



***Quantum Gravity,
from the atoms of space to cosmology***

The Universe as a Quantum Condensate

Daniele Oriti

Albert Einstein Institute

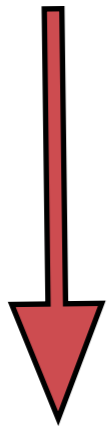
Dipartimento di Fisica e Astronomia
Universita' di Bologna
24/07/2015



Quantum spacetime: the difficult path from microstructure to macrophysics

Quantum Gravity problem:

identify microscopic d.o.f. of quantum spacetime and their fundamental dynamics



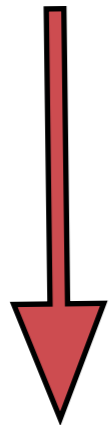
derive effective (QG-inspired) models for macroscopic continuum physics:
explain features of early Universe, obtain testable QG predictions

Quantum spacetime: the difficult path from microstructure to cosmology

Quantum Gravity problem:

identify microscopic d.o.f. of quantum spacetime and their fundamental dynamics

various approaches: group field theory, loop quantum gravity



derive effective (QG-inspired) models for (quantum) cosmology:
explain features of early Universe, obtain testable QG predictions

various models: loop quantum cosmology,

task is daunting (compare with analogue problem in condensed matter theory)

Group Field Theory: convergence of approaches

Group Field Theory: convergence of approaches

Matrix models

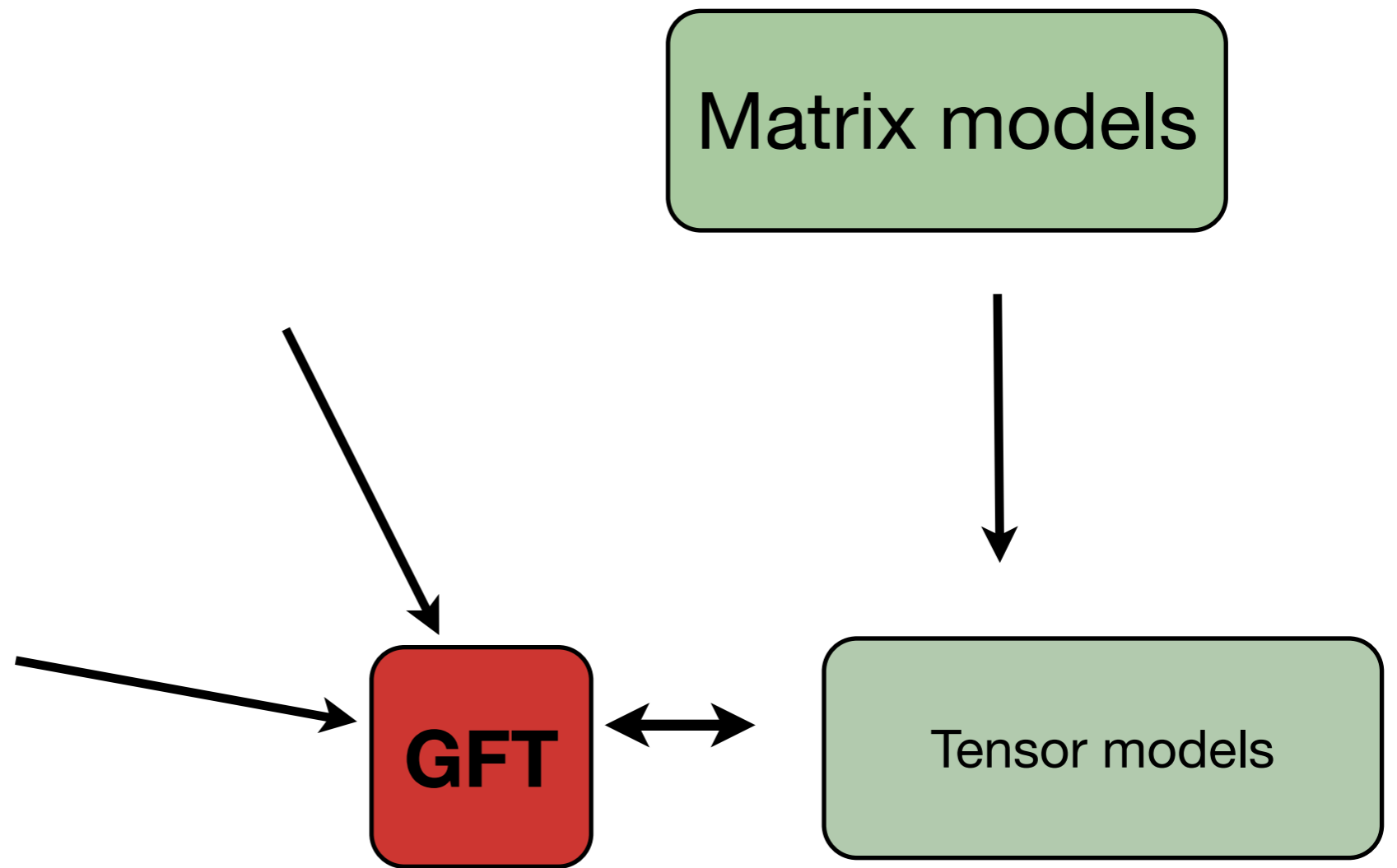
Group Field Theory: convergence of approaches

