

Nuclear matter calculations with chiral interactions

Domenico Logoteta

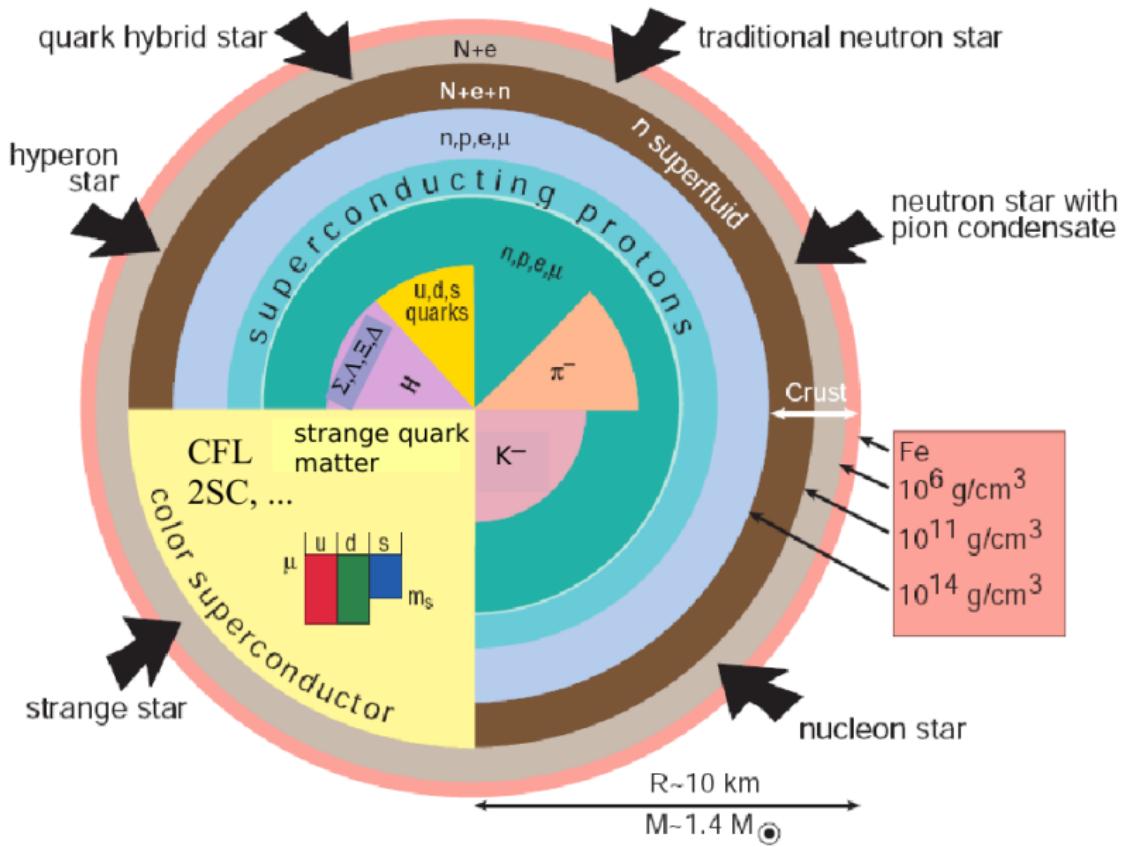
Collaborators: I. Bombaci and A. Kievsky

Frascati

27 novembre 2015

- Chiral interactions
- The Brueckner-Hartree-Fock approach in nuclear matter
- Conclusions

Neutron stars



- **V_{lowk} -approach:**

- K. Hebeler and A. Schwenk, Phys. Rev. C **82**, 014314 (2010).
K. Hebeler, S. K. Bogner, R. J. Furnstahl, A. Nogga and A. Schwenk, Phys. Rev. C **83**, (2011) 031301(R).
K. Hebeler, S. K. Bogner, R. J. Furnstahl, A. Nogga and A. Schwenk, Phys. Rev. C **83**, (2011) 031301.

- **Green's function:**

- A. Carbone, A. Polls and A. Rios Phys. Rev. C **88**, (2013) 044302.

- **Monte Carlo:**

- S. Gandolfi, A. Lovato, J. Carlson, Kevin E. Schmidt, Phys. Rev. C **90**, 061306 (2014).

- **Brueckner-Hartree-Fock:**

- D. Logoteta, I. Vidana, I. Bombaci and A. Kievsky, Phys. Rev. C **91**, (2015) 064001.

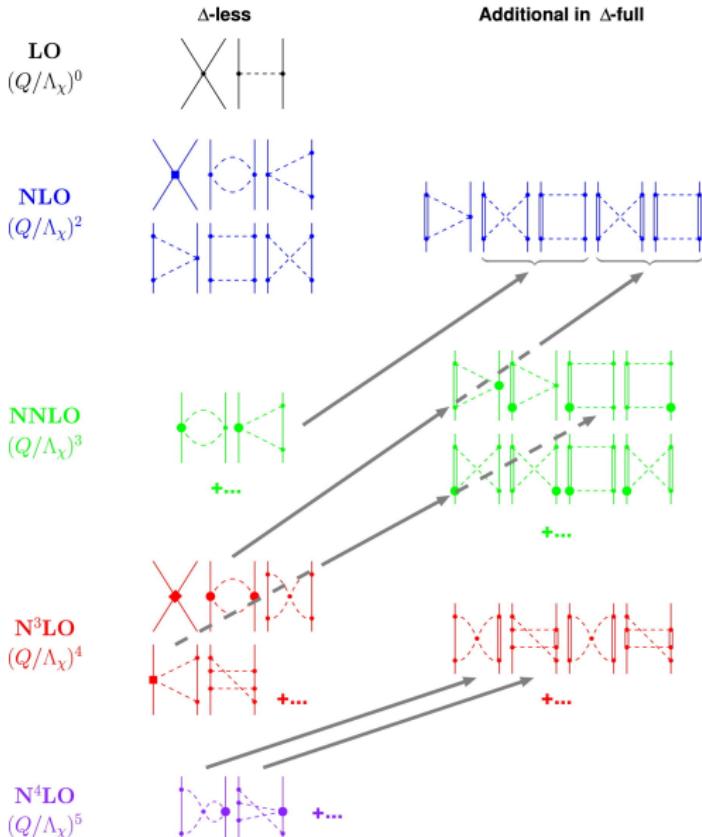
M. Kohno, Phys. Rev. C **88**, 064005 (2013).

