



Experimental studies of neutrino nuclear responses for double beta decays

Hiro Ejiri

RCNP Osaka

2015.12, Catania

Thank the organizers for the invitation.

1. DBD ν -mass sensitivity and nuclear responses
2. Nuclear, photon and muon charge exchange reactions for NME $M^{0\nu}$
3. Solar- ν interactions with DBD detectors
5. Concluding remarks

A view from the Ejiri-weekend house



1. Neutrino mass sensitivity for DBD and neutrino nuclear responses

Majorana neutrinos and $0\nu\beta\beta$ decays

$$0\nu\beta\beta \quad \mathbf{A} = \mathbf{B} + \beta + \beta$$

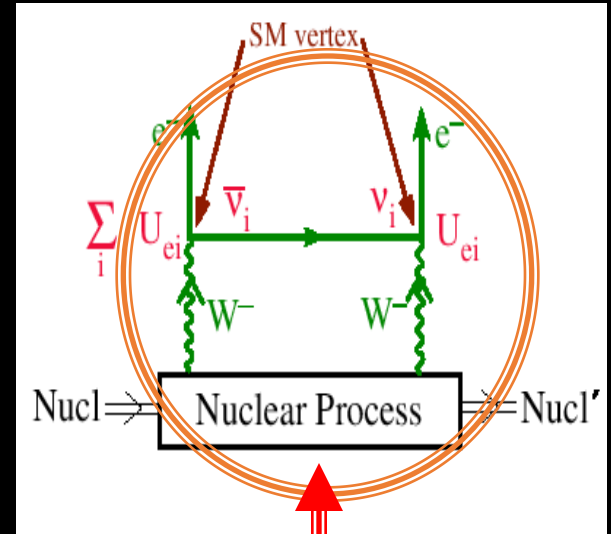
Lepton number $\Delta L=2$ beyond SM.

Particle astro physics
Majorana ν , m_ν CP

$$\mathbf{T}^{0\nu} = \mathbf{G}^{0\nu} [\mathbf{M}^{0\nu} m_\nu]^2$$

EXP

NNR Nucl. phys.
 g_A nuclear medium
 τ σ correlation



FEMTO(fm)-HC.

to enhance ν -exchange

Large luminosity
 $10^{83} \text{cm}^{-2} / \text{y}$ ton of $\beta\beta$ nuclei
 $\sigma \sim 10^{-83} \text{cm}^2$ IH

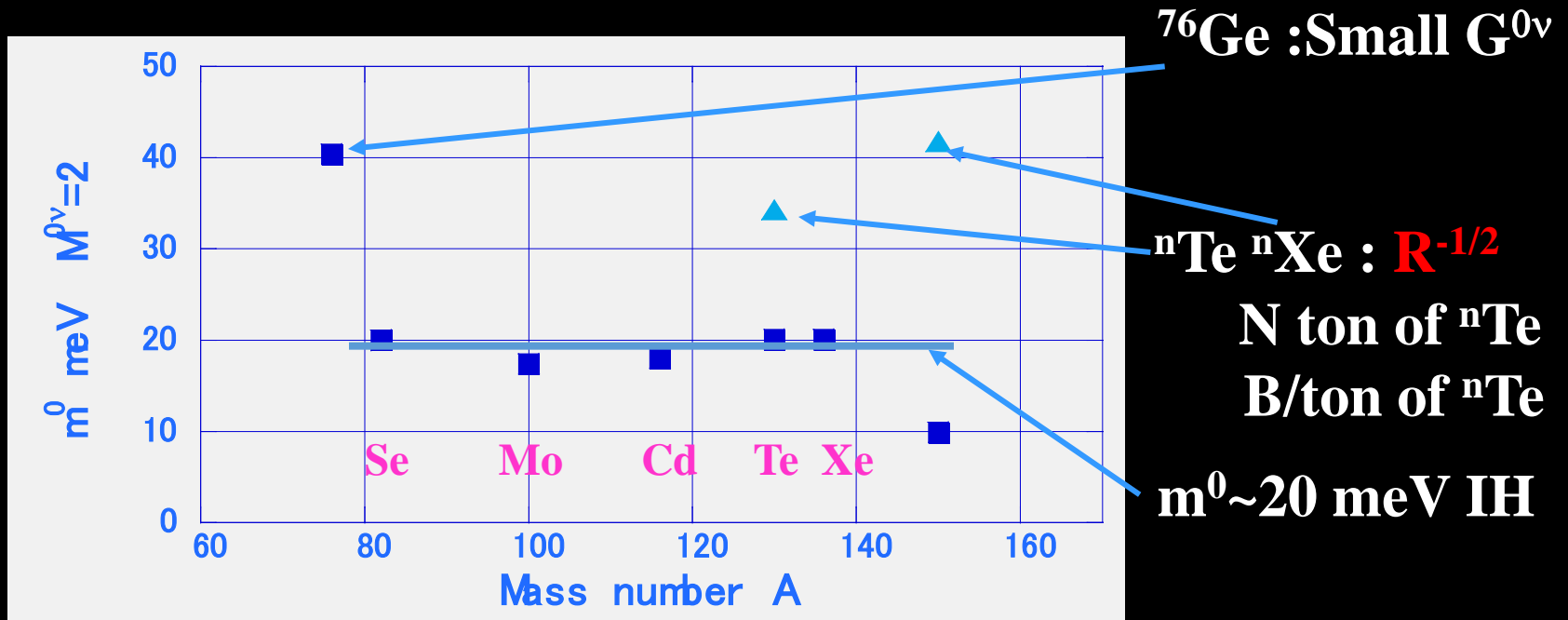
$0\nu\beta\beta$ mass sensitivity and NME

$$T = G [m_\nu M^{0\nu}]^2 \quad S = T(NT) > (BG)^{1/2} \quad BG = B NT$$

Mass sensitivity = ν -mass to be detected with 90% CL

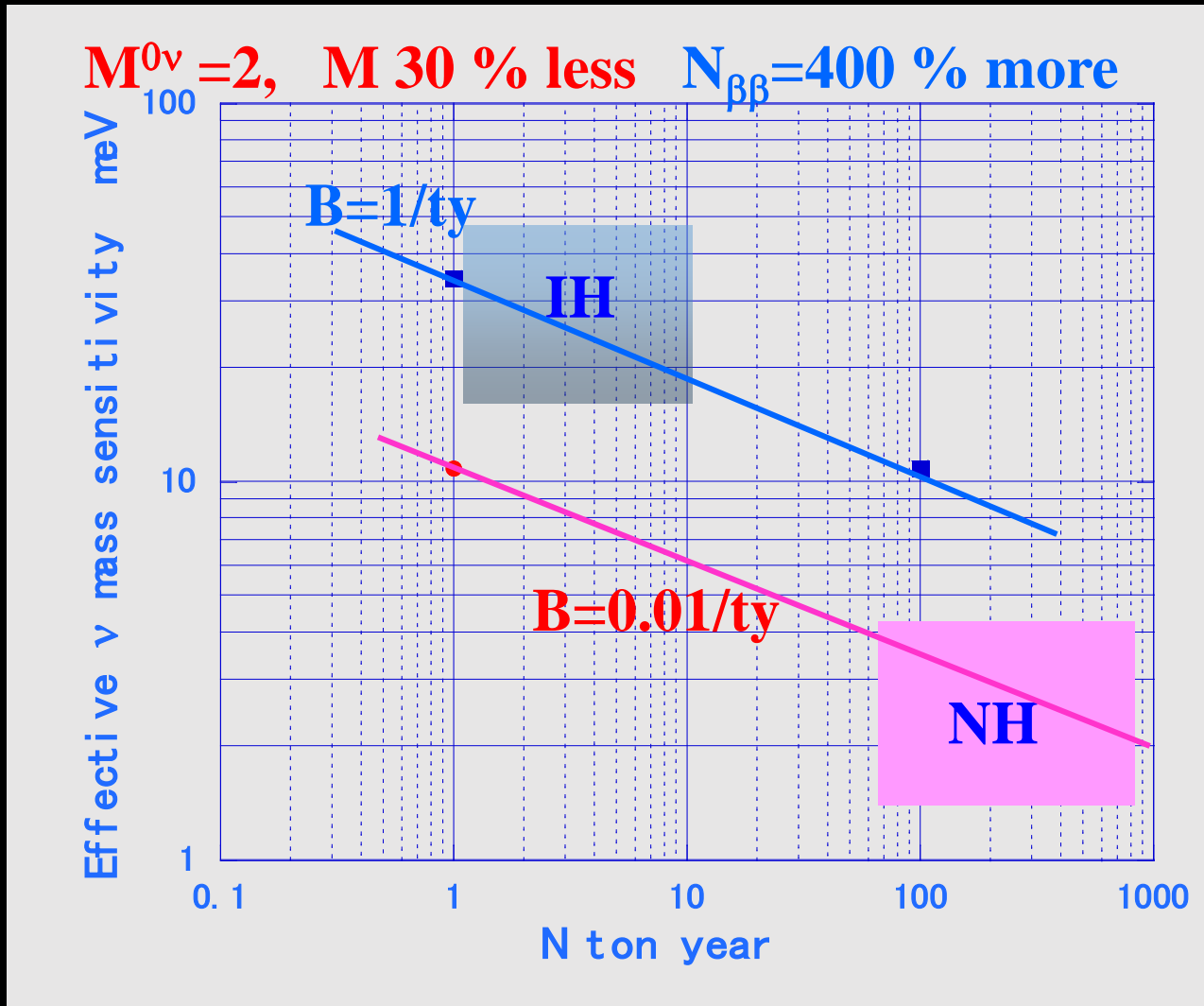
$$m_\nu = m^0 1.3 \varepsilon^{-1/2} [B/(NT)]^{1/4} \sim 2 m_0 \quad \text{if } B=1/\text{t y} \quad NT=1 \text{ t y}$$

$$m^0 = 39 \text{ meV} [2/M^{0\nu}]^{1/2} G^{-1/2} \quad G = G/0.01A$$



Mass sensitivities $\langle m_\nu \rangle = k [M^{0\nu}]^{-1} G^{-1/2} (NT)^{-1/4} (BG)^{1/4}$

Keys: $M^{0\nu} \sim 2$, $G \sim 4$, $NT \sim 10$ ty, $BG \sim 1/\text{ty}$ for IH



Nuclear matrix element for $0\nu\beta\beta$

Momentum transfer $0\nu\beta\beta$

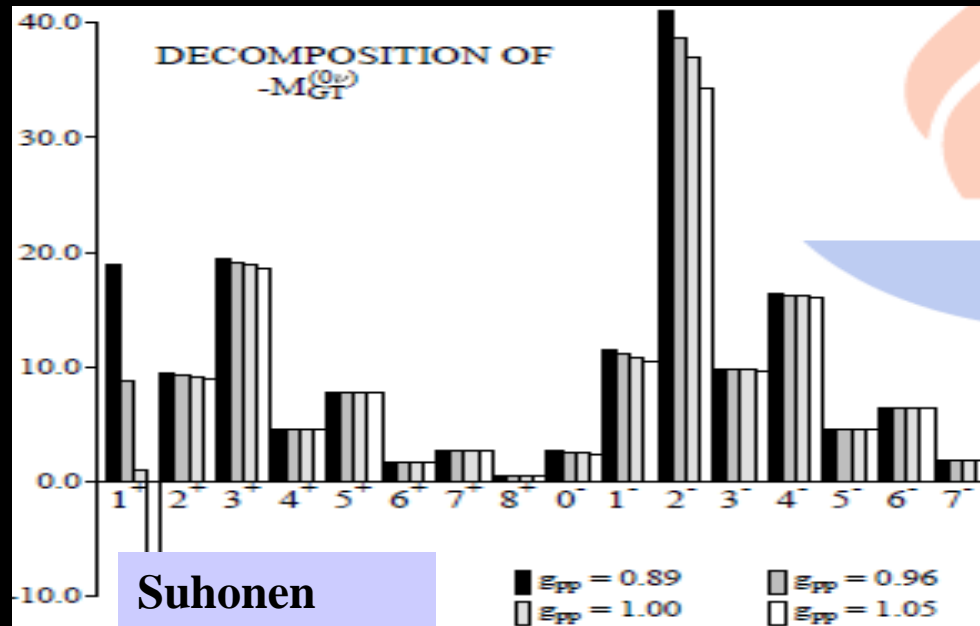
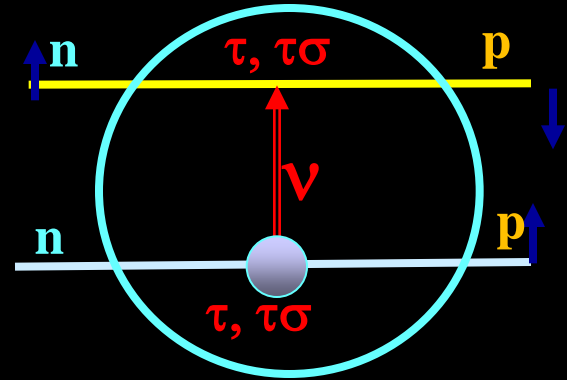
ν exchange

$$q \sim 1/\Delta r = 1-0.3 \text{ fm}^{-1}$$

$$\Delta l = qR = 1-2$$

$$J^\pi = 2^-, 2^+, 3^+$$

$$T(2^-) = g_A \tau [\sigma \times f(r) Y_1]_2$$



Nuclear matrix element for $0^+ \Rightarrow 0^+$

$$M^\beta = \langle \tau \sigma^s \rangle, \quad M^{2\nu} = \langle \tau \sigma^s \tau \sigma^s \rangle, \quad M^{0\nu} = \langle \tau \tau \sigma^s \sigma^s h(r) \rangle \quad s=0, 1$$

1. Sensitive to $\tau \sigma$ correlation and nuclear medium effects,

$$T^{\beta\beta} = [M_{\beta\beta}]^2 \sim M_\beta^4 \quad \text{Detector mass} = N_{\beta\beta} \sim M_\beta^8$$

$$g_F=1, \quad g_A=1.26 \Rightarrow 0.9 : T^{\beta\beta} \Rightarrow 1/4, \quad N_{\beta\beta} 1 \Rightarrow 16 \text{ tons}$$

2. Extremely small,

$$M_\beta \quad 10^{-1} M_\beta(\text{SP}) \sim 3 \cdot 10^{-2} M_\beta(\text{SUM})$$

$$M_\beta M_\beta \quad 10^{-2} M_{\beta\beta}(\text{SP}) \sim 1 \cdot 10^{-3} M_{\beta\beta}(\text{SUM})$$

3. $M = k^{\text{eff}} M(\text{SQP}) \quad k^{\text{eff}} = k^{\text{eff}}(\tau\sigma) k^{\text{eff}}(\Delta)$

$$k^{\text{eff}}(\tau\sigma) \sim 0.4 \quad \text{from } \tau\sigma \text{ polarization}$$

$$k^{\text{eff}}(\Delta) \sim 0.6 \quad g_A \Delta \tau\sigma \text{ polarization}$$

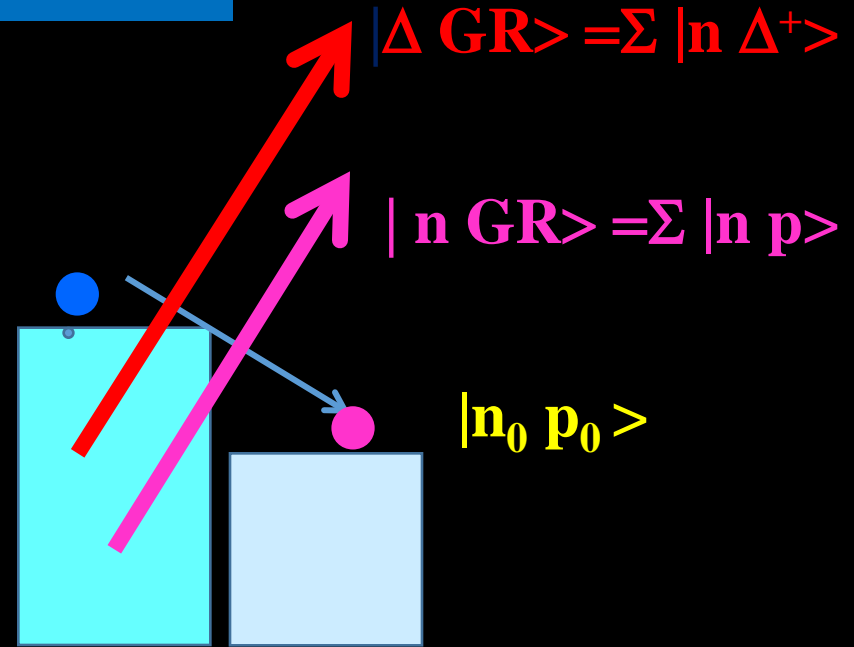
CERs provide exp. data to help/confirm $M^{0\nu}$ evaluations

Schematic view of $\beta\beta$ and GR

1. n_0 and p_0 are parts of the ground 0^+ state on the Fermi surface

2. $n\tau\sigma$ GR : coherent sum of many ($N\sim 30$) $\Sigma |n^{-1}p\rangle$

3. Δ GR: coherent sum of many ($N=60$) quark spin flip $\Sigma |n^{-1}\Delta^+\rangle$

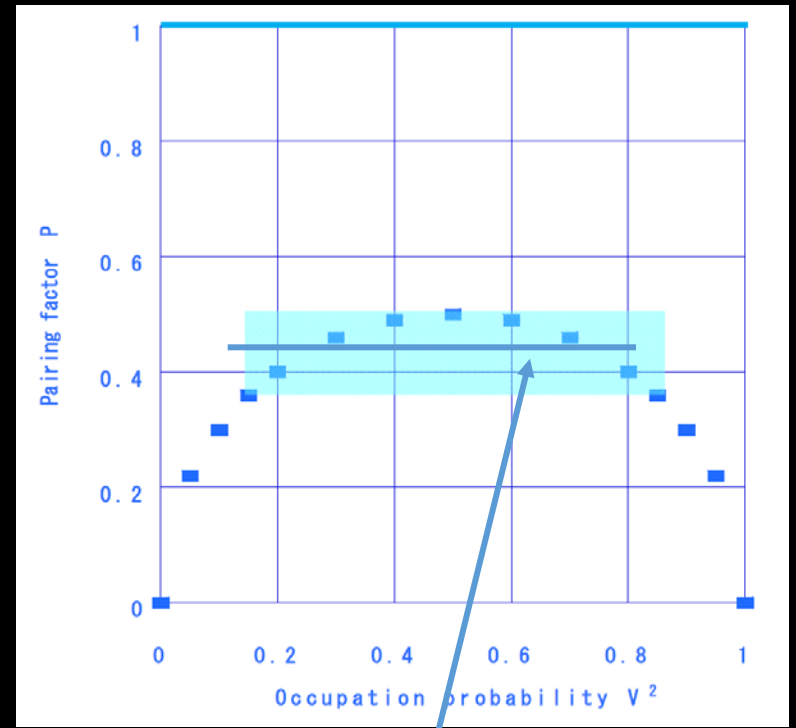
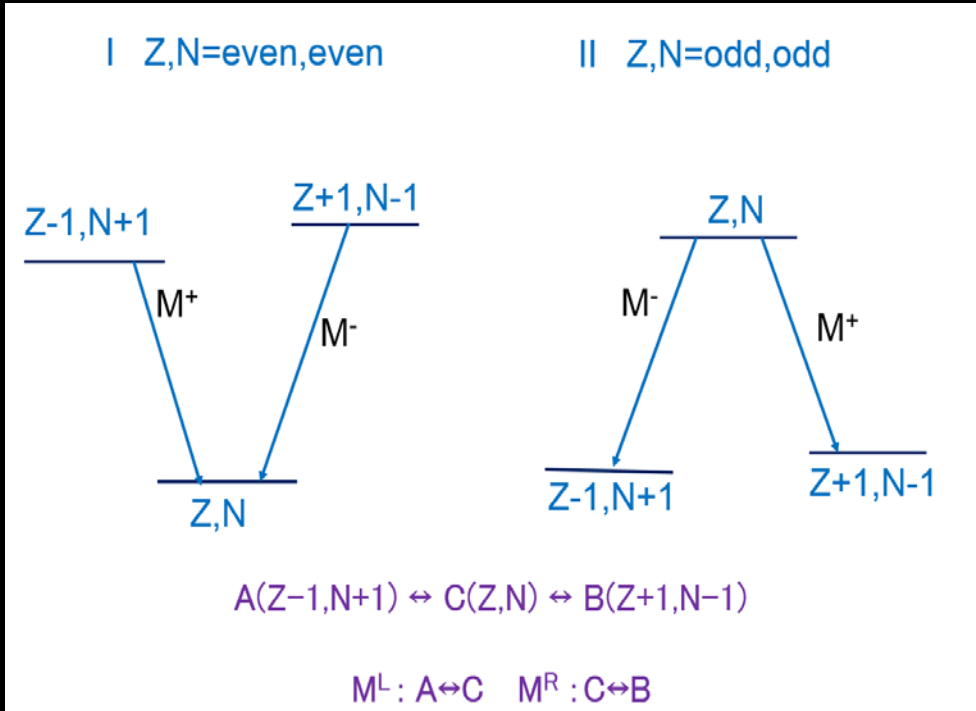


They mix destructively via repulsive interaction as

$$|np\rangle = |n_0 p_0\rangle - \varepsilon |n\tau\sigma GR\rangle - \delta |\Delta GR\rangle$$

GR and other effects are uniform, and are given by experimental renormalization of $k^{\text{eff}} = k^{\text{eff}} (\tau\sigma) \times k^{\text{eff}} (\Delta)$

$$M(GT) = \langle \tau^\pm (\sigma) \rangle \quad M(SD) = \langle \tau^\pm (\sigma Y_1) \rangle \quad \text{Ejiri Suhonen}$$



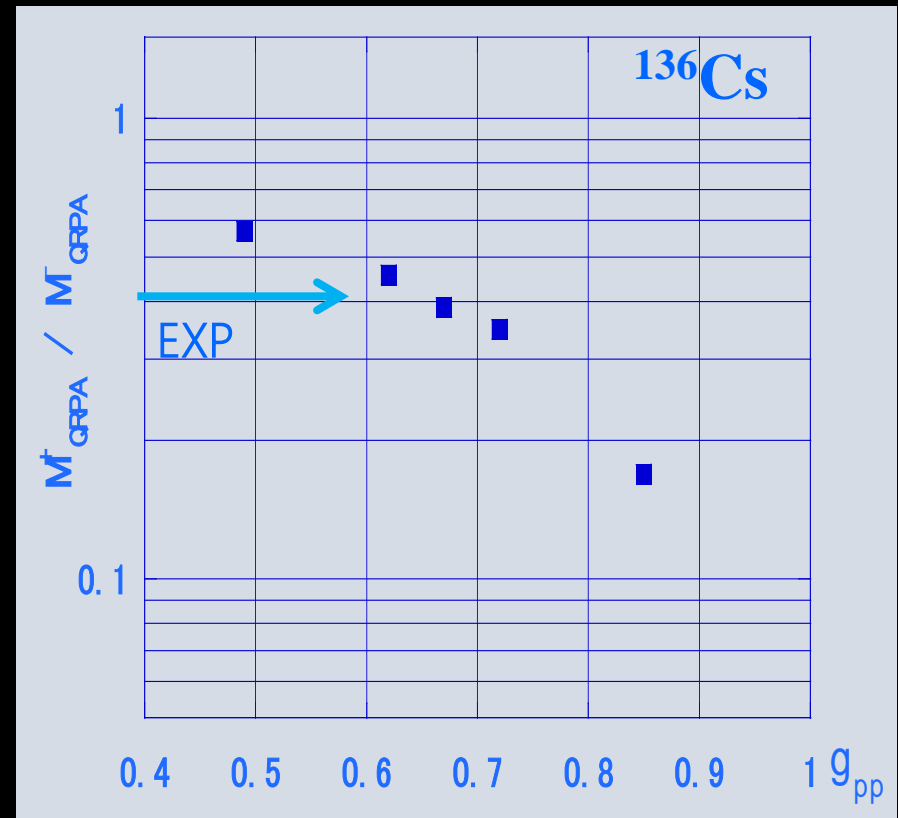
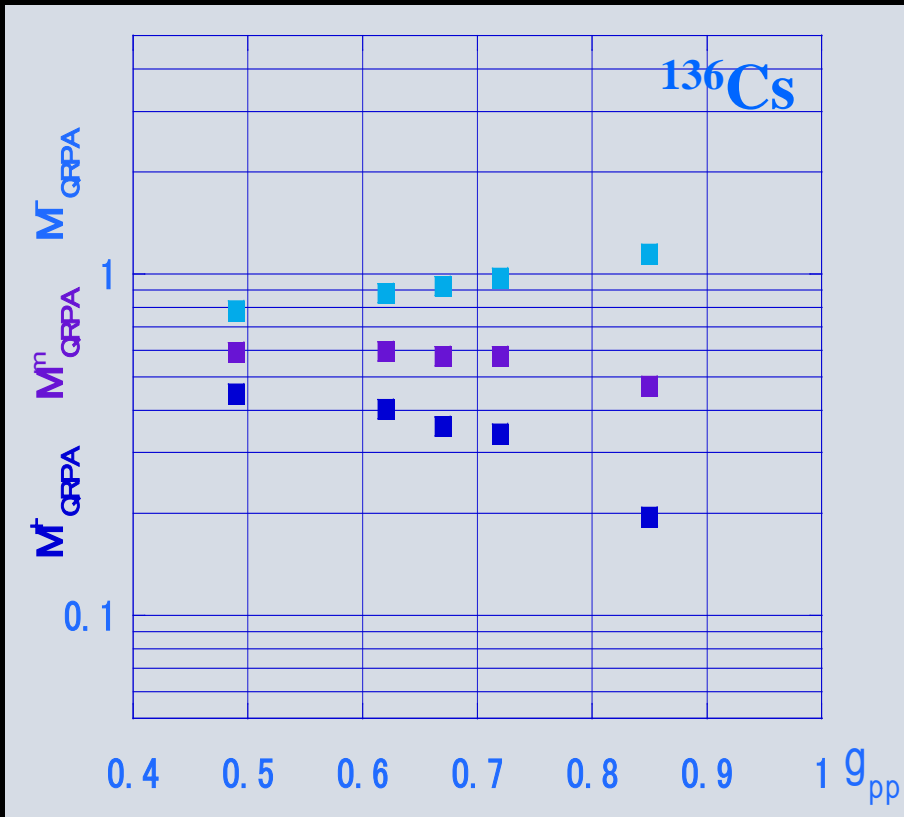
Geometrical mean $M^m = (M^+ \times M^-)^{1/2}$

1. $M^m(QP) = M^m(SP) [U_p V_n U_n V_p]^{1/2} \sim 0.43 M^m(SP)$

Insensitive to U & V nuclear surface effects and g_{pp}

2. NMEs in $\beta\beta$ are $(M^m)^2 = (M^+ \times M^-)$

QRPA M(GT) dependence on g_{pp} p-p interaction



M^m (violet) is independent of g_{pp} since the effects on M^+ & M^- cancel.

