



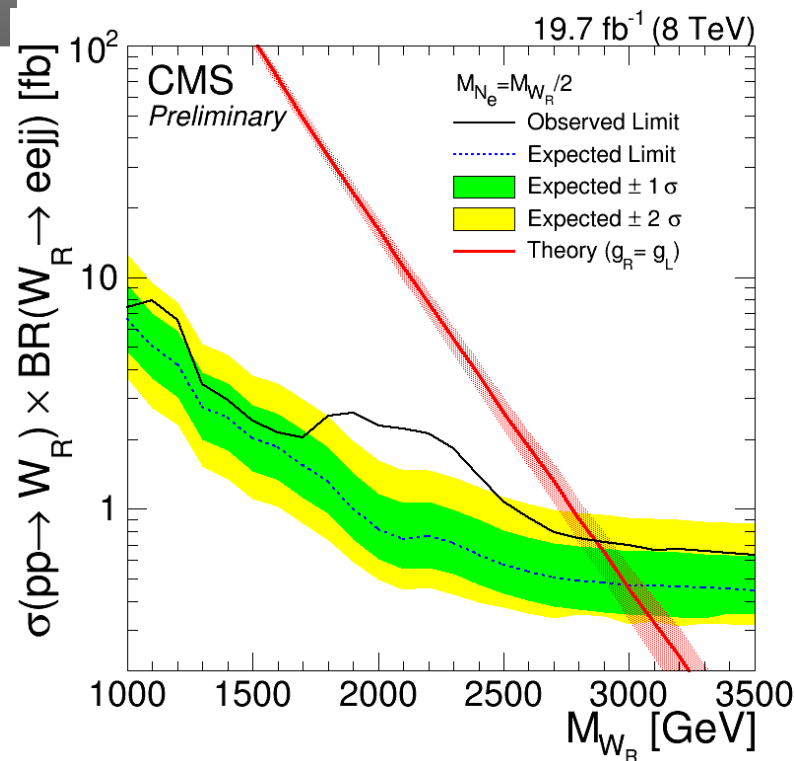
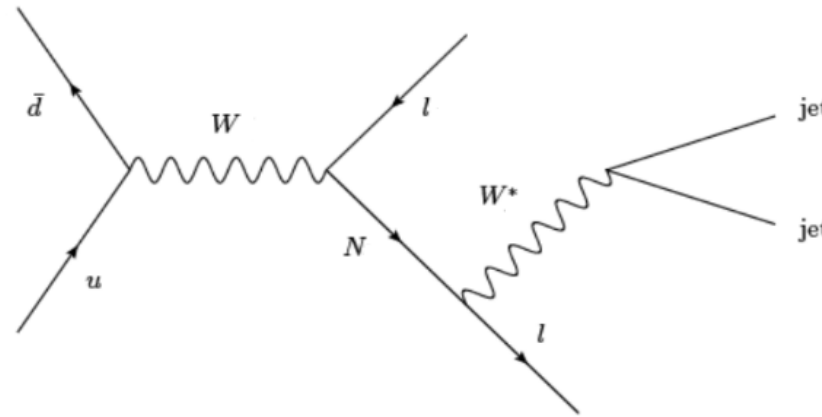
Hunting for heavy composite Majorana neutrinos at the LHC

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Search of Heavy Neutrinos in CMS



Excess over the SM expectation in the eejj channel.

The Analysis was based on 19.7 fb⁻¹ of integrated luminosity and $\sqrt{s} = 8$ TeV. It reports a 2.8 σ excess in the eejj invariant mass distribution in the interval $1.8 \text{ TeV} < M_{eejj} < 2.2 \text{ TeV}$

Composite models for quarks and leptons

- The proliferation of standard model fermions can suggest a further composition
- If quarks and leptons are composite we expect:
 - ⇒ **Excited leptons and quarks**
 - ⇒ **Contact interaction between four fermions**: a residual interaction of an interaction between the constituent particles
 - Eichten, Lane e Peskin, Phys. Rev. B 50, 811 (1983)
 - Cabibbo, Maiani e Srivastava, Phys Lett. B 149, 459 (1984)
 - Baur, Spira e Zerwas, Phys. Rev. D 42, 815 (1990)

Extended weak isospin model

Pancheri and Srivastava, Phys. Lett. B 146 (1984).

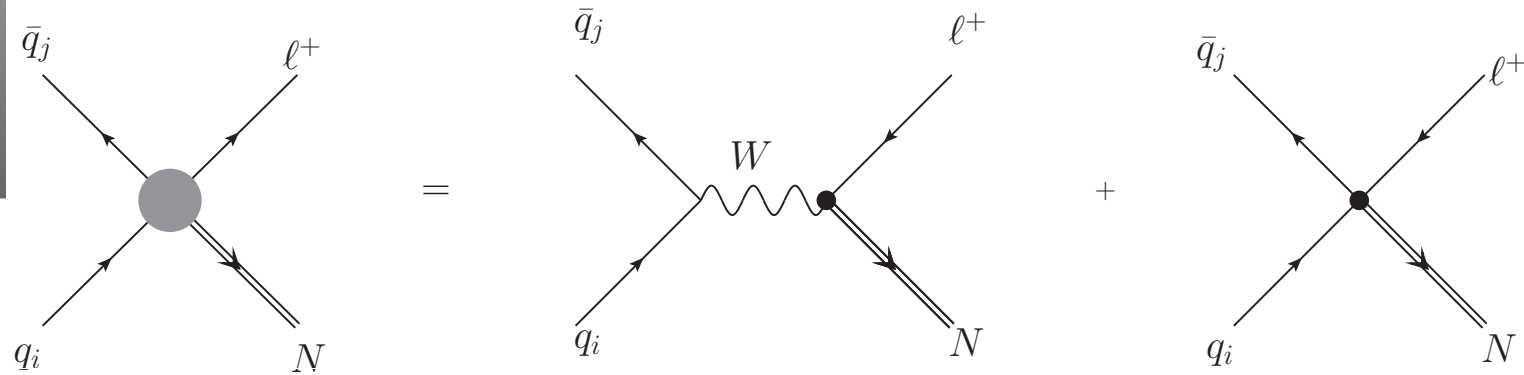
- Fermions' composition through the weak isospin symmetry
- It doesn't refer to the internal dynamics
- analogy with strong isospin \rightarrow prediction of hadronic states before the discovery of quarks and gluons
- SM $q, \ell \in I_W = 0, 1/2$ and $W^\pm, Z^0, \gamma \in I_W = 0, 1$
 \Rightarrow **excited fermions** $\in I_W \leq 3/2$

