

Probing the Intergalactic Magnetic Fields by means of high energy pair halos around extreme blazars

Paolo Da Vela

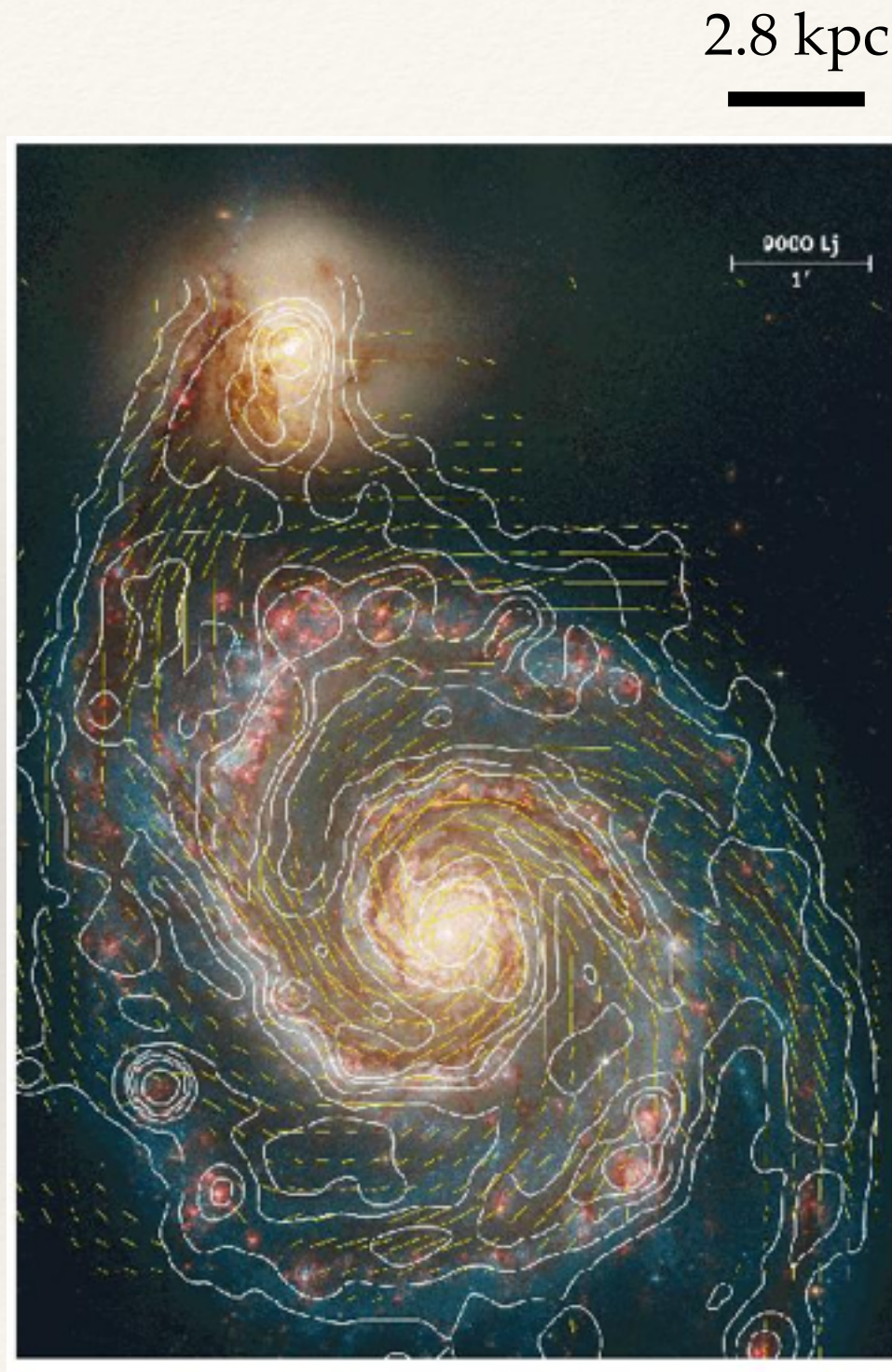
INFN, Sezione di Pisa

Summary

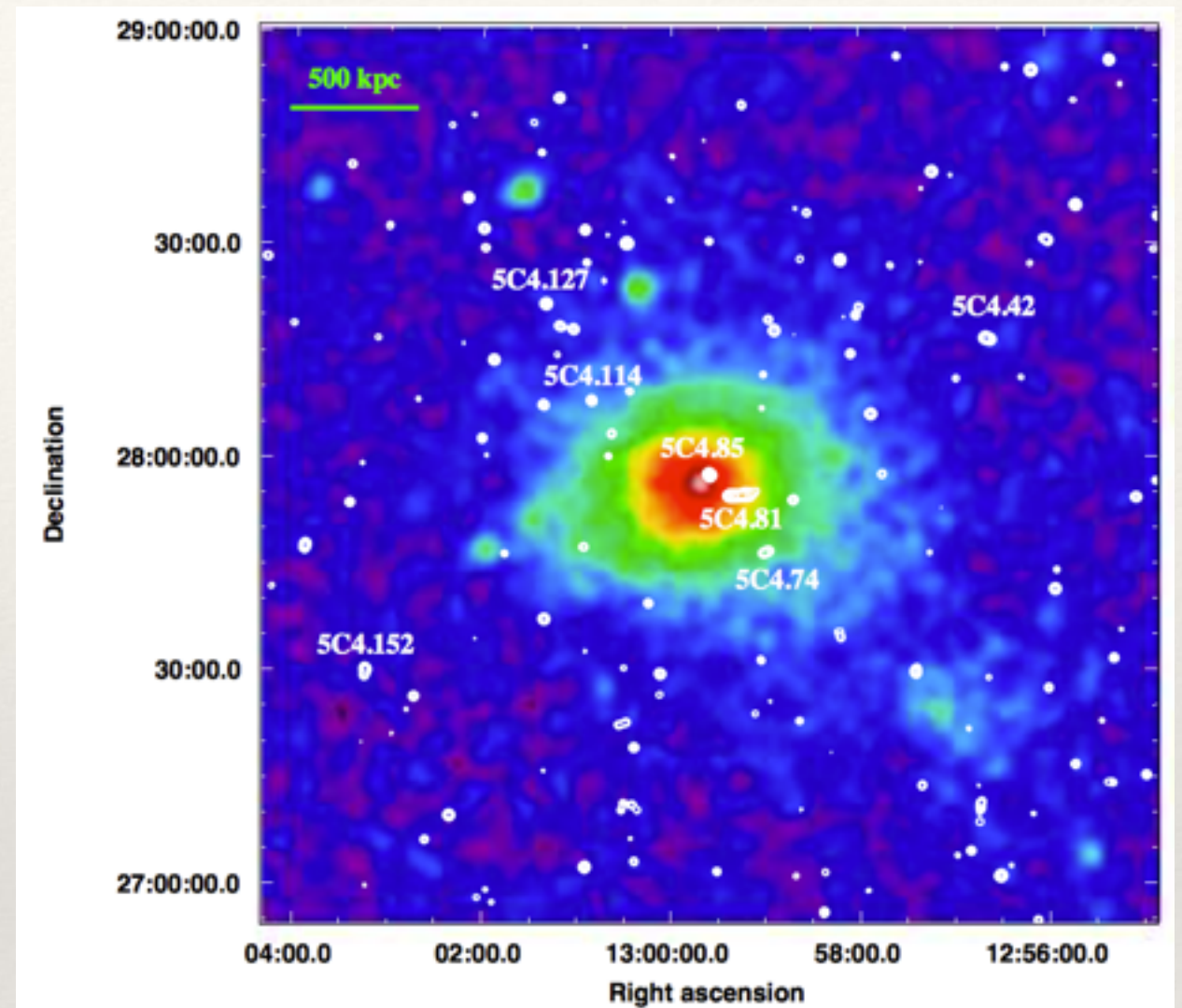
- ❖ Magnetic fields in galaxies and how to measure them
- ❖ Constraints of Inter Galactic Magnetic Fields (IGMF) with gamma ray observations of Active Galactic Nuclei (AGN)
- ❖ Search for extended gamma ray emission around TeV AGN with MAGIC telescopes
- ❖ Upper limits on halo emission
- ❖ Conclusions

Magnetic fields in galaxies and how to measure them

Magnetic fields in galaxies



M51



Coma cluster

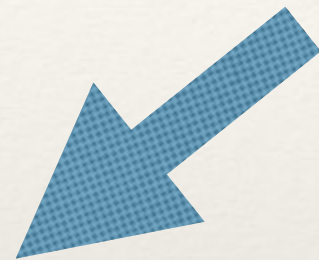
$B \approx 15 \mu\text{G}$

$B \approx 2 \mu\text{G}$

500 kpc

What is the origin of these magnetic fields?

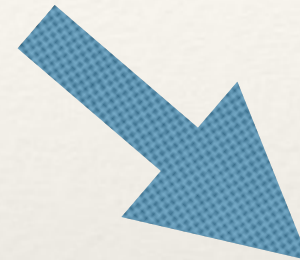
Magnetic fields in galaxies and clusters of galaxies



- galactic dynamo:

$$\frac{\partial \bar{\mathbf{B}}}{\partial t} = \nabla \times (\mathbf{U} \times \bar{\mathbf{B}}) + \nabla \times (\langle \delta \mathbf{V} \times \delta \mathbf{B} \rangle)$$

$$B_0 \simeq 10^{-14} \text{ G}$$



- adiabatic compression:

$$B_{prim,0} = B_{gal} \left(\frac{\rho_{cosmic}}{\rho_{gal}} \right)^{2/3}$$

$$B_{prim,0} \simeq 10^{-10} \text{ G}$$

- ❖ Both hypotheses assume a pre-existing magnetic field!

On the nature of the seed fields

- ❖ The nature of the seed fields is largely unknown. From theoretical point of view two main hypothesis on their origin:
 - ❖ *the astrophysical origin*
 - ❖ *the cosmological origin*

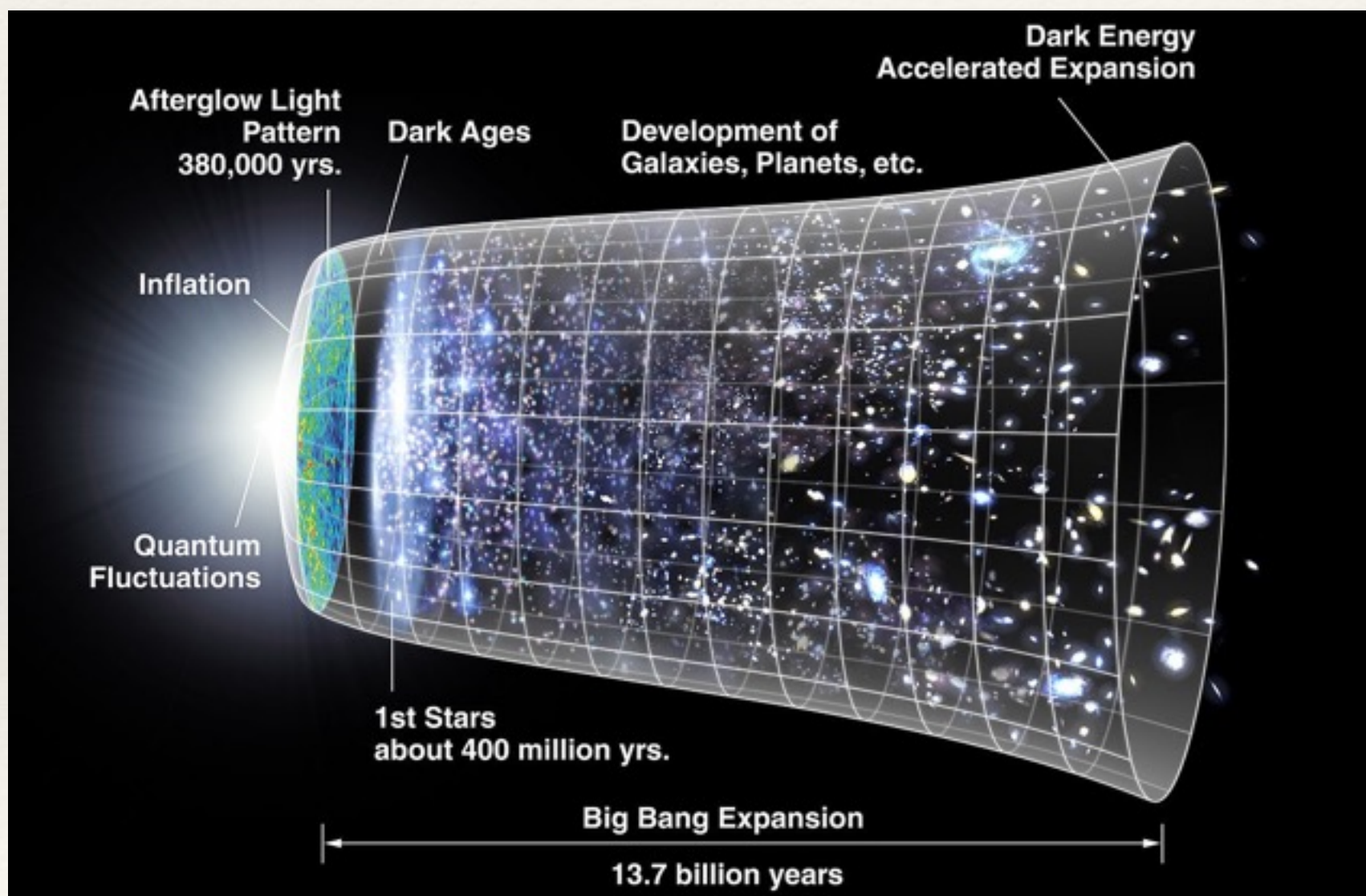
Observationally we need the measurements of the magnetic fields in the Intergalactic medium

The Astrophysical origin

- ❖ Basic idea: difference in mobility between electrons and protons.
 - ❖ seeds from Radiation-Era Vorticity
 - ❖ Biermann battery effect
 - ❖ Magnetic field produced by the stars
 - ❖ Active Galactic Nuclei

The Cosmological Origin

❖ Four main transition phases:



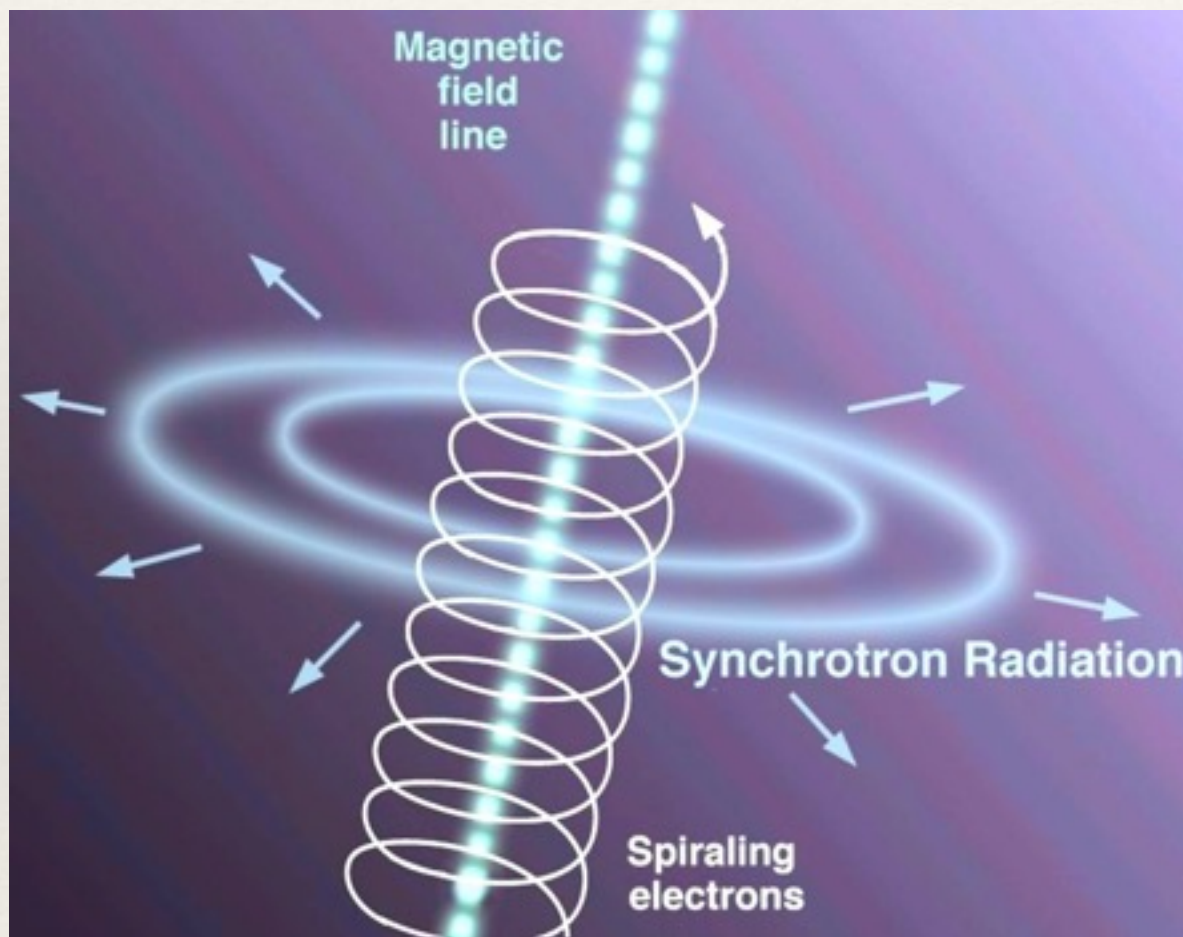
1. inflation
2. electroweak
3. QCD
4. recombination

The tools to measure magnetic fields

- ❖ Standard tools to measure magnetic fields:
 - ❖ Synchrotron emission
 - ❖ Faraday rotation
 - ❖ Zeeman splitting

