

Evolution of Transfer-like Reactions from Coulomb to Fermi Energies

ALEXIS DIAZ-TORRES



European Centre for Theoretical
Studies in Nuclear Physics and
Related Areas



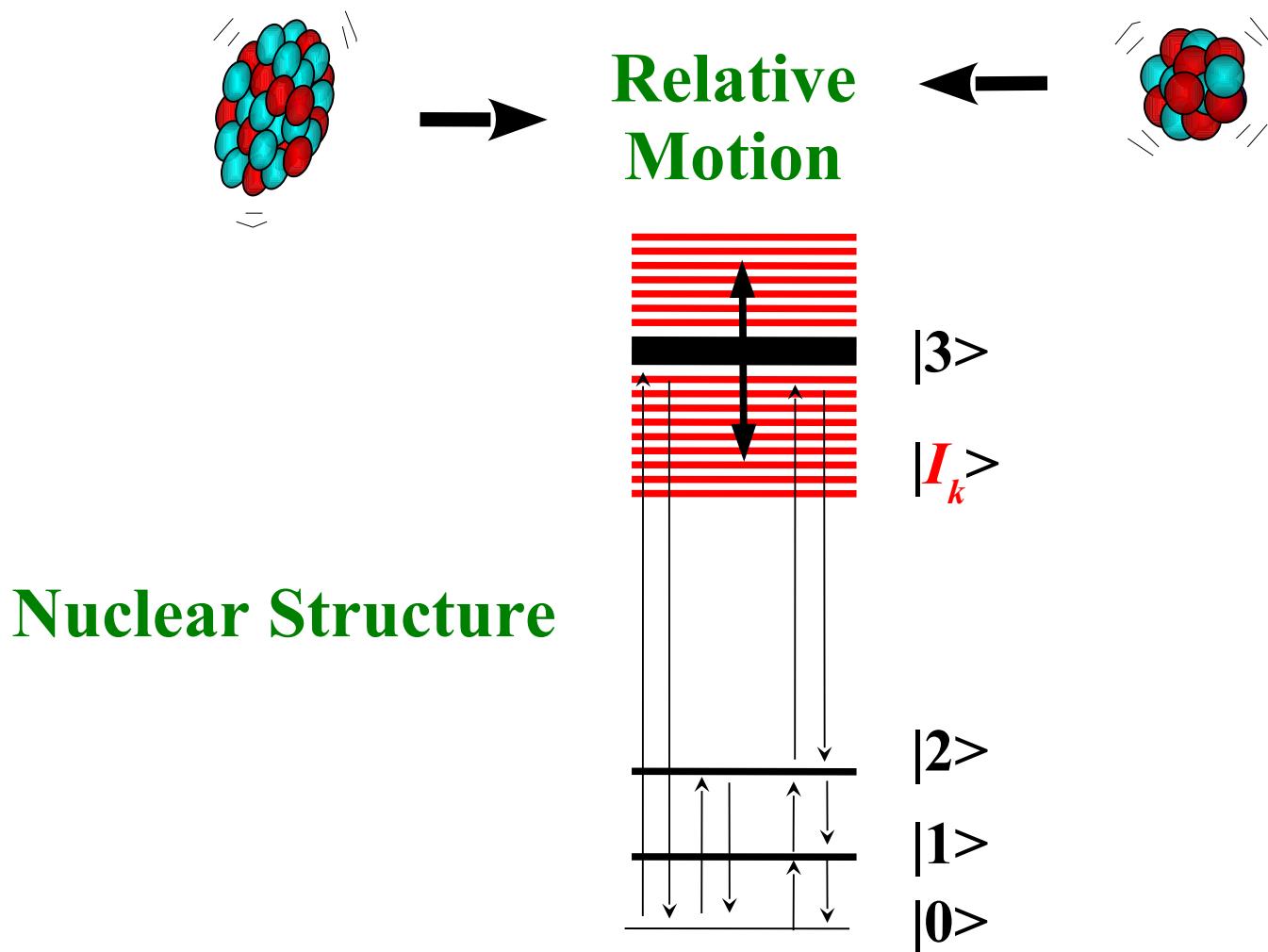
Trento, Italy

https://www.researchgate.net/profile/Alexis_Diaz-Torres2



- ★ Two-Centre Shell Model: adiabatic & diabatic states
- ★ Effects of diabaticity on quasi-fission reactions
- ★ Summary & Outlook

Low-Energy Nuclear Reaction Dynamics

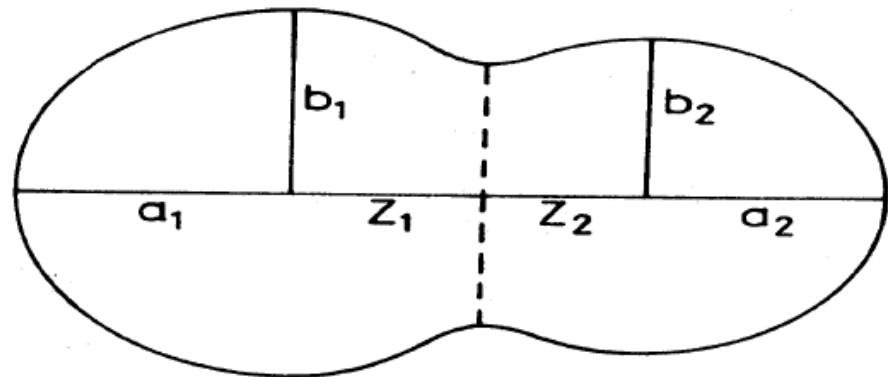


Strong interplay of **nuclear structure** and **reaction dynamics**
determines reaction outcomes (**cross sections**)

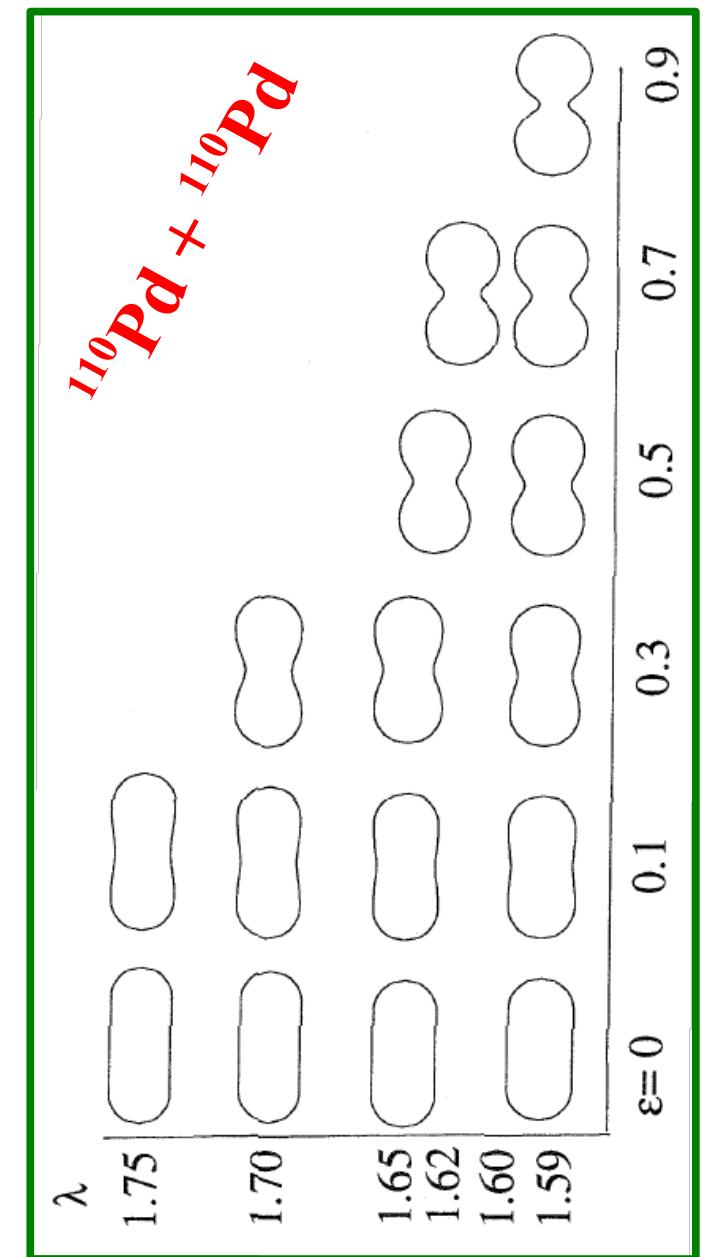
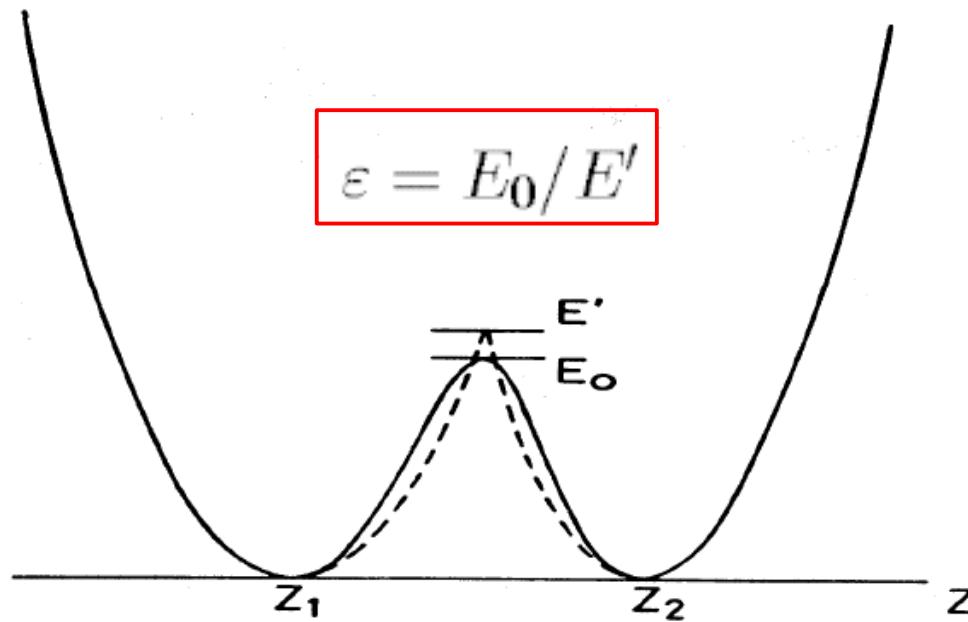
Two-Centre Shell Model (TCSM)