M^+ / M^- is sensitive to g_{pp} .
Exp. ratio gives $g_{pp} = 0.6-0.7$

GT 1⁺ $\tau\sigma$ NN & nuclear medium g_A

$$M_{\text{exp}}^m \ll M_{\text{QRPA}}^m \ll M_{\text{qp}}$$

$$M_{\text{exp}}^m = k M_{\text{qp}}$$

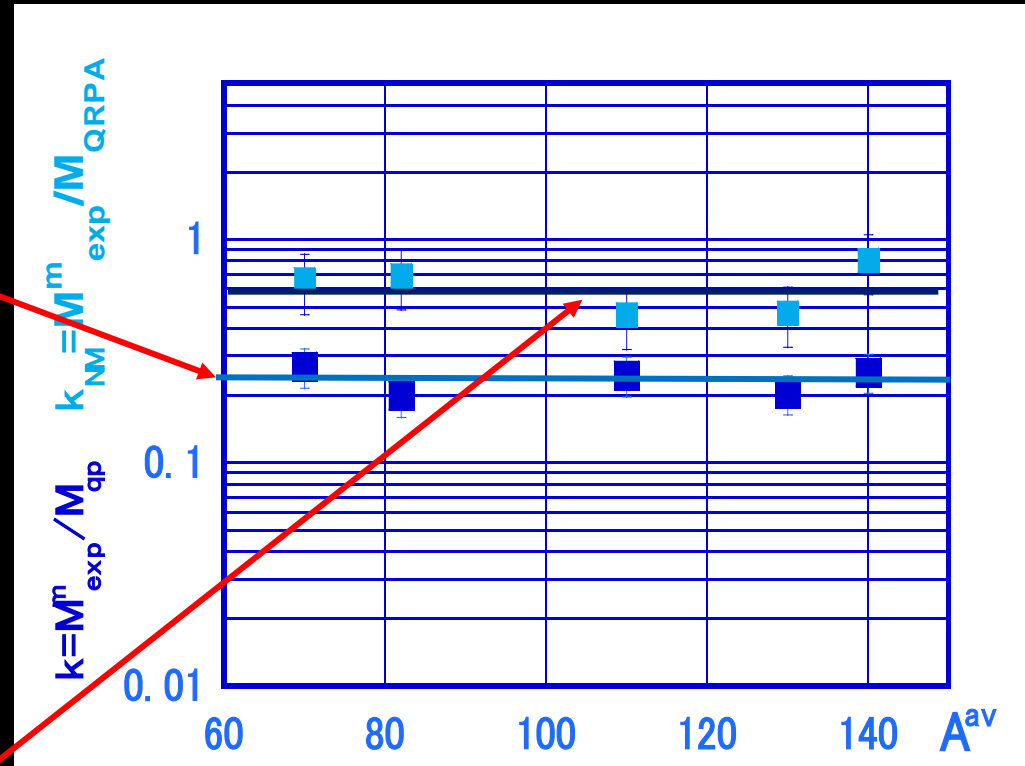
$$k = k_{\tau\sigma} k_{\text{NM}} \sim 0.23$$

$$M_{\text{QRPA}}^m = k_{\tau\sigma} M_{\text{QP}}$$

$$k_{\tau\sigma} \sim 0.4 \quad \text{NN } \tau\sigma$$

$$M_{\text{exp}}^m = k_{\text{NM}} M_{\text{QRPA}}$$

$$k_{\text{NM}} \sim 0.6 = g_A^{\text{eff}} \quad \text{N}\Delta$$



H, Ejiri J. Suhonen
 J. Phys. G. 42 2015 055201

SD 2⁻ τσ correlation & nuclear medium g_A

$$M(\text{EXP}) = k M(\text{QP})$$

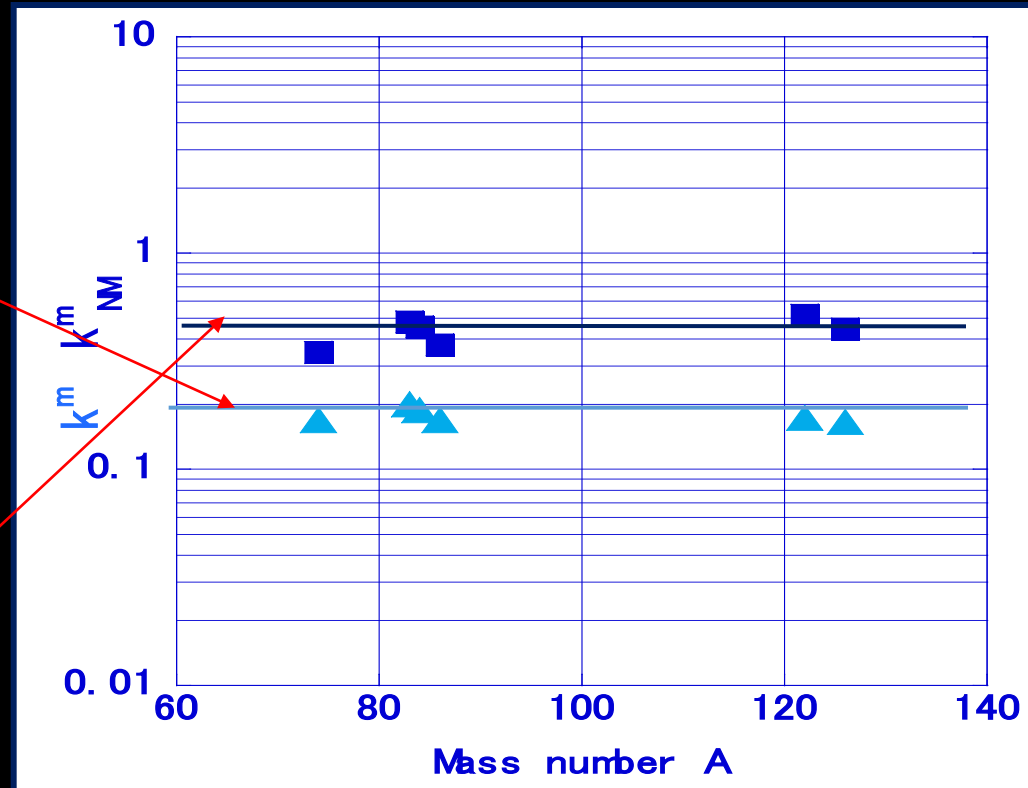
$$k = k_{\tau\sigma} k_{\text{NM}} \sim 0.2$$

$$M(\text{QRPA}) = k_{\tau\sigma} M(\text{QP})$$

$$k_{\tau\sigma} \sim 0.4 \text{ } \tau\sigma \text{ correlation}$$

$$M(\text{EXP}) = k_{\text{NM}} M(\text{QRPA})$$

$$k_{\text{NM}} \sim 0.5 = g_A^{\text{eff}} \text{ n-medium}$$

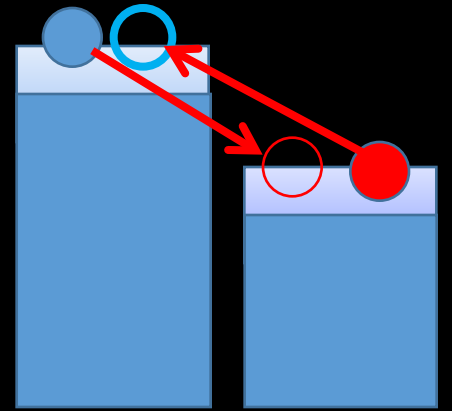


*H. Ejiri, N. Soukouti, and J. Suhonen, Phys. Lett. B 729 (2014)

FSQP : Fermi Surface Quasi Particle

Ground states $0^+ - 0^+$ QP vacuum

n p on the diffused Fermi surface

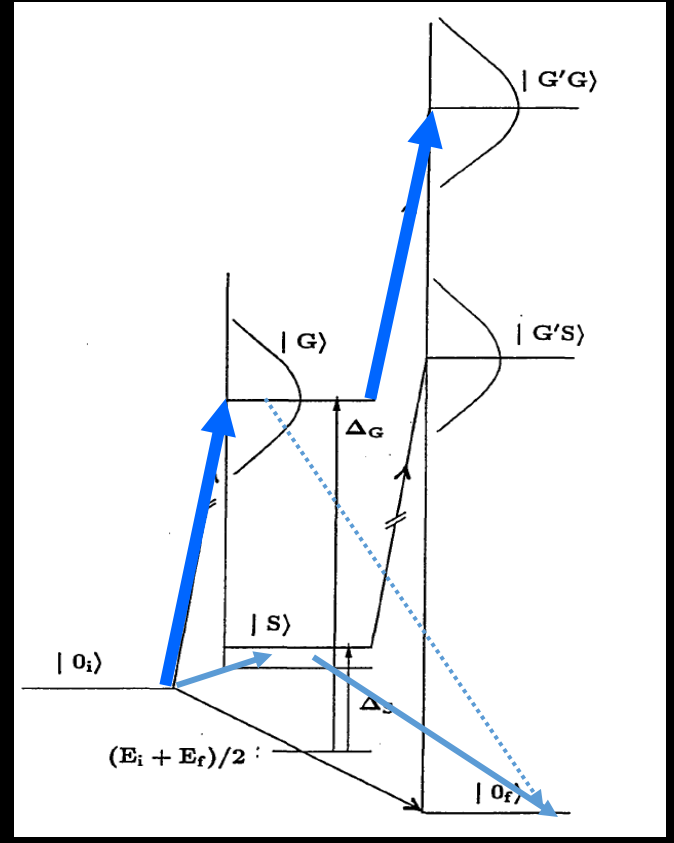


$$M(\beta\beta) = M^- M^+ / E$$

$$M^- (n \rightarrow p \quad V_n \quad U_p) \quad M^+ (p \rightarrow n \quad V_p \quad U_n)$$

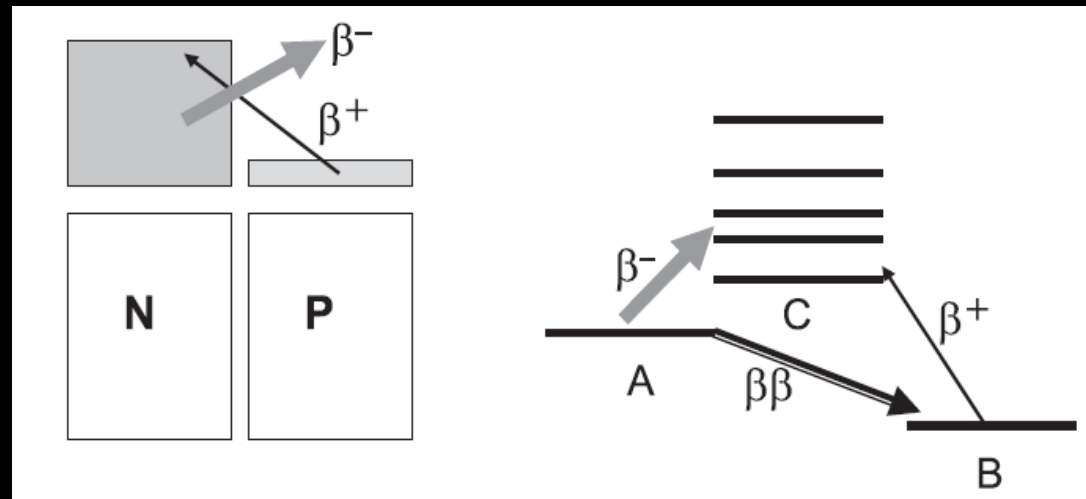
Transitions are between partially filled partially vacant n and p states. in low FSQP states within $\Delta = 2 \text{ MeV}$

GT GR \rightarrow D GT GR, and not $|0\rangle$ in $2\nu\beta\beta$.
 H. Ejiri et al. J. Phys. Soc. J. Lett. 65 (1996) 7



FSQP: Fermi Surface Quasi Particle Model

Ground state $\beta\beta$
 Fermi surface QP
 $0^+ (nn) \rightarrow 0^+ (pp)$



$$M^{2\nu\beta\beta} = \sum_{\mathbf{k}} M_{\mathbf{k}}^{-} M_{\mathbf{k}}^{+} / \Delta_{\mathbf{k}}$$

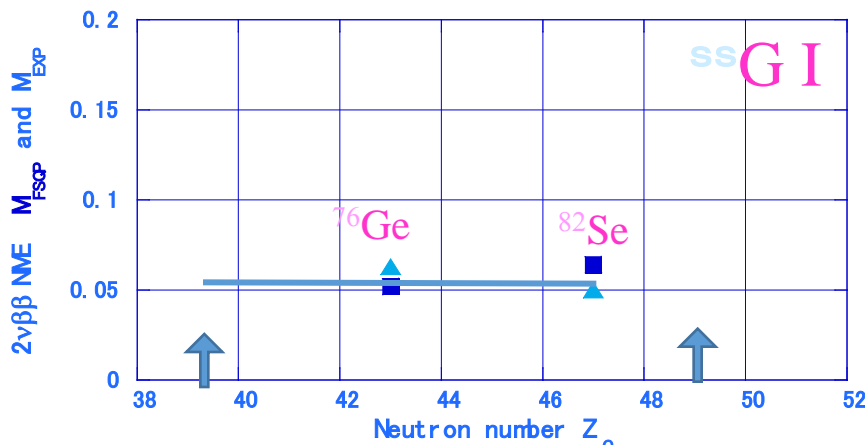
$$M_{\mathbf{k}}^{-} = (\mathbf{k}^{\text{eff}}_{\text{i}}) m_{ij} V_{\text{n}} U_{\text{p}} \quad M_{\mathbf{k}}^{+} = (\mathbf{k}^{\text{eff}}_{\text{f}}) m_{ij} U_{\text{n}} V_{\text{p}}$$

$M_{\mathbf{k}}^{-}$ and $M_{\mathbf{k}}^{+}$ are same sign
 $\mathbf{k}^{\text{eff}}_{\text{A}}$: $\tau\sigma$ & medium/isobar effects,
 derived from exp. β , EC, CER,
 $(\mathbf{k}^{\text{eff}}_{\text{A}})^2 \sim (0.23)^2 = 0.05$

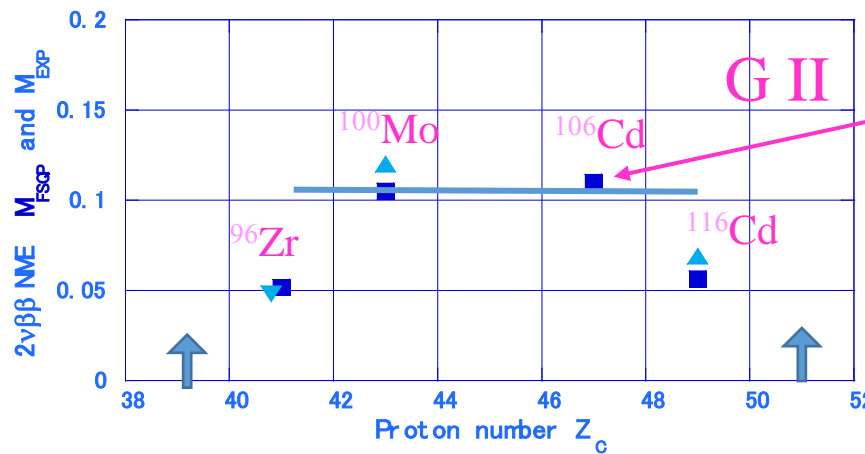
$2\nu\beta\beta$ matrix element

M(FSQP)

M(EXP)

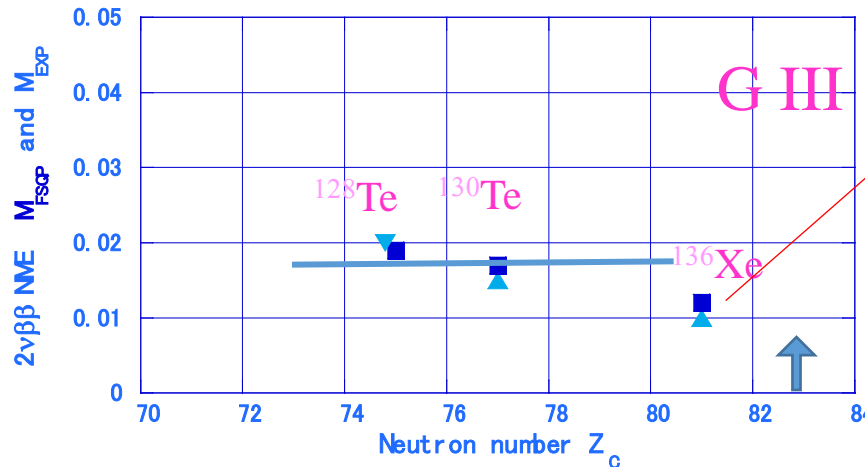


$2p_{1/2}-2p_{3/2}$



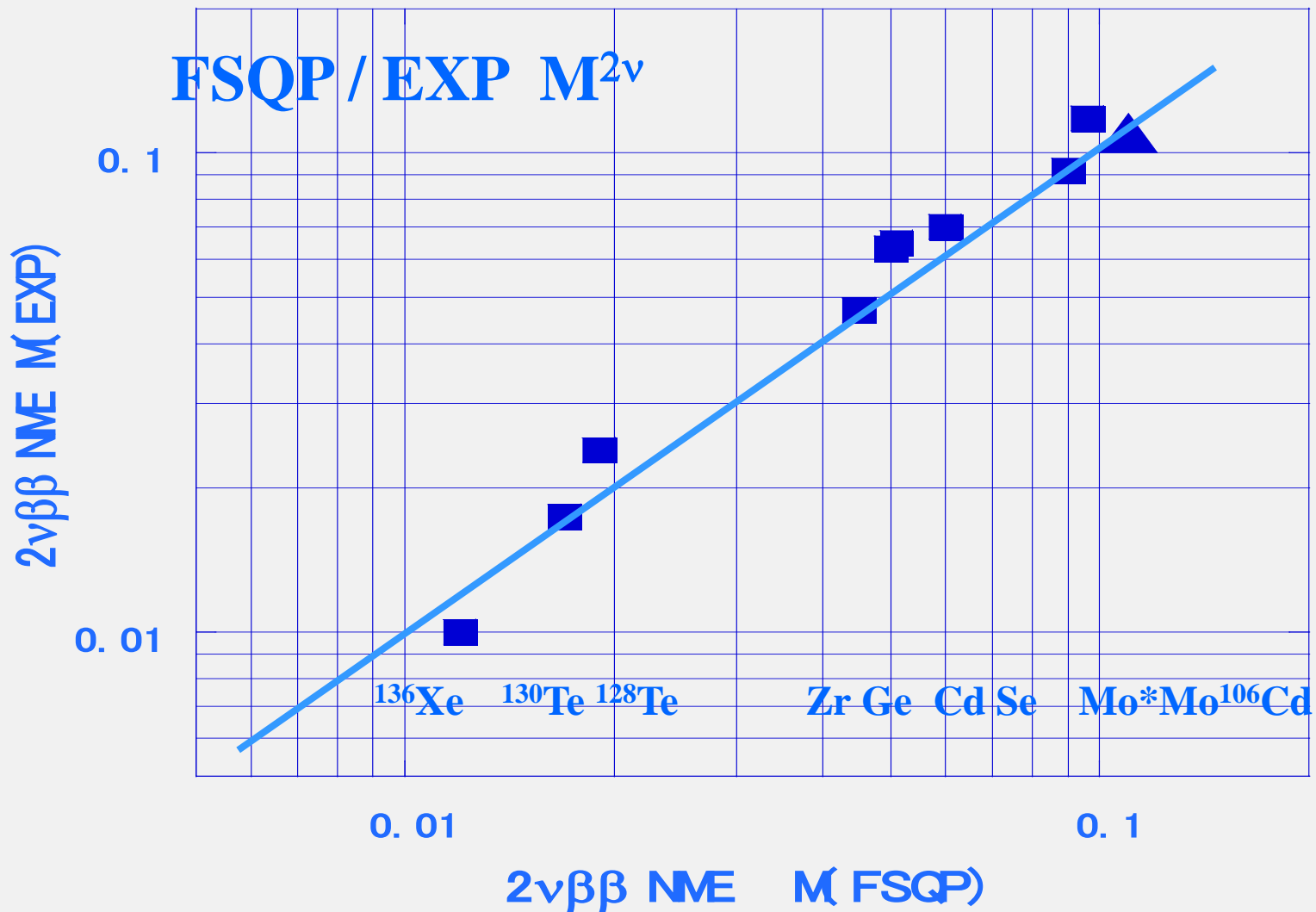
^{106}Cd predicted
 $T_{1/2}(\text{ECEC}) = 5.2 \cdot 10^{22}$

$1g_{7/2}-1g_{9/2}$



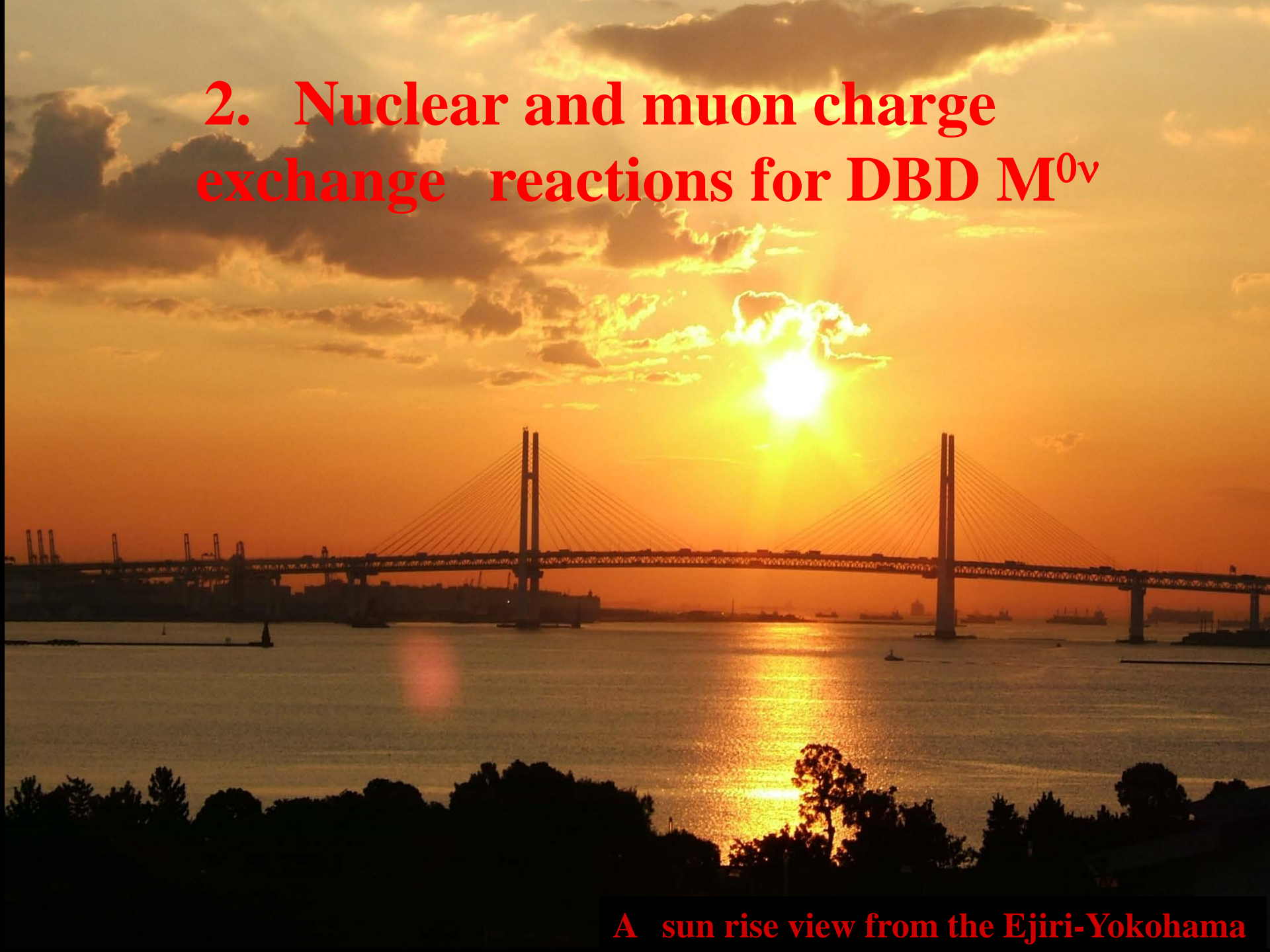
$N=82, p \rightarrow n$

$2d_{3/2}-2d_{5/2}$



$M^{2\nu}$: small by $(k^{\text{eff}})^2 \sim 0.05$, depend on VU, $E(1^+)$, not by g_{pp}
 $M^{0\nu}$, likewise, small by k^{eff} for 2^- etc, depend on V,U (N=82)

2. Nuclear and muon charge exchange reactions for DBD $M^{0\nu}$

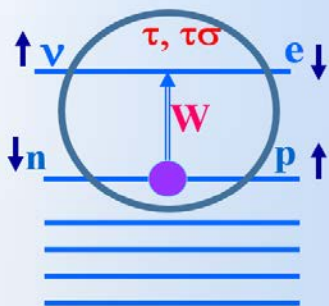


A sun rise view from the Ejiri-Yokohama

CER. Probs for ν -responses

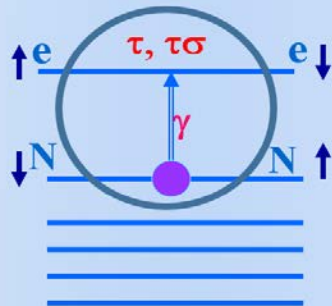
Nuclear $\tau\sigma$ responses for ν in β & $\beta\beta$

Weak probe



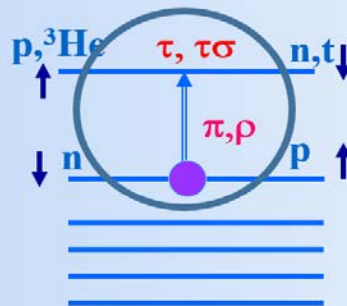
β -decay,
e capture
 ν, μ probe
J-PARC

EM probe



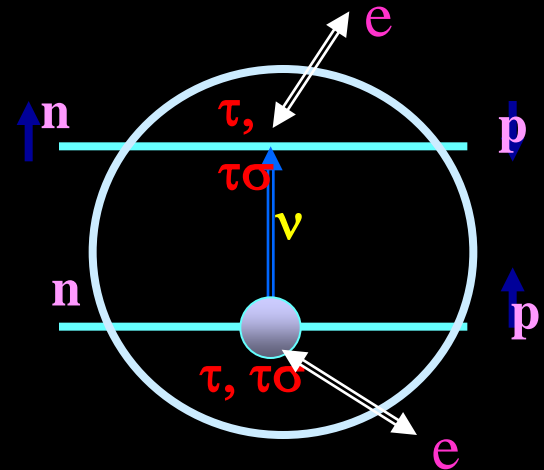
γ -capture,
e scattering
 γ from
Spring-8 HIGS

Nuclear probe

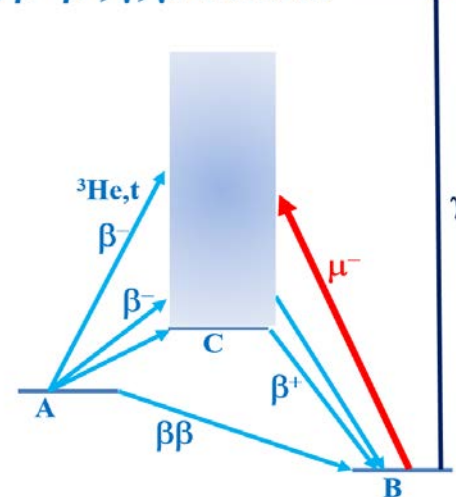


CER ${}^3\text{He}, t$
 $t, {}^3\text{He}, d, {}^2\text{He}$
N RCNP,
MSU, KVI

H. Ejiri, Prog. particle Nuclear Physics, 64 '10 249



$\beta\beta, \beta^+ \beta^-, \gamma, \mu$ scheme



$0\nu\beta\beta$: two body operator, but single β NMEs are useful for $0\nu\beta\beta$
H. Ejiri PR 338 265 2000.

