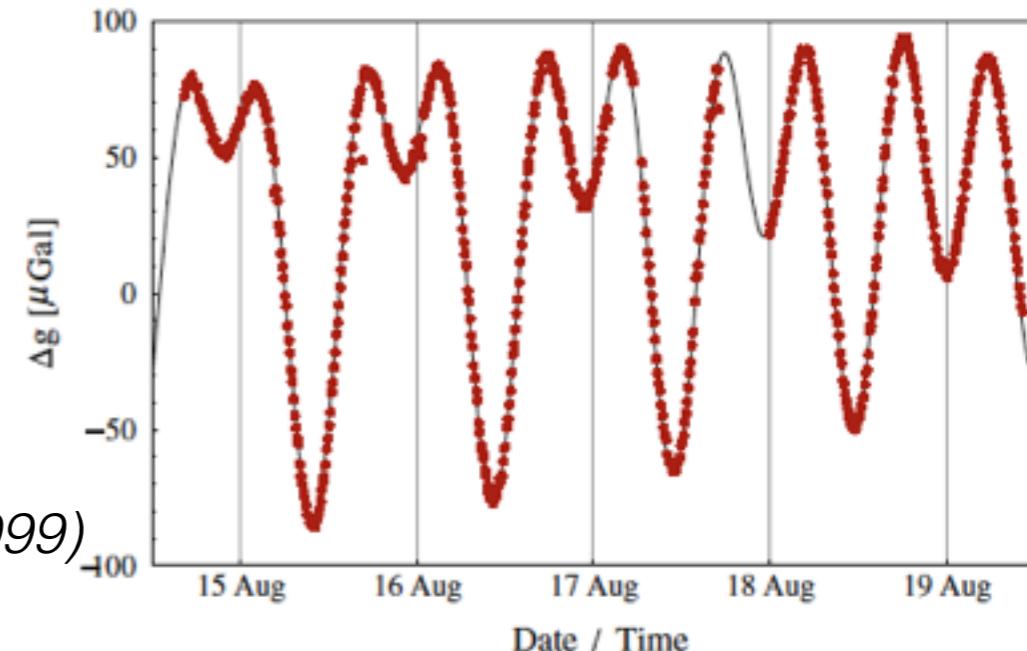
# AI & low frequency GW

- AI inertial sensors feature very good long term stability and control over systematic effects
  - might provide complementary technologies for <Hz GW detectors
- Gravity gradiometers
  - measurement/correction of NN
- Strain meters
  - free fall: vibration isolation without suspensions
  - quantum interaction of light with test mass
    - no radiation pressure noise
  - vertical orientation
- Rotation sensors

# Current performance

- Gravimeters

- resolution:  $3 \times 10^{-9}$  g in 1 second (SYRTE)
- averaging down to  $2 \times 10^{-10}$  g after 30 min (SYRTE)
- accuracy:  $10^{-9} \div 10^{-10}$  g, limited by tidal models
  - A. Peters, K.Y. Chung and S. Chu, *Nature* 400, 849 (1999)
  - H. Müller et al., *Phys. Rev. Lett.* 100, 031101 (2008)
  - M. Hauth et al., *Appl. Phys. B* 113, 49 (2013)
  - P. Gillot et al., *Metrologia* 51, L15 (2014)

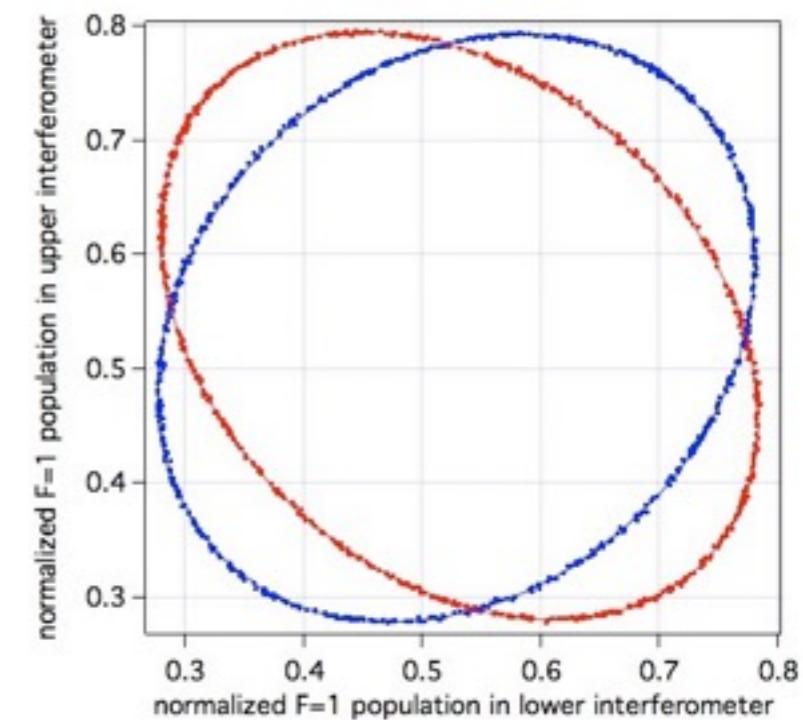


- Gravity gradiometers

- differential acceleration sensitivity:  $5 \times 10^{-9}$  g in 1 s
- $5 \times 10^{-11}$  g after  $10^4$  s
  - F. Sorrentino et al., *Phys. Rev. A* 89, 023607 (2014)

- Rotation sensors

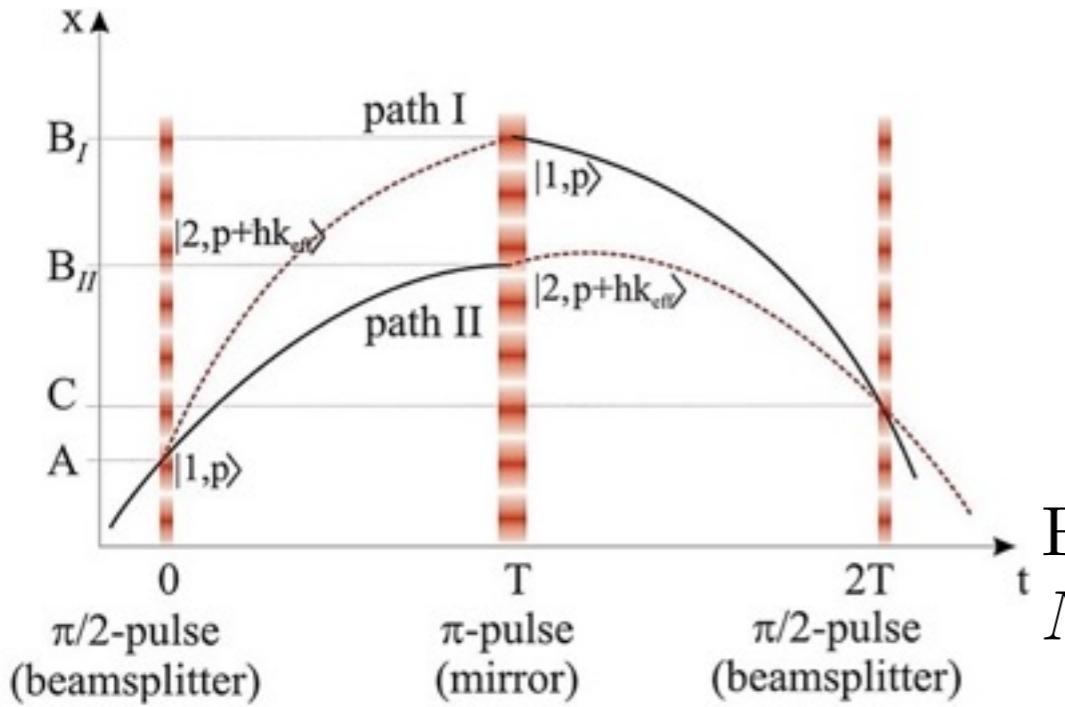
- sensitivity:  $6 \times 10^{-10}$  rad/s/√Hz
- scale factor stability <5 ppm
- bias stability <70 μdeg/h
  - T.L. Gustavson, A. Landragin and M.A. Kasevich, *Class. Quantum Grav.* 17, 2385 (2000)
  - D.S. Durfee, Y.K. Shaham, M.A. Kasevich, *Phys. Rev. Lett.* 97, 240801 (2006)



# Sensitivity limits

- Quantum noise
- Acceleration noise
- Laser wavefront
- Atomic motion
- Laser phase and frequency noise
- Aliasing
- External fields
- Other technical noise sources
- NN

# Quantum projection noise



Phase difference between the paths:

$$\Delta\Phi = k_e[z(0) - 2z(T) + z(1T)] + \Phi_e$$

$$k_e = k_1 - k_2$$

$$\text{with } z(t) = -gt^2/2 + v_0t + z_0 \text{ & } \Phi_e = 0$$

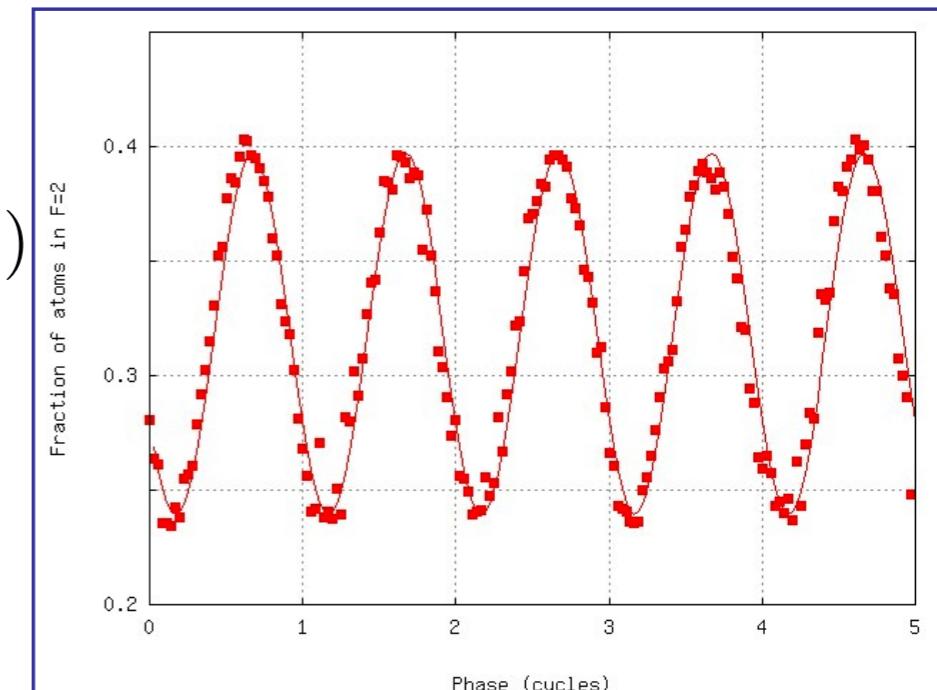
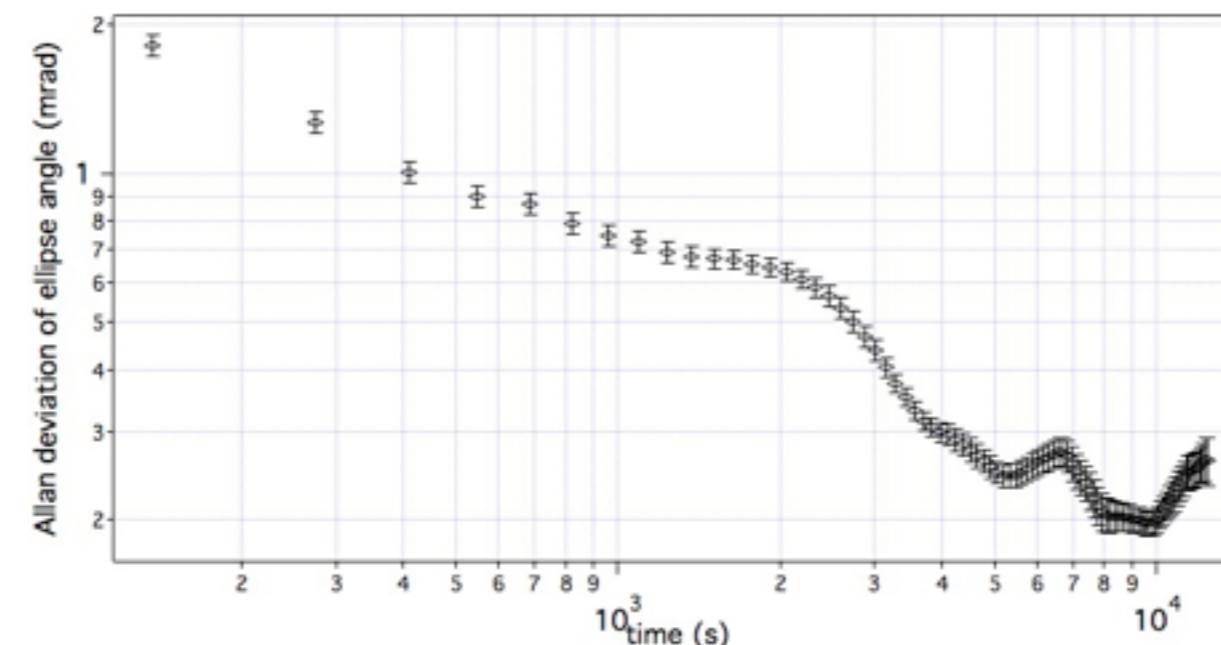
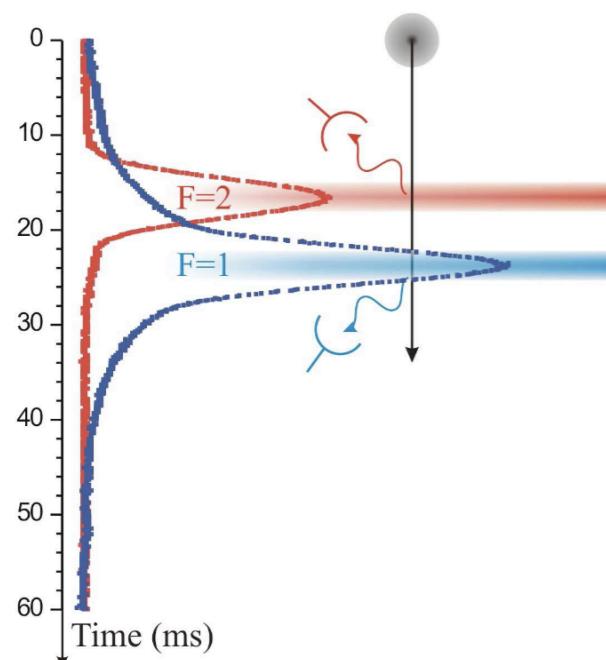
$$\rightarrow \Delta\Phi = k_e g T^2$$

Final population:

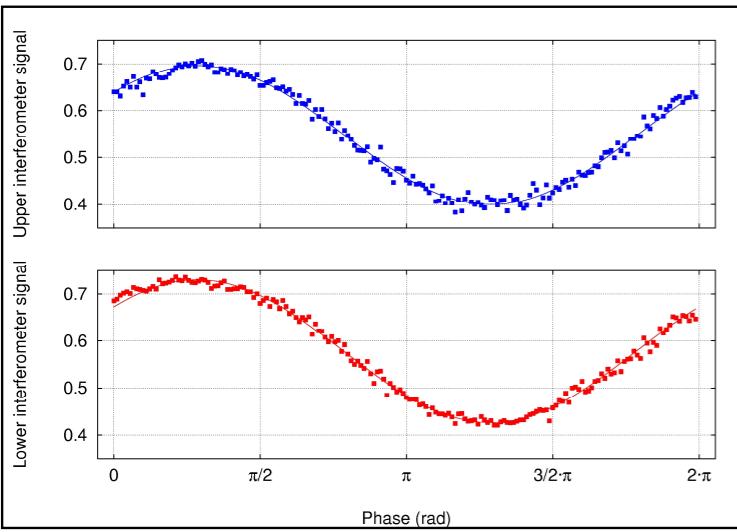
$$N_a = N/2(1 + \cos[\Delta\Phi])$$

$$T = 150 \text{ ms} \rightarrow 2\pi = 10^{-6} \text{ g}$$

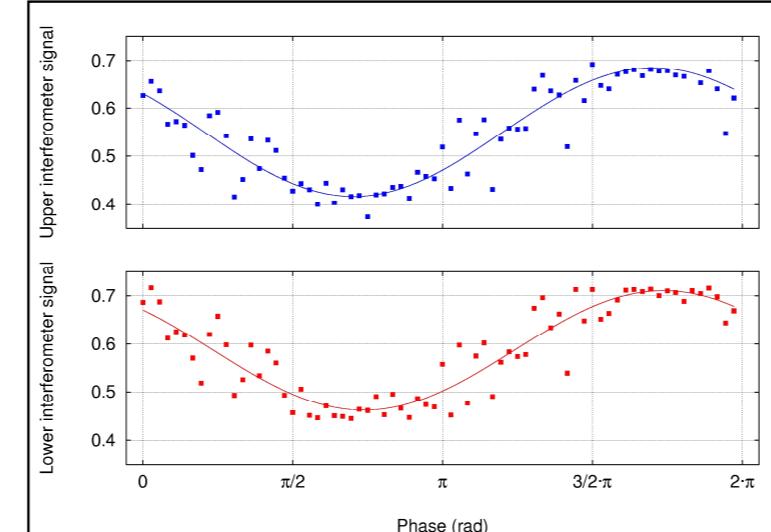
$$\text{S/N}=1000 \rightarrow \text{Sensitivity } 10^{-9} \text{ g/shot}$$