J.A. Maruhn & W. Greiner, Z. Phys. **251** (1972) 431

$$\lambda = l/(2R_0)$$



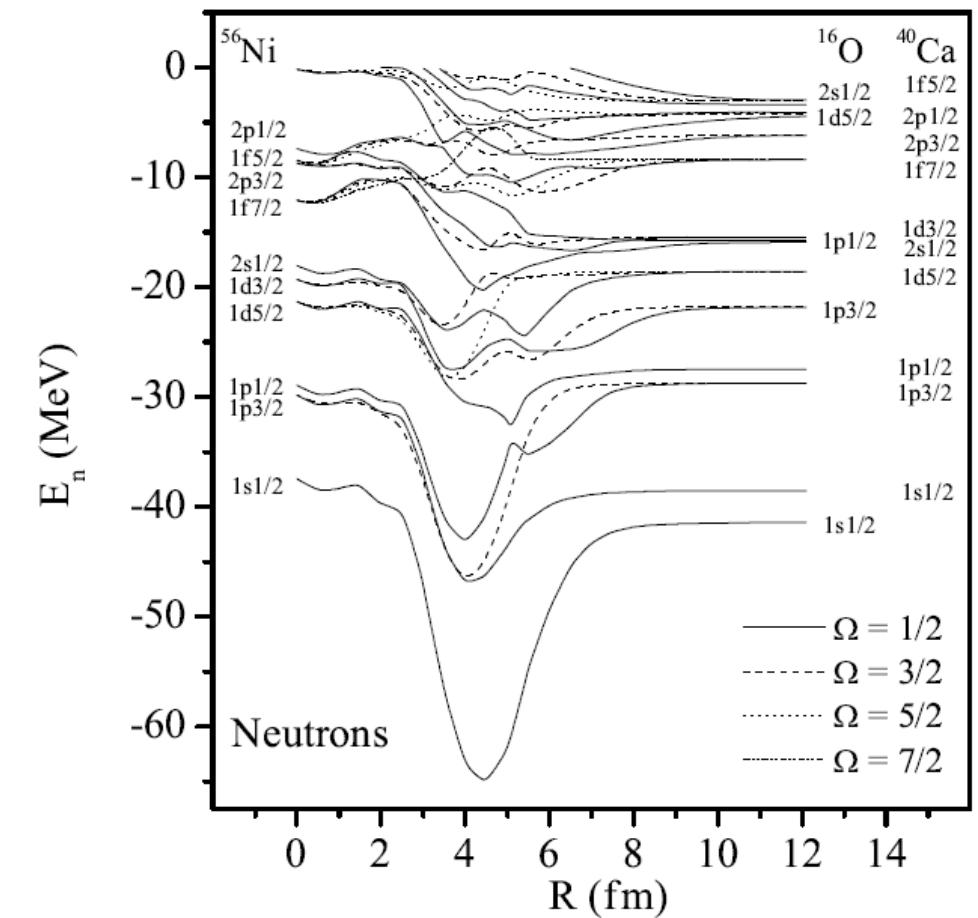
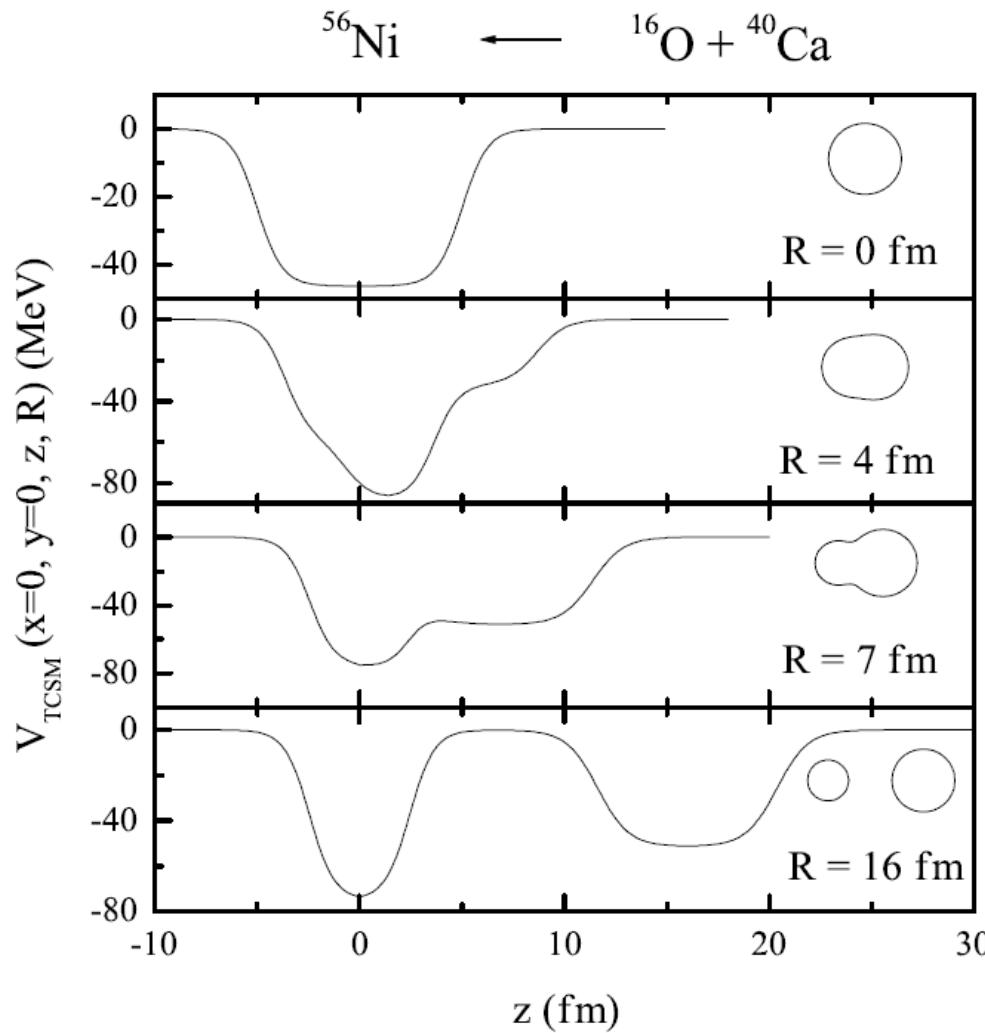
$$\varepsilon = E_0/E'$$



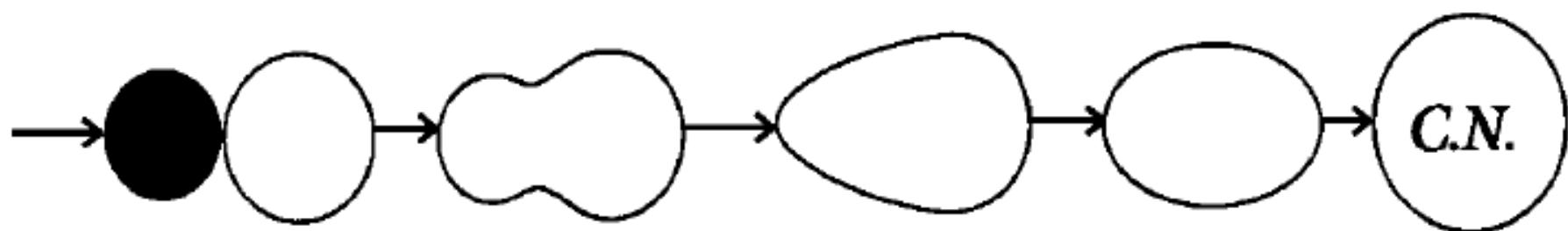
TCSM with Woods-Saxon Potentials

AD-T & W. Scheid, Nucl. Phys. A **757** (2005) 373

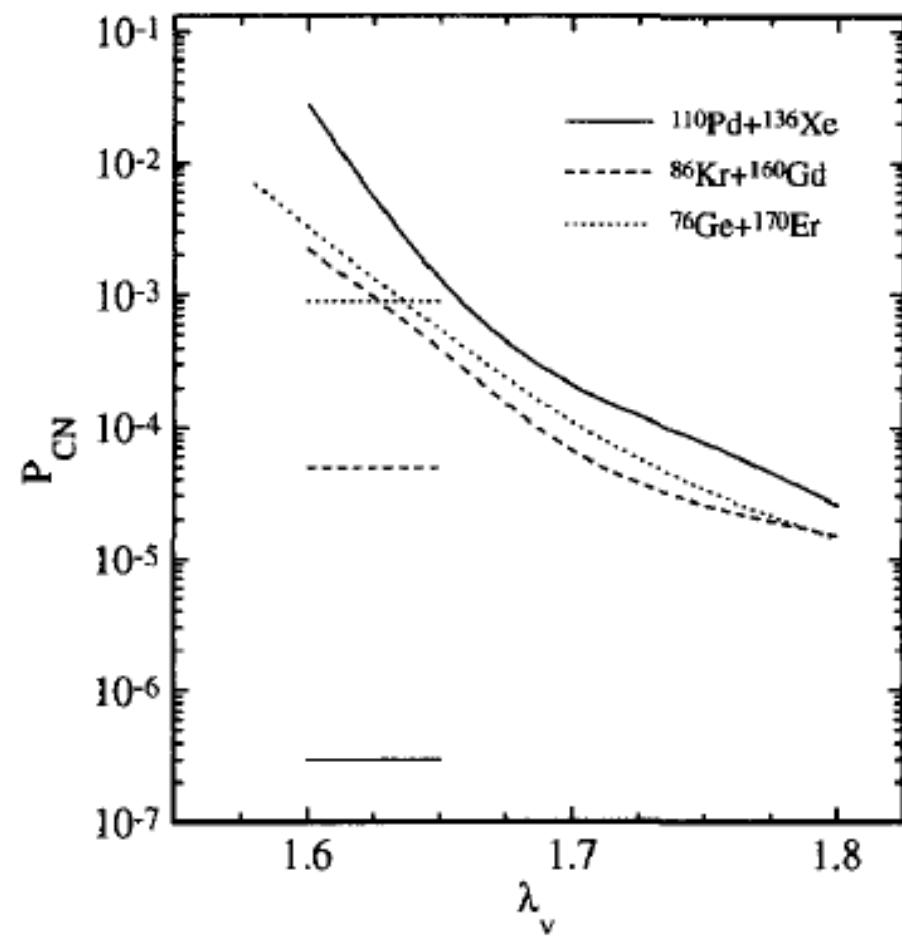
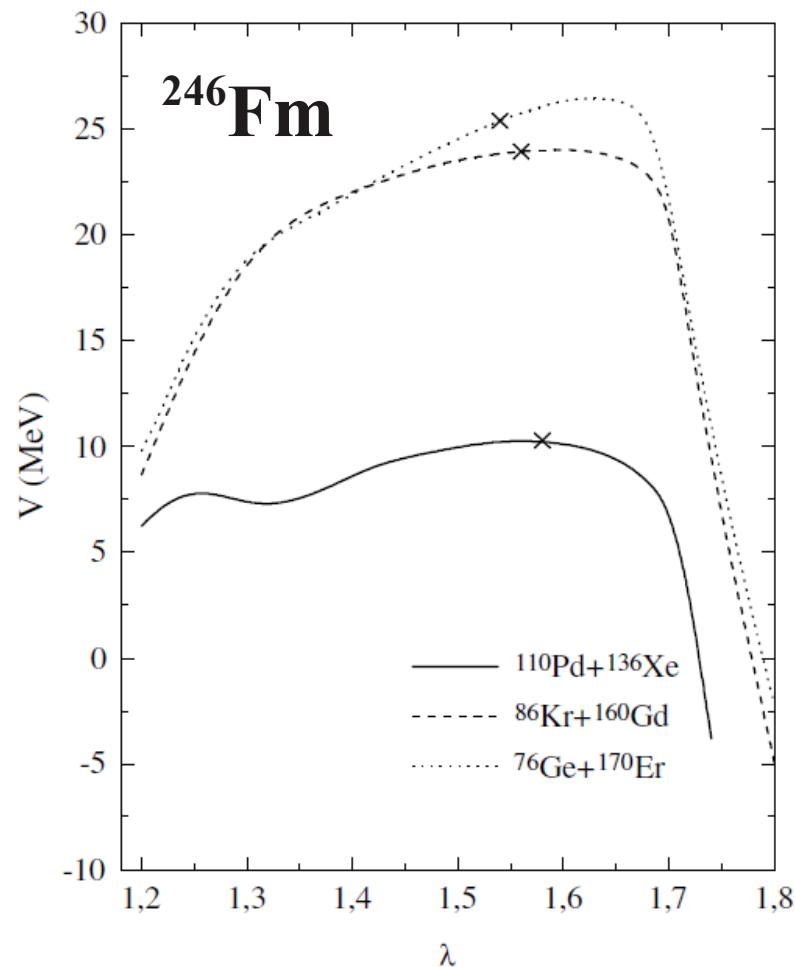
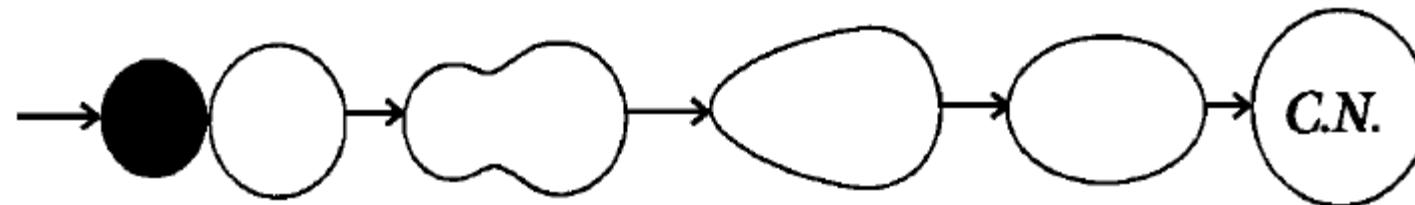
AD-T, Phys. Rev. Lett. **101** (2008) 122501