Neutrino response studies by RCNP/Osaka

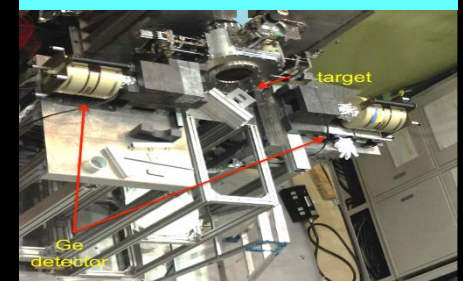
RCNP Osaka p, He, HI, μ



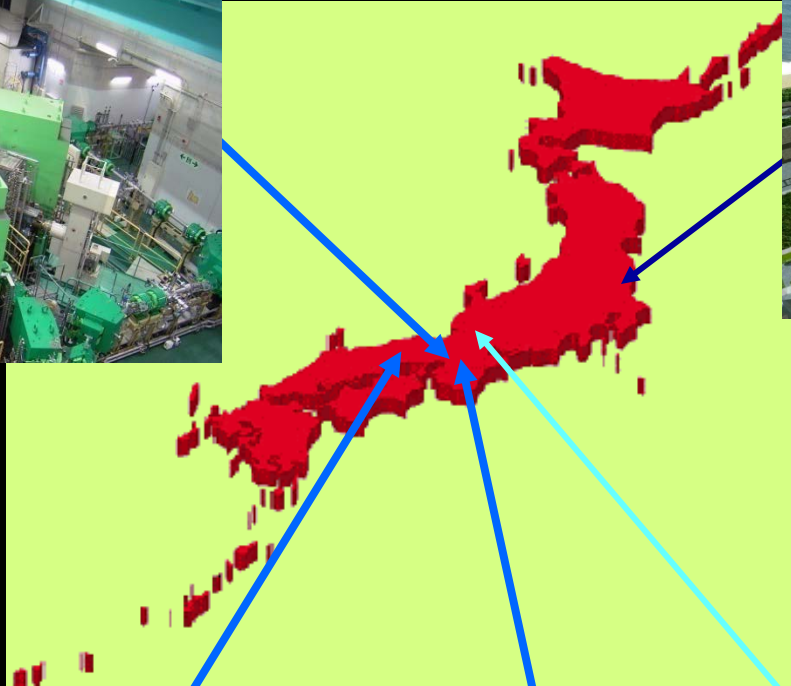
J-PARC 3-50 GeV p, ν , μ



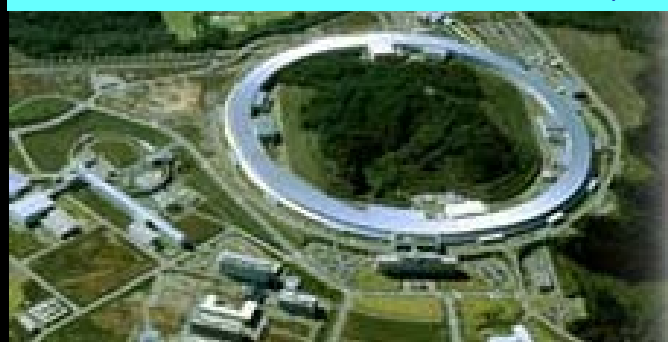
MLF D2 μ



MuSIC μ



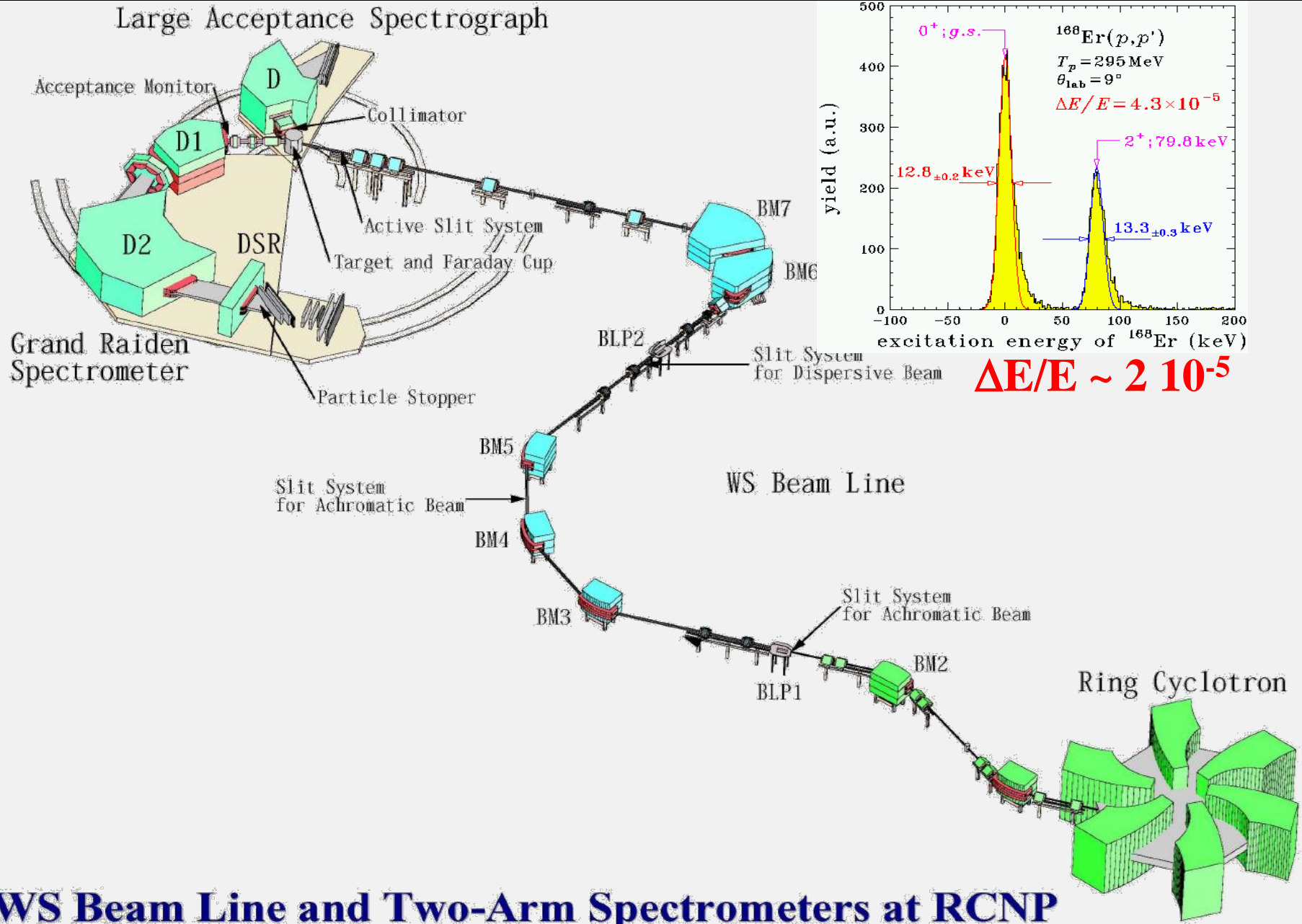
Spring-8 GeV- MeV pol. γ



Oto underground lab.
 $\beta\beta-\nu$, DM in nuclei

SK/KamLAND
Underground lab.
 ν -osc. T2K. ν / SN,
the sun and earth,

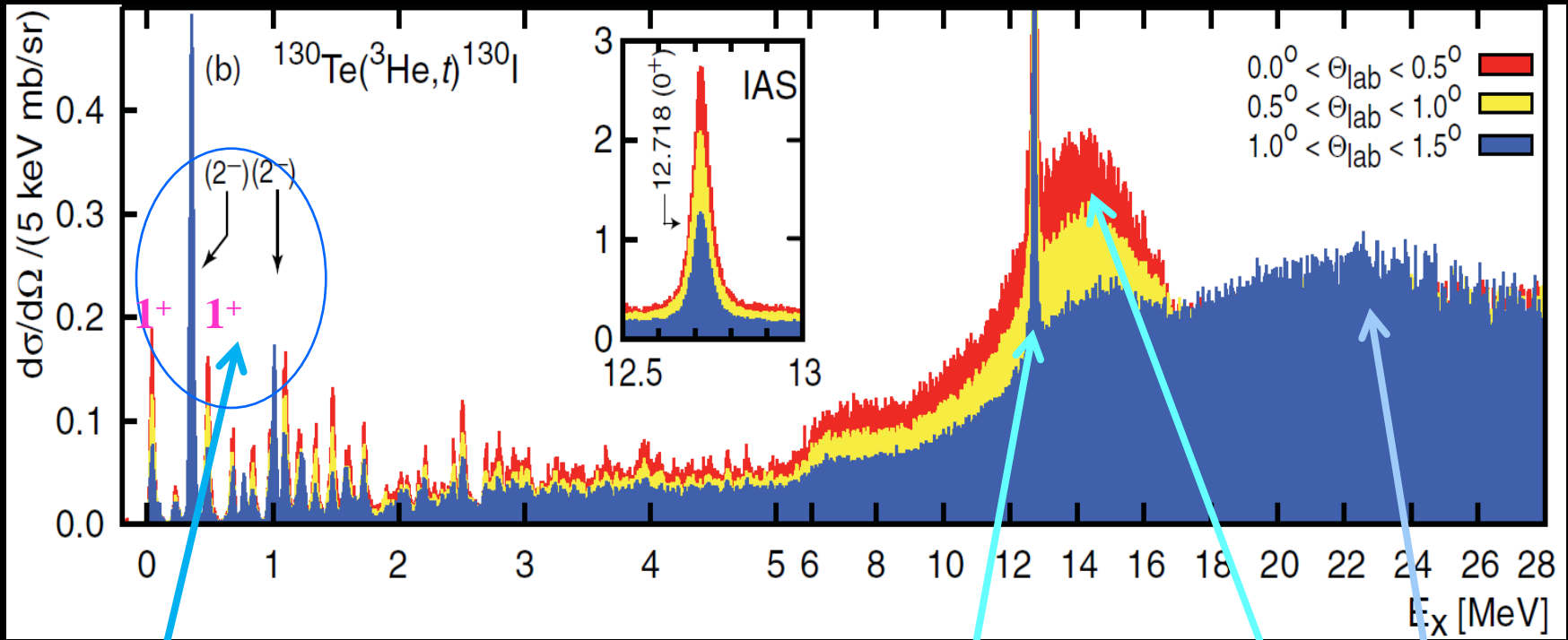
High E resolution ($^3\text{He},t$) CERs at RCNP



WS Beam Line and Two-Arm Spectrometers at RCNP

Single β ($\tau\sigma$) strengths on DBD Ge,Se, Zr,Mo,(Cd),Te,Xe

Nucl CERs Konan, KVI, MSU, Munster, RCNP/Osaka, and others. Faessler et al



Tiny GT(1⁺), 2⁻ 3⁺ at low E.
B(GT) ~ 10⁻³ B(SUM)

Strong IAS (0⁺), GT(1⁺), SD(2⁻)
GR : B(GT) ~ 0.5 B(SUM)

... H.E PR C 84 2011

P.Puppe et al PRC86 044603

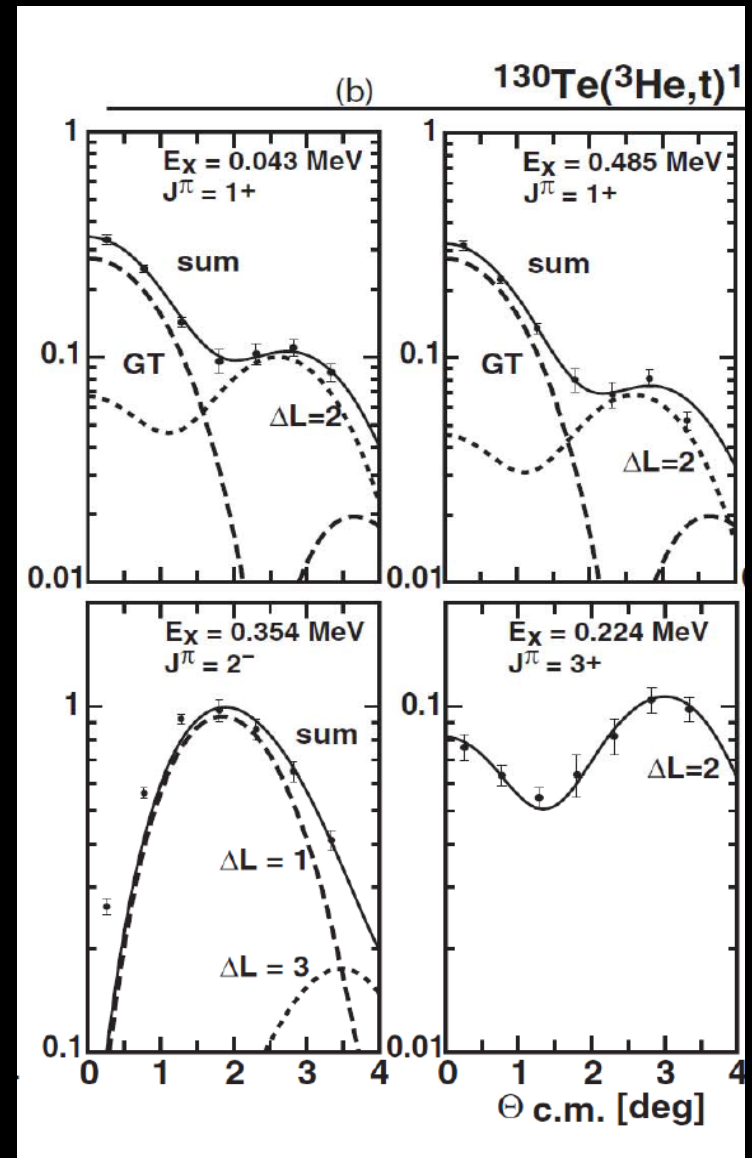
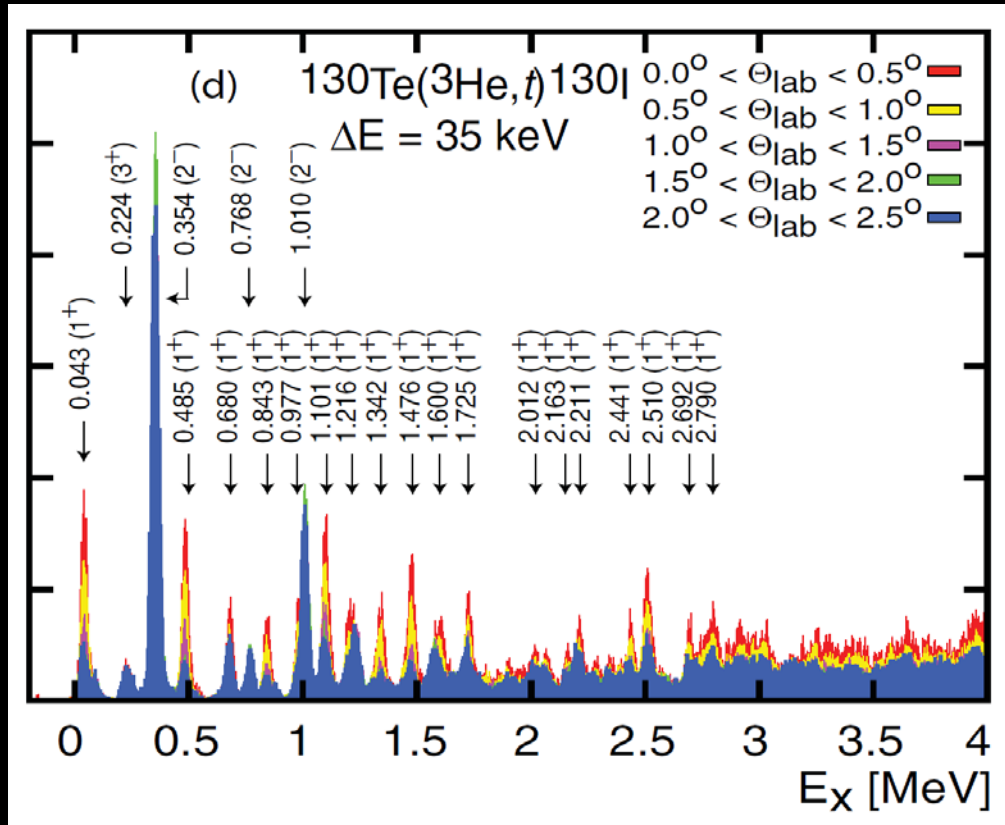
H.Ejiri et al. PR 176 (1968) 1277

D. Frekers in Dec.2nd.

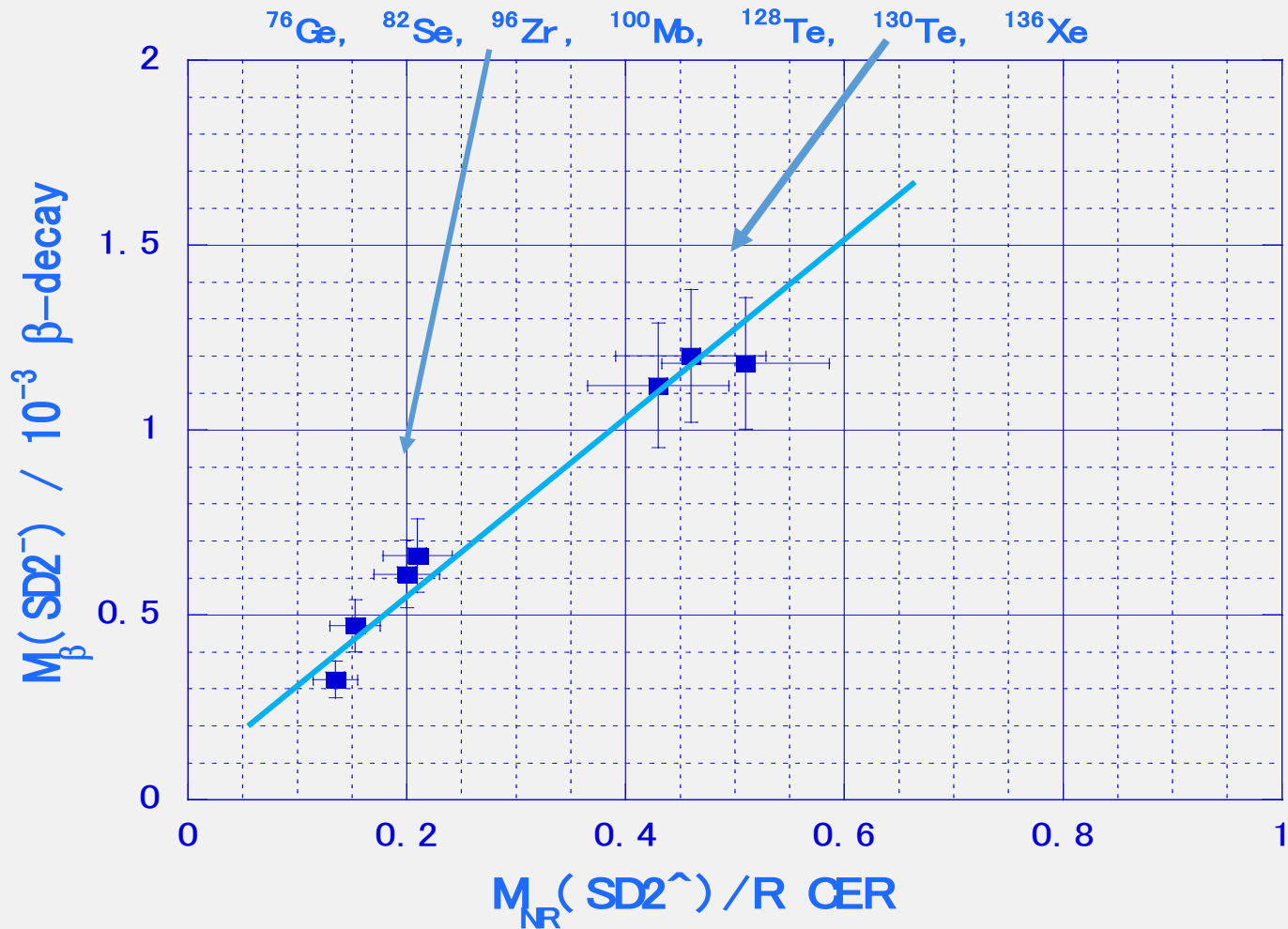
First 2⁻ GR based on reduced B(SD)

$$\frac{d\sigma_{\alpha}(0^{\circ})}{d\Omega} \frac{1}{K(E_i, 0)N_{\alpha}^D} = |J_{\alpha}|^2 B(\alpha), \quad \alpha = F, GT.$$

$$B(\alpha) = (2J+1)^{-1} M(\tau\sigma Y_1)^2$$



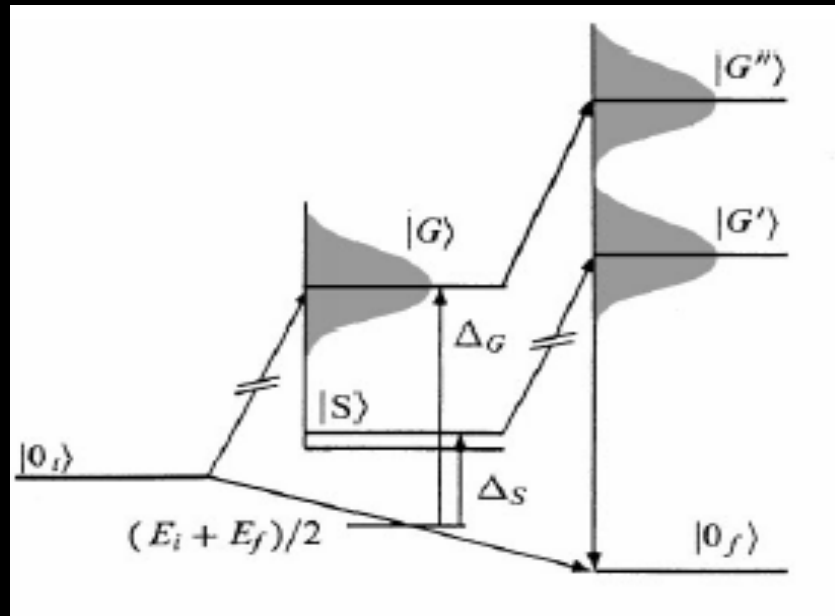
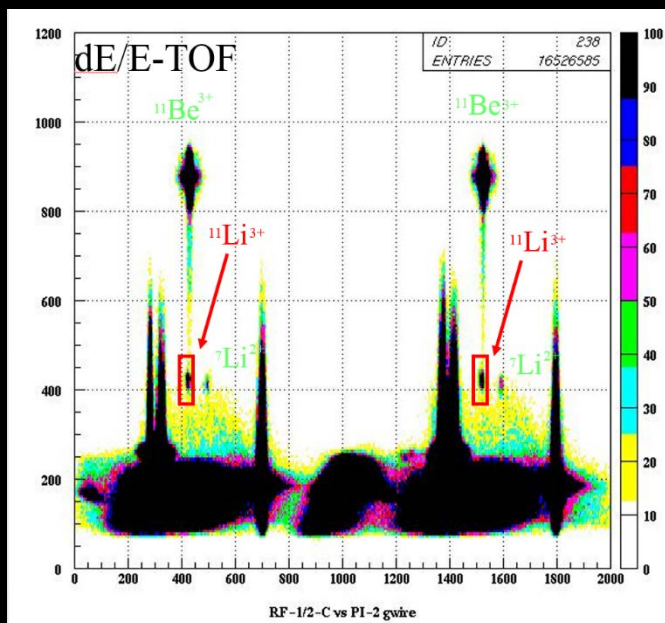
Angular distributions give $J=1^+$ ($l=0$), or $J=2^-$ ($l=1$), or $J=3^+$ ($l=2$).



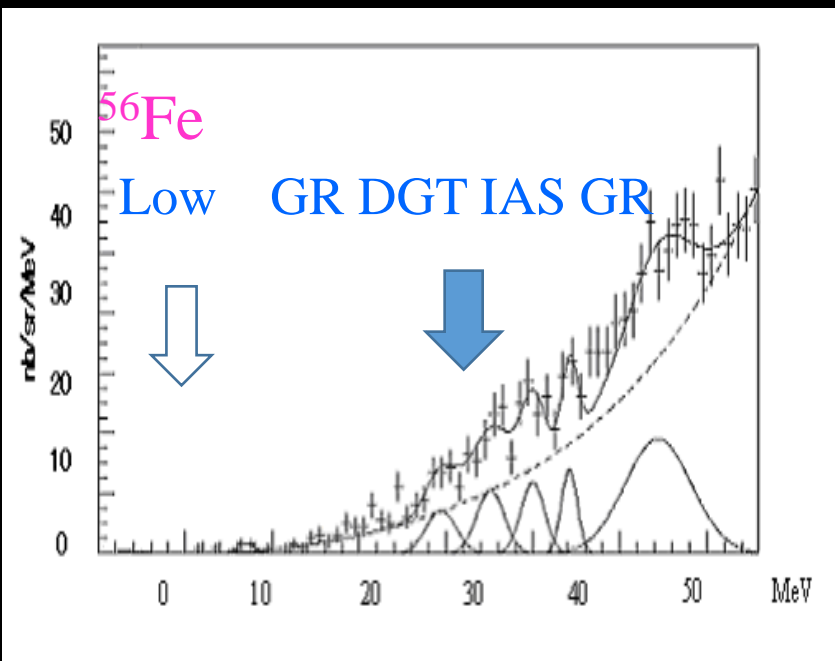
SD NMEs (Y axis) with g_A from ft data in neighboring nuclei by H.E Exps planned on SD with Agodi, Akimune, Capuzzello, Frekers et al.

Double charge exchange reaction *

RCNP 0.76 GeV ^{11}B , ^{11}Li



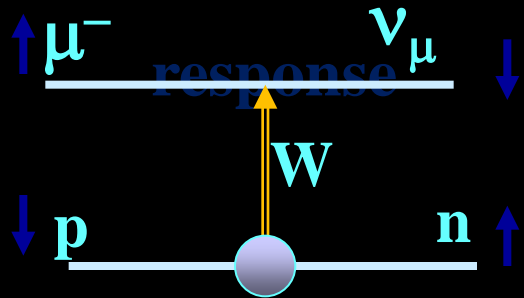
^{13}C strengths at low high states
 ^{56}Fe no low states, mostly GRs



Dec.2nd Takai HI
 NUMEN !!