- F. Sammarruca, L. Coraggio, J.W. Holt, N. Itaco, R. Machleidt, L. E. Marcucci, Phys. Rev. C **91**, 054311 (2015).

- **NN** potentials: non local N3LO (Idaho-2003), minimal non local N3LO Δ (M. Piarulli-2014 et al.)
- N3LO (Idaho-2003) \Rightarrow in \mathcal{L} included N, π
- N3LO Δ (M. Piarulli-2014 et al.) \Rightarrow in \mathcal{L} included N, π and Δ
- **NNN** potential: local N2LO (P. Navratil 2007)
- Parameters of NNN force fixed in few body calculations \Rightarrow no free parameters

Chiral forces NN

Chiral 2N Force



Chiral 3N Force

LO
 $(Q/\Lambda_\chi)^0$

Δ -less

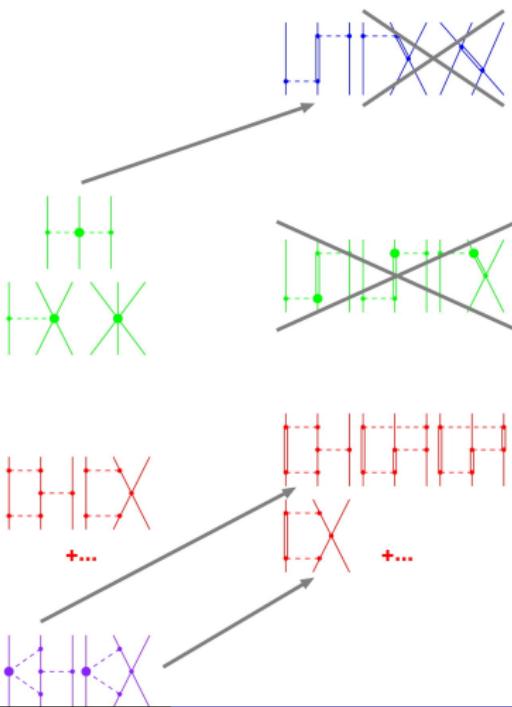
Additional in Δ -full

NLO
 $(Q/\Lambda_\chi)^2$

NNLO
 $(Q/\Lambda_\chi)^3$

N^3LO
 $(Q/\Lambda_\chi)^4$

N^4LO
 $(Q/\Lambda_\chi)^5$



The Brueckner-Hartree-Fock approach

- Starting point: the Bethe-Goldstone equation

$$G(\omega)_{B_1 B_2, B_3 B_4} = V_{B_1 B_2, B_3 B_4} + \sum_{B_i B_j} V_{B_1 B_2, B_i B_j} \times \frac{Q_{B_i B_j}}{\omega - E_{B_i} - E_{B_j} + i\eta} G(\omega)_{B_i B_j, B_3 B_4}$$

$$U_{B_i}(k) = \sum_{B_j} \sum_{\vec{k}'} n_{B_j}(|\vec{k}'|) \times \langle \vec{k} \vec{k}' | G(E_{B_i}(\vec{k}) + E_{B_j}(\vec{k}'))_{B_i B_j, B_i B_j} | \vec{k} \vec{k}' \rangle_{\mathcal{A}}$$

$$E_{B_i}(k) = M_{B_i} + \frac{\hbar^2 k^2}{2M_{B_i}} + \text{Re}[U_{B_i}(k)]$$

$$\epsilon_{BHF} = \frac{1}{V} \sum_{B_i} \sum_{k \leq k_{F_i}} \left[M_{B_i} + \frac{\hbar^2 k^2}{2M_{B_i}} + \frac{1}{2} U_{B_i}(k) \right]$$

- BHF calculations with NNN forces

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- In r-space:

$$W_{eff}(1,2) = Tr_{\sigma_3 \tau_3} \int dx_3 \sum_{cyc} W(1,2,3) n(1,2,3)(1 - P_{13} - P_{23}) \quad (A. Lovato 2011)$$

- In p-space:

$$W_{eff}(1,2) = Tr_{\sigma_3 \tau_3} \int dp_3 \sum_{cyc} W(1,2,3) n(3)(1 - P_{13} - P_{23}) \quad (J. Holt 2010)$$

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- Usually for p-space average \Rightarrow non local cutoff:

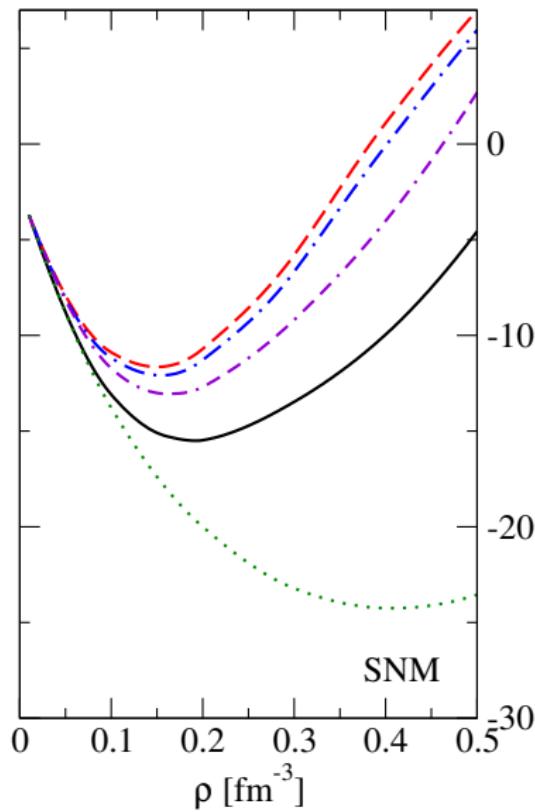
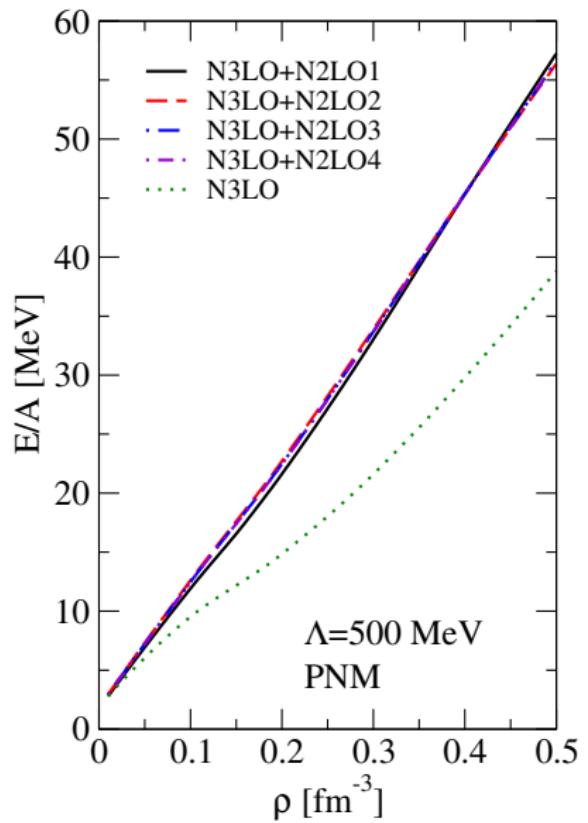
$$F_\Lambda(p', p) = e^{-(p/\Lambda)^{2n} - (p'/\Lambda)^{2n}}$$