Synchrotron emission

- ❖ It is the emission of relativistic electrons gyrating in a magnetic field



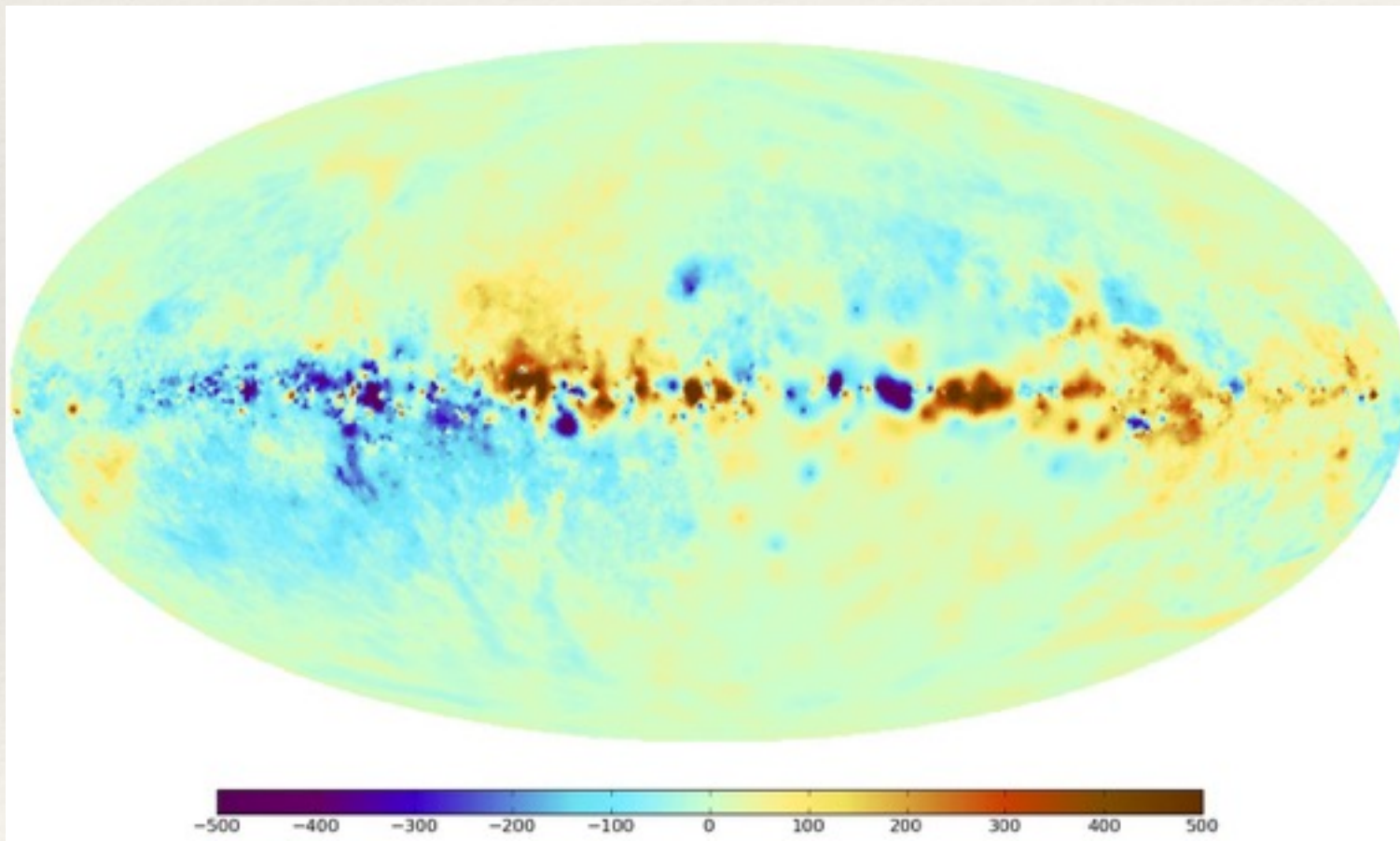
$$j_\nu \propto n_e \nu^{(1-\gamma)/2} B_\perp^{(1+\gamma)/2}$$

- ❖ Assuming equipartition between energy particle and magnetic field we can estimate B

- ❖ The degree of polarization is an indicator of the field's directionality

Faraday Rotation

- ❖ When polarized emission propagates in a region with ionized gas and magnetic field, it undergoes Faraday rotation



Oppermann et al. 2015

$$\phi_{FR} = \frac{e^3 \lambda^2}{2\pi m_e^2 c^4} \int_0^{l_s} n_e(l) B_{\parallel} dl$$



Assuming an electron distribution we can constrain B

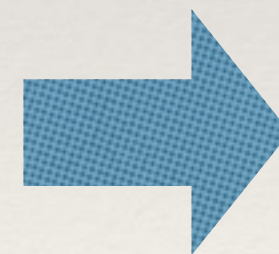
Zeeman Splitting

- ❖ The electronic energy levels of an atom are independent of its angular momentum vector. The presence of a magnetic field removes this degeneracy

$$\frac{\Delta\nu}{\nu} = 1.4g \left(\frac{B}{\mu G} \right) \left(\frac{Hz}{\nu} \right)$$

$$\frac{\Delta\nu}{\nu} = 10^{-9} g \left(\frac{B}{\mu G} \right)$$

21 cm line for
neutral hydrogen



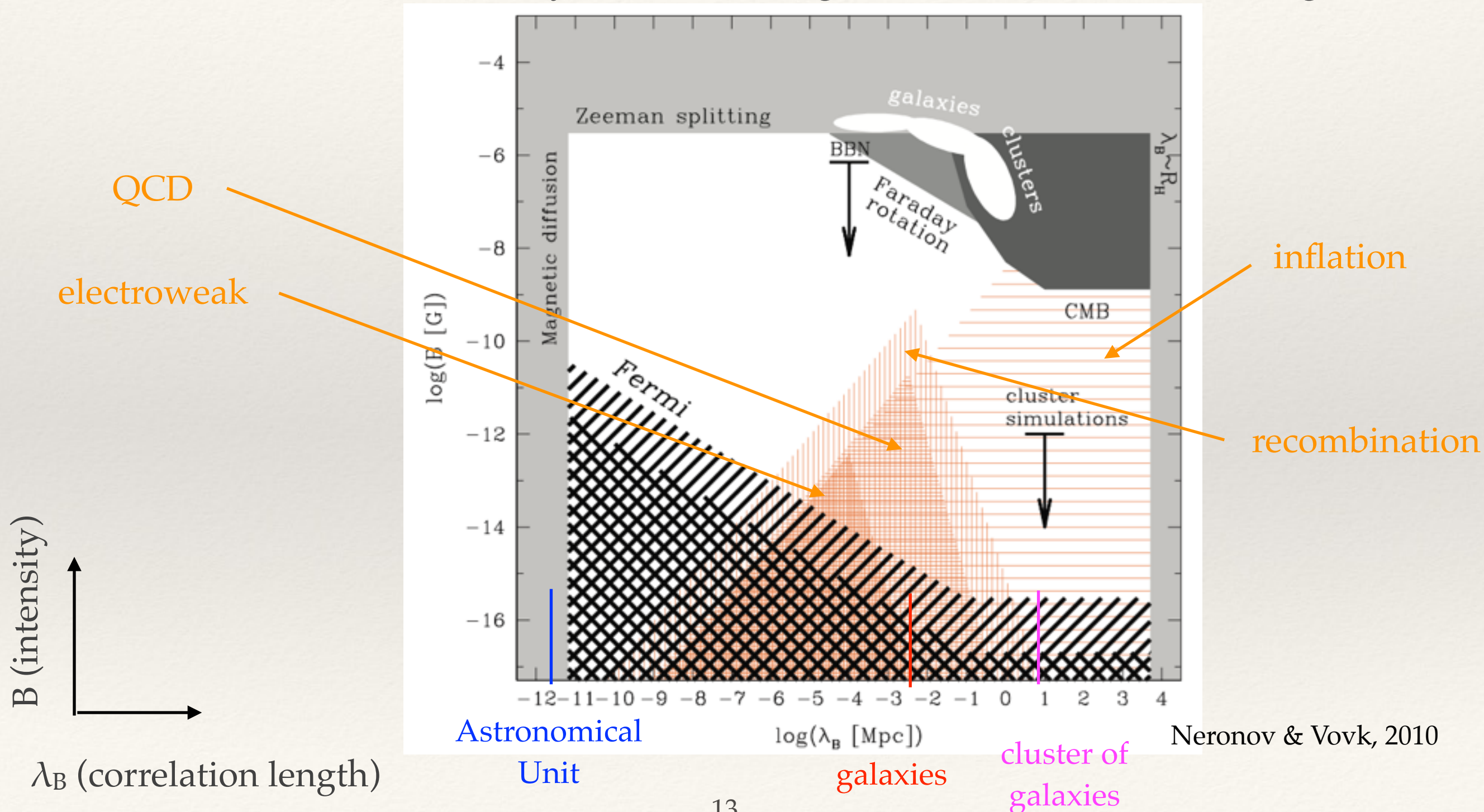
Extremely difficult to
observe

$$\frac{\Delta\nu}{\nu} = 6 \times 10^{-7} \left(\frac{T}{100K} \right)^{1/2}$$

Doppler broadening

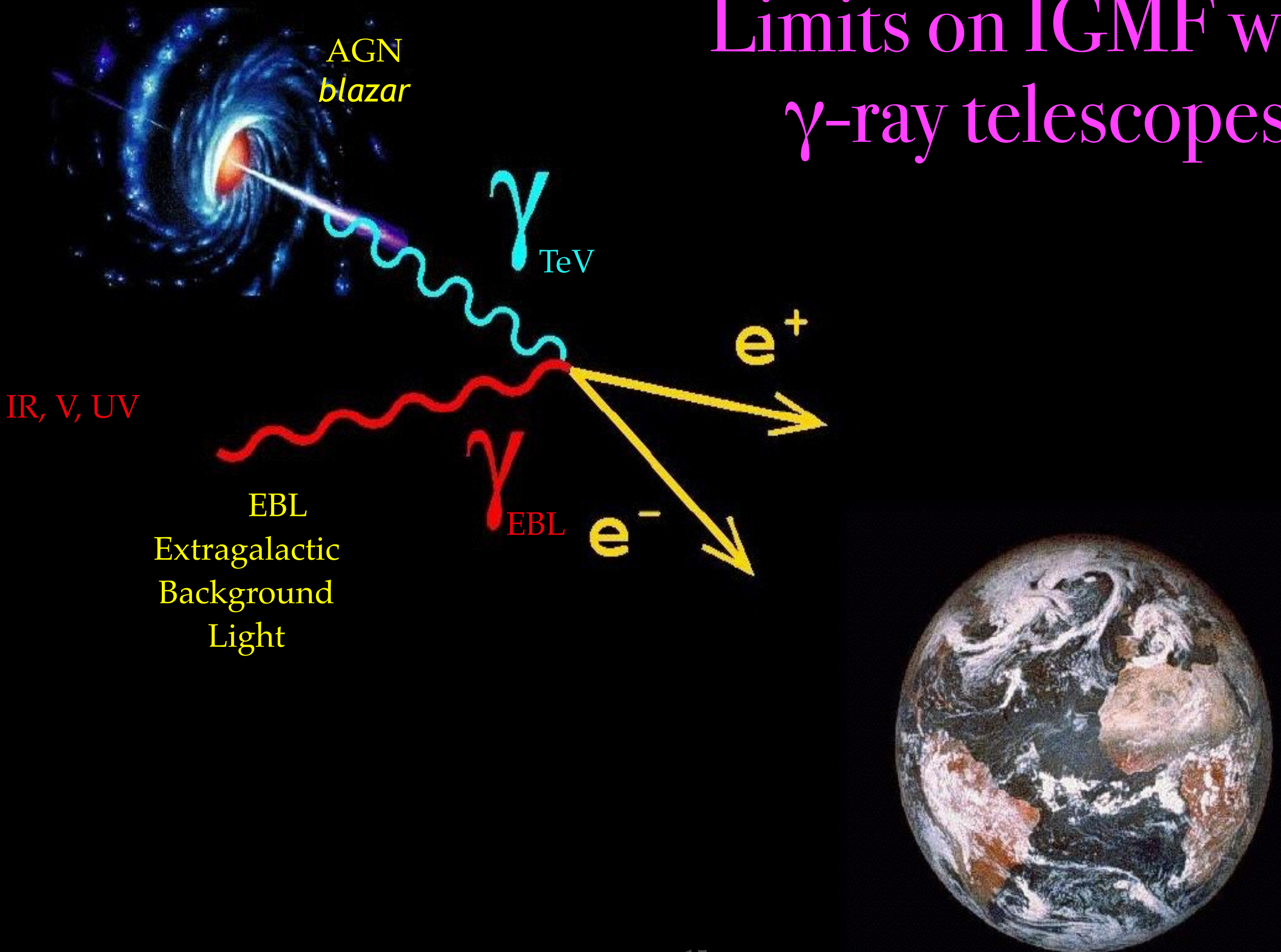
The Inter Galactic Magnetic Field

- ❖ The IGMF is characterized by the field strength B and the correlation length λ_B



Constraints of IGMF with γ -ray observations of AGN

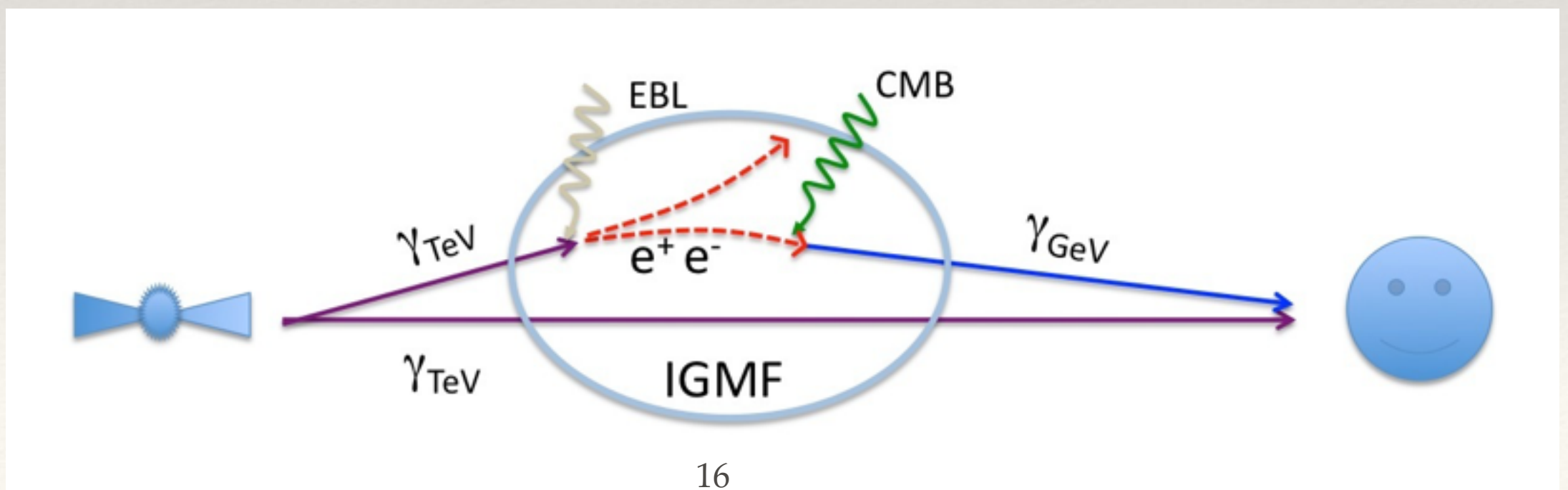
Limits on IGMF with γ -ray telescopes



Limits on IGMF with γ -ray telescopes

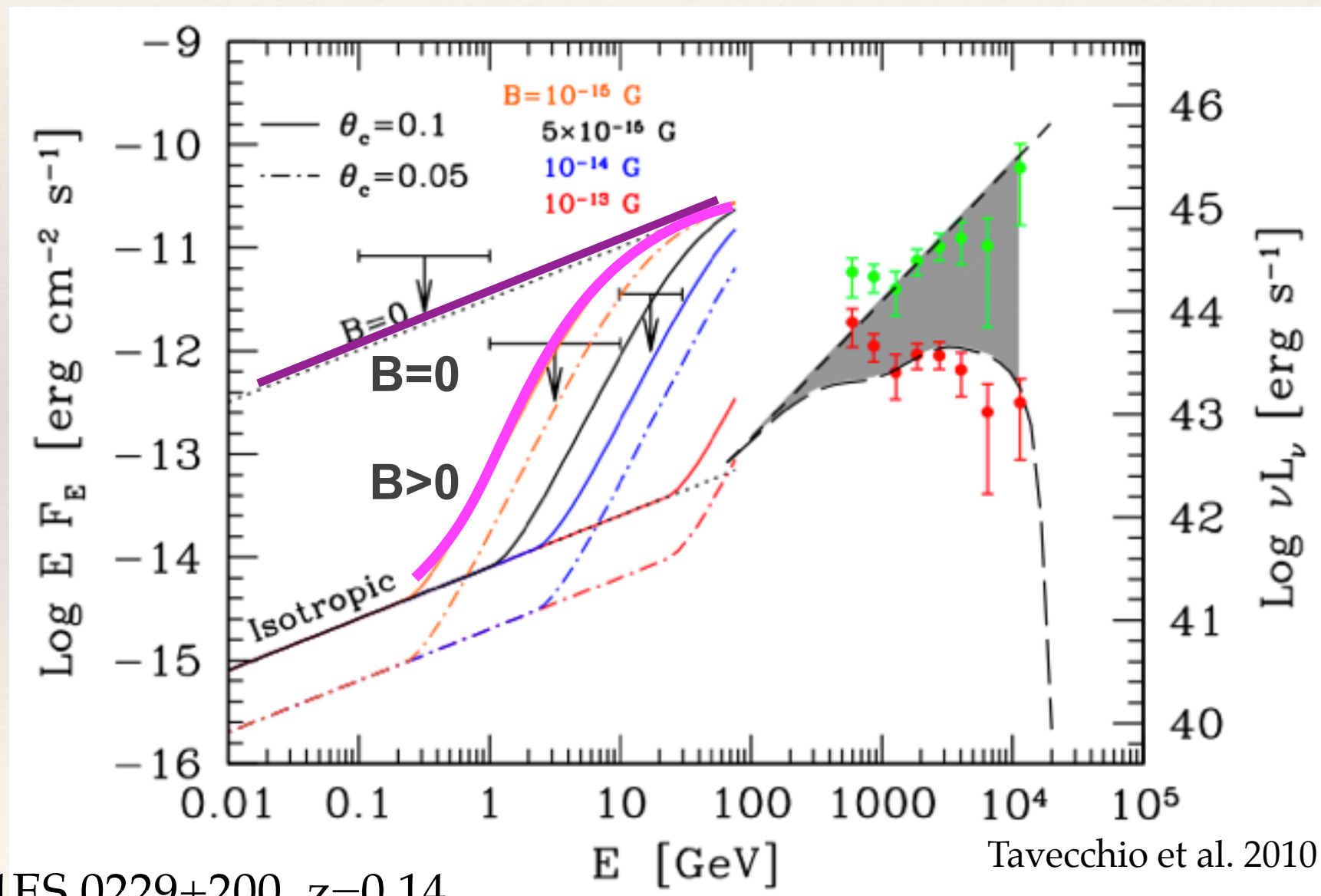
- ❖ Physical process: reprocessing of TeV photons in the GeV band
- ❖ Measurable effects:
 - ✦ spectral features
 - ✦ extended emission

$$E_{\gamma} \simeq 70 \left[\frac{E'_{\gamma 0}}{10 \text{ TeV}} \right]^2 \text{ GeV}$$



