

IDENTIFICATION OF WARPED EXTRA DIMENSIONS SIGNATURES AT LHC

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Outline

- Motivations
- Introduction to RS1 spin-2 KK graviton excitations
- Angular distributions for $p + p \rightarrow l^+l^- + X$
- Spin-2 identification potential of A_{CE}
- Extension to $p + p \rightarrow \gamma\gamma + X$
- Results for identification reach on spin-2
- Phenomenology of the radion

Motivations: hierarchies

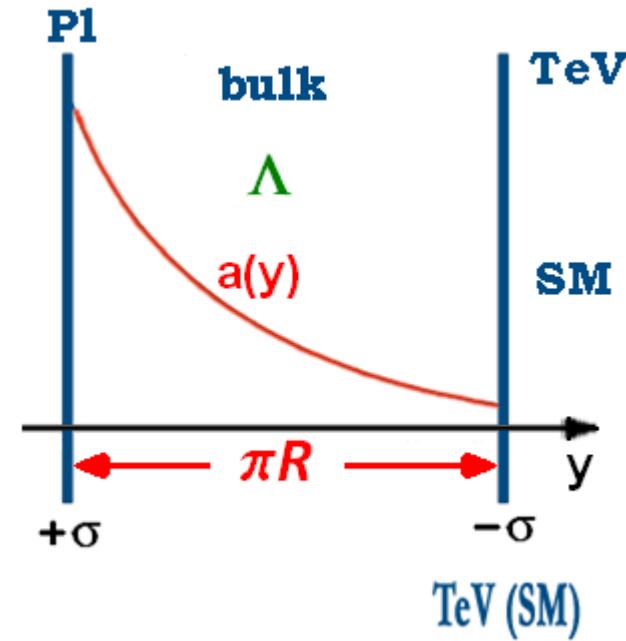
- Gauge hierarchy: $M_{EW} \sim 1 \text{ TeV} \ll M_{Pl} \sim 10^{16} \text{ TeV}$
- Fermion mass hierarchy: $m_\nu \ll m_{top}$
- Scale of supersymmetry (soft breaking)
- “Duality” to strongly coupled field theories

Simplest model [RS1] (Randall-Sundrum)

- 1 compactified “warped” extra dimension
- 5D spacetime, 2 branes set up
- Only gravity can propagate in full 5D space

RS1 model of gravity in extra dimensions

- $ds^2 = (a(y))^2 \eta_{\mu\nu} dx^\mu dx^\nu + dy^2$
warp factor
- RS solution: $a(y) = e^{-k|y|}$; $k > 0$ (5D curvature)
- Fundamental mass scales (G = Newton constant)
 $G^{(4)} \propto M_{Pl}^{-2}$; $G^{(5)} \propto M_*^{-3}$; k .
- 4D effective theory gives $M_{Pl}^2 = \frac{M_*^3}{k} (1 - e^{-2k\pi R})$
- “Natural”: $M_{Pl} \sim M_* \sim k$
- Mass spectrum on Planck brane $\propto M_* \sim 10^{16} \text{ TeV}$
- Mass spectrum on TeV brane exponentially warped down: $M_H \propto M_* e^{-k\pi R} \sim \text{TeV}$ for $kR \sim 10!$ (automatic)



$$x^M \equiv (x^\mu, y)$$

$$ds^2 = G_{MN} dx^M dx^N$$

$$\eta_{\mu\nu} = \begin{pmatrix} -1 & . & . & . \\ . & 1 & . & . \\ . & . & 1 & . \\ . & . & . & 1 \end{pmatrix}$$

- Exponential twist of hierarchy problem
- Junction conditions at $y=0, y=\pi R \Rightarrow$
tower of KK, spin-2, graviton excitations h_n
- Masses and couplings to SM matter (TeV brane):
 - Zero mode: $M_0 = 0$ (ordinary gravitation) $\text{Coupling} \propto \frac{1}{M_{Pl}}$
 - Higher modes: $M_n = x_n k e^{-k\pi R}, J_1(x_n) = 0$
 $(x_1 \approx 3.8; x_2 \approx 7.0; x_3 \approx 10.2; \dots)$ $\text{Coupling} \propto \frac{1}{\text{TeV}}$
- Interaction: $\int d^4x \left[\frac{1}{M_{Pl}} h_{\mu\nu}^{(0)}(x) T^{\mu\nu} + \frac{1}{\Lambda_\pi} \sum_{n=1}^{\infty} h_{\mu\nu}^{(n)}(x) T^{\mu\nu} + \dots \right]$
- $\Lambda_\pi = M_{Pl} e^{-k\pi R}$ [$kR \approx 10$] \Rightarrow gravity effects at TeV scale
- Of interest for colliders: narrow peaks in dilepton/diphoton production

- RS-model parameters and notations:

- $M_1 \equiv M_G$ mass of first KK graviton
(expected $M_G \approx \text{TeV}$)
- $c = k/M_*$ coupling to matter [universal]

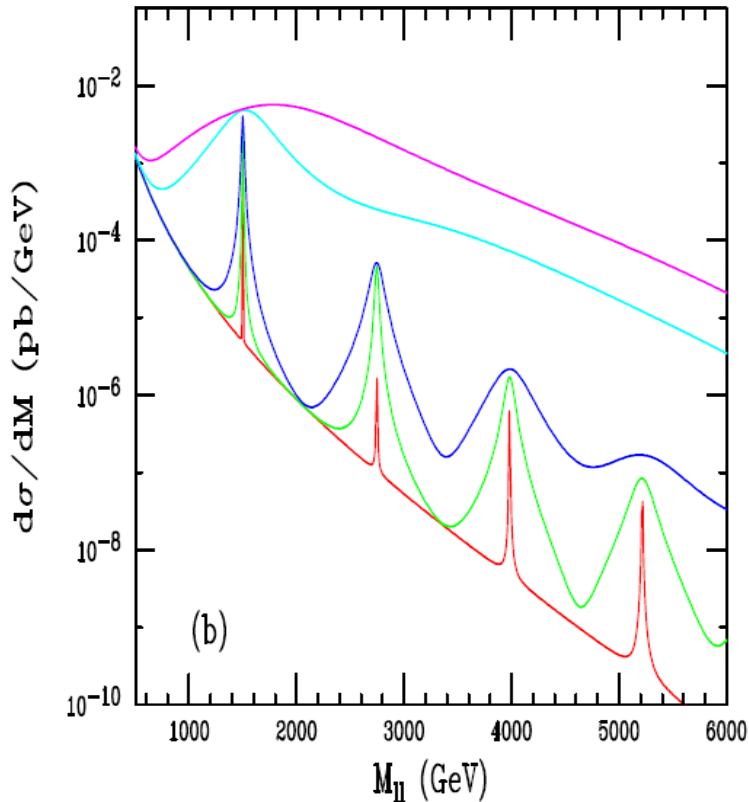
- Unevenly spaced spectrum:

$M_n = (x_n/x_1) M_G$ distinctive of
the model by itself

- Constraints:

- $c < 0.1$ [$k \sim 5\text{D}$ curvature]
- $c \geq 0.01$ [string arguments]
- $\Lambda_\pi < 10 \text{ TeV}$ [no extra hierarchy]

Width $\Gamma_n = \rho M_n x_n^2 c^2$:
narrow resonances for small c



DY production, $M_1=1.5 \text{ TeV}$ and subsequent tower states at LHC ($c = 1, 0.5, 0.1, 0.05, 0.2$, and 0.01 , from top to bottom). [*Davoudiasl et al.*]

Experimental (TEVATRON):

$M_G > 560 \text{ GeV}$ ($c = 0.01$)

$M_G > 1050 \text{ GeV}$ ($c = 0.1$)

[arXiv:1004.1826v1 [hep-ex], 2010]

Extensions of RS1:

- SM fermions and gauge bosons in the bulk (but not Higgs) → fermion mass hierarchy.
- Stability of R: massless scalar “graviton” → spin-0 “radion”.
- Complicated phenomenology predicted.

