

The future of $0\nu\beta\beta$





Majorana conjecture

$$\nu = \bar{\nu}$$

Main consequence :
Lepton Number Violation

Majorana vs. Dirac



$$\begin{array}{ccc} \mathbf{V}_L^M & \xleftrightarrow{\text{CPT}} & \mathbf{V}_R^M \\ & \xleftrightarrow{\text{Lorentz}} & \end{array}$$

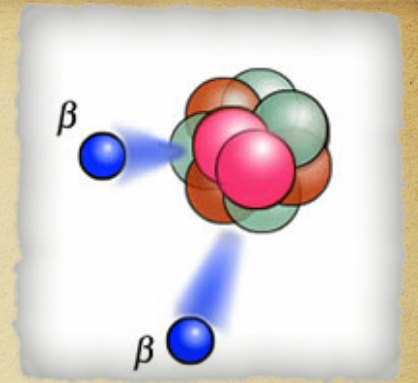
Majorana



$$\begin{array}{ccc} \mathbf{V}_L^D & \xleftrightarrow{\text{Lorentz}} & \mathbf{V}_R^D \\ \text{CPT} \updownarrow & & \updownarrow \text{CPT} \\ \overline{\mathbf{V}}_R^D & \xleftrightarrow{\text{Lorentz}} & \overline{\mathbf{V}}_L^D \end{array}$$

Dirac

Neutrino-less DBD ($0\nu\beta\beta$)

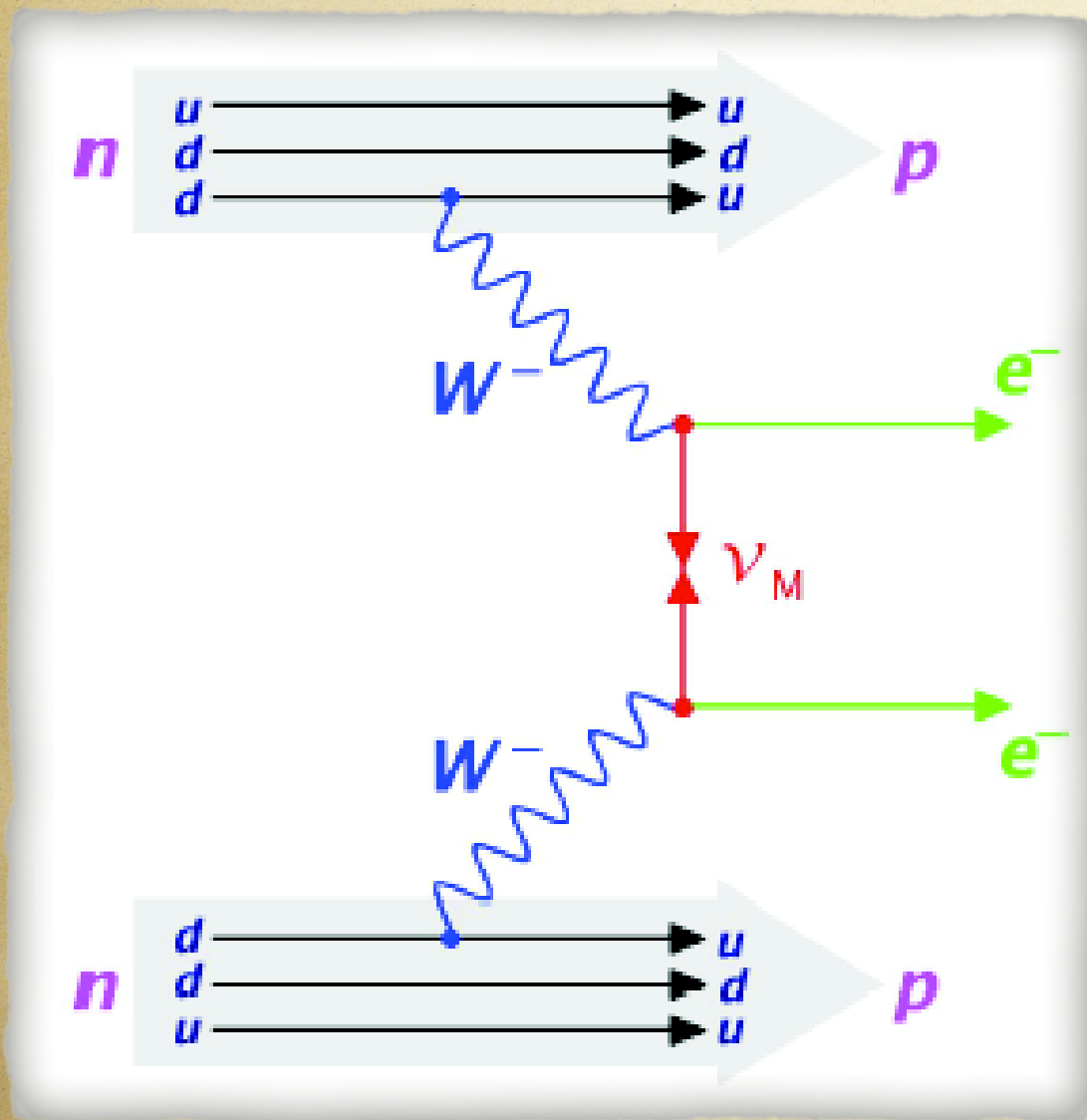


Only if:

Majorana Neutrinos

If observed:

Proof of the Majorana
nature of Neutrino



how patients should we be ?

parameter containing
the **physics**

$$1/\tau = G(Q,Z) |M_{\text{nuc}}|^2 \langle M_{\beta\beta} \rangle^2$$

what the **experimentalists**
try to measure

what the **nuclear theorists**
try to calculate

just on the back
of the envelope

$$[T_{1/2}^{0\nu}]^{-1} = C \cdot \frac{\langle m_{\beta\beta} \rangle^2}{m_e^2}$$

$$C \sim 10^{-12} \text{ y}^{-1}, \quad m_e \sim 500 \text{ keV}, \quad m_{\beta\beta} \sim 10 \text{ meV}$$

$$\tau_{1/2}^{0\nu} > 10^{26} \text{ y}$$

$$[\text{universe life } 15 \cdot 10^9 \text{ y}, \text{ Avogadro number } 6 \cdot 10^{23}]$$

what are we looking at ?

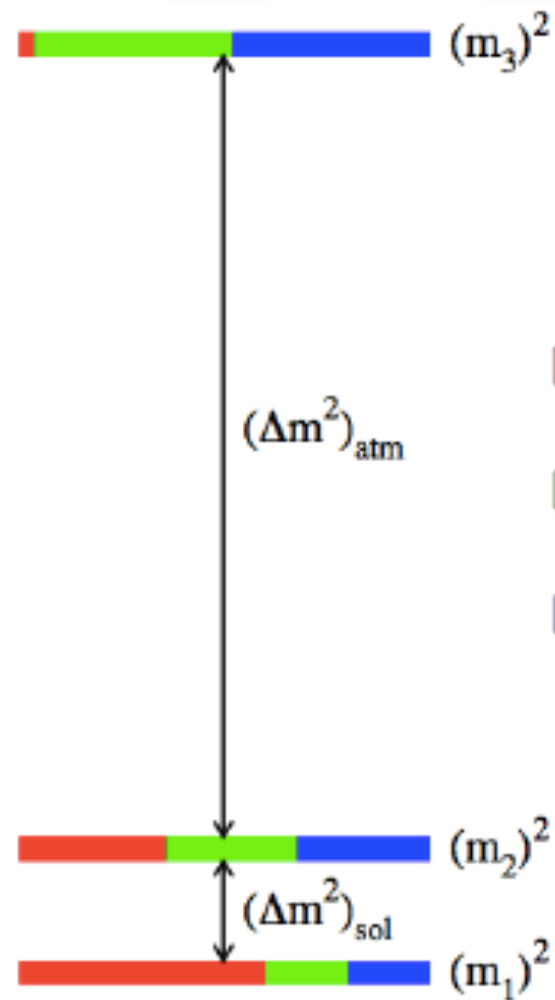
$$m_{\beta\beta} = \sum m_{\nu_k} U_{ek}^2 = \cos^2 \theta_{13} (m_1 \cos^2 \theta_{12} + m_2 e^{2i\alpha} \sin^2 \theta_{12}) + m_3 e^{2i\beta} \sin^2 \theta_{13}$$

The observable comes as a combination of the three neutrino masses, the mixing angles and the Majorana phases.

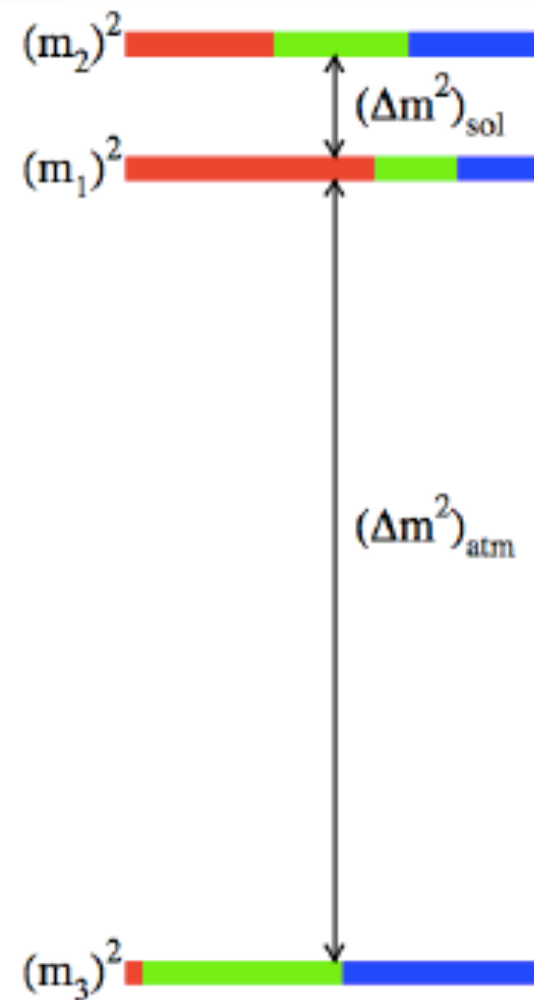
Let's parameterize as a function of the known parameters

$$m_{\beta\beta} = f(U_{ek}, m_{\text{lightest}}, \delta m_{\text{sol}}, \Delta m_{\text{atm}})$$

Two possibilities:



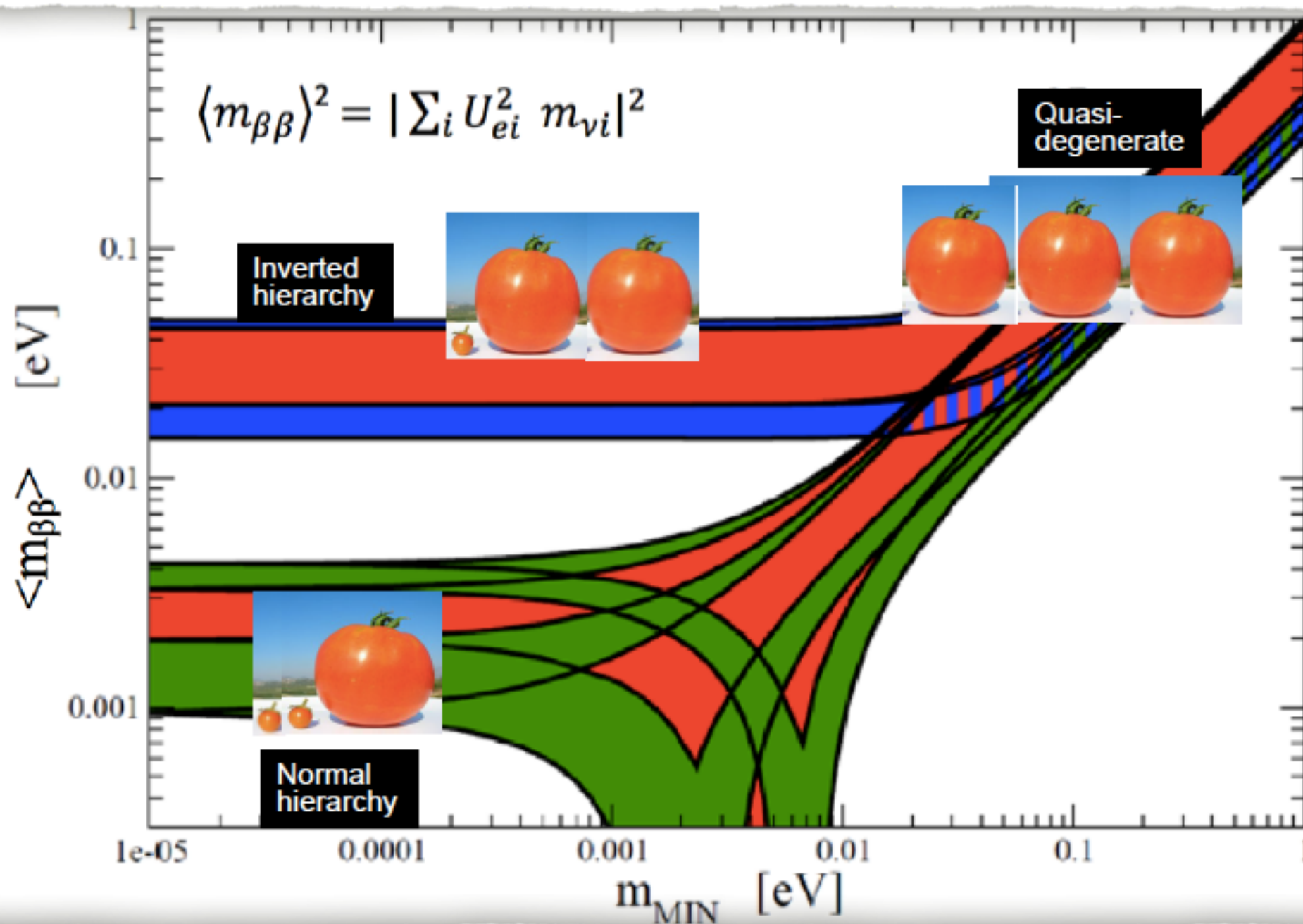
normal hierarchy



inverted hierarchy



the final result is :