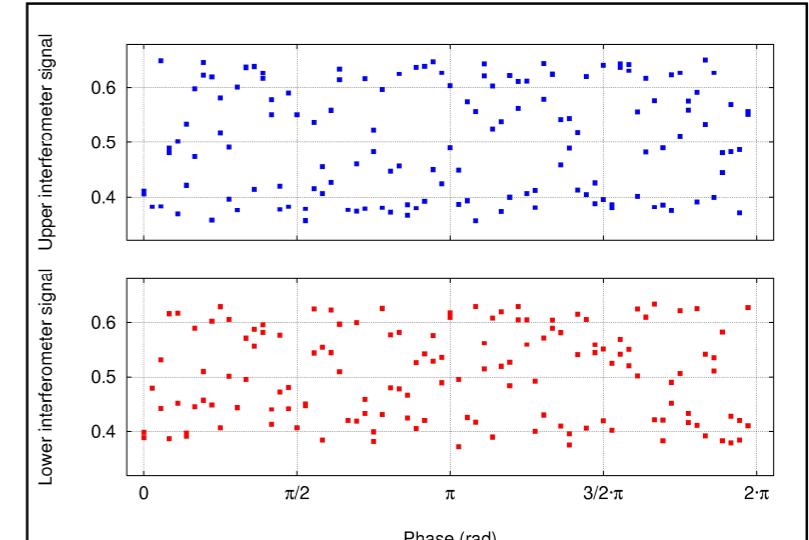
# Acceleration noise



$T = 5 \text{ ms}$   
resol. =  $2.3 \times 10^{-5} \text{ g/shot}$

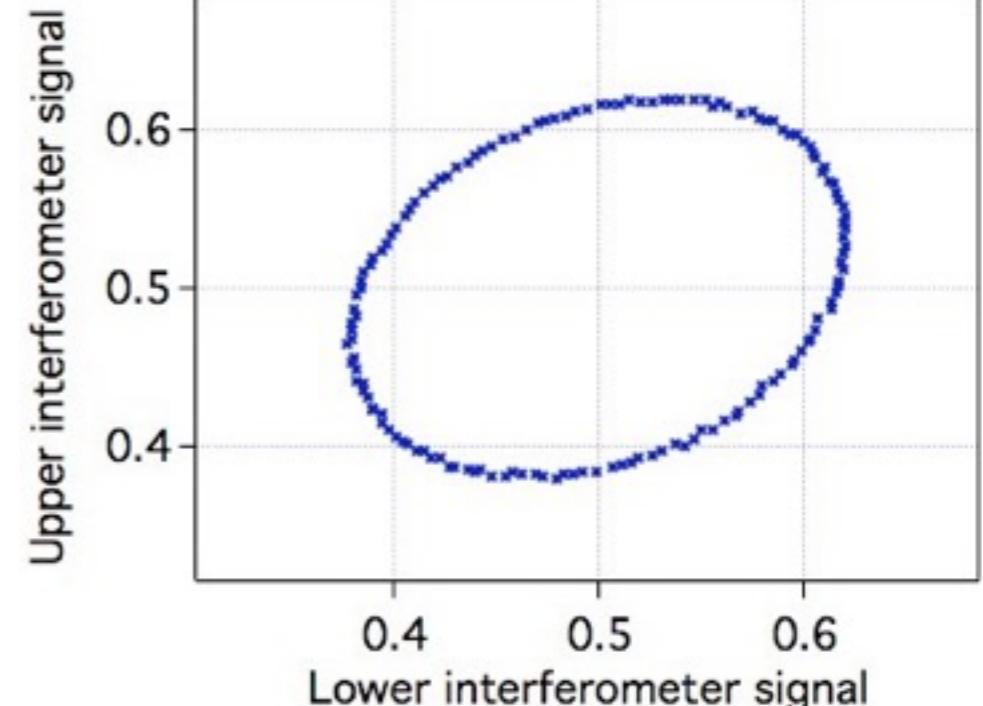
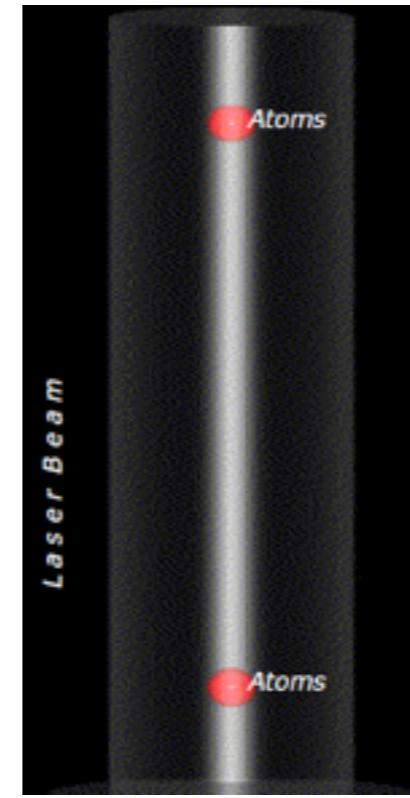
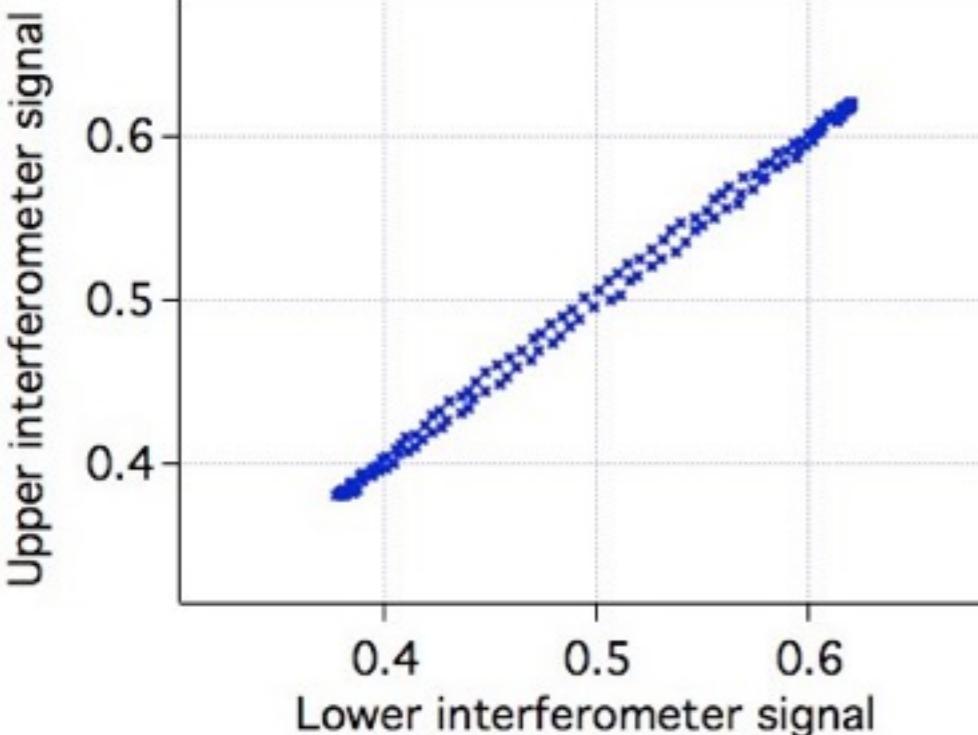


$T = 50 \text{ ms}$   
resol. =  $1.0 \times 10^{-6} \text{ g/shot}$



$T = 150 \text{ ms}$   
resol. =  $3.2 \times 10^{-8} \text{ g/shot}$

$$\Delta\Phi = k_e g T^2$$



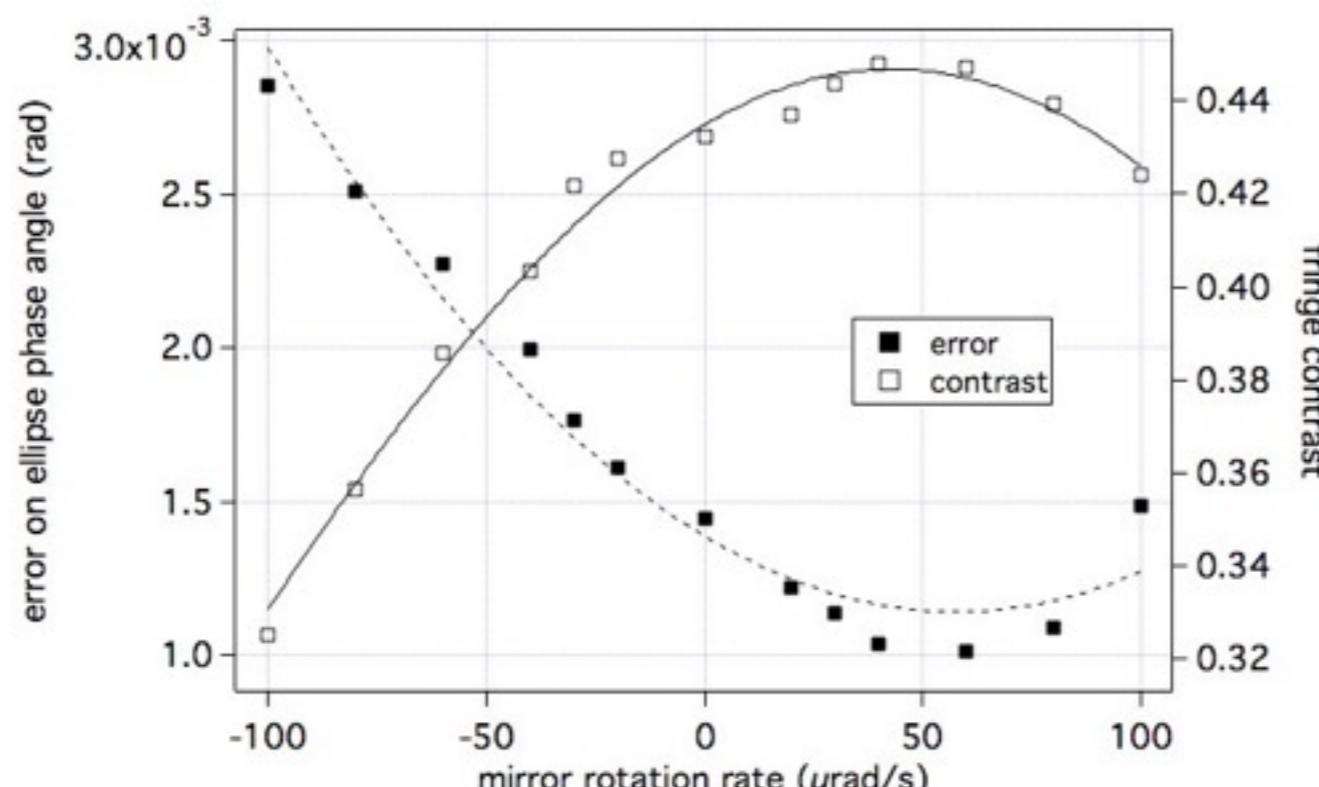
up to 140 dB CMRR with 1 m separation

G. T. Foster & *et al.*, Opt. Lett 27, 951 (2002)

Sensitivity limits...

# Atomic motion

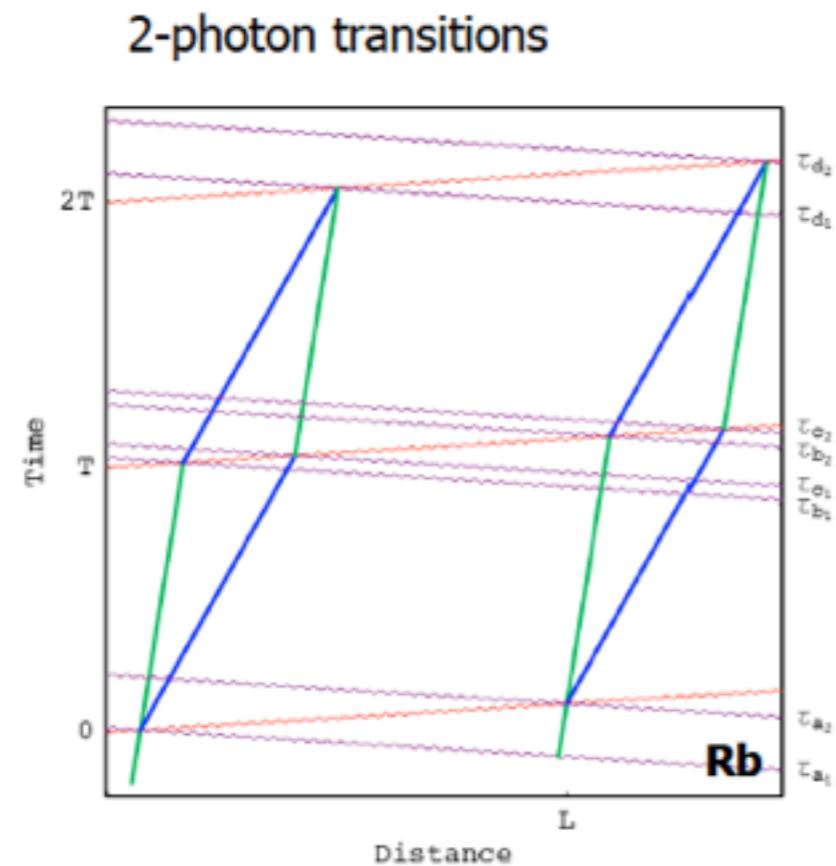
- AI phase is affected by residual atomic motion
- Velocity spread due to finite temperature
- Shot to shot jitter in initial position and mean velocity
- Atomic motion couples with
  - laser wavefront distribution
  - Coriolis acceleration & higher order inertial effects
  - gravitational and magnetic gradients
- Coupled constraints on atomic temperature, launching velocity, and wavefront jitter
- Limits are particularly stringent for very long baseline (space GW detectors)
- For ground applications, nK temperatures and lambda/100 wavefront errors would be ok



