

Macro dark matter selfgravitating halos around galaxies

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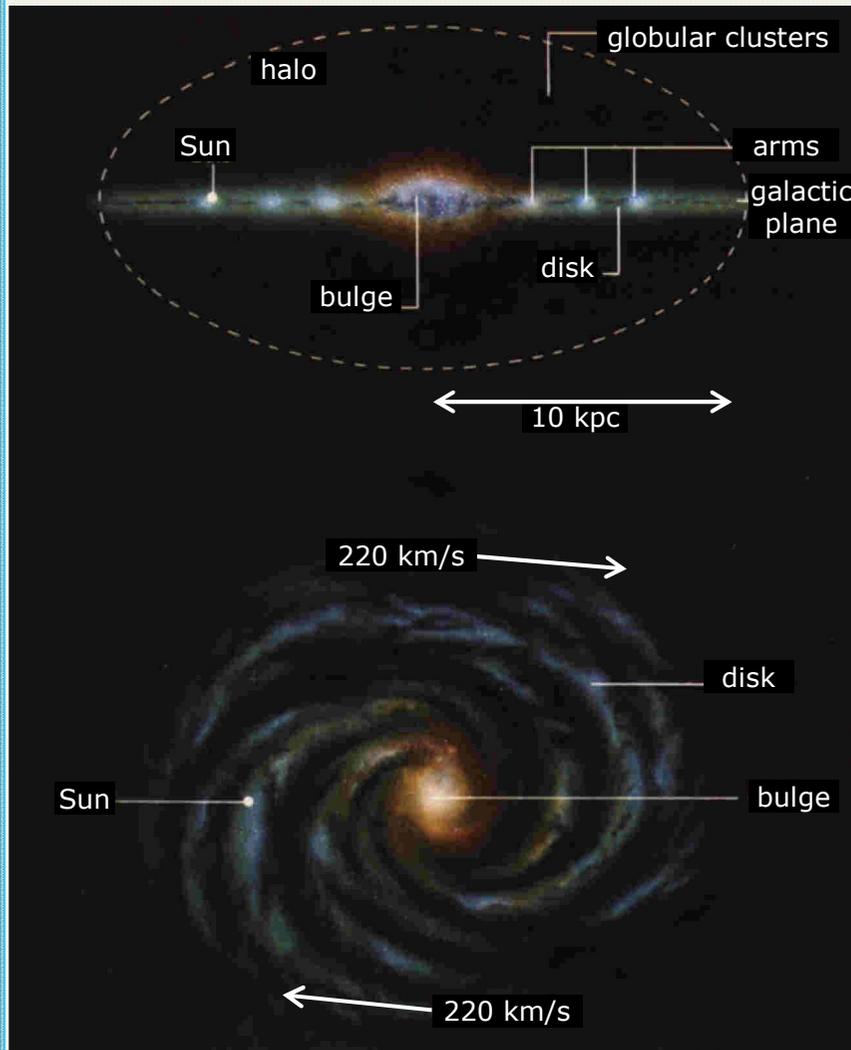
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*Quest for visible and invisible
strange stuff in the Universe
LNF - Frascati, 27 Nov 2015*



The Milky Way



radial velocity

$$v(r) = \sqrt{\frac{GM(r)}{r}}$$

simplified model with mass concentrated in the bulge

- bulge: uniform density ρ
- disk: negligible density



$$M(r) = \frac{4}{3} \pi \rho r^3 \quad \text{for } r \leq R_b$$

$$M(r) = M \quad \text{for } r > R_b$$

The rotation curve

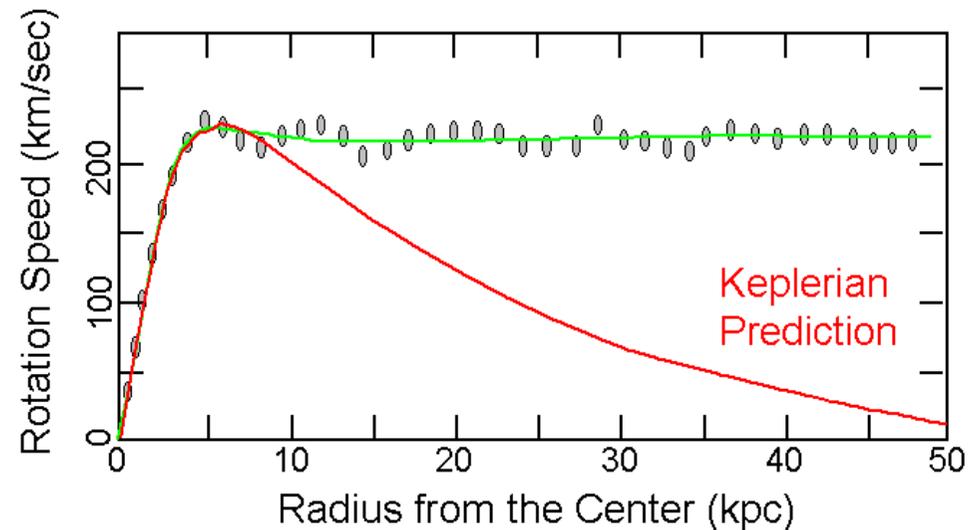
rotation velocity

Bulge: $v(r) = 2 \left(\frac{\pi}{3} G \rho \right)^{1/2} r$ for $r \leq R_b$ (rigidly rotating body)

