Ultra-Magnetized White Dwarfs and High-Energy Emission

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"Compact Stars in the QCD phase diagram V" and "Working Group 2 Meeting of COST Action MP1304"

Monday, 23 May 2016

Outline

Motivation

2 Models of electromagnetic emission

3 Results

- ④ Conclusion
- 6 Acknowledgments

6 References

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Motivation

- Why only few SGRs/AXPs emit in radio?
- Electromagnetic emission models in neutron stars and white dwarfs pulsars: the effect of the radius.
 - Polar Cap Model and Outer Gap Model
- Application to SGRs/AXPs.
- White Dwarfs Pulsars: Sources of ultra high energetic photons.

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Models of electromagnetic emission

• There are two basic standard models associate with pair e^{\pm} production and emission of radiation in compact stars.



Figure: (REA; TORRES, 2011)

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Polar Cap Model

• The first process is the polar-cap model, where electrons are accelerated on the stellar surface emitting γ -rays by curvature radiation

$$\hbar\omega \simeq \frac{3}{2} \frac{\hbar\gamma^3}{r_c},\tag{1}$$

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where r_c is the curvature radius

• In this model the condition to pair production is given by

$$\left(\frac{e\Delta V}{mc^2}\right)^3 \frac{\hbar}{2mcr_c} \frac{h}{r_c} \frac{B_s}{B_c} \ge \frac{1}{15}.$$
(2)

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Polar Cap Model



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• The second model of emission consider the production of the photons far away from the stellar surface, where the magnetic field lines are open. In this model, the energy is limited by a combination of curvature radiation and inverse Compton scattering.

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- The magnetic field responsible for the pair production is given by their dipole expression

$$B(r) = B_{\rm p} R^3 / r^3 \tag{3}$$

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• There is a minimum photon frequency given by the frequency of infrared photons IR

$$\omega_c \approx \gamma_{\parallel}^3 \frac{m^5 c^9}{e^7} \Omega^2 B^{-3}, \tag{4}$$

and a maximum frequency

$$\omega_{\max} \approx \frac{e^{15}}{\gamma_{\parallel} \hbar m^9 c^{15}} \Omega^{-4} B^7$$
(5)



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SGRs/AXPs as pulsars of neutron stars, $M = 1.4 M_{\odot}$ e radius $R = 10^6$ cm

• The value of magnetic field is calculate using the expression

$$B_{\rho}/2 = \left(\frac{3c^3I}{8\pi^2 R^6} P \dot{P}\right)^{1/2} = B_{\rm p}^{\rm NS} \sim 10^{14} \text{ G.}$$
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$$\dot{E} = I\Omega\dot{\Omega} = \dot{E}_{\rm NS} \sim 10^{33} \text{ erg/s.}$$
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• The value of $gap h_{max}$ and the curvature radius are given in terms of period and the radius of the star

$$h_{\rm max} = R_{\rm p} \approx \left(\frac{R^3\Omega}{c}\right)^{1/2} \sim 1 \ {\rm cm}; \qquad r_c \sim (Rc/\Omega)^{1/2} \sim 10^8 \ {\rm cm}.$$
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 We highlight the importance of the light cylinder comparing to neutron star radius.

$$R_{\rm L} \equiv c/\Omega \simeq 5 \times 10^9 P \sim 10^{10} \ {\rm cm} \tag{9}$$

SGRs/AXPs as pulsars of neutron stars, $M=1.4M_{\odot}$ e radius $R=10^{6}$ cm

• The difference of potential is

$$\Delta V = \frac{B_{\rm p}\Omega h^2}{2c} \sim 10^{13} \,\,\mathrm{V},\tag{10}$$

and the associate Lorentz factor

$$\gamma = \frac{e\Delta V}{mc^2} \sim 10^7. \tag{11}$$

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Polar Cap Model for SGRs/AXPs as pulsars of neutrons stars

• Using the condition

$$\left(\frac{e\Delta V}{mc^2}\right)^3 \frac{\hbar}{2mcr_c} \frac{h}{r_c} \frac{B_s}{B_c} \ge \frac{1}{15}.$$
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• M1 - Pure Dipole, $B_s = B_p$ e $r_c = (Rc/\Omega)^{1/2}$ e $h \approx R(R\Omega/c)^{1/2}$

$$4 \log B_p - 7.5 \log P + 9.5 \log R = 106.37$$

$$4 \log B_p - 7.5 \log P = 49.37, \qquad (13)$$

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• M2 - Lines of magnetic field *B* very curved, such that $r_c \sim R = 10^6$ cm, and polar cap area similar to the anterior case and $B_s = B_p$

$$4 \log B_p - 6.5 \log P + 10.5 \log R = 108.70$$

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• M3 - Lines of magnetic field *B* very curved on the Polar Cap $r_c \sim R$, but now we consider the value of *h* with magnetic field dependence, $h \sim (B_{\rm p}/B_s)^{1/2} R(R\Omega/c)^{1/2}$, where $B_s = 2 \times 10^{13}$ G is a fix value

$$7 \log B_p - 13 \log P + 21 \log R = 204.08$$

$$7 \log B_p - 13 \log P = 78.08$$
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Polar Cap Model for SGRs/AXPs as pulsars of neutrons stars