Extended weak isospin model

Multiplets of the model

I_W	Multiplet	Q	Y	Coupled to	I_W	Multiplet	Q	Y	Coupled to
0	E^-	-1	-2	e_R through B^μ	0	(i) U	2/3	4/3	u_R through B^μ and G^μ
1/2	$\epsilon \equiv \begin{pmatrix} E^0 \\ E^- \end{pmatrix}$	0	-1	$\ell_L = \begin{pmatrix} \nu_e \\ e \end{pmatrix}_L$ through W^μ and B^μ	1/2	(ii) D	-1/3	-2/3	d_R through B^μ and G^μ
1	$\epsilon \equiv \begin{pmatrix} E^0 \\ E^- \\ E^{--} \end{pmatrix}$	0	-2	e_R through W	1/2	$\Psi \equiv \begin{pmatrix} U \\ D \end{pmatrix}$	2/3	1/3	$q_L = \begin{pmatrix} u \\ d \end{pmatrix}_L$ through W B^μ and G^μ
3/2	$\epsilon_M \equiv \begin{pmatrix} E^+ \\ E^0 \\ E^- \\ E^{--} \end{pmatrix}$	1	-1	$\ell_L = \begin{pmatrix} \nu_e \\ e \end{pmatrix}_L$ through W	1	(i) $U \equiv \begin{pmatrix} U_+ \\ U \\ D \end{pmatrix}$	5/3	+4/3	u_R through W
		0				(ii) $D \equiv \begin{pmatrix} U \\ D \\ D_- \end{pmatrix}$	2/3	-2/3	d_R through W
		-1					-1/3		
		-2			3/2	$\Psi_M \equiv \begin{pmatrix} U_+ \\ U \\ D \\ D_- \end{pmatrix}$	5/3	1/3	$q_L = \begin{pmatrix} u \\ d \end{pmatrix}_L$ through W
							2/3		
							-1/3		
							-4/3		

Contact and gauge interactions



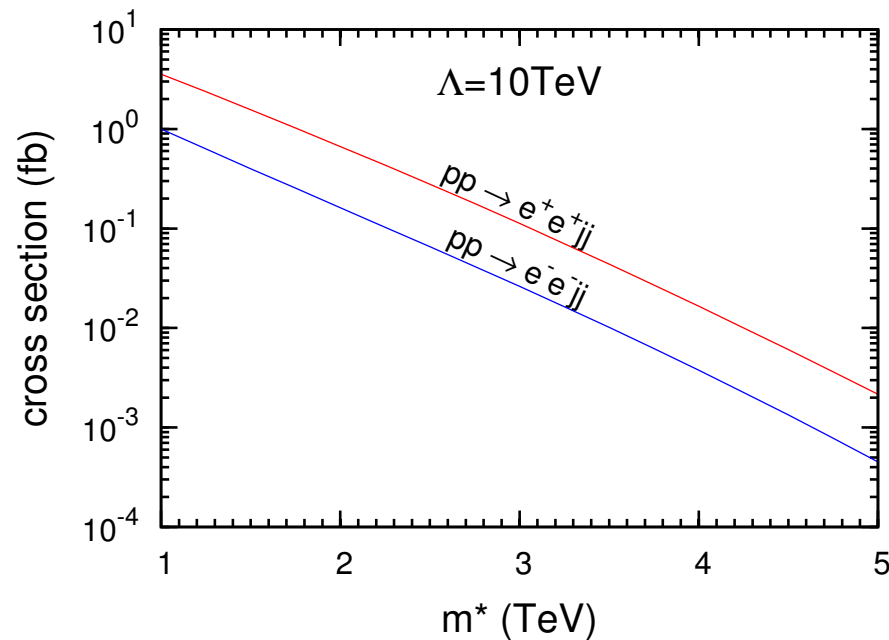
$$\mathcal{L}_G = \frac{gf}{\sqrt{2}\Lambda} \bar{N} \sigma_{\mu\nu} \ell_L \partial^\nu W^\mu + h.c.$$

$$\mathcal{L}_{CI} = \frac{g_*^2}{\Lambda^2} \bar{q}_L \gamma^\mu q'_L \bar{N}_L \gamma_\mu \ell_L$$