Disk: $v(r) = \sqrt{\frac{GM}{r}}$ for $r > R_b$ (Keplerian velocity)

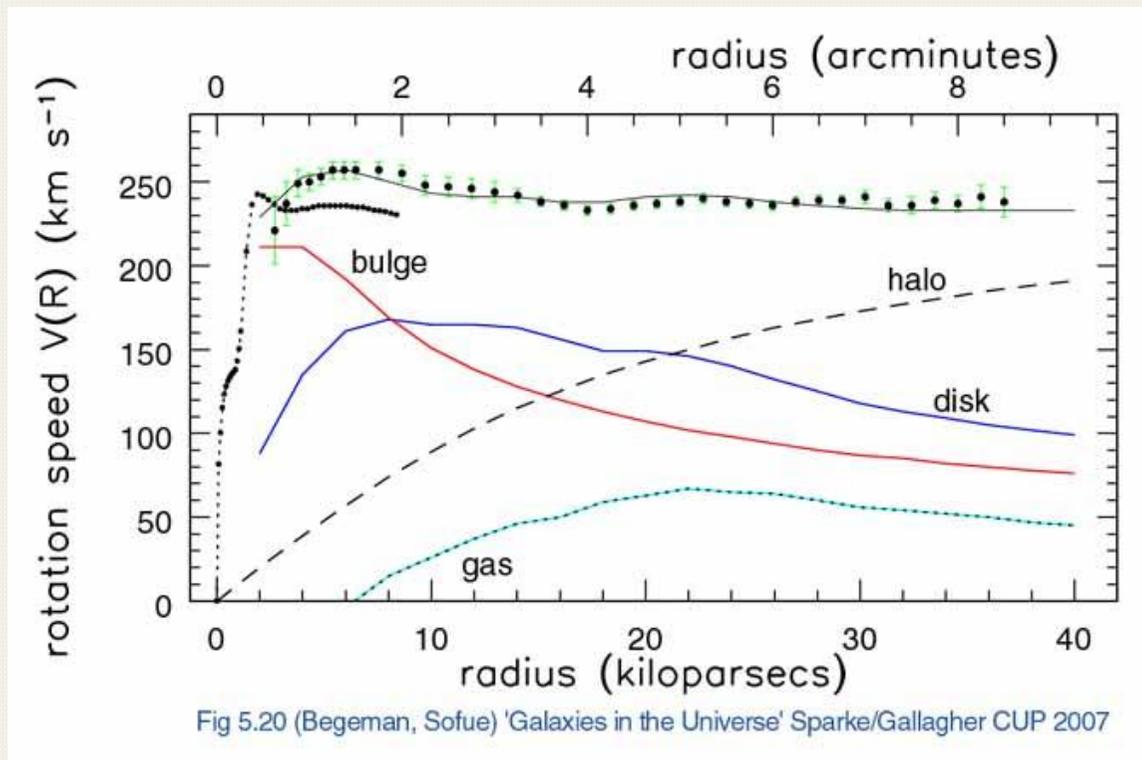
**theoretical predictions
in contrast
with observational data**

Observed vs. Predicted Keplerian



Dark matter halo

The existence of dark matter halo can explain the rotation speed of the Galaxy (Zwicky, 1933)



$$M_{halo} \approx 10 M_{gal}$$

$$R_{halo} \approx 10 R_{gal}$$

**two big questions
what particle ?
what mass ?**

First possibility: WIMPs

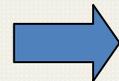
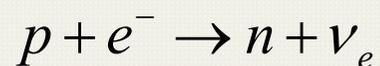
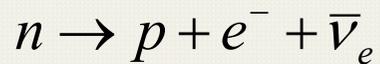
A VERY SIMPLE CASE

degenerate gas of particles constituting the galactic halo
 - polytropic model with $n=3/2$ -

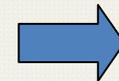
$$M = \frac{3}{2} \left(\frac{\pi}{2} \right)^{3/2} (2.71406) \frac{\hbar^3}{G^{3/2} m^4} \rho_0^{1/2}$$

$$R = \frac{(9\pi)^{1/6}}{2\sqrt{2}} (3.65375) \frac{\hbar}{G^{1/2} m^{4/3}} \rho_0^{-1/6}$$

a possible hypothesis due to importance of β decay in stellar equilibrium



massive neutrino



$$m = m_\nu \approx 10 \text{ eV}$$

Neutrino hypothesis

with $M_{halo} \approx 10 M_{gal}$ $R_{halo} \approx 10 R_{gal}$

➔ $\rho_0 \approx 10^{-25} \text{ g/cm}^3$

➔ $R \cong 90 \left(\frac{M}{10^{12} M_{\odot}} \right)^{-1/3} \text{ kpc}$

critical density

$$\rho_{cr} = \frac{m_{\nu}^4 c^3}{3\pi^2 \hbar^3} = 7.8 \cdot 10^{-17} \text{ g/cm}^3 \gg \rho_0$$