• The three curves and the SGRs/AXPs with their respective magnetic field inferred for neutron stars and their periods



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SGRs/AXPs as white dwarfs pulsars, $M = 1.4 M_{\odot}$ e radius $R = 3 \times 10^8$ cm

• Considering the SGRs/AXPs as white dwarfs, we inferred the news values for the quantities:

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$$B_{\rho}/2 = \left(\frac{3c^3I}{8\pi^2 R^6} P\dot{P}\right)^{1/2} = B_{\rm p}^{\rm WD} \sim 10^9 \,\,{\rm G}.$$
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- $\, \bullet \,$ this potential can produce curvature photons with a energy up to $\sim 10^{21}$ eV.
- SGRs/AXPs are within the GZK limit (\approx 10 Mpc for photons with a energy of $10^{19}~\text{eV})$

Polar Cap Model for SGRs/AXPs as white dwarfs pulsars

Using the condition

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we achieved the next curves:

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we achieved the next curves:

• m1 -

$$4 \log B_p - 7.5 \log P + 9.5 \log R = 106.37$$

$$4 \log B_p - 7.5 \log P = 25.84, \qquad (22)$$

m2 -

$$4\log B_{\rho} - 6.5\log P = 23.68 \tag{23}$$

m3 -

$$7 \log B_{\rho} - 13 \log P = 34.06 \tag{24}$$

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Polar Cap Model for SGRs/AXPs as white dwarfs pulsars

• m1, m2 e m3 and SGRs/AXPs with its respective magnetic fields like WD and their periods



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Outer Gap Model para SGRs/AXPs like neutron star

• In this model the magnetic field is given by its dipolar expression

$$B = B_{\rm p} R^3 / r^3 = B_{\rm p} (\Omega R / c)^3, \qquad (25)$$

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• O1 combining $\omega_c \approx \omega_{\rm max}$, we have the next death-line

$$5 \log B_{\rm p} - 12 \log P + 15 \log R = 161.54$$

$$5 \log B_{\rm p} - 12 \log P = 71.53$$

where $\gamma_{\parallel} \sim 10$ fitting by Vela pulsar.

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 \bullet O2 - If the tertiary photons has a synchrotron frequency $\omega_s \approx \omega_B$ we have the next curve

$$2 \log B_p - 5 \log P + 6 \log R = 67.41$$

$$2 \log B_p - 5 \log P = 31.40$$
(27)

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Outer Gap Model para SGRs/AXPs like neutron star

• The two curves of the Outer Gap Model and the SGRs/AXPs



Outer Gap Model for SGRs/AXPs as white dwarfs pulsars

• Considering the Outer Gap Model we have the anterior conditions ω_s , ω_c e ω_{max} used to WD, we have,

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Outer Gap Model for SGRs/AXPs as white dwarfs pulsars

• The curves o1, o2 and the all SGRs/AXPs



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• The dependency of the death-lines with the radius of the star.

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- The dependency of the death-lines with the radius of the star.
- The radius of SGRs/AXPs seems to be out of scale in the comparison with the light cylinder radius $R_{\rm NS}/R_{\rm L} \sim 10^6/10^{10}$, for white dwarfs $R_{\rm WD}/R_{\rm L} \sim 10^8/10^{10} \sim 10^{-2}$, the same value of ordinary pulsars.

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- In the Outer Gap Model "o2" for white dwarfs the SGRs/AXPs are very near or below the death-line. This could explain the absence of radio emission observed in these sources.
- The four SGRs/AXPs that emit in radio seems to be neutron stars pulsars with emission explained by the Polar Cap Model.

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Acknowledgments

- Organizing Committee of "Compact Stars in the QCD phase diagram V" and "Working Group 2 Meeting of COST Action MP1304"
- Coordenação de Aperfeiçoamento de Pessoal de Nível Superior (CAPES)
- Conselho Nacional de Desenvolvimento Científico e Tecnológico (CNPq)
- FAPESP thematic project 2015/26258-4

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Grazie!

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