The question is which, if any, part of this phase space can be attained by a realistic experiment

set a goal of exploring IH. Get down to 10-20 meV

The name of the game

expected
number of
 $\beta\beta 0\nu$ events

$$S = \frac{\overset{\text{detector mass}}{M} \cdot N_A \cdot \overset{\text{isotopic abundance}}{a}}{\underset{\text{molecular mass}}{W}} \cdot \ln(2) \cdot \frac{\overset{\text{live time}}{t}}{\underset{\text{\(\beta\beta 0\nu\) half-life}}{T_{1/2}^{0\nu}}} \cdot \overset{\text{efficiency}}{\varepsilon}$$

mean number of
background counts
around the Q-value

$$B = \overset{\text{background rate in counts/keV/kg/y}}{b} \cdot \underset{\text{detector mass}}{M} \cdot \overset{\text{energy resolution (detector FWHM)}}{\Delta E} \cdot \underset{\text{live time}}{t}$$

Sensitivity is S/\sqrt{B}

Sensitivity

$$\propto K \sqrt{\frac{M \cdot t}{B \cdot \Delta E}}$$

(i.a. • ε)

$$m_{\beta\beta} \propto \sqrt{1/\tau}$$

which way ?

- increase abundance of the right isotope (linear)
- increase M a lot (square root)
- decrease B (ideally get to mythical zero background and get rid of the square root)
- get an extraordinary good energy resolution
(remember we are talking of a signal of a few MeV but still gaining only by a square root)

brutal consideration

$$\text{Sensitivity} \propto K \sqrt{\frac{M \cdot t}{B \cdot \Delta E}} \quad (\text{i.a.} \bullet \varepsilon)$$

$$m_{\beta\beta} \propto \sqrt{1/\tau}$$

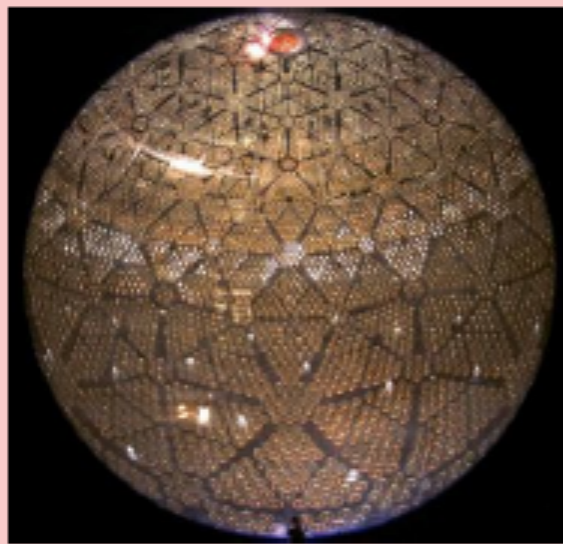
To get a factor 10 in $m_{\beta\beta}$ you have a choice :

M	100 Ton	instead of 1 Ton
t	500 y	instead of 5 y
ΔE	50 eV	instead of 5 keV

B 0.001 instead of 0.1

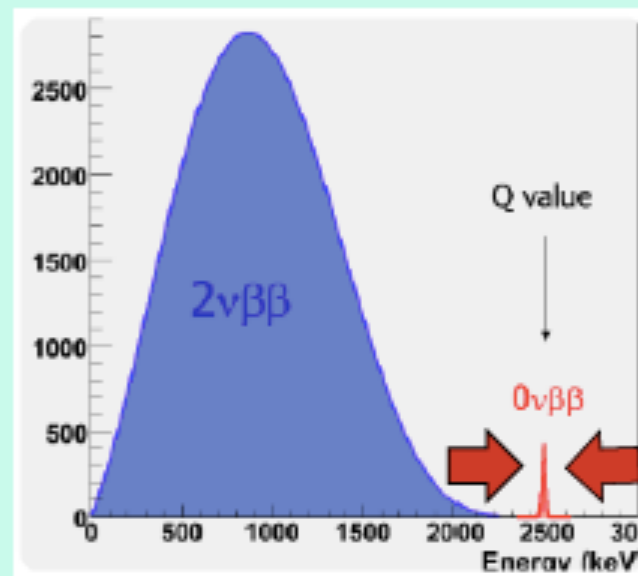
meaning :

The “Brute Force” Approach



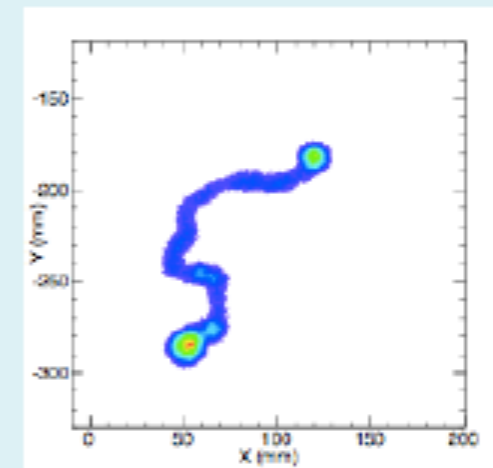
focus on the numerator
with a **huge amount
of material**
(often sacrificing
resolution)

The “Peak-Squeezer” Approach



focus on the denominator
by **squeezing down ΔE**
(various technologies)

The “Final-State Judgement” Approach

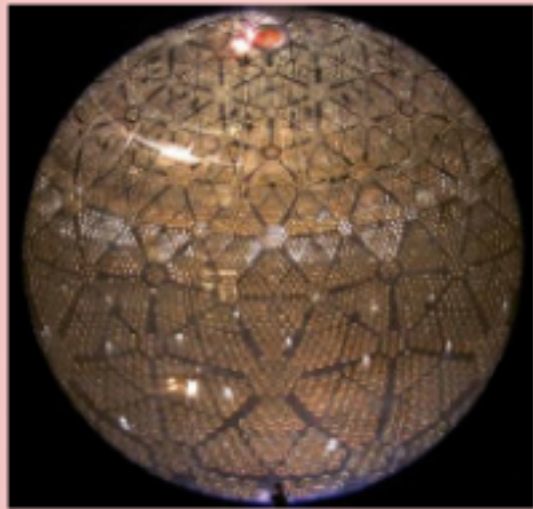


try to make the
background zero by
**tracking or
tagging**

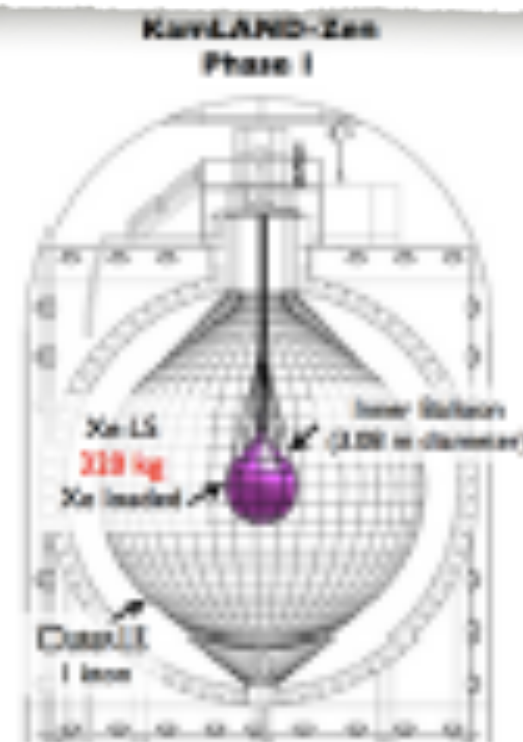
or better make the right cocktail of all of the above

the state of the art: brute force

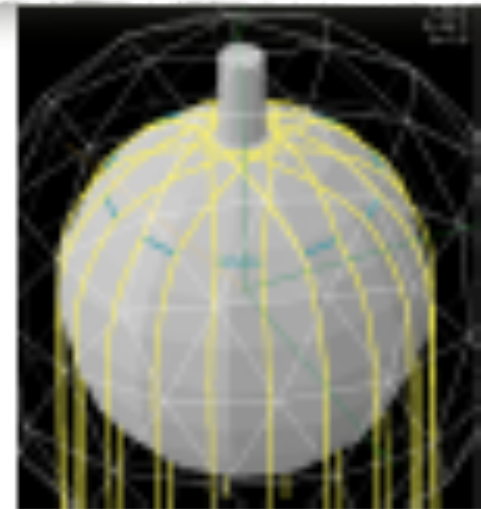
The “Brute Force” Approach



focus on the numerator
with a **huge amount
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(often sacrificing
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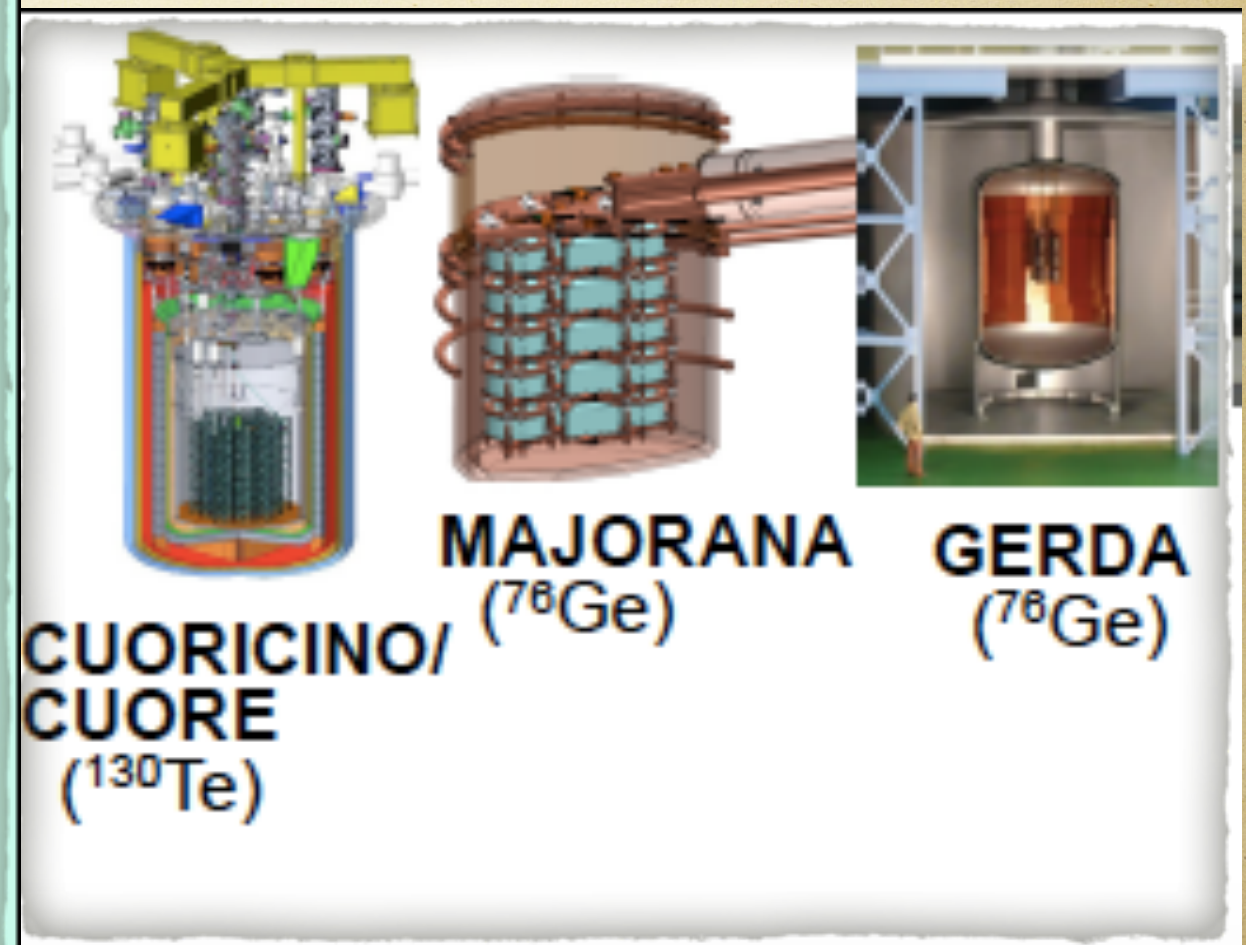
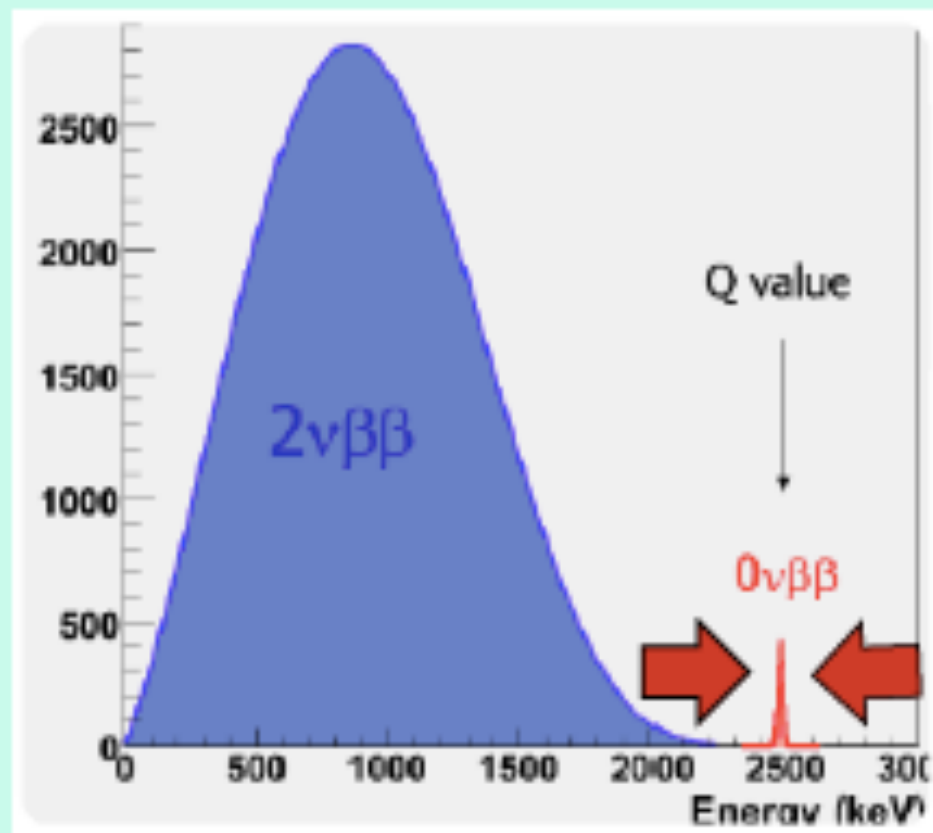
KamLAND-Zen
(^{136}Xe)