“Principal” signatures

- Spin-2 KK graviton excitations – cleanest events:

$$p + p \rightarrow l^+ l^- + X \quad \left[\begin{array}{l} p + p \rightarrow R + X \\ \gamma \gamma + X \end{array} \right] \quad \xrightarrow{\text{red arrow}} \quad l^+ l^-, \gamma \gamma$$

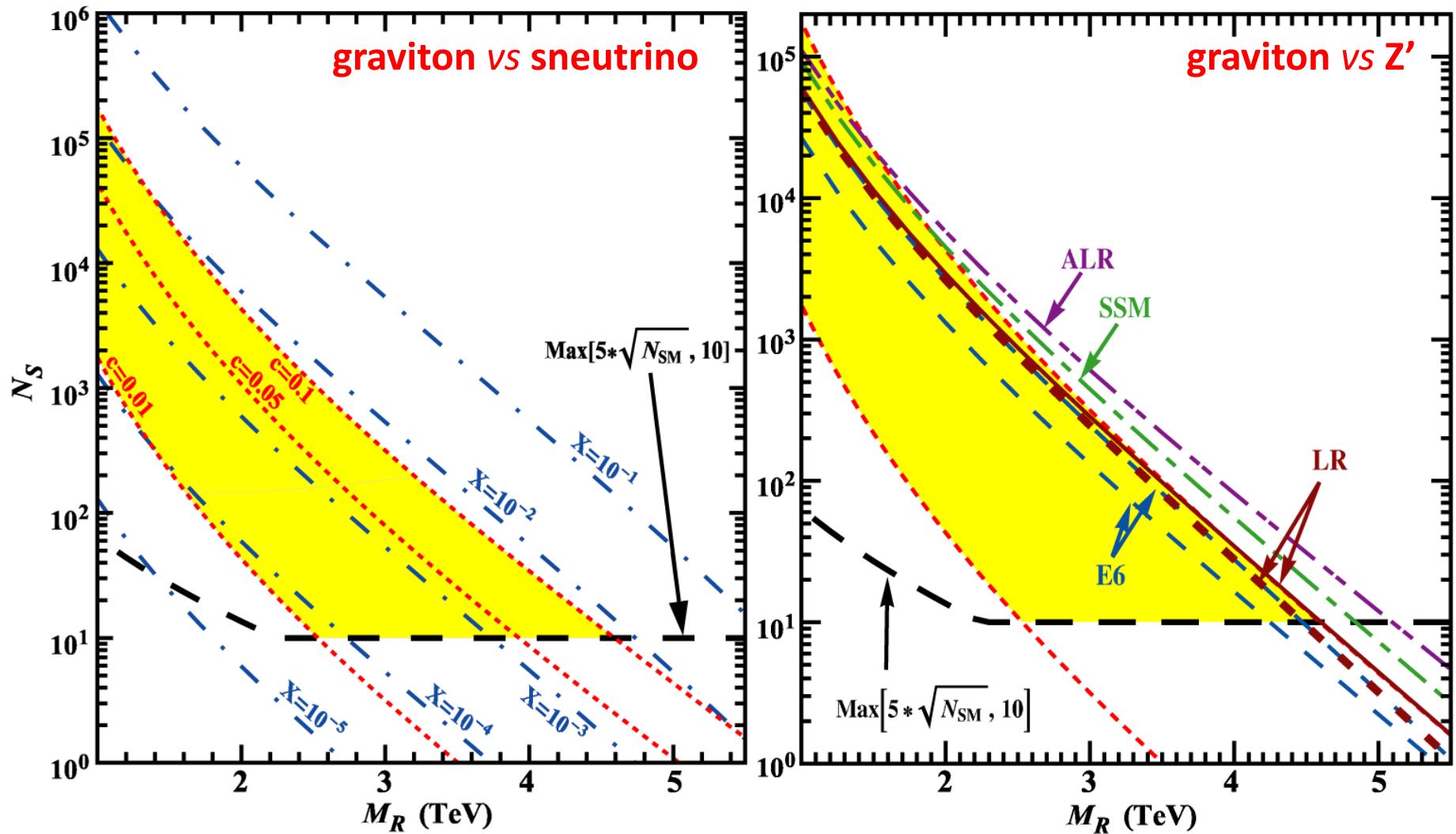
- Discovery vs. identification



- Discovery reach: peak observation at some M_R over SM background → allowed region on model parameters ($R \rightarrow l^+ l^-, R \rightarrow \gamma\gamma, \dots$)
- Identification reach: identify source of observed peak (different models can give same M_R and same number of events under the peak) \leftrightarrow test the peak quantum numbers → more stringent conditions on parameters.

Model examples: production of narrow peaks with same n. of events

- Spin:
 - Model of gravity with extra warped dimension (**spin-2**)
 - Z' models (**spin-1**)
 - SUSY with R_p sneutrinos (**spin-0**)
 - Identify spin of peak → discriminate among models
- “CONFUSION REGIONS”



$$pp \rightarrow l^+l^- + X \ (l = e, \mu)$$

$$0.01 < c < 0.1, \mathcal{L}_{\text{int}} = 100 \text{ fb}^{-1}.$$

Cross section for $p + p \rightarrow R \rightarrow l^+ l^-$

I) Production cross section: $\sigma(R_{ll}) = \int_{-z_{\text{cut}}}^{z_{\text{cut}}} dz \int_{M_R - \Delta M/2}^{M_R + \Delta M/2} dM \int_{-Y}^Y dy \frac{d\sigma}{dM dy dz}$.

\rightarrow number of events under the peak

Experimental inputs & cuts

- Minimal number of signal events over SM background (**discovery**):

$$[\text{signal} \equiv \max(5 \cdot \sqrt{N_B}, 10)]$$
- Inputs: reconstruction efficiency $\varepsilon_l = 90\%$; $L_{\text{int}} = 10; 100 \text{ fb}^{-1}$;
- cuts $|y| \leq 2.5$; $P_T^l > 20 \text{ GeV}$; $|\eta_{l,\bar{l}}| < 2.5$; $z_{\text{cut}} = 0.987$
- ATLAS detector resolution: $\Delta M = 24 (0.625 M + M^2 + 0.0056)^{1/2} \text{ GeV}$,
dilepton mass **M** measured in TeV units.

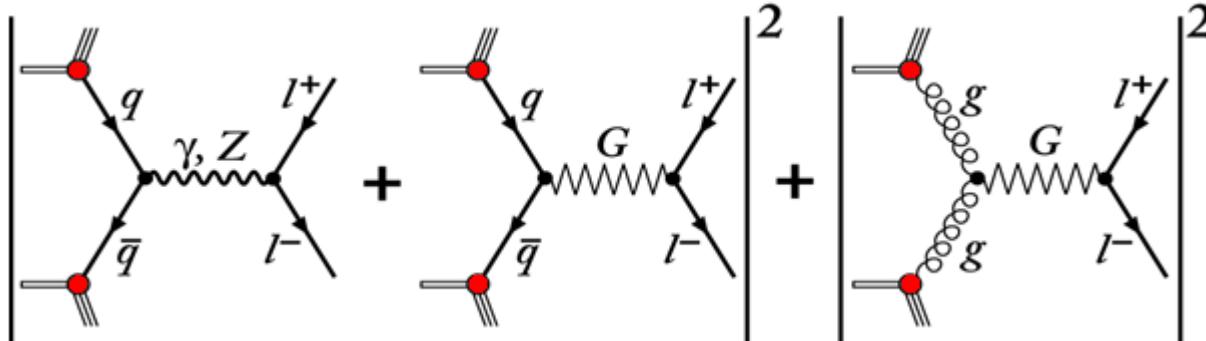
II) Angular distribution $\frac{d\sigma}{dz} = \int_{M_R - \Delta M/2}^{M_R + \Delta M/2} dM \int_{-Y}^Y \frac{d\sigma}{dM dy dz} dy,$

$$Y = \log(\sqrt{s}/M). \quad (y \equiv \text{rapidity of lepton pair}, M \equiv \text{dilepton mass}, z \equiv \cos\theta^*)$$

Hadron level

$$p p \rightarrow l^+ l^- + X$$

- **SM:** $q\bar{q} \rightarrow \gamma, Z \rightarrow l^+ l^-$
- **Graviton exchange** signatures in RS scenario:

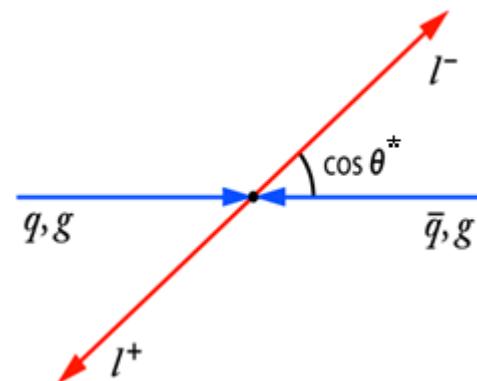


- Parton distribution functions $f_q(\xi, M)$ – CTEQ6
- Partonic cross sections $\hat{\sigma}$ [*Giudice et al., 1999; Han et al., 1999*]

$$\bar{q}q, gg \rightarrow \gamma, Z, G \rightarrow l^+ l^-$$

C.M. $l^+ l^-$ system:

$$z \equiv \cos(\theta^*)$$



Summary of angular distributions

$$\frac{d\sigma(G_{ll})}{dz} = \frac{3}{8}(1+z^2)\sigma_{q\bar{q}}^{\text{SM}} + \frac{5}{8}(1-3z^2+4z^4)\sigma_{q\bar{q}}^G + \frac{5}{8}(1-z^4)\sigma_{gg}^G,$$

$$\frac{d\sigma(V_{ll})}{dz} = \frac{3}{8}(1+z^2)(\sigma_{q\bar{q}}^{\text{SM}} + \sigma_{q\bar{q}}^V),$$