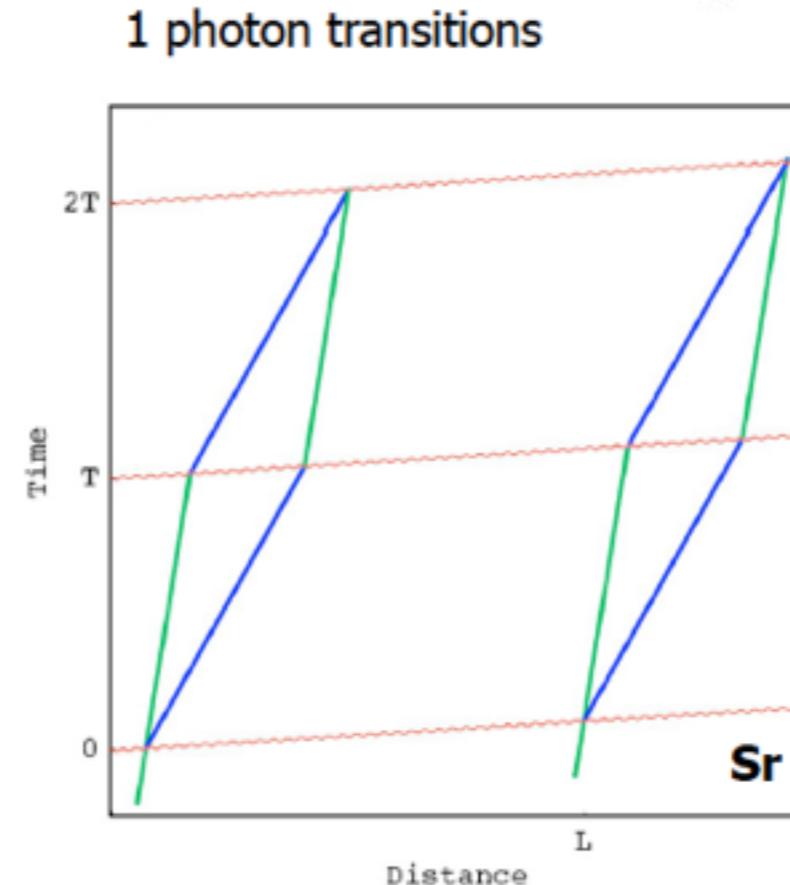
# Laser frequency noise

- Simultaneous AI sensors (gradiometer, strain meter) sharing the same atom optics provide excellent CMRR for phase noise of laser field
- However, for long baseline the differential measurement is affected by common mode laser frequency noise due to high propagation delay
- Requirements on laser frequency noise are dramatically released if the beam splitters are driven by one-photon transitions (e.g. Sr clock transition @ 689 nm)
  - N. Yu and M. Tinto, Gen Relativ Gravit (2011) 43:1943 (2011)
  - P. W. Graham et al., PRL 110(17):171102 (2013)

$$\sigma^2 \simeq 4 \frac{\pi^4}{\tau} \frac{d^2}{c^2} S_\nu$$



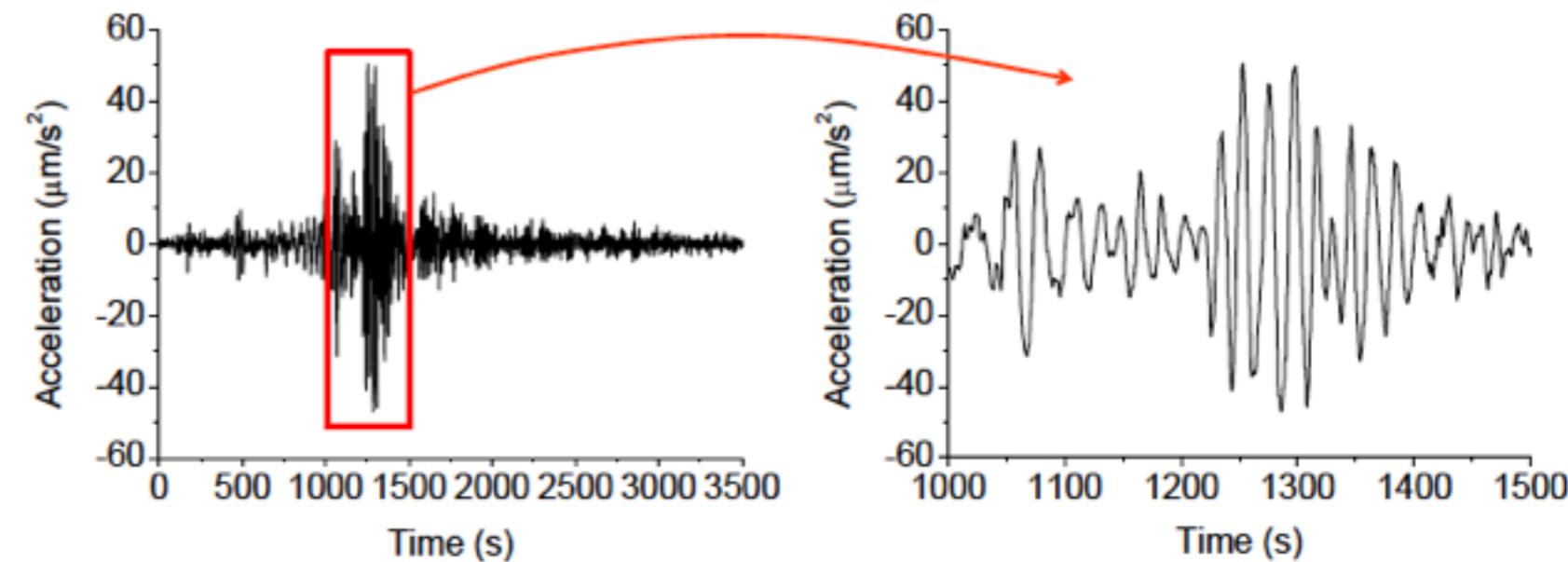
GW signal from relative positions of atom ensembles with respect to optical phase fronts.



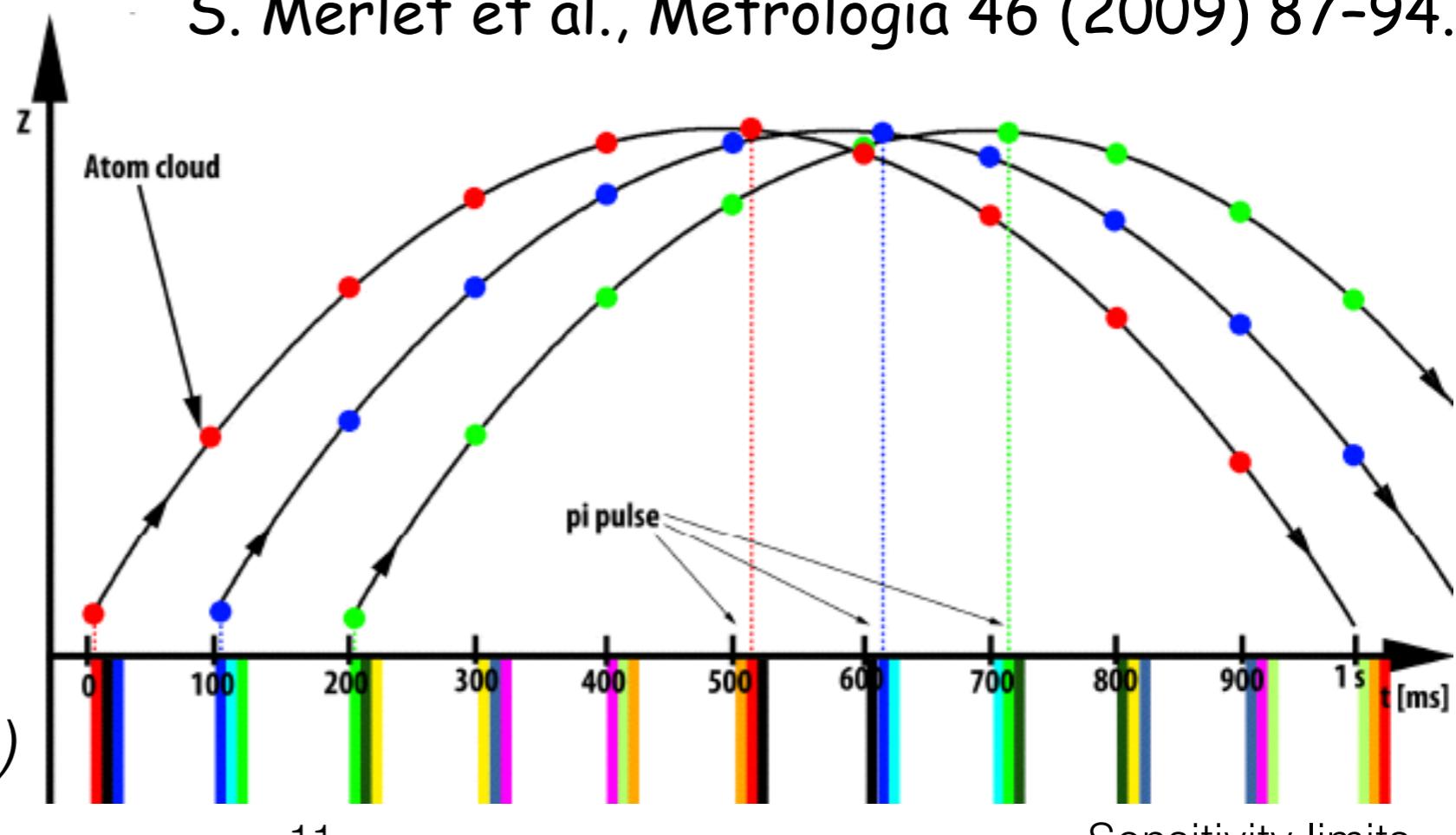
GW signal from light propagation time between atom ensembles.

# Bandwidth & aliasing

- Sensitivity  $\sim T^2$
- Bandwidth  $\sim T^{-1}$
- Highly sensitive gravimeters detect earthquakes
  - $T \sim 50$  ms
  - rep rate  $\sim 5$  Hz
- Improving  $T$  without loosing bandwidth
  - interleaved interferometers
  - main limit: Doppler separation of different channels
  - see e.g. *I. Dutta et al., PRL 116, 183003 (2016)*