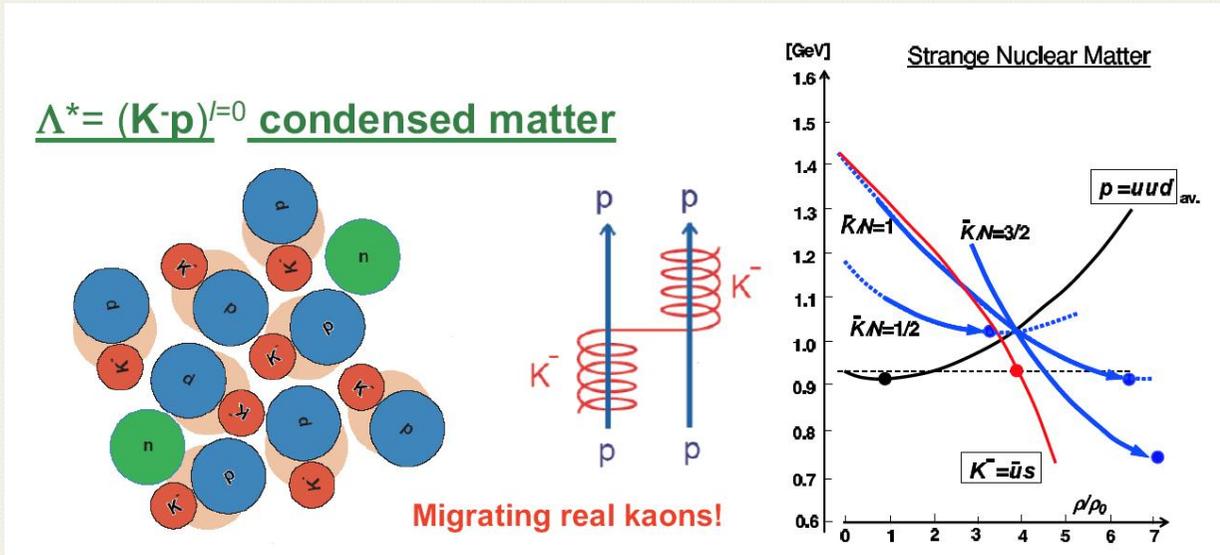
General Relativity? NO

$$\frac{GM}{Rc^2} = 4.8 \cdot 10^{-7} \ll 1$$

neutrino dark matter halos are nonrelativistic and also Newtonian

- **SMPs as alternative candidates to WIMPs for dark matter**
- **very massive particles ($m \sim \text{GeV}$), low number density**
 - low effective interaction rate in spite of a not small cross section (dark matter in big bang standard model ?)
- **massive particle lifetime sufficiently large ? stability ?**
 - big bang relics, background ?
- **the role of strangeness**
 - quark configuration with the same (approximate) number of u, d, s
 - chemical potential due to Pauli exclusion principle favourable to stable configurations (strange quark matter conglomerates)
- **quark matter configuration**
 - $\Lambda^*(1405)$ as a possible candidate for dark matter (also in neutron stars?)

Neutron stars and ... dark matter ?



Akaishi & Yamazaki, 2015 (for neutron stars)

one possibility among different $\bar{K}N$ states:

$$\Lambda^*(1405) \equiv K^-p$$

strong decay,
single Λ^* is unstable



Λ^* clusters:

$$Nm_{\Lambda^*} \equiv m^* \sim \text{GeV}$$

Neutron stars: $\bar{K}^0 n \rightarrow K^- p$? $N \ll N_n$? **Hyperon stars (cores) ?**

Dark matter: $\Lambda^*(1405) \equiv K^- p$ $N \leq 10$ may be stable? $m^* \sim 5 \div 10 \text{ GeV}$

Ultra-dense kaonic nuclear states as partial constituent of dark matter ?

Statistics: semidegenerate gas?

- For calculating selfgravitating equilibrium configurations of dark matter halos, the internal structure of the single Λ^* cluster is not relevant. Nevertheless we could expect $\rho > 10^{15} \text{ g/cm}^3$ in the internal structure of the Λ^* cluster.
- We consider Λ^* cluster like a massive particle of mass m^* only gravitationally interacting with the other Λ^* clusters composing the halo.
- The existence of stable Λ^* clusters is an *open question*, nevertheless we may explore the possibility of having halos composed by strange dark matter.
- The first possibility is to consider a *semidegenerate gas* of particles with mass $m^* = 5 \div 10 \text{ GeV}$.
- We search for halos with masses $M \sim 10^{12} M_\odot$ and $R \sim 100 \text{ kpc}$, then the mean density $\langle \rho \rangle$ is of the order of 10^{-26} g/cm^3 .