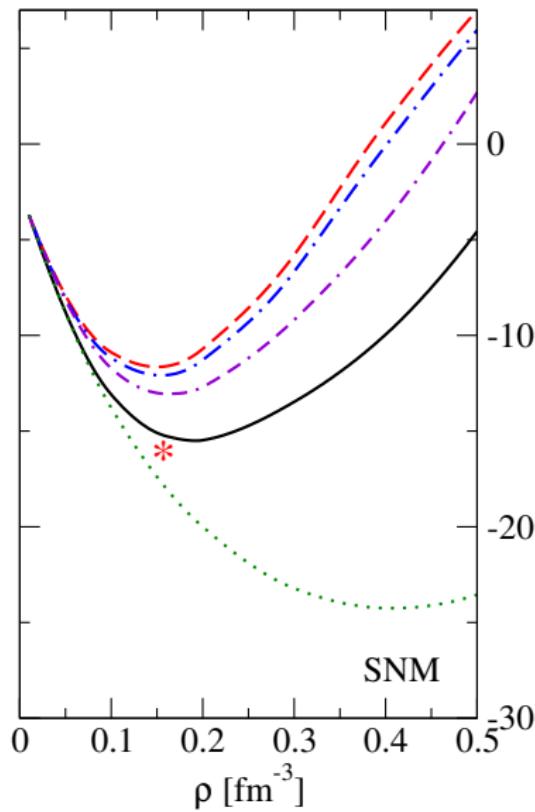
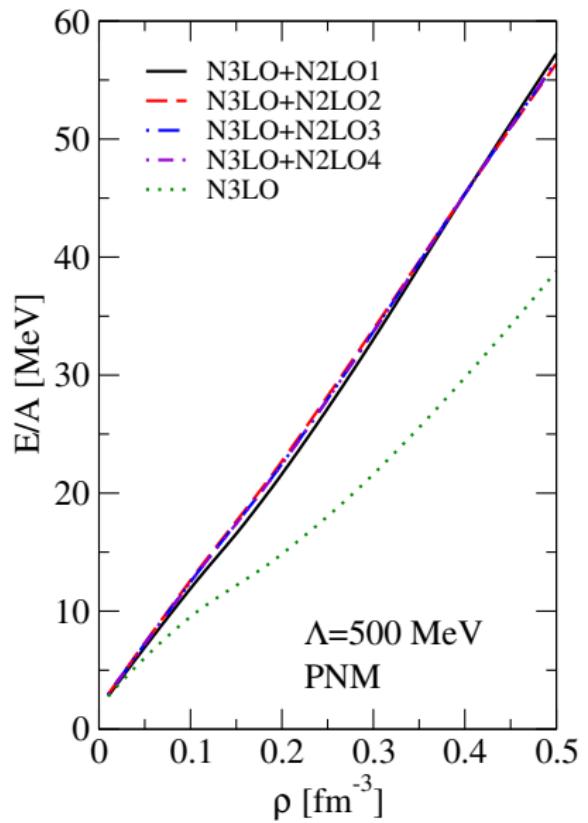
- Following: L. E. Marcucci, A. Kievsky, S. Rosati, R. Schiavilla and M. Viviani Phys. Rev. Lett. **108**, (2012) 052502.
L. Coraggio, J. W. Holt, N. Itaco, R. Machleidt, L. E. Marcucci and F. Sammarruca, Phys Rev. C **89**, (2014) 044321.