The Adiabatic Picture of Fusion



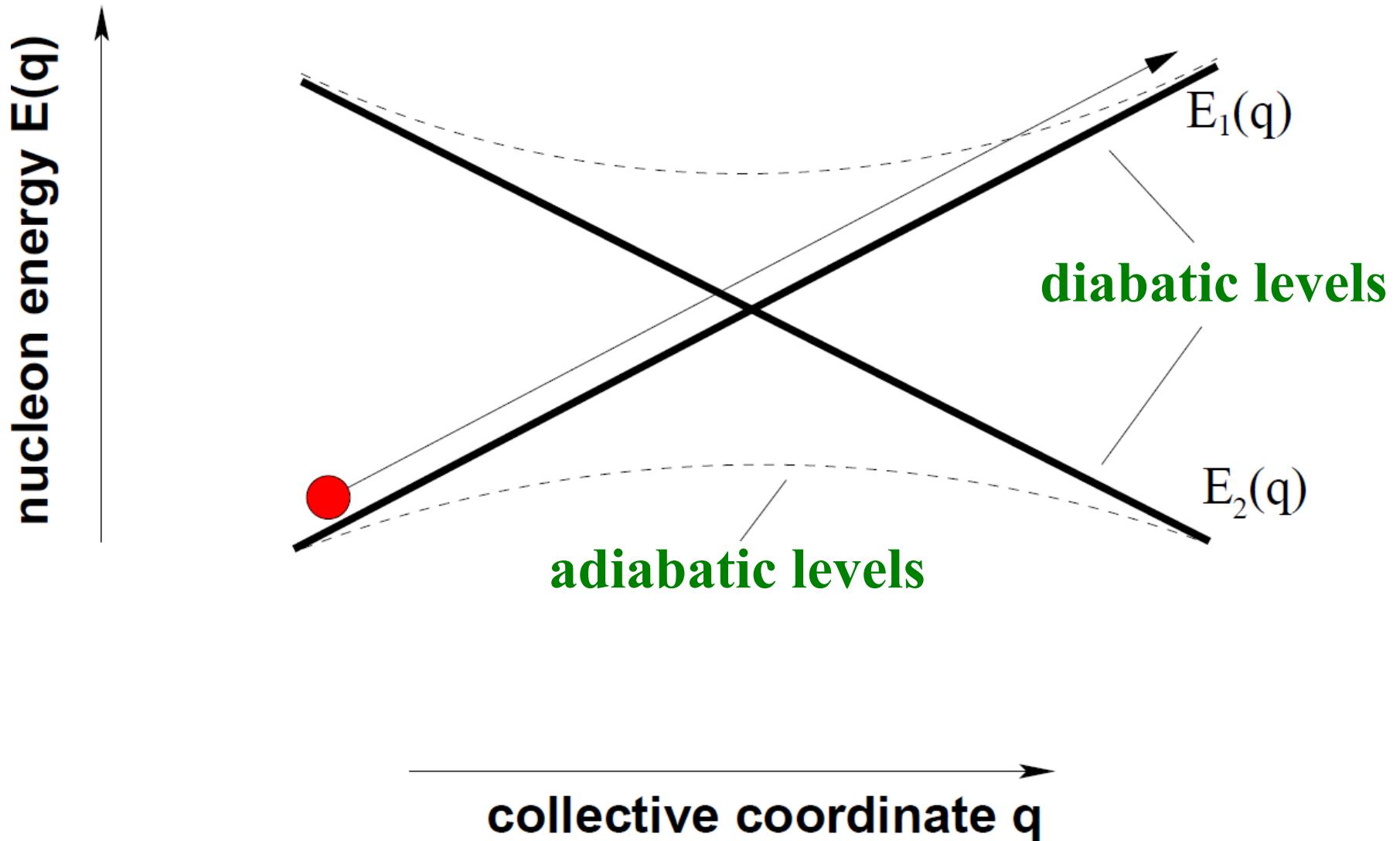
Some Problems with the Adiabatic Picture of Fusion



Adiabatic and Diabatic Single-Particle Motion

W. Noerenberg, Phys. Lett. B **104** (1981) 107

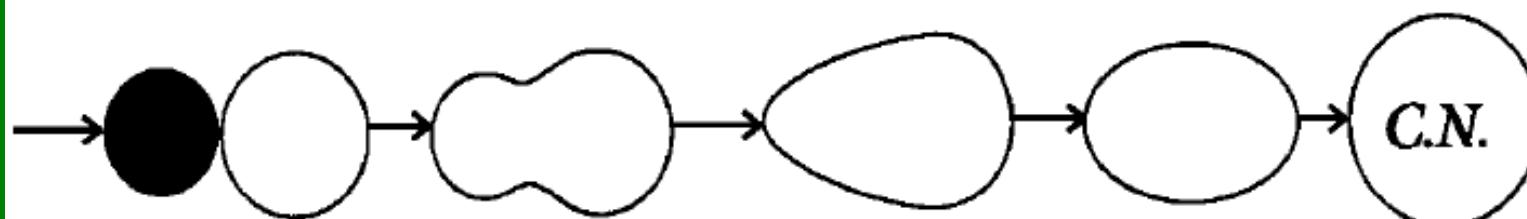
AD-T & W. Scheid, Nucl. Phys. A **757** (2005) 373



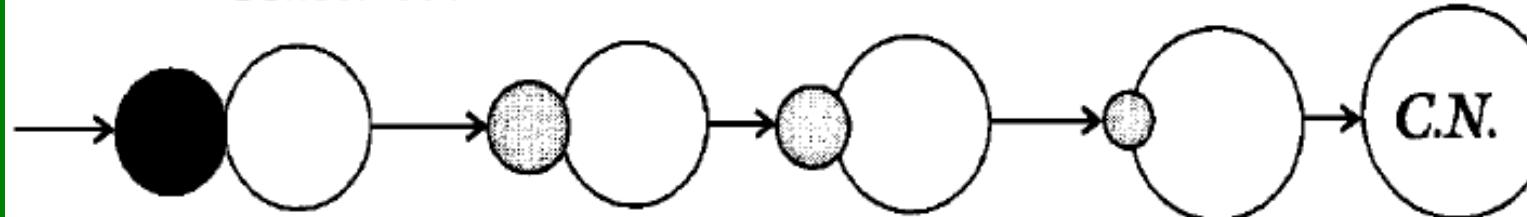
Fusion of Heavy Nuclei

The **diabatic picture** supports the dinuclear system model for complete fusion of heavy nuclei at Coulomb energies.

The Macroscopic Dynamical Model:
Fusion of Two Viscose Liquid Drops



The Dinuclear System Concept:
Conservation of Nuclear Individualities



Effects of diabaticity on quasi-fission reactions

- ◆ Dynamical collective potential-energy landscape plays a crucial role in determining the reaction outcomes.
- ◆ Elasto-plasticity should be important as energy increases from Coulomb to Fermi energies.