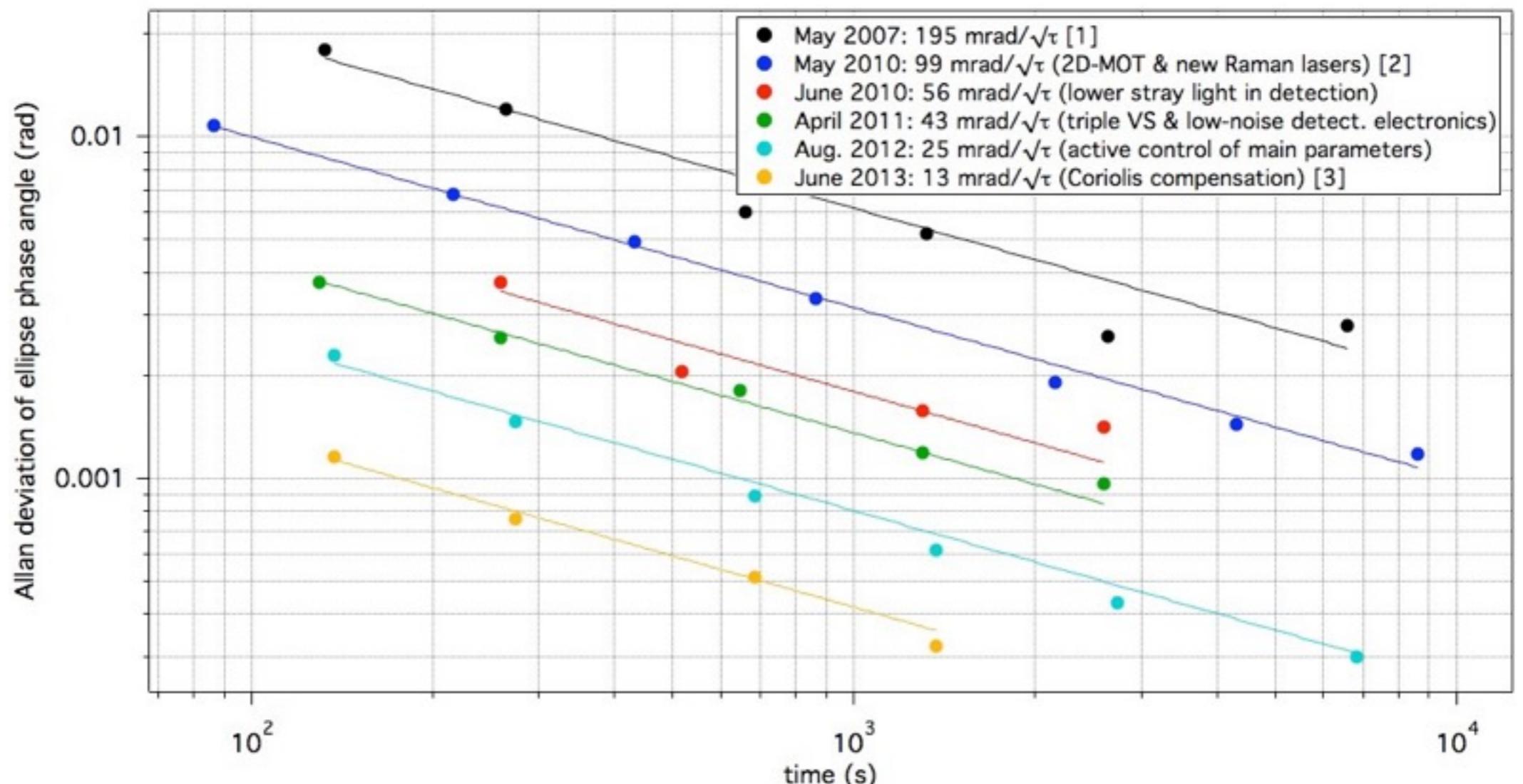
S. Merlet et al., Metrologia 46 (2009) 87–94.



# Technical noise sources

- magnetic fields
- stray light
- photon shot noise
- tilt noise

- [1] G. Lamporesi et al., Phys. Rev. Lett 100, 050801 (2008)
- [2] F. Sorrentino et al., New J. Phys. 12, 095009 (2010)
- [3] F. Sorrentino et al., Phys. Rev. A 89, 023607 (2014)

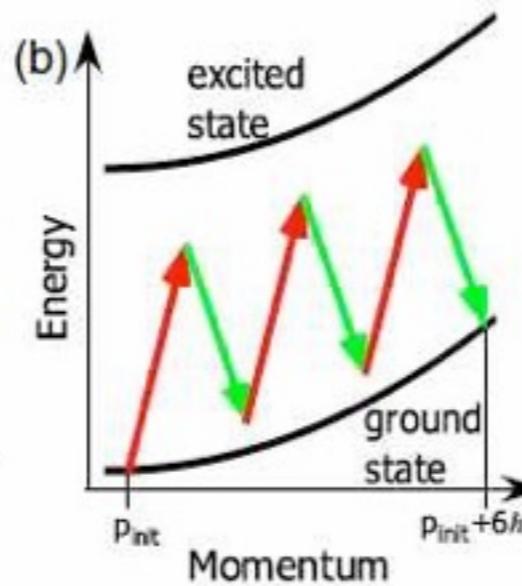
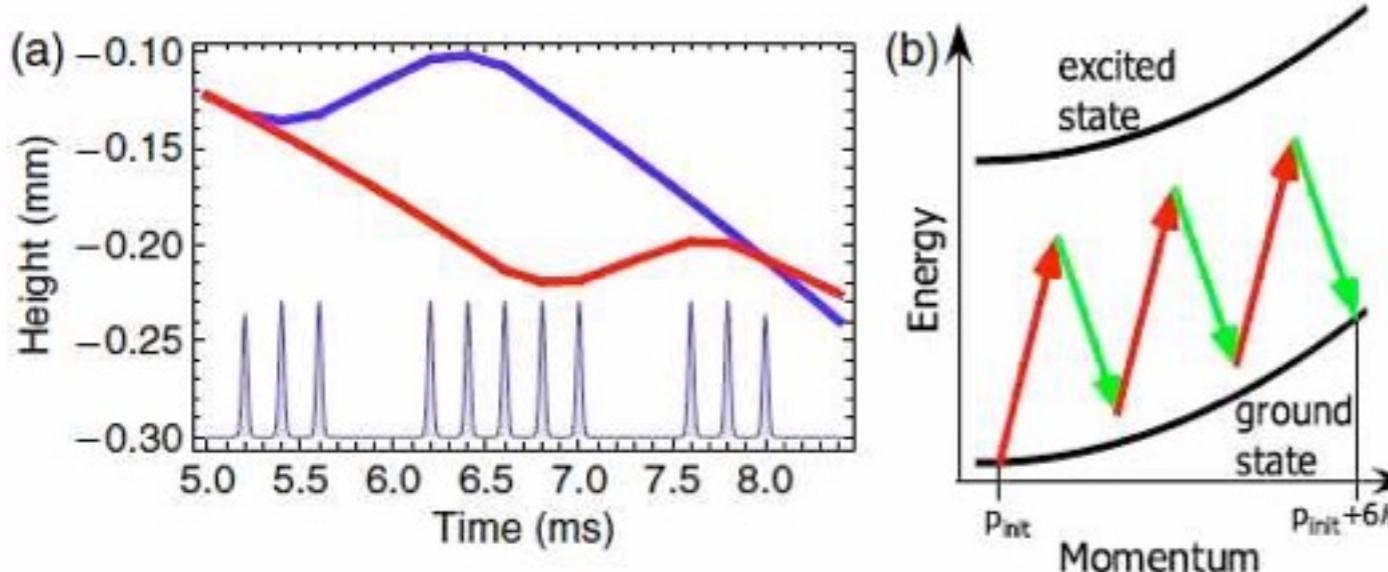


# Potential improvements

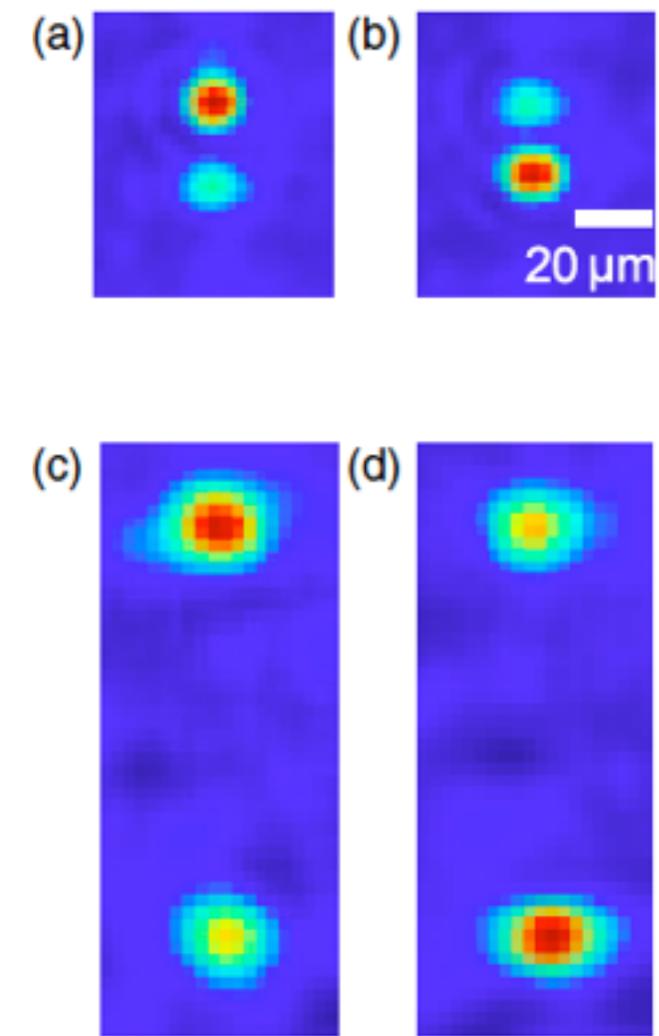
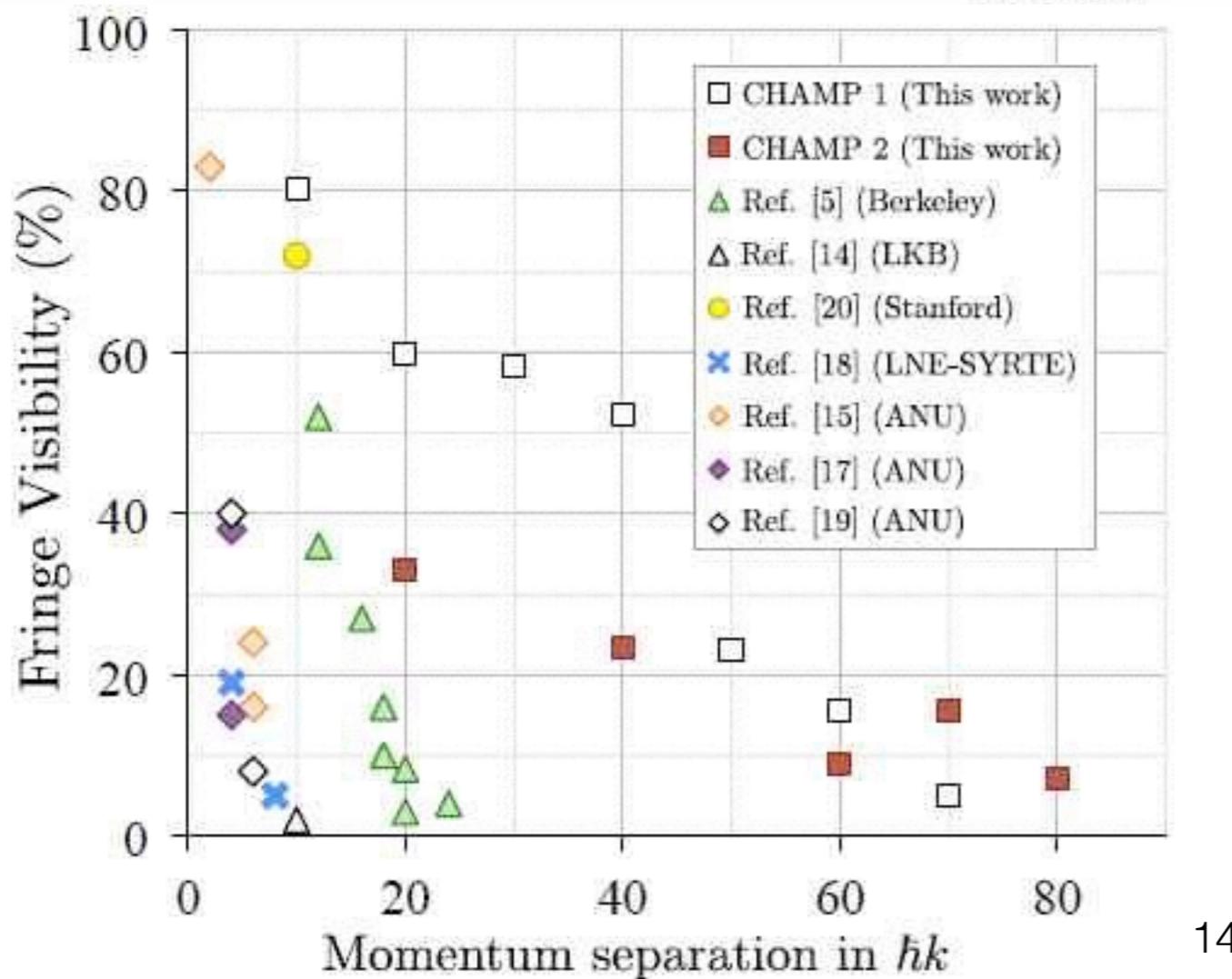
$$\Delta\Phi = nkgT^2$$

