SURPRISES ON THE X(3872) AD Polosa INFN Roma

CHARMONIUM LEVELS

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... but then (2004/5) the red dots came about

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Discovery of the X(3872): B decays



::Belle, Phys. Rev. Lett. 91 (2003) 262001 [arXiv:hep-ex/0309032] ::

Discovery of the X(3872): pp̄ collisions





AFTER THE X MANY MORE

A Chevrolin Contractions and the Chevroline

state	$M ({ m MeV})$	Γ (MeV)	J^{PC}	Seen In	Observed by:	Comments
$Y_{s}(2175)$	2175 ± 8	58 ± 26	$1^{}$	$(e^+e^-)_{ISR}, J/\psi \to Y_s(2175) \to \phi f_0(980)$	BaBar, BESII, Belle	
X(3872)	3871.4 ± 0.6	< 2.3	1^{++}	$B \to KX(3872) \to \pi^+\pi^- J/\psi, \gamma J/\psi, D\bar{D^*}$	Belle, CDF, D0, BaBar	Molecule?
X(3915)	3914 ± 4	28^{+12}_{-14}	?++	$\gamma\gamma ightarrow \omega J/\psi$	Belle	
Z(3930)	3929 ± 5	29 ± 10	2^{++}	$\gamma\gamma ightarrow Z(3940) ightarrow Dar{D}$	Belle	$2^3P_2(car c)$
X(3940)	3942 ± 9	37 ± 17	$0^{?+}$	$e^+e^- \to J/\psi X(3940) \to D\bar{D^*} \text{ (not } D\bar{D} \text{ or } \omega J/\psi)$	Belle	$3^1S_0(car c)?$
Y(3940)	3943 ± 17	87 ± 34	$?^{?+}$	$B \to KY(3940) \to \omega J/\psi \text{ (not } D\bar{D^*})$	Belle, BaBar	$2^3P_1(car c)?$
Y(4008)	4008^{+82}_{-49}	$226\substack{+97\\-80}$	$1^{}$	$(e^+e^-)_{ISR} \rightarrow Y(4008) \rightarrow \pi^+\pi^- J/\psi$	Belle	
Y(4140)	4143 ± 3.1	$11.7^{+9.1}_{-6.2}$	$?^{?}$	$B \to KY(4140) \to J/\psi\phi$	CDF	
X(4160)	4156 ± 29	$139\substack{+113 \\ -65}$	$0^{?+}$	$e^+e^- \rightarrow J/\psi X(4160) \rightarrow D^*\bar{D^*} \pmod{D\bar{D}}$	Belle	
Y(4260)	4264 ± 12	83 ± 22	$1^{}$	$(e^+e^-)_{ISR} \rightarrow Y(4260) \rightarrow \pi^+\pi^- J/\psi$	BaBar, CLEO, Belle	Hybrid?
Y(4350)	4324 ± 24	172 ± 33	$1^{}$	$(e^+e^-)_{ISR} \rightarrow Y(4350) \rightarrow \pi^+\pi^-\psi'$	BaBar	
Y(4350)	4361 ± 13	74 ± 18	$1^{}$	$(e^+e^-)_{ISR} \rightarrow Y(4350) \rightarrow \pi^+\pi^-\psi'$	Belle	
Y(4630)	$4634\substack{+9.4\\-10.6}$	92^{+41}_{-32}	$1^{}$	$(e^+e^-)_{ISR} ightarrow Y(4630) ightarrow \Lambda_c^+ \Lambda_c^-$	Belle	
Y(4660)	4664 ± 12	48 ± 15	$1^{}$	$(e^+e^-)_{ISR} ightarrow Y(4660) ightarrow \pi^+\pi^-\psi'$	Belle	
$Z_1(4050)$	4051^{+24}_{-23}	82^{+51}_{-29}	?	$B ightarrow KZ_1^{\pm}(4050) ightarrow \pi^{\pm} \chi_{c1}$	Belle	
$Z_2(4250)$	4248^{+185}_{-45}	177^{+320}_{-72}	?	$B \to KZ_2^{\pm}(4250) \to \pi^{\pm}\chi_{c1}$	Belle	
Z(4430)	4433 ± 5	45^{+35}_{-18}	?	$B \to KZ^{\pm}(4430) \to \pi^{\pm}\psi'$	Belle	
$Y_b(10890)$	$10,890\pm3$	55 ± 9	$1^{}$	$e^+e^- ightarrow Y_b ightarrow \pi^+\pi^-\Upsilon(1,2,3S)$	Belle	