$(l = e, m)$

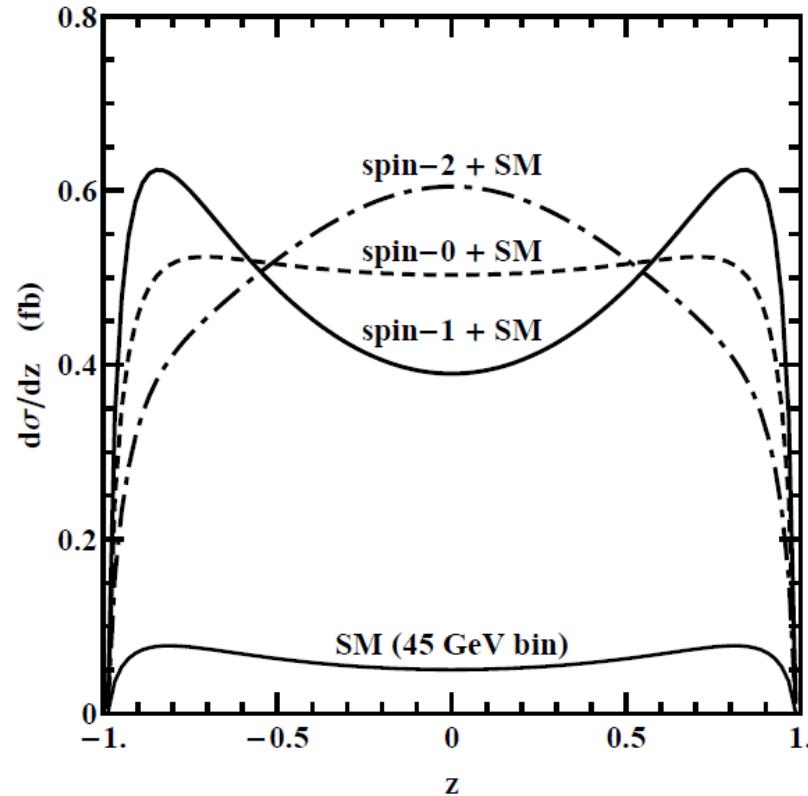
$$\frac{d\sigma(S_{ll})}{dz} = \frac{3}{8}(1+z^2)\sigma_{q\bar{q}}^{\text{SM}} + \frac{1}{2}\sigma_{q\bar{q}}^S.$$

$$\sigma(G_{ll}) = \sigma_{q\bar{q}}^{\text{SM}} + \sigma_{q\bar{q}}^G + \sigma_{gg}^G,$$

$$\sigma(V_{ll}) = \sigma_{q\bar{q}}^{\text{SM}} + \sigma_{q\bar{q}}^V,$$

$$\sigma(S_{ll}) = \sigma_{q\bar{q}}^{\text{SM}} + \sigma_{q\bar{q}}^S.$$

z-even



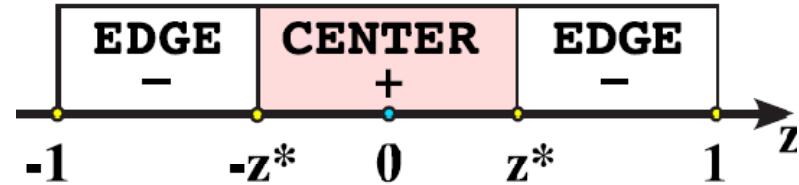
Process	Normalized density
$q\bar{q} \rightarrow (\gamma, Z) \rightarrow l^+ l^-$, $q\bar{q} \rightarrow Z' \rightarrow l^+ l^-$	$\frac{3}{8}(1+z^2)$
$gg \rightarrow G \rightarrow l^+ l^-$	$\frac{5}{8}(1-z^4)$
$q\bar{q} \rightarrow G \rightarrow l^+ l^-$	$\frac{5}{8}(1-3z^2+4z^4)$
$q\bar{q} \rightarrow \tilde{\nu} \rightarrow l^+ l^-$	$\frac{1}{2}$ (flat)

$c = 0.01$, $M_R = 1.6$ TeV, equal numbers of DY resonant events

Identification reach on spin-2

- Difficulty of angular analysis: boost to CM frame and sign of z
- Centre-Edge Asymmetry $A_{CE} = \frac{\sigma_{CE}(R_{ll})}{\sigma(R_{ll})}$: **z-symmetric integration!**

$$\sigma_{CE}(R_{ll}) = \left[\int_{-z^*}^{z^*} - \left(\int_{-1}^{-z^*} + \int_{z^*}^1 \right) \right] \frac{d\sigma(R_{ll})}{dz} dz,$$



$$A_{CE}^G = \epsilon_q^{\text{SM}} A_{CE}^V + \epsilon_q^G \left[2(z^*)^5 + \frac{5}{2} z^* (1 - (z^*)^2) - 1 \right] + \epsilon_g^G \left[\frac{1}{2} z^* (5 - (z^*)^4) - 1 \right],$$

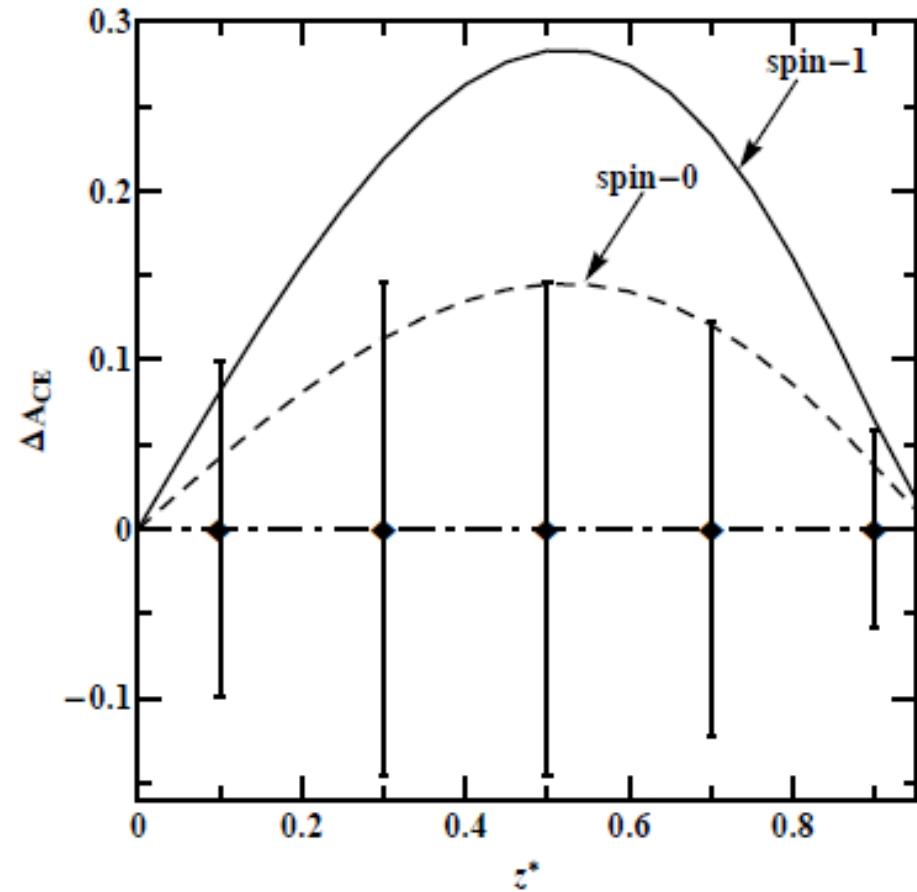
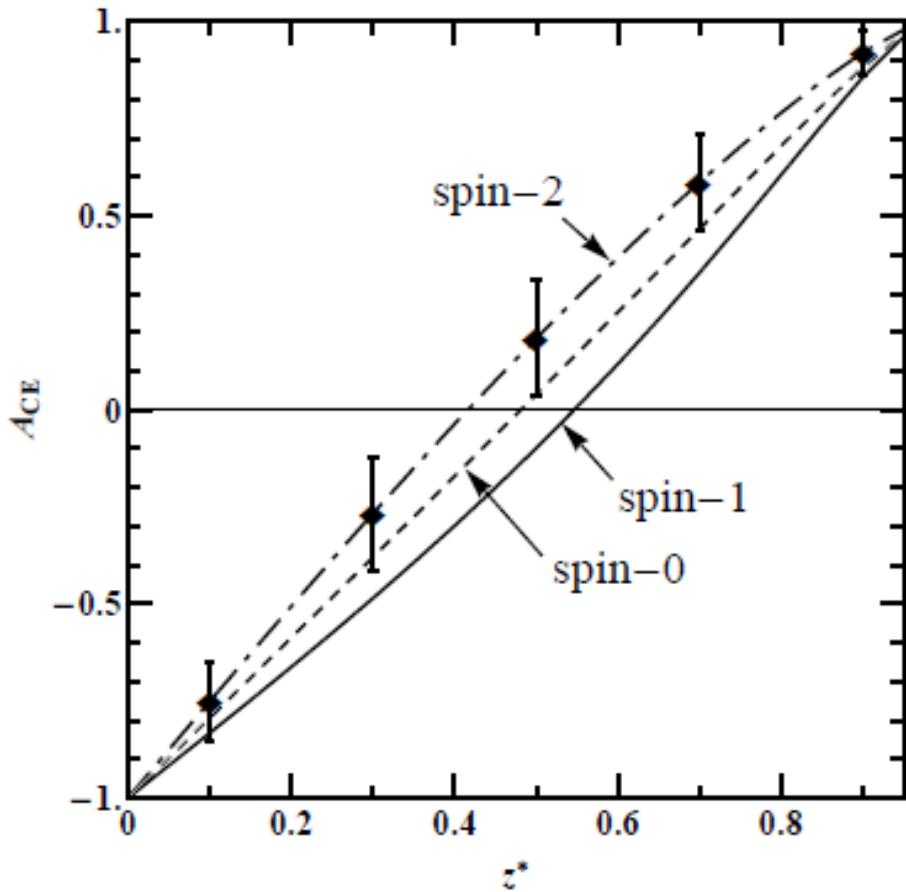
$$A_{CE}^V \equiv A_{CE}^{\text{SM}} = \frac{1}{2} z^* (z^*)^2 + 3 - 1, \quad A_{CE}^S = \epsilon_q^{\text{SM}} A_{CE}^V + \epsilon_q^S (2 z^* - 1).$$