We consider a Majorana neutrino

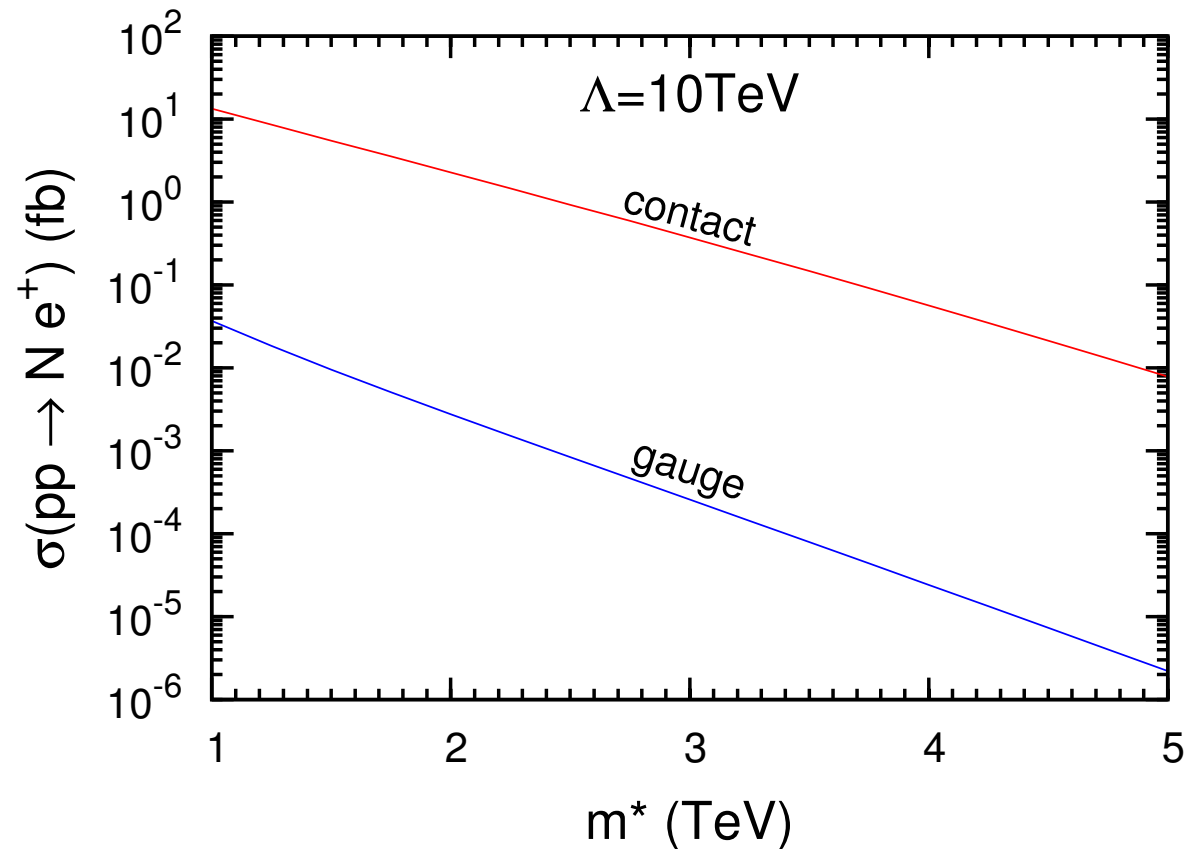
The same-sign dilepton is a peculiar final state for Majorana neutrinos

We choose the positive same-sign dilepton due to its larger cross section



Production of the heavy Majorana neutrino

$$pp \rightarrow N \ell$$



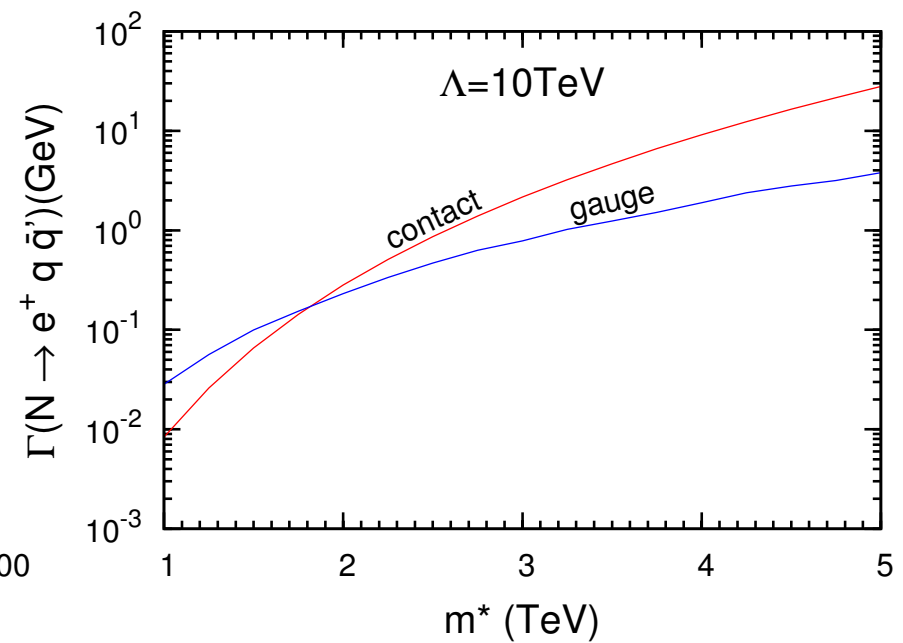
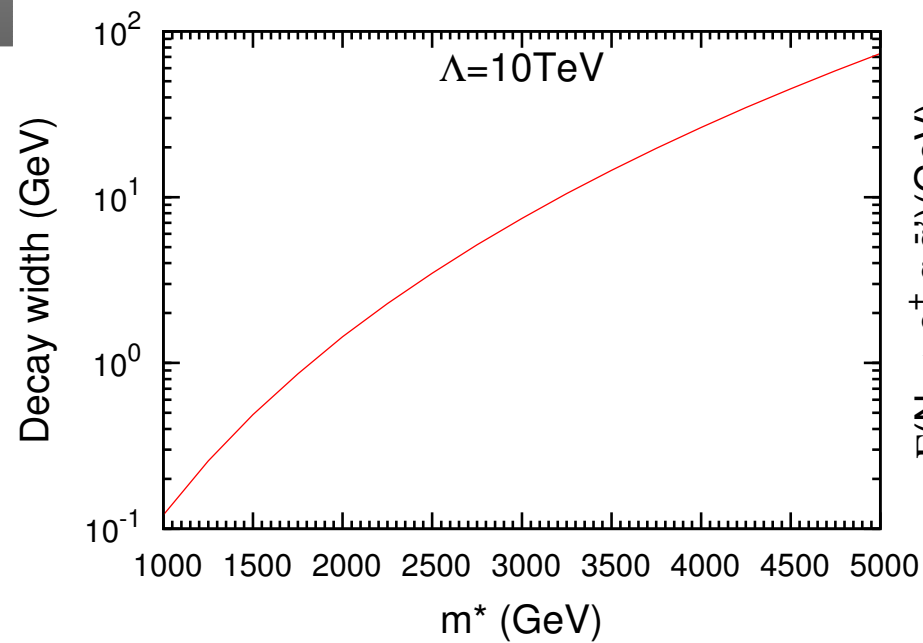
The contact interaction is the dominant one