SNO+
(^{130}Te)

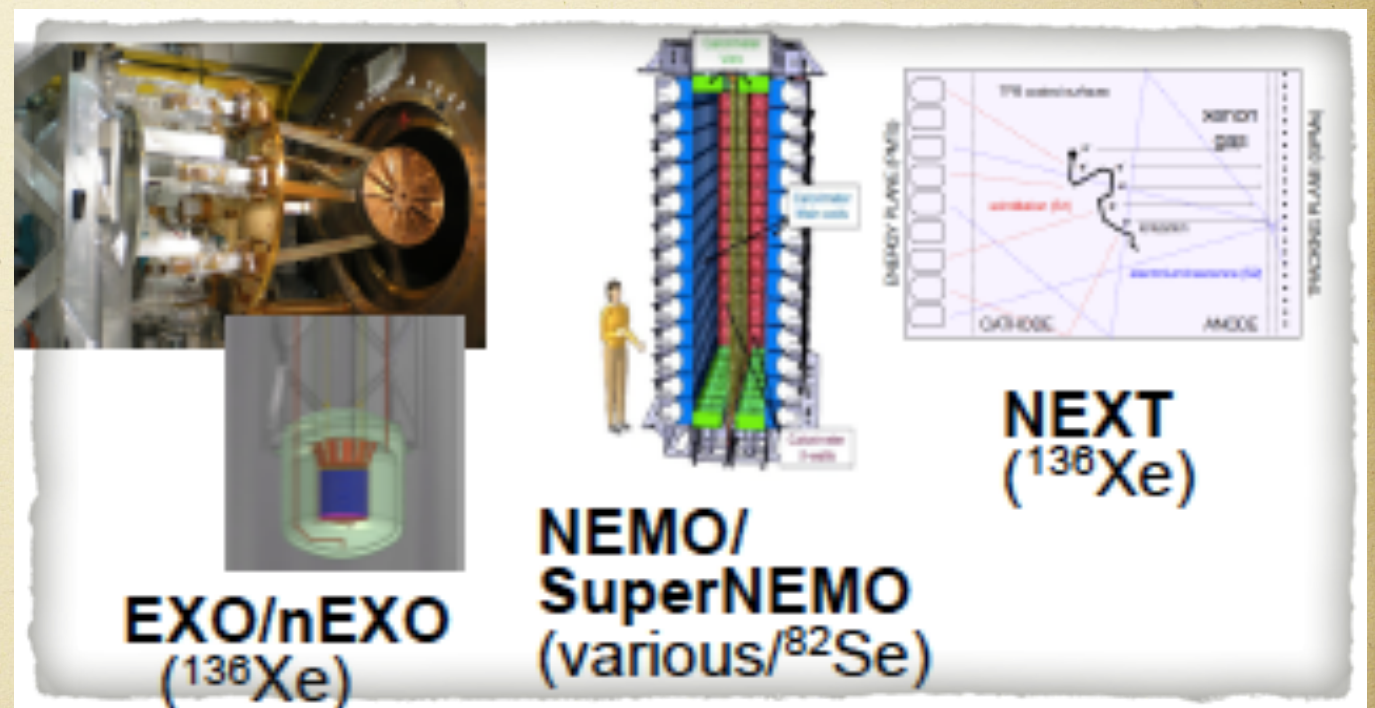
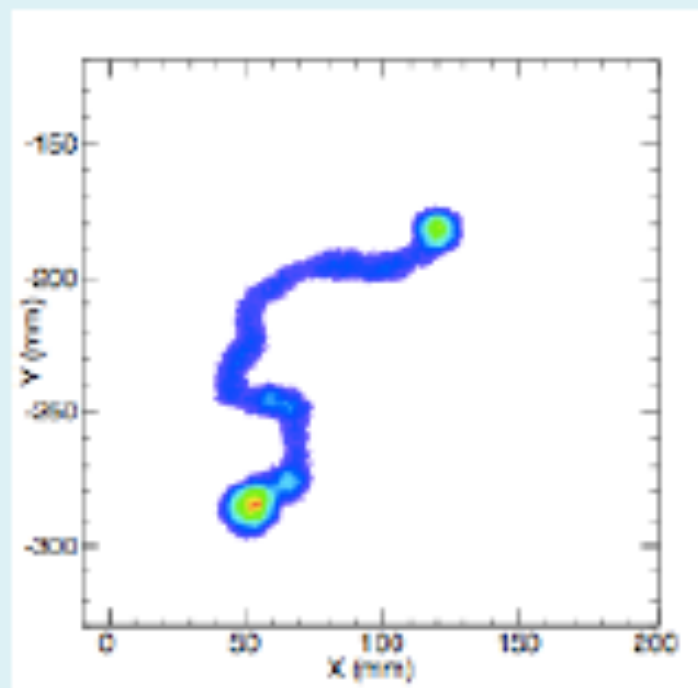
the state of the art: peak squeezer

The “Peak-Squeezer” Approach



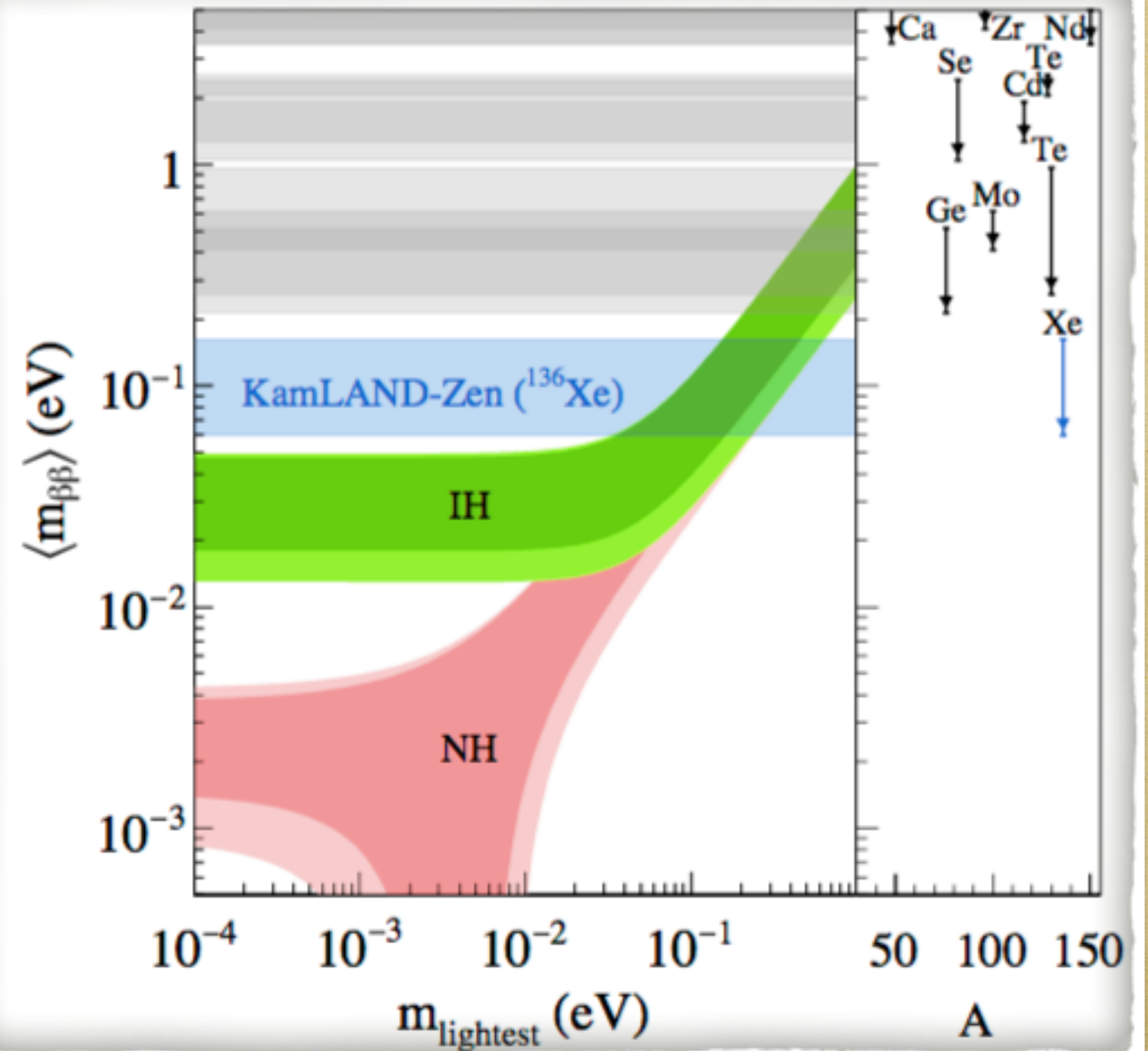
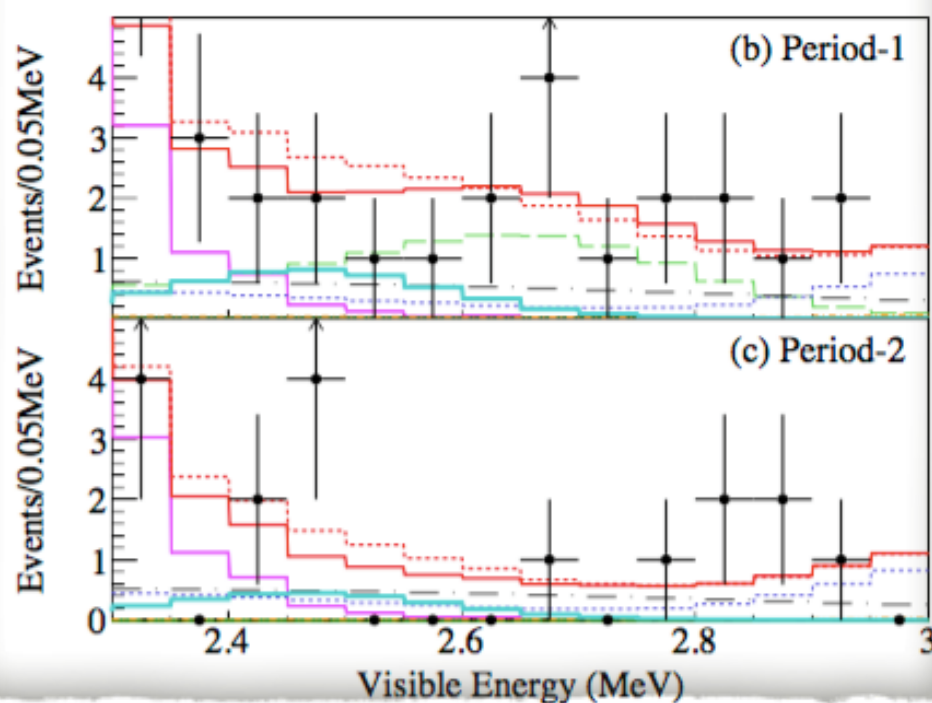
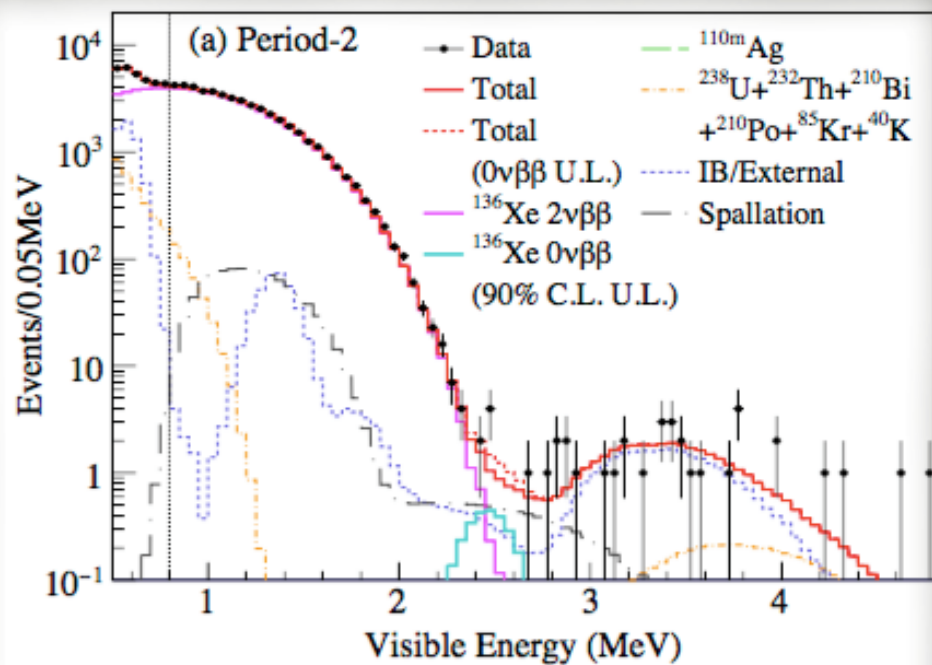
the state of the art: tracking