From Godfrey arXiv:0910.3409

 $X \to \gamma J/\psi \mapsto C = +1$ and $X \to \rho^0 J/\psi \to (\pi^+ \pi^-)_S J/\psi \mapsto P = 1$

Radiative decays

:: BaBar, Phys.Rev.Lett.102:132001,2009. ::

 $\mathcal{B}(B^{\pm} \to X(3872)K^{\pm}) \times \mathcal{B}(X(3872) \to J/\psi\gamma) = (2.8 \pm 0.8(\text{stat}) \pm 0.1(\text{syst})) \times 10^{-6}$

:: BaBar, Phys. Rev.D, 77 (2008) 111101 ::

 $\mathcal{B}(B^{\pm} \to X(3872)K^{\pm}) \times \mathcal{B}(X(3872) \to J/\psi\pi^{+}\pi^{-}) = (8.4 \pm 1.5(\text{stat}) \pm 0.7(\text{syst})) \times 10^{-6}$

 $\frac{\mathcal{B}(X(3872) \to J/\psi\gamma)}{\mathcal{B}(X(3872) \to J/\psi\pi^+\pi^-)} = (0.3 \pm 0.1)$

 $\Gamma(2^{3}P_{1}\rightarrow\psi\gamma)/\Gamma(2^{3}P_{1}\rightarrow\psi\pi\pi)\sim40$

Charmonium predictions: :: Eichten, Lane and Quigg, Phys. Rev. D 69, 094019 (2004) :: :: Barnes and Godfrey, Phys. Rev. D, 69, 054008 (2004) ::



THE MOLECULE EXPLANATION

Contraction and a second a first state

A way for accomodating isospin violations is to suppose that X(3872) is a DD* molecule

$$|X> = \frac{|D^0 \bar{D}^{0*}> + |\bar{D}^0 D^{0*}>}{\sqrt{2}}$$

So that one can produce both the I=0 or I=1 combinations. In $B \rightarrow KX$ decay indeed can have both

Braaten; Mehen; Swanson; Hanhart; Tornqvist; Voloshin (hadrocharmonium)

THE MOLECULE EXPLANATION

B(X→ψρ)/B(X→ψω)~1

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HOW LARGE IS A MOLECULE?

Contractions and a Character

Using the indetermination principle

$$\frac{\hbar^2}{2\mu r_0^2} - \frac{g^2}{4\pi} \frac{e^{-\frac{m_\pi c}{\hbar}r_0}}{r_0} = \mathcal{E} = M_D + M_{D^*} - M_X \sim 0.25 \text{ MeV}$$

and the fact that $g^2/4\pi \sim 10$ we have a characteristic size

 $r_0 \sim 8 \text{ fm}$

which could either be infinity(!) as

 $\mathcal{E} = 0.25 \pm 0.40 \text{ MeV}$

SOME PROBLEMS

- Such a large state must spend part of its time in a very tight configuration to allow J/ψ ρ and J/ψ ω decays - one can model this but it is difficult to 'reliably' calculate...
- D* has a width of about 70 keV. Why the X→DDπ decay rate is ~3 MeV then? In other words, which is the mechanism accelerating the D* decay into the molecule?

X PROMPT CROSS SECTION

At the moment CDF (which has performed the most precise determination of the X(3872) mass) measures a prompt X production cross section of about 30-70 nb.