AD-T, Phys. Rev. C **69** (2004) 021603(R)

AD-T, Phys. Rev. C **74** (2006) 064601

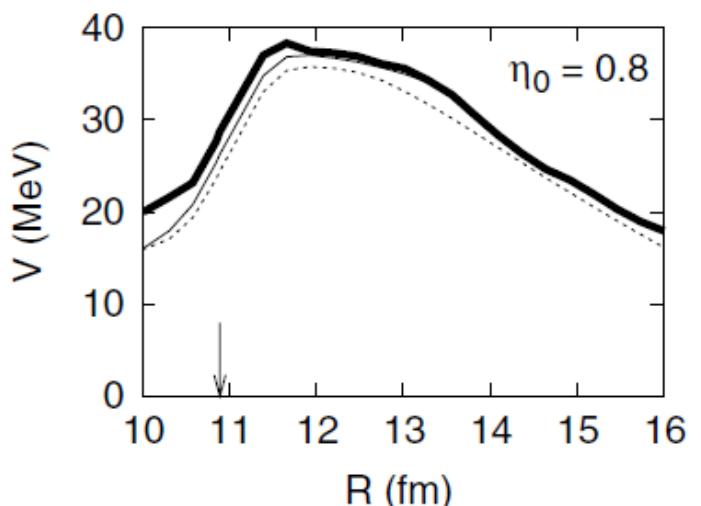
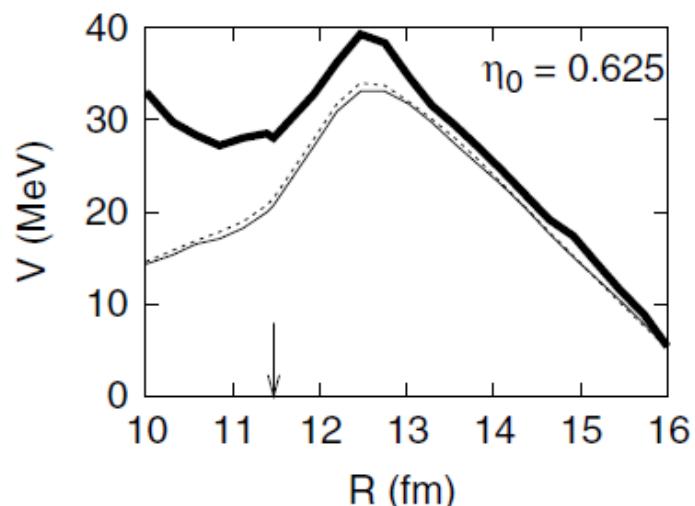
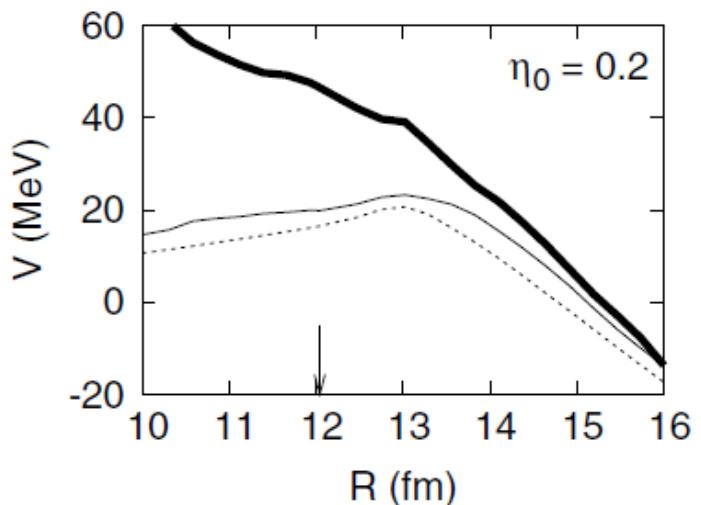
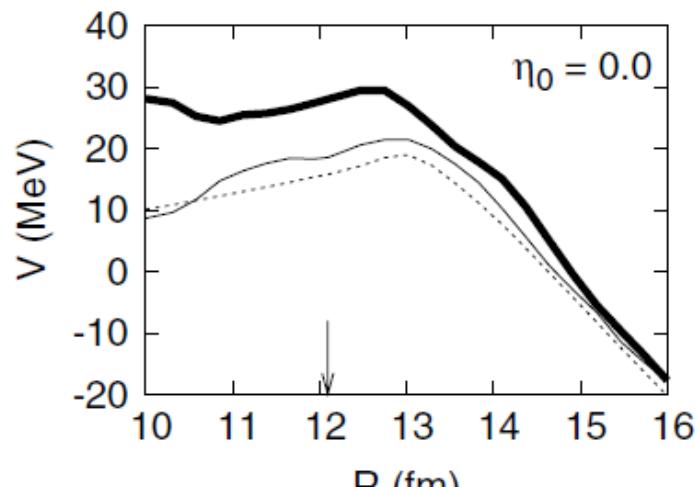
AD-T, Phys. Rev. C **82** (2010) 054617

Diabatic Contribution to the Collective PES

AD-T, Phys. Rev. C 74 (2006) 064601

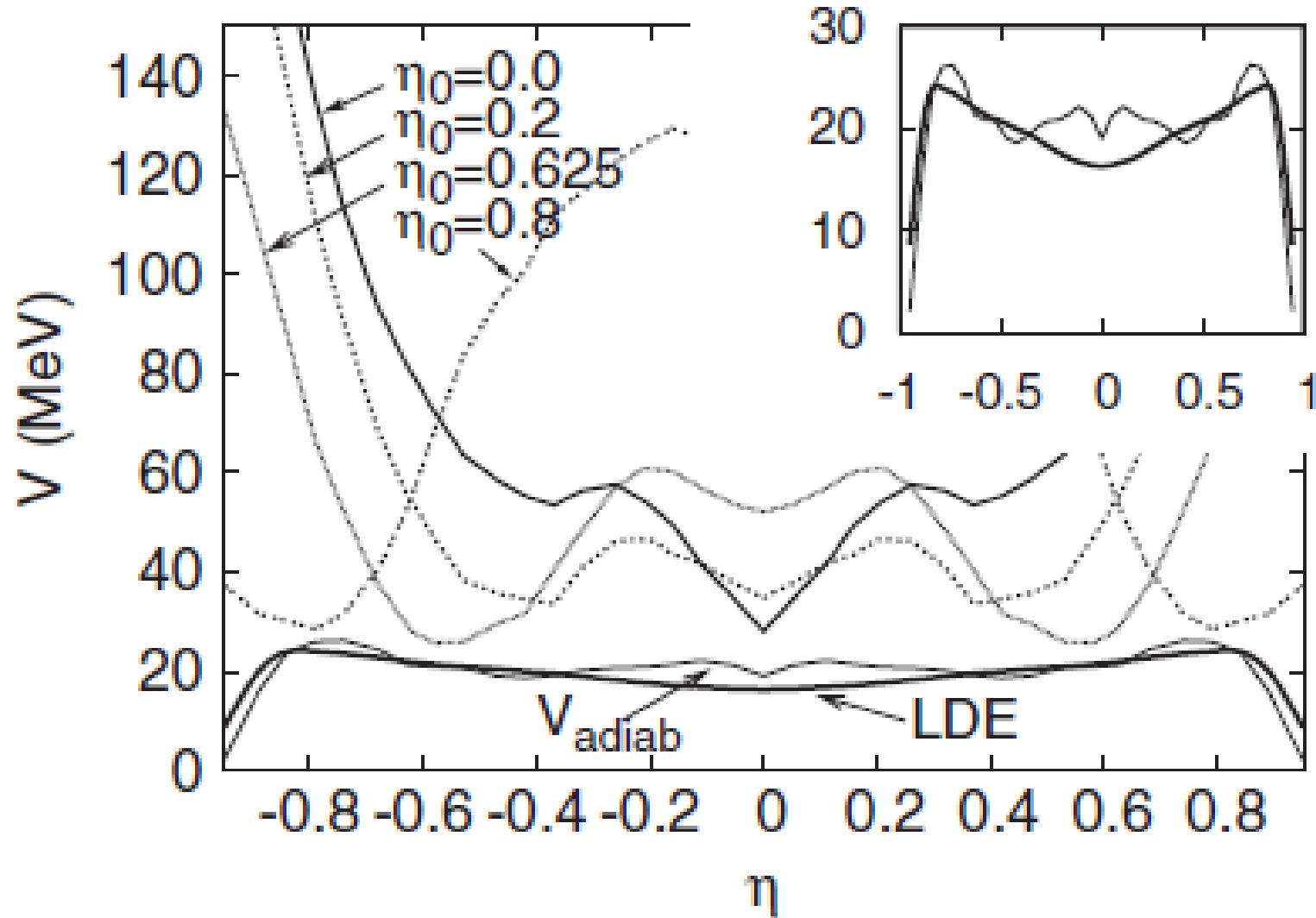
$$\Delta V_{diab}(q) \approx \sum_{\alpha} E_{\alpha}^{diab}(q) (n_{\alpha}^{diab} - n_{\alpha}^{adiab})$$