The “Final-State Judgement” Approach



so far the winner is

Search for Majorana Neutrinos near the Inverted Mass Hierarchy region with KamLAND-Zen



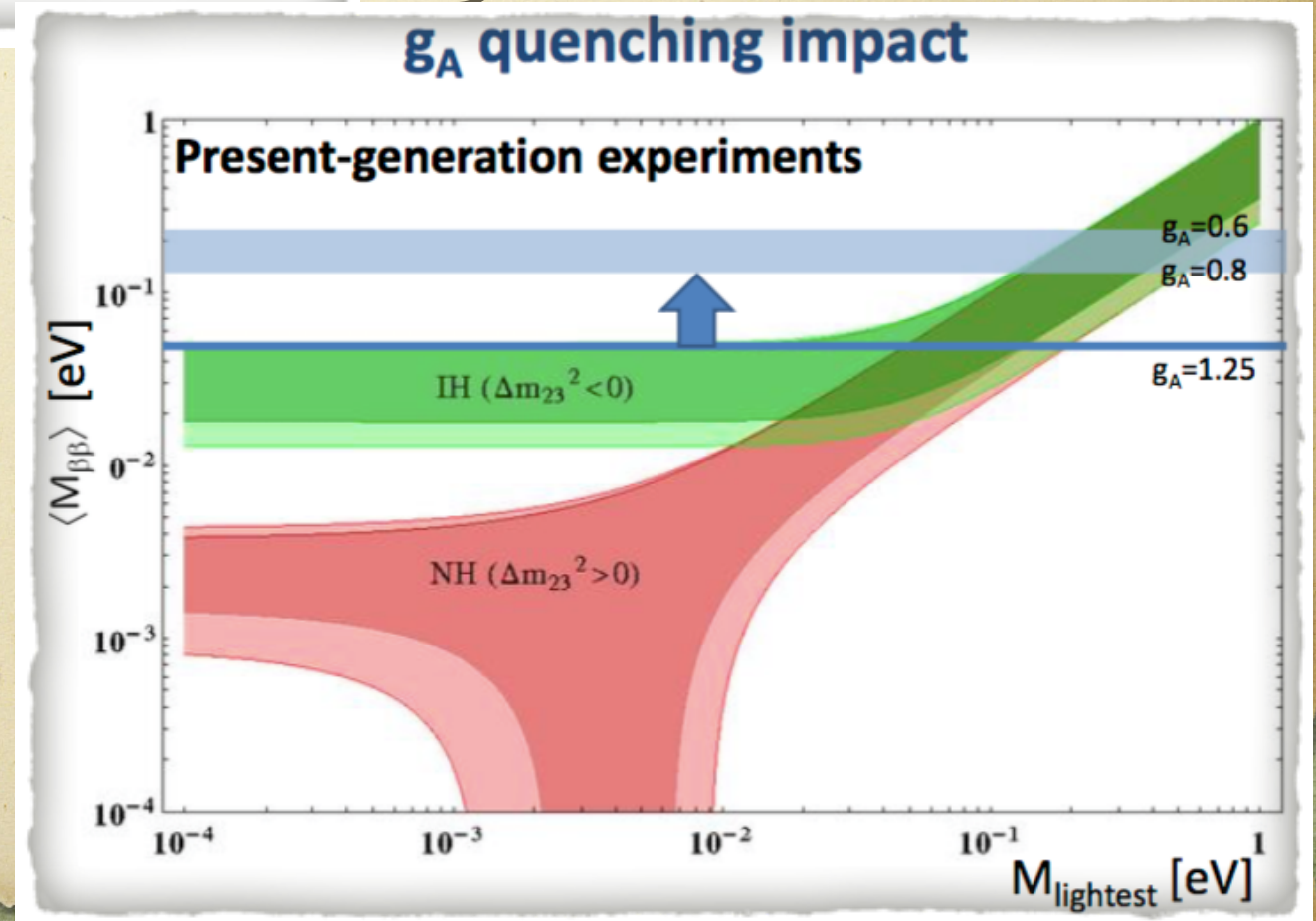
$$T_{1/2}^{0\nu} > 1.1 \times 10^{26} \text{ yr (90\% C.L.)}$$

something worrisome

$$\mathcal{M} \equiv g_A^2 \mathcal{M}_{0\nu} = g_A^2 \left(M_{GT}^{(0\nu)} - \left(\frac{g_V}{g_A} \right)^2 M_F^{(0\nu)} + M_T^{(0\nu)} \right)$$

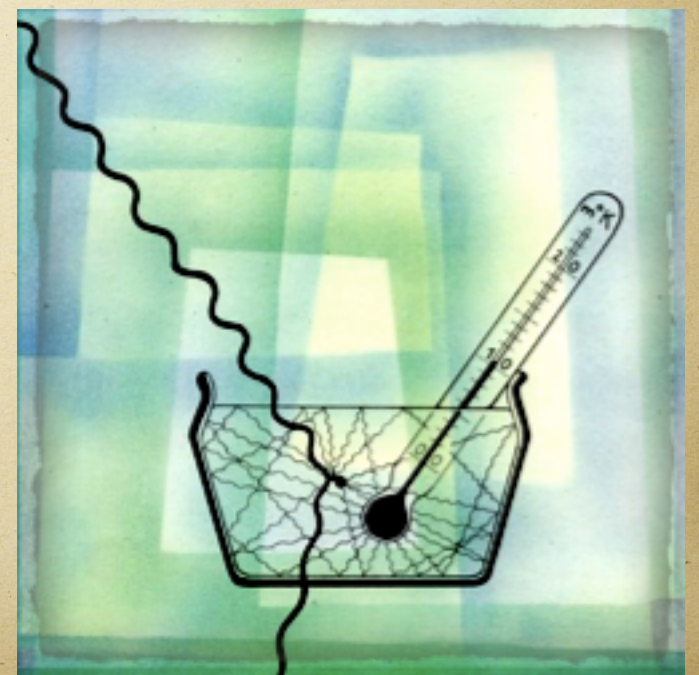
$$g_A = \begin{cases} g_{\text{nucleon}} & = 1.269 \\ g_{\text{quark}} & = 1 \\ g_{\text{phen.}} & = g_{\text{nucleon}} \cdot A^{-0.18} \end{cases}$$

$2\nu\beta\beta$
who knows ?



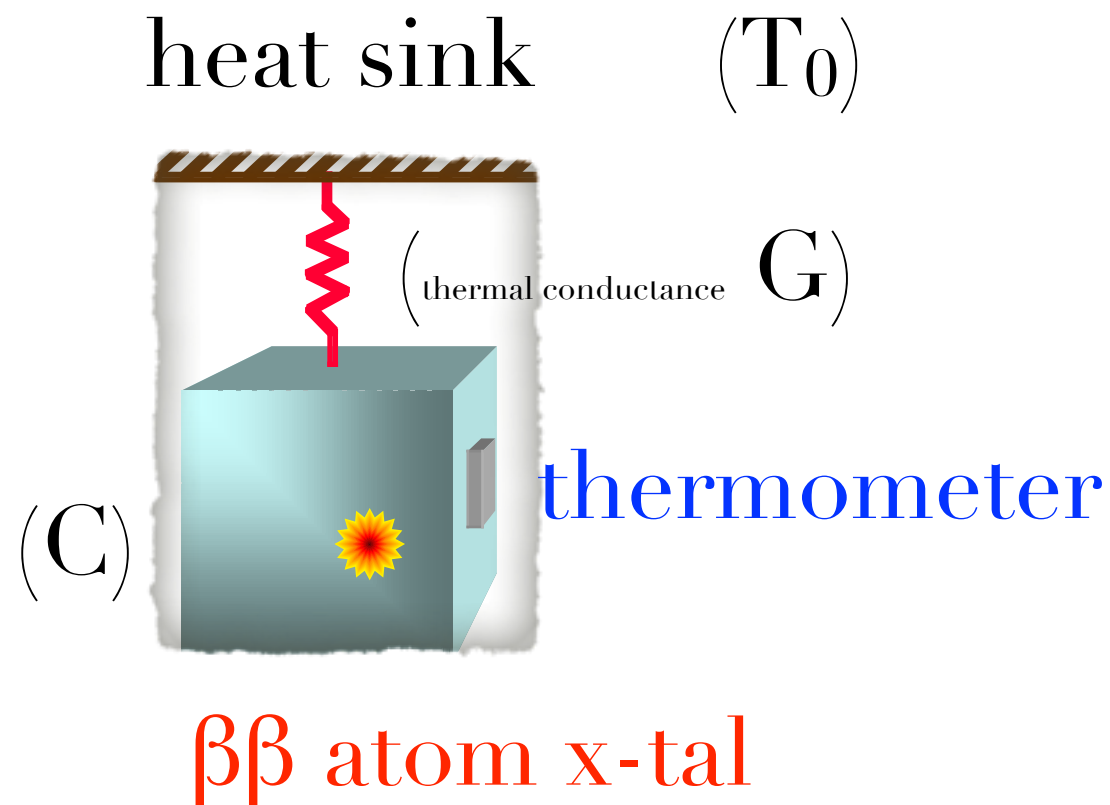
Bolometric technique as a competitor

- from MiBeta to CUORE via Cuoricino and Cuore-0
- Scintillating bolometers as an evolution toward Zero Background



(very) Low Temperature Calorimeter

A True Calorimeter



Basic Physics: $\Delta T = E/C$
(Energy release/ Thermal capacity)

Implication: $\text{Low } C \Rightarrow \text{Low } T$

Bonus: (almost) No limit to ΔE
($k_B T^2 C$)