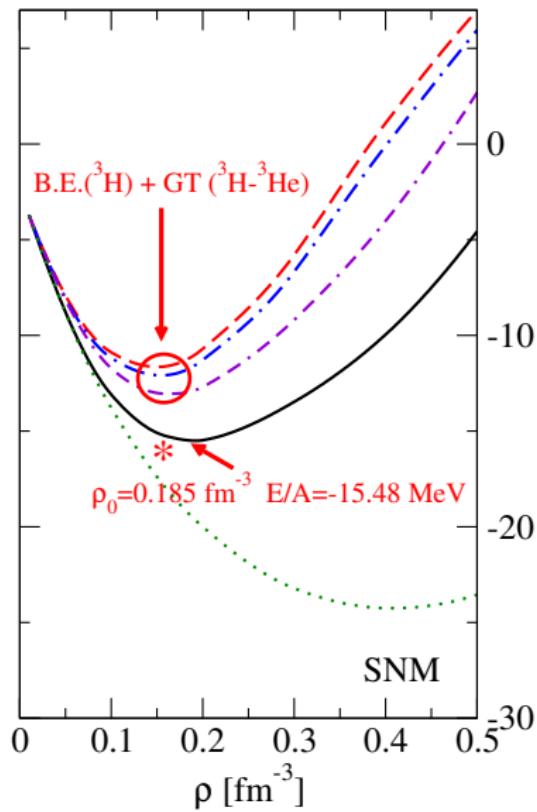
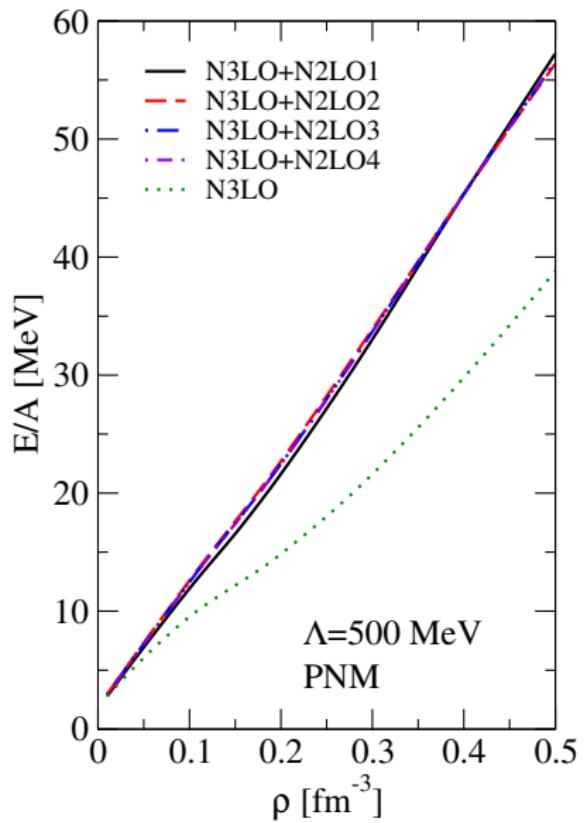


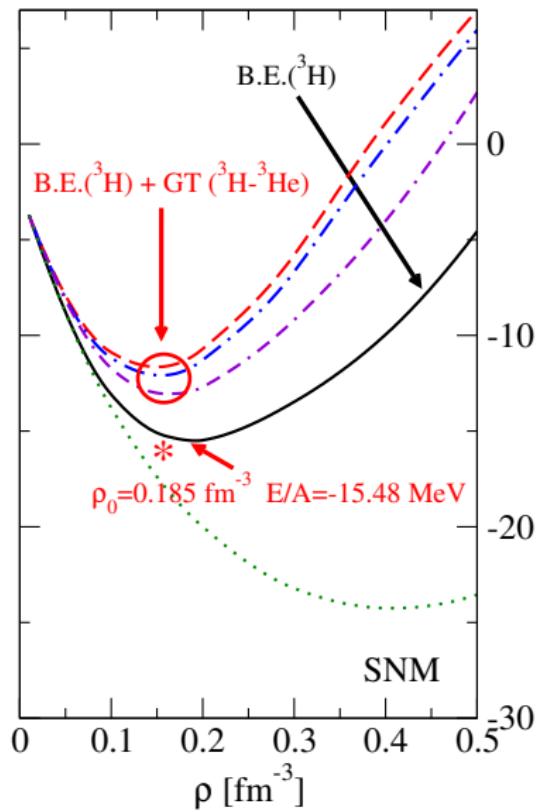
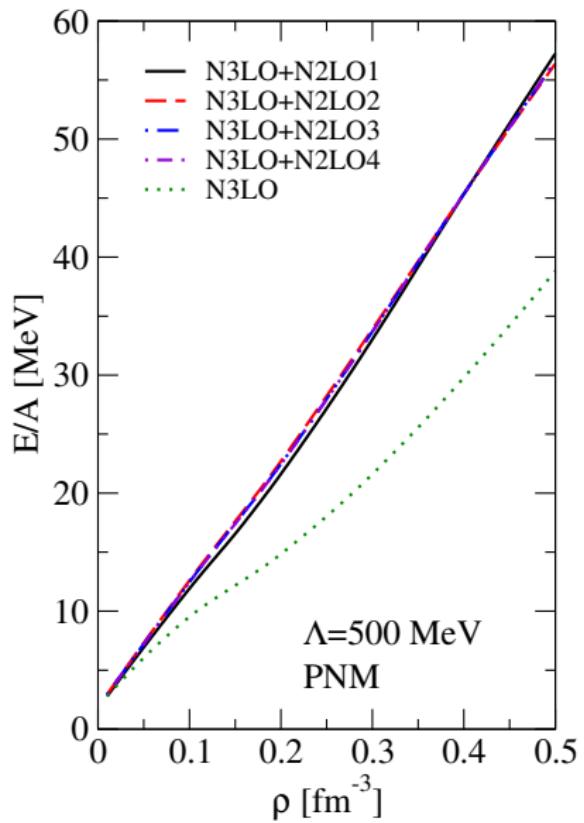
- Low energy constants (c_D, c_E) fixed to reproduce the ${}^3\text{H}$ binding energy + $({}^3\text{H}-{}^3\text{He})$ GT transition matrix element.
- In our many-body calculations we use the same cutoff $F_\Lambda(q^2) = e^{-q^2/\Lambda^2}$ employed in the few-body ones \Rightarrow fully consistent calculation.

	c_D	c_E
N2LOL1 ($\Lambda = 500$)	1.00	-0.029
N2LOL2 ($\Lambda = 500$)	-0.20	-0.208
N2LOL3 ($\Lambda = 500$)	-0.04	-0.184
N2LOL4 ($\Lambda = 500$)	0.0	-0.18