We find that the prompt production cross section is an important table test of the molecule interpretation.

Bignamini, Grinstein, Piccinini, Polosa, Sabelli Phys Rev Lett 2009

How can a loosely bound state (almost zero binding energy) be poduced in the wild environment of hadronic collisions?

Beppe Nardulli once told me: "Sometimes it is more profitable to have important enemies than important friends..."

PROMPT PRODUCTION

$$\sigma(p\bar{p} \to X(3872))$$

 \sim



$$\begin{split} & \left| \int d^{3}k \langle X | D\bar{D}^{*}(\mathbf{k}) \rangle \langle D\bar{D}^{*}(\mathbf{k}) | p\bar{p} \rangle \right|^{2} \\ & \left| \int_{\mathcal{R}} d^{3}k \langle X | D\bar{D}^{*}(\mathbf{k}) \rangle \langle D\bar{D}^{*}(\mathbf{k}) | p\bar{p} \rangle \right|^{2} \\ & \int_{\mathcal{R}} d^{3}k | \psi(\mathbf{k}) |^{2} \int_{\mathcal{R}} d^{3}k | \langle D\bar{D}^{*}(\mathbf{k}) | p\bar{p} \rangle |^{2} \\ & \int_{\mathcal{R}} d^{3}k | \langle D\bar{D}^{*}(\mathbf{k}) | p\bar{p} \rangle |^{2} \end{split}$$

Using Pythia & Herwig we can compute

$$\sigma_{\max}(p\bar{p} \to X(3872)) = \int_{\mathcal{R}} d^3k |\langle D\bar{D}^*(\mathbf{k})|p\bar{p}\rangle|^2$$

where $\mathcal{R} \sim [0.40]$ MeV
as $k \sim \sqrt{2\mu(-0.25+0.40)} \simeq 17$ MeV

TUNING MC'S



[The D⁰ D^{*-} pair cross section as function of $\Delta \phi$ at CDF Run II. We find that we have to rescale the Herwig cross section values by a factor K= 1.8 to best fit the data on open charm production. As for Pythia we need K=0.74]



COUNTING PAIRS OVER 5*10**9 SIMULATED EVENTS



Bignamini, Grinstein, Piccinini, Polosa, Sabelli Phys Rev Lett 2009

 $\sigma_{\rm th} \sim 1/300 \; \sigma_{\rm exp}$

One needs to integrate cross section up to about 205 MeV with Herwig and 130 MeV with Pythia in order to reach the experimental value. We thus EXCLUDE any molecular interpretation of X(3872).

ARE XYZ DIQUARK-ANTIDIQUARK PARTICLES?

Suppose that X(3872) is a new type of meson with a body-plan different from standard ones, namely a diquark-antidiquark meson

Maiani, Piccinini, Polosa, Riquer Phys Rev D 2005

A diquark qq has the same color as an antiquark (there is attraction between two quarks in this color channel)

• This would explain easily $J/\psi \rho$ and $J/\psi \omega$ decays

With a bit more effort it would also explain the observed isospin violations

 $J/\psi \gamma$ is understood! (only with 4q!)

THE MOST CHALLENGING PREDICTIONS

TWO NEUTRAL X are PREDICTED with a separation in mass of few MeV (of the order of m_u-m_d)

All the mesons (red dots) found in the first table are NEUTRAL particles whereas the tetraquark model predicts also the existence of CHARGED particles like

 $[cu][\bar{c}\bar{d}] \quad Q = +1$ $[cu][\bar{d}\bar{s}] \quad Q = +2$

Light states like

 $[uu][\bar{d}\bar{s}] \quad Q = +2$

are disfavored as the spin one light diquark is itself disfavored

BUT THEN A NEW PAPER BY BABAR CAME...

arXiv:1005.5190

The quantum numbers of the X are not 1⁺⁺ but 2⁻⁺ !!