Not for all apps : $\tau = C/G \sim 1\text{s}$

$$C(T) = \beta \frac{m}{M} \left(\frac{T}{\Theta_D} \right)^3$$

$$\Delta T(t) = \frac{\Delta E}{C} \exp \left(-\frac{t}{\tau} \right)$$

Why a bolometer

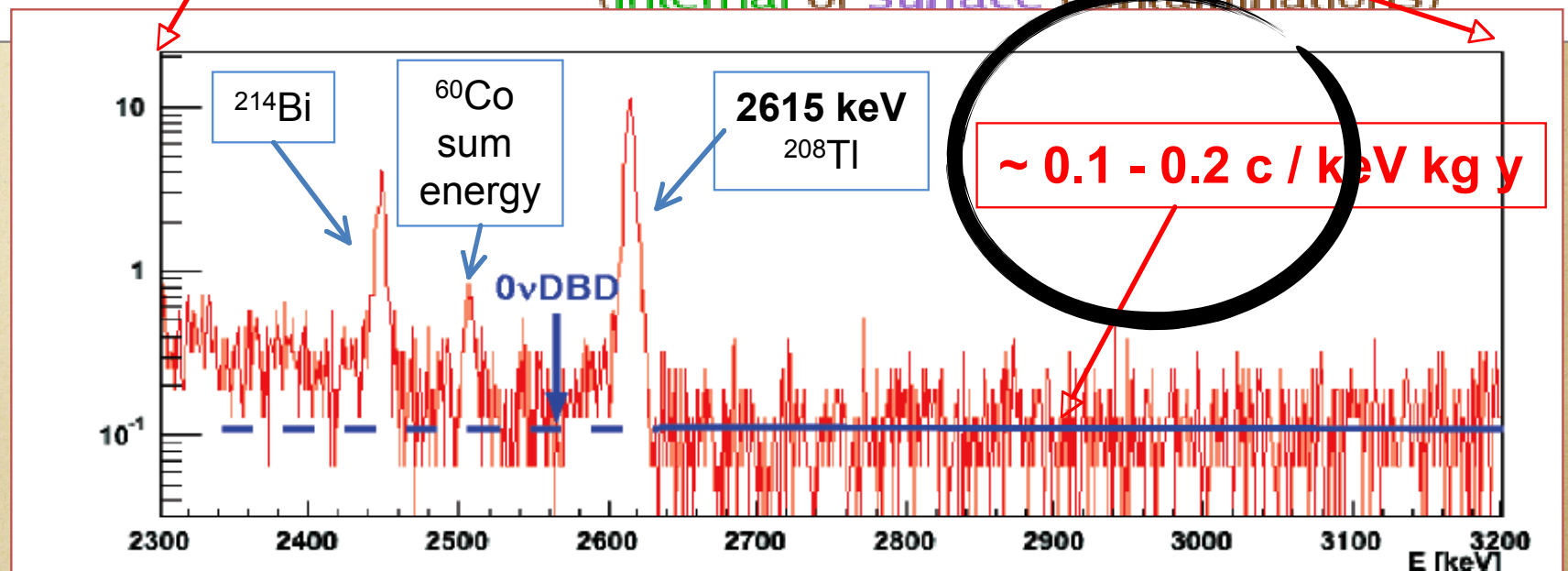
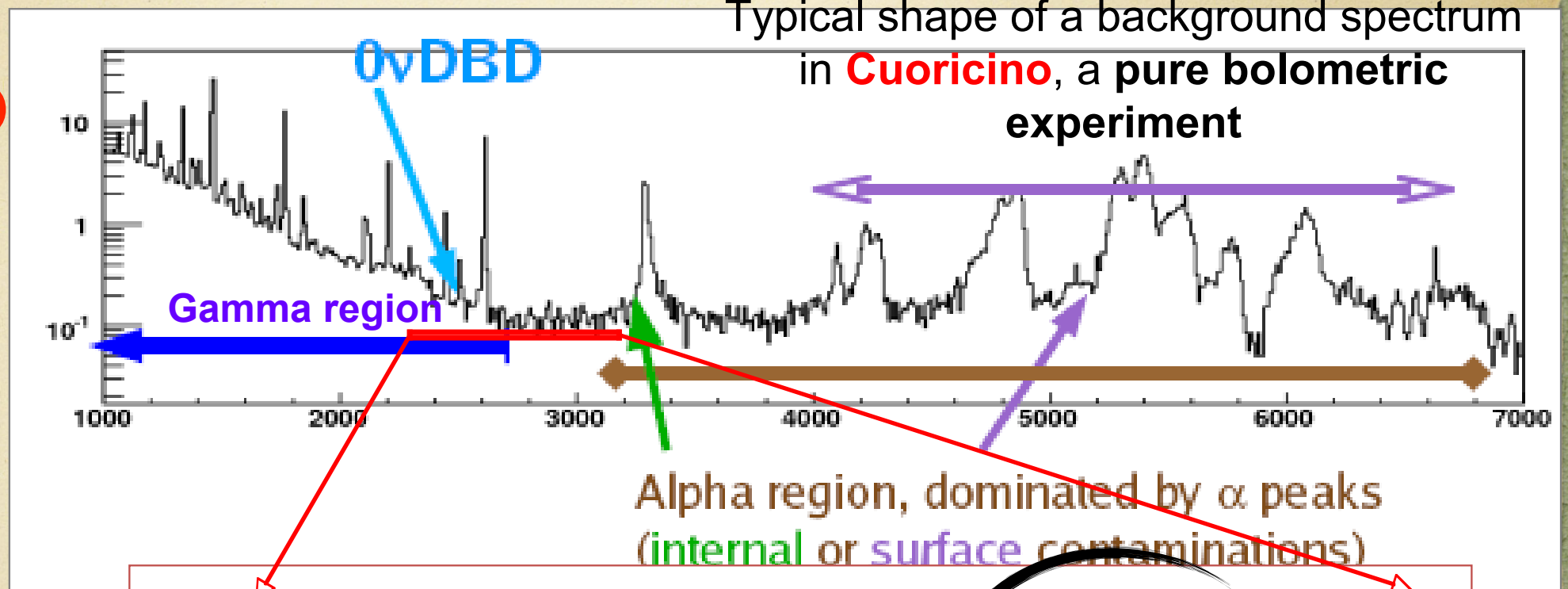
- M , t , B , ΔE are the parameters of the game
- t is irrelevant
- M is 'easy' with a calorimeter
- ΔE is a definite bonus
- B is what this part of the talk is mostly about

it is time to deal with the enemy:
what is the background?

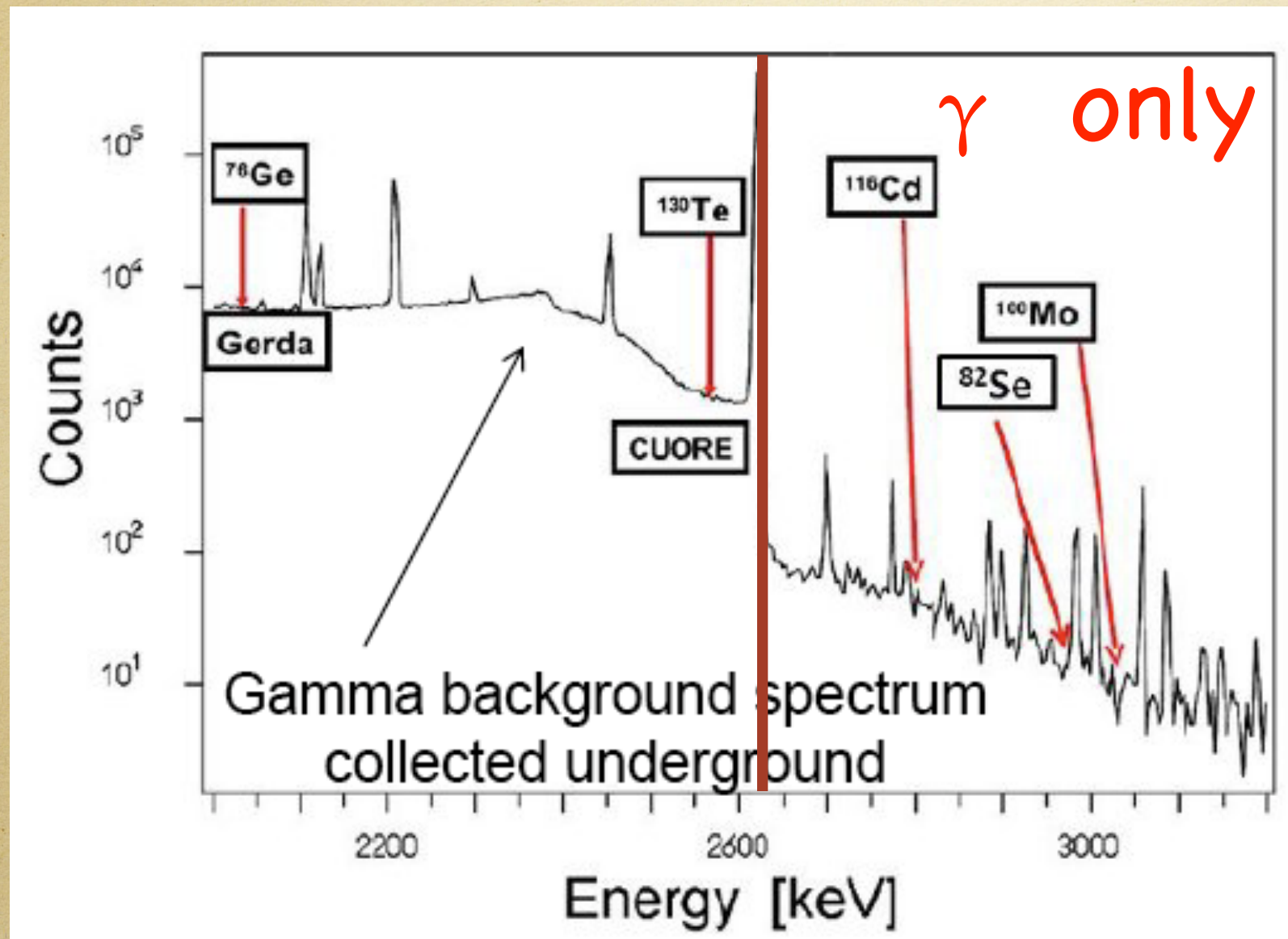
Cuoricino
 $b = 0.18 \pm 0.02$
 c/keV/kg/y

B
is
experiment
dependent.
Cuoricino
as an
example

Typical shape of a background spectrum
in **Cuoricino**, a pure bolometric
experiment