Decays of the heavy Majorana neutrino

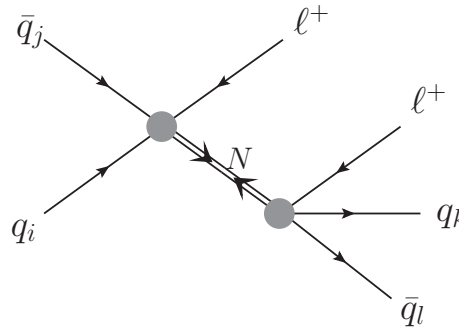
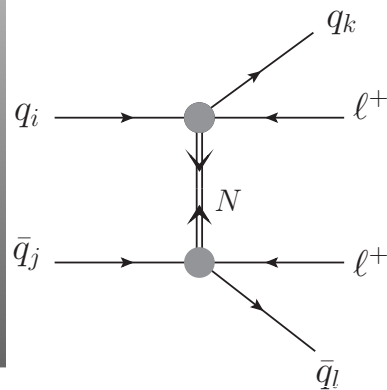
$$N \rightarrow \ell q \bar{q}'$$

$$N \rightarrow \ell^+ \ell^- \nu(\bar{\nu})$$

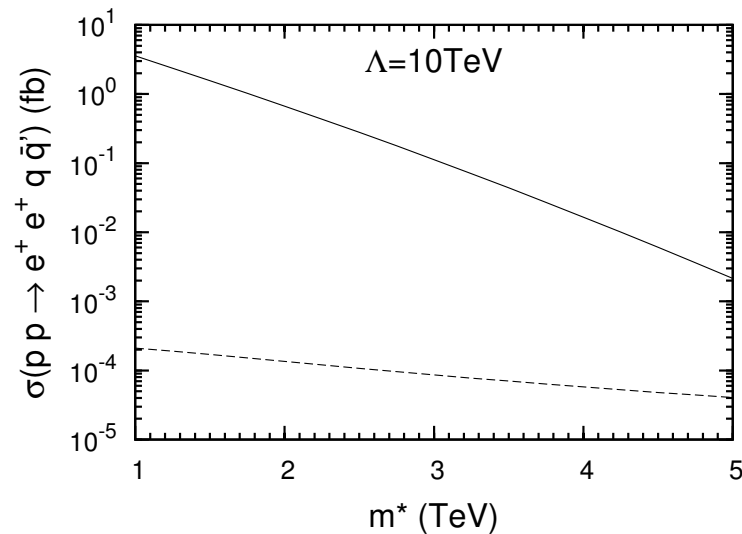
$$N \rightarrow \nu(\bar{\nu}) q \bar{q}'$$



Processes under examination



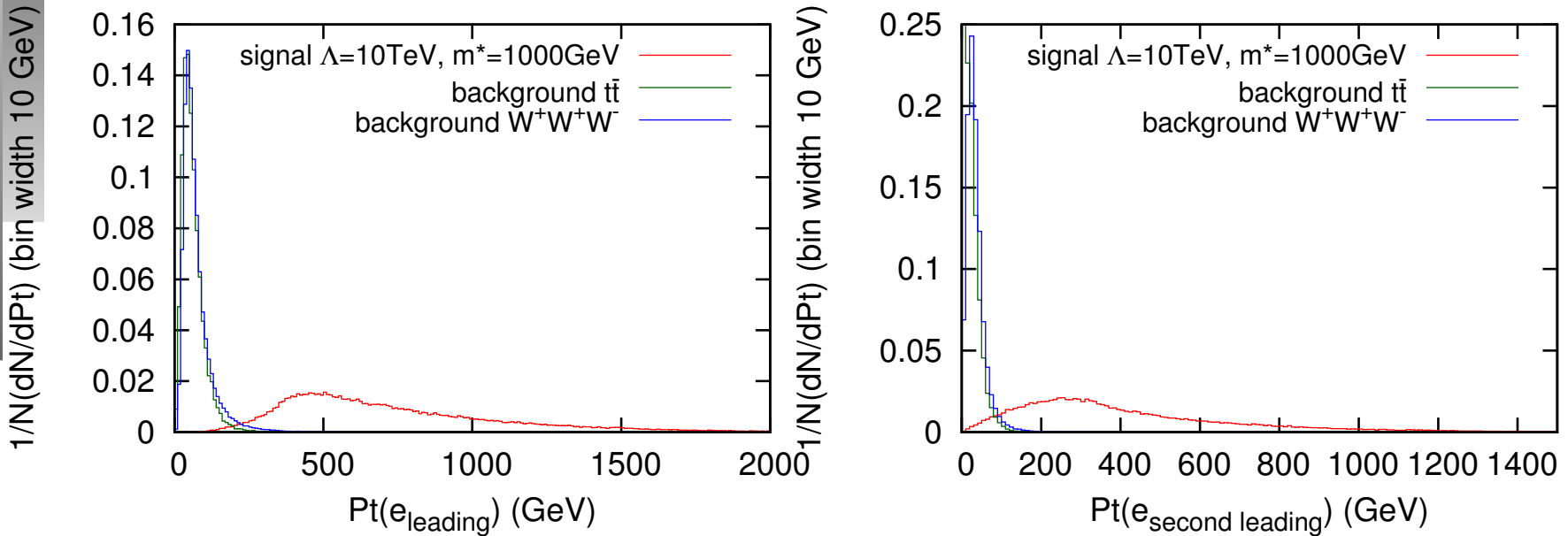
Exchange of virtual heavy Majorana neutrino (left),
resonant production of heavy Majorana neutrino (right)