Takahisa Ejiri et al 2010

Muon charge exchange reactions

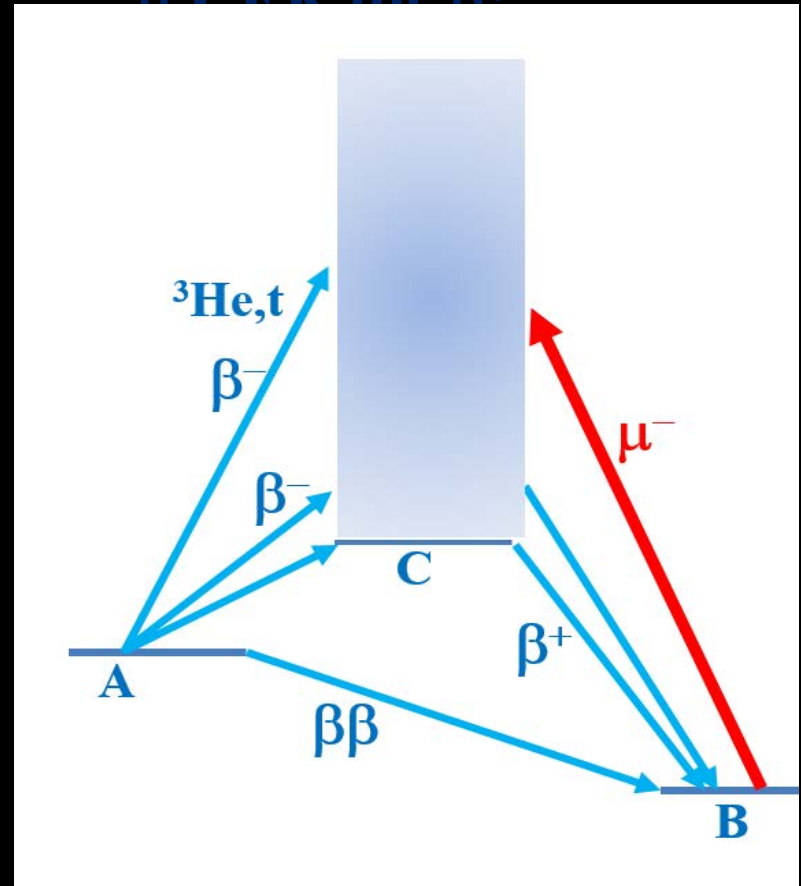


$E \sim NME$ in 0-100 MeV

β^+ anti- ν response

for $\beta\beta-\nu$ SN ν

Large μ capture rate ~ 1



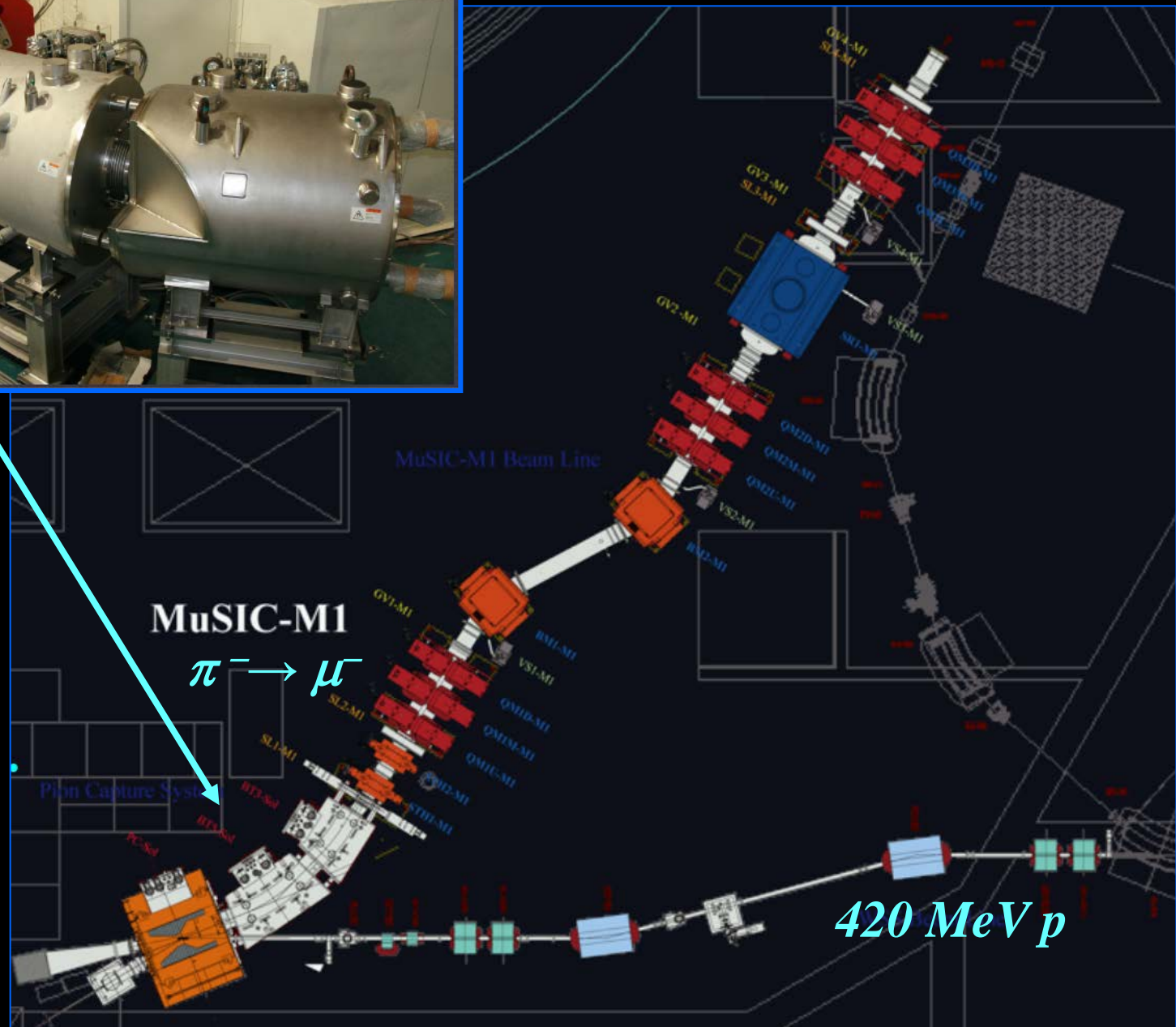
H. Ejiri Proc. EM interactions 1972. AIP Conf. MEDEX11 2011.

J. Suhonen, M. Kortelainen Czech. J. Phy. 56 2006 519, EXP 453.

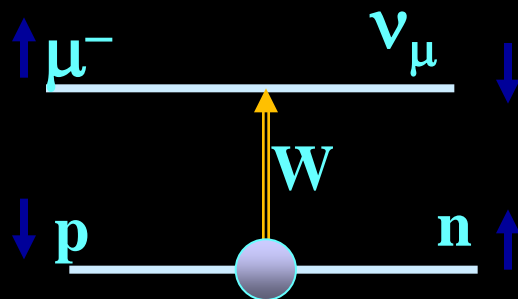
V. Egorov ($\mu, \gamma, n \gamma, p \gamma$) on ${}^{48}\text{Ti}, {}^{76}\text{Se}$ 2004 γ D.R. Zinatalina et al. .

MUSIC RCNP Osaka

DC μ beam

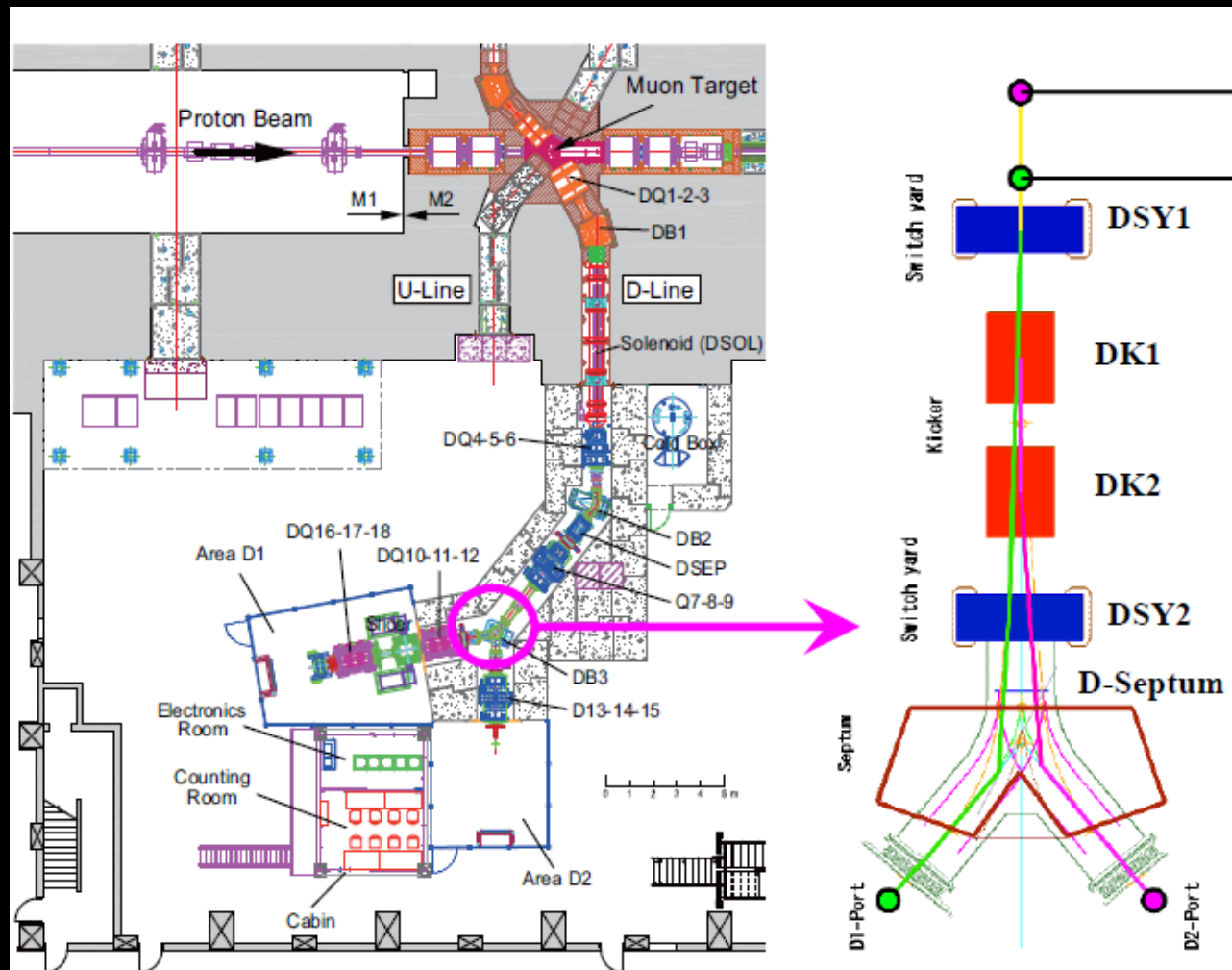


MLF μ probe



MLF (pulsed)

Enriched
 $^{100}\text{Mo}(\mu, \nu_\mu \text{ xn } \beta\gamma)$

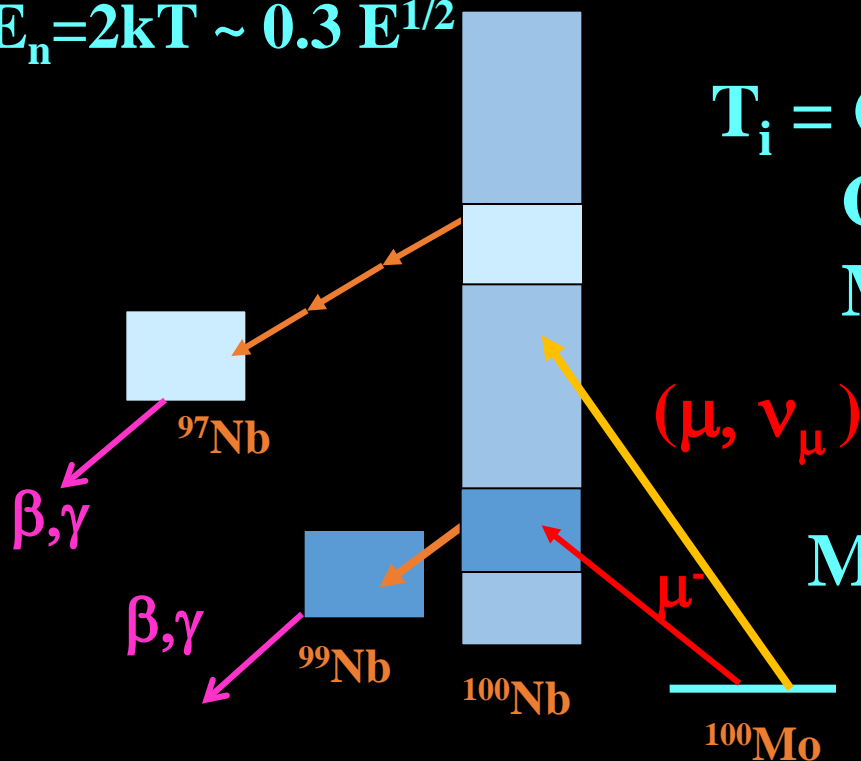


Present CER $^{100}\text{Mo}(\mu, \nu_\mu, xn \gamma)^{100-x}\text{Nb} \nu-\tau(\beta)+$ responses

$F(E_n) = E \exp(-E_n/kT)$ $T = \sum T_i$ transition at i^{th} excited state

$$E_n = 2kT \sim 0.3 E^{1/2}$$

$T_i = G_i M_i^2$ with
 $G_i = k(Q - E_0)^2$ phase space
 M_i : NME



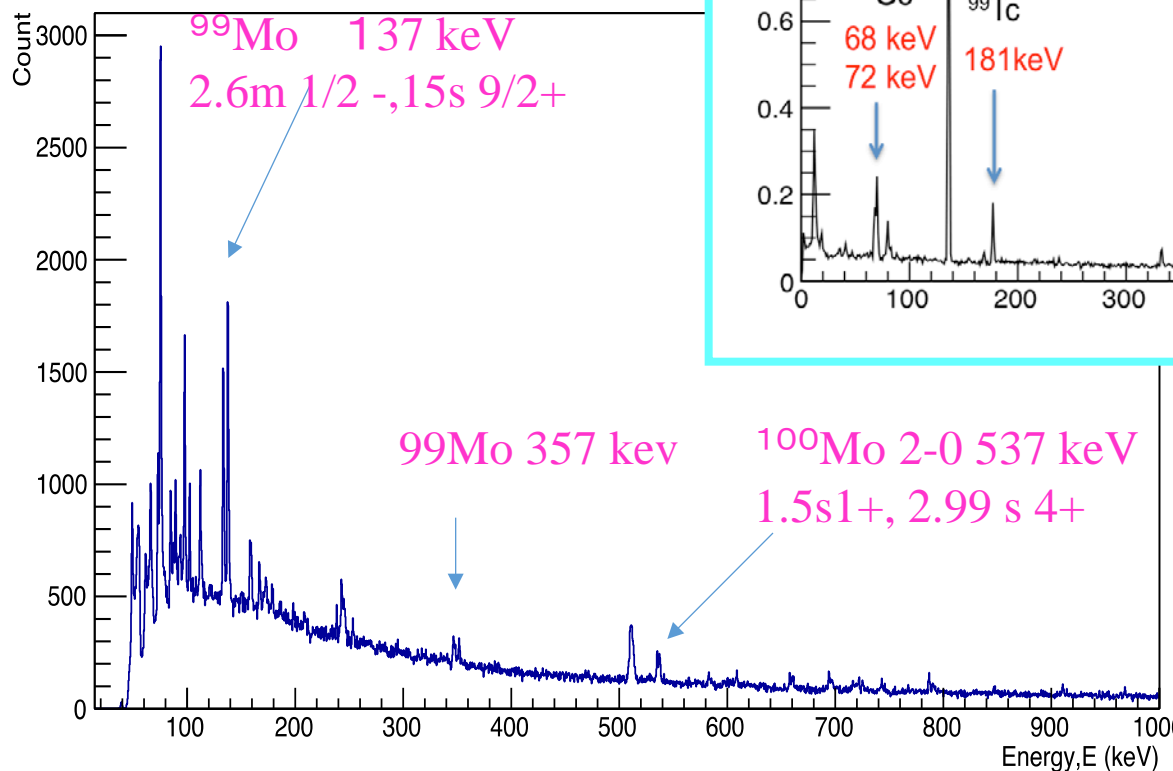
$M_i = g_A M_A + g_V M_V$ for $i = 0, 1, 2, \dots$
 Effective g_A for $i = 0, 1, 2, \dots$

γ_i from $^{100-i}\text{Nb}$ gives the isotope i distribution, relative strength in the whole excitation region. Life time gives the absolute strength

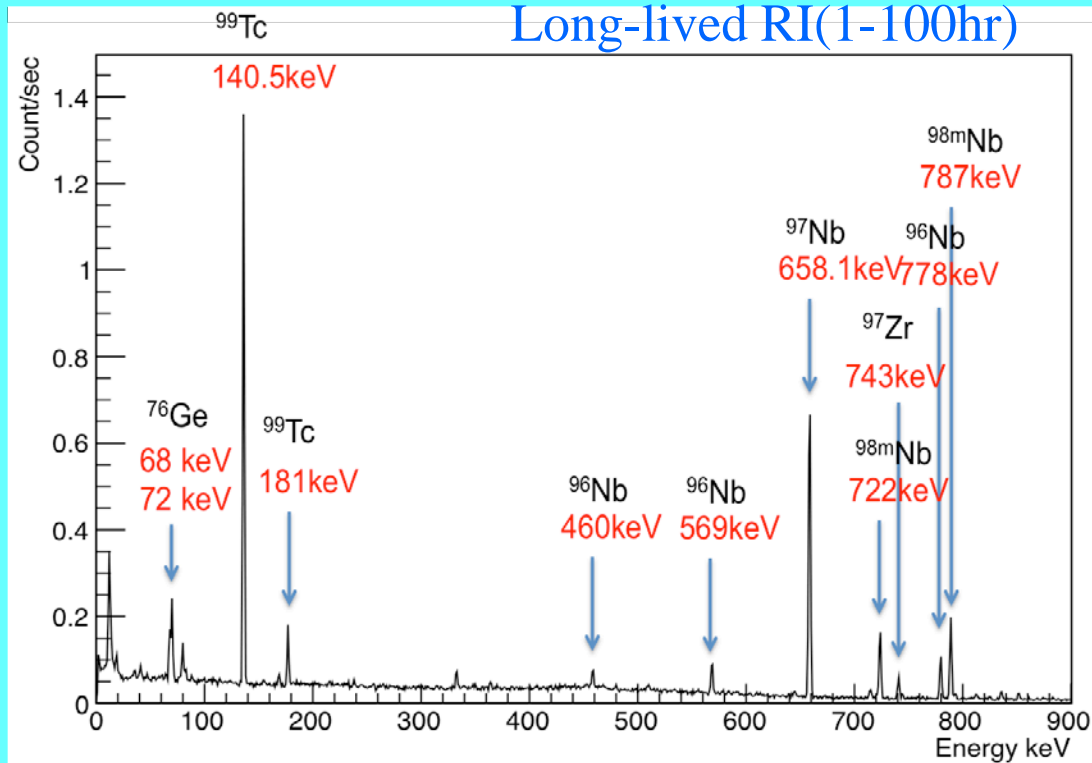
MLF (pulsed)

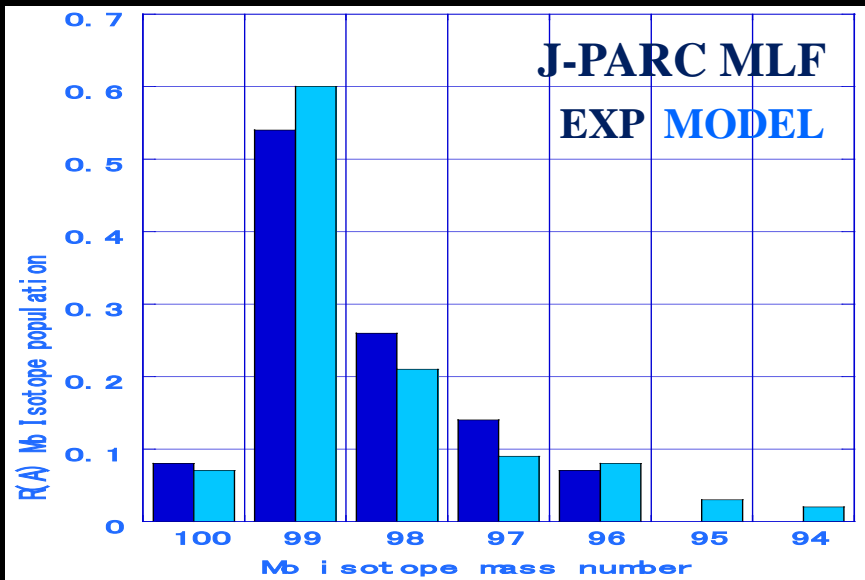
$^{100}\text{Mo}(\mu, \nu_{\mu} \text{ xn } \beta\gamma)$

Short-lived RI(1-200s)



Long-lived RI(1-100hr)

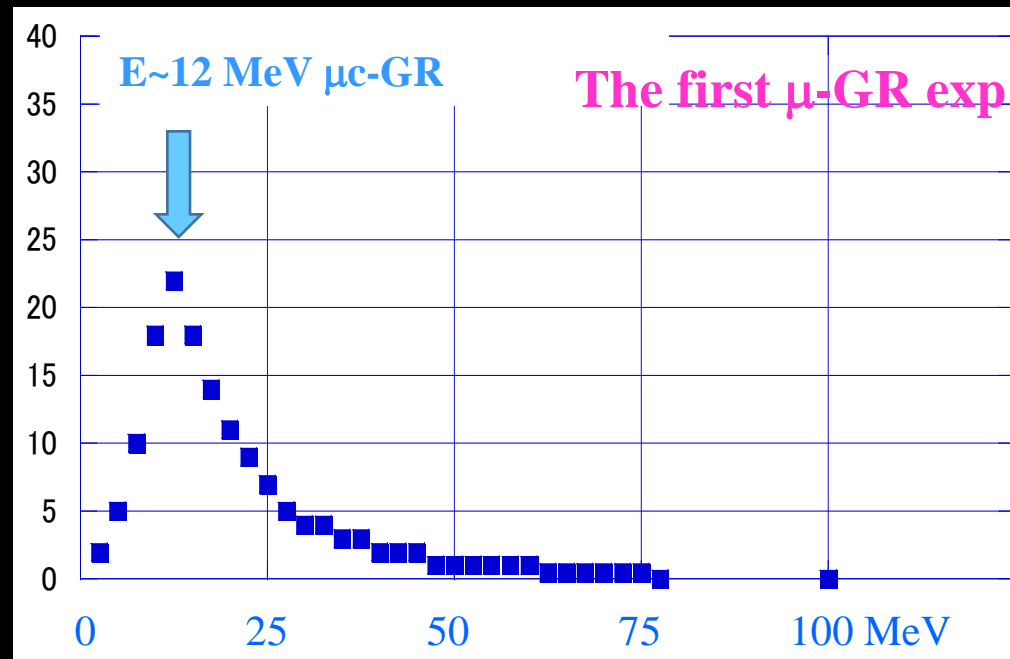




Observed isotope population agrees with calculation with μ -GR as given below.

H. Ejiri et al. JPSJ 84 044202 2013

I. Hashim PhD Thesis 2015

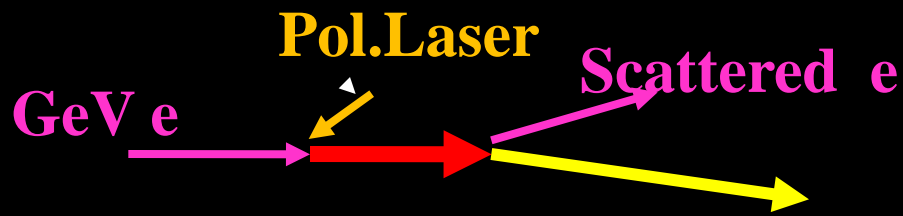


LEPS Photon probe

H. Ejiri PRL 21 '68, PR 38 '78

H. Ejiri, A. Titov PR C 88 054610 2013

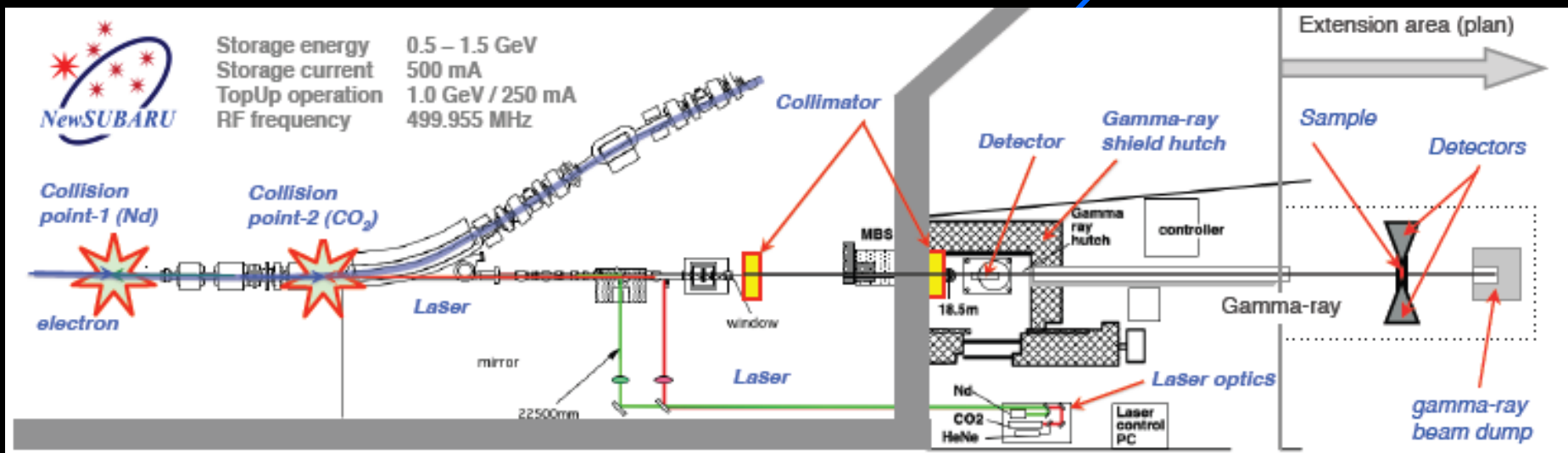
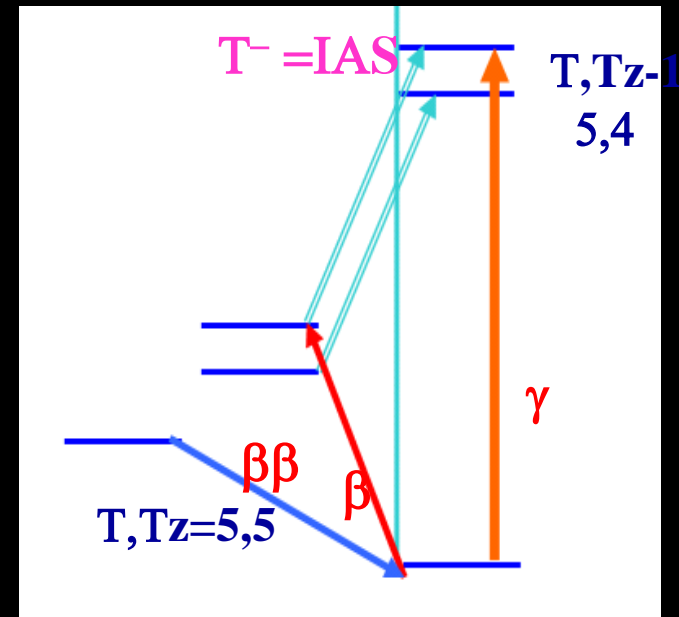
Laser electron photon Sources