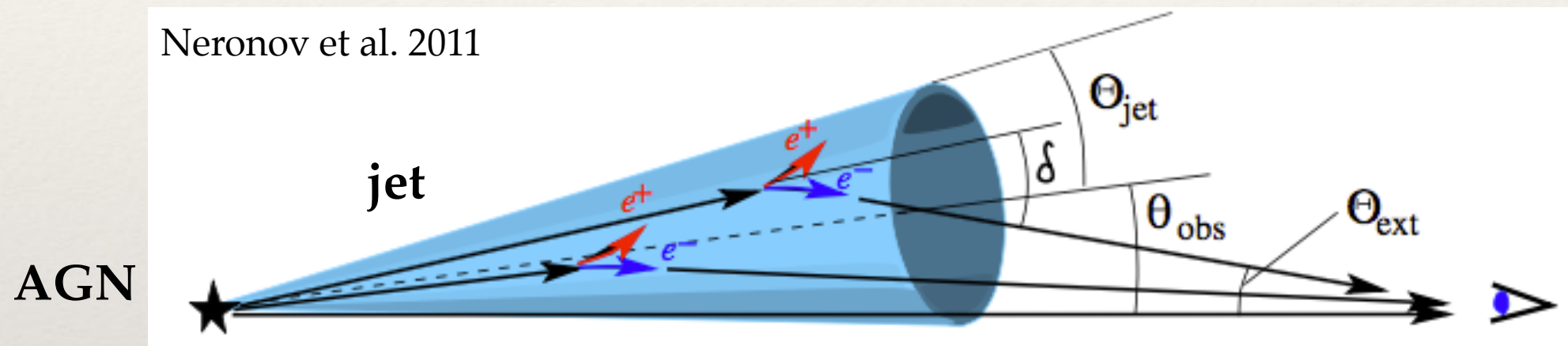
Spectral features

- ❖ Measuring the amount of absorbed flux of a TeV blazar we can predict the the amount of cascade emission. Its suppression depends on the IGMF strength



Extended emission

- ❖ Observable effect: extended emission around the point source. The angular extension grows with increasing IGMF.



- ❖ Two regimes:

- ❖ $\lambda_B \gg D_e$

$$\Theta_{\text{ext}} \simeq 0.5^\circ (1+z)^{-2} \left[\frac{\tau}{10} \right]^{-1} \left[\frac{E_\gamma}{0.1 \text{ TeV}} \right]^{-1} \left[\frac{B_0}{10^{-14} \text{ G}} \right]$$

- ❖ $\lambda_B \ll D_e$

$$\Theta_{\text{ext}} \simeq 0.07^\circ (1+z)^{-1/2} \left[\frac{\tau}{10} \right]^{-1} \left[\frac{E_\gamma}{0.1 \text{ TeV}} \right]^{-3/4} \left[\frac{B_0}{10^{-14} \text{ G}} \right] \left[\frac{\lambda_{B0}}{1 \text{ kpc}} \right]^{1/2}$$

Our project to search for extended γ -ray emission

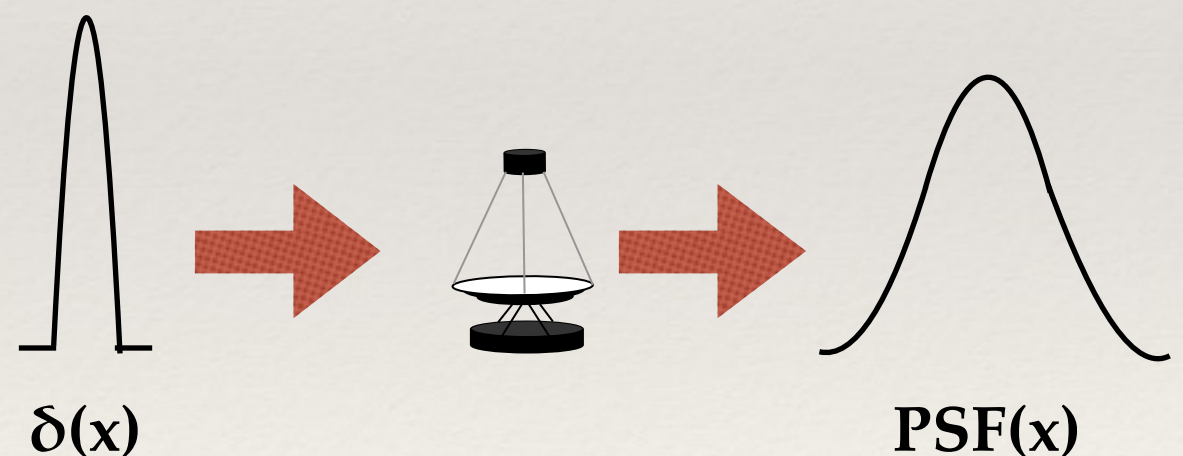
- ❖ Through the study of the emission profiles of several TeV AGN we looked for halo emission around 100 GeV.
- ❖ We used the VHE data of MAGIC telescopes to derive an analytical description of the Point Spread Function (PSF)
- ❖ Four steps have been considered:

- ❖ the background

- ❖ model of PSF

- ❖ stability of the PSF

- ❖ comparison of PSF with profiles of AGN



Search for extended γ -ray emission around TeV AGN with MAGIC telescopes

Imaging Air Cherenkov Telescopes

IACT

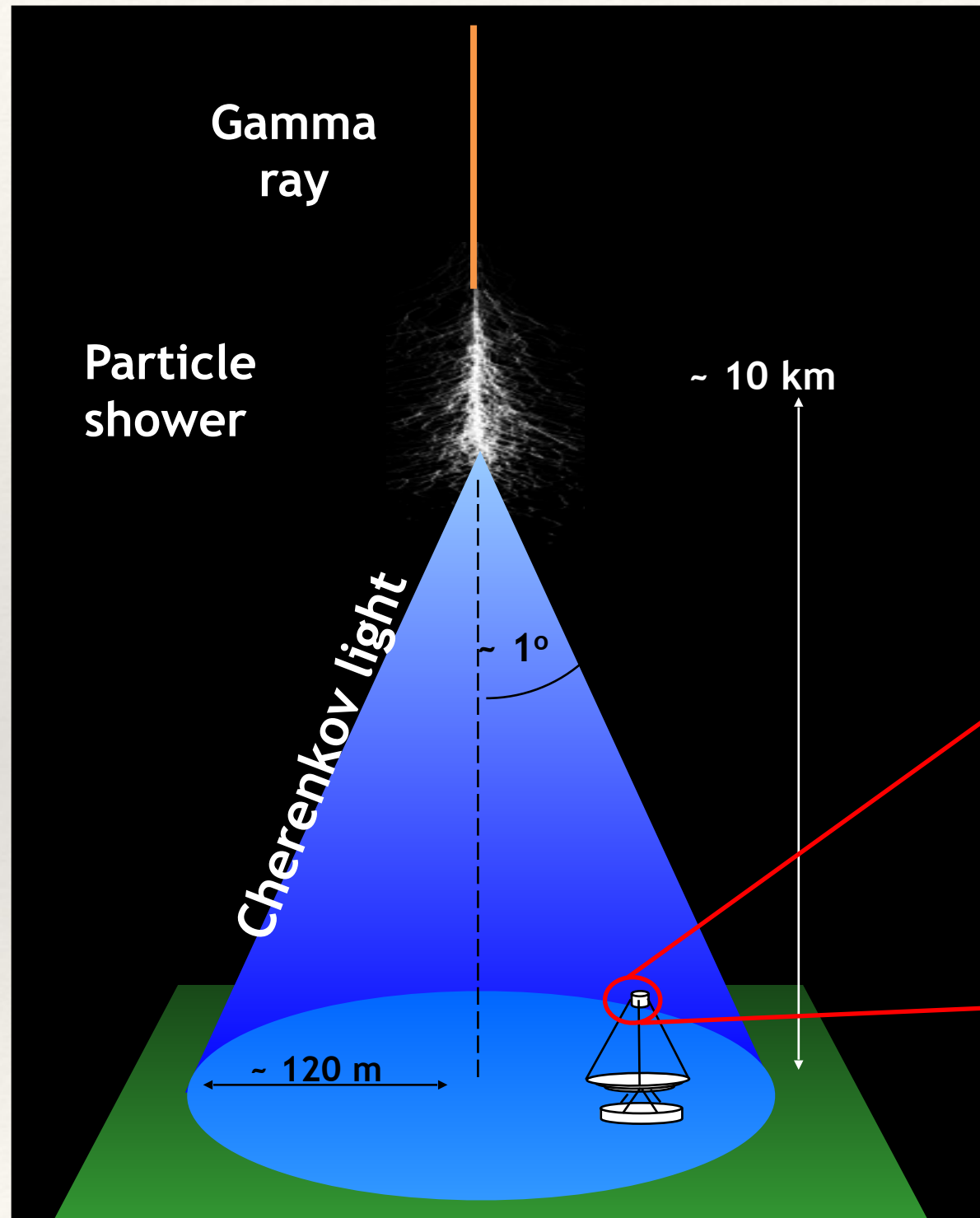
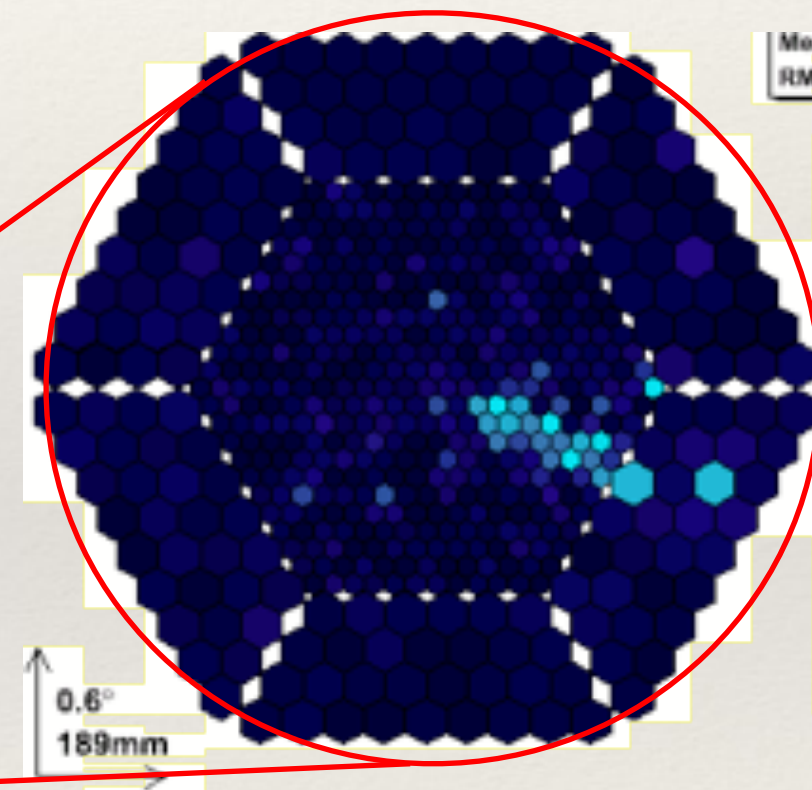


Image of Cherenkov flash on the camera



- ❖ gamma/hadron discrimination
- ❖ reconstruction of the energy and direction of the primary photon

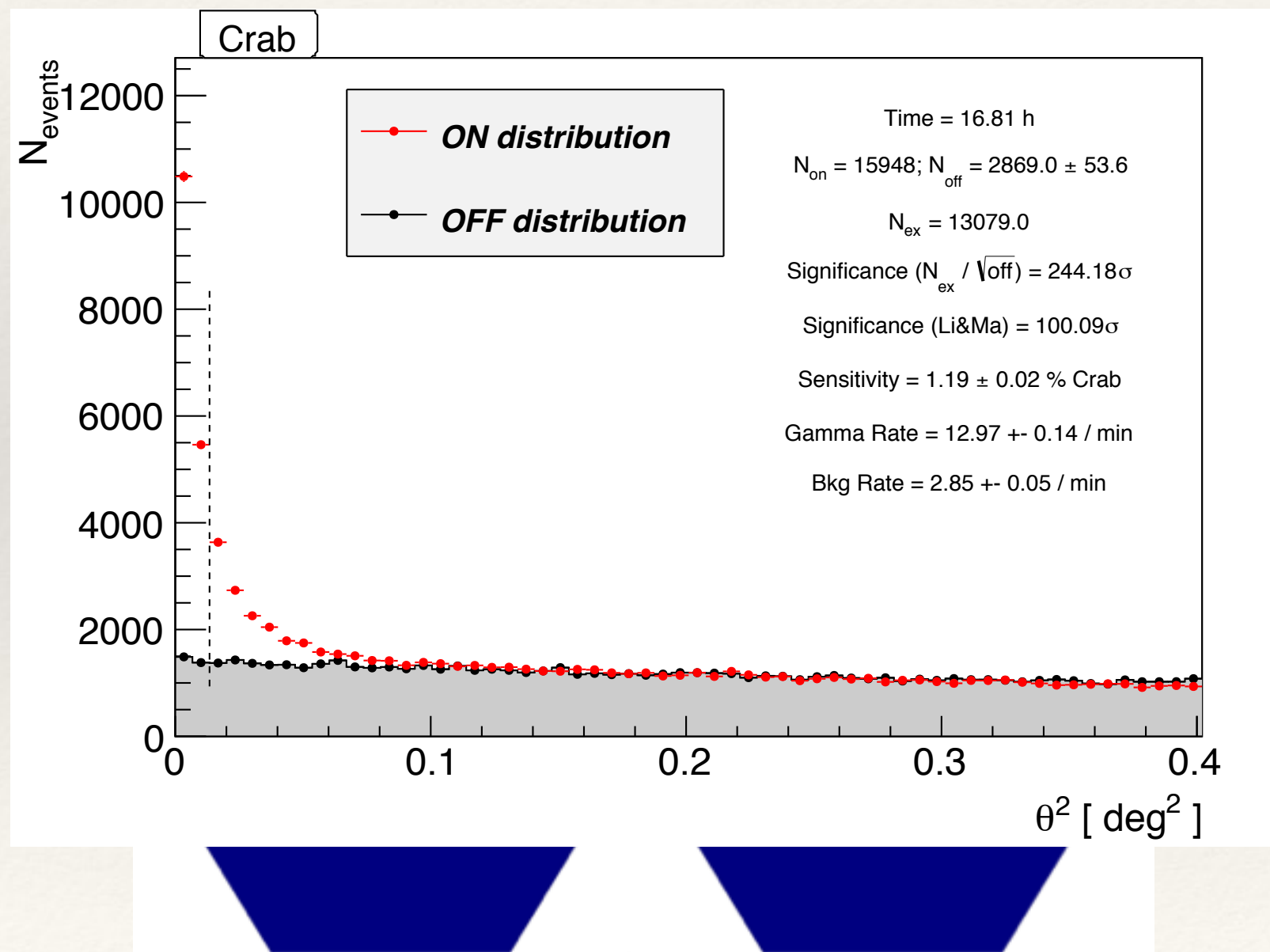
The MAGIC telescopes

- Energy threshold ~ 50 GeV
- FOV 3.5°
- Energy Resolution $\sim 16\%$ ($E > 300$ GeV)
- *Angular Resolution $\sim 0.06^\circ$ ($E > 300$ GeV)*
- Sensitivity (5σ in 50 hours) $\sim 0.8\%$ Crab Nebula flux (> 250 GeV)



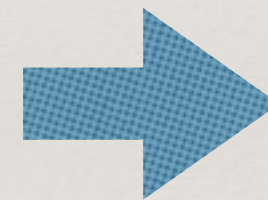
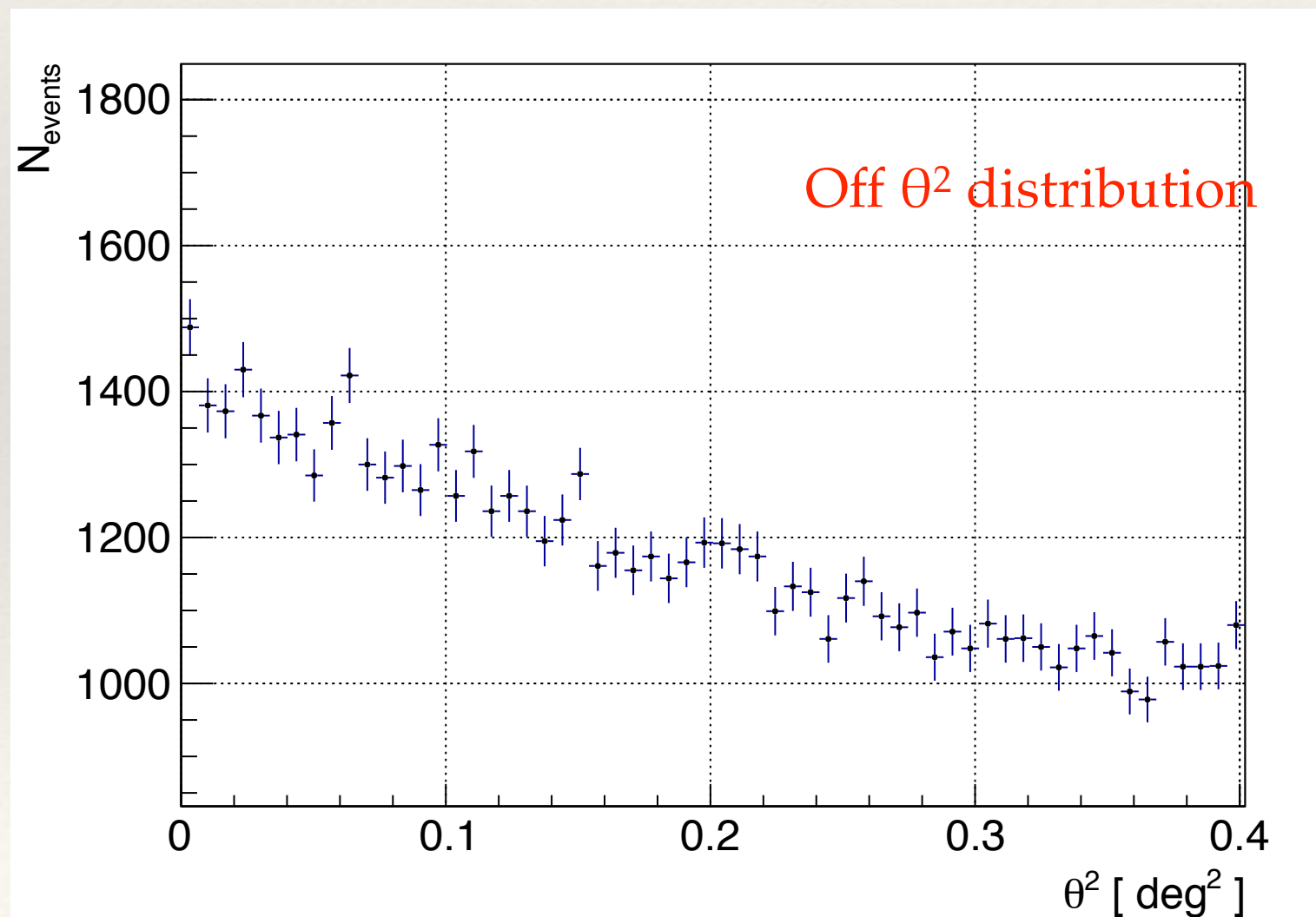
γ -ray direction reconstruction