ϵ_q^G , ϵ_g^G s and ϵ_q^{SM} : fractions of peak events from $q\bar{q}, gg \rightarrow G \rightarrow l^+l^-$ and SM background, respectively. $\epsilon_q^G + \epsilon_g^G + \epsilon_q^{\text{SM}} = 1$.

A_{CE} is the same for Z' models and SM.

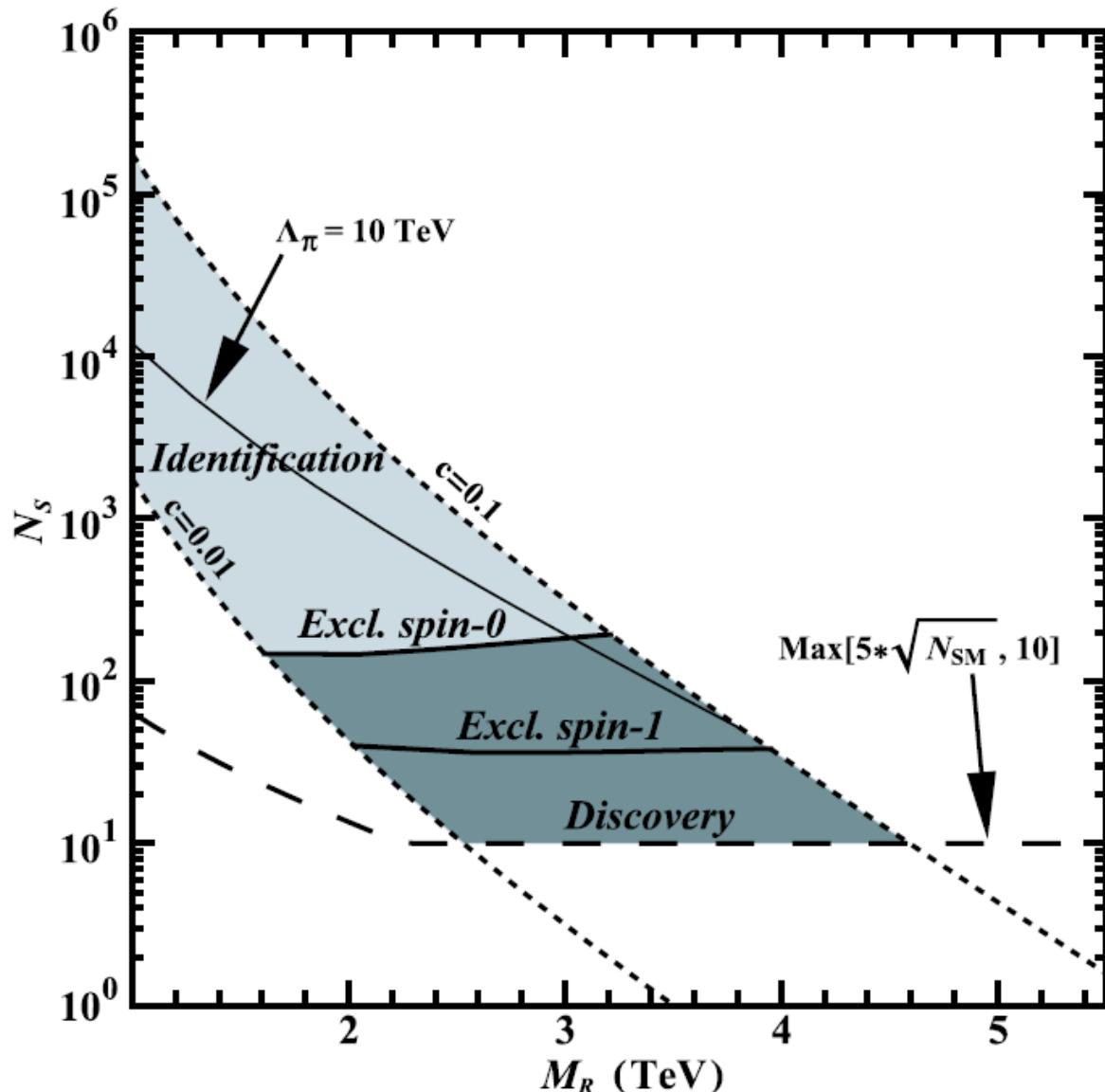
χ^2 analysis using A_{CE}

- Constraints for identification reach on spin-2:
 - Exclusion spin-0: $\Delta A_{CE} = A_{CE}^G - A_{CE}^S$
 - Exclusion spin-1: $\Delta A_{CE} = A_{CE}^G - A_{CE}^V$
- $\chi^2 = \left[\frac{\Delta A_{CE}}{\delta A_{CE}} \right]^2, \quad \delta A_{CE} = \sqrt{\frac{1 - (A_{CE}^G)^2}{\epsilon_l \mathcal{L}_{int} \sigma(G_{ll})}}.$
- Condition: $\chi^2 \geq \chi^2_{C.L.} \Rightarrow$ determine range in (c, M_G)



$c = 0.01$; $M_R = 1.6$ TeV; $L_{int} = 100$ fb $^{-1}$; same number of events;
error bars correspond to 2σ ; with experimental cuts