$$E_\gamma \sim 4 \gamma^2 E_1 = 2 \cdot 10^8 E_1 \sim 20 \text{ MeV,}$$

β^+ responses via IAS γ

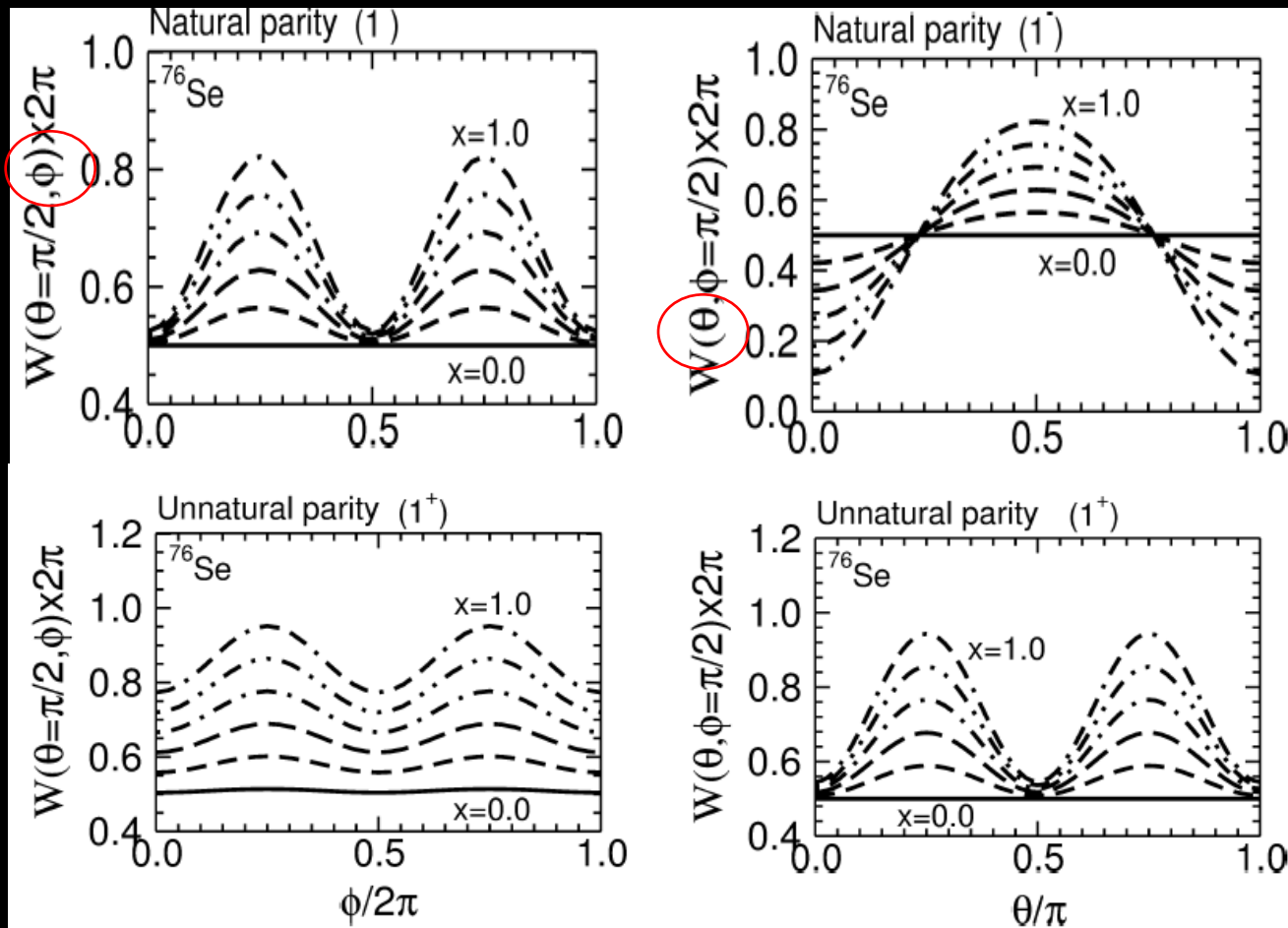
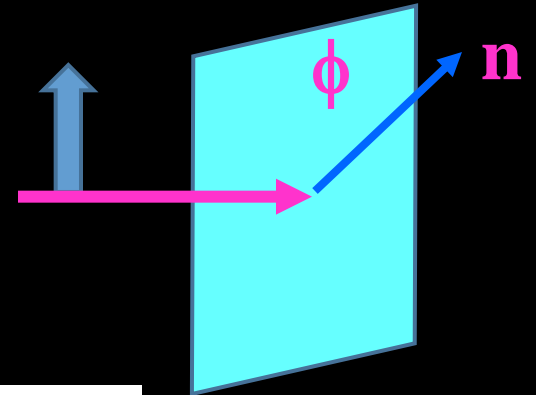
$$\langle f | g M_\beta | i \rangle = g/e (2T)^{1/2} \langle f | e m_\gamma | I \rangle$$



Polarized photon

$$\int \sigma(\gamma, n) dE = 2.9 \times 10^{-3} \text{ MeV fm}^2 \quad E1 \ 0^+ \rightarrow 1^-,$$

$$\int \sigma(\gamma, n) dE = 2.7 \times 10^{-3} \text{ MeV fm}^2 \quad M1 \ 0^+ \rightarrow 1^+.$$



$\gamma + ^{76}\text{Se} = \text{IAS } J^\pi$
 $n + ^{75}\text{Se } 3/2^-$
 $x = d/s, f/p$
fraction

$V \tau \ 1^- \ 2^+$
 $AV \tau \sigma \ 1^+ \ 2^-$

H. Ejiri, A. Titov
 PR C 88 054610
 2013

Hadrons (Δ , π) *

$N \Rightarrow \Delta$ quark $\tau\sigma$ flip GR
 0.3 GeV high excitation
 $0\nu\beta\beta$ in Δ, π enhanced

Effect on low $\beta\beta$ $0^+ - 0^+$

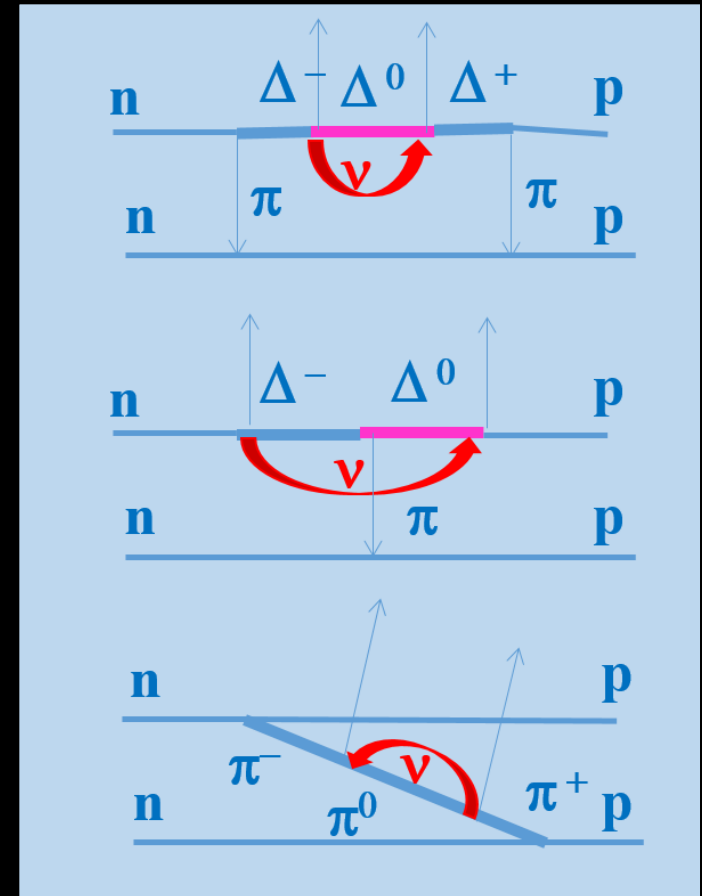
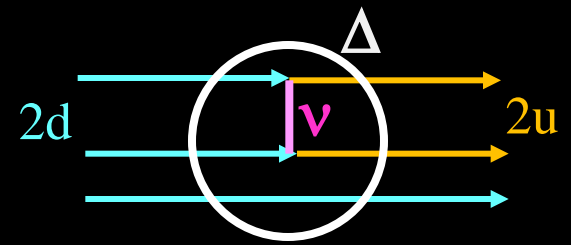
$$P(\Delta)^2 \sim (10^{-2})^2 \sim 10^{-4} \sim (k=0.2)^4$$

$|n \Delta\rangle$ interferes with $|n p\rangle$

$$|I\rangle = |n p\rangle - \varepsilon |n, \Delta^+\rangle$$

$$M^\beta \sim k^{\text{eff}} M_0 \quad k^{\text{eff}}(\Delta) \sim 0.6$$

Exp by β deays and CER



*Pontecorvo; Haxton, Stephenson, Kotani Doi .

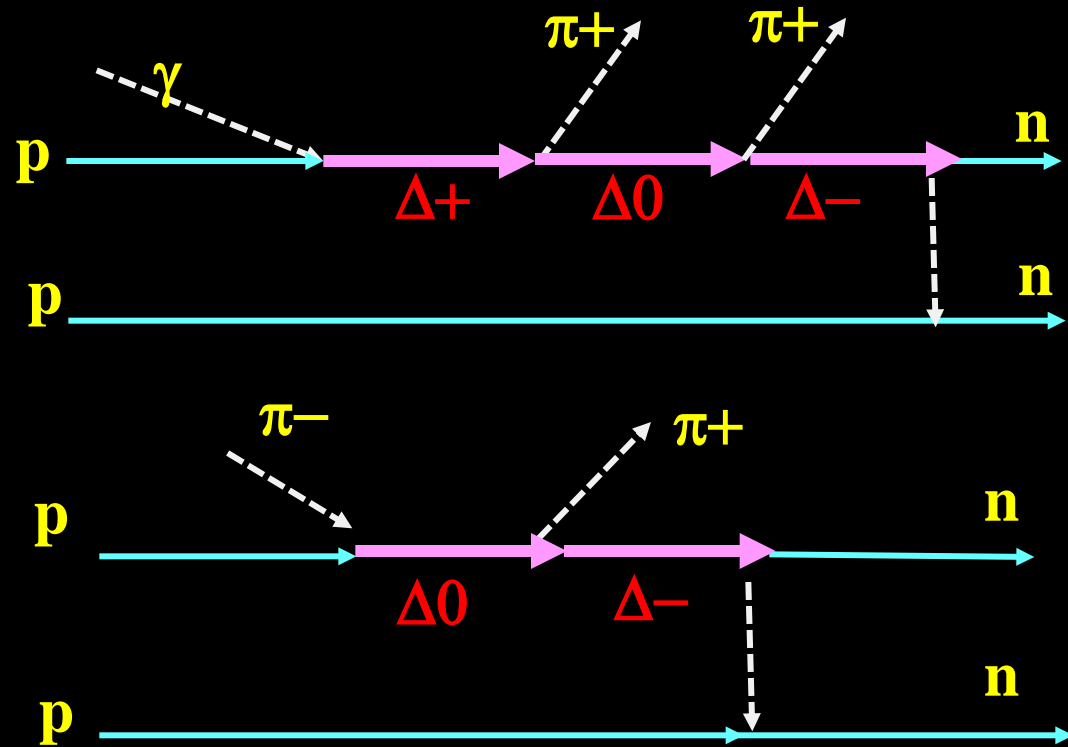
DCE of hadrons

$$\gamma + 2p = 2\pi^+ + 2n$$

Δ^+ Δ^- DCE

$$\pi^- + 2n = \pi^+ + 2p$$

π^+ π^- DCE



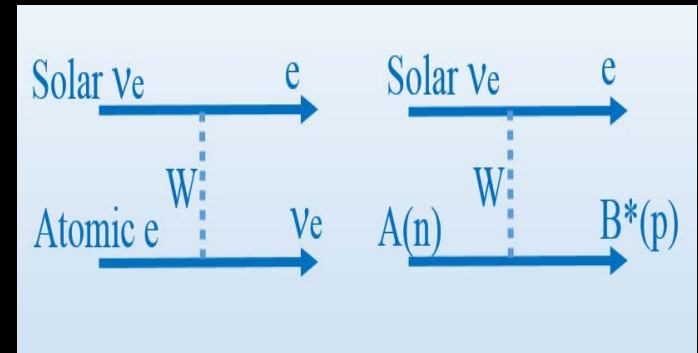
Hodron contributions



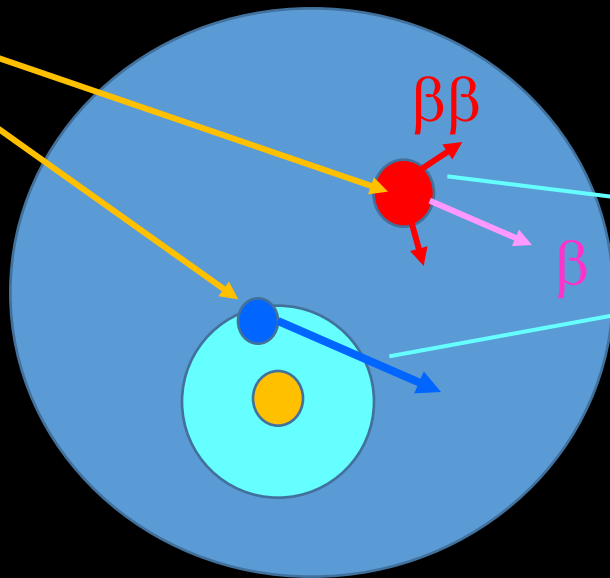
3. Solar- ν interactions with DBD detectors and BG contributions

Solar- ν interactions with nuclei and atomic electrons in DBD detectors are serious BGs

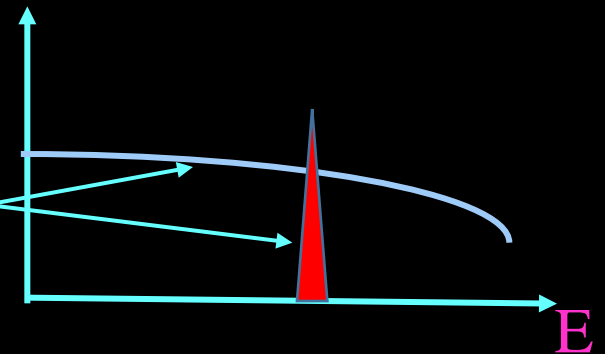
- Solar ν unavoidable.
- BG rate \sim or $>$ or $<$ $\beta\beta$ signal rate
- ν cross section on DBD nucleus,
- E-resolution, ratio of DBD nuclei to scintillation detector atoms



Solar ν



νe scattering.



BG rate and solar- ν interaction by CER

Mass sensitivities $\langle m_\nu \rangle = (2 m^0 \sim 40 \text{ meV}) (\text{NT})^{-1/4} (\text{B})^{1/4}$

with $M=2$, $\epsilon=0.6$ $B=\text{BG}/t \text{ y}$ $\text{NT}=t \text{ y}$

$\text{BG}/t \text{ y} \leq 1$

$\text{NT}/t \text{ y} \sim 4 \text{ y t}$

to get IH 30 meV

$\leq 0.01/t \text{ y}$

$\text{NT} \sim 100 \text{ ty}$

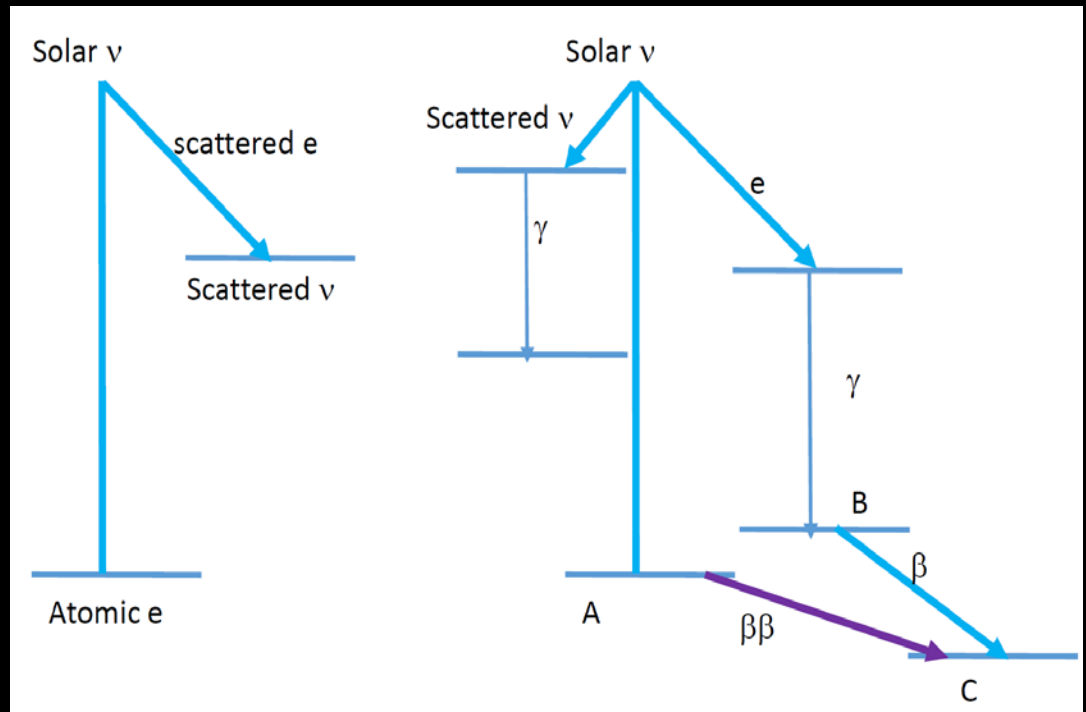
to get NH 3 meV

Solar ν interactions

i. DBD nuclei

ii. Other isotopes
if enrichment
 $\ll 100\%$

iii. Liquid scintillator
atoms and nuclei
if concentration $\ll 1$



Solar- ν CC int.with DBD nuclei

Isotope	pp	pep	^7Be	^8B	^{13}N	^{15}O	Total
^{76}Ge	0	0.7	0	5.0	0.06	0.4	6.2
^{100}Mo	390	8.2	126	6.0	6.7	8.3	545
^{130}Te	0	3.0	23.2	6.1	1.3	2.4	35.9
^{136}Xe	0	6.0	50.9	9.8	2.3	4.7	73.4

H. Ejiri et al *
PR L 85 2000 2917

H. Ejiri S. Elliott
PR C 89 2014

Solar ν GT responses from the GT responses by CER ($^3\text{He}, t$)

1. Large pp- ν to 1^+ ground state for pp- ν study * Ejiri Zuber

2. ^{130}Te , $B=1.3/\text{ty}$ at $\epsilon'=5\%$ (SNO+)

SSTC 12.4 hr and/or E-resolution of 0.2%

$B=0.05$ with $\epsilon'=0.2\%$ (CUORE) $\ll B(\text{RI})=10^3/\text{t y}$ at 2014

3. ^{136}Xe , $B=1.5/\text{ty}$ at $\epsilon'=5\%$ Large $Q_{\beta}\sim Q_{\beta\beta}$ SSTC 13.2d. Tracking

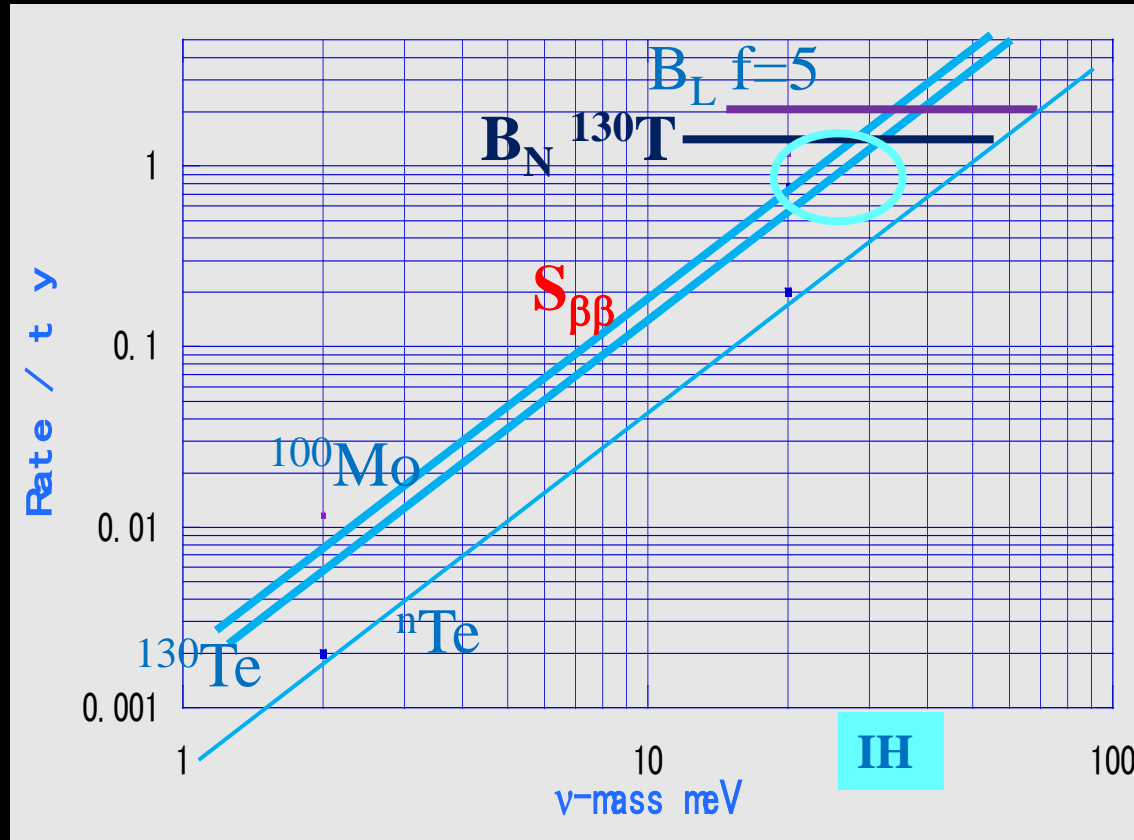
$\ll B\sim 100/\text{t y}$ at 2012 EXO

Solar- ν BG rates for liquid scintillator and $\beta\beta$ signal rate

Solar- ν $B_L=0.4 f$ with $f=\epsilon'(E\text{-resolution}) / R(\text{concentration})$

$$S_{\beta\beta} = 0.7 [M/2]^2 [m_n/20 \text{ meV}]^2$$

H. Ejiri
K. Zuber
2015



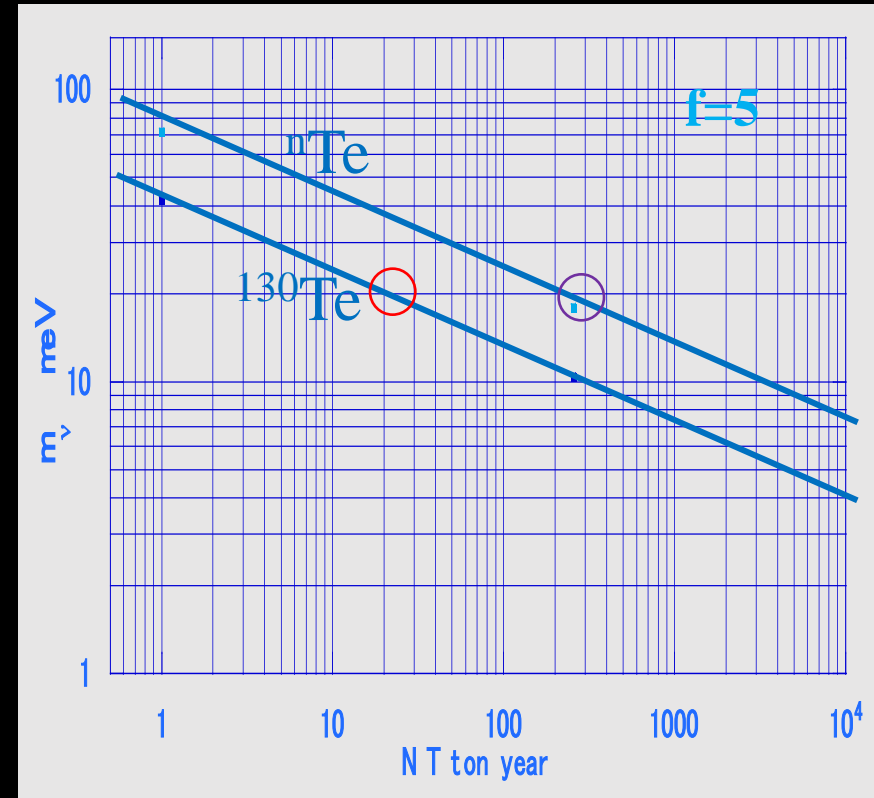
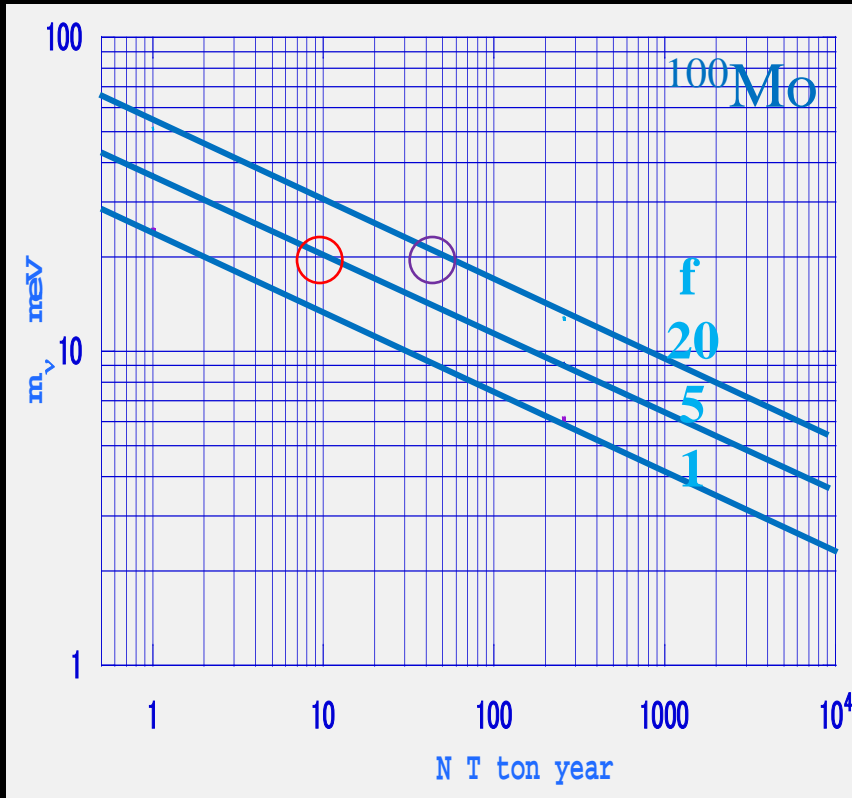
^{130}Te $B_L=2.2$ /ty, ^{136}Xe $B_L=3.0$ /ty in case of $f=5$, $\epsilon'=5\%$ $R=1\%$

Solar- ν in liquid scinti. with DBD isotope*

- H. Ejiri,
- K. Zuber 2015

$B = \text{Solar-}\nu \text{ e scat. in liq. scinti.}$ $m = m^0 k [f^{1/4} / (NT)^{1/4}]$

$f = \varepsilon' / R$, $\varepsilon' = \text{Energy resolution}$ $R = N(\beta\beta) / (N(\text{Liq.}) \text{ Concentration})$



$M=2, f=5$ ($\varepsilon'=5\%$, $R=1\%$) 10 ty, $f=20$ ($\varepsilon'=10\%$, $R=0.5\%$) $NT=40$ t y for IH.