- ❖ θ : angular distance between real and reconstructed position of the source
- ❖ Distribution after selection of gamma-like events



Study of MAGIC PSF

- ❖ For the proper characterization of the PSF the acceptance on the camera must be evaluated. The key point is the **off-axis response**.



the shape of the background is not flat, correction for the camera acceptance

Fit of MAGIC PSF

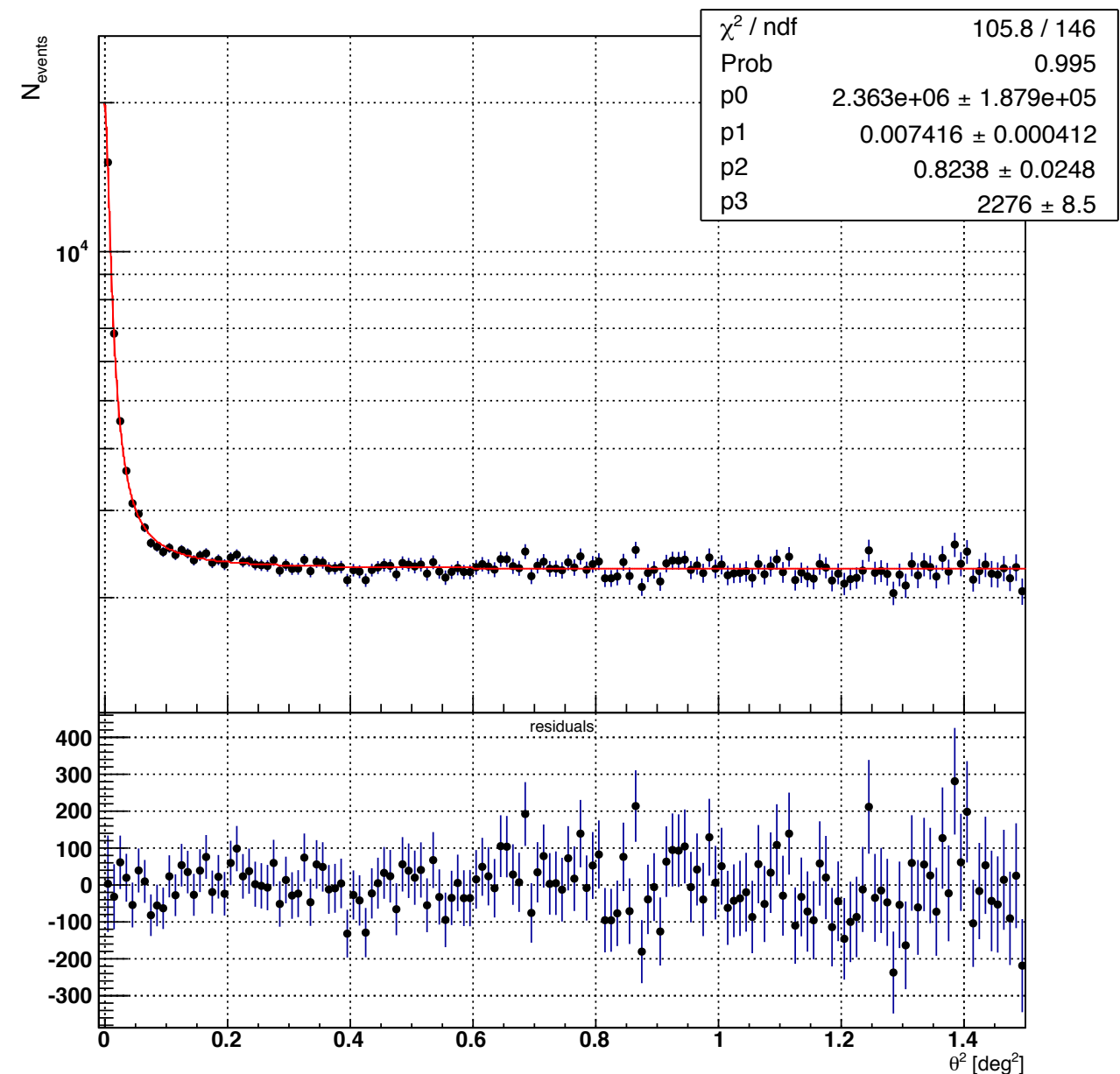
- ❖ We used the King function to fit the MAGIC PSF

$$f(\theta^2) = p0 * p1 * \left[1 + \left(\frac{\theta^2}{p1} \right)^2 \right]^{-p2} + p3$$

King function

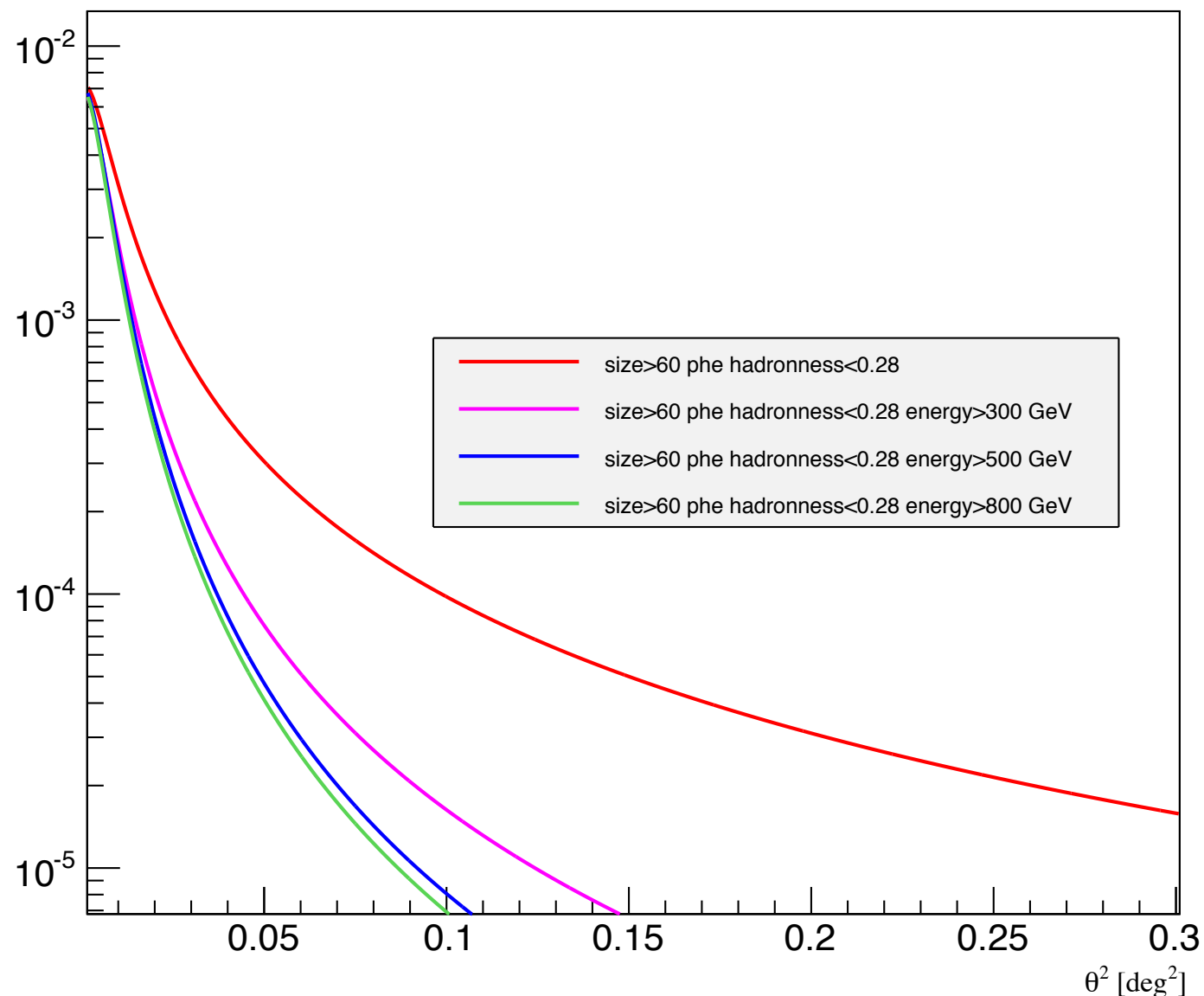
corrected background

This function is able to describe the data also at high values of θ^2



The dependence on energy

- ❖ We studied the behavior of the PSF adding an energy cut

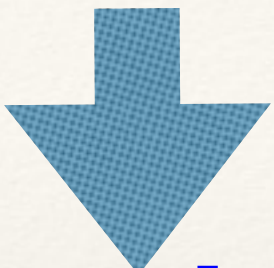


Increasing the energy cut
the King function becomes
narrower

Search for extended γ -ray emission around TeV AGN

- ❖ We compared the PSF with the emission profiles of five AGN:
 - ❖ Markarian 421 (10 hours, $z=0.03$)
 - ❖ 1ES 1011+496 (11 hours, $z=0.212$)
 - ❖ PKS 1424+240 (28 hours, $z\geq 0.6$)
 - ❖ RX J1136.5+6737 (33 hours, $z=0.13$)
 - ❖ 1ES 0229+200 (40 hours, $z=0.14$)

Comparison of the PSF with AGN



PSF model: $f(\theta^2) = p_0 * p_1 * \left[1 + \left(\frac{\theta^2}{p_1} \right)^2 \right]^{-p_2} + p_3$

two independent methods

❖ **Method 1:**

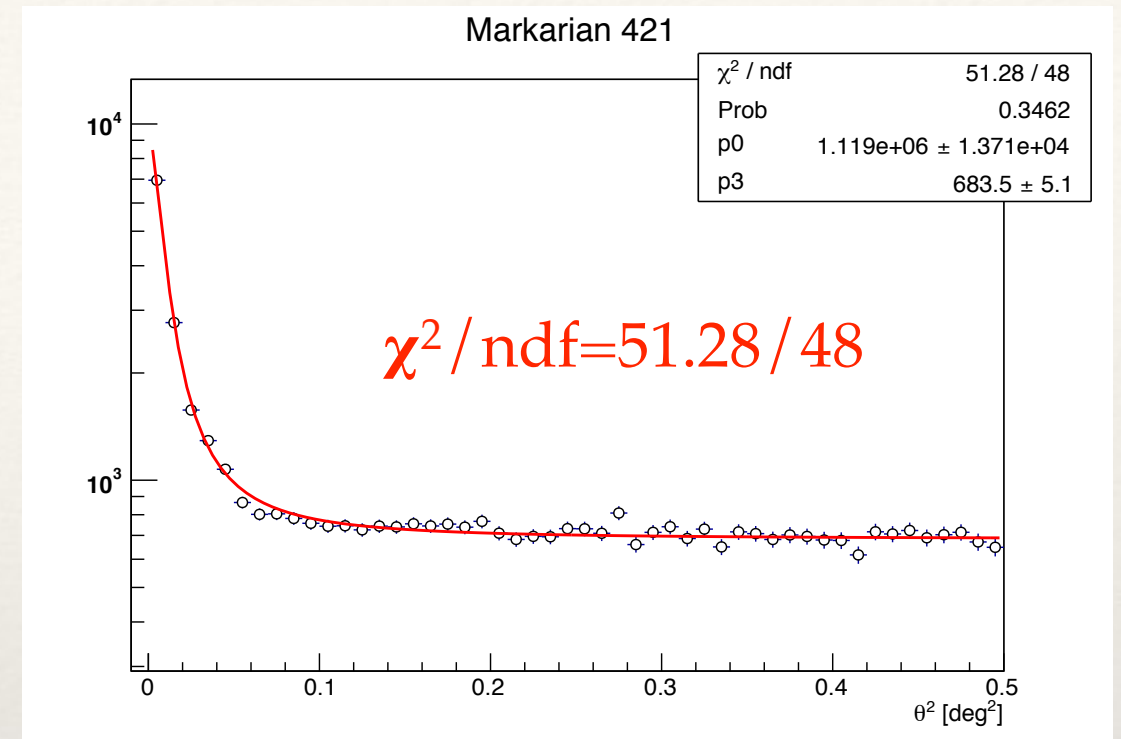
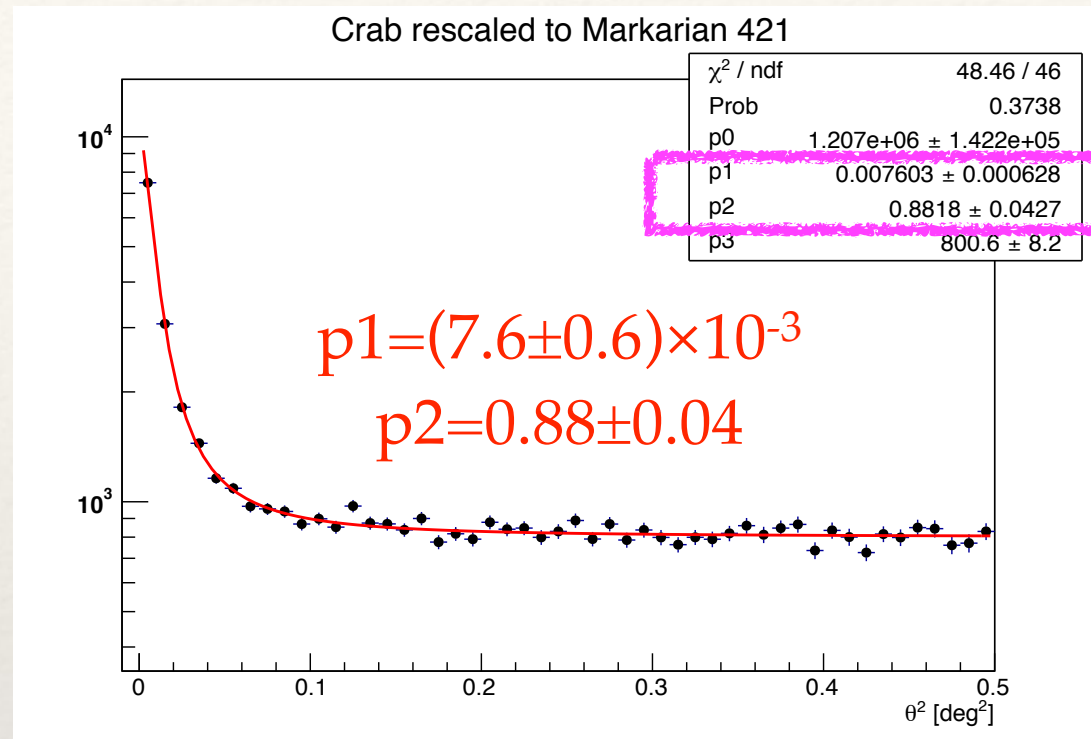
- ★ Crab: fit with 4 free parameters
- ★ AGN: p1 and p2 fixed to Crab

❖ **Method 2:**

- ★ Crab: fit with 4 free parameters
- ★ AGN: fit with 4 free parameters
- ★ confidence contours on p1 and p2 for both Crab and AGN

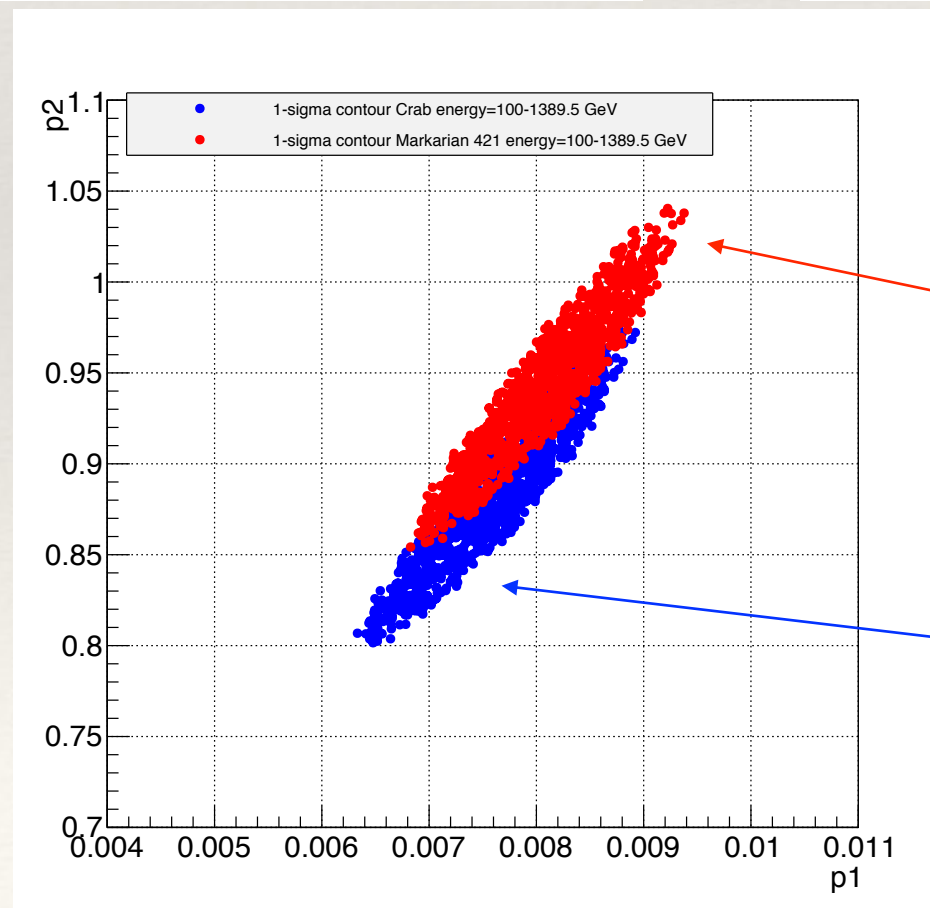
Markarian 421

Method 1:



Method 2:

1σ contour plots



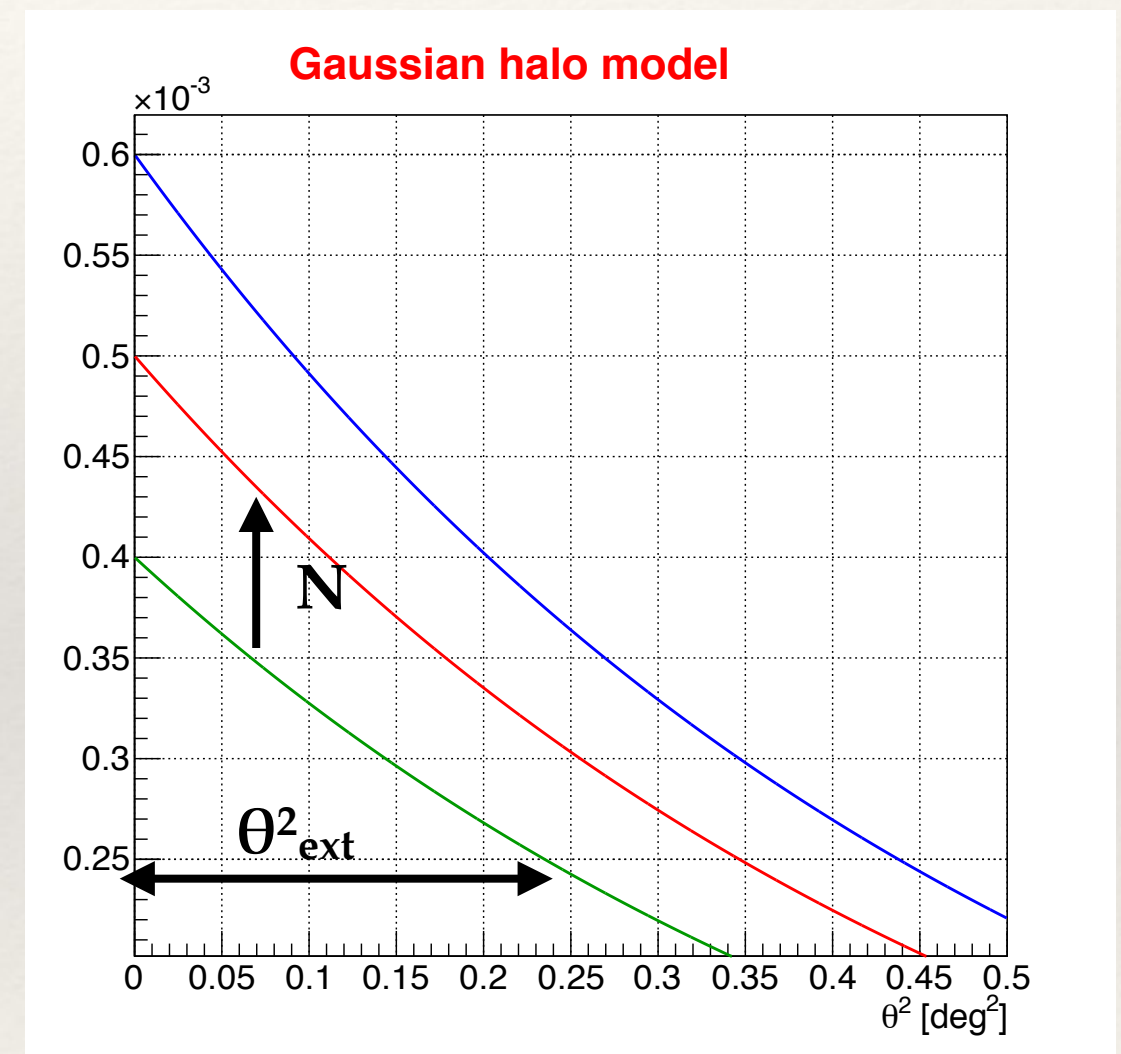
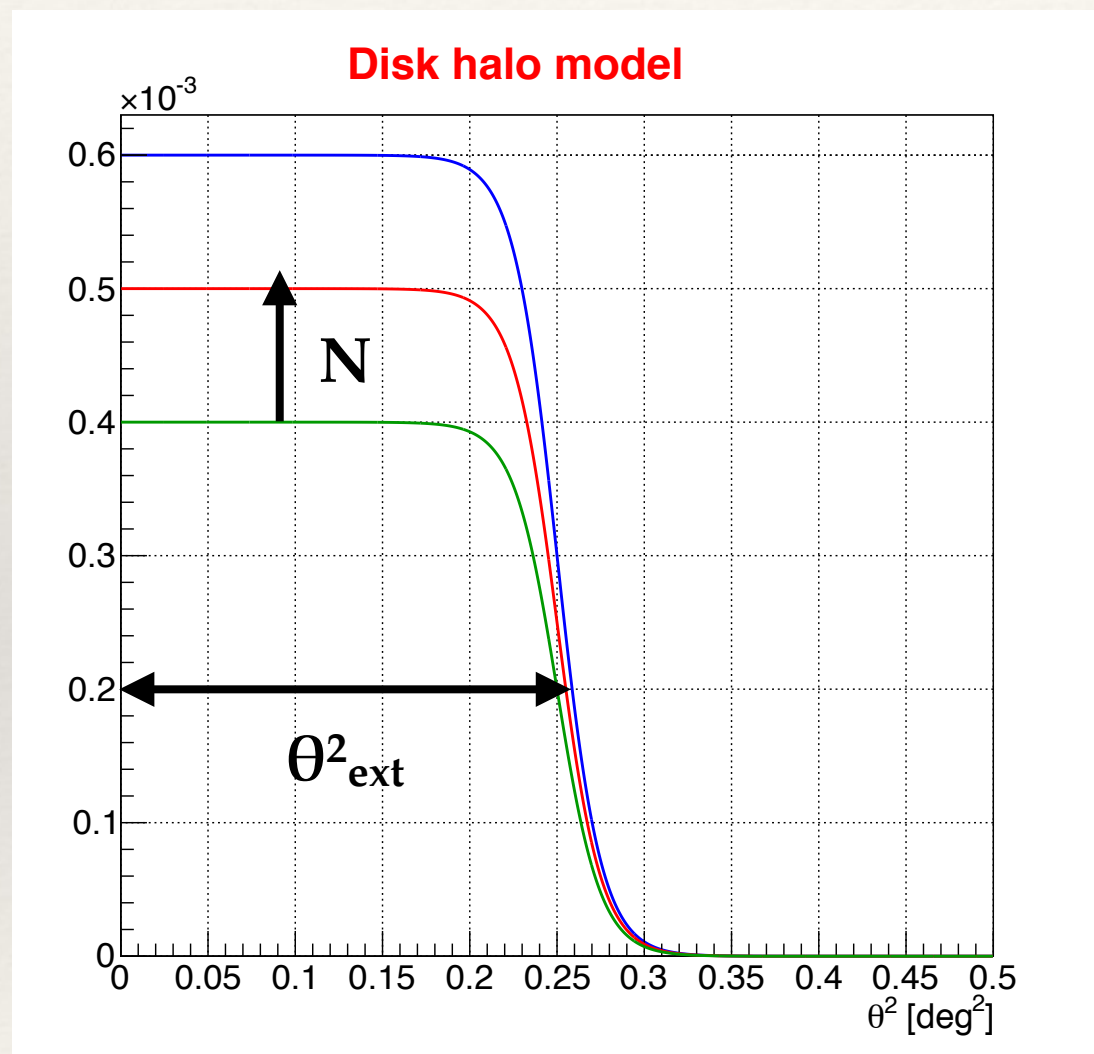
Mrk421

Crab

Upper limits on halo emission

Upper limits on halo

- ❖ We added a halo model to the PSF to compute upper limits



$$\text{halo}_{\text{disk}}(N, \theta^2_{\text{ext}}; \theta^2) = \frac{N}{1 + e^{k*(\theta^2 - \theta^2_{\text{ext}})}}$$

$$\text{halo}_{\text{gaus}}(N, \theta^2_{\text{ext}}; \theta^2) = N e^{-\frac{\theta^2}{2\theta^2_{\text{ext}}}}$$

Upper limits on halo

- ❖ The total model is:

$$f(\theta^2) = halo(N, \theta_{ext}^2; \theta^2) + p0 * p1 * \left[1 + \left(\frac{\theta^2}{p1} \right)^2 \right]^{-p2} + p3$$

- ❖ For a given halo extension (θ_{ext}^2) we increased the normalization (N) until we got a worsening of the fit correspondent to 2σ level

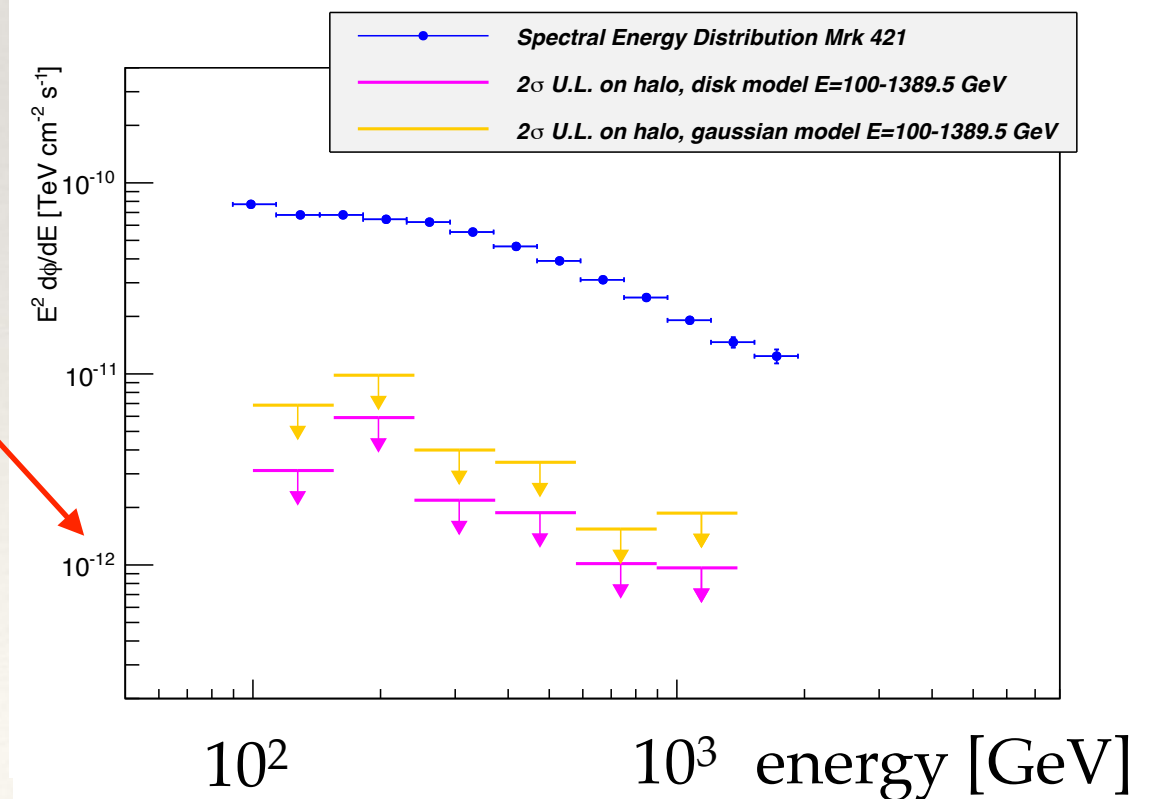
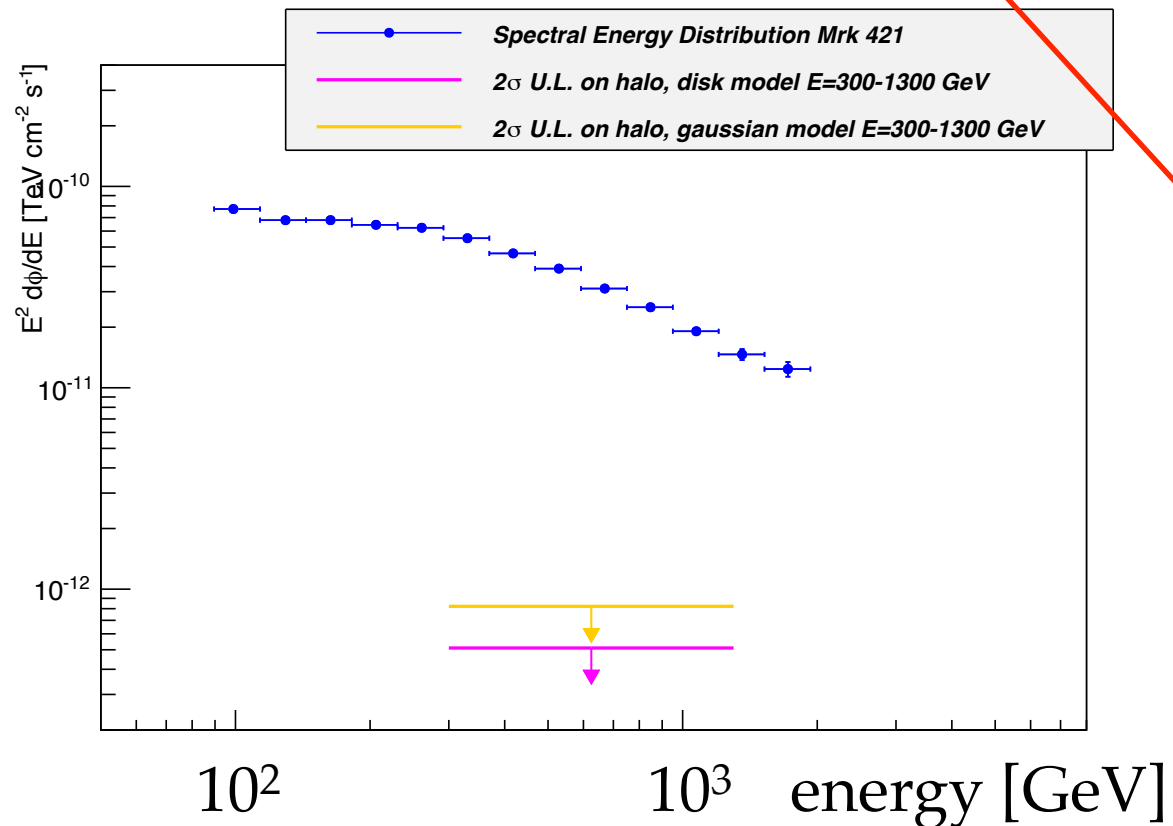
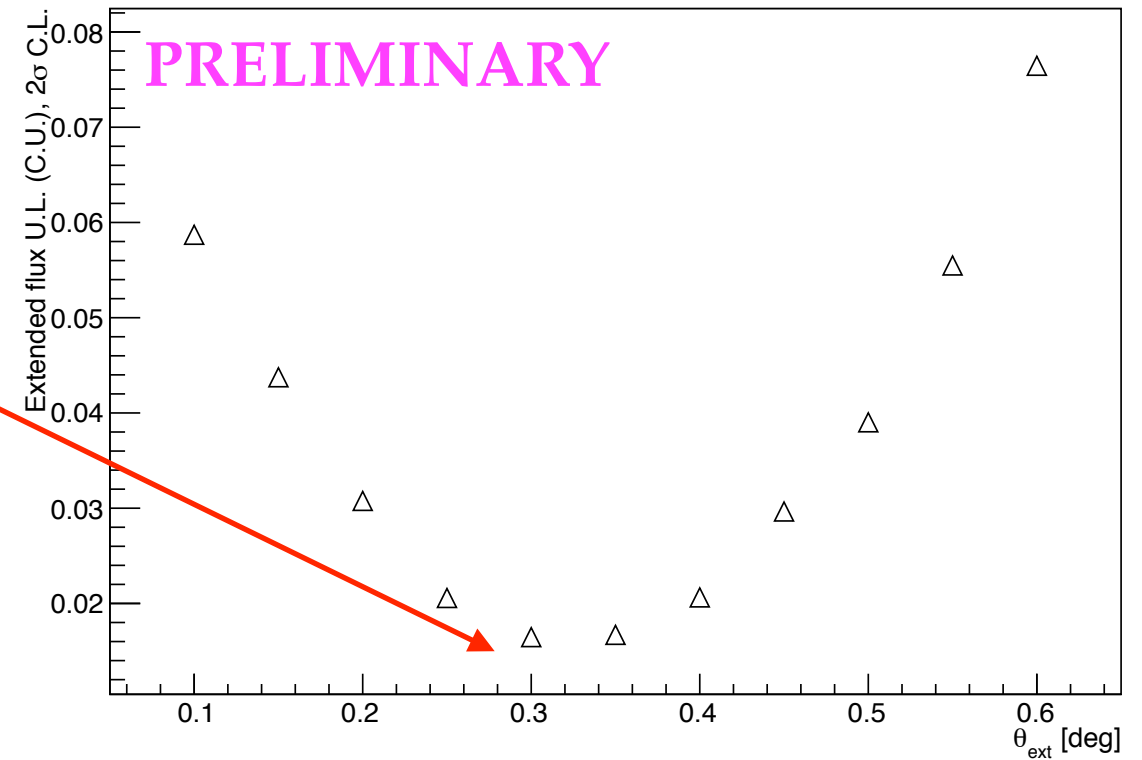
$$U.L. (2\sigma source units) = \frac{\int_0^\infty halo(N, \theta_{ext}^2; \theta^2) d\theta^2}{\int_0^\infty p0 * p1 * \left[1 + \left(\frac{\theta^2}{p1} \right)^2 \right]^{-p2} d\theta^2}$$