- Large momentum transfer splitters (increase  $nk$ )
- Large scale AI (increase  $T$ )
- High flux atomic sources (reduce quantum noise)
- Choice of atomic species
- Interferometer topology
- Squeezing

# Large momentum transfer



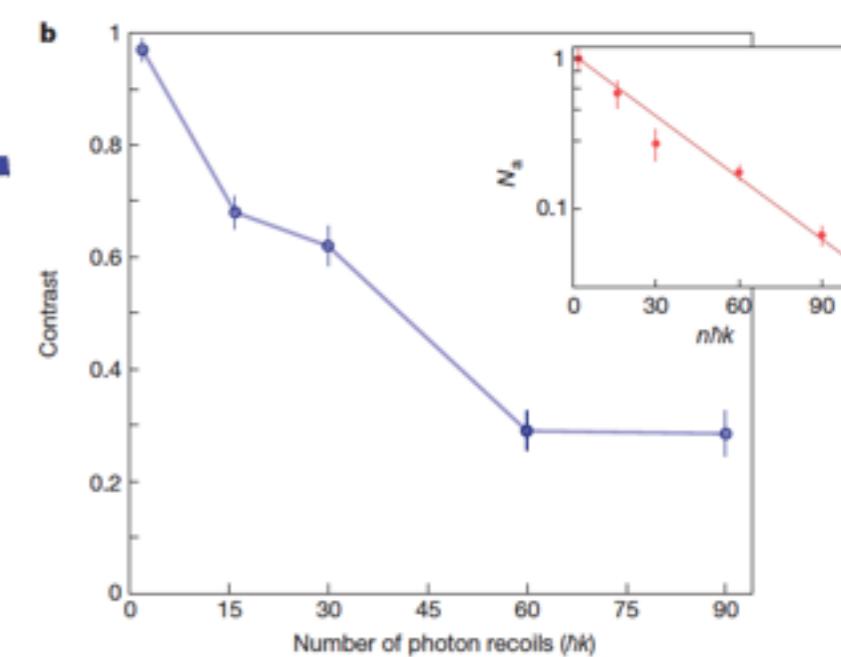
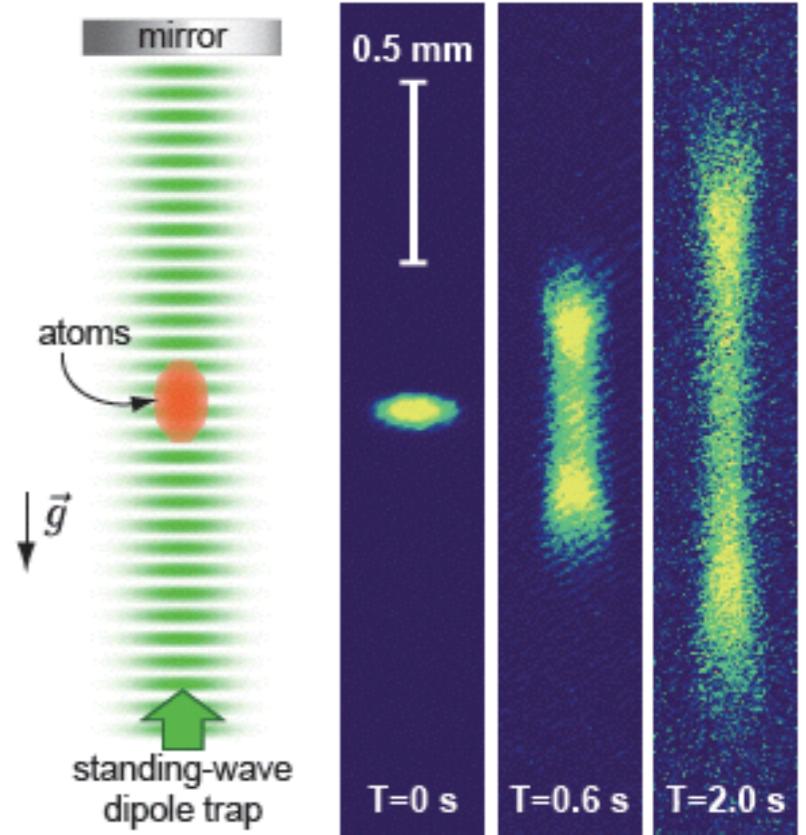
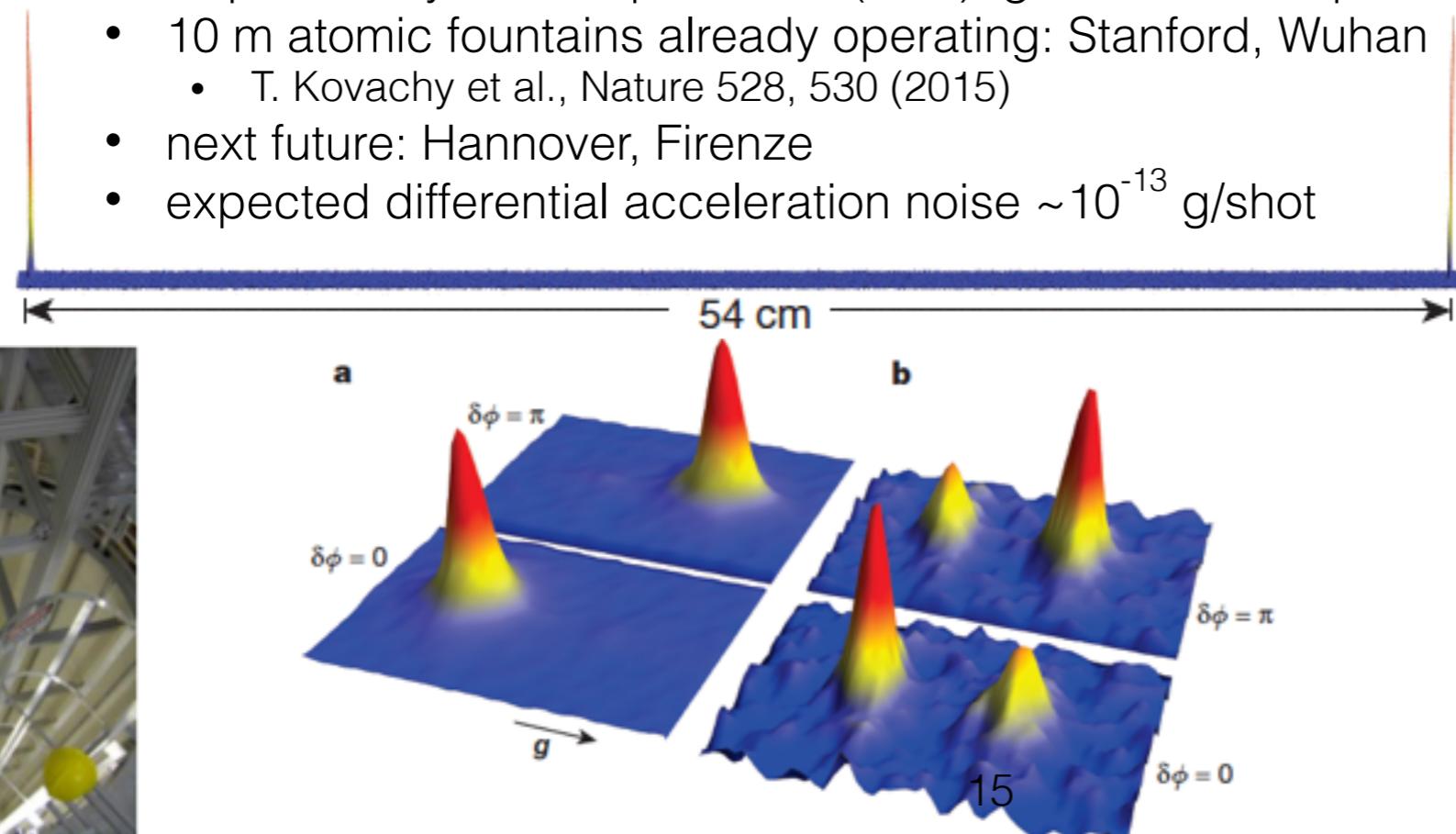
- H. Müller et al., PRL **102**, 240403 (2009)
- S.-W. Chiow et al., PRL **107**, 130403 (2011)
- G. D. McDonald et al., PRA **88**, 053620 (2013)



Sensitivity limits...