○ $M=2, f=5$ ($\varepsilon'=5\%$, $R=1\%$), $NT=20$ t y for IH, 200 ty if ^{n}Te .

Summary

1. Single β NMEs with $i=1^+, 2^-, (3^+)$ **low-states** and **GRs** by **high ΔE (He,t) CERs**, and β^+ **GR** ($1^+, 1^-, 2^-$) by μ CER ($\mu, \nu_\mu, \nu_n \beta\gamma$) help/confirm theories for $g_A^2 M_{\beta\beta}$.
2. Experimental $M^\beta(1^+, 2^-, 3^+)$ for low states by CERs are used to evaluate $M^{\beta\beta}$ on the basis of FSQP
3. $M^\beta(1^+, 2^-, 3^+)$ are reduced from QP by $k^{\text{eff}} \sim 0.2$,
 $k_{\tau\sigma} \sim 0.4$ due to nuclear $\tau\sigma$, and
 $k_m \sim 0.5$ due to medium isobar effects ($g_A \sim 0.65$),
which are not included in pnQRPA
4. Solar ν CC & NC with DBD nuclei/atoms and atoms in liquid scintillator with DBD are, if $\varepsilon' \sim 5\%$ resolution, **BG \sim a few/t y $\sim S(\beta\beta)$** , serious BG for IH exps.

Remarks

It is timely to discuss realistic and coordinated studies of IH-DBD EXPs and DBD NMEs.

- 1. EXPs with $N_{\text{tons}} > 80\%$, $B < 1/t$ y, $\Delta E < \text{a few } \%$**
- 2. NMEs theories and CER experiments
if $M^{0\nu} > \text{or } < 2$, if $g_A \sim 0.9 \sim 0.6$**
- 3. The NUMEN workshop/project will
contribute much to studies of DBD NMEs**

Thank you for your attention

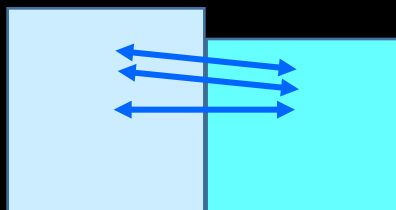


Ejiri-weekend house at Shounan

Group I

$$28 < N < 50$$

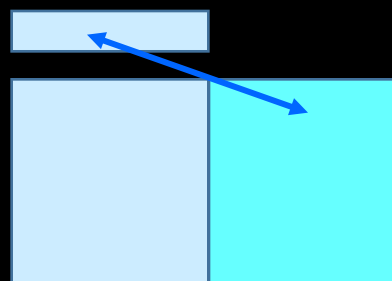
$$28 < Z < 50$$



Group II

$$50 < N < 82$$

$$28 < Z < 50$$



Group III

$$50 < N < 82$$

$$50 < Z < 82$$

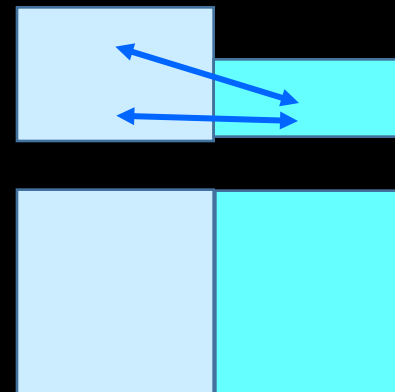


Table 1. FSQP shell configurations for DBD nuclei.

Group	Z, N	FSQP GT configuration	DBD nucleus
I	$28 \leq Z, N \leq 50$	$1g_{9/2}^p 1g_{9/2}^n, 1f_{5/2}^p 1f_{5/2}^n, 2p_{3/2}^p 2p_{1/2}^n$	$^{76}\text{Ge}, ^{82}\text{Se}$
II	$40 \leq Z \leq 50, 50 \leq N \leq 70$	$1g_{9/2}^p 1g_{7/2}^n$	$^{96}\text{Zr}, ^{100}\text{Mo}, ^{106}\text{Cd}, ^{116}\text{Cd}, ^{128}\text{Xe}$
III	$50 \leq Z, N \leq 82$	$2d_{5/2}^p 2d_{3/2}^n, 1h_{11/2}^p 1h_{11/2}^n$	$^{128}\text{Te}, ^{130}\text{Te}, \text{ and } ^{136}\text{Xe}$

Double $\beta\beta$ symmetry and GR

$$T^{2\nu} = \tau\tau\sigma\sigma/\Delta, \quad T^{0\nu} = \tau\sigma\tau\sigma / |\mathbf{r}_1 - \mathbf{r}_2|,$$

$$[H, T] \sim E_{GG} T$$

$T|i\rangle$ Double (β -phonon) GR,
 GR($\beta\beta$) absorbs most $T(\beta\beta)$,
 Little for other; $\langle f|T|i\rangle \sim 0$

$$2\nu\beta\beta \quad T = \tau\sigma\tau\sigma,$$

$T|i\rangle = |DGT \text{ GR}\rangle$, most strength

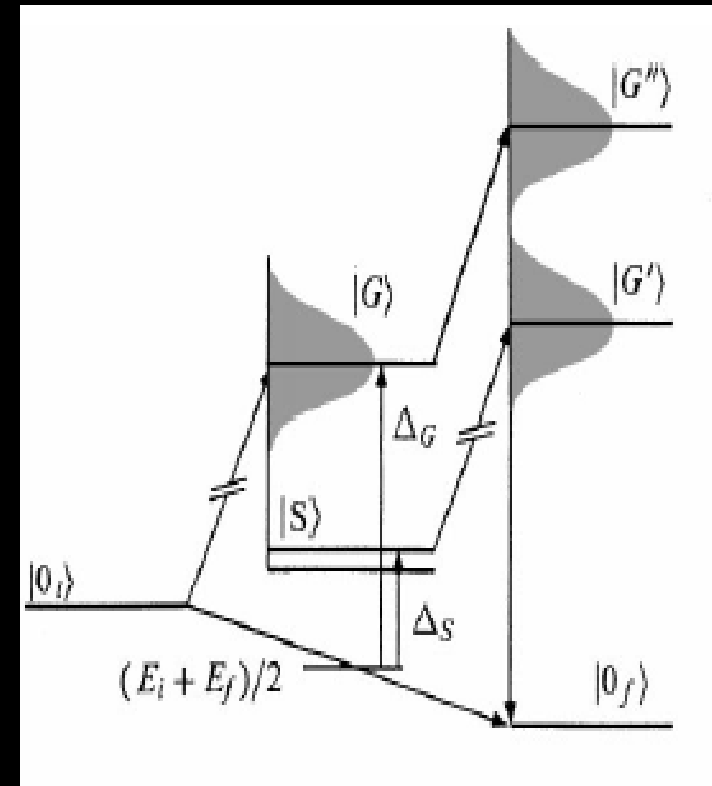
Little GT $\beta\beta$ to low states

$$M^{2\nu} \sim 2 \cdot 10^{-2} M^{2\nu}(\text{SUM}) \quad B(GT^2) \sim 4 \cdot 10^{-4} B(GT^2 \text{ SUM})$$

although $M^{2\nu}$ to low states are favored by the E denominator

$$0\nu\beta\beta \quad T = \tau\sigma\tau\sigma H(\mathbf{r}, \mathbf{r}), \quad T|i\rangle = |\beta\beta \text{ GR}\rangle,$$

Little $0\nu\beta\beta$ strength left to low states as in $2\nu\beta\beta$??.



5. Concluding remarks



Concluding remarks

It is crucial and timely to discuss realistic and coordinated IH ν -mass DBD exps.

Key points are

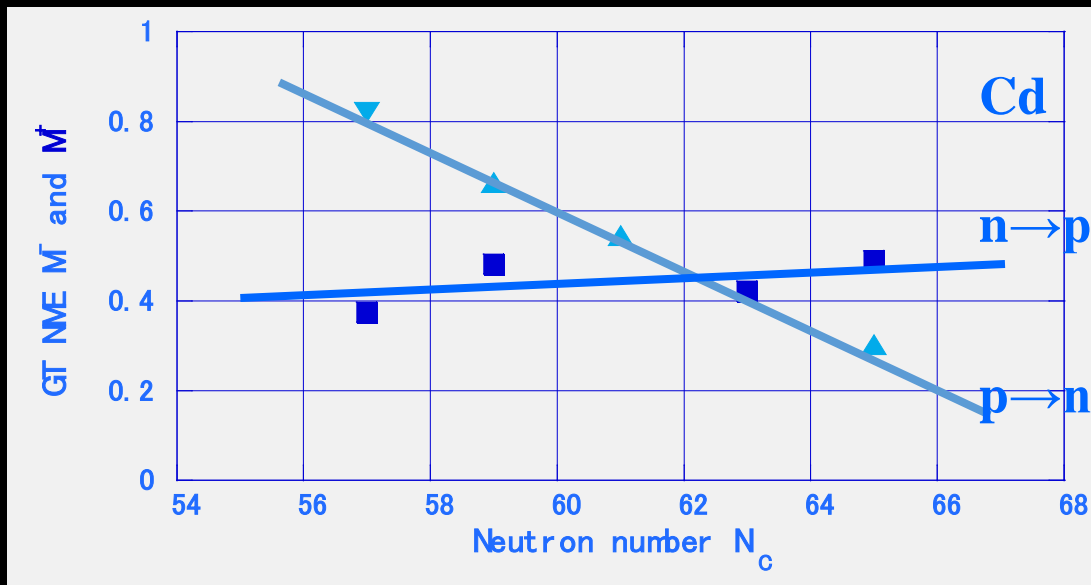
$S(\beta\beta) > 1-2/y$ for 25 meV
 $M^{0\nu}(g_A) \sim 2$ or ? \longleftarrow Present workshop
 $N \sim$ multi-ton $\sim 80\%$ enrichment,

BG $\sim 0.3/t$ y or less
 $\varepsilon =$ resolution $<$ a few %,
 $R = N(\text{DBD})/N'(\text{others}) > 0.1$

$$M(\beta\beta) = M^- M^+ / E$$

$$M^- (n \rightarrow p \quad V_n \quad U_p) \quad M^+ (p \rightarrow n \quad V_p \quad U_n)$$

$$P = (P^- = V_n \quad U_p) \quad (P^+ = V_p \quad U_n) \sim \text{constant}$$



A. Muon capture nuclear γ at SIN

Ordinary muon capture (OMC) studies for the matrix elements in $\beta\beta$ decay.

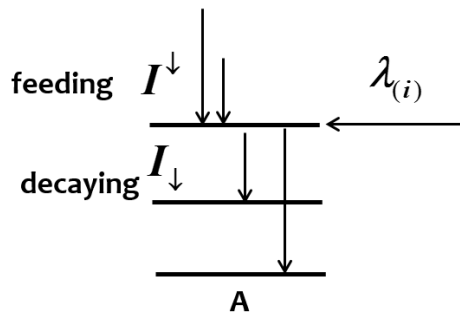
What do we gain from μ -capture for the $\beta\beta$ decay?

Experiment: PSI, beamline μ -E4

K.Ya. Gromov, D.R. Zinatulina, D. Frekers, C. Briançon, V. Egorov, R. Vasiliev, M. Shirchenko, I. Yutlandov, C. Petitjean, J. Deutsch.

14.06.2013
JINR DLNP

Extraction of the partial rates



$$\lambda_{(i)} = \frac{\sum I_{\downarrow} - \sum I_{\downarrow}}{\varepsilon \sum I(nK)}$$

detailed balance

$$\lambda_{cap} = \lambda_{total} - Q\lambda_{decay} \quad Q \rightarrow \text{Huff-factor}$$

$$\lambda_{(i)} [\%] = \frac{\lambda_{(i)}}{\lambda_{cap}}$$

(μ, ν_{μ}) low states,

No ν_{μ} measurement, but many γ from and to the state

Very hard to get the strength

Limited to low states.

J. Suhonen, M. Kortelainen Czech. J. Phy. 56 2006 519, EXP 453.

V. Egorov ($\mu, \gamma, n\gamma, p\gamma$) on ^{48}Ti , ^{76}Se 2004 γ -exp.

SD 2- $\tau\sigma$ NN & nuclear medium g_A

$$M(\text{EXP}) = k M(\text{QP})$$

$$k = k_{\tau\sigma} k_{\text{NM}} \sim 0.2$$

$$M(\text{QRPA}) = k_{\tau\sigma} M(\text{QP})$$

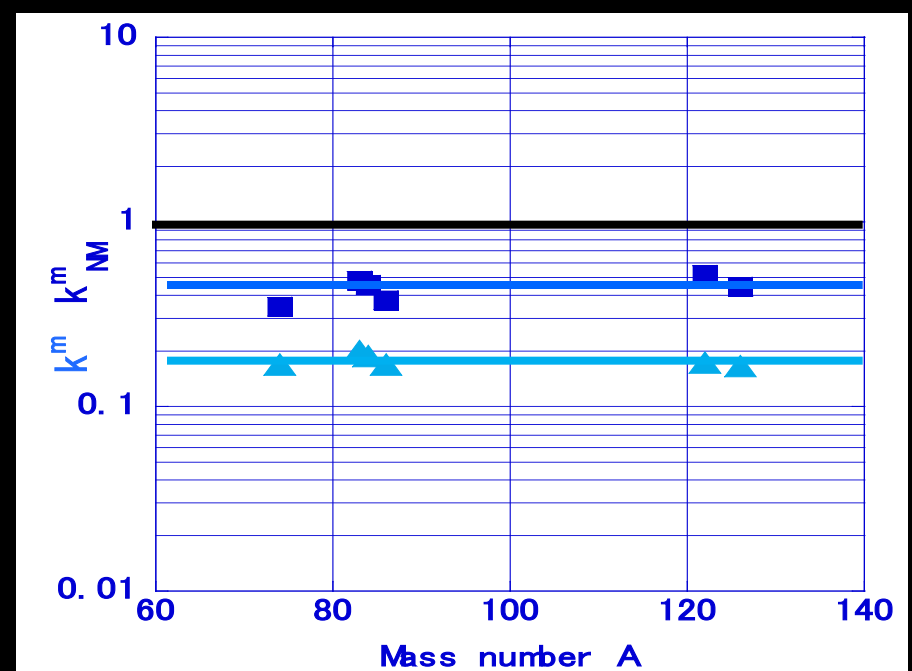
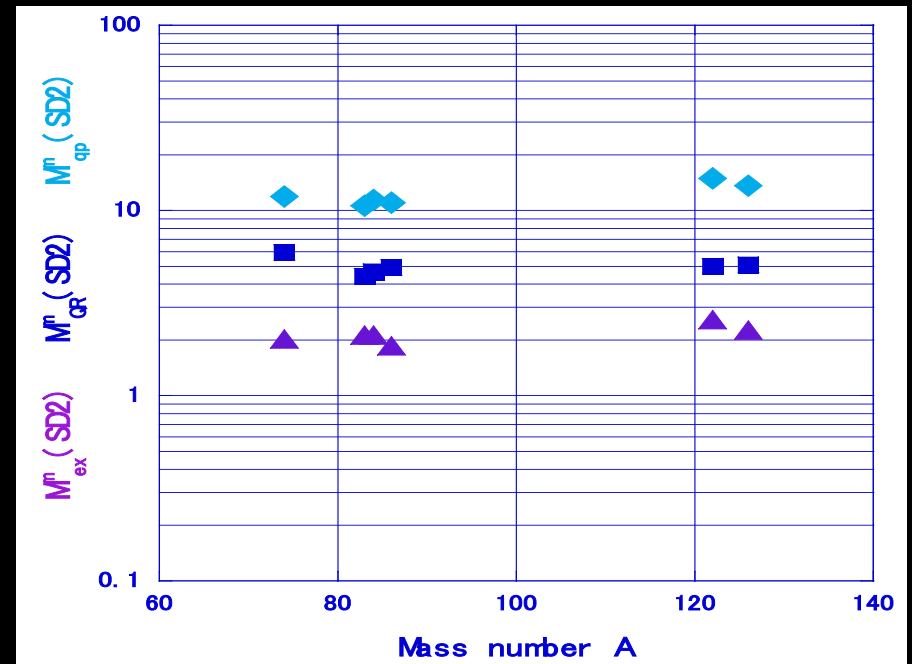
$$k_{\tau\sigma} \sim 0.4 \quad \tau\sigma \text{ correlation}$$

$$M(\text{EXP}) = k_{\text{NM}} M(\text{QRPA})$$

$$k_{\text{NM}} \sim 0.5 = g_A^{\text{eff}} \text{ n-medium}$$

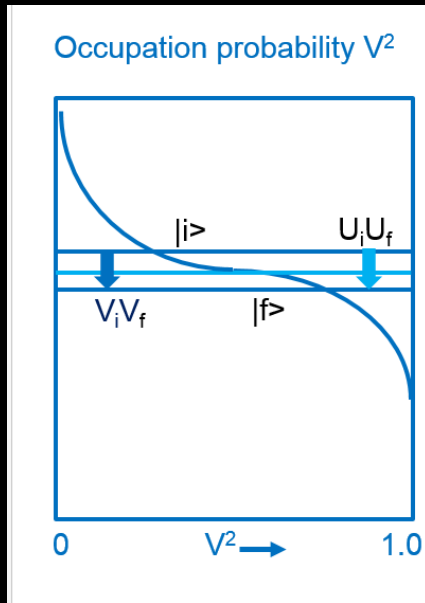
H. Ejiri, N Soucouthi, J. Suhonen,
PL B729, 27 2014 .

Similar g_A in J. Suhonen O. Civitarese
PLB 725 (2013) 153



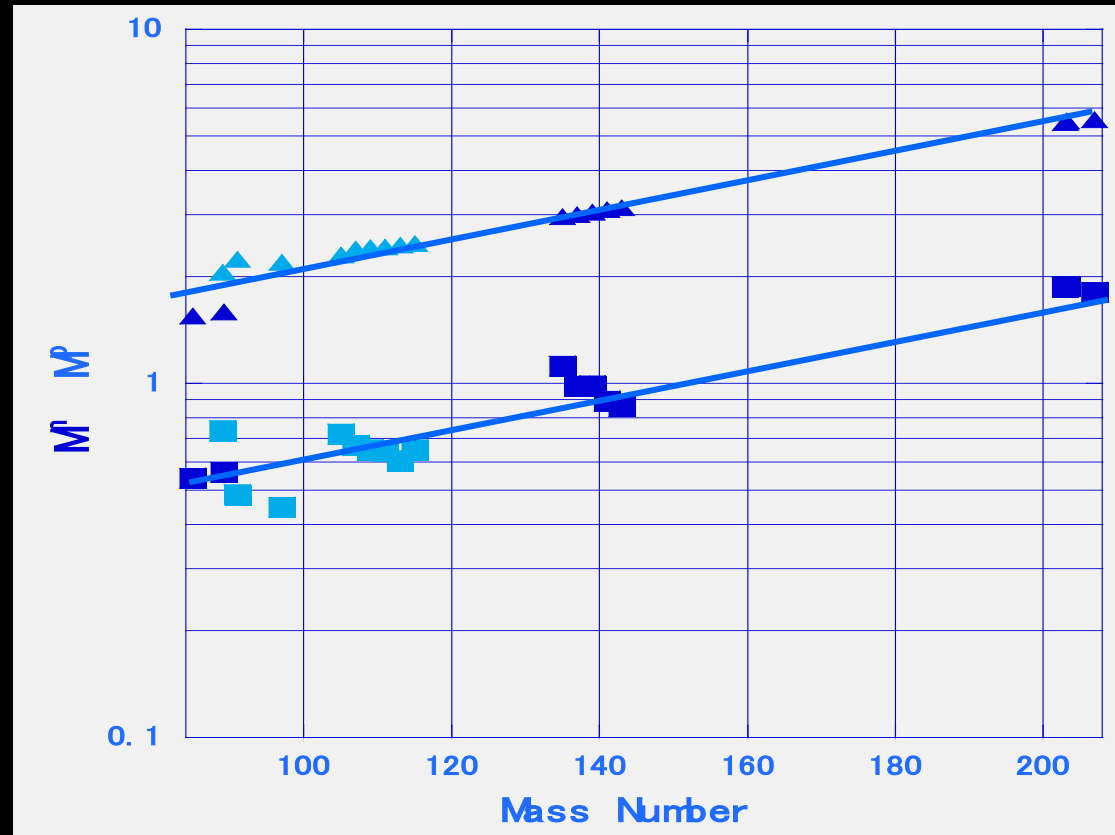
SO (Spin Octupole) ν -responses $=(\text{NME}(\text{SO}))^2$

$$M(\text{SO}) = \langle \tau^\pm (\sigma \times r^3 Y_3) \rangle \sim M4 \gamma$$



$$M(M4) = M_{\text{sp}} P$$

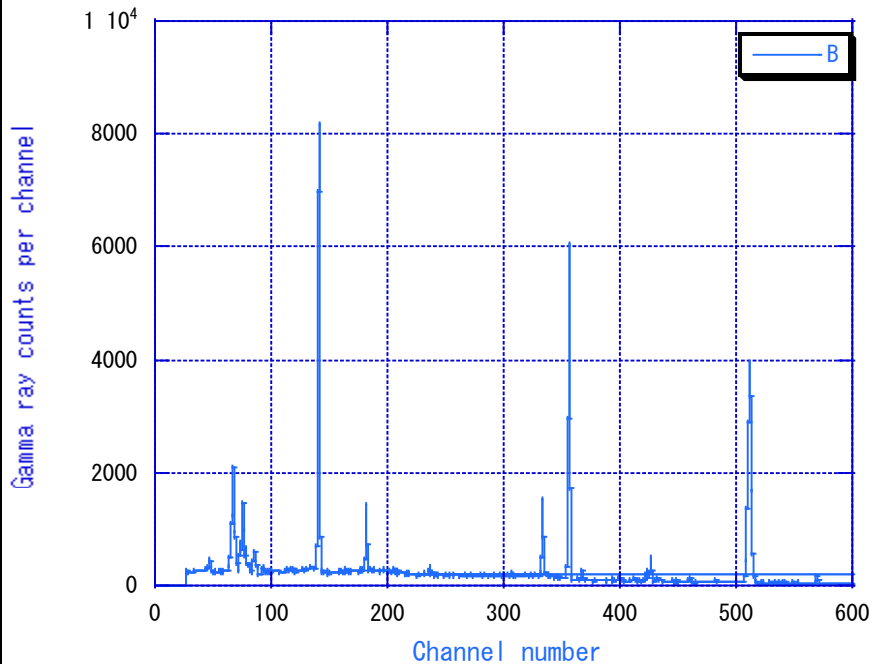
$$P = V_i V_f + U_i U_f \sim 1$$



$$M4 \gamma \quad M(\text{EXP}) = k M(\text{QP}) \quad k \sim 0.33 \quad \text{for proton and neutron}$$

Non destructive nuclear isotope detections

RPID Mo Au (γ, n)

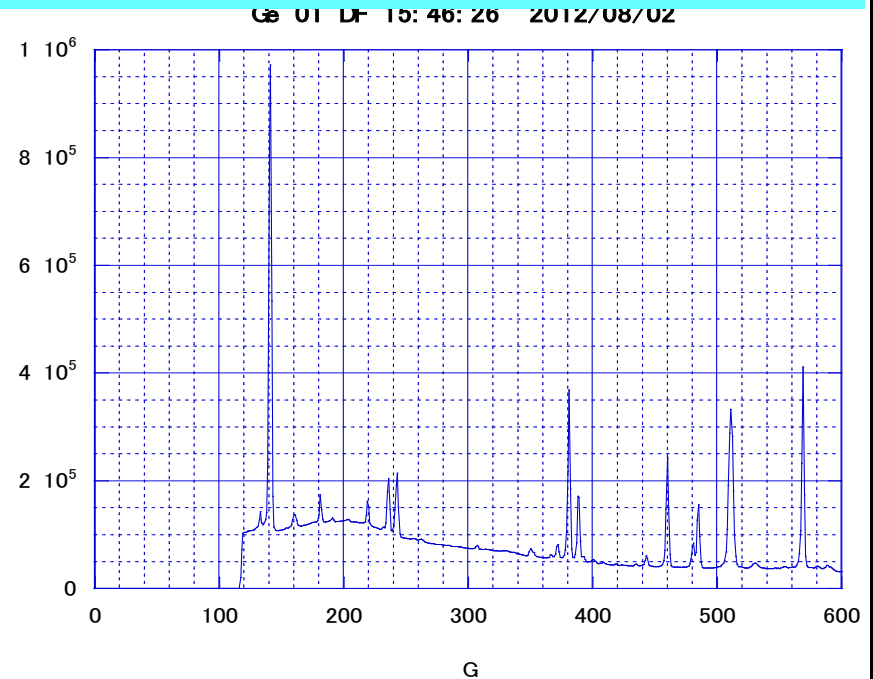


^{197}Au 7.3 % in Mo \Rightarrow 5 ppm \rightarrow ppb
 $3 \times 10^{10} \gamma$ at NewSUBARU

^{99}Tc 140 keV γ for medical applications

H.Ejiri, T. Shima et al., JPSJ 80 (2011),
 Phy Rev ST 15 (2012)

MCID Mo (μ, n)



^{100}Mo 9.6 % in Mo \Rightarrow 7 ppm \rightarrow ppb
 $3 \times 10^9 \mu$ (0.7 hr I μA p) at RCNP

MuSIC ^{99}Tc 140 keV γ for medical

H.Ejiri, I.Hashim, Y.Kuno, K.Ninomiya,
 A. Sato, T.Shima, K.Takahisa, et al..
 JPSJ 82 2013 044202 (2013).

