





Does Analog Computation Exist?

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"Operating ICT basic switches below the Landauer limit"

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Landauer's Principle

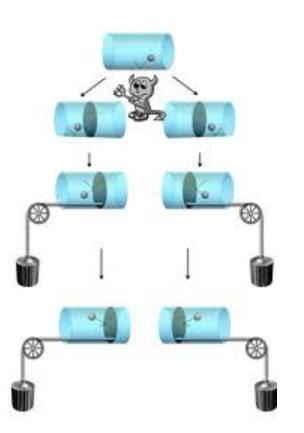
- Landauer principle: information is physical
 - →information theory
 - thermodynamics (second law)
- Relevant for the engineering of devices used to process information: Landauer principle= fundamental bound on energy dissipation ⇒ energy consumption of a device is:
 - Landauer bound + how good we are in engineering devices

Landauer bound + how good we are in engineering devices?

(at least in principle)

Szilard engine

One-particle gas in a partitioned box (Szilard 1929)



Initial probability ½ of the atom being on either side of the barrier

The demon performs a measure to determine the position of the particle (no heat production): information is gained

The piston is moved to one side or the other: extract work W= kT ln 2 that can be used to lift a weight

Apparent violation of the II law

Resolution of the puzzle:

forgetting is costly!

(Landauer 1961)



- Szilard engine: the Demon has a single memory register (0,1), initially 0; after measure it represents the location of the particle, 0 or 1
- To complete the cycle the memory of the Demon must be erased (information lost)
- Landauer principle: erasure of information requires a minimun heat production: ⟨Q⟩ ≥ ⟨Q_{Landauer}>= kT ln2 per bit ⇒ the conversion of kTln2 work into heat compensates the work extracted!
 - Information is physical

information is physical ⇒ system with i= 1....M possible states initial state: systems can be in any one of the possible states with probability p_i: S_S = -∑_{i∈M} p_i In p_i final state: system is in a specific state with probability 1, S_S =0 variation of Shannon entropy
 ΔS_S = ∑_{i∈M} p_i In p_i ≡ Boltzmann (thermodynamic) entropy (k=1)

$$\Delta S_S = \sum_{i \in M} p_i \ln p_i \equiv Boltzmann (thermodynamic) entropy (k=1) since information can only be processed by physical systems (computers)$$

 -∆S ≥ k ln2 per bit for perfect erasure, what happens if we admit errors?

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p = error probability 0 \le p \le \frac{1}{2}
Landauer bound: -\Delta S/k \ge \ln 2 + p \ln p + (1-p) \ln (1-p)
-\Delta S/k = \ln 2 + p \ln p + (1-p) \ln (1-p) if the erasure is efficient for p=0 -\Delta S \ge k \ln 2; for p=\frac{1}{2} -\Delta S \ge 0
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Landauer principle for analog computing systems:

we want to show that the continuous generalization

$$\Delta S_S = \int_{x \in M} p(x) \ln p(x) \equiv Boltzmann (thermodynamic) entropy generated when the erasure is realized in a physical system$$



dimensionless probability density p(x) cannot be defined by dimensionful physical degrees of freedom without an appropriate regularisation

 define a density (note: not probability density) that regularizes the entropy (Jaynes)

Landauer's principle and phase transitions

Idea: continuous phase transitions with order parameter m ≡ information erasure by resetting to standard value ⇒ entropy change must satisfy Landauer bound

0≤ m ≤ 1 plays the role of an error probability p

- SSB: degenerate vacua, the system "chooses" a state ≡ reset
- T= 0 ⇒ m =1 perfectly ordered phase ≡ reset with no errors
- 0 <T < T_C ⇒ 0 < m < 1 partially ordered phase ≡ reset with errors
- T = T_C ⇒ m = 0 (disordering procedure efficient) →
 Landauer bound saturated
- phase transition at T=T_C

How do we prove this conjecture for classical systems?

(C.A. Trugenberger, MCD PhysRevE.89 (2014) 052138)

- associative memory: (Hopfield model) bits are stochastic neurons s_i = ± 1, mean field
- phase transition from a disordered phase to an ordered phase with order parameter increasing from the value 0 at T=T_c to 1 at T=0

P=
$$\frac{1}{2}$$
 (1-m) ; $\frac{-\Delta S(J,T)}{kN} = \frac{1}{kN} (S_{T_c} - S_T)$

 $\frac{-\Delta S}{kN} = \ln 2 + p \ln(p) + (1-p) \ln(1-p),$

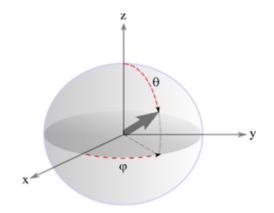
thermodynamic entropy completely scales in terms of this error probability alone

• Z_2 , easy generalization to $Z_{(2n+1)}$ n=1/2, 1,3/2,.....

Continuous Symmetries

discrete → continuous symmetry digital →analog

• $O(3) \rightarrow O(2)$, Heisenberg ferromagnet $S^2 = 1$



infinite many possible orientation

mean field: classical limit of the discrete case

- L(x) = coth x 1/x Langevin function,
 L(x)→1 x→∞; L(x) →0 x=0
- T= T_C m \rightarrow 0 \Rightarrow $S(T_C)$ /kN = In 4π = volume of the phase space \equiv area of the 2-sphere
- What happen when T→0 ?

$$\lim_{T\to 0} \frac{S(T,J)}{kN} = \ln(0) \to -\infty$$

why the entropy is negative and diverges?
 classical systems e.g. system of N classical harmonic oscillator:

$$\frac{S_{\text{ho}}}{kN} = \left[\ln \frac{kT}{\hbar\omega} + 1\right] \to -\infty \text{ for } T \to 0$$

- discrete case (quantum): entropy not divergent, can we use this to regularize the entropy for the continuous symmetry?
- Spin s, (2s +1) components ≡ representation of O(3);
 s→∞ classical spin on the unit sphere

(Millard and Leff, J. Math Phys 12 (1971), 1000; Lieb, Commun. Math Phys 31 (1973) 327)

- classical limit: $s \rightarrow s/s_{max}$ $\Delta(2s_{max}+1) \rightarrow 4\Pi$ for $s_{max} \rightarrow \infty$, $\Delta \rightarrow 0$
- Heisenberg principle:

$$\Delta = \frac{1}{2}\hbar^2 s_{\max} ; \hbar s_{\max} \to s_{\text{clas}} \text{ for } s_{\max} \to \infty , \hbar \to 0$$
 $S_{\text{clas}} = 1$

$$\frac{S}{kN} = \ln \frac{4\pi}{(2s_{\text{max}}+1)} \left[\frac{\sinh(1+\frac{1}{2s_{\text{max}}})\beta mJ}{\sinh\frac{\beta mJ}{2s_{\text{max}}}} \right] - \beta mJB_{s_{\text{max}}}(\beta mJ)$$

- $S(T=T_C)/kN = In 4\Pi$ for $s_{max} \rightarrow \infty$
- $S(T \rightarrow 0) / kN = ln [4\Pi / (2s_{max} + 1)] = ln \frac{1}{2} \hbar^2 s_{max} = ln \frac{1}{2} \hbar$
- $\Delta S/kN = (S(T_C) S(T=0))/kN = ln (8\Pi / \hbar)$
- Analog Landauer bound: entropy production during erasure process = available configuration volume measured in units of the minimum quantum of configuration volume ⇒ regularization
- even if we start with continuous, analog information, only a finite countably amount of information can be encoded in a physical system (Shannon, Bekenstein)

Grazie!