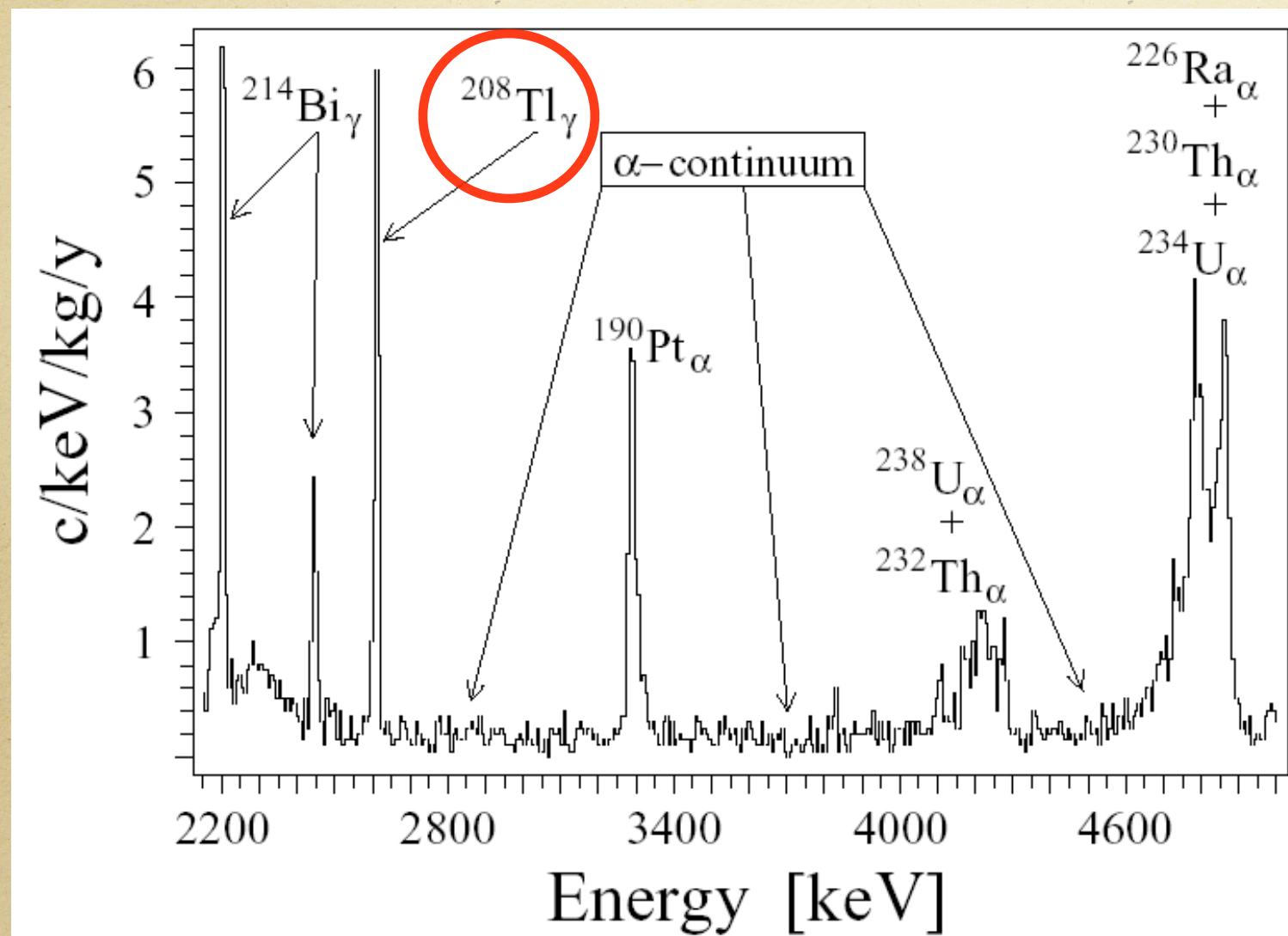
we have two enemies then



Photons

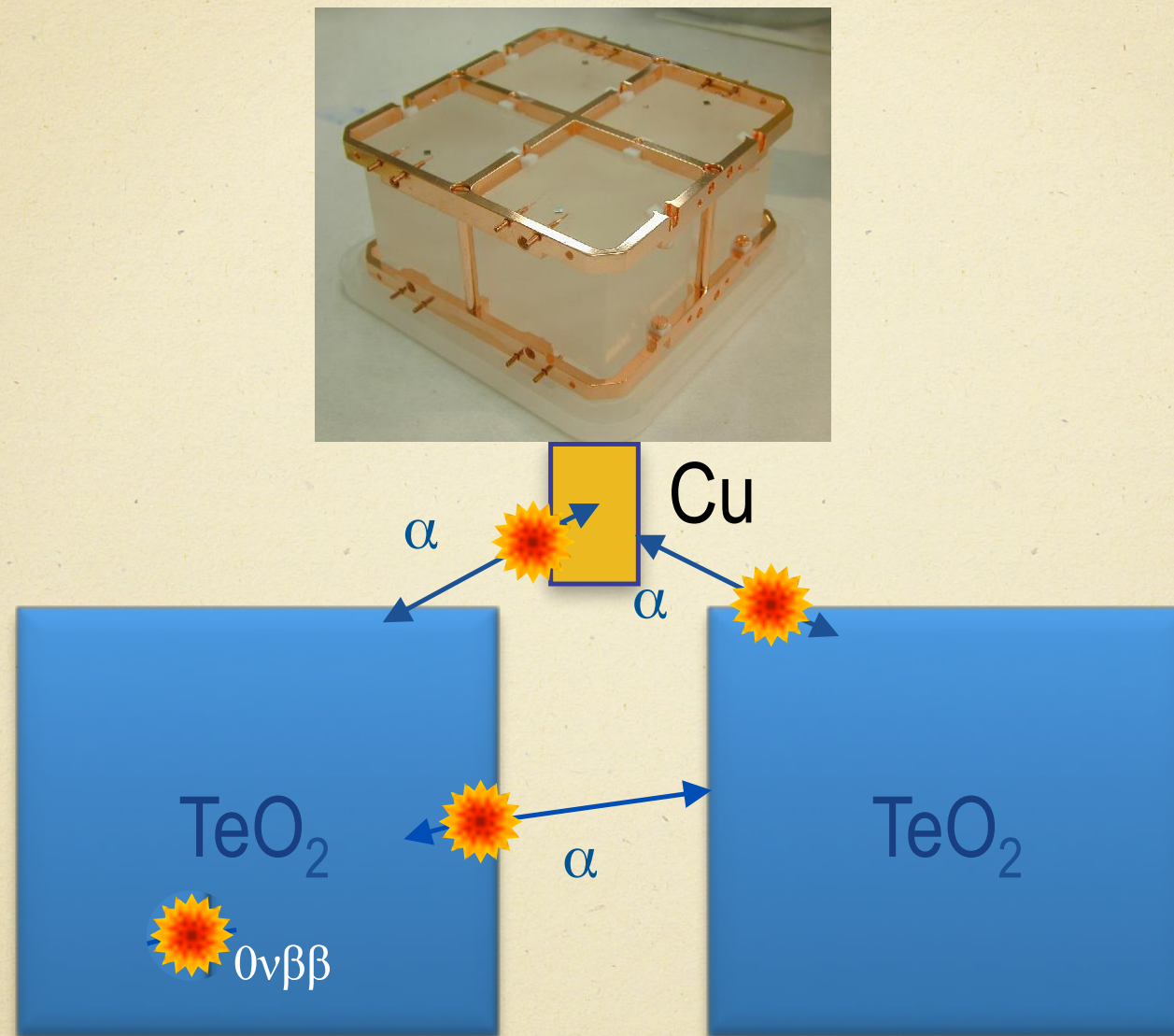
^{208}Tl is where photons start to disappear
Keep it in mind !

and...



the α land

the nasty α background



what is measured is part of the α energy
(if it were an internal emission...no problem !)
that induces a flat background

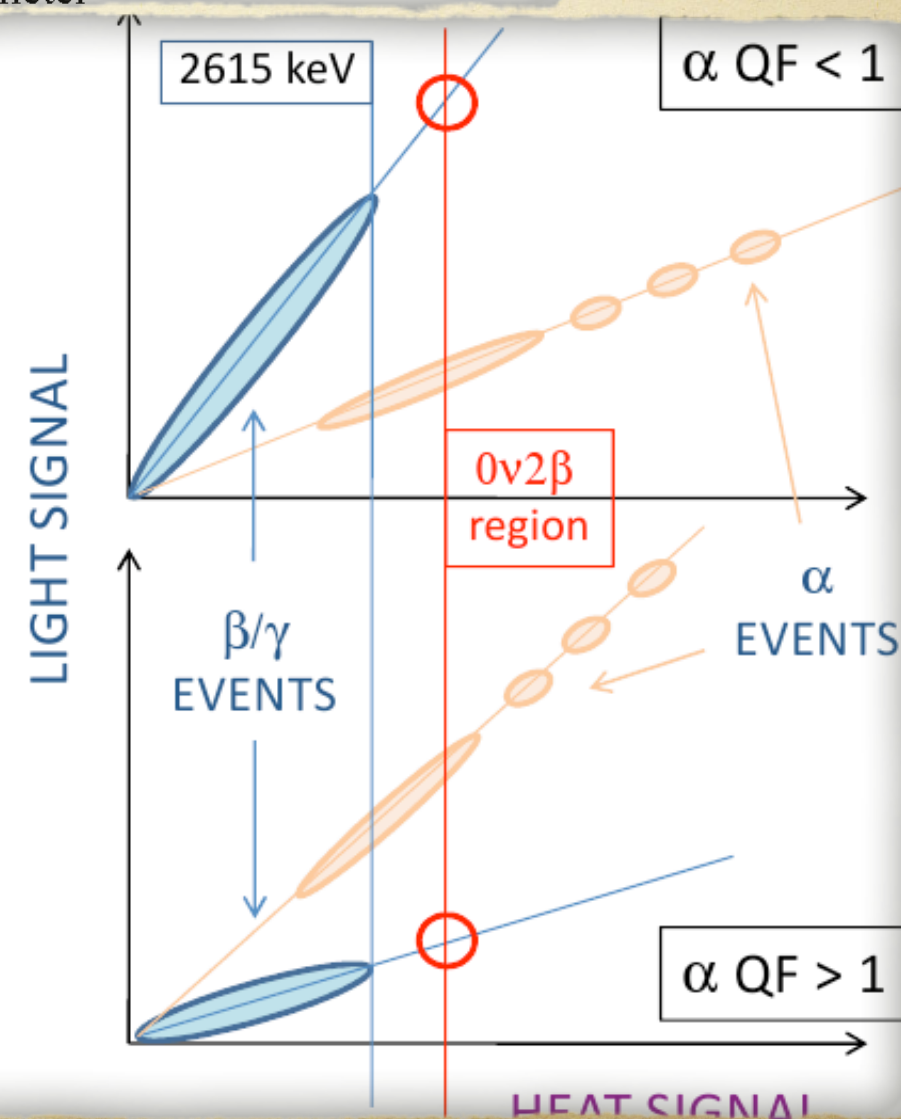
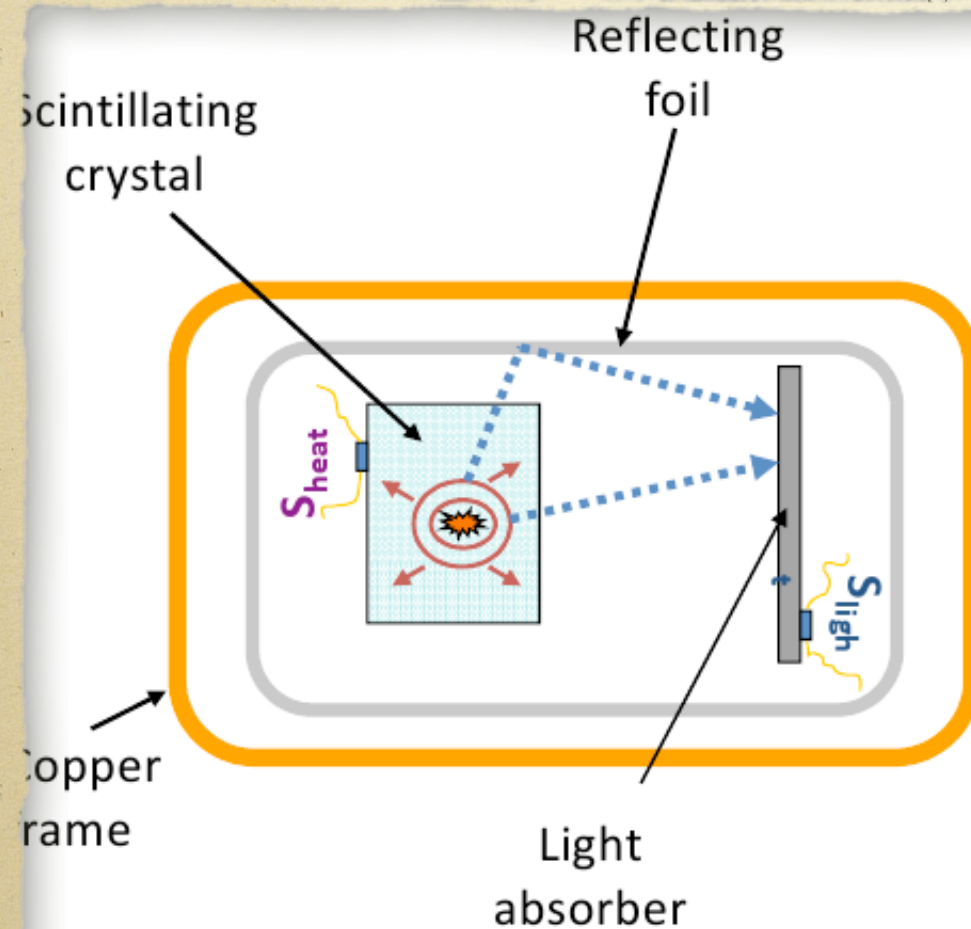
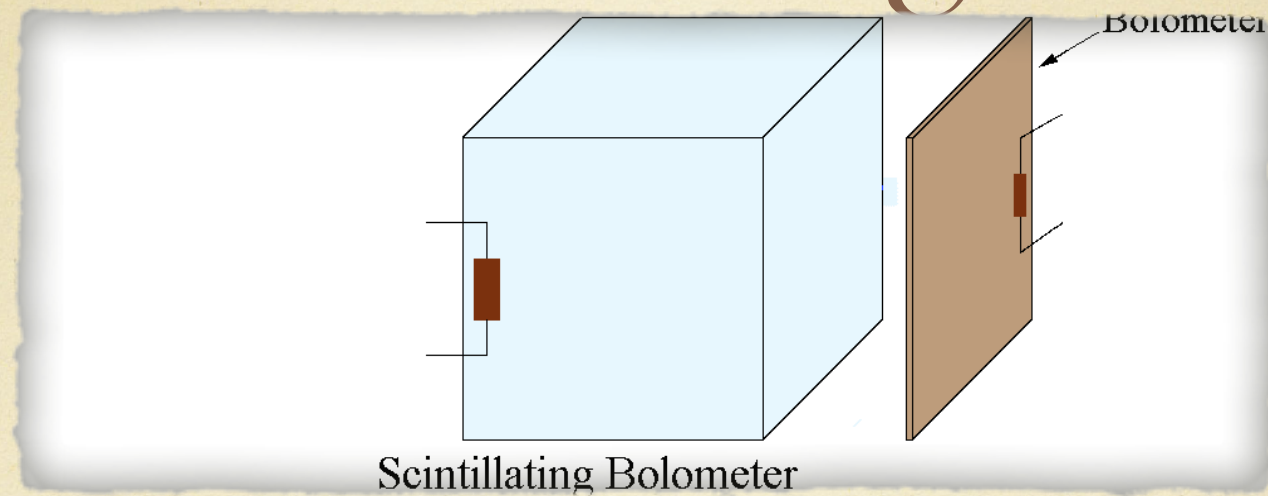
The LUCIFER concept