GR / Medium Effect on single $\beta/\tau\sigma$ responses

P-space ; Quasi-particle

$$M(\text{SQP}) = k_A m_{sp} V_n U_p$$

occupied $n \Rightarrow$ vacant p

Q=1-P space $M = k^{\text{eff}} M(\text{SQP})$

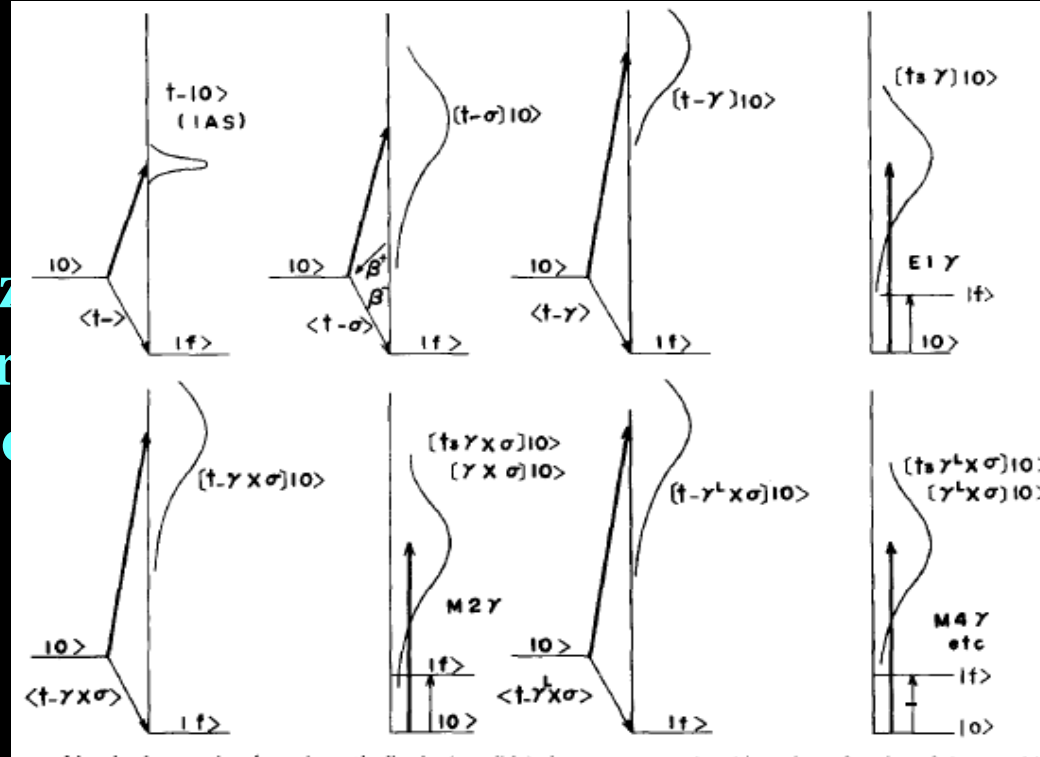
$$k^{\text{eff}} = k^{\text{eff}}(\tau\sigma) k^{\text{eff}}(\Delta)$$

$k^{\text{eff}}(\tau\sigma) \sim 0.33$ from $\tau\sigma$ polarization

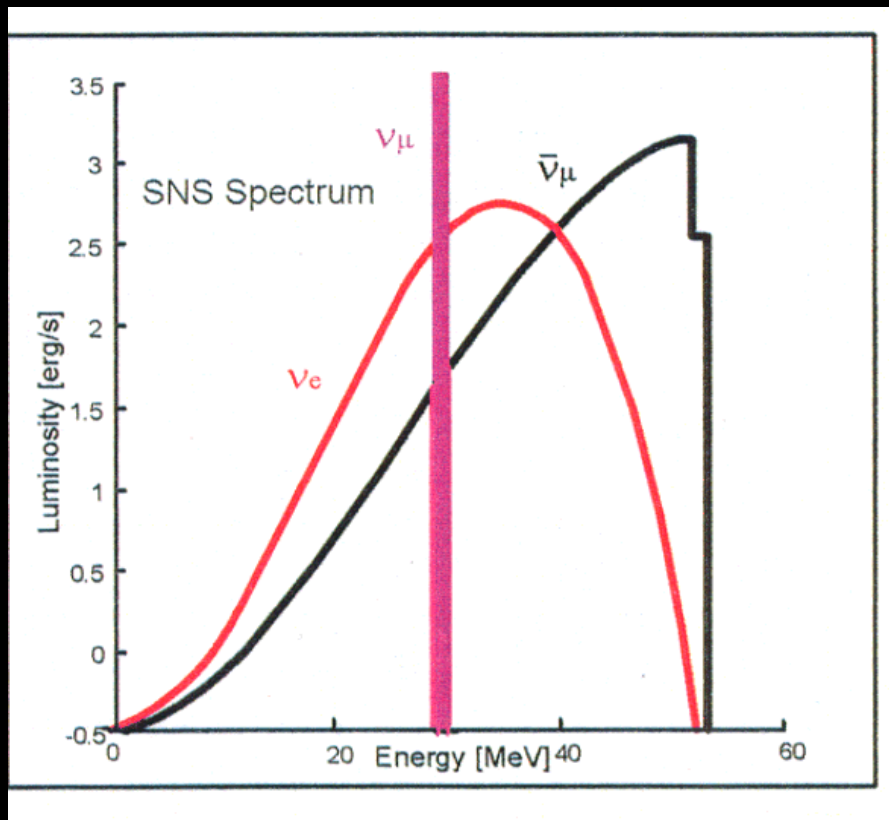
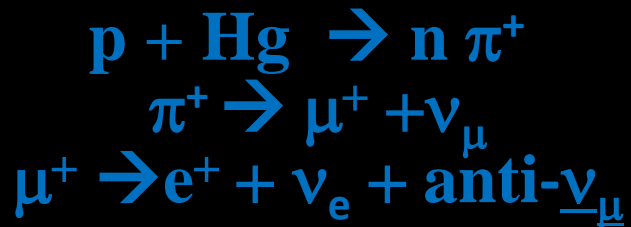
$k^{\text{eff}}(\Delta) \sim 0.7$ $\Delta \tau\sigma$ polarization

Δ effect on GT Bohr M

Small admixtures of GR and Δ/π , nuclear medium effects are hard to calculate. k^{eff} : uniform effect is given by exp.



Low energy ν probes



J-PARC

$$E(\nu_e) E(\nu_\mu) \sim 20\text{-}50 \text{ MeV}$$

$$\sim E(\nu \text{ SN}), E(\nu \beta\beta)$$

$$I(\nu) \sim 10^{14} / \text{MW}$$

$$\sigma \sim 10^{-40} \text{ cm}^2$$

Detectors ~ 1 ton

J-PARC 0.4 – 3 GeV,

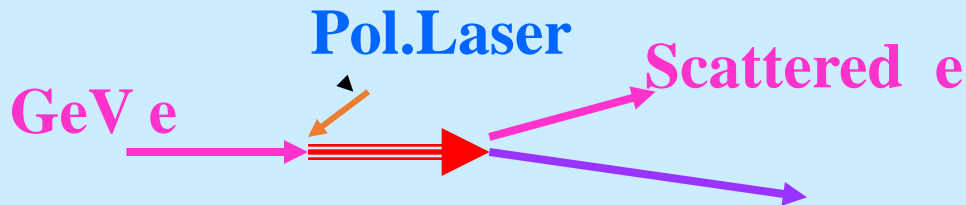
SNS 1 GeV

H. Ejiri NIM. 503 (2003) 276 – 278.

RCNP 1MW under discussion

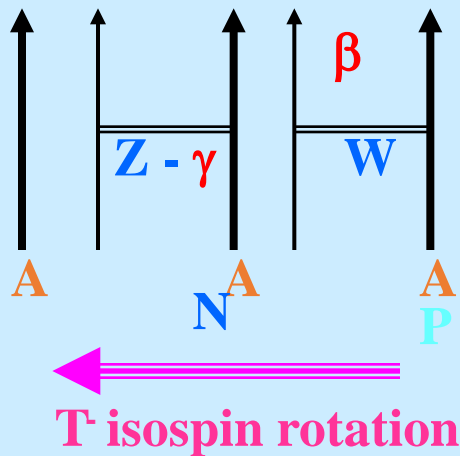
Photon Probes

LEPS Laser electron photon Sources



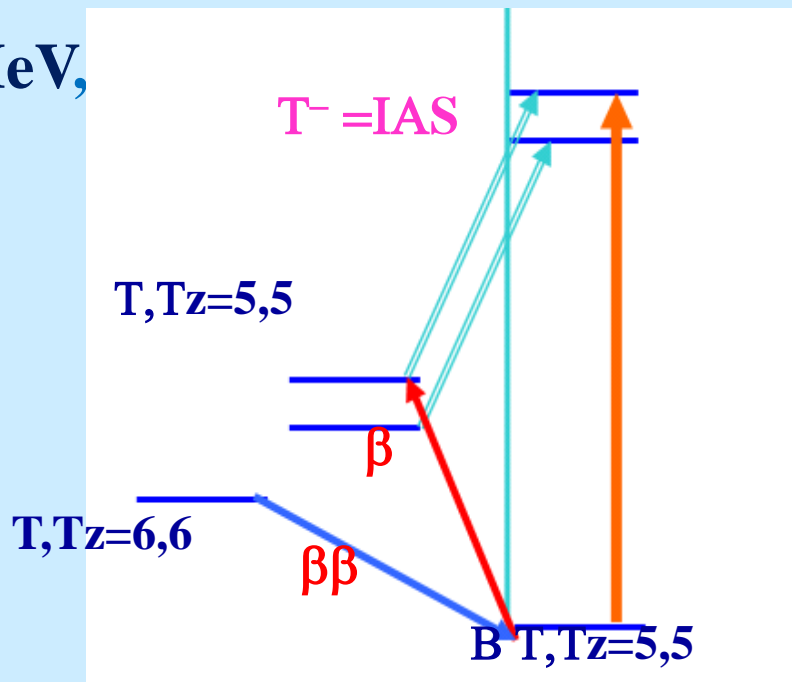
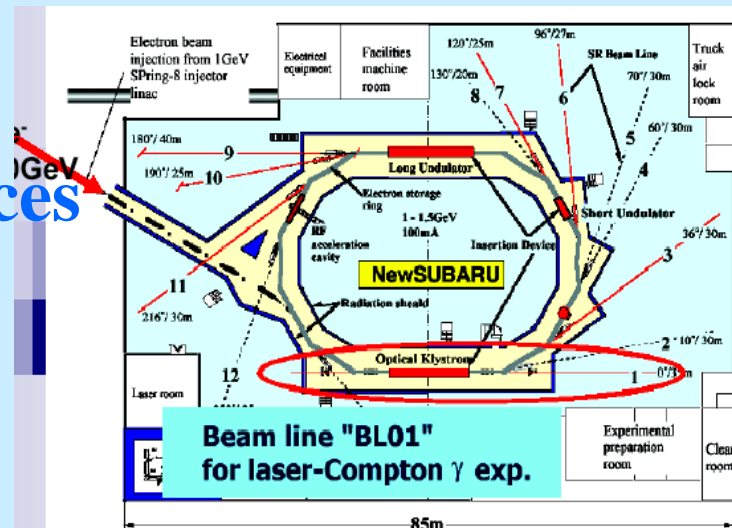
Pol. Laser electron photon

$$E_\gamma \sim 4 \gamma^2 E_l = 2 \cdot 10^8 E_l \sim 20 \text{ MeV,}$$



Responses via IAS

$$\langle f | \mathbf{g} M_\beta | i \rangle = \mathbf{g}/e (2T)^{1/2} \langle f | \mathbf{e} m_\gamma | I \rangle$$

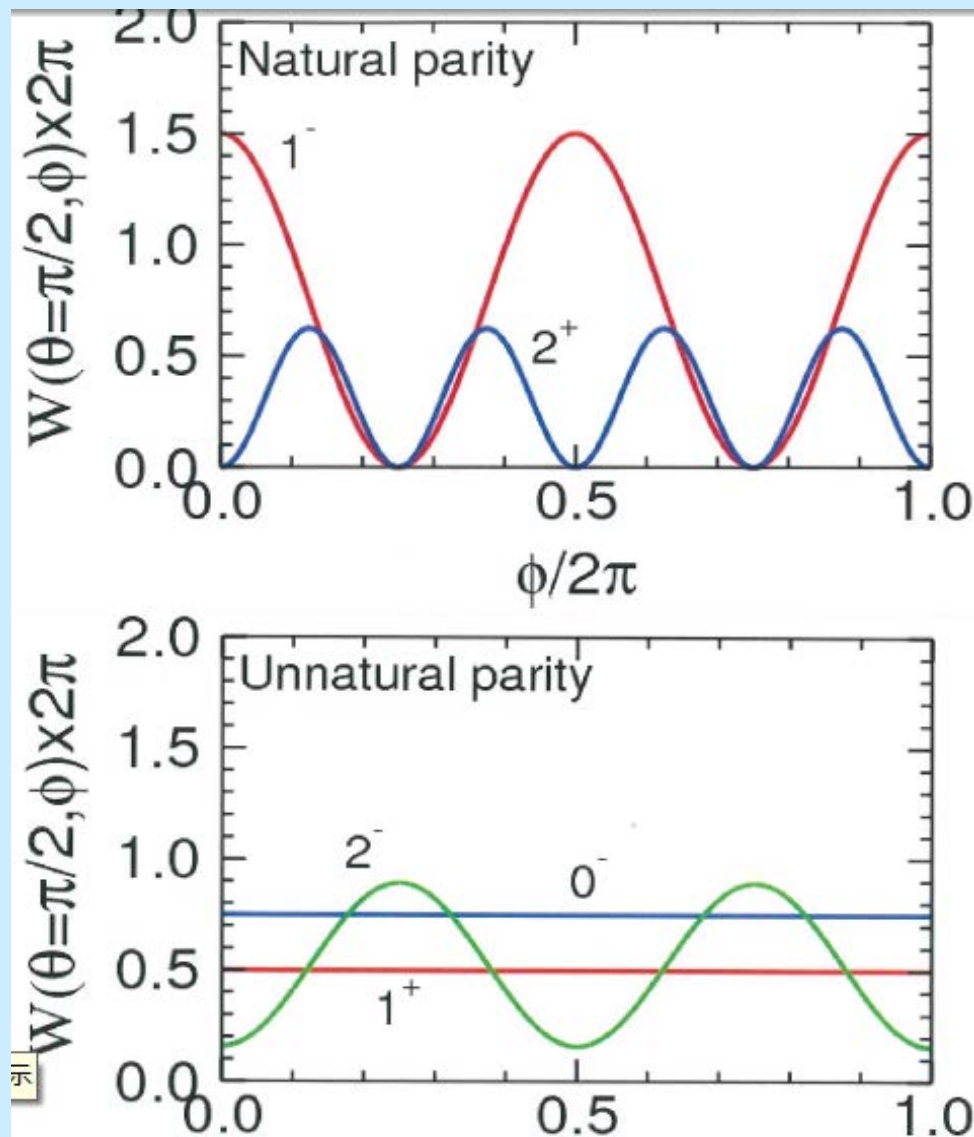
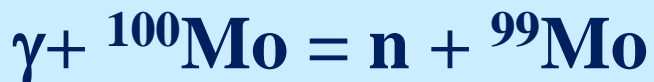
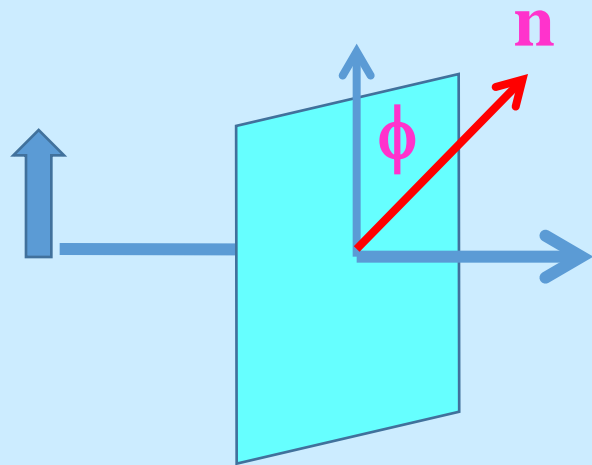


H. Ejiri PRL 21 '68, PR 38 '78

Polarized photon

Vector C τ 1^-

Axial vector C $\tau\sigma$ $1^+ 2^-$



Single β $\tau\sigma$ symmetry, $\tau\sigma$ GR, $\tau\sigma$ polarization

1. $T = \beta$ transition/phonon operator

$$T = \tau Y_l, \tau\sigma Y_l$$

2. $[H, T] \sim E_G T$

$T|i\rangle$; T GR, most T strengths, and little $\langle f|T|i\rangle \sim 0$

$T = \tau$ $T|i\rangle = IAS$ No τ Fermi strength

$T = \tau\sigma$, $T|i\rangle = GT$ GR, little ($\sim 10^{-2}$) GT strength to low states

$T = \tau\sigma r Y$, $T|i\rangle = SD$ GR, little 2^- strength to low states

Civitarese, Muto, T symmetry

3. T polarization

$$|f\rangle = |f\rangle_0 - \varepsilon |GR\rangle \quad M = M_0 [1 - \varepsilon M_{GR}/M_0] = k^{\text{eff}} M_0$$

$\varepsilon \sim 0.05$, $\varepsilon^2 \sim 0.003$ admixture of GR makes $k^{\text{eff}} = 0.2 - 0.1$, as expts.

Experiments for $\beta/\beta\beta$ responses/ ME

To help evaluate/confirm $M^{2\nu}$, $M^{0\nu}$ theories.

1. Absolute/relative values for $M^{0\nu}$, $g_A^{\text{eff}} M^{0\nu}$ 20 %

Absolute values for $[m_\nu, \text{SUSY ...}]$ within 20%

Hierarchy : $m(\text{IH})/m(\text{NH}) > 4$, depending on α, δ

Mechanisms of light ν -mass /SUSY .

2. Experimental design

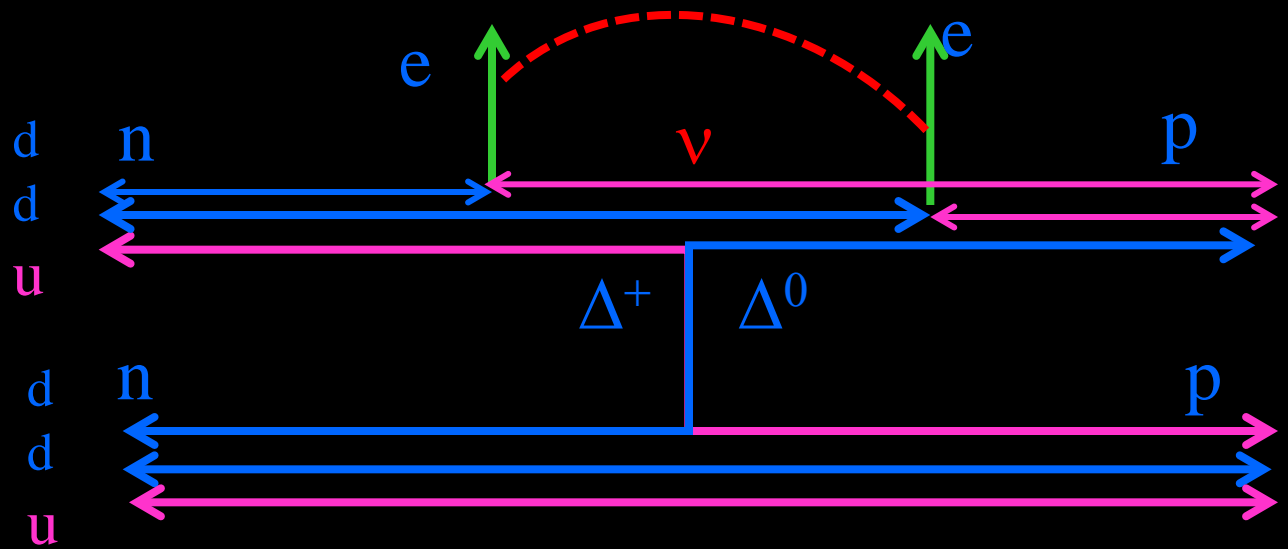
Detector mass factor 4 more if $M^{0\nu}$ is 30 % less

Choice of detector ($\beta\beta$) nuclei and detector volume

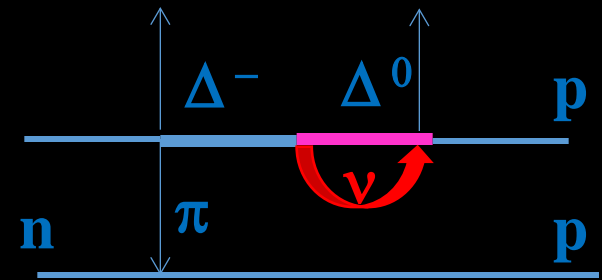
3. Experimental consistency if light ν -mass term.

Non nucleonic
degree Δ

ν -exchange
in $\Delta 1\sigma$
enhanced



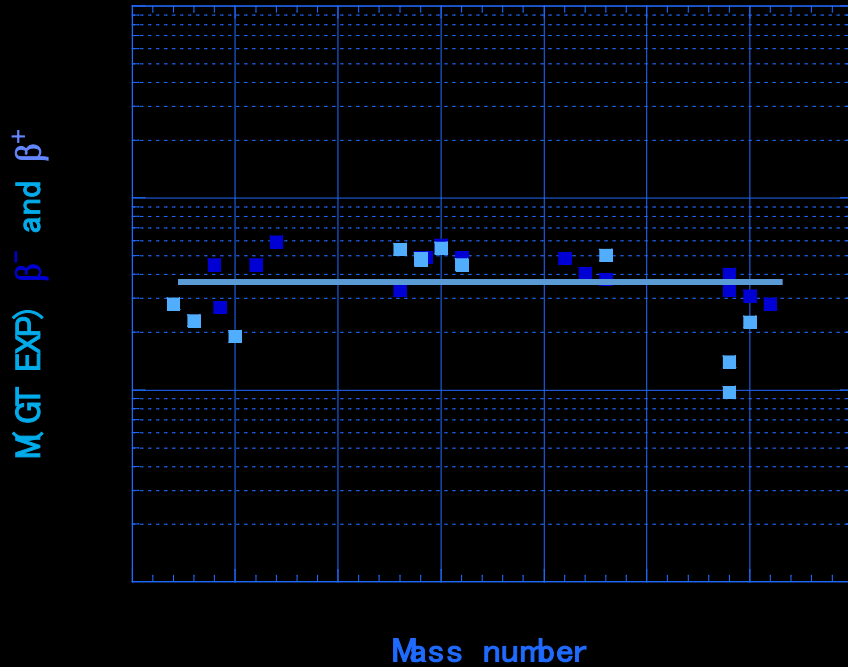
Primacoff & Rosen process
forbidden in $0^+ - 0^+$ Doi, Kotani



Intermediate $\phi(pn) + \varepsilon\phi(\Delta n) + \varepsilon\phi(\Delta p)$ $M^\beta \sim k^{\text{eff}}(D)M^\beta_0$

Nuclear Matrix elements from 0^+ to 1^+ and 2^-

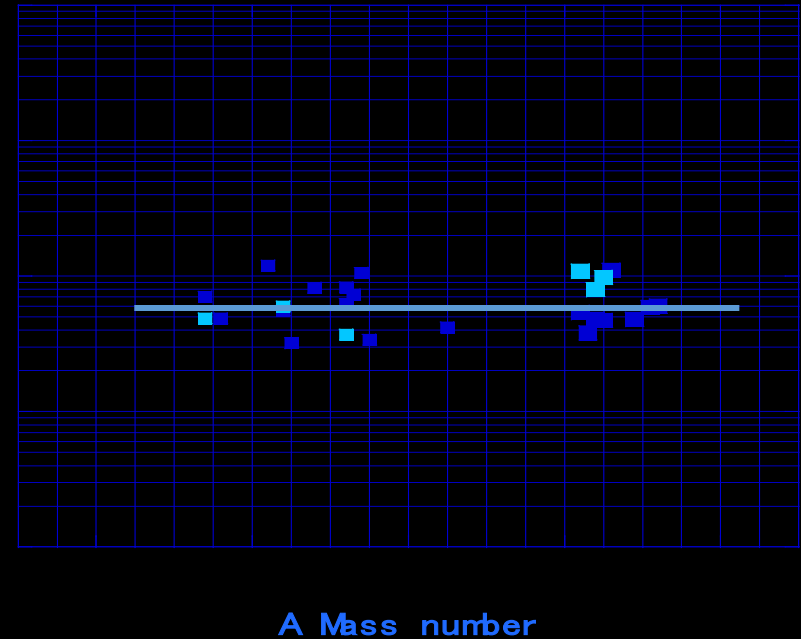
$J=1^+$ GT M (EXP) β^- β^+



$$M(1^+) \sim 0.4$$

$$k^{\text{eff}} = M(1^+)/M(\text{SQP}) \sim 0.25$$

$J=2^-$ $[\sigma r Y_1]$ M (EXP) 10^{-3} β^- β^+



$$M(2^-) \sim 0.6 \cdot 10^{-3} \text{ n.u}$$

$$k^{\text{eff}} = (M(2^-)/M(\text{SQP})) \sim 0.1$$

E keV	T _{1/2}	100Mo	98	97	96	95	94	92	RI
									A RT
140.5	66+6h	1n							99 Nb-Mo-Tc
181	66h	1n							99 Nb Mo Tc
658.1	72.1h	3n	1n						97 Nb Mo
743.4	17h	3n	1n						97 Nb Mo
459.9	23h	4n	2n	1n	0n				96 Nb Mo
778.2	23h	4n	2n	1n	0n				96 Nb Mo
849.9	23h	4n	2n	1n	0n				96 Nb Mo
1200.2	23h	4n	2n	1n	0n				96 Nb Mo
204.1	87h	5n	3n	2n	1n				95 Nb Mo
934.5	10.2d		6n	5n	4n	3n	2n		92 Nb Zr
141.2	14.6h				6n	5n	4n	2n	90 Nb Mo
1129.2	14.6h				6n	5n	4n	2n	90 Nb Zr
2319	14.6h				6n	5n	4n	2n	90 Nb Zr
908.9	78h					6n	5n	3n	89 Nb Zr Y
388.5	80h					6n	5n	3n	87 Nb Zr Y Sr
484.8	80h						7n	5n	87 Nb Zr Y Sr
380.8	13.4h						7n	5n	87 Nb Zr Y
1076.7	4.7						8n	6n	86 Nb Zr Y St