^{256}No ←



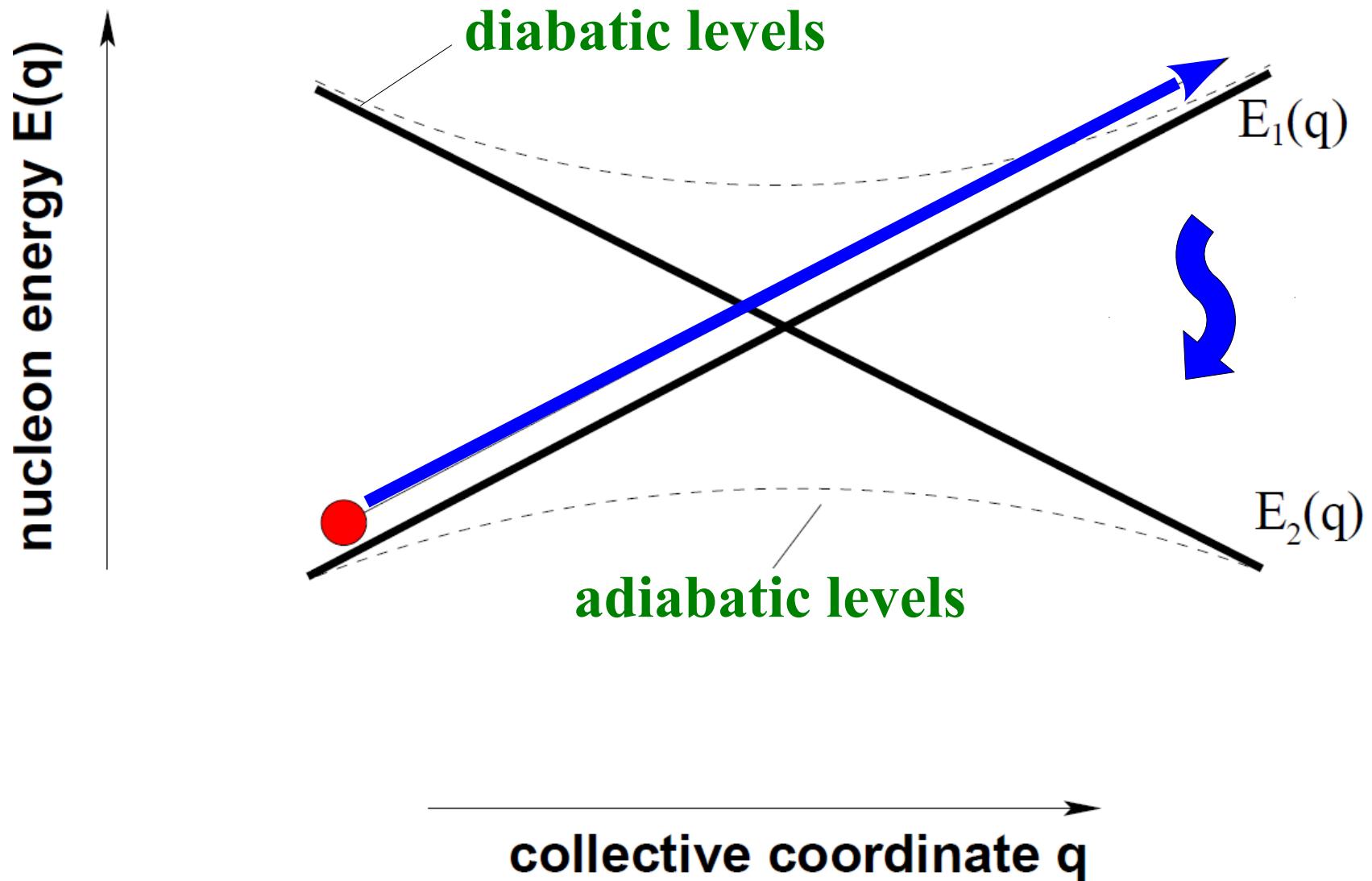
Driving Potential

AD-T, Phys. Rev. C 74 (2006) 064601

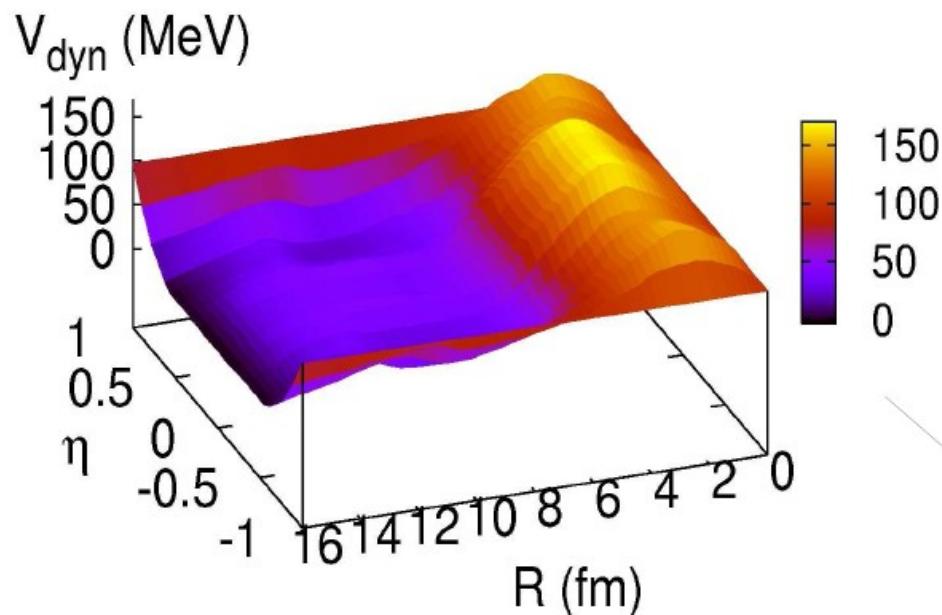


Diabatic Dissipative Dynamics

W. Noerenberg, Phys. Lett. B 104 (1981) 107



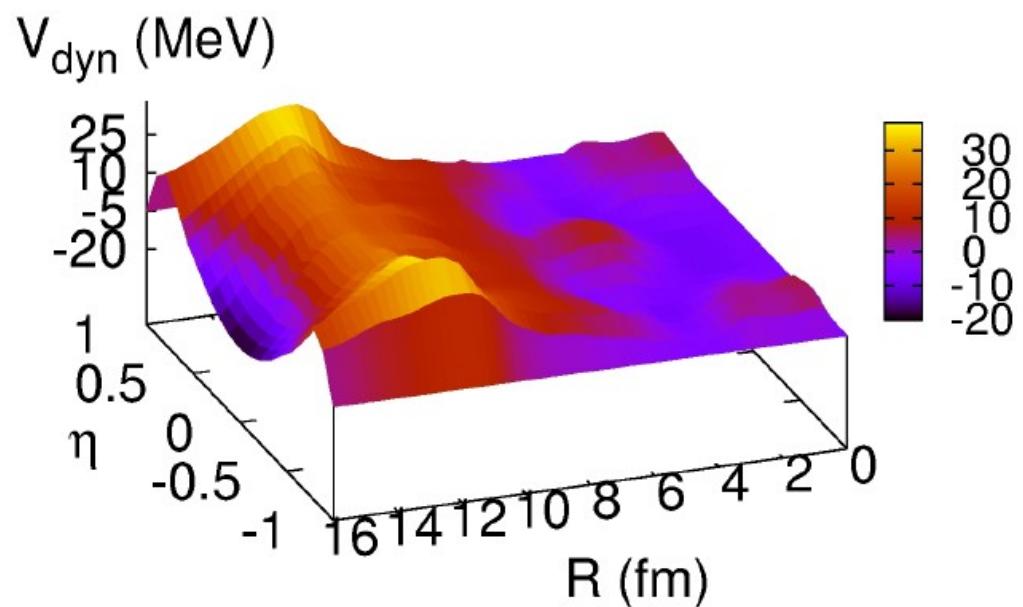
Time = 5×10^{-23} s



TIME

Time = 2×10^{-21} s

**Dynamical
Landscape !**



Equations of Motion

$$\frac{dP_\nu(t)}{dt} = \sum_{\mu \neq \nu} [\Lambda_{\nu\mu}(t)P_\mu(t) - \Lambda_{\mu\nu}(t)P_\nu(t)], \quad (1)$$

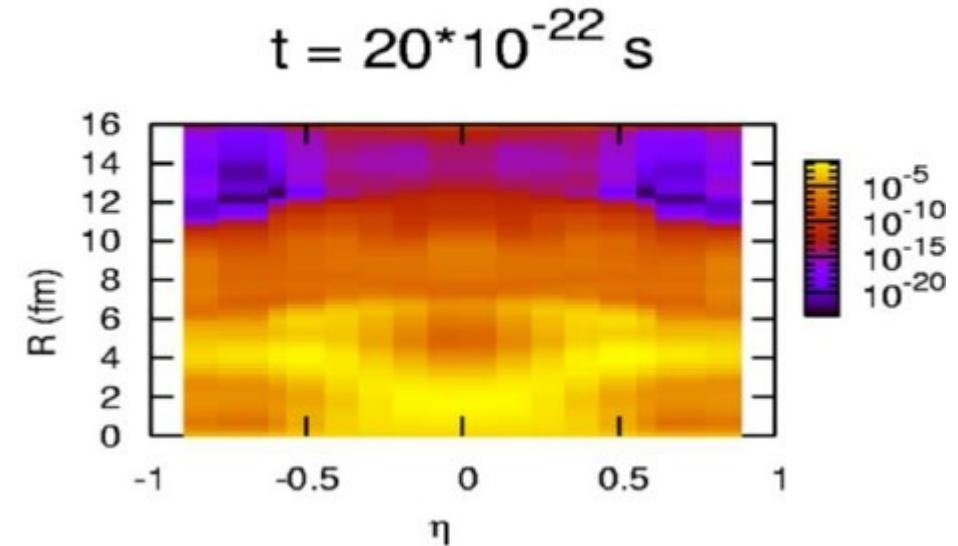
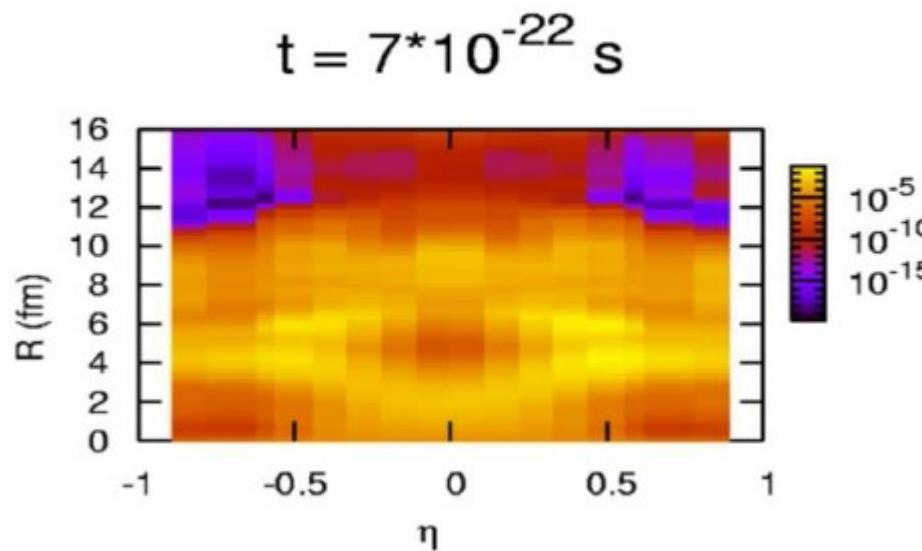
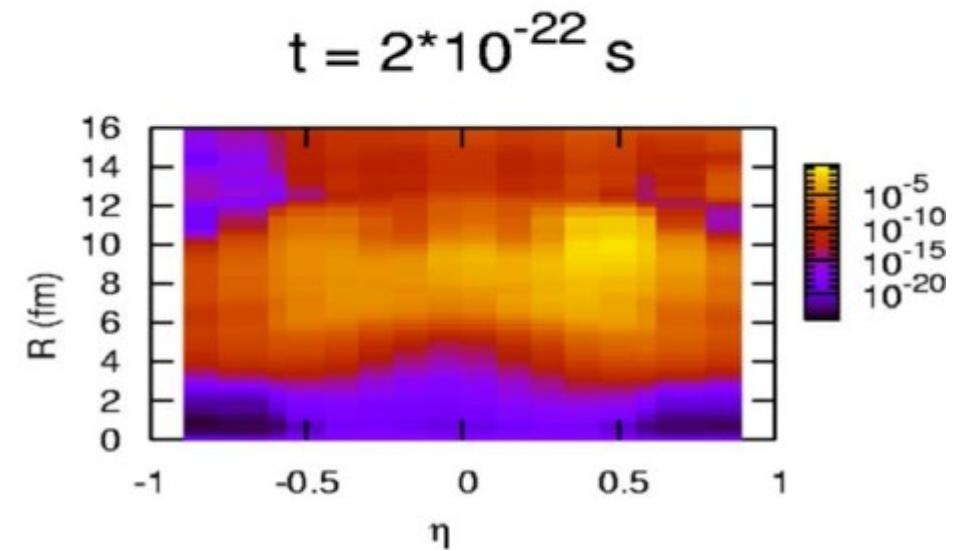
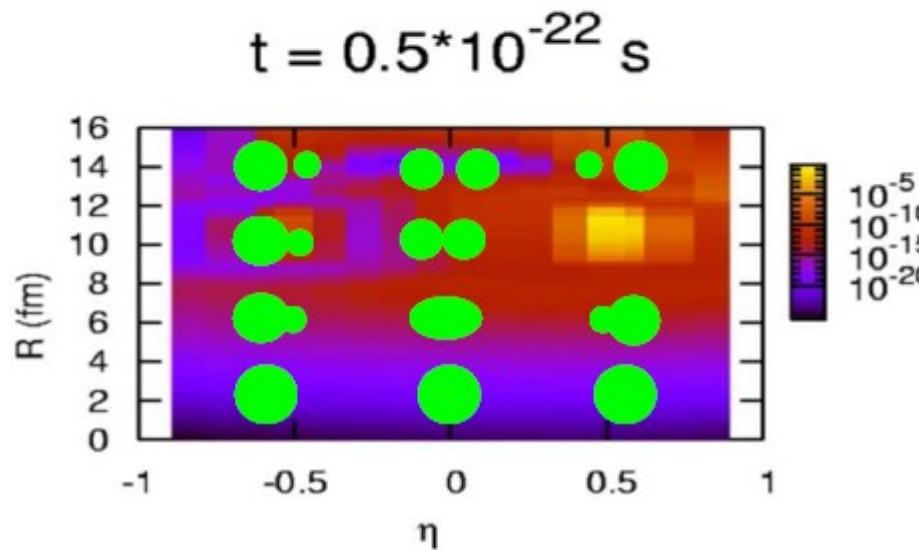
$$\Lambda_{\nu\mu}(t) = \lambda_0^{\nu\mu}(t) \exp[\sqrt{aE_\nu^*(t)} - \sqrt{aE_\mu^*(t)}], \quad (2)$$