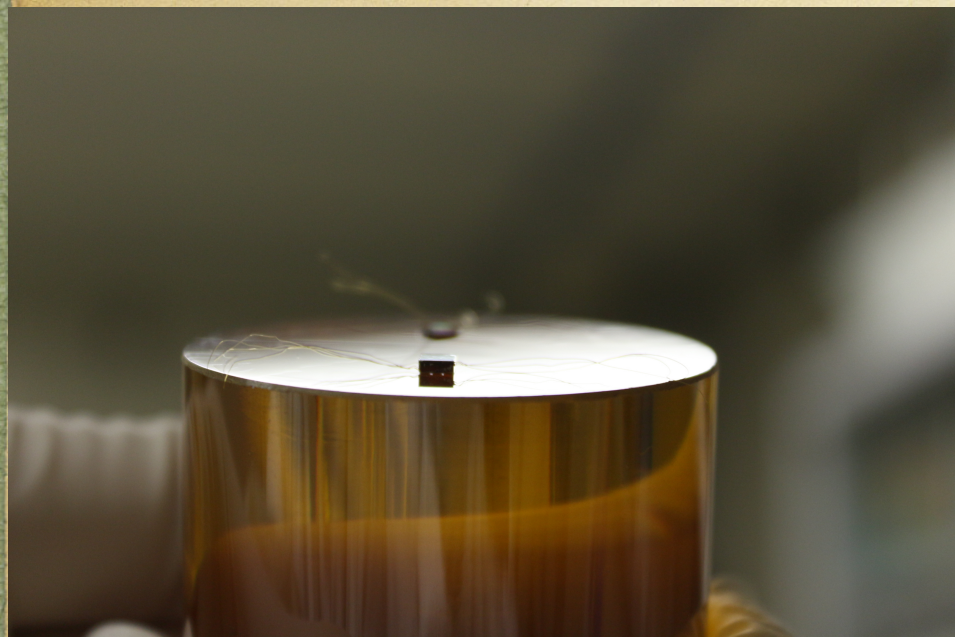
Lucifer is a Latin word (from the words *lucem* *ferre*), literally meaning "light-bearer", which in that language is used as a name for the dawn appearance of the planet Venus, heralding daylight.



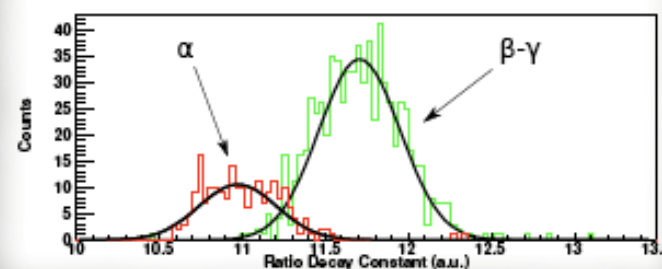
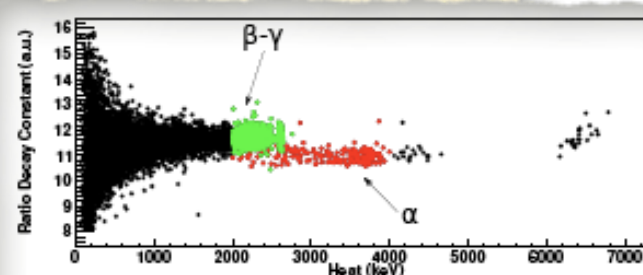
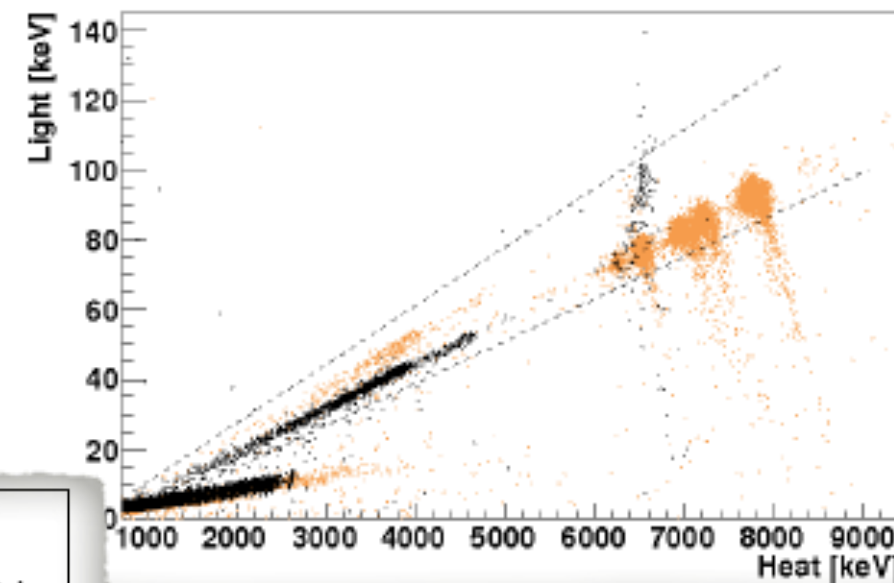
Bringing
light
underground

Heat & Light

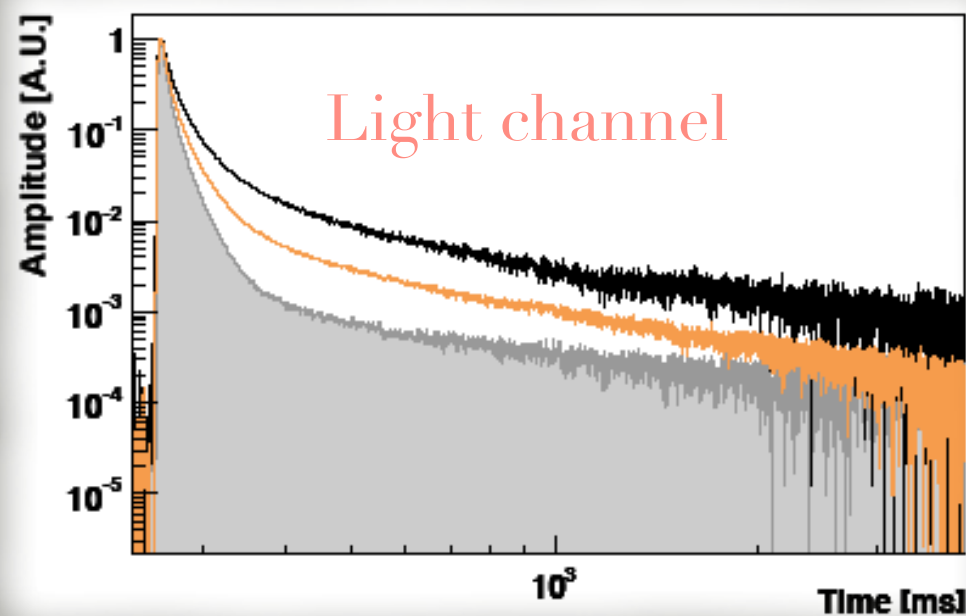




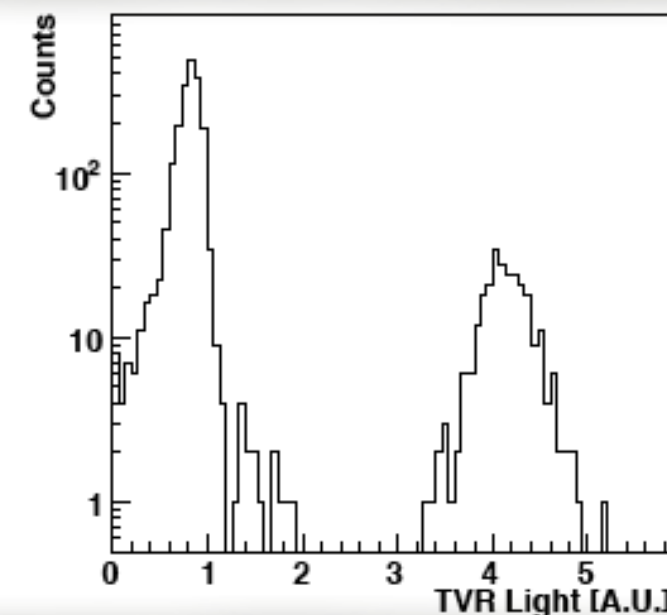
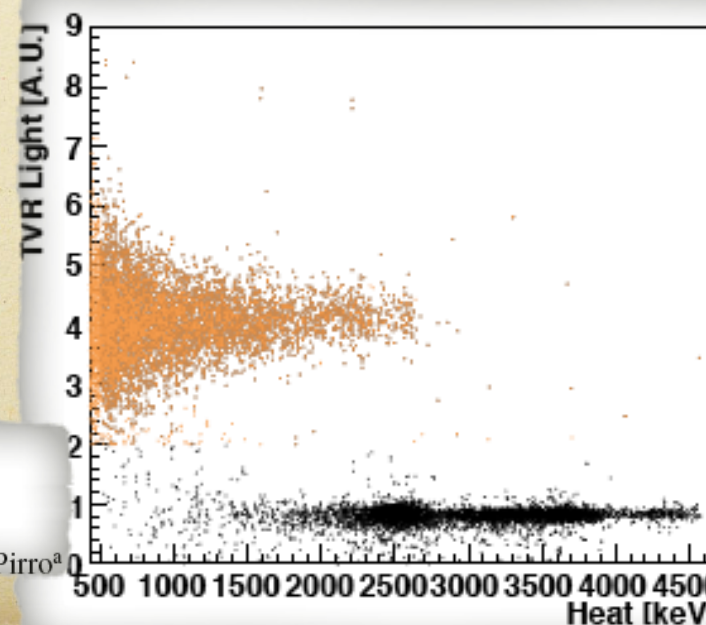
ZnSe