This is why we cannot justify the prompt production of X at CDF. There is no I++ (DD*) molecule.

O(200) papers on the molecular interpretation...

THE SIMPLEST QCD STRING

Selem and Wilczek hep-ph/0602128 (on the Chew-Frautschi model)

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Loosely bound (S-wave) molecules cannot have orbital or radial excitations. Two diquarks bound by a string can.





$$\mathcal{E} = \frac{m_1}{\sqrt{1 - (\omega r_1)^2}} + \frac{m_2}{\sqrt{1 - (\omega r_2)^2}} + \frac{\sigma}{2\pi\omega} \int_0^{\omega r_1} \frac{dv}{\sqrt{1 - v^2}} + \frac{\sigma}{2\pi\omega} \int_0^{\omega r_2} \frac{dv}{\sqrt{1 - v^2}}$$
$$\ell = \frac{\omega r_1^2 m_1}{\sqrt{1 - (\omega r_1)^2}} + \frac{\omega r_2^2 m_2}{\sqrt{1 - (\omega r_2)^2}} + \frac{\sigma}{2\pi\omega^2} \int_0^{\omega r_1} \frac{dv v^2}{\sqrt{1 - v^2}} + \frac{\sigma}{2\pi\omega^2} \int_0^{\omega r_2} \frac{dv v^2}{\sqrt{1 - v^2}}$$

 $\sigma \sim 1 \text{ GeV}^2$ from Regge slopes and $d\mathcal{E}'/dr' = T = \sigma/2\pi$ and $\frac{m_i \omega^2 r_i}{1 - (\omega r_i)^2} = T$

THE SIMPLEST QCD STRING II

Cotugno, Faccini, Polosa, Sabelli Phys Rev Lett 2010

In the limit of infinite (heavy-light) quark mass we find

$$\mathcal{E}(r)_{\text{TOT}} \simeq 2M + \frac{3}{(16\pi^2 M)^{1/3}} (\sigma \ell)^{2/3} + A\left(\vec{S} \cdot \vec{\ell}\right) \frac{1}{r} \frac{d}{dr} \mathcal{E}(r)$$

 $\mathcal{E}(r) \text{ as } T/(M\omega) \rightarrow 0$

No spin-spin because of large r and M



(charm)

(beauty)

THE NEXT-TO-SIMPLEST QCD STRING

Burns, Piccinini, Polosa, Sabelli -w.i.p.-

$$\mathcal{E}(r)_{\text{TOT}} \simeq 2\mathbf{M} + \frac{3}{(16\pi^2 \mathbf{M})^{1/3}} (\boldsymbol{\sigma}\ell)^{2/3} + \mathbf{A} \left(\vec{S} \cdot \vec{\ell}\right) \frac{1}{r} \frac{d}{dr} \mathcal{E}(r) + \mathbf{B} \left[\vec{S}^2 - 3\left(\vec{S} \cdot \vec{n}\right)^2\right]$$

 $\mathcal{E}(r)$ as $T/(M\omega){\rightarrow}0$



D-WAVE BOTTOMONIUM

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BaBar :: Moriond 2010

CHARMONIUM

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The situation for charmonium is a bit more tricky since the `infinite` mass limit is less appropriate here.



IS THE X(3872) A $^{1}D_{2}$ (2⁻⁺) CHARMONIUM?!