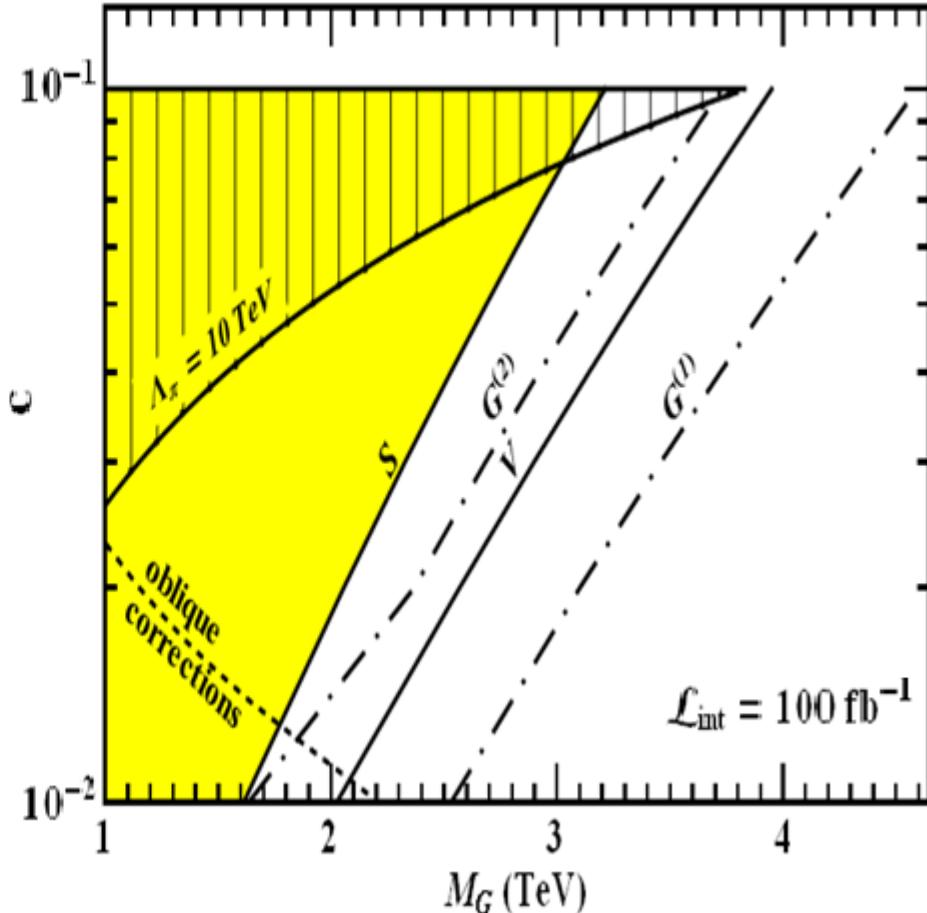
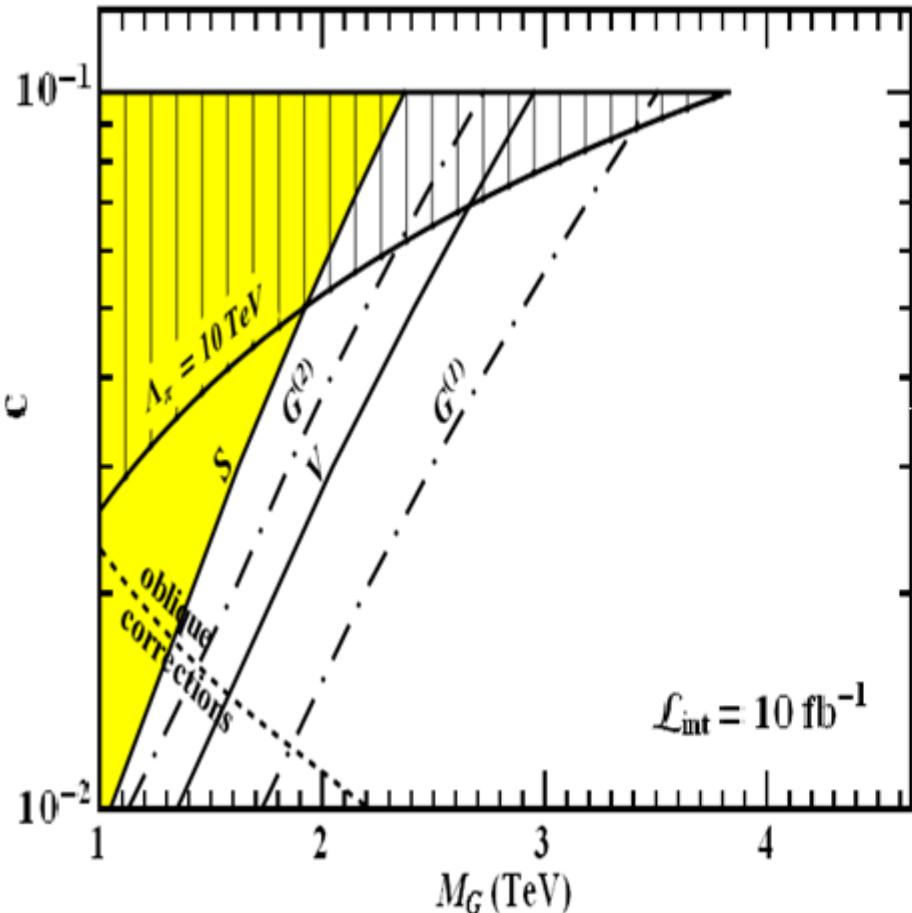
- $z^* \approx 0.55$ is “optimal” – and χ^2 smooth
- $\Delta A_{CE}^{spin-1} > \Delta A_{CE}^{spin-0} \Rightarrow$ Excl. spin-0 => Excl. spin-1



Exclusion limits and identification reach at 95% C.L. and $\mathcal{L}_{\text{int}} = 100 \text{ fb}^{-1}$. The channels $l = e, \mu$ are combined. The theoretically favored region, limited by the $\Lambda_\pi = 10 \text{ TeV}$ and $c = 0.1$ lines, is also indicated.

RS parameter plane

■ – domain allowing for spin-2 identification
vs. domain allowed for discovery (95% C.L.)



- Potential advantage of A_{CE}
 - “transparency” to spin-1 Z’ (model independence)
 - ratio of cross sections --> minimized systematic uncertainties
 - “free” from sign ambiguity in $\cos\theta^*$ (cfr FB asymmetry)
 - Boost to CM must still be performed ...
- Attempt with pseudorapidity [Diener *et al.*, 2009]

$$y = \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right) \rightarrow \eta = -\ln \tan \left(\frac{\theta}{2} \right)$$

Then: $\Delta\eta_{lab} = \Delta\eta^* = \ln \left(\frac{1 + \cos\theta^*}{1 - \cos\theta^*} \right)$

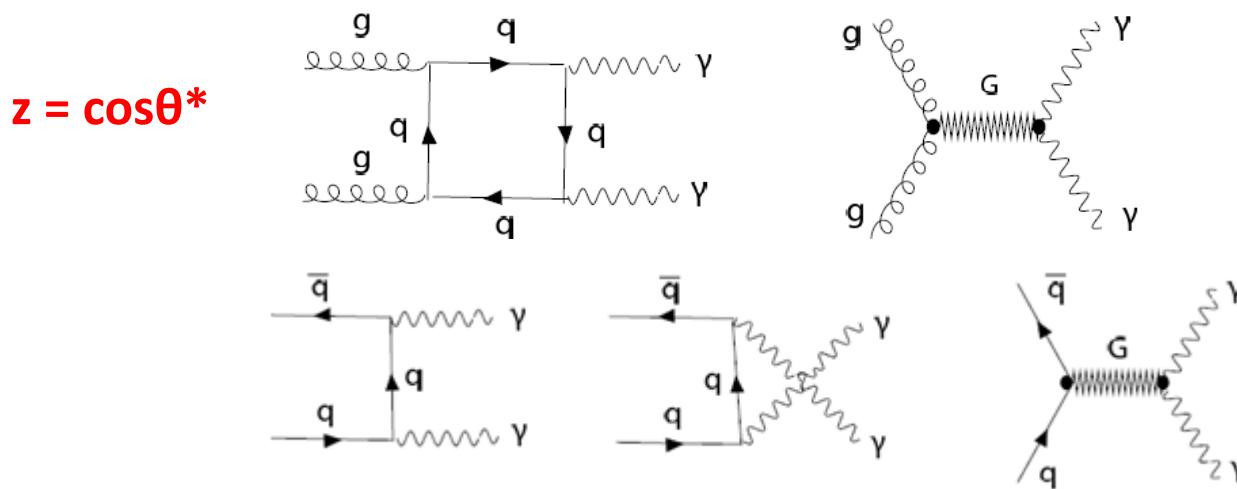
$$\tilde{A}_{CE} = \frac{\left(\int_{-\Pi}^{\Pi} - \int_{-\infty}^{-\Pi} - \int_{\Pi}^{\infty} \right) \frac{d\sigma}{d\Delta\eta} d\Delta\eta}{\int_{-\infty}^{\infty} \frac{d\sigma}{d\Delta\eta} d\Delta\eta}$$

TABLE I: \tilde{A}_{CE} values with corresponding statistical uncertainties for 100 fb^{-1} integrated luminosity, $p_{T_l} > 20 \text{ GeV}$, $|\eta| < 2.5$, within one bin $\Delta M_{l+l-} = 42.9 \text{ GeV}$ and $M_R = 1.5 \text{ TeV}$. Also shown are the expected number of total events for each model assuming 100 fb^{-1} integrated luminosity.