Gravitational equilibrium

For a mass $m^*=5\text{GeV}$ we have

$$\rho_{cr} = \frac{m^{*4} c^3}{3\pi^2 \hbar^3} = 4.9 \cdot 10^{18} \text{ g/cm}^3 \gg \langle \rho \rangle; \quad \frac{GM}{Rc^2} = 4.8 \cdot 10^{-7} \ll 1$$

ALSO

strange dark matter halos are nonrelativistic and Newtonian

Semidegenerate Fermi distribution function with cutoff in energy:

$$\begin{cases} f(\varepsilon) = \frac{g}{h^3} \frac{1 - e^{(\varepsilon - \varepsilon_c)/kT}}{e^{(\varepsilon - \mu)/kT} + 1} & \text{for } \varepsilon \leq \varepsilon_c \\ f(\varepsilon) = 0 & \text{for } \varepsilon > \varepsilon_c \end{cases}$$

cutoff:

$$\varepsilon_c = m(\phi_R - \phi)$$

mass density:

$$\rho = m \int f d^3q$$

Poisson equation for gravitational equilibrium:

$$\frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{d\phi}{dr} \right) = 4\pi G\rho \quad \text{with} \quad \phi'(0) = 0; \quad \phi(0) = \phi_0$$

Dimensionless quantities

by imposing $r = \eta x$ with $\eta = \left(\frac{gm^4 \sigma G}{h^3} \right)^{-1/2}$ and $\sigma^2 = \frac{2kT}{m}$

$\rightarrow \frac{1}{x^2} \frac{d}{dx} \left(x^2 \frac{dW}{dx} \right) = -8 \pi \hat{\rho}$ with $W'(0) = 0; W(0) = W_0$

$$R = \eta \hat{R}; \quad M = \frac{\sigma^2 \eta}{G} \hat{M}; \quad \rho_0 = \frac{\sigma^2}{G \eta^2} \hat{\rho}_0; \quad W = \frac{\varepsilon_c}{kT}; \quad g = 2s + 1$$

dimensionless quantities depend on W_0 and θ_R

$\rightarrow \hat{\rho} = 2\pi \int_0^W g_s(z, W, \theta_R) z^{1/2} dz; \quad \hat{M} = 4\pi \int_0^{\hat{R}} \hat{\rho} x^2 dx = -\frac{1}{2} \left(x^2 \frac{dW}{dx} \right)_{x=\hat{R}}$

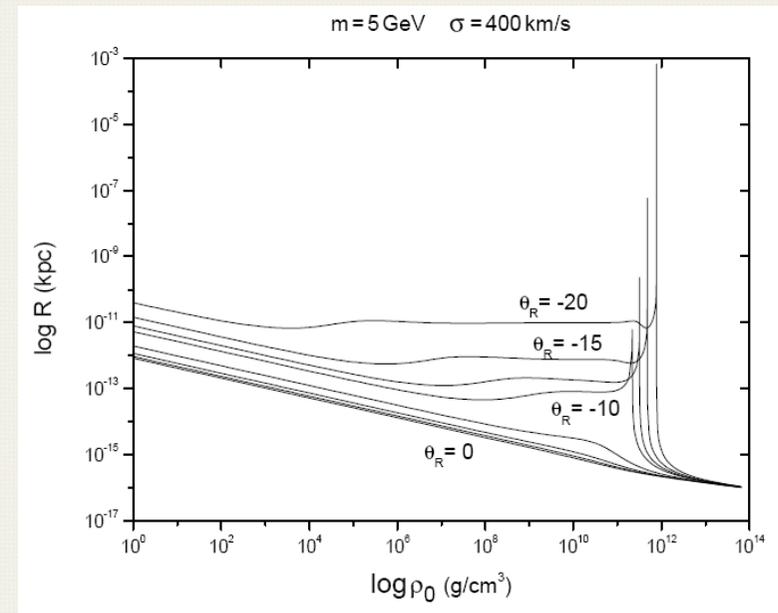
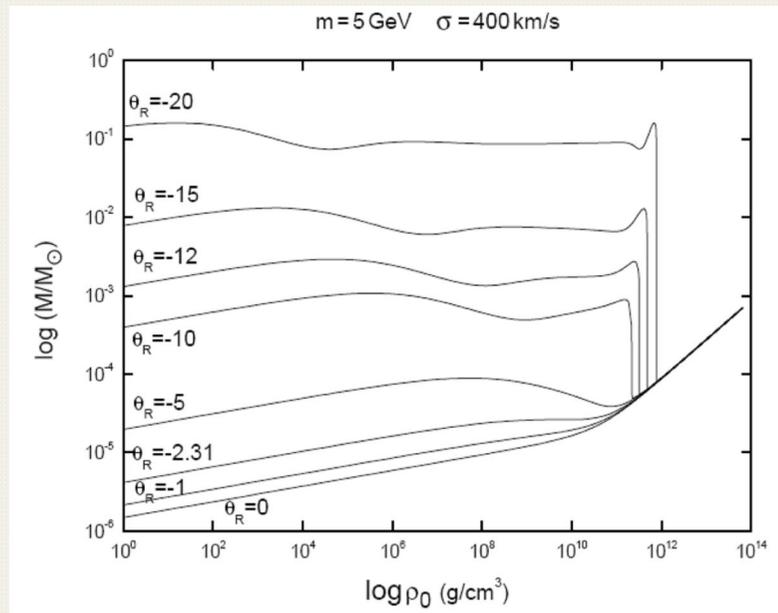
where