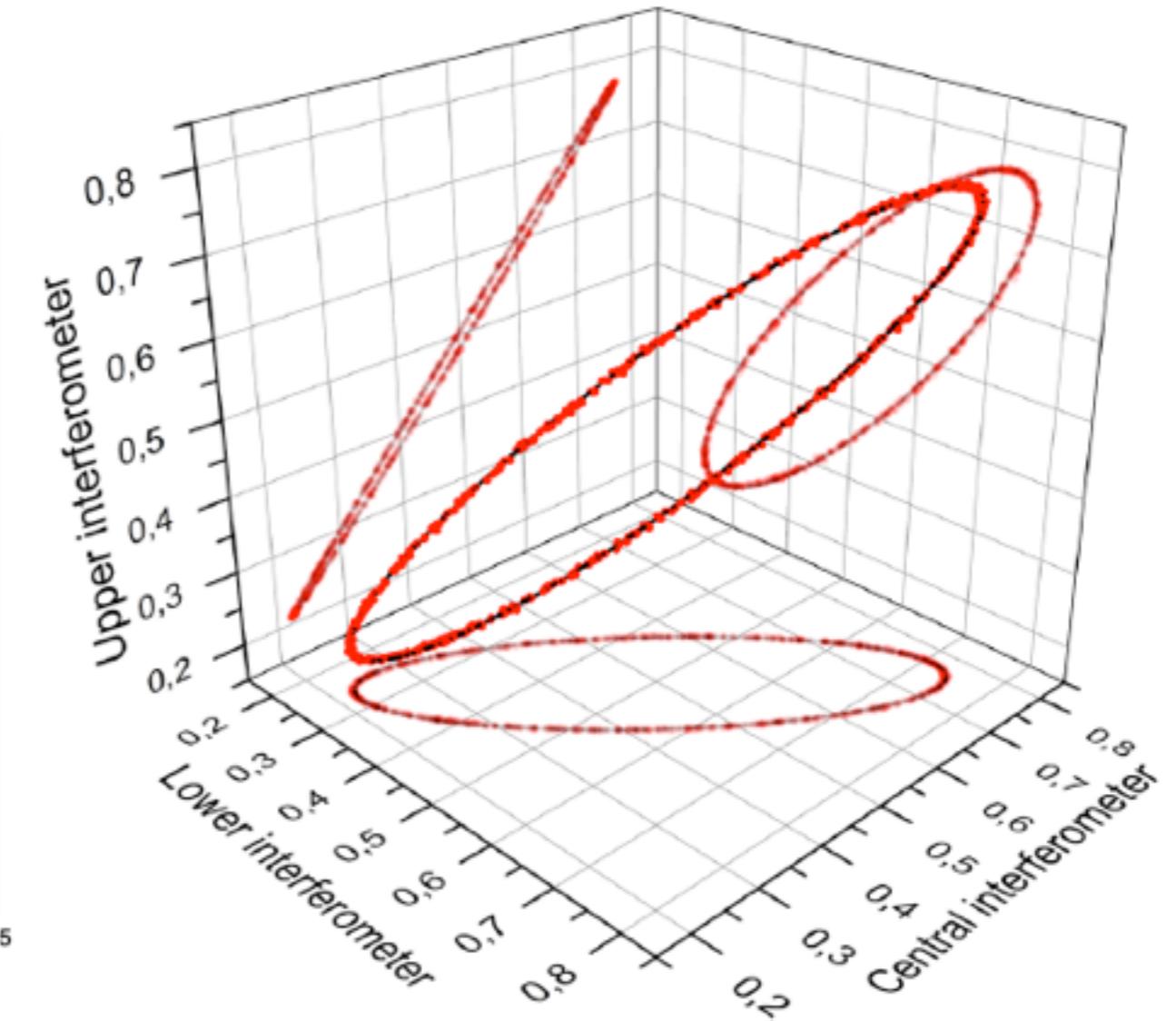
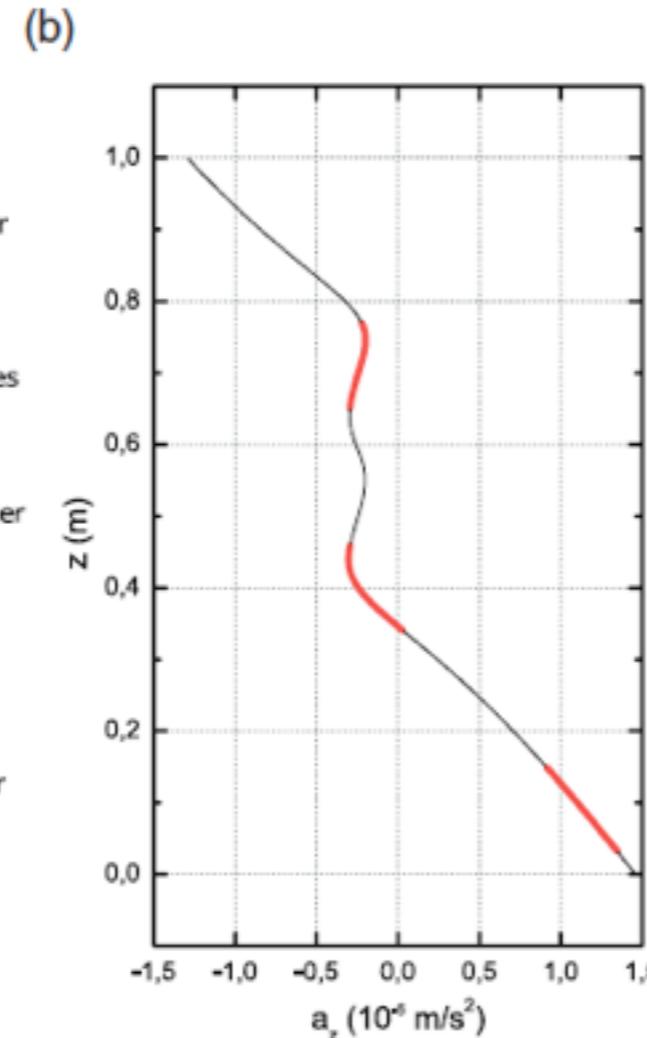
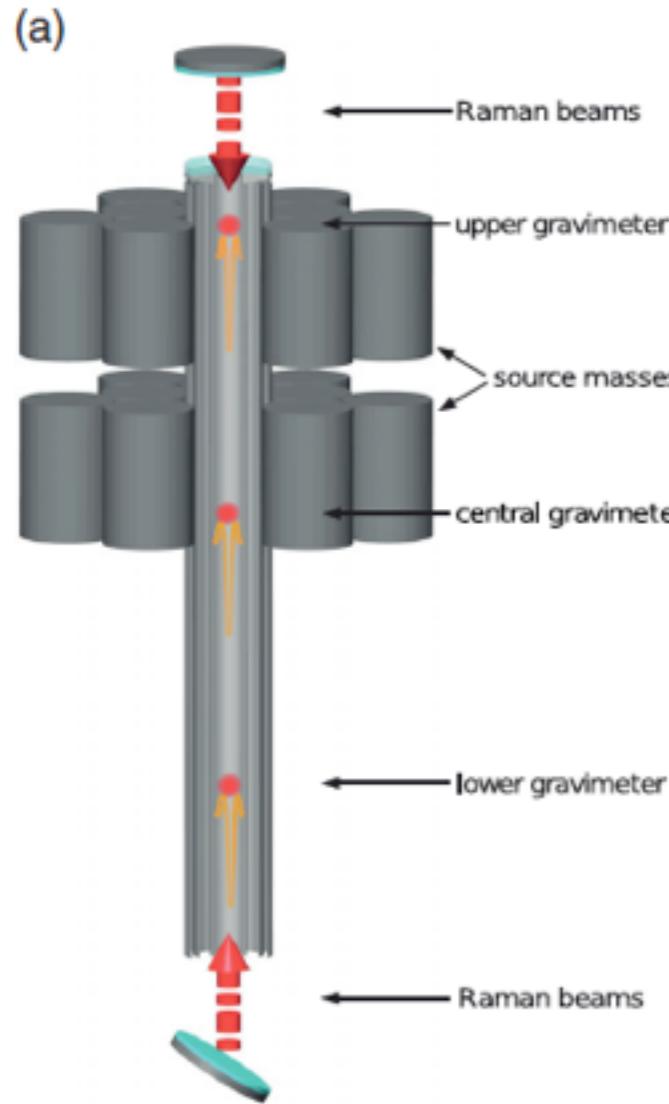
# Long interaction time

- Microgravity: in principle  $T > 10$  s
- Increasing  $T$  on Earth
  - trapped atoms (Sr)
    - up to 15000 coherent photon recoils
    - decoherence time  $> 500$  s
      - G. Ferrari et al., PRL 97, 060402 (2006)
      - V. V. Ivanov et al., PRL 100, 043602 (2008)
      - F. Sorrentino et al., Phys. Rev. A 79, 013409 (2009)
      - M. Tarallo et al., PRA 86, 033615 (2012)
  - free fall
    - requires large vertical size
    - requires very low temperatures ( $\sim nK$ ): good for LMT splitter
    - 10 m atomic fountains already operating: Stanford, Wuhan
      - T. Kovachy et al., Nature 528, 530 (2015)
    - next future: Hannover, Firenze
    - expected differential acceleration noise  $\sim 10^{-13}$  g/shot



Sensitivity limits...