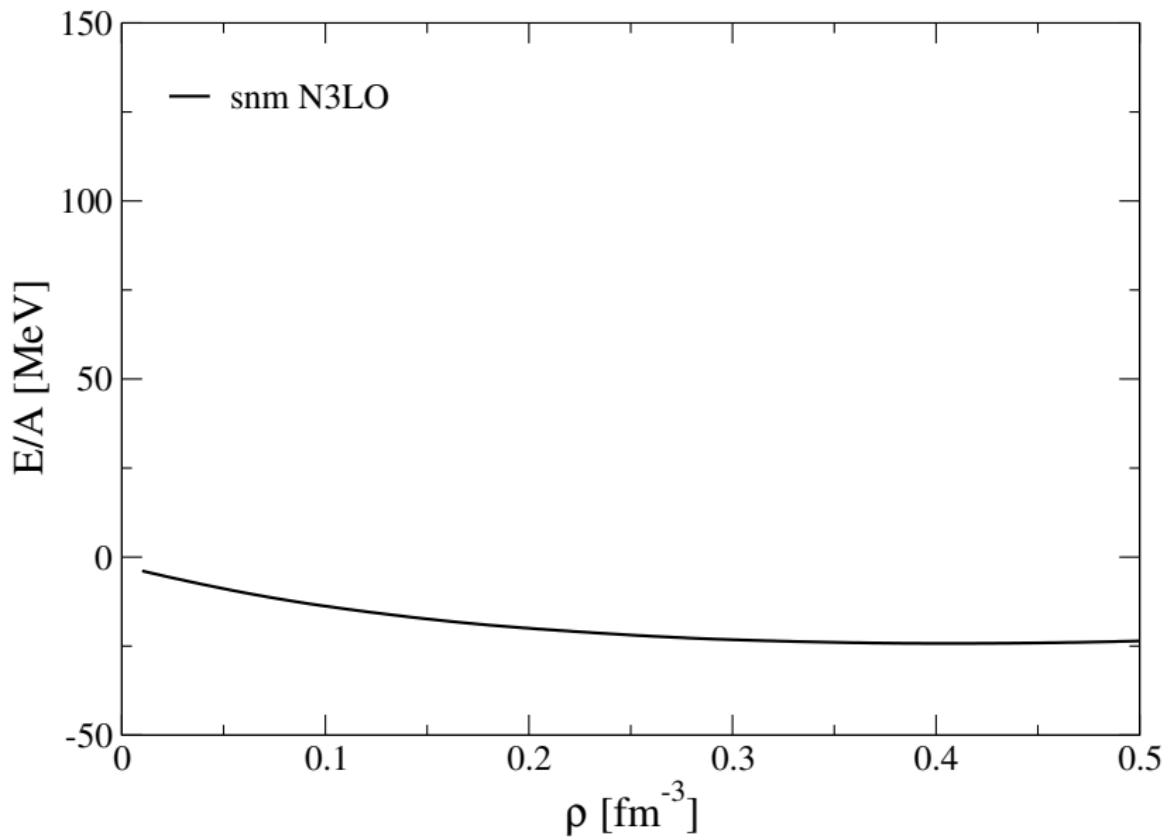




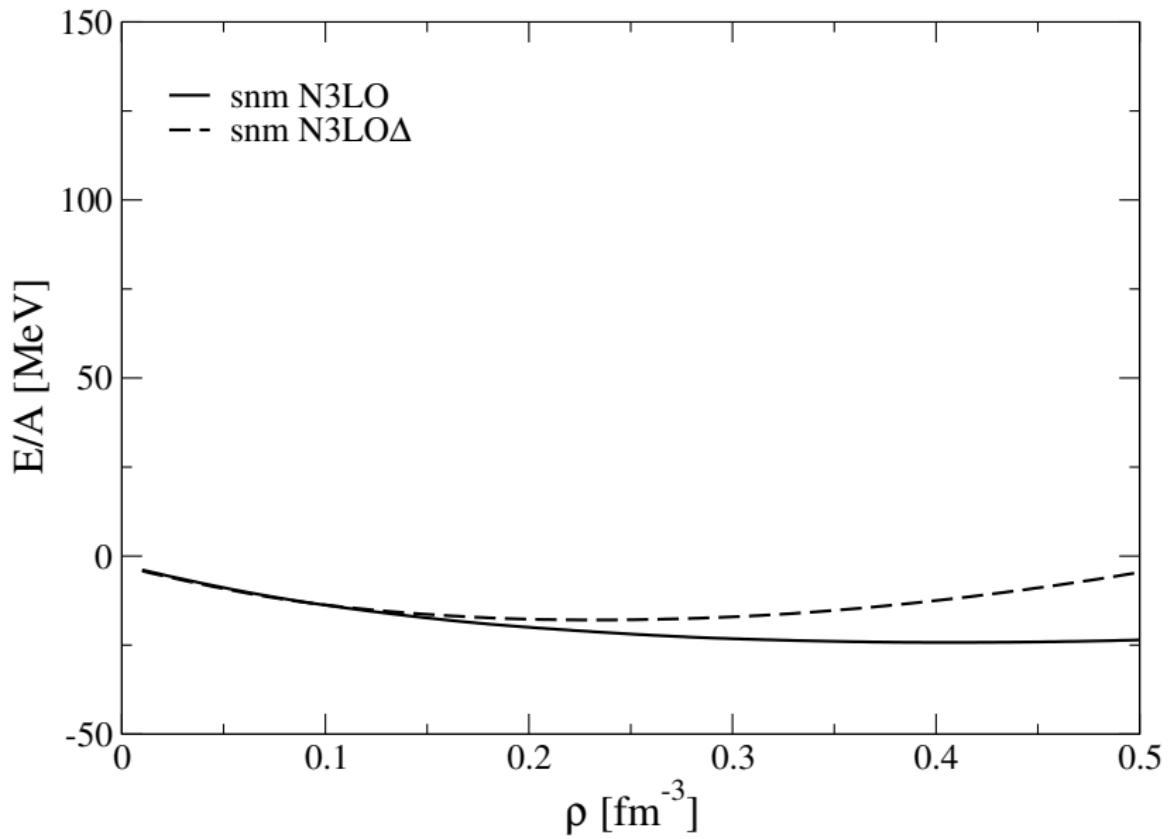




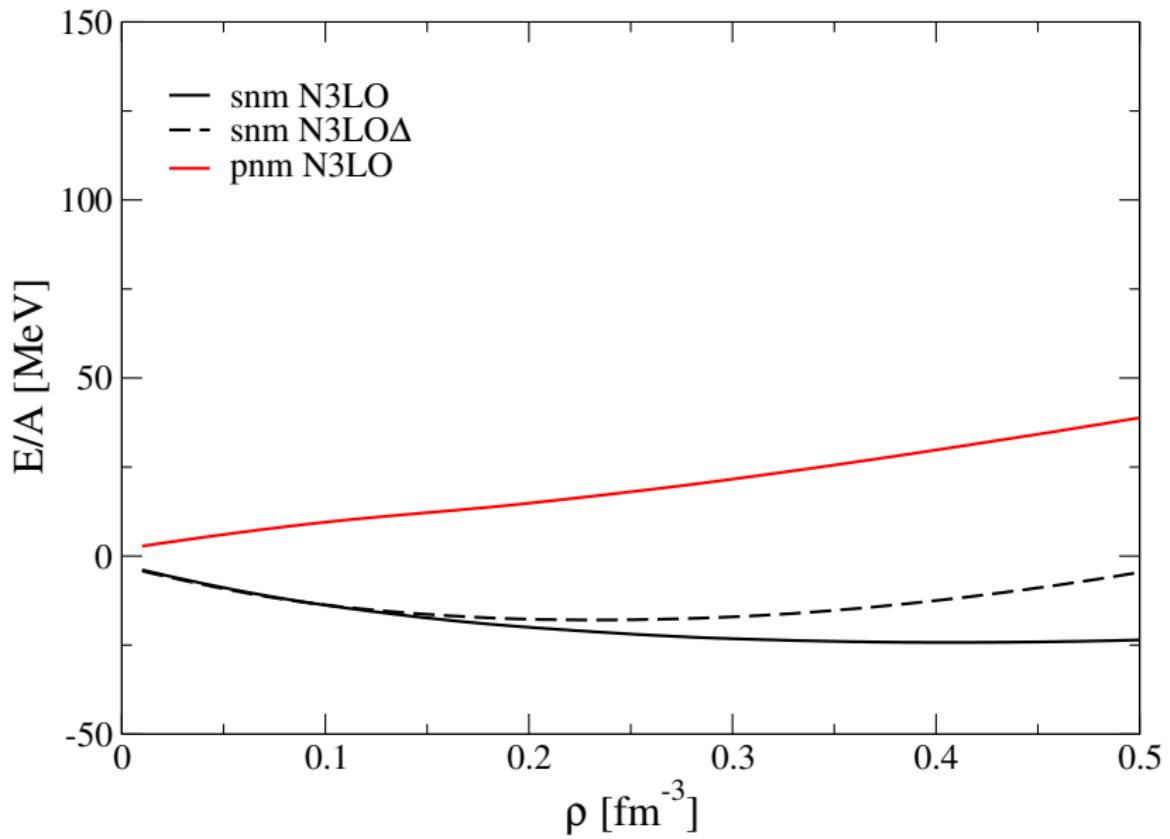
Energy per particle E/A