Markarian 421

Most stringent upper
limit for $\theta_{\text{ext}}=0.3^\circ$. U.L.
(C.U.)=1.6%

Upper limits performed in 6
energy bins in the range
100-1389.5 GeV

Upper limits on halo (disk model) for Markarian 421, E=300-1300 GeV

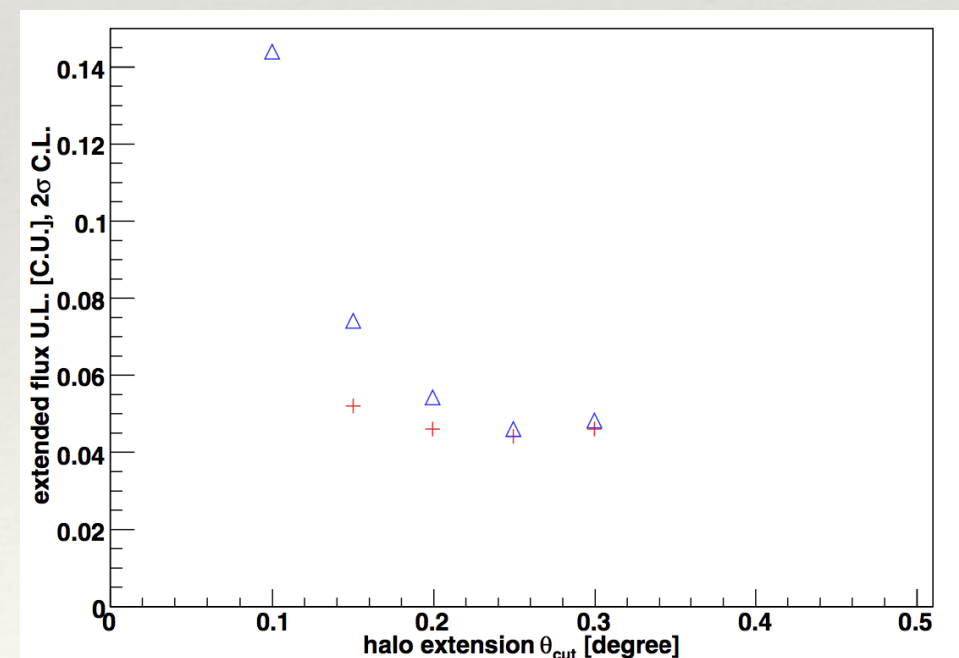
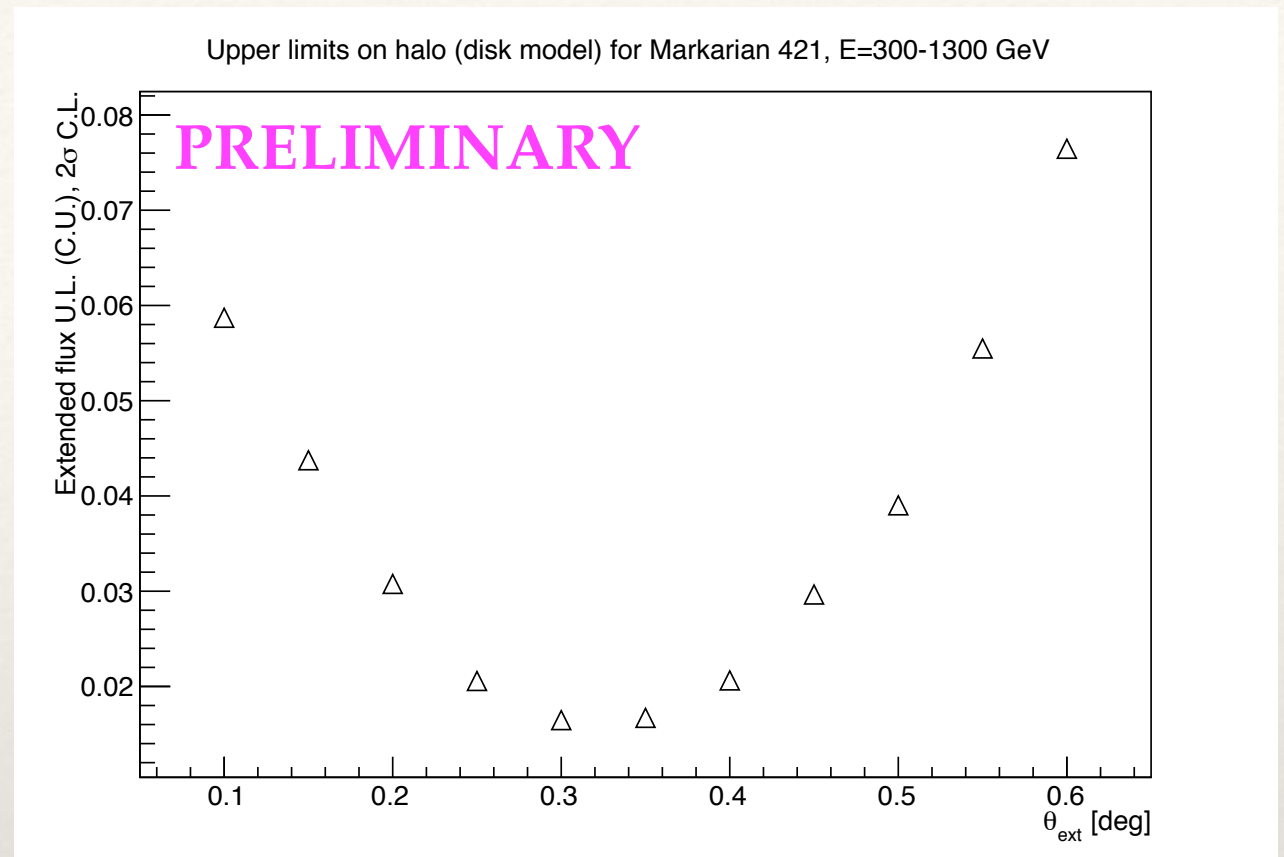


Markarian 421

- ❖ Our procedure provided an upper limit which is more than 3 times better than previous measurement

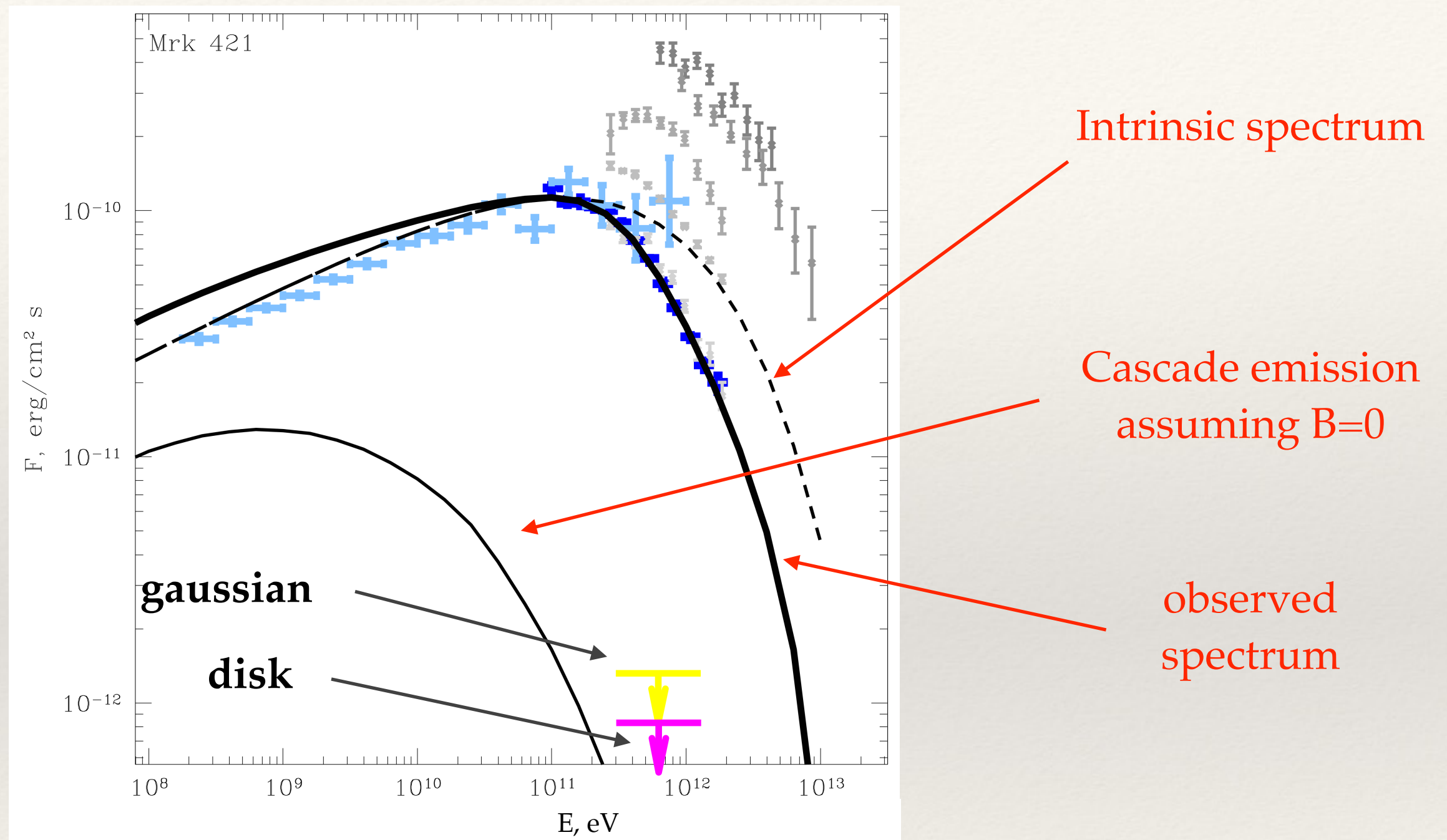
Most stringent upper
limit for $\theta_{\text{ext}}=0.3^\circ$. U.L.
(C.U.)=1.6%

Most stringent upper
limit for $\theta_{\text{ext}}=0.25^\circ$.
U.L.(C.U.)=5%



Aleksić et al. 2010

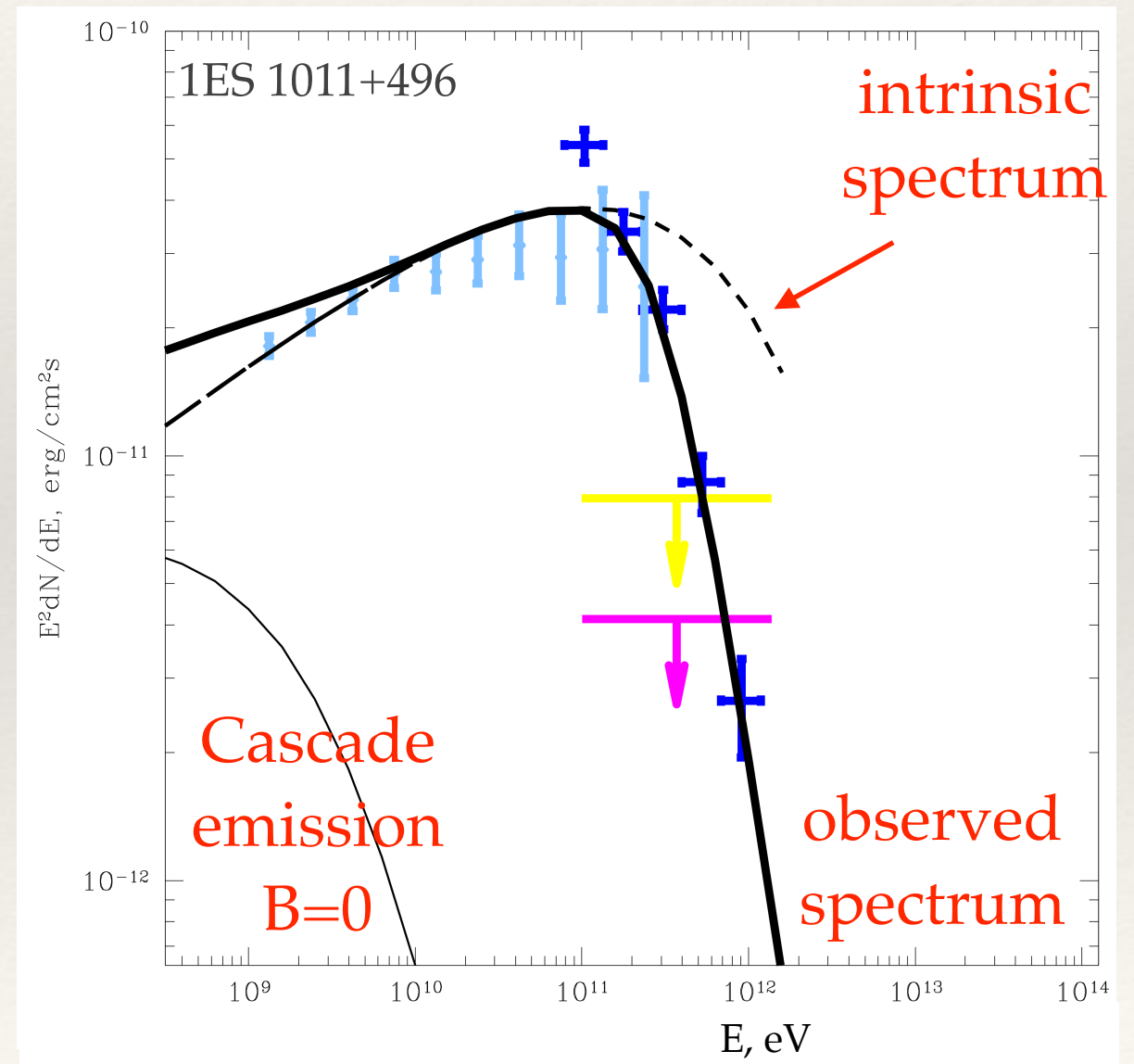
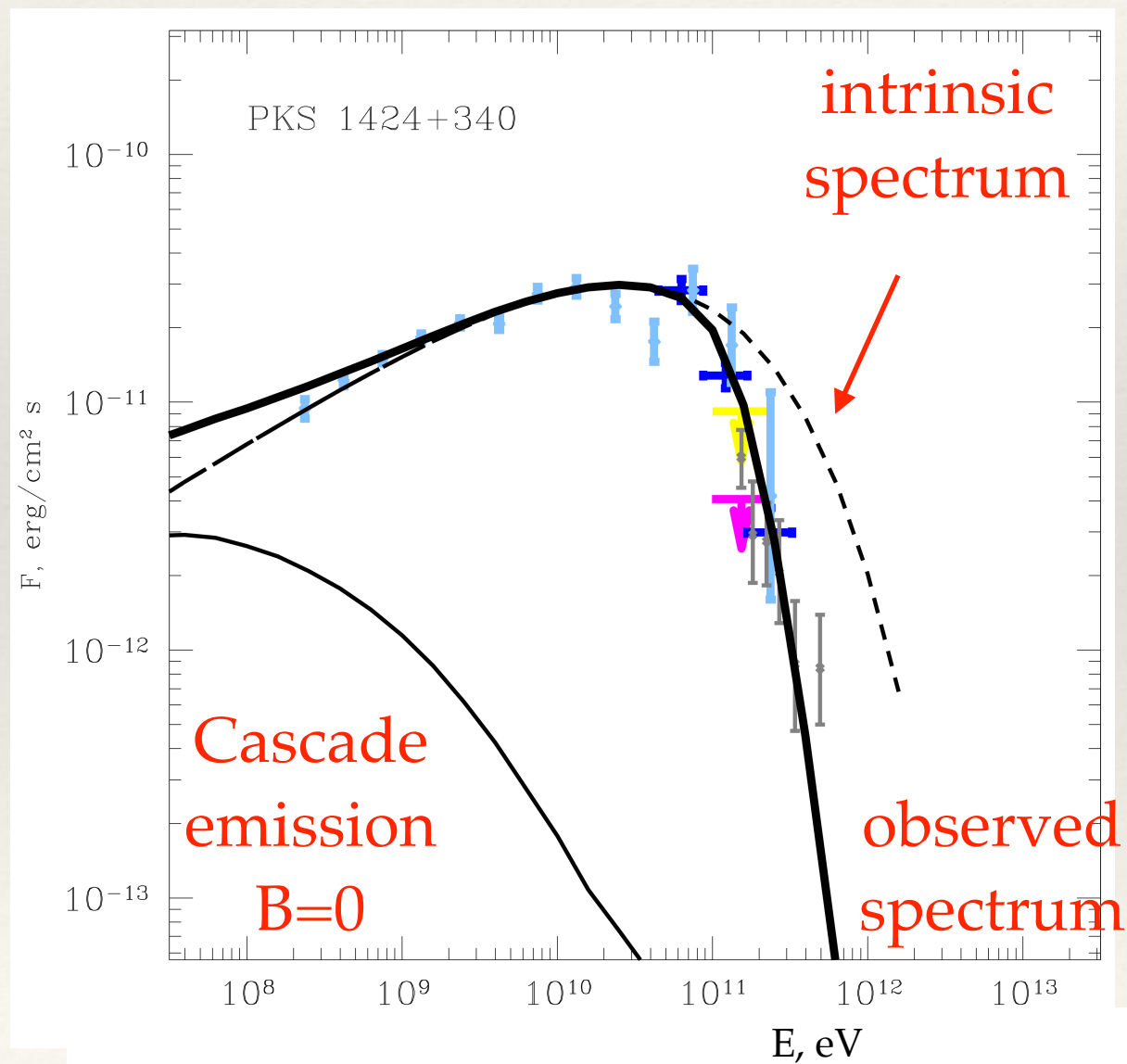
Markarian 421: spectral emission of cascade



- ❖ Due to the vicinity of the source and the cut-off in the intrinsic spectrum the contribution of the cascade emission to the total spectrum is too low