The resonant production is the dominant process

- In SM the lepton number is conserved, so processes like the our ($\Delta L = 2$) are not allowed
- However in the SM there are several processes that can produce same sign leptons in association with jets
- the main backgrounds are:
 - $pp \rightarrow t\bar{t} \rightarrow \ell^+\ell^+\nu\nu jets$
 - $pp \rightarrow W^+W^+W^- \rightarrow \ell^+\nu\ell^+\nu jj$

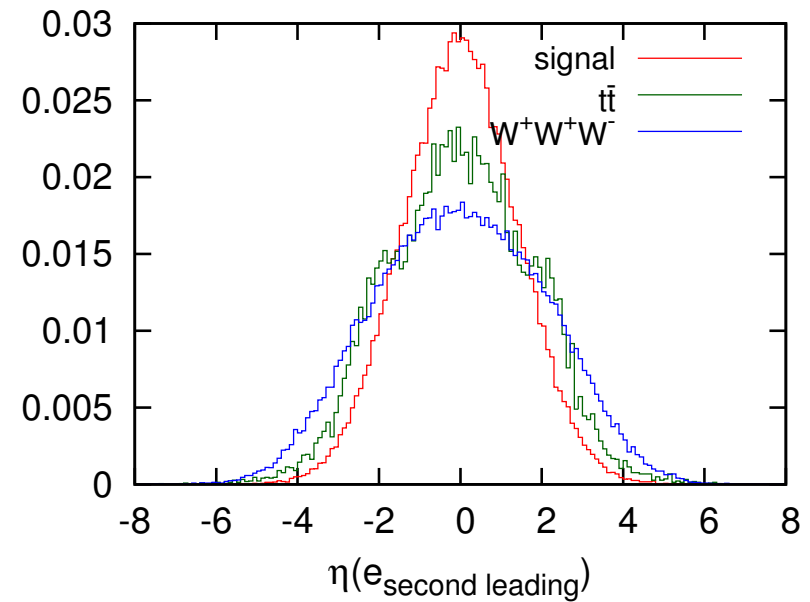
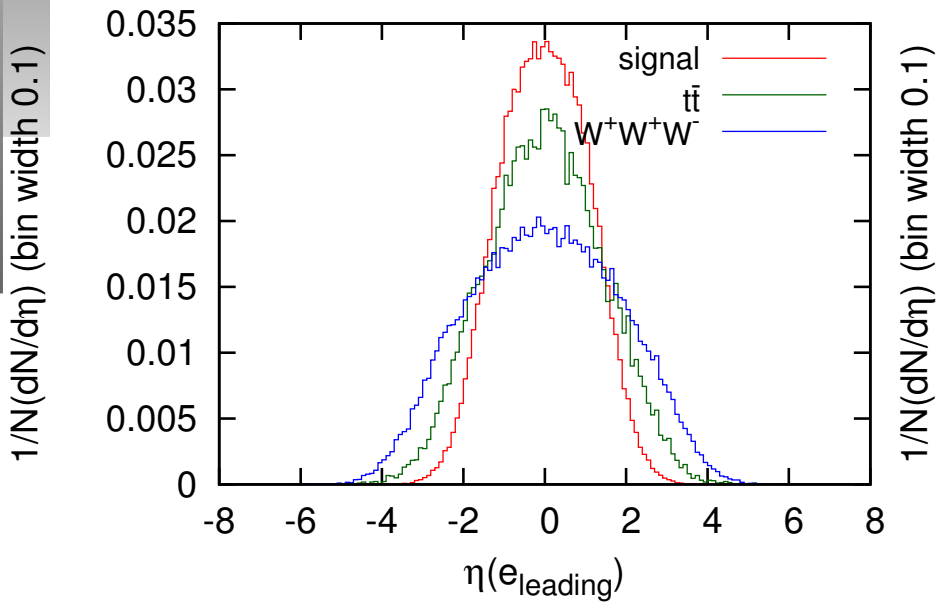
Kinematical distributions: Transverse momentum



We can reduce drastically the background with the cuts:

- $p_T(e_{\text{leading}}^+) \geq 200 \text{ GeV}$
- $p_T(e_{\text{second-leading}}^+) \geq 100 \text{ GeV}$