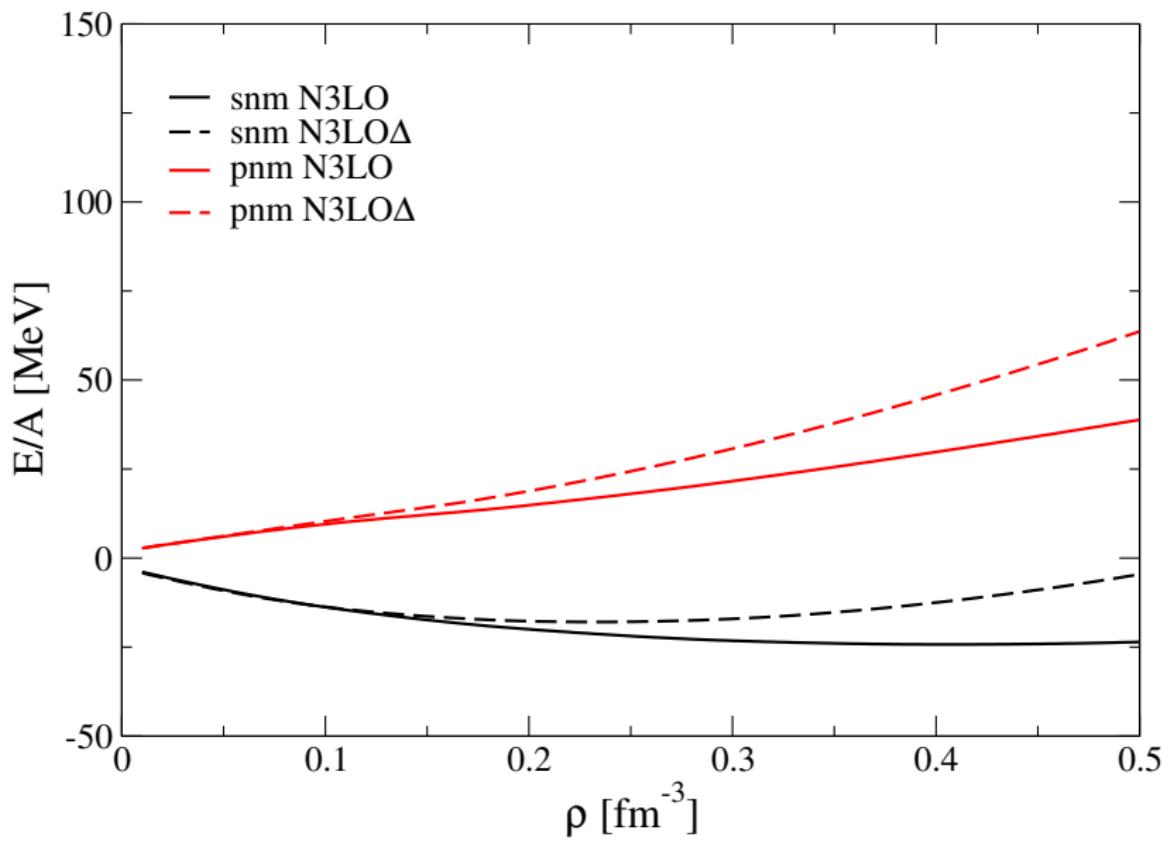
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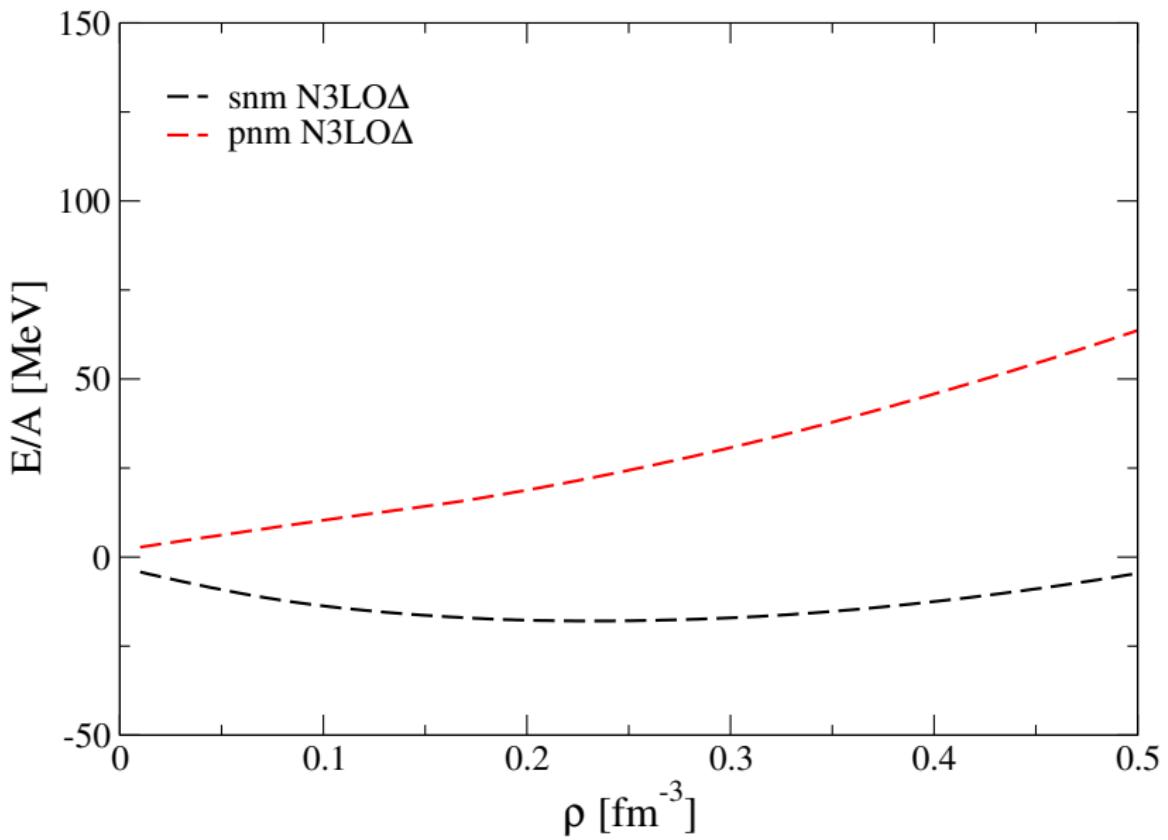
Energy per particle E/A