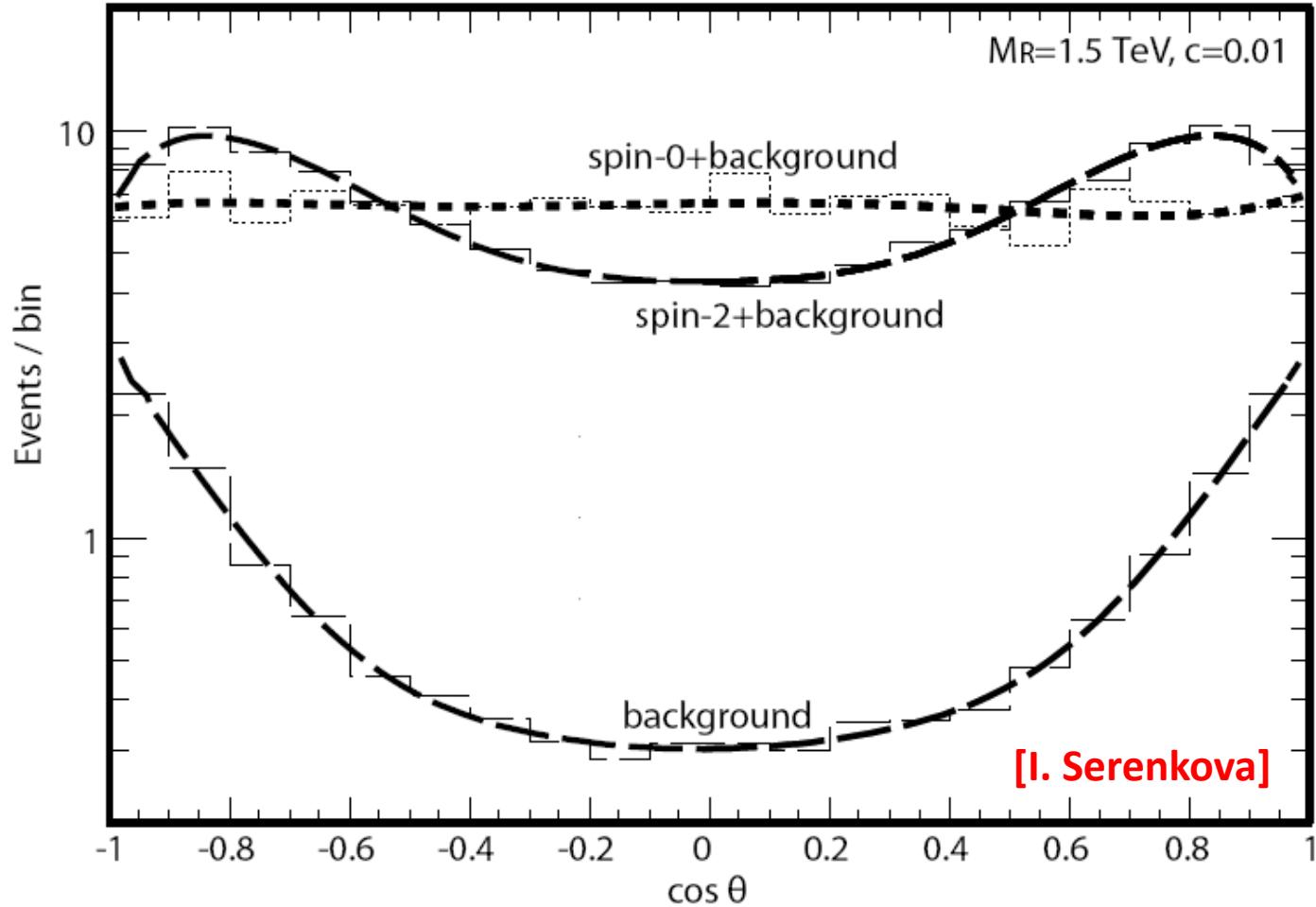
Model	$\tilde{A}_{CE} \pm \delta\tilde{A}_{CE}$	N Events
$E_6 \chi$	-0.106 ± 0.017	3875
$E_6 \psi$	-0.095 ± 0.022	2223
$E_6 \eta$	-0.092 ± 0.021	2480
LR Symmetric	-0.099 ± 0.018	3350
Sequential SM	-0.097 ± 0.016	4162
Littlest Higgs	-0.095 ± 0.001	6217
Simplest Little Higgs	-0.094 ± 0.017	3542
RS Graviton	$+0.228 \pm 0.011$	8208
R-parity violating $\tilde{\nu}$	$+0.055 \pm 0.066$	251

$$p + p \rightarrow G \rightarrow \gamma\gamma \quad (\text{work in progress})$$

- $V \rightarrow \gamma\gamma$ excludes spin-1: $\frac{G \rightarrow \gamma\gamma}{G \rightarrow l^+l^-} \sim 2$ – predicted => strong test (Randall – Wise)
- Can exclude spin-0? A_{CE} -based analysis applicable.



- $g g \rightarrow G \rightarrow \gamma\gamma \propto 1 + 6z^2 + 4z^4$: dominant, peaked at $z = \pm 1$
- $q \bar{q} \rightarrow G \rightarrow \gamma\gamma \propto 1 - z^4$
- Potential background from initial brehmsstrahlung – not for $M_G \sim \text{TeV}$



Angular distribution of photons in the diphoton c.m. system for spin-2 graviton resonant production in the RS model with $c=0.01$; spin-0 resonant production, and prompt photon background. We take $M_R = 1.5$ TeV. The spin-0 distribution is normalized so that the total number of events is the same for the graviton and the scalar.

Results for spin-2 identification

- Discovery reach [5 σ] and spin-2 identification reach on lowest RS graviton from ACE analysis (95% CL)

L_{int}	Discovery		Identification	
	c=0.01	c=0.1	c=0.01	c=0.1
10 fb ⁻¹	1.7 TeV	3.5 TeV	1.1 TeV	2.4 TeV
100 fb ⁻¹	2.5 TeV	4.6 TeV	1.6 TeV	3.2 TeV

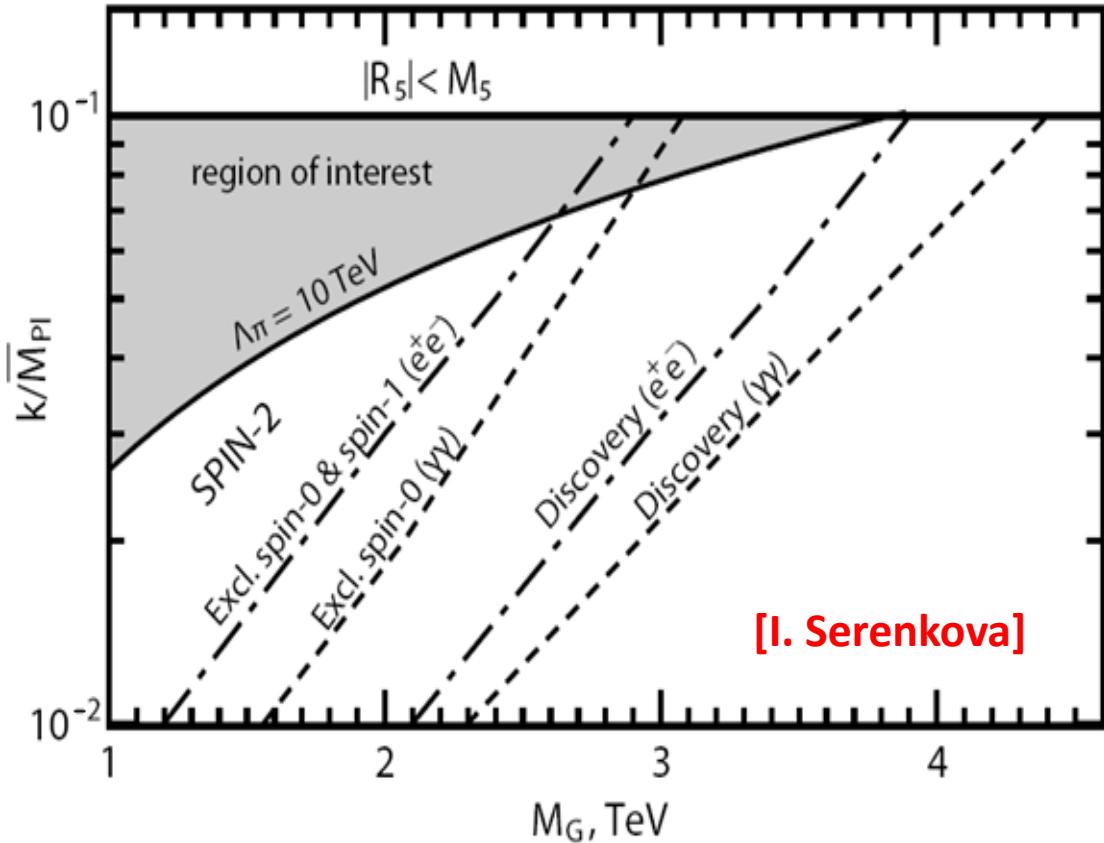
- Allowed region for discovery substantially reduced by spin-0 **and** spin-1 rejection
- Dramatic restriction from $\Lambda_\pi < 10$ TeV (order of magnitude)

- Discovery of $G^{(n=2)}$ in addition to lowest $G^{(n=1)}$:

Feasible (5σ) for $M_G =$

L_{int}	c=0.01	c=0.1
10 fb $^{-1}$	1.1 TeV	2.7 TeV
100 fb $^{-1}$	1.6 TeV	3.7 TeV

- Recall $\frac{M_2}{M_1} = \frac{x_2}{x_1}$ predicted.
- Relevant (c, M_G) pairs always stay within the spin-0 (\Rightarrow spin-1) rejection domain.
- RS resonance doubly clinched: spin-2 from angular analysis & mass spectrum.