$$z = \frac{\varepsilon}{kT}; \quad f(\varepsilon) \Rightarrow \frac{g}{h^3} g_s(z, W, \theta_R); \quad g_s(z, W, \theta_R) = \begin{cases} \frac{1 - e^{z-W}}{e^{z-W-\theta_R} + 1} & \text{for } z \leq W \\ 0 & \text{for } z > W \end{cases}$$

with $\theta = \frac{\mu}{kT}$ and $\theta_R = \theta - W \leq 0$ (MM & Alberti, 2014)

Equilibrium configurations

- By integrating the Poisson equation, we obtain different equilibrium configurations at different values of W_0 and θ_R .
- The solutions also depend on m (mass of the particle) and σ (surface velocity dispersion) through scaling laws.
- The results are summarized in M vs ρ_0 and R vs ρ_0 diagrams for $m=5\text{GeV}$ and $\sigma=400\text{km/s}$.



Statistics: classical ?

The mass $m^*=5\text{GeV}$ doesn't allow to obtain the expected values of central density, mass and radius for a galactic halo

In fact we have: $\rho_0 \sim \sigma^3 m^4$; $M \sim \sigma^{3/2} m^{-2}$; $R \sim \sigma^{-1/2} m^{-2}$
 The densities are too large, masses and radii too small



Semidegenerate regime is not appropriate to describe strange dark matter halos: we need θ_R values much more negative, typical for a classical regime

➔ Boltzmann (King) distribution function with cutoff in energy

On the other hand, for $-\theta_R \gg 1$ $f(\varepsilon) \rightarrow \frac{g}{h^3} e^{\mu/kT} \left(e^{-\varepsilon/kT} - e^{-\varepsilon_c/kT} \right)$ for $\varepsilon \leq \varepsilon_c$
 and $g_s(z, W, \theta_R) \Rightarrow e^{\theta_R} g_K(z, W)$ with $g_K(z, W) = \left(e^{W-z} - 1 \right)$ for $z \leq W$

THEN

**strange dark matter halos are nonrelativistic,
Newtonian and do not follow quantum statistics**

In order to obtain halos with appropriate densities, masses and radii, we calculate equilibrium configurations at fixed central density ($\rho_0 = 10^{-24} \text{ g/cm}^3$) and particle mass ($m^* = 5 \text{ GeV}$), while increasing the value of $-\theta_R$ until to reach $M \sim 10^{12} M_\odot$ and $R \sim 100 \text{ kpc}$.

Using the expression of surface velocity dispersion σ in function of ρ_0

$$\sigma = \left(\frac{1}{\hat{\rho}_0} \frac{h^3}{g m^4 e^{\theta_R}} \right)^{1/3} \rho_0^{1/3},$$