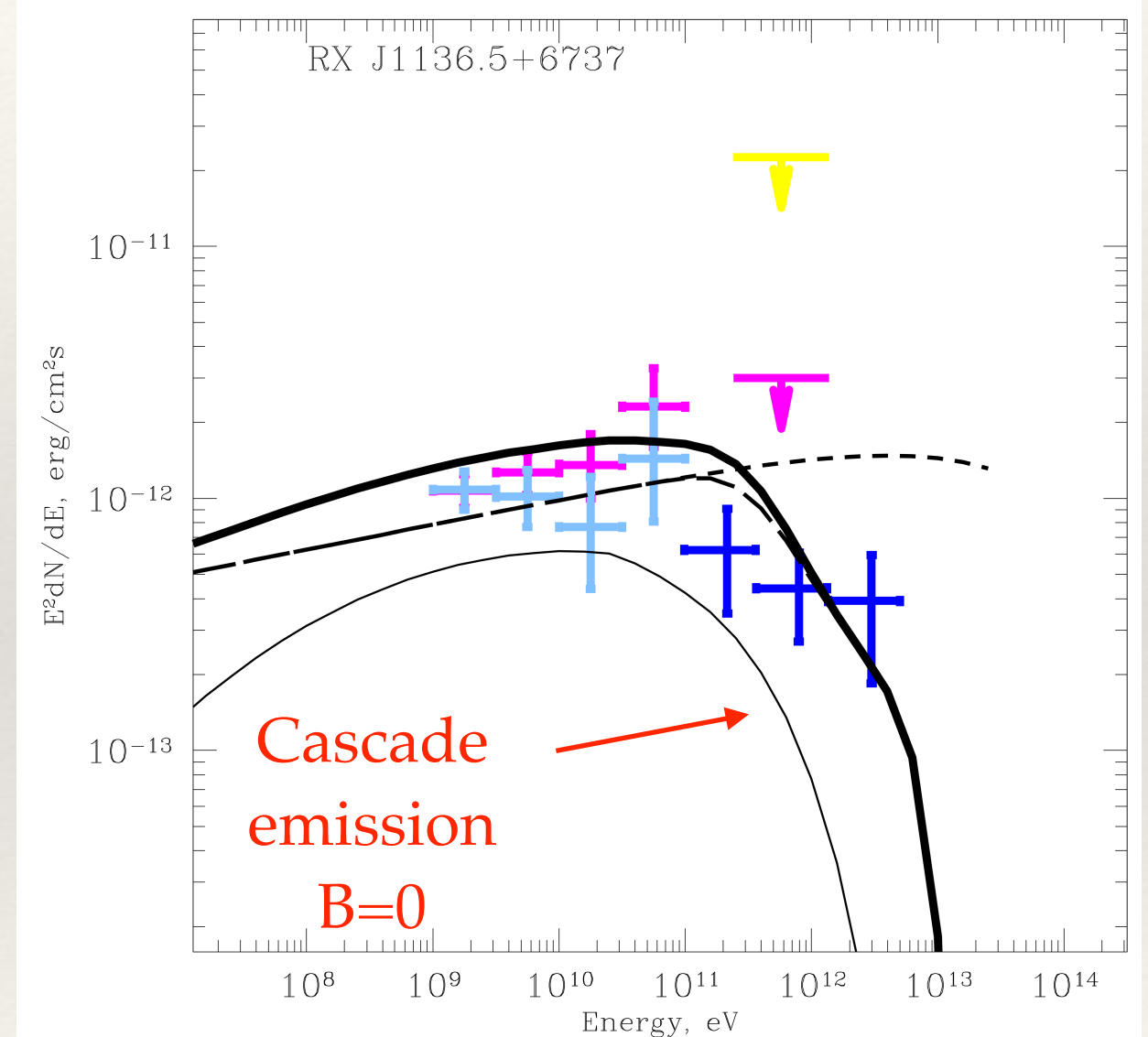
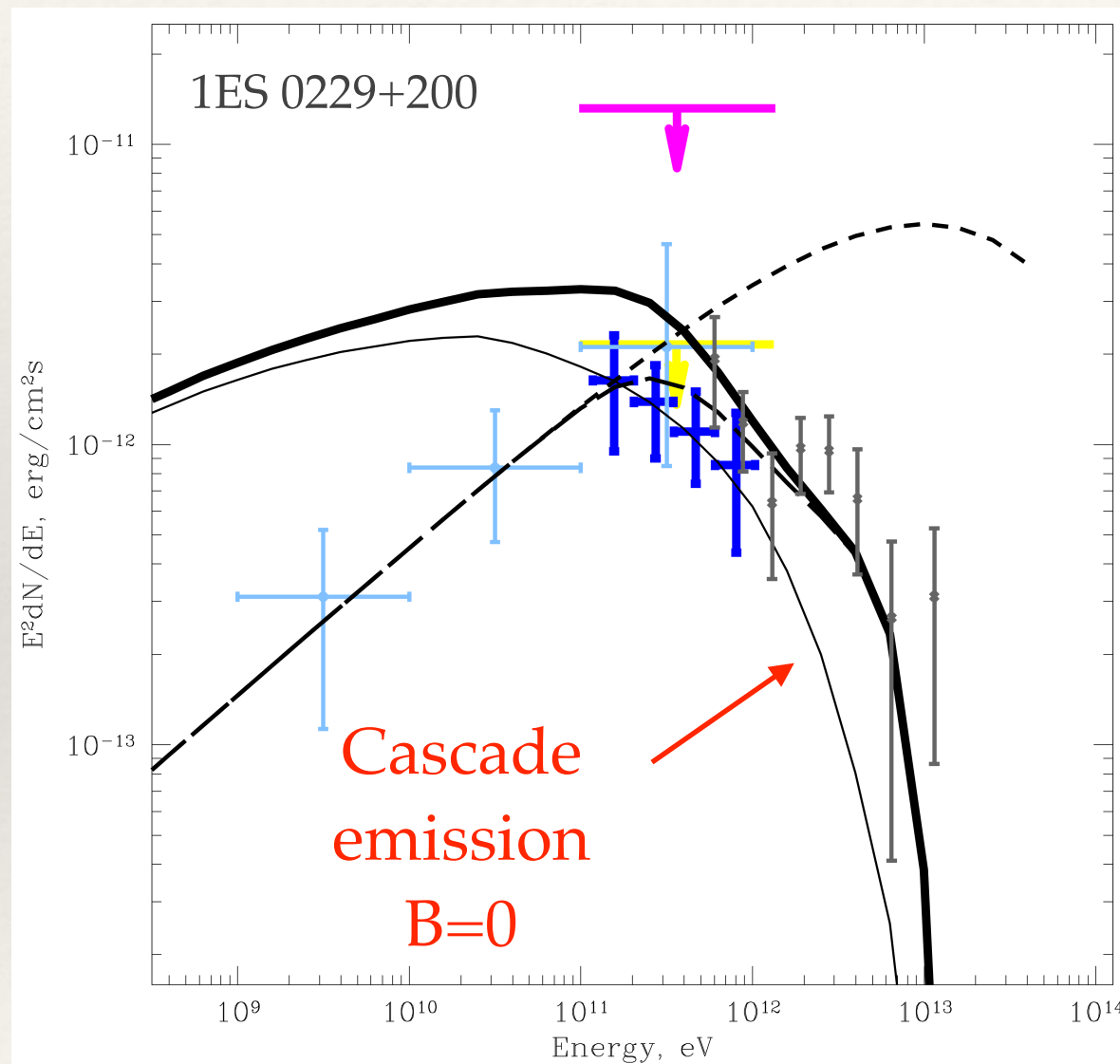
PKS 1424+240 & 1ES 1011+496

- ❖ Also for these sources the contribution of cascade emission is too low. **To be able to constrain the IGMF we need sources with hard intrinsic spectra.**



1ES 0229+200 & RX J1136.5+6737

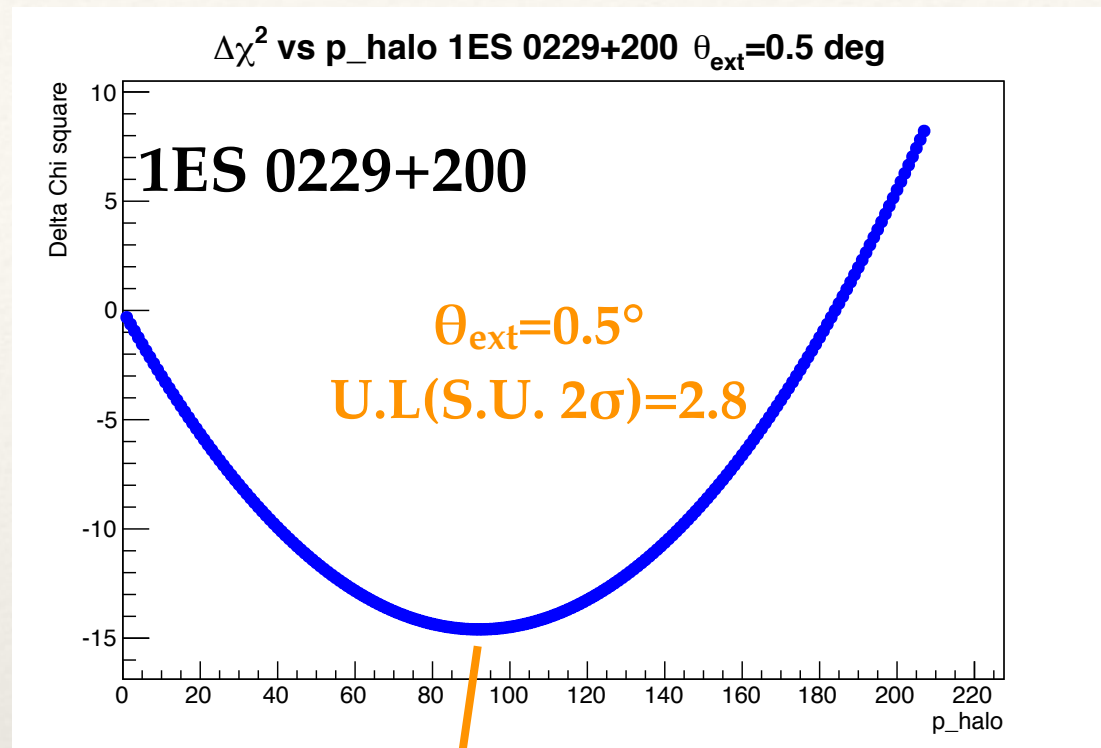
- ❖ These sources have intrinsic spectral index less than 1.8



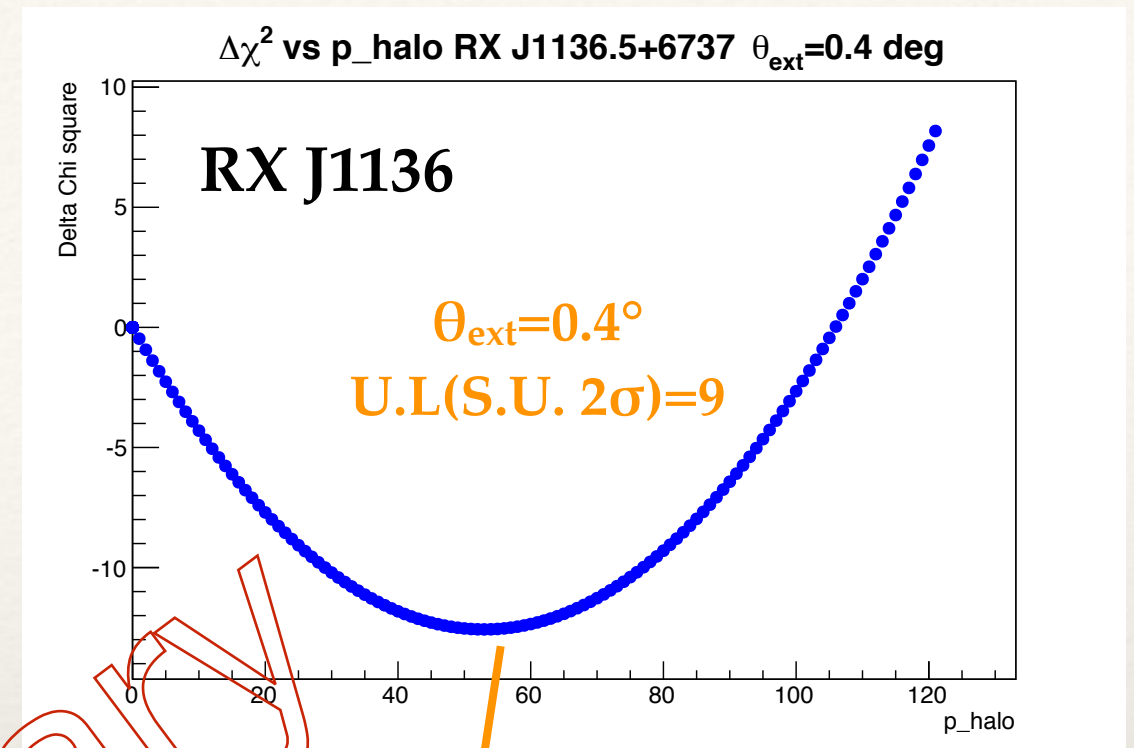
In these cases our upper limits are above the source level...

1ES 0229+200 & RX J1136.5+6737

$\Delta\chi^2$



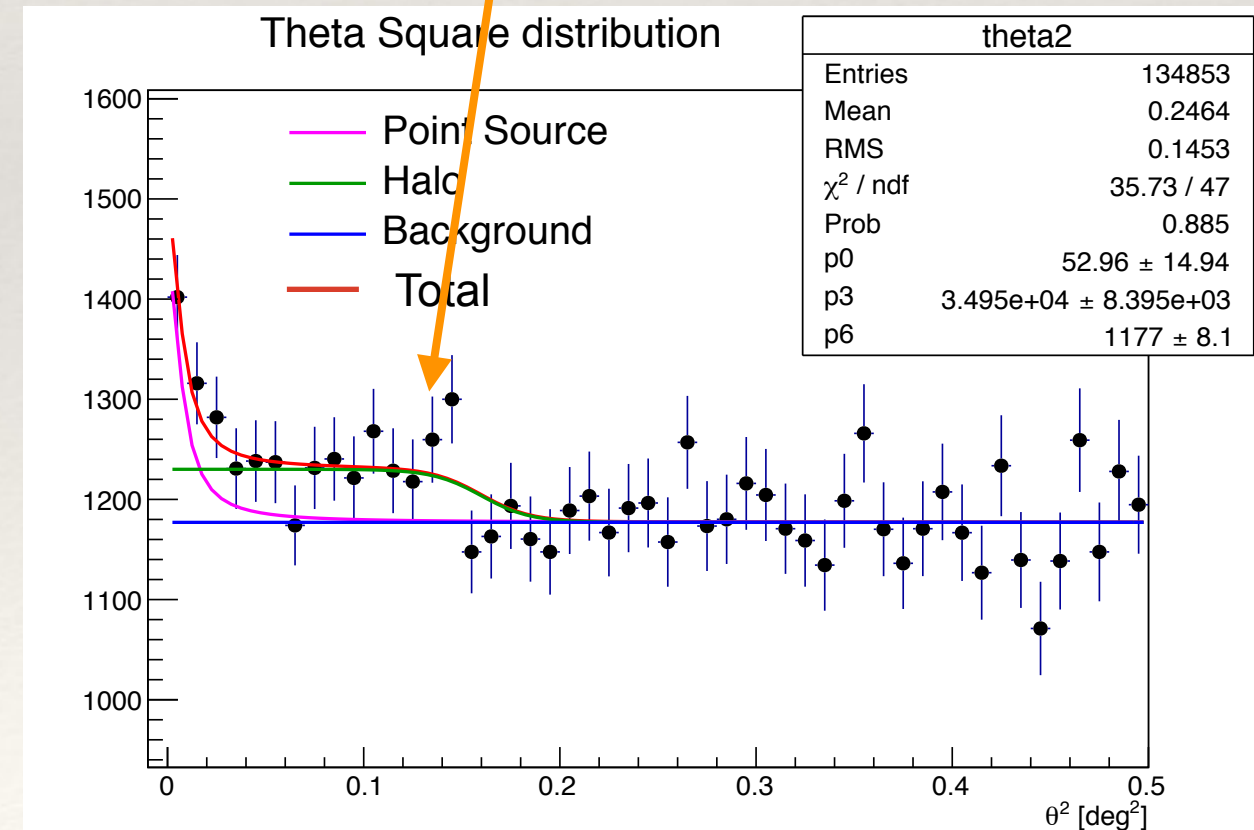
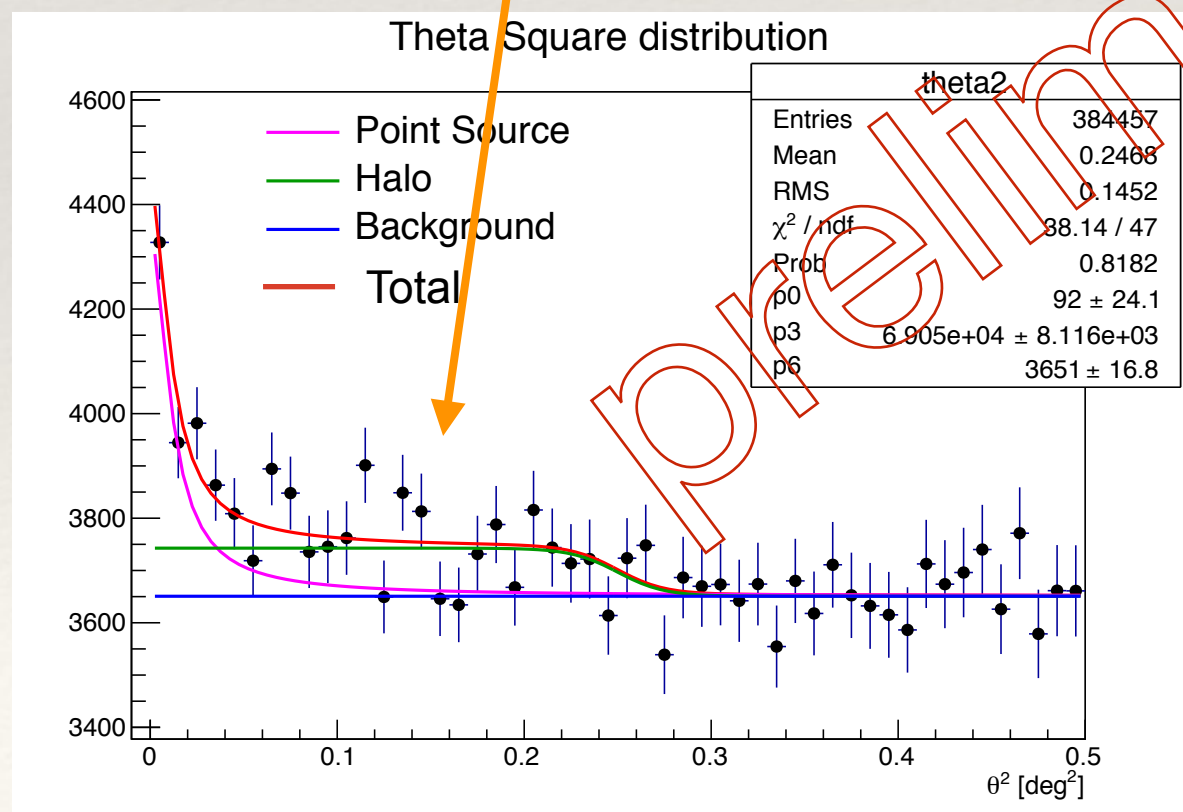
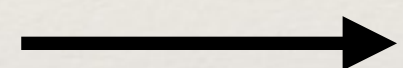
$\Delta\chi^2$