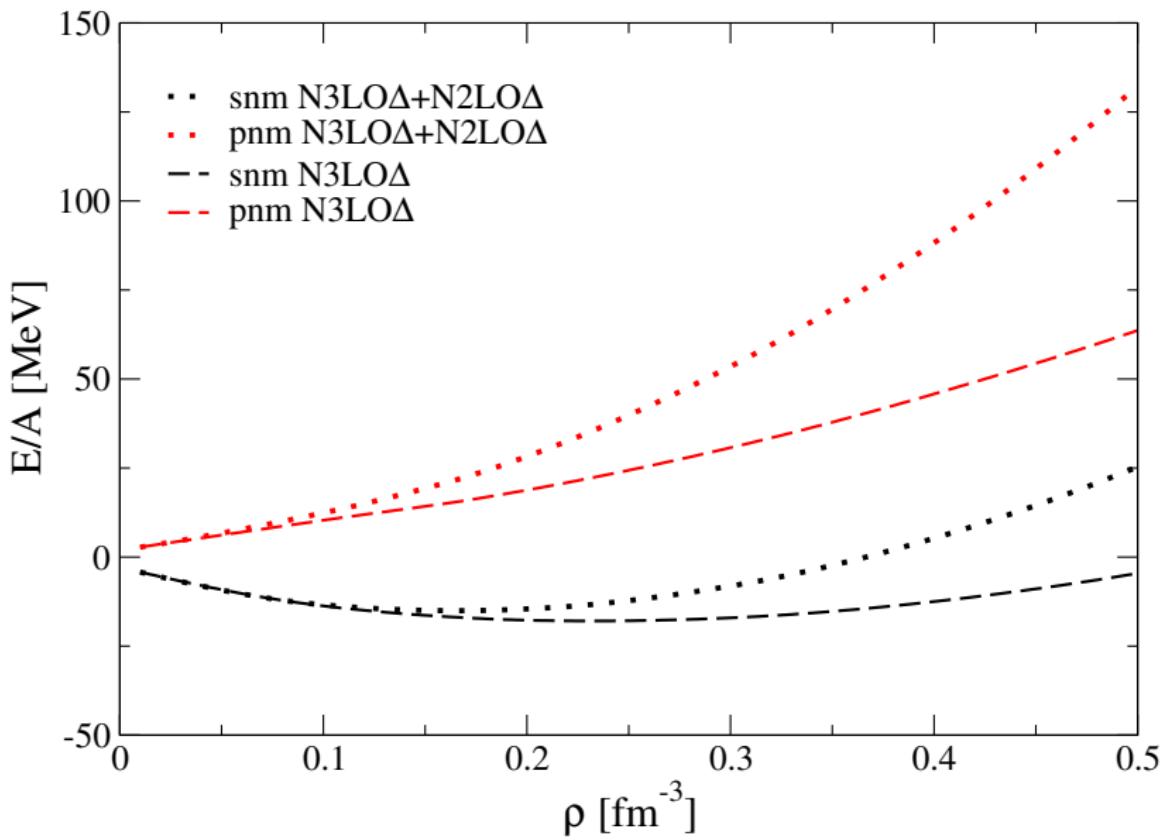
Energy per particle E/A



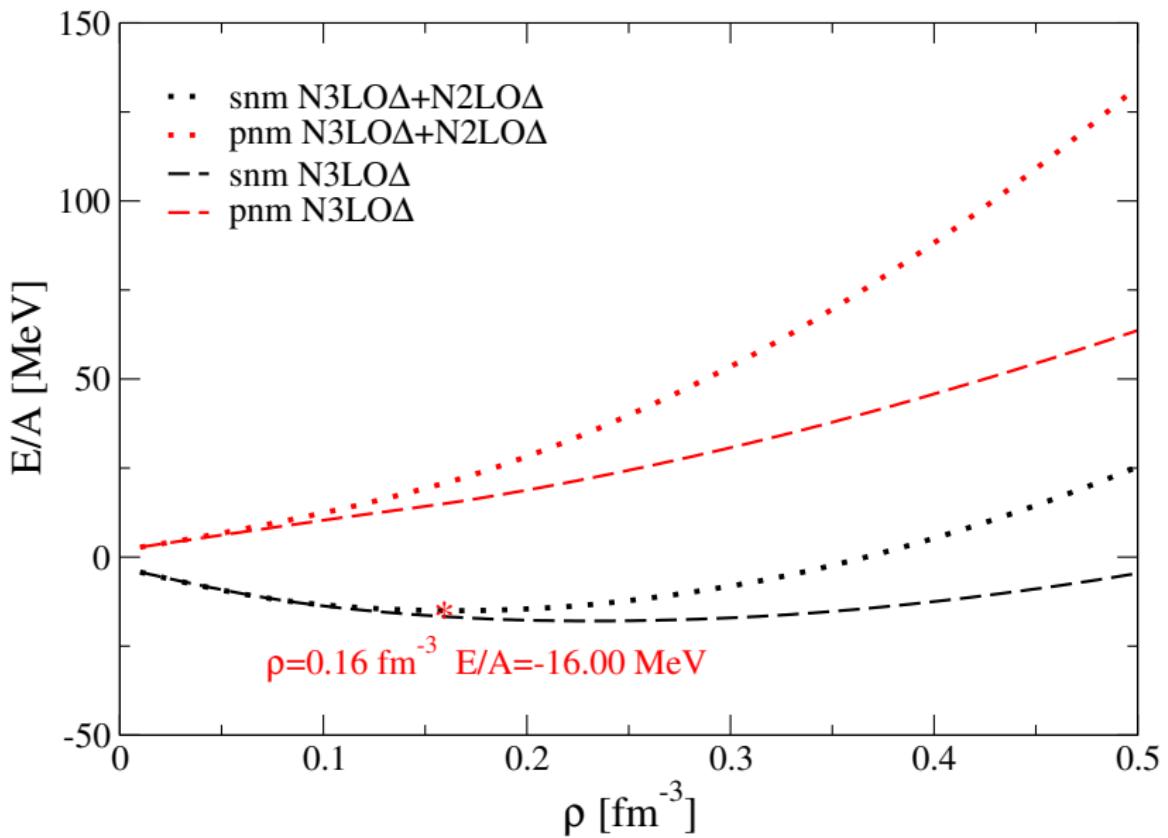
Energy per particle E/A *p*-space average



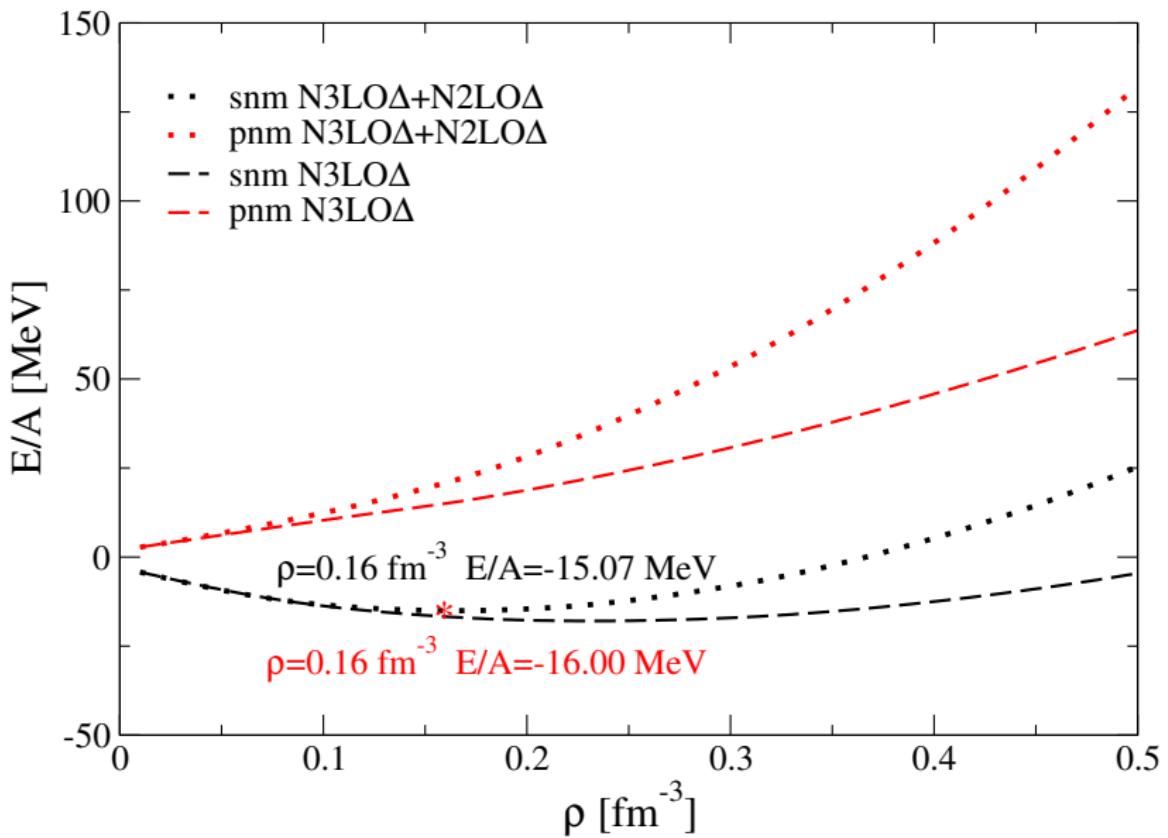
Energy per particle E/A *p*-space average



Energy per particle E/A *p*-space average



Energy per particle E/A *p*-space average



- Microscopic calculations of nuclear matter based on realistic interaction can help us to understand discrepancies between many-body and few-body nuclear physics.
- New generation of interactions based on chiral perturbation theory provide very interesting results.
- Next to do ⇒ study of hyperonic matter based on chiral forces.



- Problem of maximum mass on neutron stars.