$$\lambda_0^{\nu\mu}(t) \approx \kappa_0^{\nu\mu}(t) \exp \left\{ -\frac{(\mathbf{q}_\mu - \mathbf{q}_\nu)^2}{\tilde{\sigma}_\nu^2} \right\}, \quad (3)$$

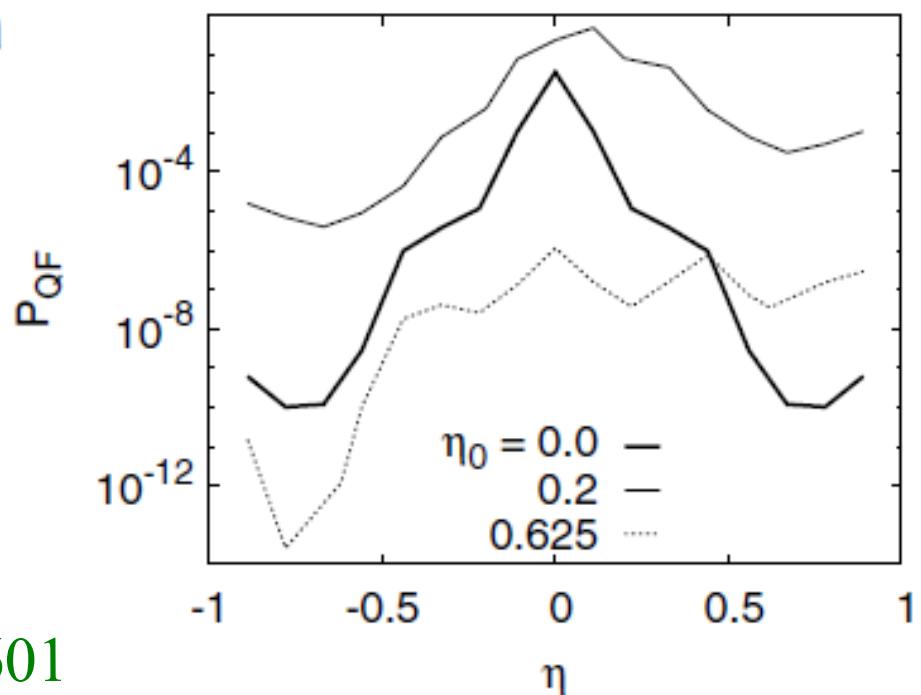
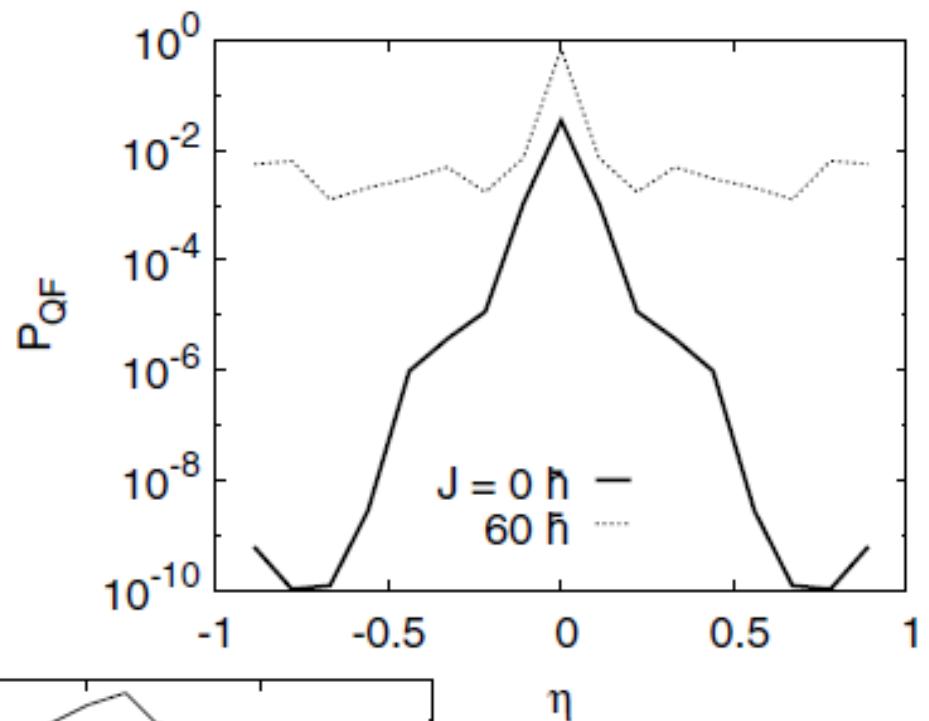
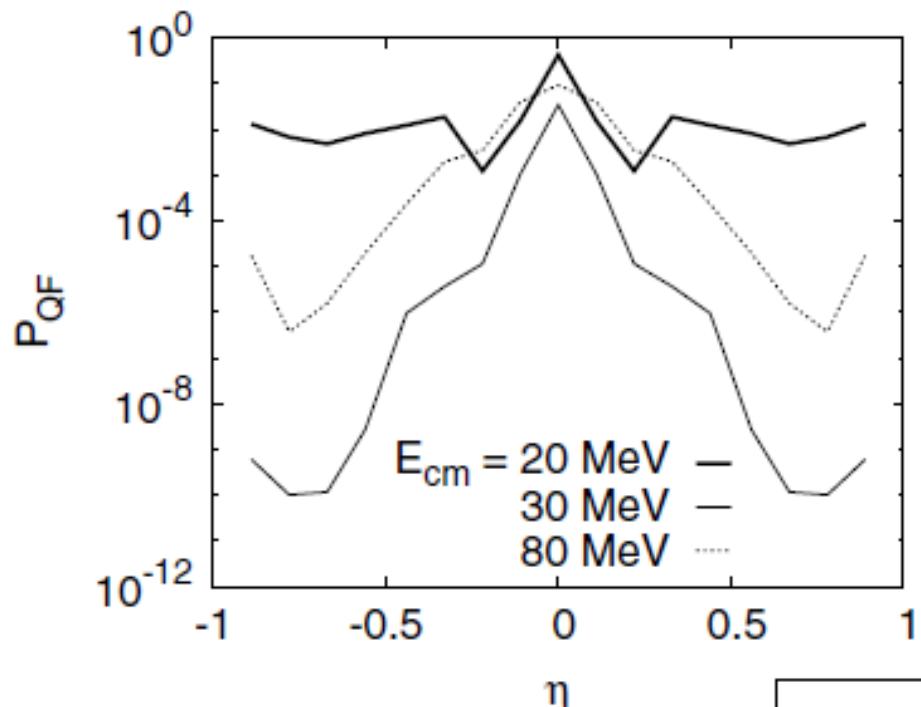
$$E_\nu^*(t) = \text{Maximum} \left[\epsilon_\nu^*(t), E_{\text{c.m.}} - \frac{\hbar^2 J(J+1)}{2\Theta_\nu} - V_{\text{dyn}}^\nu(t) \right], \quad (4)$$

$$V_{\text{dyn}}^\nu(t) = V_{\text{adiab}}^\nu[\epsilon_\nu^*(t)] + \Delta V_{\text{diab}}^\nu(t), \quad (5)$$

Probability distribution of the nuclear shapes:

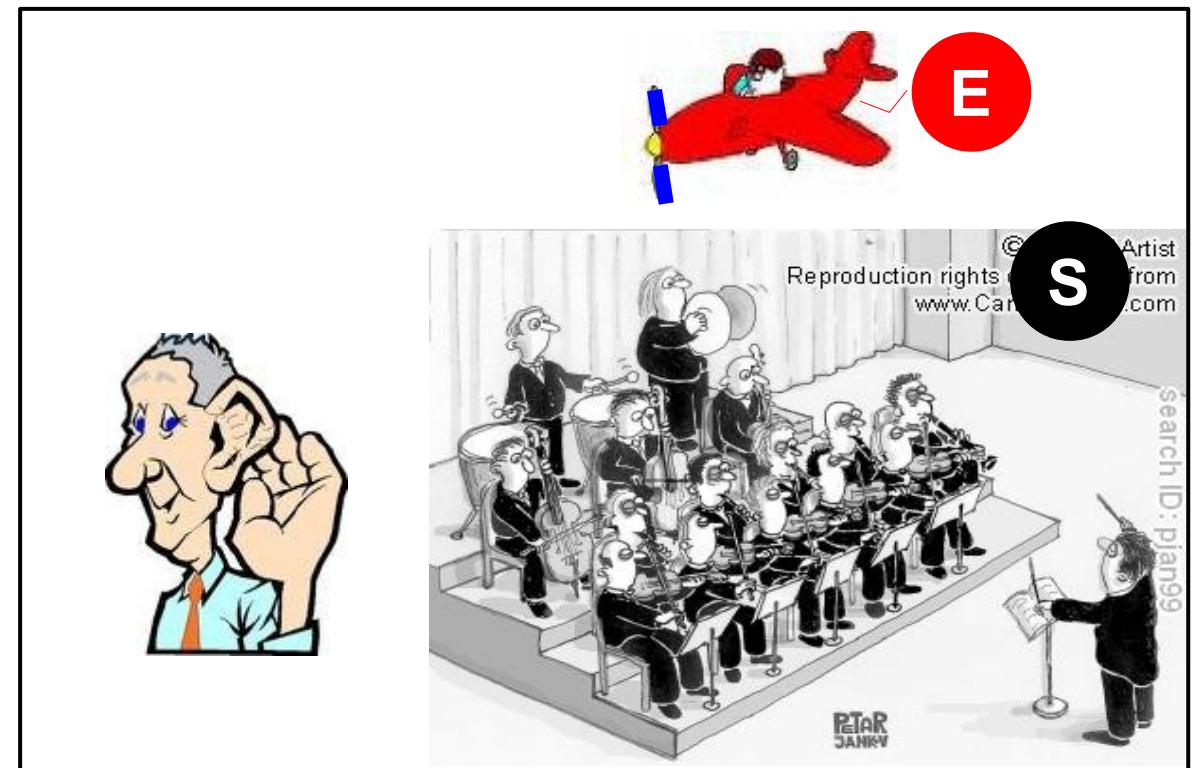
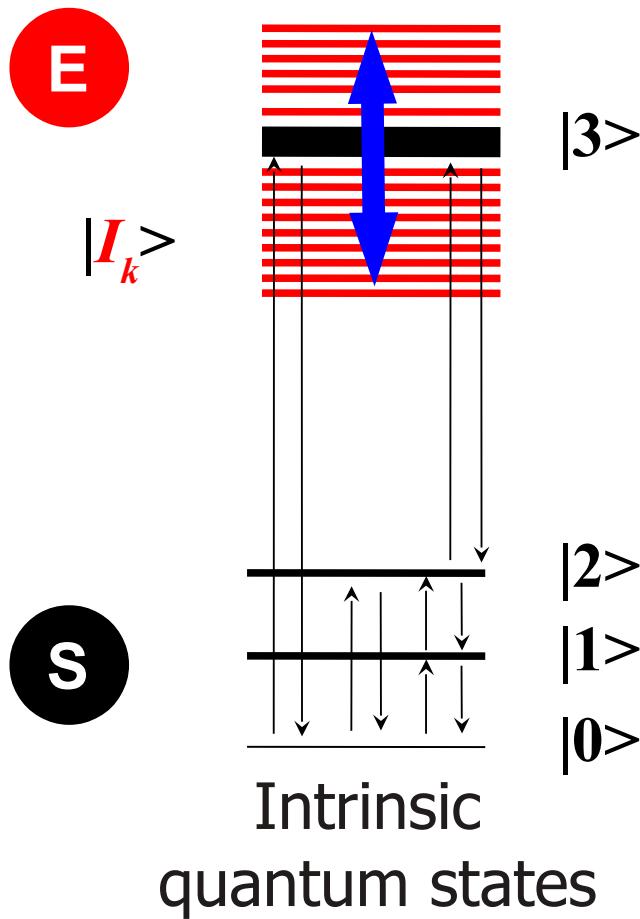
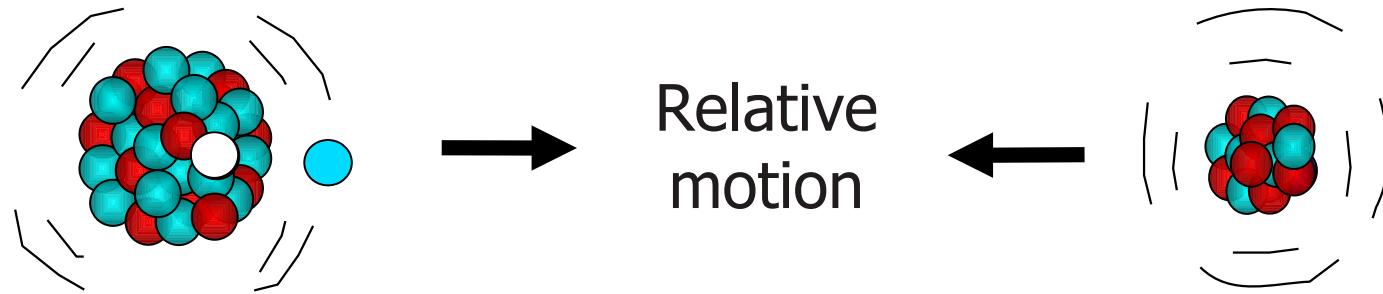


Quasi-fission Mass Distributions



Quantum Decoherence in Nuclear Collisions

AD-T, PRC 82 (2010) 054617



Can decoherence & dissipation be quantified?

Coupled-Channels Density-Matrix Framework

AD-T, PRC **82** (2010) 054617

**Hamiltonian
Dynamics**

**Decoherence
&
Dissipation**

$$i\hbar \frac{\partial \rho_{\alpha\alpha'}^{rs}}{\partial t} = \sum_t (T^{rt} \rho_{\alpha\alpha'}^{ts} - \rho_{\alpha\alpha'}^{rt} T^{ts}) + [U_\alpha(r) - U_{\alpha'}(s)] \rho_{\alpha\alpha'}^{rs} + \sum_\beta [V_{\alpha\beta}(r) \rho_{\beta\alpha'}^{rs} - \rho_{\alpha\beta}^{rs} V_{\beta\alpha'}(s)] + (\varepsilon_\alpha - \varepsilon_{\alpha'}) \rho_{\alpha\alpha'}^{rs} + i\hbar \{ \delta_{\alpha\alpha'} \sum_\mu \sqrt{\Gamma_{\alpha\mu}^{rr}} \rho_{\mu\mu}^{rs} \sqrt{\Gamma_{\alpha\mu}^{ss}} - \frac{1}{2} \sum_\mu (\Gamma_{\mu\alpha}^{rr} + \Gamma_{\mu\alpha'}^{ss}) \rho_{\alpha\alpha'}^{rs} \}$$

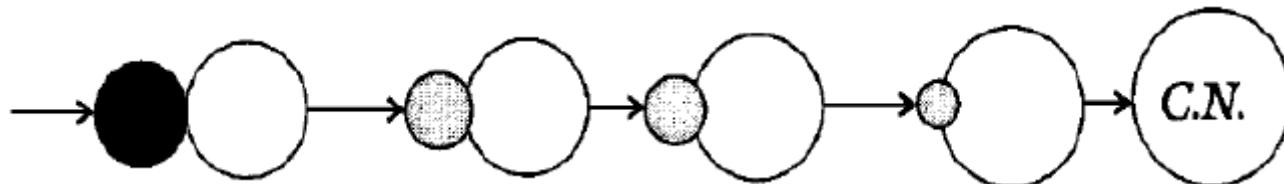
Expectation value of an observable: $\langle \hat{\mathcal{O}}(t) \rangle = \text{Tr}[\hat{\mathcal{O}} \hat{\rho}(t)]$

What do I expect as energy increases?

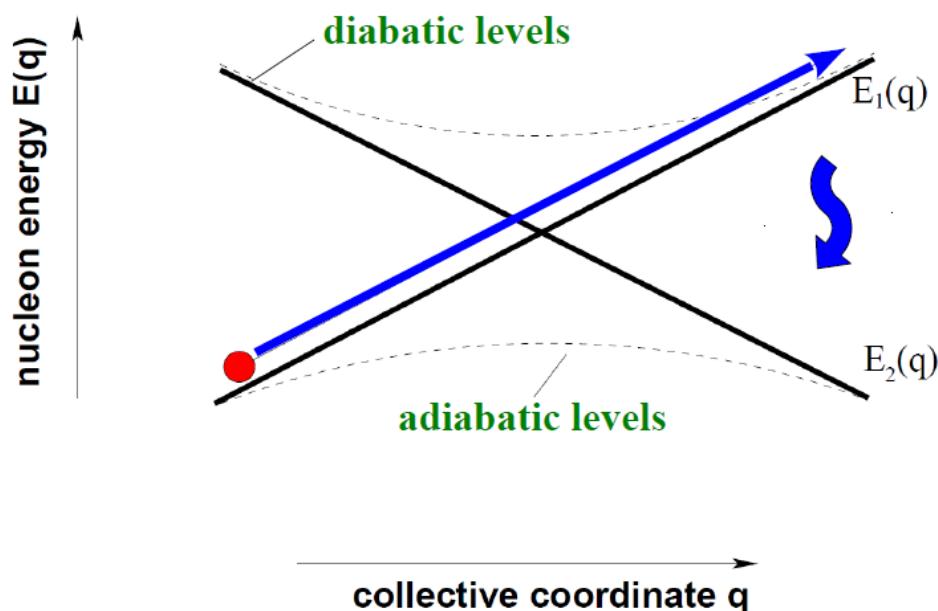
- ♦ There is **experimental** evidence of **elasto-plastic** behaviour.
See e.g., M. Rhein *et al.*, Phys. Rev. Lett. **69** (1992) 1340
 $^{208}\text{Pb} + ^{208}\text{Pb}$ @ 12 MeV/nucleon
 δ -electron spectra → Fast stopping & Long contact times
- ♦ **Elasto-plasticity** seems **not** to be included in current models of collisions at **Fermi** energies.
 - DIT model [e.g., Tassan-Got & Stephan, NPA **524** (1991) 121]
 - DNS model [e.g., Adamian *et al.*, PRC **78** (2008) 024613]
 - CoMD model [e.g., Papa, Bonasera *et al.*, PRC **64** (2001) 024612]
- ♦ **Elasto-plasticity** should be **important**.
 - Pre-equilibrium emission.
 - Interaction time *vs.* relaxation time of the dynamical potential.
 - Thermalization of reaction products by two-body collisions.

Summary

- ◆ The **diabatic picture** supports the dinuclear model for heavy-ion complete fusion in competition with quasi-fission at energies near the Coulomb barrier.