Halo intensity



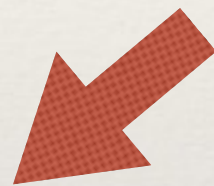
Halo intensity



IGMF from 1ES 0229+200 and RX J1136.5+6737

- ❖ Assuming that the shape of the θ^2 distributions are due to the presence of extended emission we estimate the IGMF strength

$$\Theta_{ext} \simeq 0.5^\circ (1+z)^{-2} \left[\frac{\tau}{10} \right]^{-1} \left[\frac{E_\gamma}{0.1 \text{ TeV}} \right]^{-1} \left[\frac{B_0}{10^{-14} \text{ G}} \right]$$



1ES 0229+200

RX J1136.6+6737

$\Theta_{ext}=0.3^\circ\text{-}0.55^\circ$ (2σ level)

$E_\gamma \simeq 100 \text{ GeV}$

$z=0.14$

$\tau \simeq 10$

$\Theta_{ext}=0.35^\circ\text{-}0.45^\circ$ (2σ level)

$E_\gamma \simeq 200 \text{ GeV}$

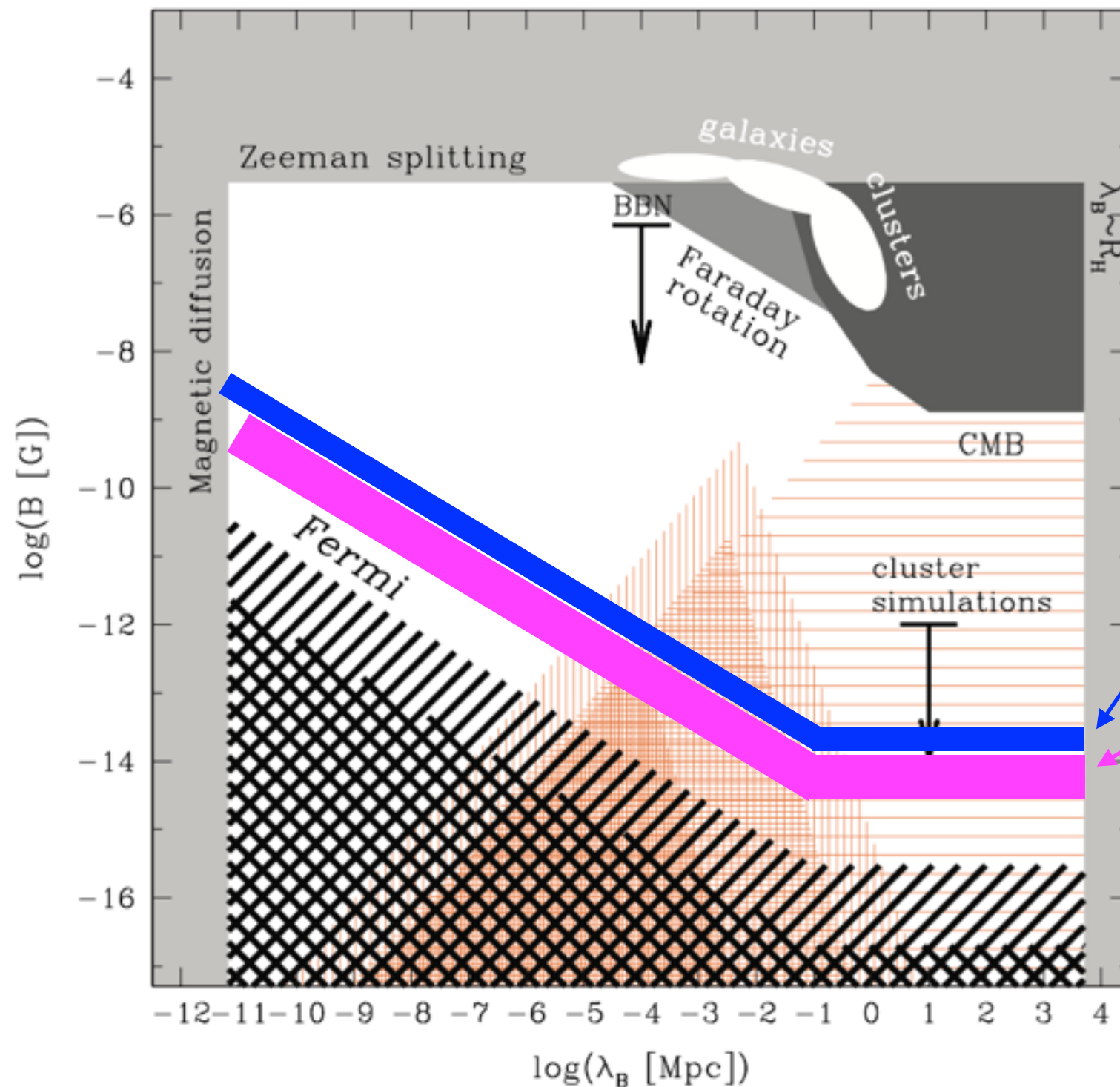
$z=0.13$

$\tau \simeq 15$

$B \simeq (0.5\text{-}1.3) \times 10^{-14} \text{ G}$

$B \simeq (3.2\text{-}4.2) \times 10^{-14} \text{ G}$

IGMF from 1ES 0229+200 and RX J1136.5+6737



RXJ 1136.5+6737

1ES 0229+200

Conclusions

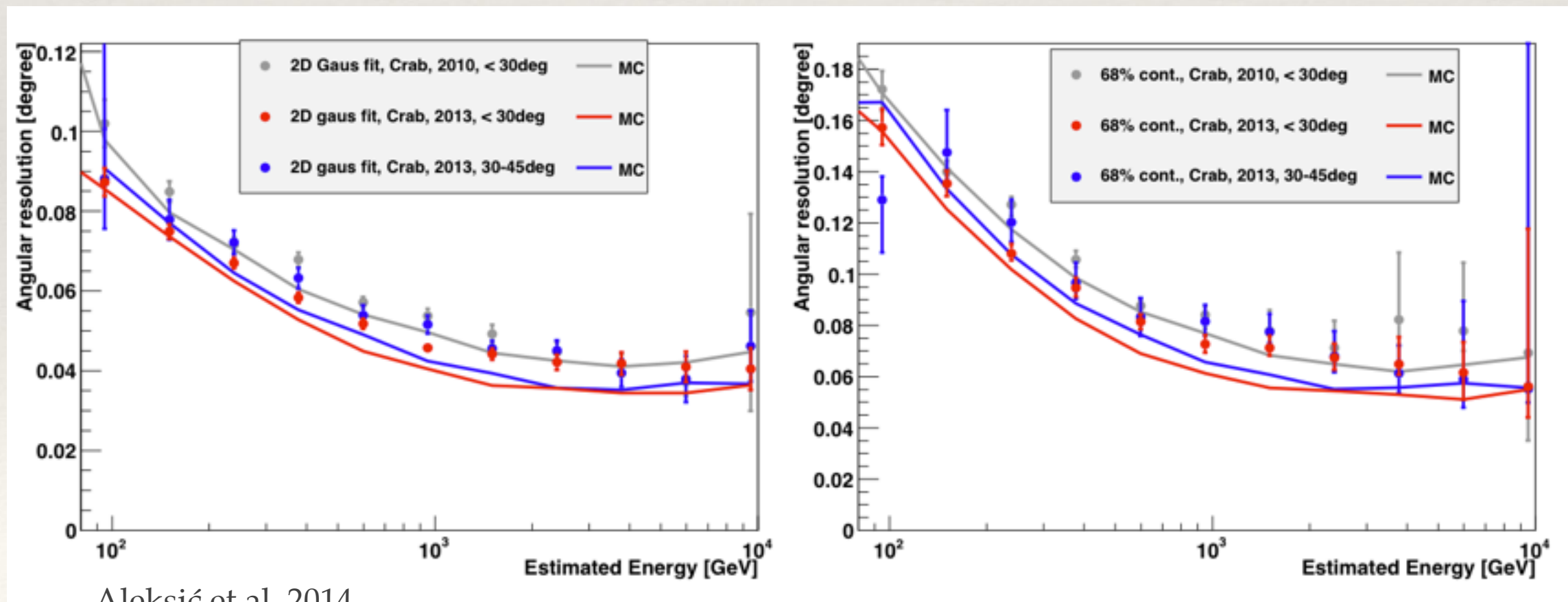
- ❖ We studied the emission profiles of several AGN in order to look for extended emission in VHE domain.
- ❖ We characterized the MAGIC PSF and gave a good analytical description (the King function).
- ❖ We computed upper limits on halo emission for two different halo models, improving previous published values on Markarian 421.
- ❖ The emission profiles of 1ES 0229+200 and RX J1136.5+6737 are compatible with the presence of extended emissions. The implied strengths of IGMF of the order of 10^{-14} G
 - ❖ Results needs to be confirmed with higher statistics with MAGIC.
- ❖ Expected improvements with CTA due to better sensitivity and angular resolution

Back up

Angular resolution

❖ Two techniques:

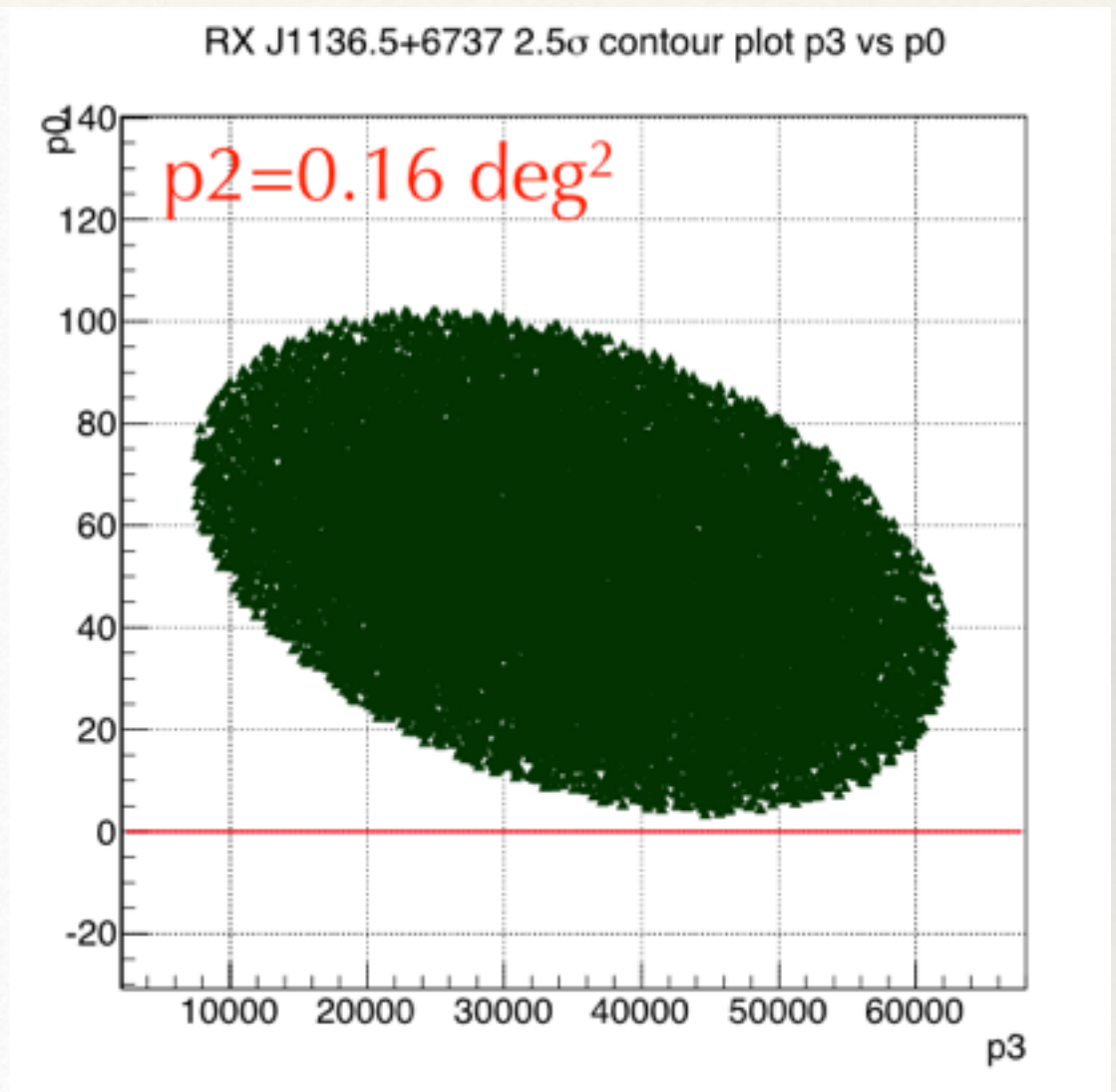
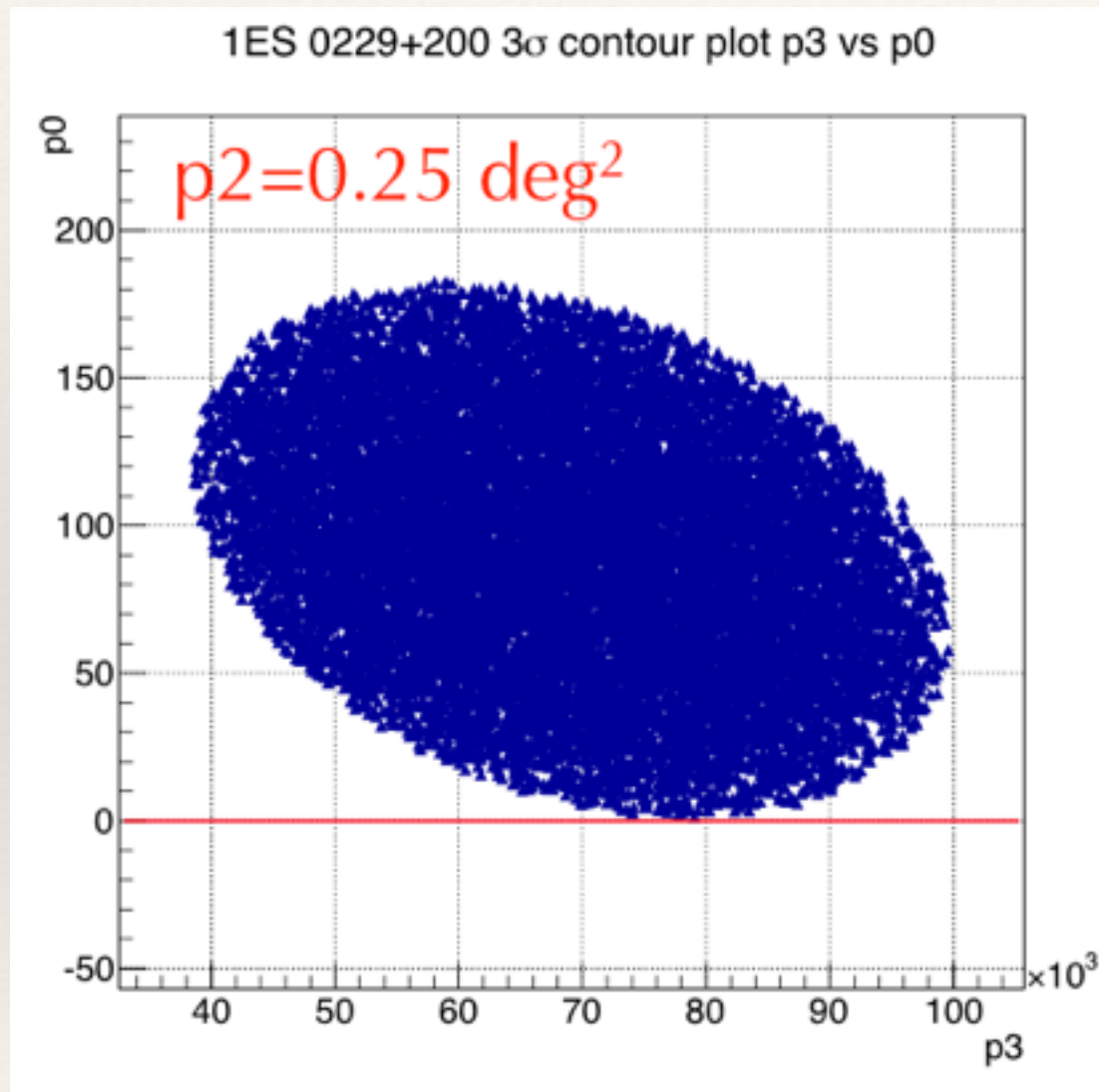
1. fit of the distribution of the γ -ray excesses with a 2 dimensional Gaussian
2. angular region that contains 68% of the excesses



Aleksić et al. 2014

1ES 0229+200 & RX J1136.5+6737

❖ Significance of the model:



$$f(\theta^2) = halo(p_0, p_1, p_2; \theta^2) + p_3 * p_4 * \left[1 + \left(\frac{\theta^2}{p_4} \right)^2 \right]^{-p_5} + p_6$$