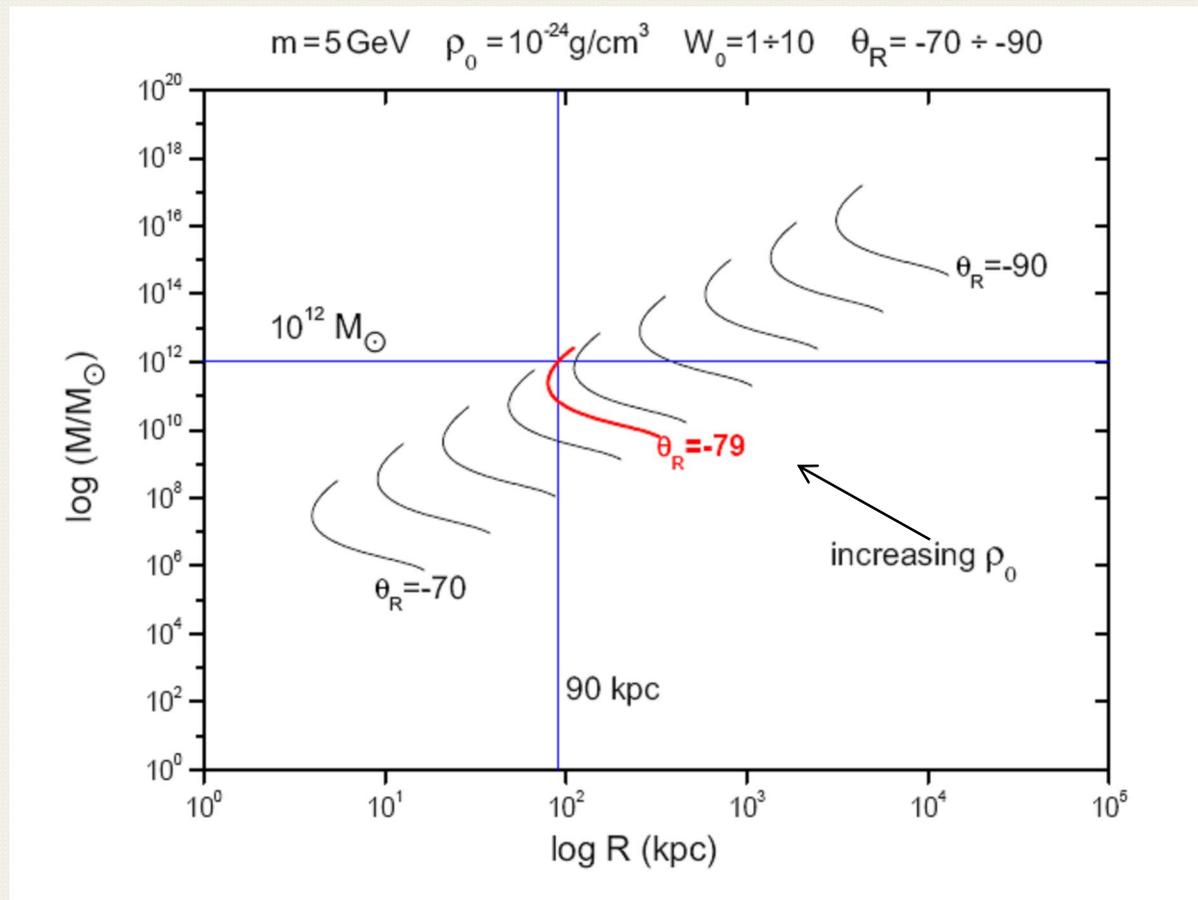
the expressions of M and R become

$$M = \frac{\hat{M}}{\hat{\rho}_0^{1/2}} \frac{h^3}{g m^4 e^{\theta_R} G^{3/2}} \rho_0^{1/2}; \quad R = \hat{R} \hat{\rho}_0^{1/6} \frac{h}{g^{1/3} m^{4/3} e^{\theta_R/3} G^{1/2}} \rho_0^{-1/6}$$

now dimensionless quantities depend on W_0 only

Individuating the value of θ_R

We calculated solutions in the range $W_0=1\div 10$ (for globular clusters the most significant values are between 4 and 8; for galactic halos we expect even less)



In this regime, the dependence on θ_R become a scaling law.

It is possible to make a tuning by varying the central density ρ_0 and the parameter θ_R in order to match the requested values in M and R, also at different values of W_0

For $m^* = 5\text{GeV}$ and $\rho_0 = 10^{-24}\text{ g/cm}^3$ we get

$$\theta_R = -79; \quad W_0 = 1.9; \quad M = 9.78 \cdot 10^{11} M_\odot; \quad R = 90.17 \text{ kpc}$$

$$\bar{\rho} = \frac{3M}{4\pi R^3} = 2.16 \cdot 10^{-26} \text{ g/cm}^3; \quad \sigma = 391 \text{ km/s}$$

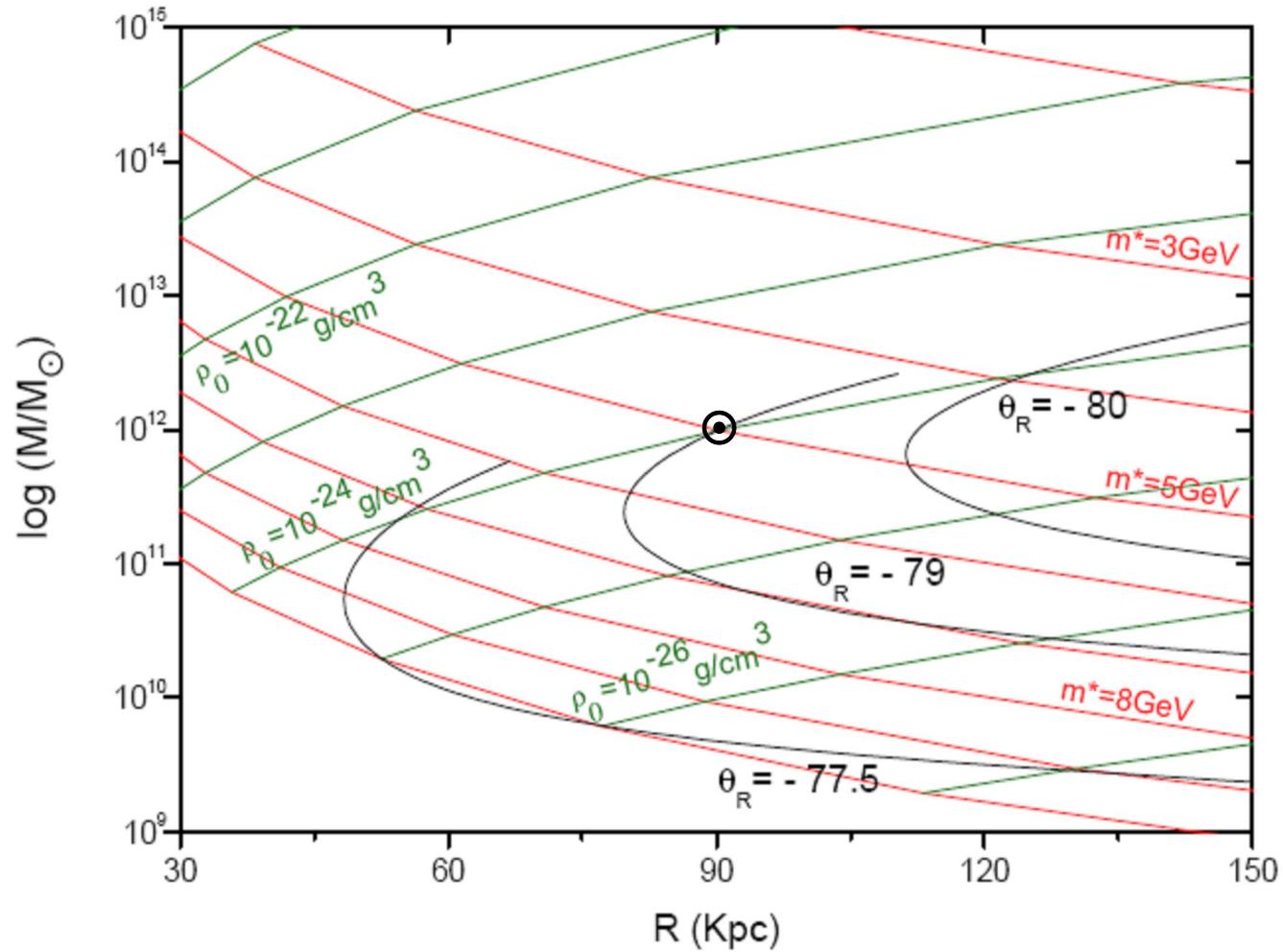
The obtained values are very satisfying !