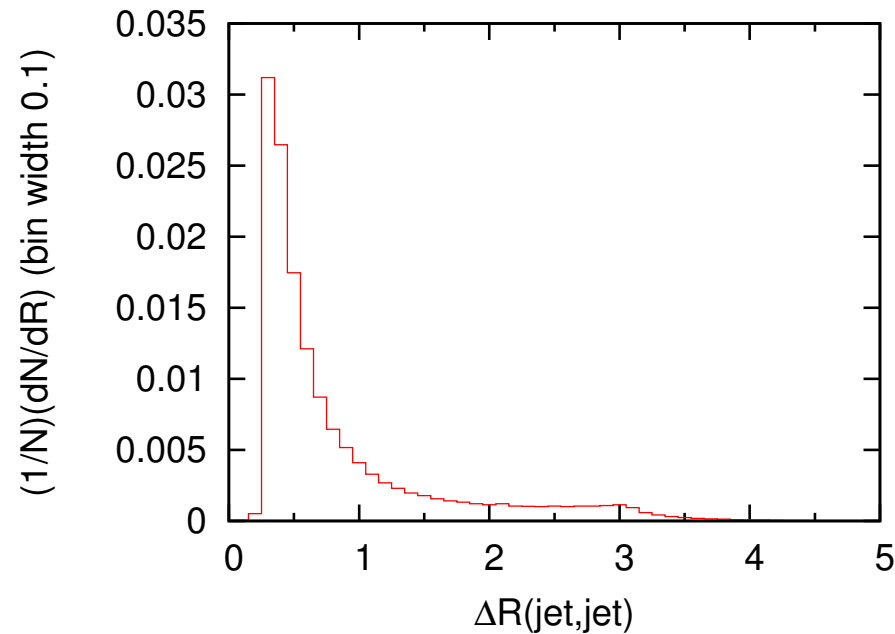
Kinematical distributions: *pseudorapidity*



The pseudorapidity is not very selective

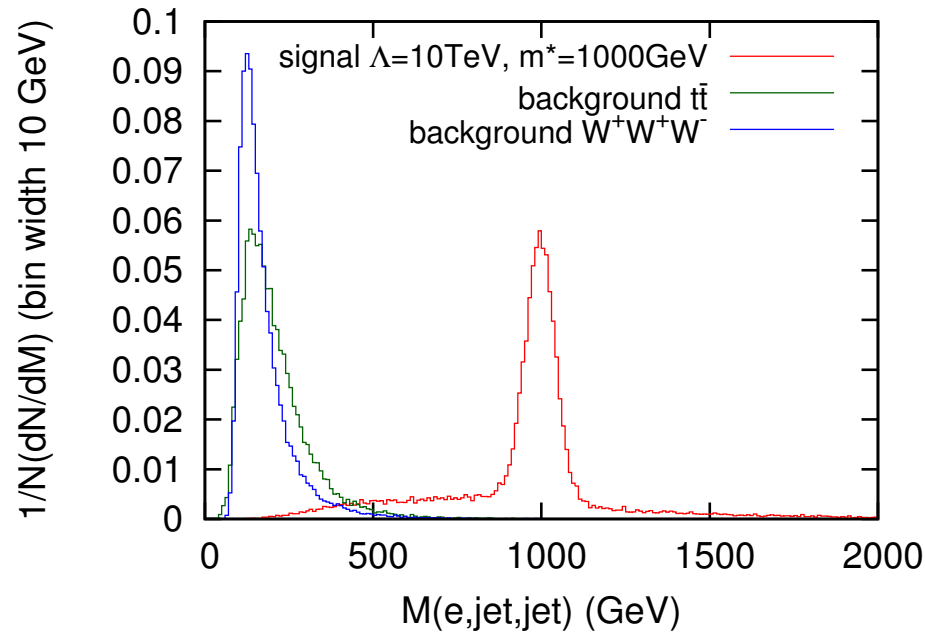
Kinematical distributions:

$$\Delta R(\text{jet}, \text{jet})$$



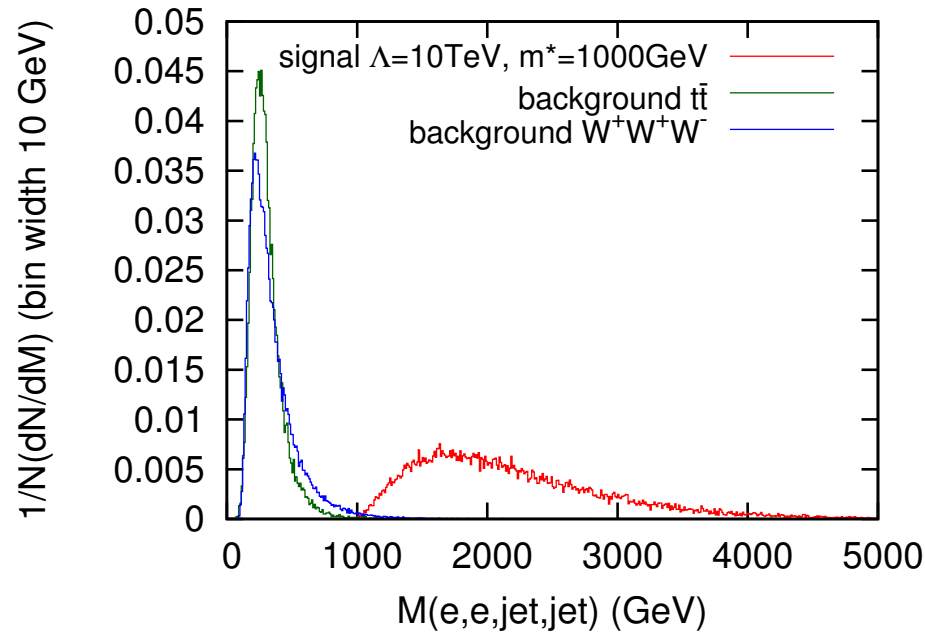
- A large fraction of the events have two jets with very small separation
- In the reconstruction process we will have merging