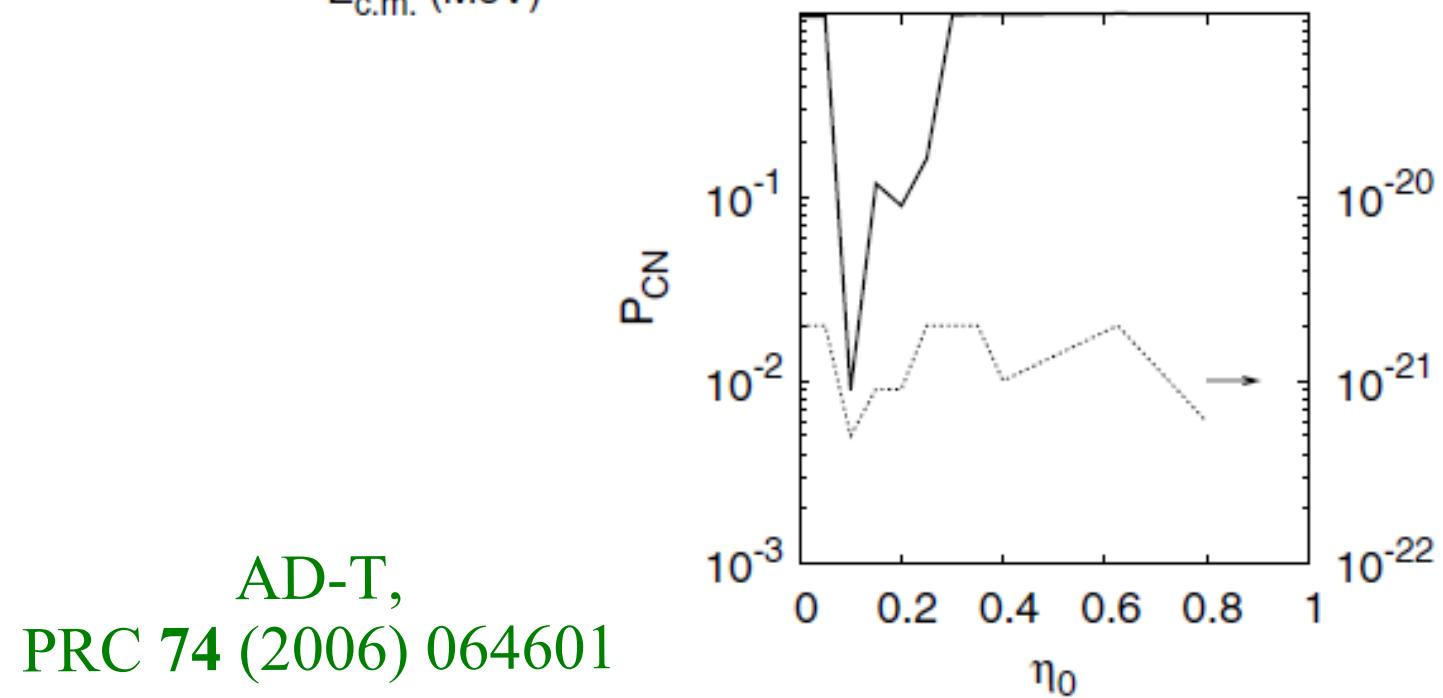
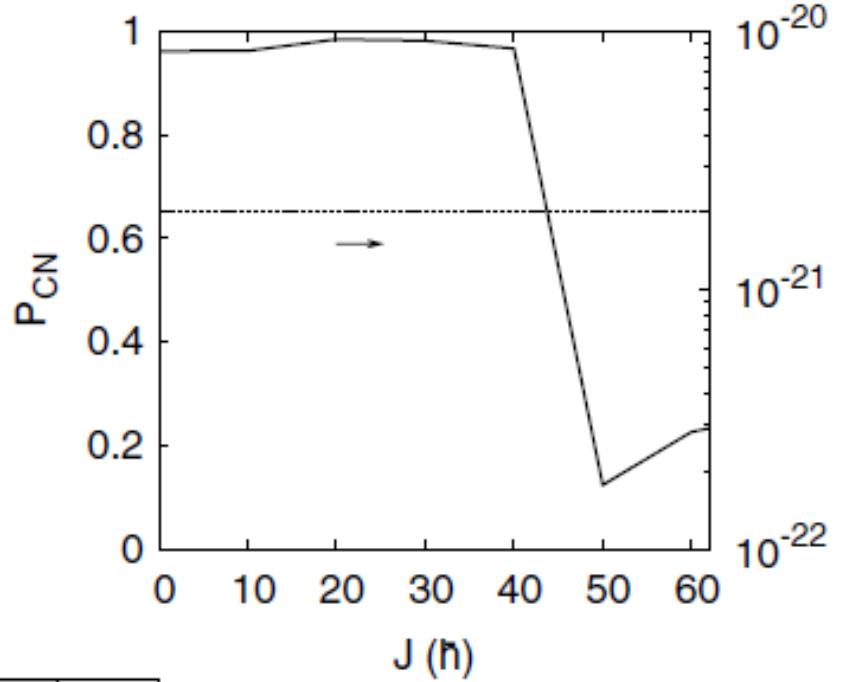
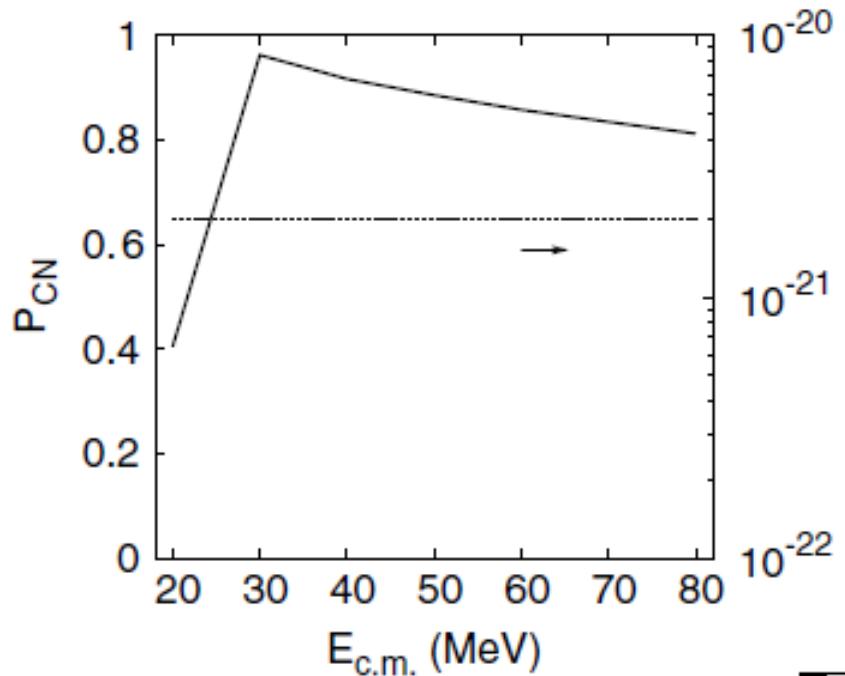
- ◆ Elasto-plastic effects on the **mass & charge distributions** of the reaction products should be important as energy increases from **Coulomb** to **Fermi** energies.



How does elasto-plasticity impact on observables?

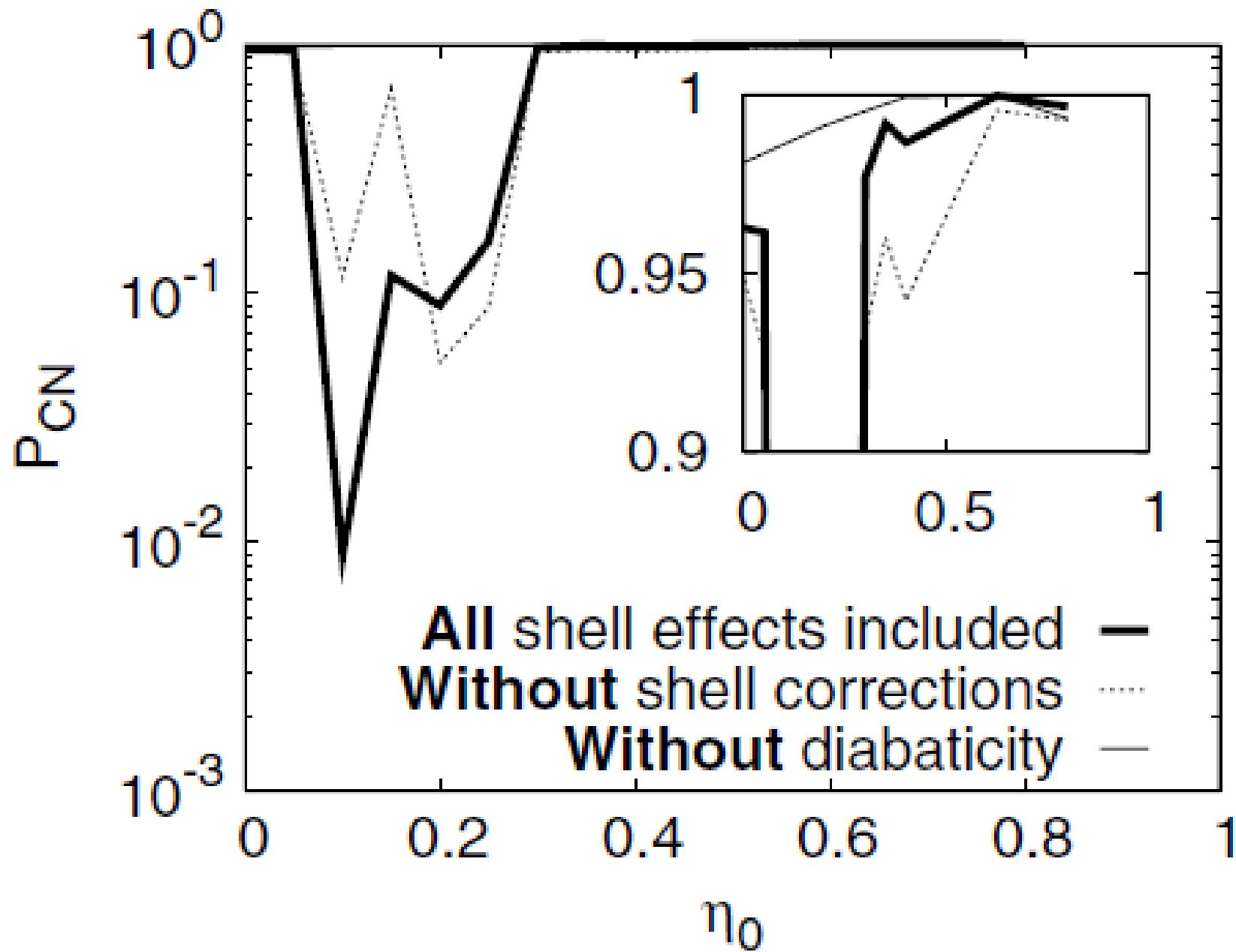
EXTRA SLIDES

Probability of Compound-Nucleus Formation



AD-T,
PRC 74 (2006) 064601

Probability of Compound-Nucleus Formation



Initial Density Matrix

$$\hat{\rho}_{\alpha\alpha'}^{rs}(0) = \underbrace{<\alpha|\hat{\sigma}(0)|\alpha'>}_{\text{Internal states}} \otimes \underbrace{g^*(z_r) g(z_s)}_{\text{Radial motion}}$$

$$\hat{\sigma}(0) = \sum_i w_i |i><i|$$

$$w_i = \delta_{ij} \quad \text{Pure state}$$

$$e^{(-\frac{E_i}{k_B T})}$$

$$w_i = \frac{e^{(-\frac{H_0}{k_B T})}}{\text{Tr}[e^{(-\frac{H_0}{k_B T})}]} \quad \text{Mixed state}$$