The other solutions are obtained by scaling laws involving the total mass M and the radius R

$$M = 9.78 \cdot 10^{11} \left(\frac{\rho_0}{10^{-24} \text{ g/cm}^3} \right)^{1/2} \left(\frac{m^*}{5 \text{ GeV}} \right)^{-4} M_\odot$$

$$R = 90.17 \left(\frac{\rho_0}{10^{-24} \text{ g/cm}^3} \right)^{-1/6} \left(\frac{m^*}{5 \text{ GeV}} \right)^{-4/3} \text{ kpc}$$

Strange dark matter halos

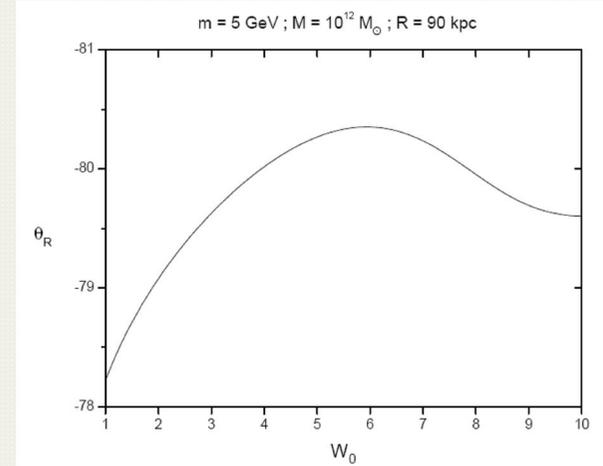
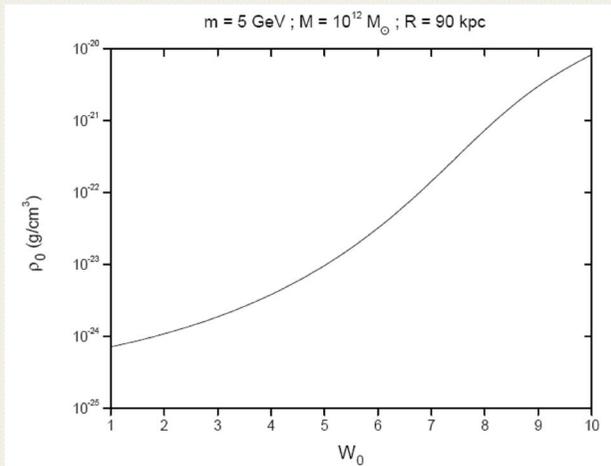