Experimental cuts:

Detector response has been taken into account: the photons considered are the two isolated highest transverse momentum photons with $|n| \leq 2.4$ and $p_T \geq 40$ GeV.

Discovery limits and identification reaches on the spin-2 graviton parameters in the plane (M_G, c), using the $\gamma\gamma$ and e^+e^- production cross sections and A_{CE} at the LHC with $L_{int}=100 \text{ fb}^{-1}$. *Discovery* – 5σ . *Exclusion* – rejection of the spin-0 hypothesis with $\gamma\gamma$ channel and rejection of both spin-1 and spin-0 hypotheses with e^+e^- channel at 95% C.L. Both of them determine the regions lying above the corresponding curves and marked as spin-2 that correspond to the identification of spin-2 RS graviton. The theoretically favored region, $\Lambda_\pi < 10$ TeV (shaded), is also indicated.

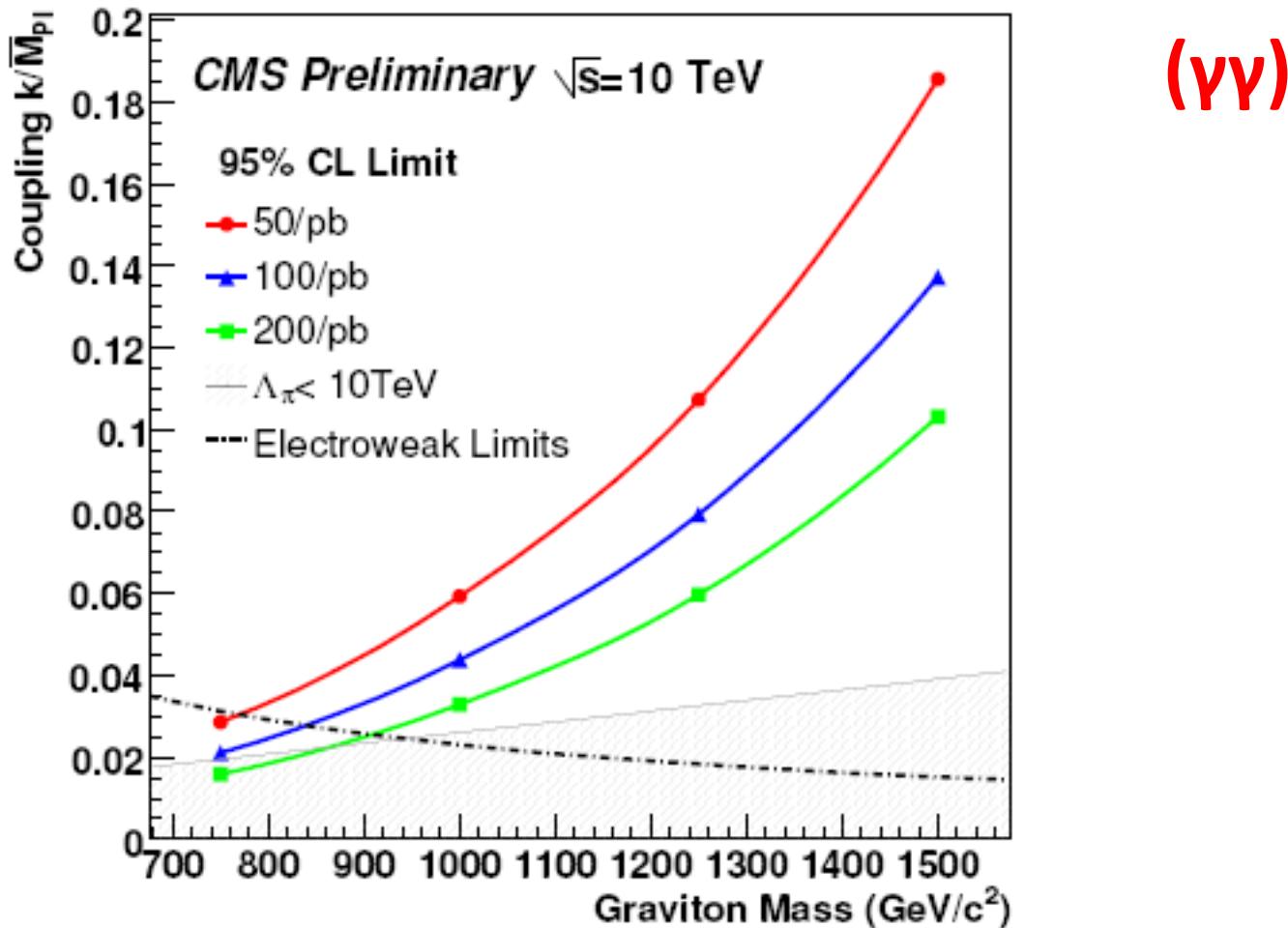


Figure 3: Limit on RS parameters (M_1, \tilde{k}) , extrapolated from the results of the large ED diphoton search for 100/pb. The area to the left of the curves is excluded. The gray shaded region shows the area excluded for $\Lambda_\pi < 10 \text{ TeV}/c^2$. The area below the dash-dotted line is excluded by precision electroweak data [11].

Radion phenomenology

- 4D massless scalar: $d s^2 = e^{-2k|y|T(x)} g_{\mu\nu}(x) d x^\mu d x^\nu + (T(x))^2 d y^2$
- Ruled out (tests of general relativity) & R must be stabilized
- Massive scalar field Φ in the bulk with VEV and “brane” conditions (Goldberger – Wise)
⇒ minimum of potential at $y = \pi R \approx 10$ (stable 2-brane set-up)

On the TeV brane:

- $m_\Phi \sim \varepsilon \frac{1}{\sqrt{kR}} M_{Pl} e^{-\pi kR} \equiv \varepsilon \frac{\Lambda_\pi}{\sqrt{kR}} \sim \text{TeV}$, but $m_\Phi < G_1(KK)$!
→ "back reaction", $\varepsilon \leq 1$

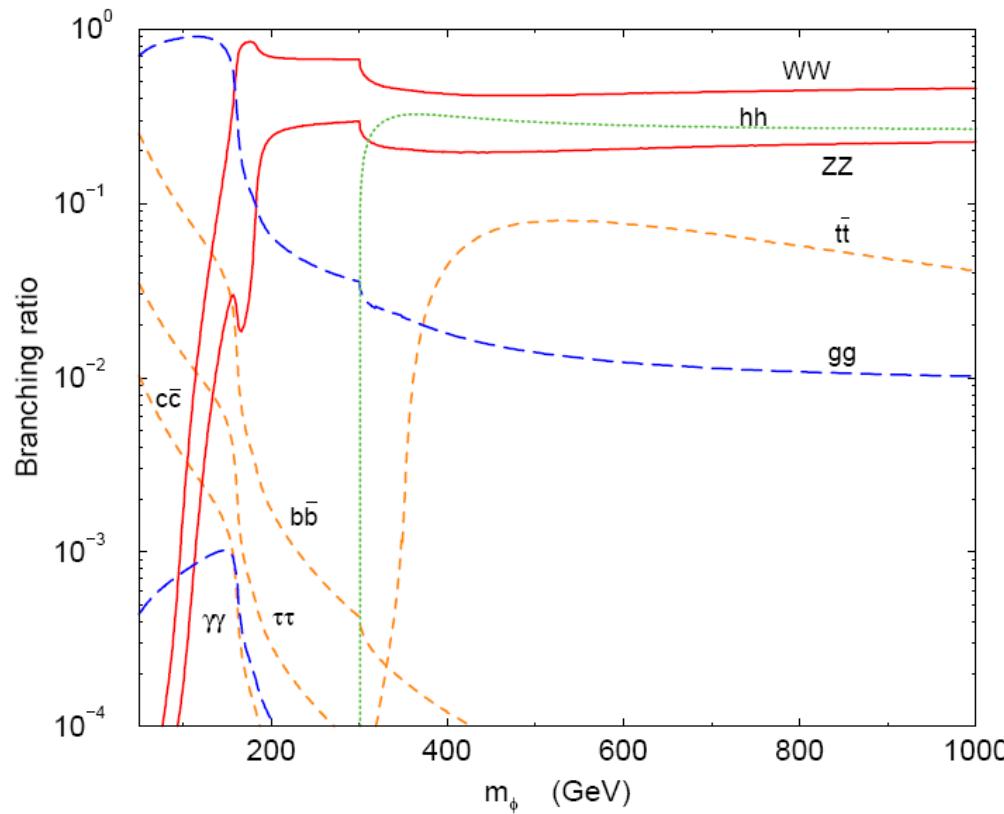
- Lightest mode in RS1, most important excitation to study at LHC?