Nuclear weak/ ν responses for $\beta\beta-\nu$

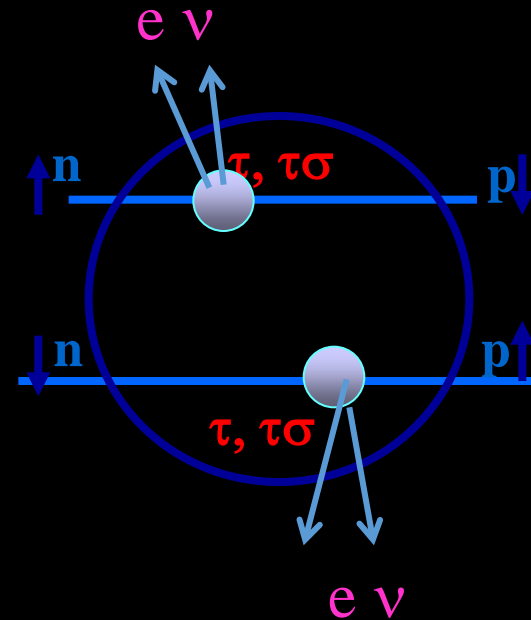
$$2\nu\beta\beta \quad A = B + 2\beta + 2\nu$$

$$T^{2\nu} = G^{2\nu} [M^{2\nu}]^2$$

$$M^{2\nu} = (g_A)^2 \langle \tau\sigma\tau\sigma \rangle$$

GT isospin spin $\tau\sigma 1^*$

F $\tau\tau$ is only to double $|IAS\rangle$



T($\beta, \beta\beta$) mode GR quenches β and $\beta\beta$ to low states,
 $k^{\text{eff}} = 0.3$ for β , $k^{\text{eff}} = 0.3^2 \sim 0.1$ for $\beta\beta$, but
GR itself contributes little to $\beta\beta$ to low states

$$|S\rangle = Q_S |0_i\rangle, \quad Q_S = Q_S^0 - \varepsilon Q_G^0,$$

$$|G\rangle = Q_G |0_i\rangle, \quad Q_G = Q_G^0 + \varepsilon Q_S^0,$$

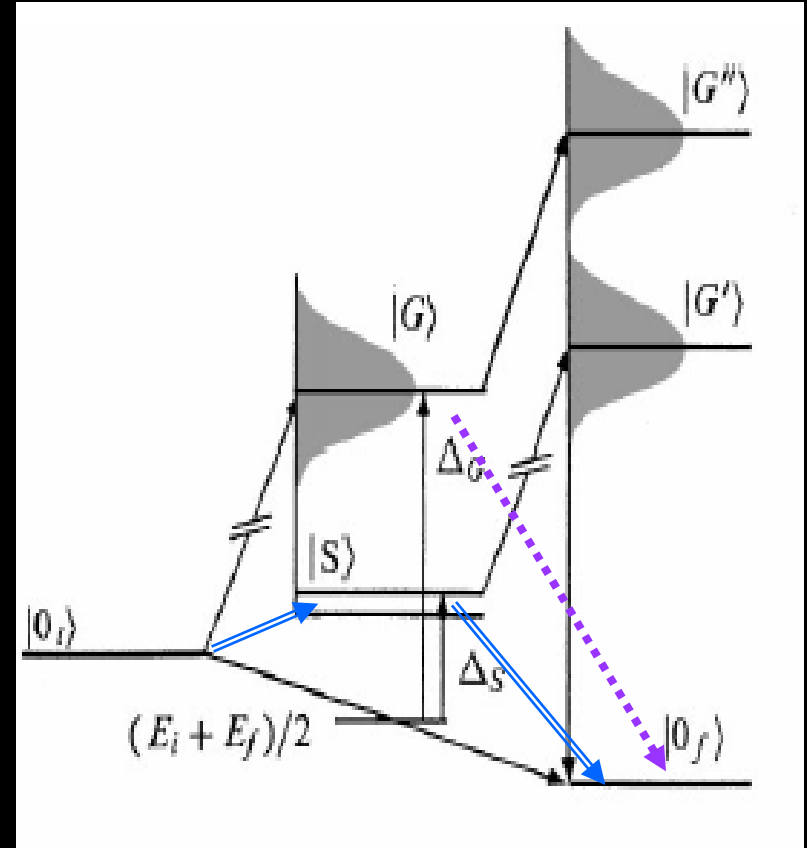
$$|0_f\rangle = Q_{S'} |S\rangle = Q_{S'} Q_S |0_i\rangle \quad \text{with } Q_{S'} = Q_{S'}^0 - \varepsilon' Q_G^0$$

$$M^{2\nu}(GT) \simeq \frac{M_S^\nu M_{S'}^\nu}{\Delta_S} + \frac{\varepsilon M_S^\nu M_G^\nu}{\Delta_G} \left(1 - \frac{\varepsilon'}{\varepsilon}\right).$$

$$M^{0\nu} \sim M_s M_{s'}, [1 + \delta]$$

$$\delta = \varepsilon \frac{M_G}{M_S} (1 - \varepsilon/\varepsilon') \ll 1$$

0.1 10 0.1



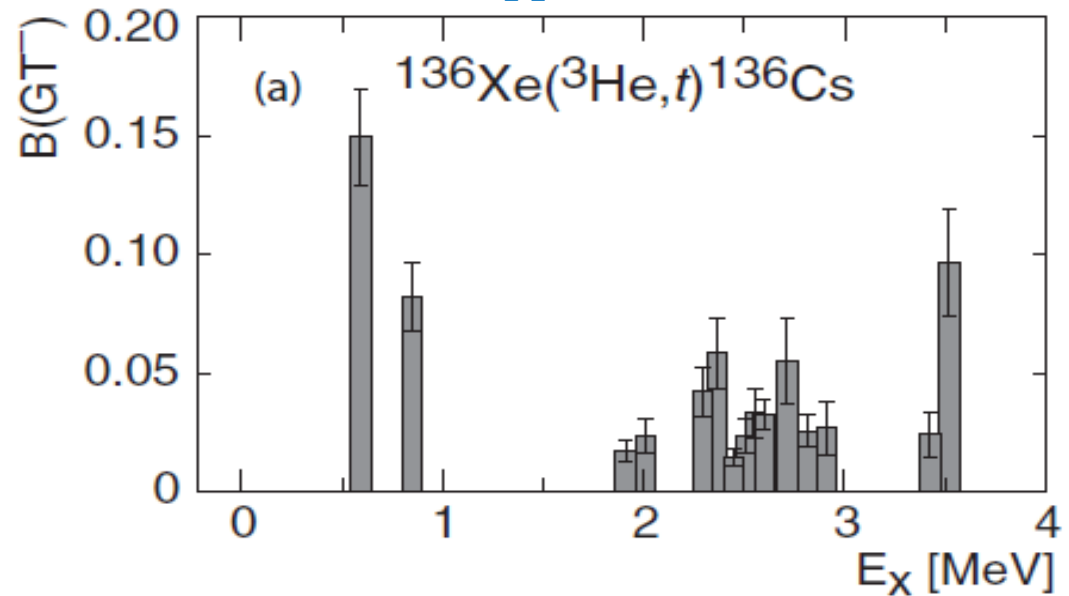
Ejiri Phys. Rep. 338 (2000)

Ejiri NP A577 (1994), JPSJ 65(199

Ericson et al., PL B328 259 '94

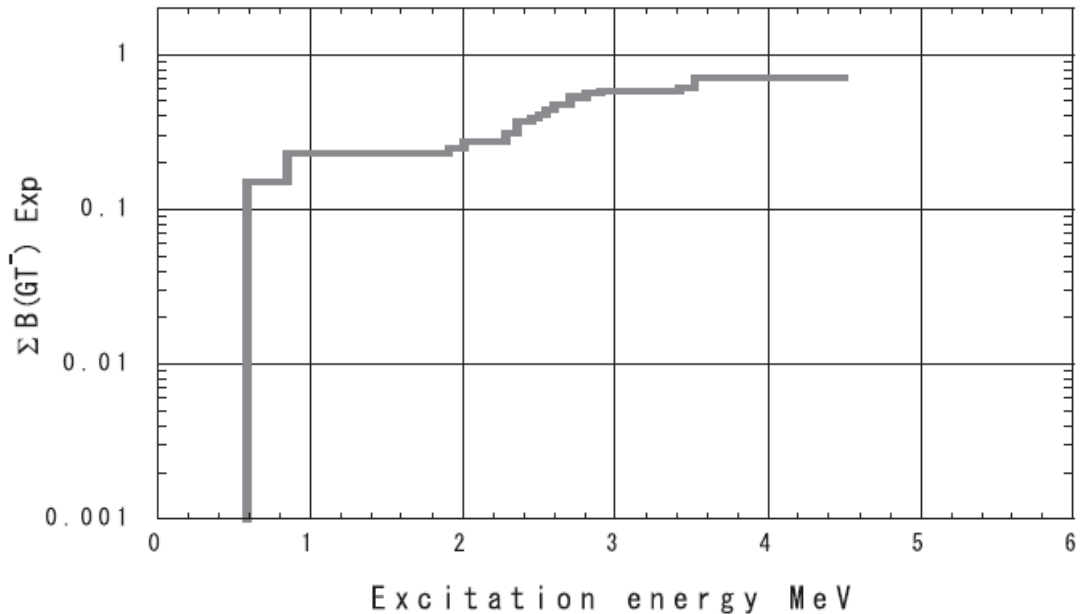
Spin strength

$$B(GT) = |\sigma|^2$$



Sum spin strength

$$\Sigma B(GT) = \Sigma |\sigma|^2$$



$$\sum |M_{\mathbf{k}}^{-}|^2$$

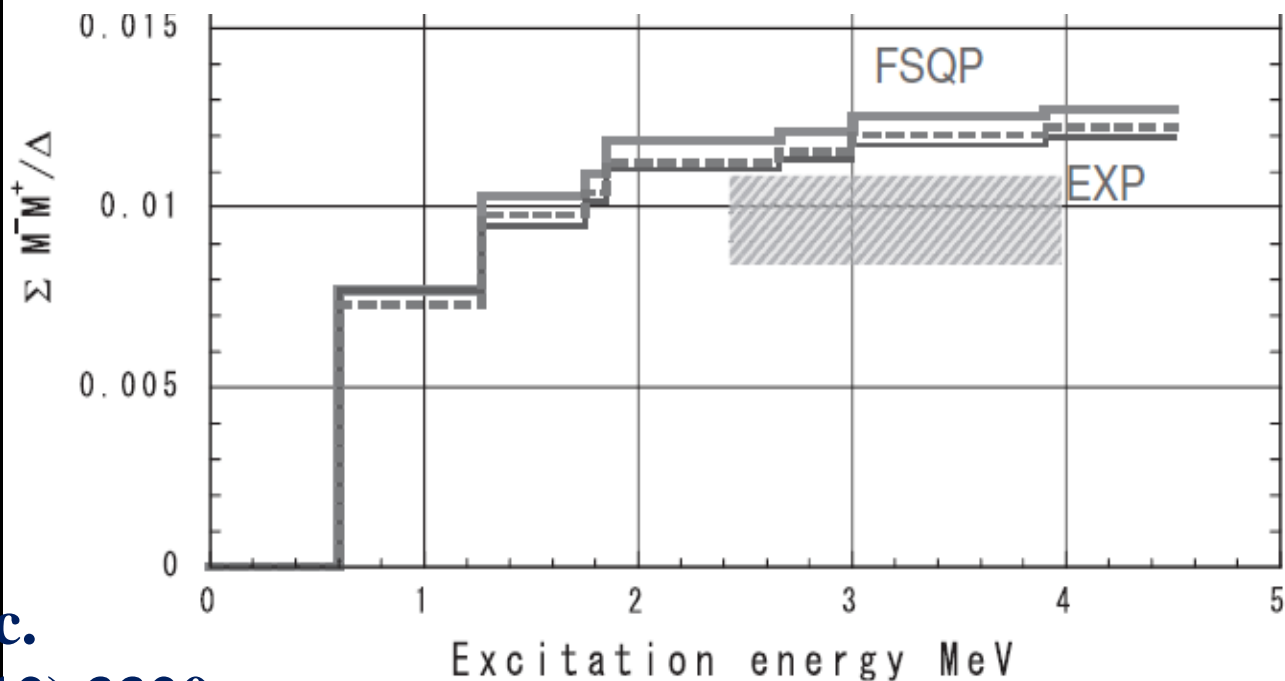
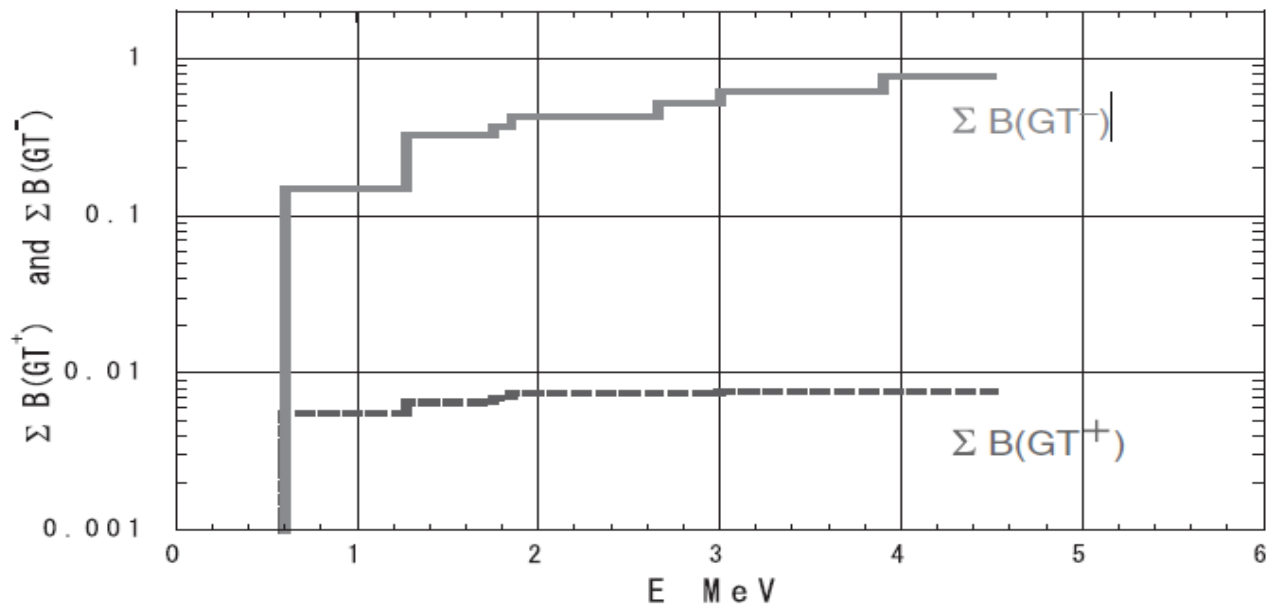
$$\sum |M_{\mathbf{k}}^{+}|^2$$

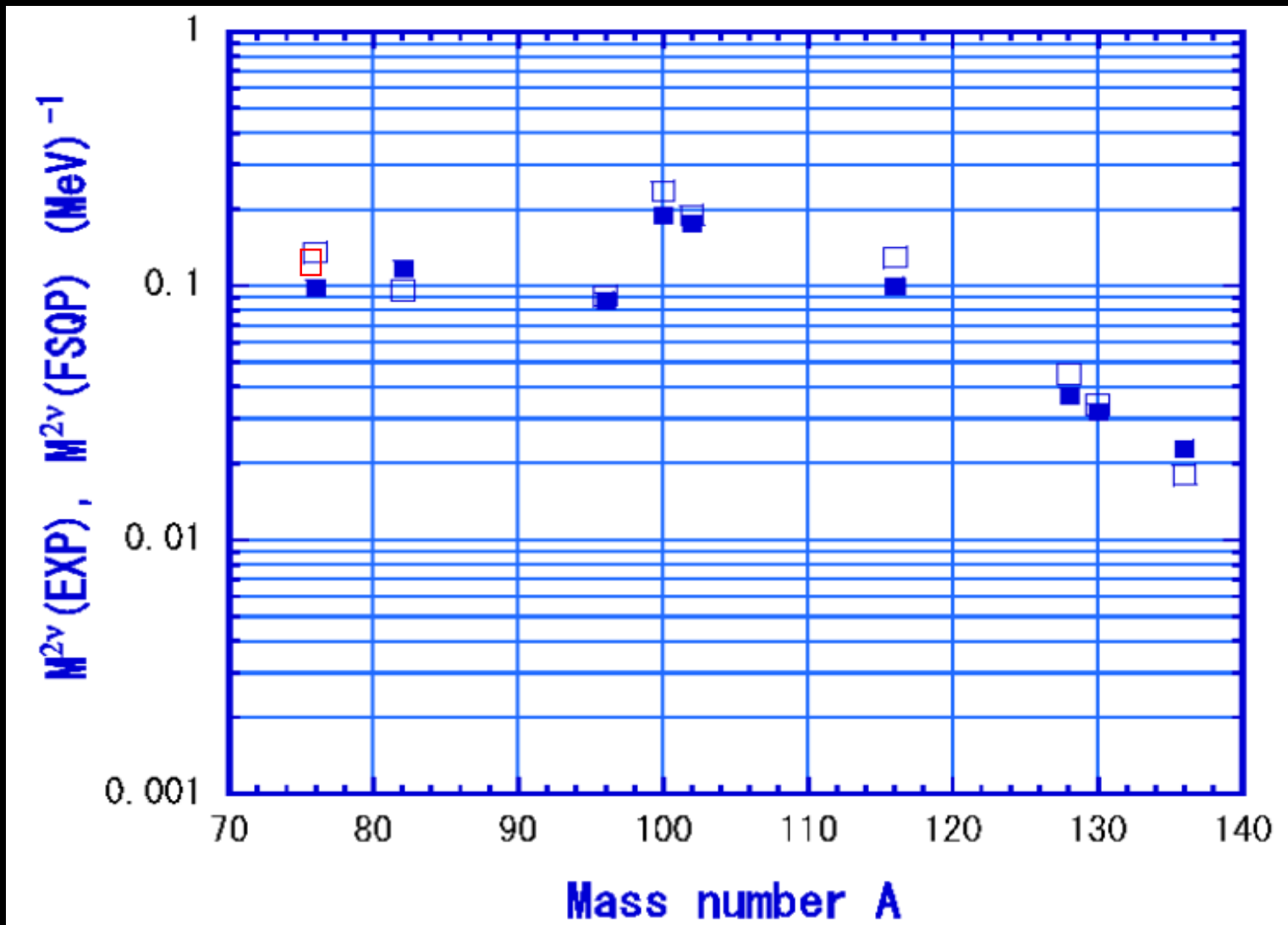
$$M^{2\nu\beta\beta} =$$

$$\sum_{\mathbf{k}} M_{\mathbf{k}}^{-} M_{\mathbf{k}}^{+} / \Delta_{\mathbf{k}}$$

EXP: EXO KamLAND

H. Ejiri J Phys. Soc.
Japan Lett. 81 (2012) 3320



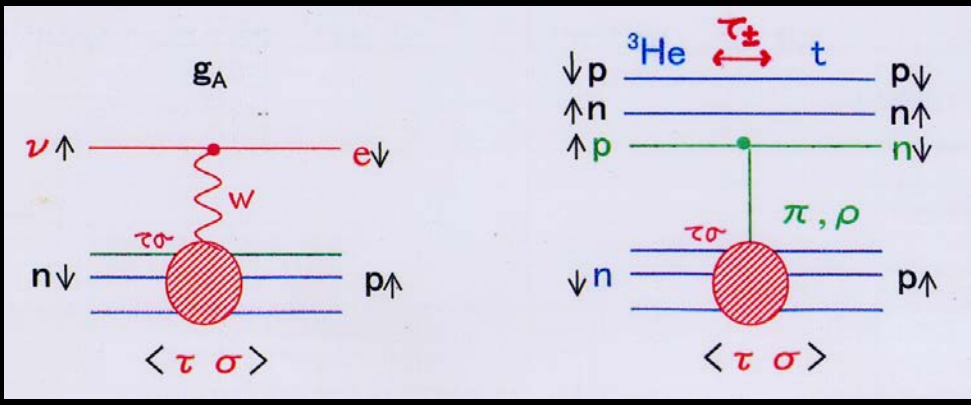


All FSQP 1+ states contribute. In case of ^{96}Zr , ^{100}Mo , ^{116}Cd , where n & p are in N=4 & 3 major shells, only 1 g7/2 g7/2 1+ state. SSD The ground 1+ state (^{100}Mo , ^{128}Te) dominate $2\nu\beta\beta$.

CER: Charge exchange reactions for ν -responses

RCNP ($^3\text{He}, t$) ($^7\text{Li}, ^7\text{Be}$), MSU ($t, ^3\text{He}$), KVI ($d, ^2\text{He}$)

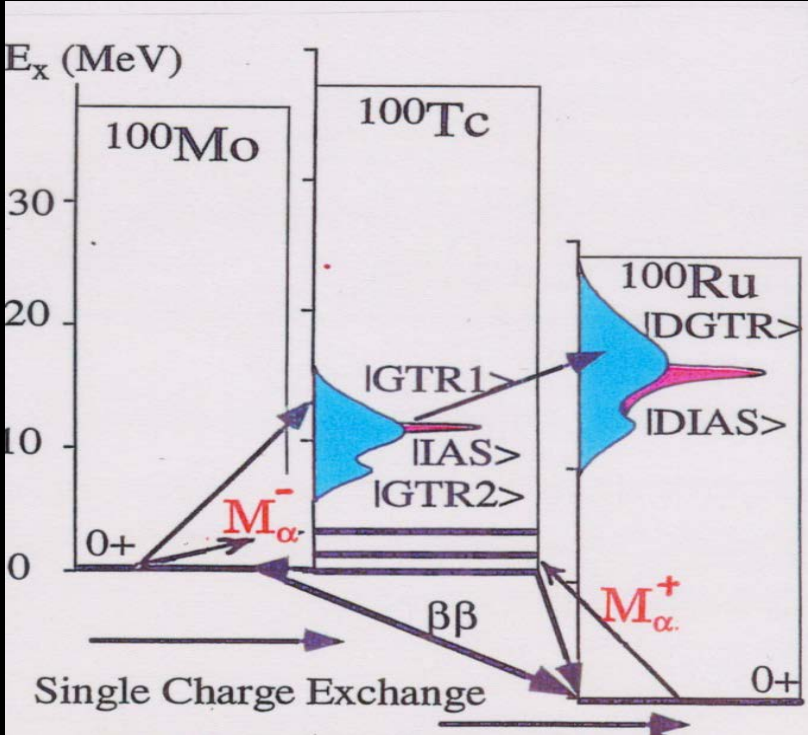
RCNP/Osaka, KVI, Munster, MSU Frekers



$$\frac{d\sigma_\alpha(0^\circ)}{d\Omega} \frac{1}{K(E_i, 0)N_\alpha^D} = |J_\alpha|^2 B(\alpha), \quad \alpha = F, GT.$$

Non-central V

CER q-dependence and EC/ β rates are used to get weak responses



($^3\text{He}, t$), ($^{11}\text{B}, ^{11}\text{Li}$), ($t, ^3\text{He}$)
 β^- $\beta\beta$ β^+

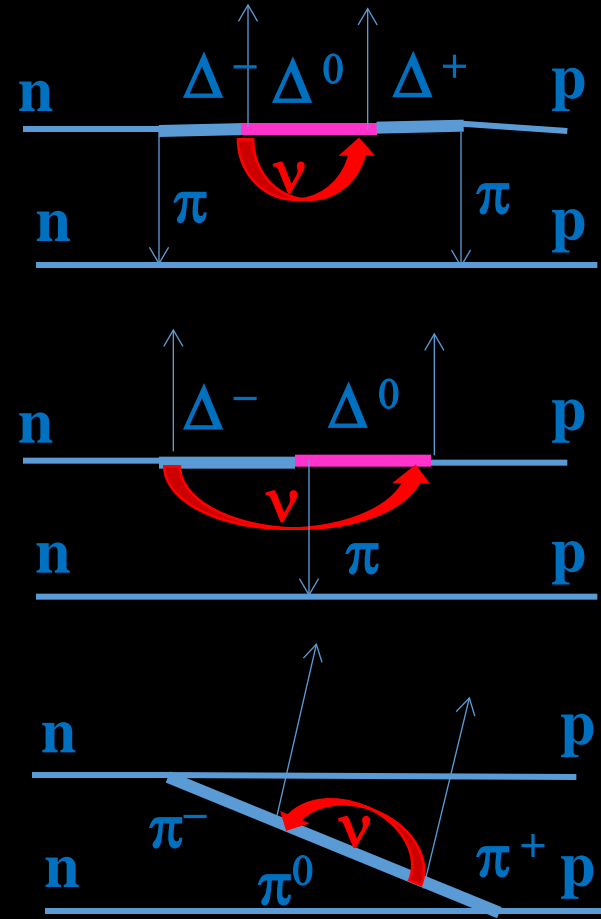
Hadronic (Δ, π) *

$N \Rightarrow \Delta$ quark $\tau\sigma$ flip GR
 0.3 GeV high excitation
 $0\nu\beta\beta$ in Δ, π enhanced

Effect on low $\beta\beta$ $0^+ - 0^+$
 $P(\Delta)^2 \sim (10^{-2})^2 \sim 10^{-4}$

$|n \Delta\rangle$ interferes with $|n p\rangle$
 $|I\rangle = |n p\rangle - \varepsilon |n, \Delta^+\rangle$
 $M^\beta \sim k^{\text{eff}} M_0 \quad k^{\text{eff}}(\Delta) \sim 0.7$

Exp by β deays and CER



*Pontecorvo; Haxton, Stephenson, Kotani Doi .

Nuclear Matrix Elements for Two-Neutrino Double Beta Decays

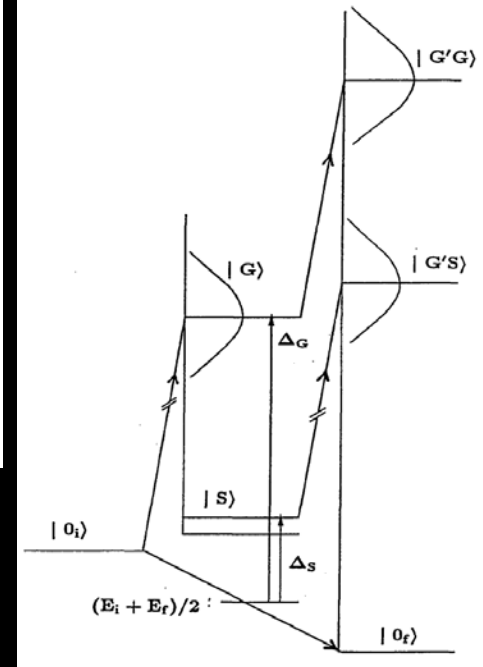
JPSJ Lett 65 '96 7

Hiroyasu EJIRI* and Hiroshi TOKI**

Research Center for Nuclear Physics (RCNP), Osaka University, Mihogaoka, Ibaraki, Osaka 567

(Received October 16, 1995)

Nuclear matrix elements $M^{2\nu}$ for two-neutrino double beta decays in medium-heavy even nuclei are found to be given by a product of matrix elements $M_S^\nu M_{S'}^\nu$ for successive single beta decays through low-lying single particle-hole 1^+ states in intermediate nuclei. The empirical systematics is explained theoretically by coupling of single particle-hole 1^+ states with GT (Gamow Teller) giant resonances.



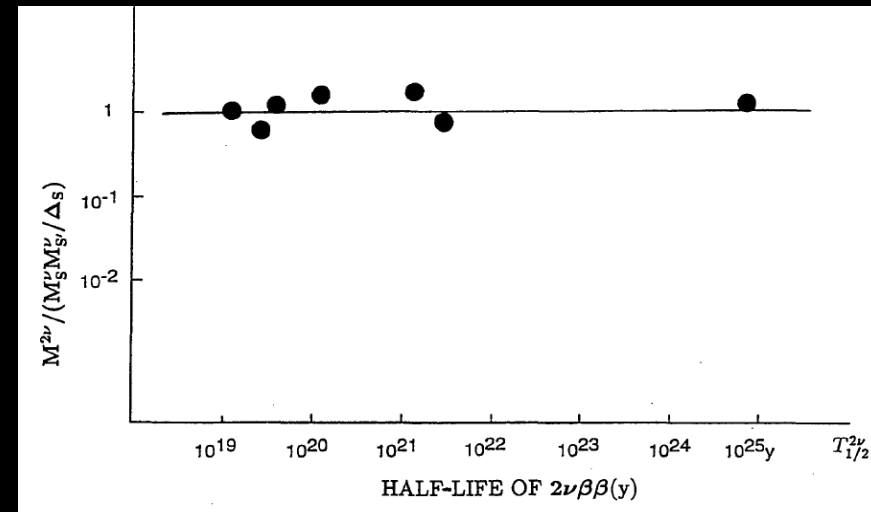
Now let us introduce the $\tau\sigma$ -type interaction H_I between $|S\rangle_0$ and $|G\rangle_0$. Then $|S\rangle$ and $|G\rangle$ are expressed as

$$|S\rangle = Q_S |0_i\rangle, \quad \text{with} \quad Q_S = Q_S^0 - \varepsilon Q_G^0, \quad (11)$$

$$|G\rangle = Q_G |0_i\rangle, \quad \text{with} \quad Q_G = \varepsilon Q_S^0 + Q_G^0. \quad (12)$$

$$M^{2\nu}(1^+) = \left(\frac{M_S^\nu M_{S'}^\nu}{\Delta_S} \right) (1 + \delta),$$

$$\delta = \frac{\Delta_S}{\Delta_G} \cdot \frac{\varepsilon M_G^\nu}{M_{S'}^\nu} \left(1 - \frac{\varepsilon'}{\varepsilon} \right). \quad (25)$$



Start CER to get low-lying GT 1^+ , SD 2^- , ... for $2\nu\beta\beta$ and $0\nu\beta\beta$.