Changing parameters

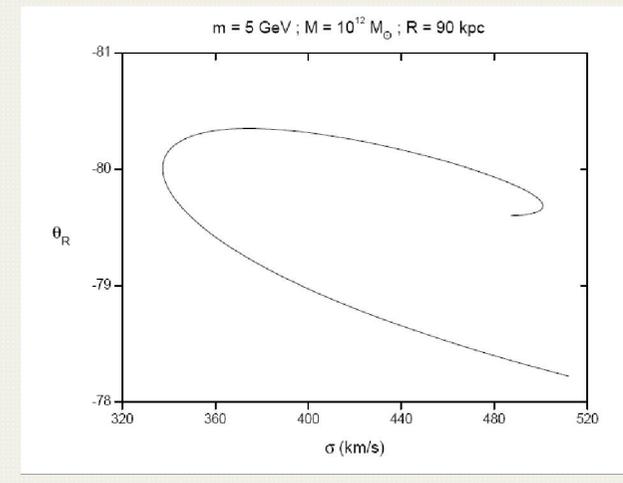
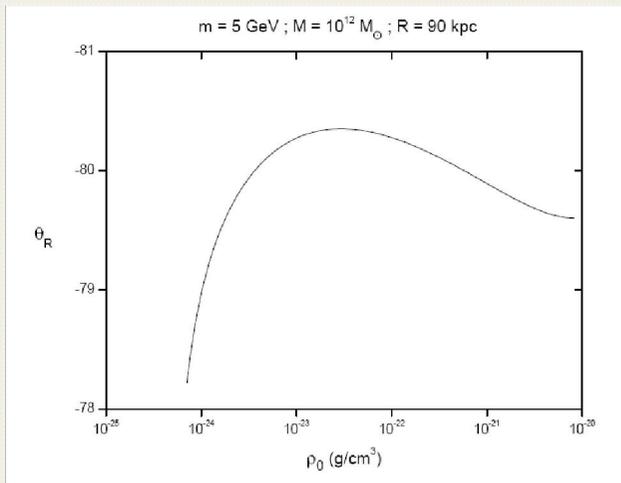
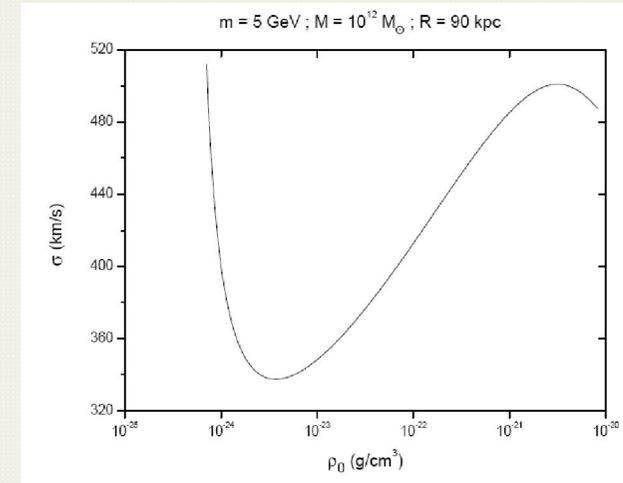
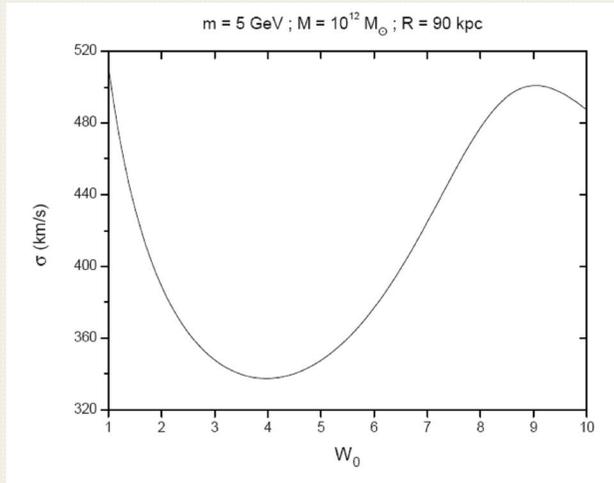
We can fix m , M and R and study the behavior of the other parameters at different values of W_0

$$\rho_0 = \frac{\hat{R}^3 \hat{\rho}_0}{\hat{M}} \frac{M}{R^3}; \quad \theta_R = \frac{1}{2} \ln \left(\frac{\hat{M} \hat{R}^3}{M R^3} \right) + \ln \left(\frac{h^3}{g m^4 G^{3/2}} \right)$$

and, consequently,
$$\sigma = \left(\frac{1}{\hat{\rho}_0} \frac{h^3}{g m^4 e^{\theta_R}} \right)^{1/3} \rho_0^{1/3}$$



Other diagrams



- We obtained the relevant parameters for constructing equilibrium configurations of selfgravitating halos composed by macro component ($m^* \sim \text{GeV}$) deriving from strange dark matter.
- If Λ^* clusters are stable, galactic halos reproducing the same rotation velocity curve in galaxies are possible, in alternative to WIMP-composed halos.
- The halos have a mass $M \sim 10^{12} M_\odot$ and a radius $R \sim 90 \text{kpc}$ in complete accordance with the expected values.
- Galactic halos are completely Newtonian (only Poisson equation is needed), non relativistic (velocity dispersion $\sigma \sim 400 \text{km/s}$) and do not follow quantum statistics ($\theta_R \sim -80$).
- The existence of stable Λ^* clusters, if confirmed, may have strong implications in the standard big bang model.
- The possibility to consider the presence of gravitational field of the galaxy together with visible mass luminosity data can give more stringent constraints to values of the particle mass and W_0 . Finally, a dynamical stability analysis of halos will be considered.



Thank you

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