- Interaction with SM fields:

$$\mathcal{L} \sim \frac{1}{\sqrt{6}\Lambda_\pi} \Phi T_\mu^\mu \equiv \frac{1}{\Lambda_\Phi} \Phi T_\mu^\mu$$

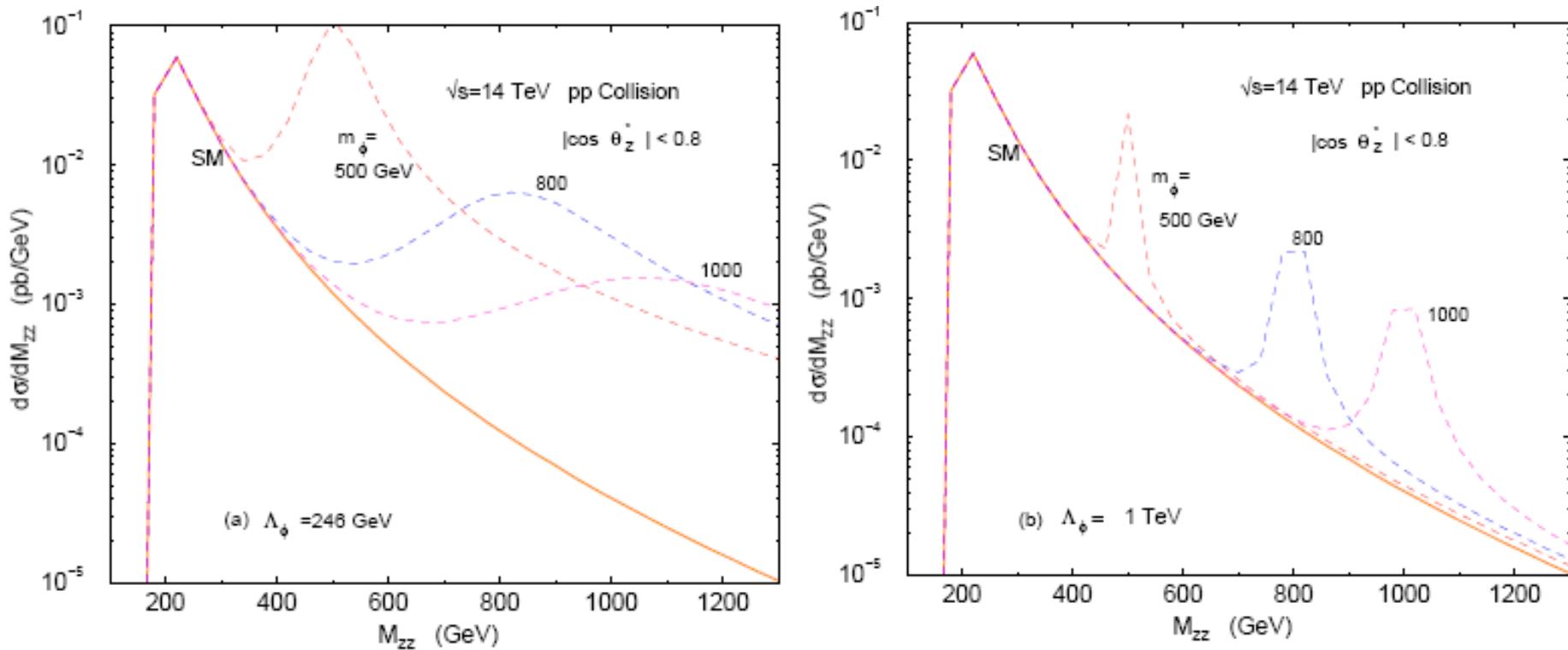
- Λ_Φ new scale at $\sim \text{TeV}$, roughly $\frac{\Lambda_\Phi}{4\pi} < m_\Phi < \Lambda_\Phi$

- Similar couplings as SM Higgs (scale factor $\frac{v_{SM}}{\Lambda_\Phi}$)
- Except $g_{\Phi \rightarrow gg} \gg g_{H \rightarrow gg}$ (loop & trace anomaly)
- $\Phi \leftrightarrow H$ interference for lower m_Φ → accompanies Higgs searches



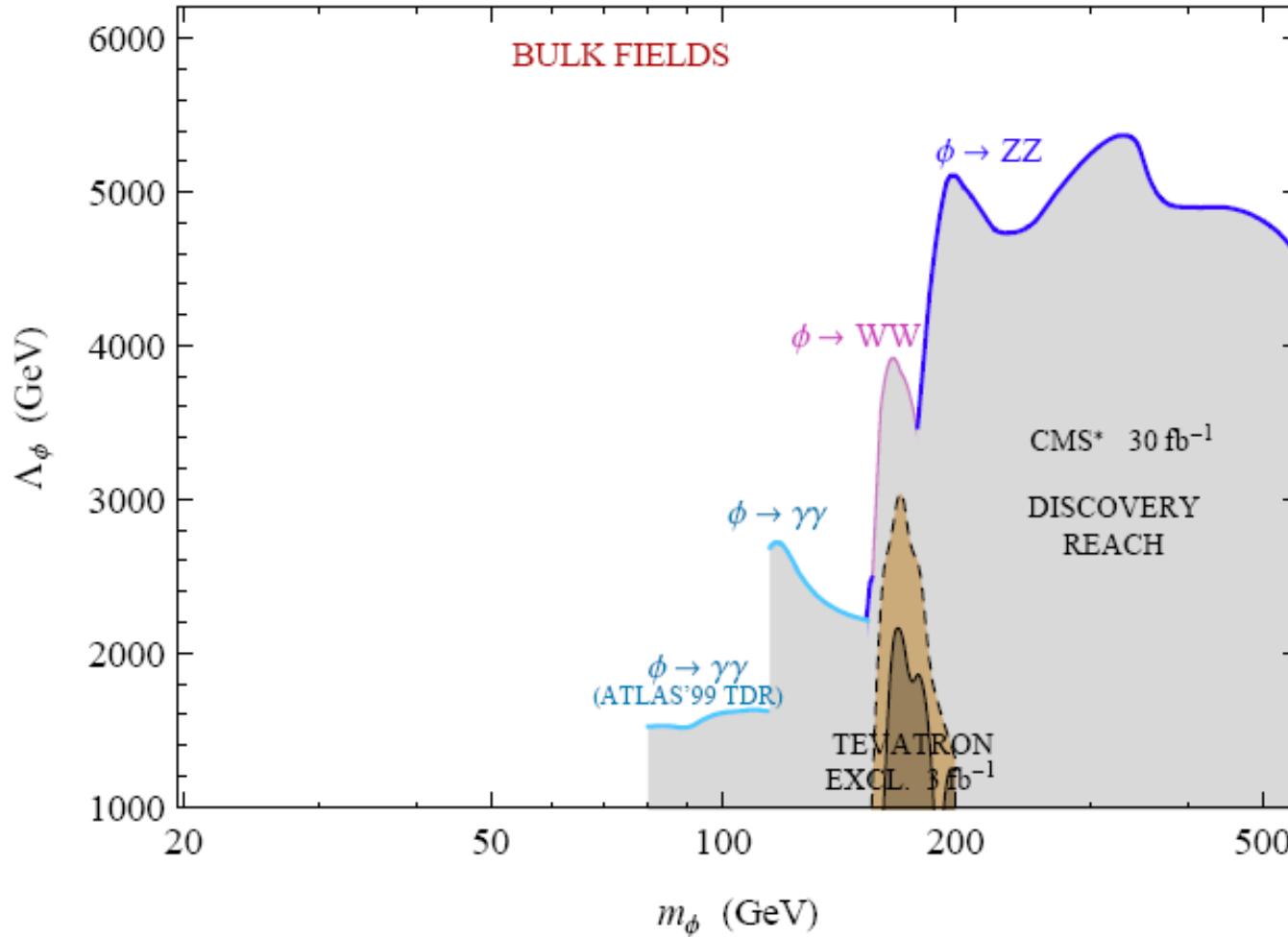
Branching ratios of the radion versus m_ϕ . ($m_H = 150$ GeV.) [K. Cheung , arXiv:hep-ph/0009232]

Expected scenario at the LHC



Invariant mass distribution $d\sigma/dM_{zz}$ for the radion signal with $m_\phi = 500, 800, 1000$ GeV and the SM background $q\bar{q} \rightarrow Z Z$ at the LHC, for (a) $\Lambda_\phi = 246$ GeV and (b) $\Lambda_\phi = 1$ TeV. A cut of $|\cos \theta_z^*| < 0.8$ is imposed.

[K. Cheung , arXiv:hep-ph/0009232]



LHC discovery reach for the radion using “translated” Higgs projections from CMS (and ATLAS in the lower mass region) for 30 fb^{-1} of luminosity.
[P. Nath et al., arXiv:1001.2693v1 [hep-ph]]