Maybe isospin violations mentioned are not the main problem, but what about radiative decays? J/ $\psi\gamma$ and J/ $\psi\rho$ would be P-wave decays - but then why $\frac{\mathcal{B}(X(3872) \rightarrow J/\psi\gamma)}{\mathcal{B}(X(3872) \rightarrow J/\psi\pi^+\pi^-)} = (0.3 \pm 0.1)$

FRAGMENTATION OF A GLUON IN A ¹D₂

Cho and Wise hep-ph/9410214

$$\frac{d\sigma}{dp_{\perp}}(p\bar{p} \to 1^{1}D_{2} + \text{All}) = \sum_{h=0}^{2} \int_{0}^{1} dz \frac{d\sigma}{dp_{\perp}}(p\bar{p} \to g(p_{\perp}/z) + \text{All};\mu) \times D_{g \to 1^{1}D_{2}^{(h)}}(z;\mu)$$

 $x \simeq \sqrt{M_{\perp}^2/1960^2} \simeq 0.02$ $p_{\perp} \gtrsim 5 \text{ GeV}$ $|y| \leq 6$ factorization scale $\mu \simeq M_{\perp}$

Updating the pdf's we find

$$\sigma(p\bar{p} \to 1^1 D_2) = 0.6 \text{ nb}$$



Still very small w/ respect to the prompt production at CDF (??)

OUTLOOK

- There are about <u>20 newly found narrow resonances</u> resembling standard charmonia but often evading the well known charmonium features.
- Belle has observed a Z⁺(4430) decaying into charmonium + charged pion! This is a highly exotic state. What is it? (A Tetraquark?)
- Are we really observing hadrons with a different `body plan`? We think so.
- More urgently: is the X just D-wave charmonium?

THREE MORE (NEW) STATES

The second distribution of the

	state	$M ({ m MeV})$	Γ (MeV)	J^{PC}	Seen In	Observed by:	Comments
	$Y_{s}(2175)$	2175 ± 8	58 ± 26	$1^{}$	$(e^+e^-)_{ISR}, J/\psi \to Y_s(2175) \to \phi f_0(980)$	BaBar, BESII, Belle	
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	X(3940)	3942 ± 9	37 ± 17	$0^{?+}$	$e^+e^- \to J/\psi X(3940) \to D\bar{D^*} \text{ (not } D\bar{D} \text{ or } \omega J/\psi X(3940)$	ψ) Belle	$3^1S_0(car c)?$
	Y(3940)	3943 ± 17	87 ± 34	$?^{?+}$	$B \to KY(3940) \to \omega J/\psi \text{ (not } D\bar{D^*})$	Belle, BaBar	$2^{3}P_{1}(c\bar{c})?$
	Y(4008)	4008^{+82}_{-49}	$226\substack{+97\\-80}$	$1^{}$	$(e^+e^-)_{ISR} \rightarrow Y(4008) \rightarrow \pi^+\pi^- J/\psi$	Belle	
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• • • •	X(4160)	4156 ± 29	139^{+113}_{-65}	$0^{?+}$	$e^+e^- \to J/\psi X(4160) \to D^*\bar{D^*} \pmod{D\bar{D}}$	Belle	
	Y(4260)	4264 ± 12	83 ± 22	1	$(e^+e^-)_{ISR} \rightarrow Y(4260) \rightarrow \pi^+\pi^- J/\psi$	BaBar, CLEO, Belle	Hybrid?
	Y(4350)	4324 ± 24	172 ± 33	$1^{}$	$(e^+e^-)_{ISR} \rightarrow Y(4350) \rightarrow \pi^+\pi^-\psi'$	BaBar	
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	$Y_b(10890)$	$10,890\pm3$	55 ± 9	1	$e^+e^- \rightarrow Y_b \rightarrow \pi^+\pi^-\Upsilon(1,2,3S)$	Belle	

From Godfrey arXiv:0910.3409

TWO ARE THE SAME: YB

Cotugno, Faccini, Polosa, Sabelli Phys Rev Lett 2010

We just redo the fits using Belle data under the hypothesis that Y(4660) and Y(4630) are the same (1⁻⁻) particle. This hypothesis improves the fit







What we find surprising is that

$$\frac{\mathcal{B}(Y_B \to \Lambda_c \bar{\Lambda}_c)}{\mathcal{B}(Y_B \to \psi(2S)\pi^+\pi^-)} = 24.6 \pm 6.6$$