Kinematical distributions: $M(e,j,j)$



The invariant mass of the second leading electron plus the two jets give informations about the mass of heavy neutrino

Kinematical distributions: $M(e,e,j,j)$



The eejj invariant mass can easily accommodate the excess in the interval observed by CMS

⇒ LHE files generation by CalcHEP for signal and background

- Scan in parameters space:
 - $\Lambda \in [8, 40]$ TeV with step of 1 TeV
 - $m^* \in [500, 5000]$ GeV with step of 250 GeV
- 100000 events for each LHE file

⇒ DELPHES simulates the particles reconstruction by the detector considering

- efficiency
- geometrical acceptance

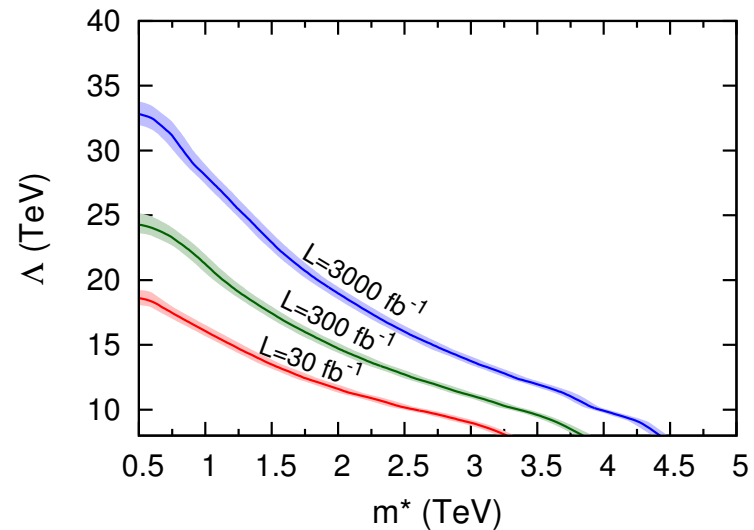
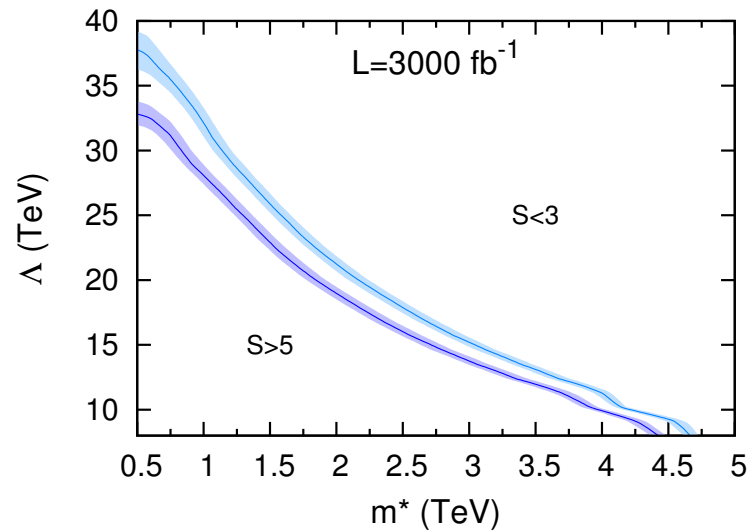
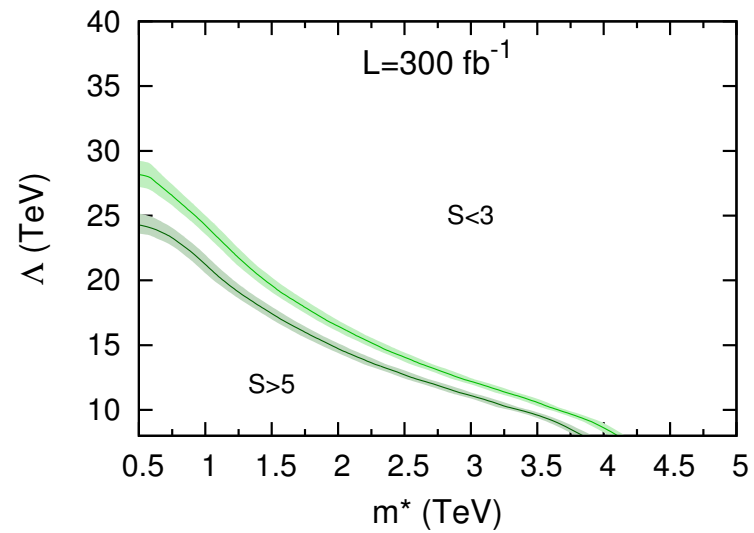
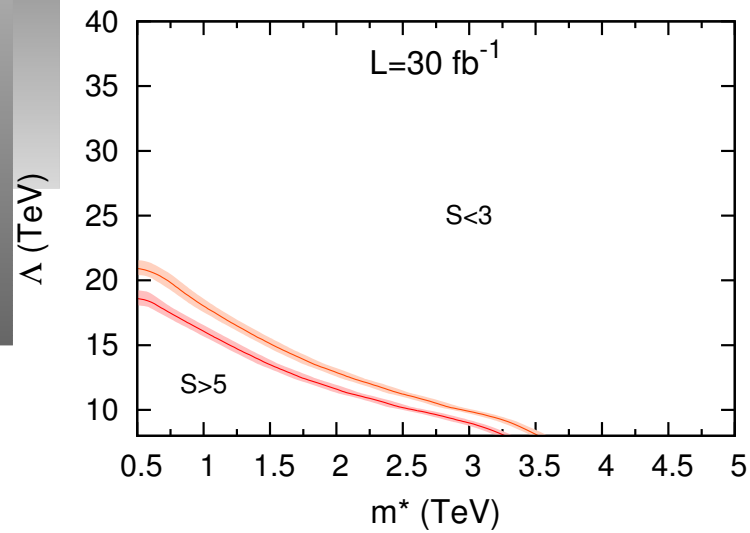
- Selection criteria:

- Presence of two e^+
- Kinematical cuts:

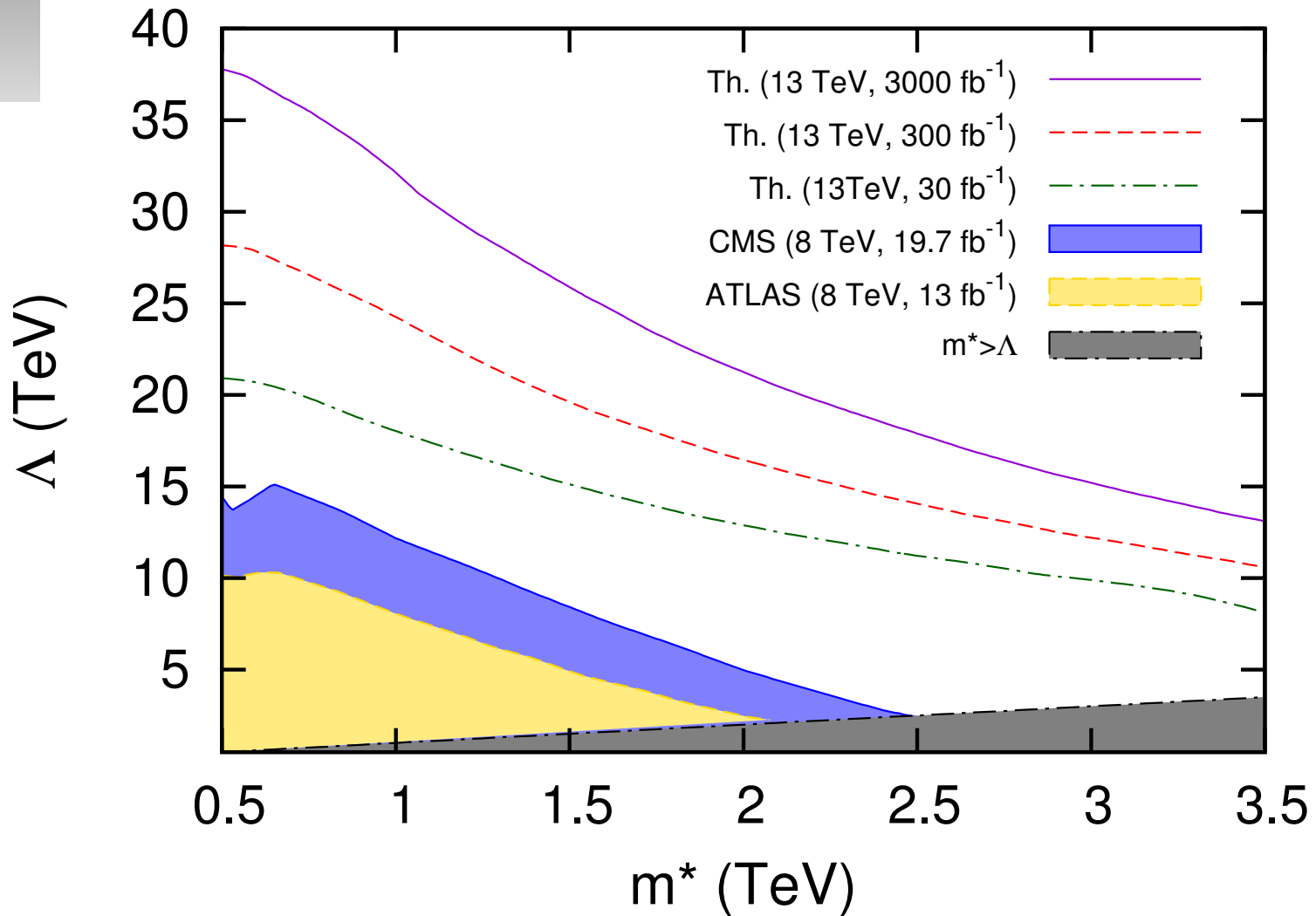
$$p_T(e_{\text{leading}}^+) \geq 200 \text{ GeV}, p_T(e_{\text{second-leading}}^+) \geq 100 \text{ GeV}$$

- Determination of reconstruction efficiencies
- Determination of expected number of events for signal and background: $N_s = L\sigma_s\epsilon_s$, $N_b = L\sigma_b\epsilon_b$
- Determination of statistical significance: $S = \frac{N_s}{\sqrt{N_b}}$
- Determination of contour plots in the parameter space (Λ, m^*) at $S = 3$ and $S = 5$

Contour plots



Contour plots



Conclusions

- We have performed a phenomenological study about a heavy neutrino in view of the recent observation by CMS of the excess in $eejj$ channel
- We considered the composite model scenario and an excited Neutrino of Majorana type
- We found that the invariant mass distribution of the second leading electron and the two jets is highly correlated to the heavy neutrino mass
- We provide the contour plots of Statistical significance at 3- and 5-sigma and compare them to the experimental data of Run-I, showing a great potential of discovery or improving the current bounds in $eejj$ signature from a heavy composite Majorana neutrino