# Multiple samples

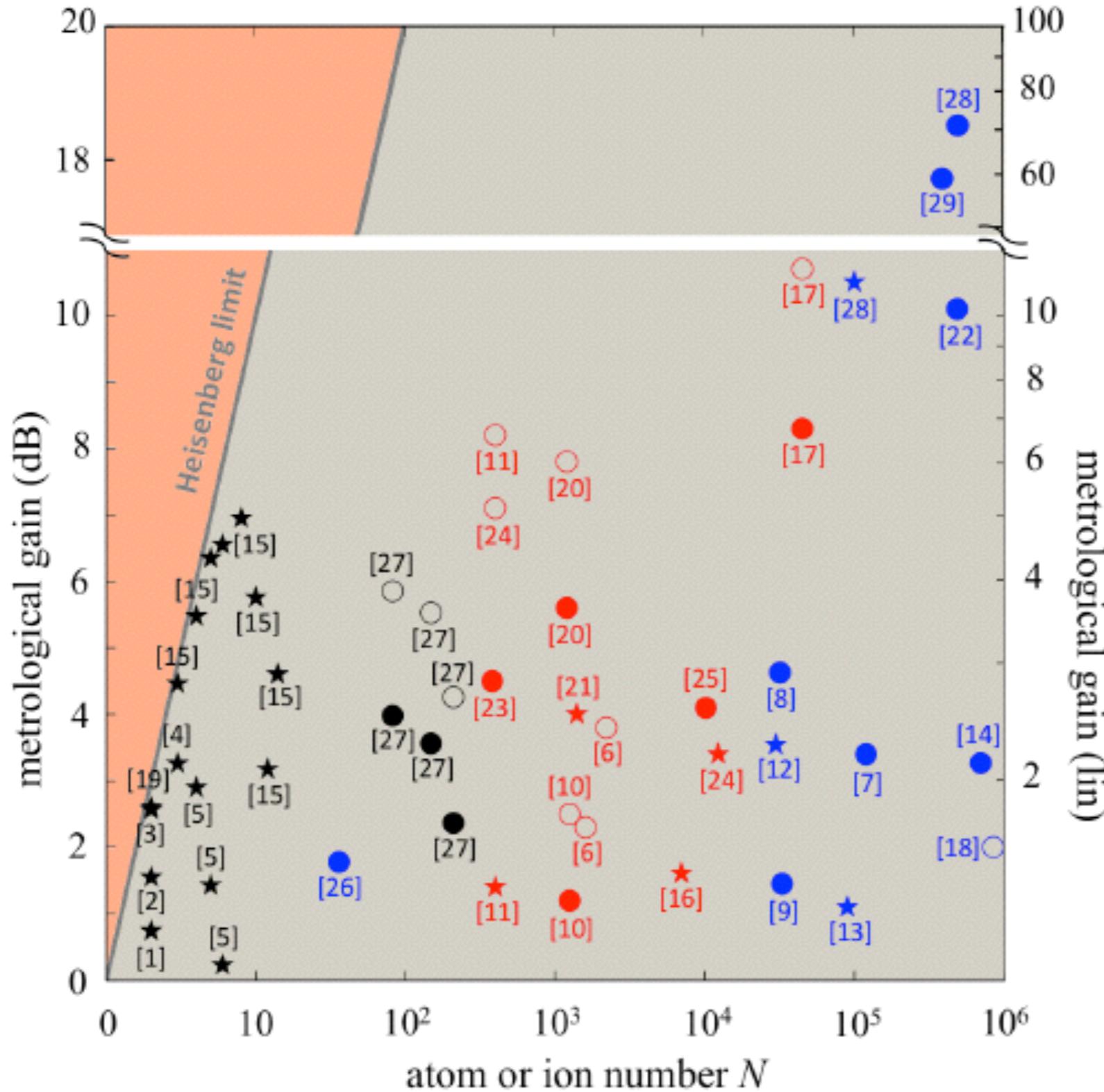


- Use three atomic clouds to measure the vertical derivative of vertical gravity gradient
  - *G. Rosi et al., Phys. Rev. Lett. 114, 013001 (2015)*
- Scalable to arbitrary large number of samples
- Simultaneous, correlated Al can improve g measurements as well
  - *F. Sorrentino et al., Appl. Phys. Lett. 101, 114104 (2012)*
- Multiple, correlated Als to mitigate NN
  - *W. Chaibi et al, Phys. Rev. D 93, 021101 (2016)*

# Atomic species

- Most sensitive AI have been realised with alkali atoms (Rb, Cs)
  - Simple laser cooling
  - Splitter based on Raman transitions between hyperfine states -> simple detection of AI output ports
- Some interferometric schemes would work better with other species (e.g. Sr)
  - low sensitivity to stray fields
  - fast cooling to sub- $\mu\text{K}$  temperatures
  - intercombination lines allow single-photon splitters to reduce the effect of laser frequency noise with long  $T$

# Squeezing



## TRAPPED IONS

- [1] Sackett *et al.*, 2000 NIST
- [2] Meyer *et al.*, 2001 NIST
- [3] Leibfried *et al.*, 2003 NIST
- [4] Leibfried *et al.*, 2004 NIST
- [5] Leibfried *et al.*, 2005 NIST
- [15] Monz *et al.*, 2011 Innsbruck
- [19] Noguchi *et al.*, 2012 Tokyo
- [27] Bohnet *et al.*, 2015 NIST

## BOSE-EINSTEIN CONDENSATES

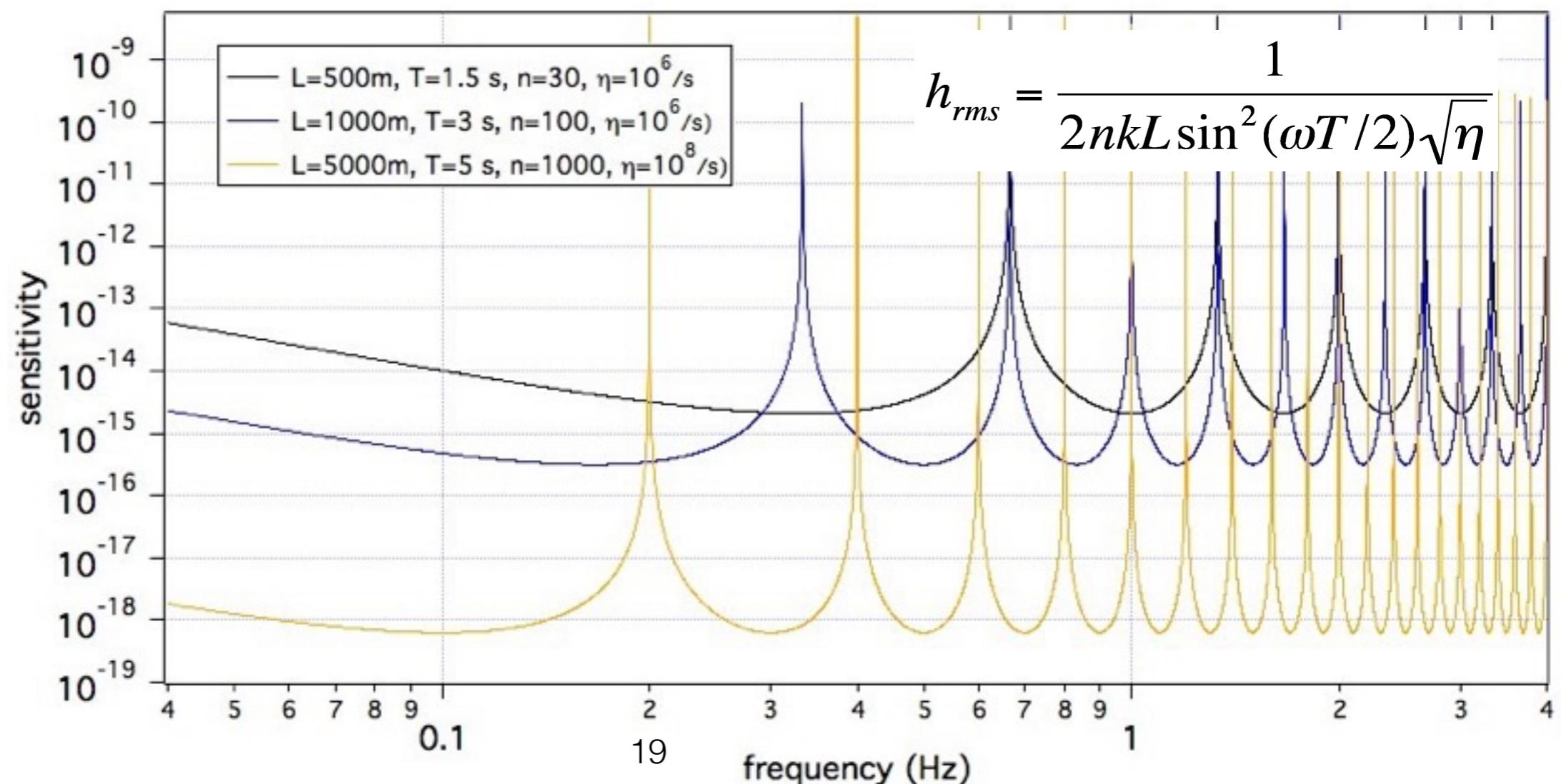
- [6] Esteve *et al.*, 2008 Heidelberg
- [10] Reidel *et al.*, 2010 Basel
- [11] Gross *et al.*, 2010 Heidelberg
- [16] Lucke *et al.*, 2011 Hannover
- [17] Hamley *et al.*, 2012 GeorgiaTech
- [20] Berrada *et al.*, 2013 Wien
- [21] Okeleon *et al.*, 2013 Basel
- [23] Strobel *et al.*, 2014 Heidelberg
- [24] Muessel *et al.*, 2014 Heidelberg
- [25] Muessel *et al.*, 2015 Heidelberg

## COLD ATOMS

- [7] Appel *et al.*, 2009 Copenhagen
- [8] Leroux *et al.*, 2010a MIT
- [9] Schleier-Smith *et al.*, 2010 MIT
- [12] Leroux *et al.*, 2010b MIT
- [13] Louchet-Chauvet *et al.*, 2010 Copenhagen
- [14] Chen *et al.*, 2011 JILA
- [18] Sewell *et al.*, 2012 ICFO
- [22] Bohnet *et al.*, 2014 JILA
- [26] Barontini *et al.*, 2015 ENS
- [28] Hosten *et al.*, 2016 Stanford
- [29] Cox *et al.*, 2016 JILA

# Summary

- Differential acceleration sensitivity of a  $\sim 10$  m size vertical gradiometer
  - interaction time  $T$
  - $n$  photon recoils
  - $10^6$  atoms at QPN
- $n=30, T=1.5$  s  $\rightarrow 6 \times 10^{-13}$  g
- $n=100, T=3$  s  $\rightarrow 4 \times 10^{-14}$  g
- Vertical strainmeter made of two samples with separation  $L$

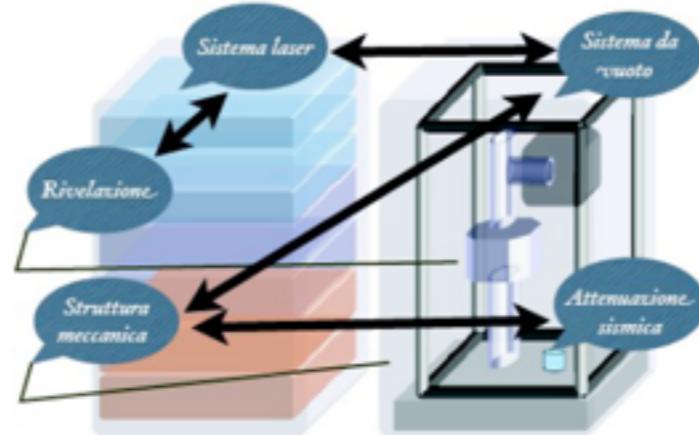


# Technology readiness

- Commercial atomic gravimeters & gradiometers
  - MuQuans (France)
  - AOSense (USA)
  - AtomSensors (Italy)



**AOSense**



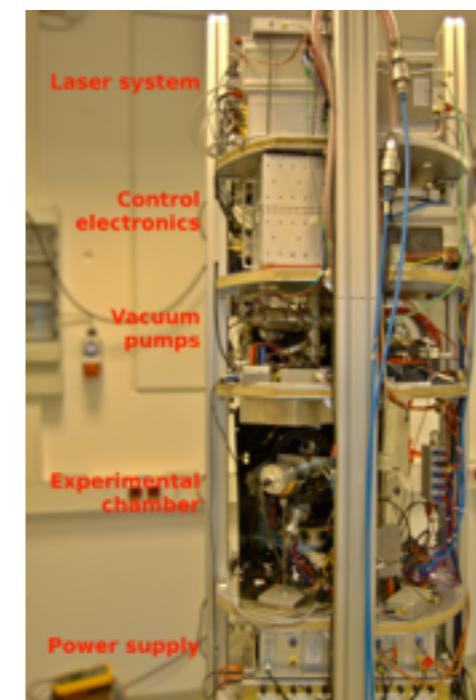
Atom Sensors



**μQuans**

- Prototypes for microgravity
  - I.C.E. (CNES)
  - QUANTUS/MAIUS (DLR)
  - S.A.I. (ESA)

G. Tino et al., Precision Gravity Tests with  
Atom Interferometry in Space, Nuclear  
Physics B 243, 203 (2013)



**QUANTUS**



**I.C.E.**  
Atom Interferometry in Microgravity

# Thank you