Heat channel



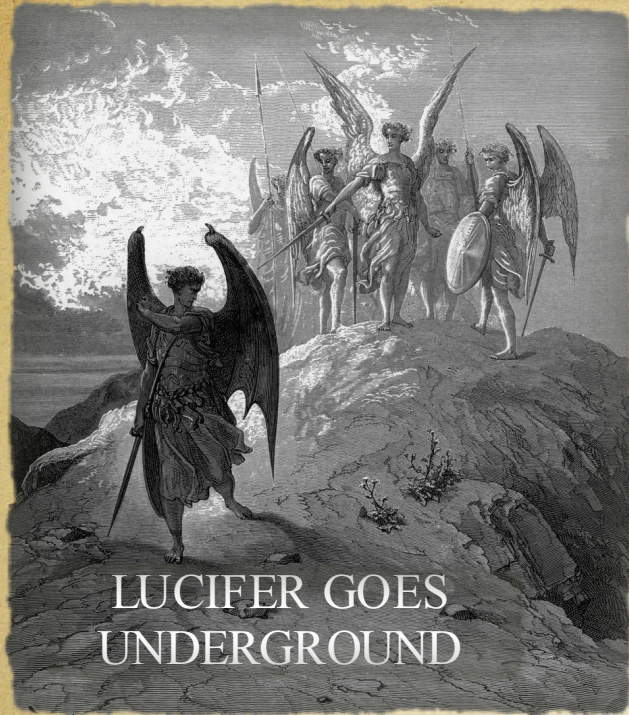
Light channel



ZnSe scintillating bolometers for Double Beta Decay

Astropart.Phys. 34 (2011) 344-353

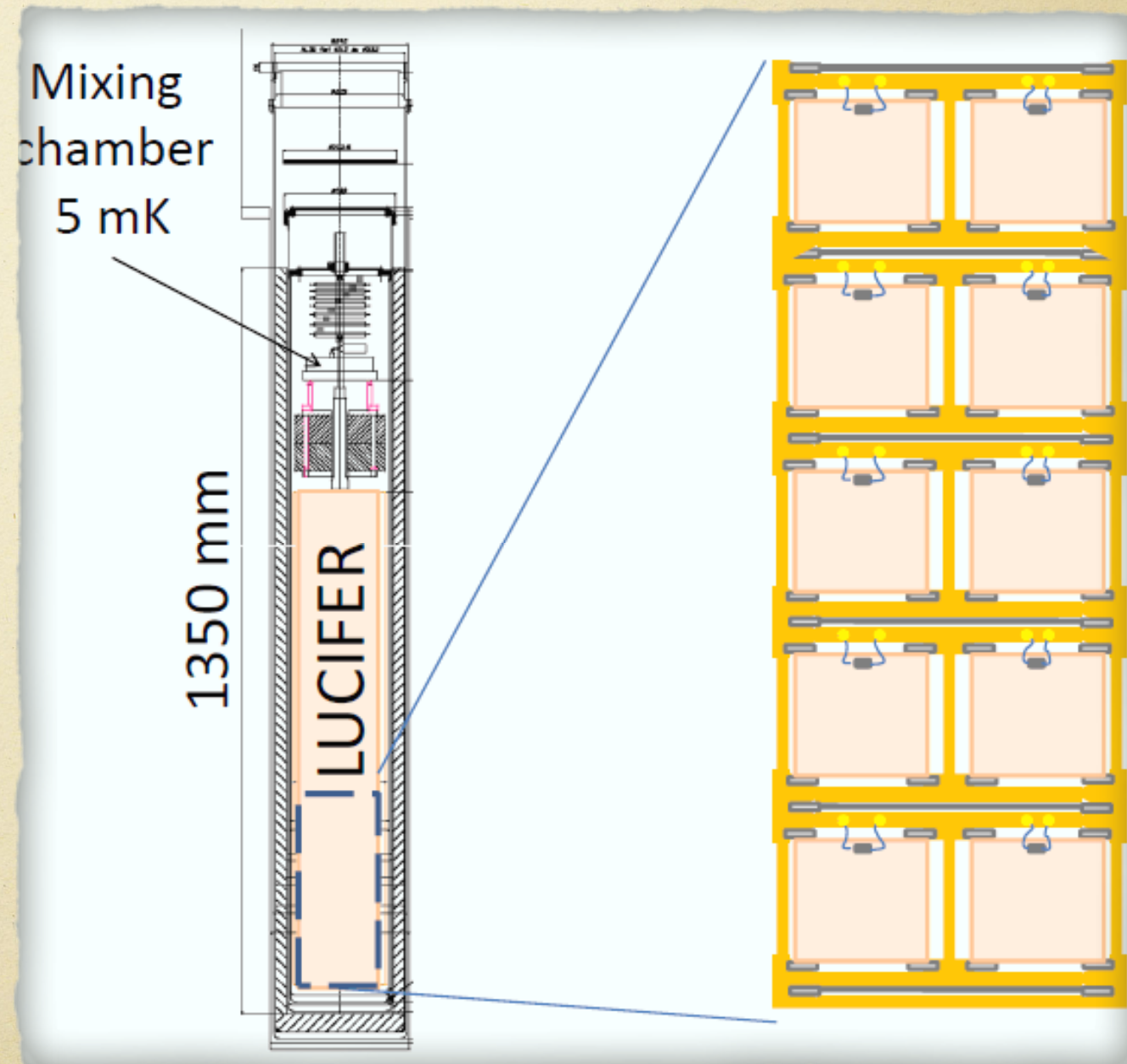
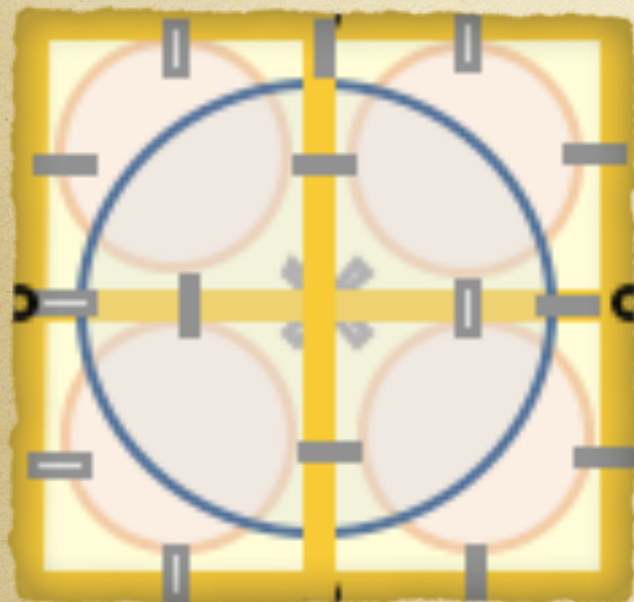
C. Arnaboldi^a, S. Capelli^{b,a}, O. Cremonesi^a, L. Gironi^{b,a}, M. Pavan^{b,a}, G. Pessina^a, S. Pirro^a



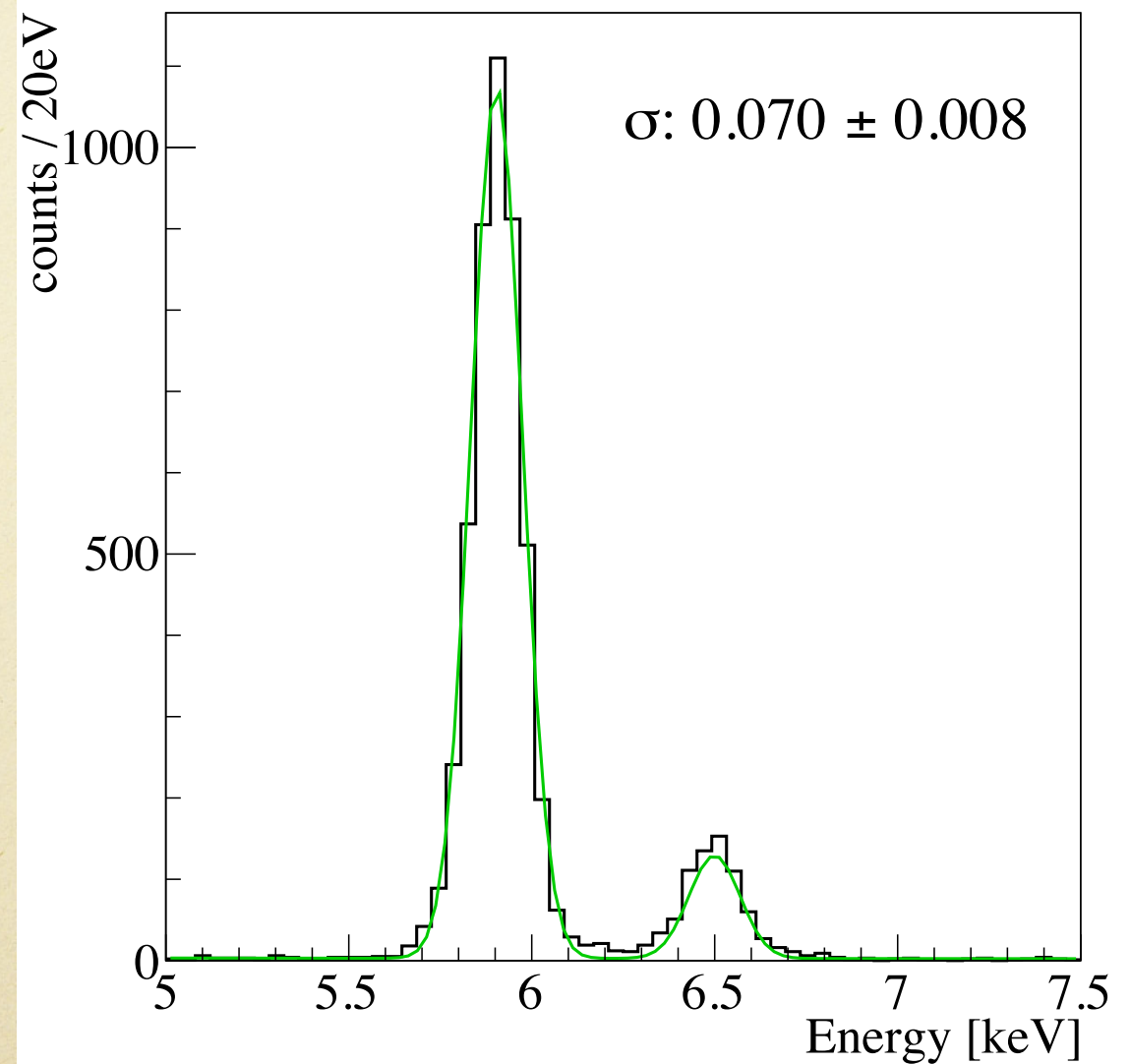
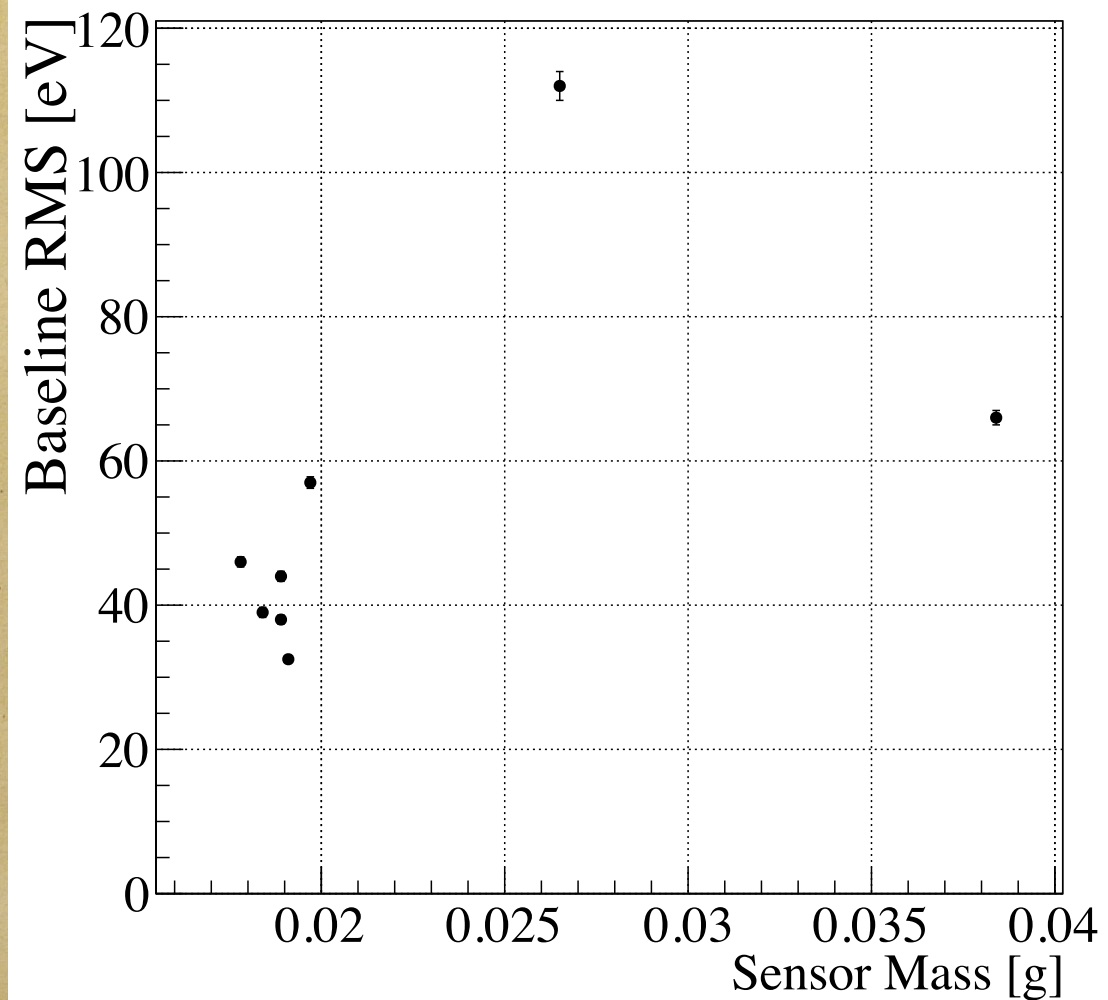
LUCIFER demonstrator

Low-background Underground Cryogenic Installation For Elusive Rates

ERC-2009-AdG 247115



with spectacular L.D.



goals

- assess background (has to be lower than 10^{-3})
- decide amongst ZnSe, ZnMo or LiMo
- project lifetime and mass limits to a realistic 1 Ton experiment that
- should be able to span the entire inverted hierarchy

if you want to know all...
pls. read:

<http://arxiv.org/pdf/1601.07512.pdf>

Neutrinoless double beta decay: 2015 review

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(Dated: January 28, 2016)

The discovery of neutrino masses through the observation of oscillations boosted the importance of neutrinoless double beta decay ($0\nu\beta\beta$). In this paper, we review the main features of this process, underlining its key role both from the experimental and theoretical point of view. In particular, we contextualize the $0\nu\beta\beta$ in the panorama of lepton-number violating processes, also assessing some possible particle physics mechanisms mediating the process. Since the $0\nu\beta\beta$ existence is correlated with neutrino masses, we also review the state-of-art of the theoretical understanding of neutrino masses. In the final part, the status of current $0\nu\beta\beta$ experiments is presented and the prospects for the future hunt for $0\nu\beta\beta$ are discussed. Also, experimental data coming from cosmological surveys are considered and their impact on $0\nu\beta\beta$ expectations is examined.

Conclusions

- Neutrino Physics is one of the leading field in HEP today
- Dirac or Majorana nature of neutrino mass is a fundamental question that needs to be answered at (almost) all cost(s)
- Neutrino-less DBD might possibly be the sole chance to give a measure of neutrino mass
- The second generation experiments might not be enough to win.
- We have to prepare for third generation. **Toward 0 background.**