Matrix models

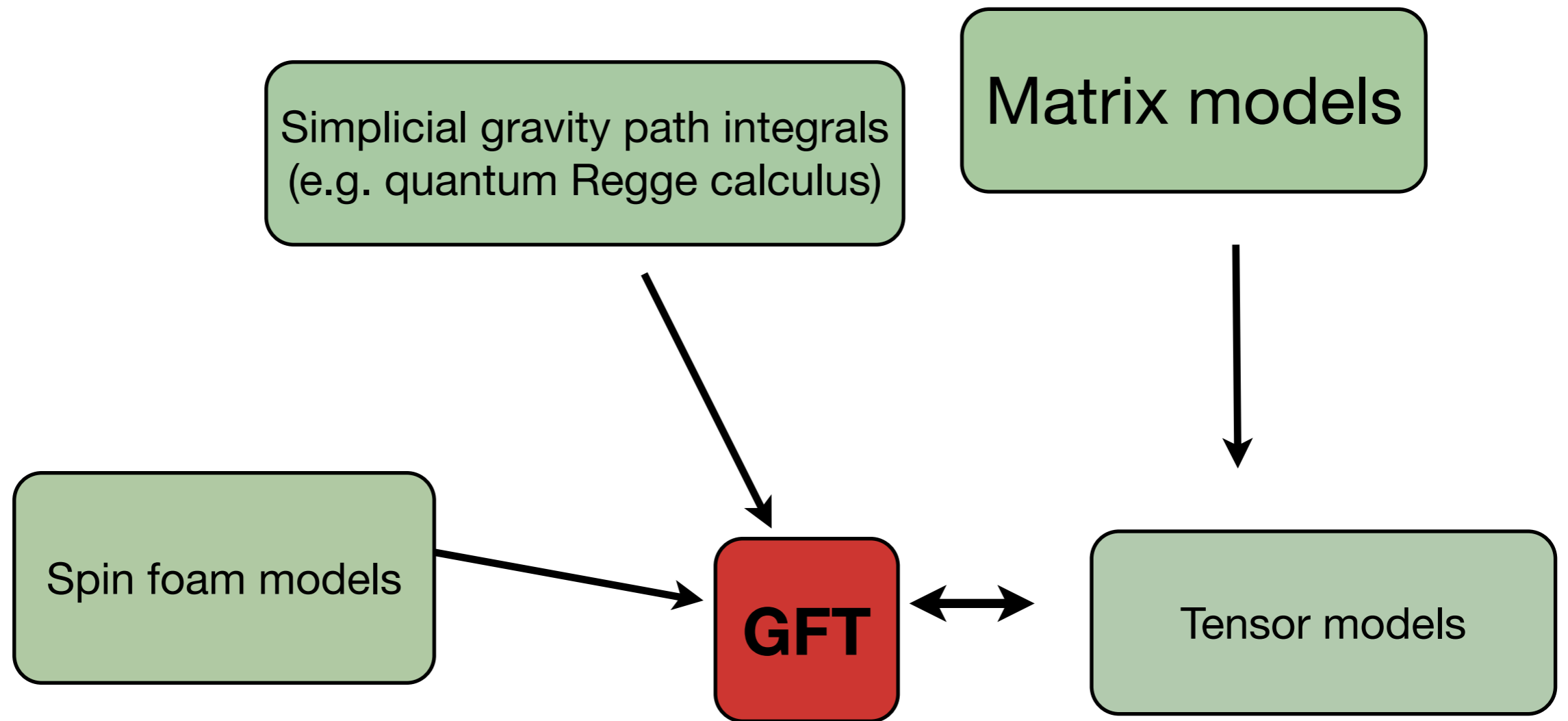


Tensor models

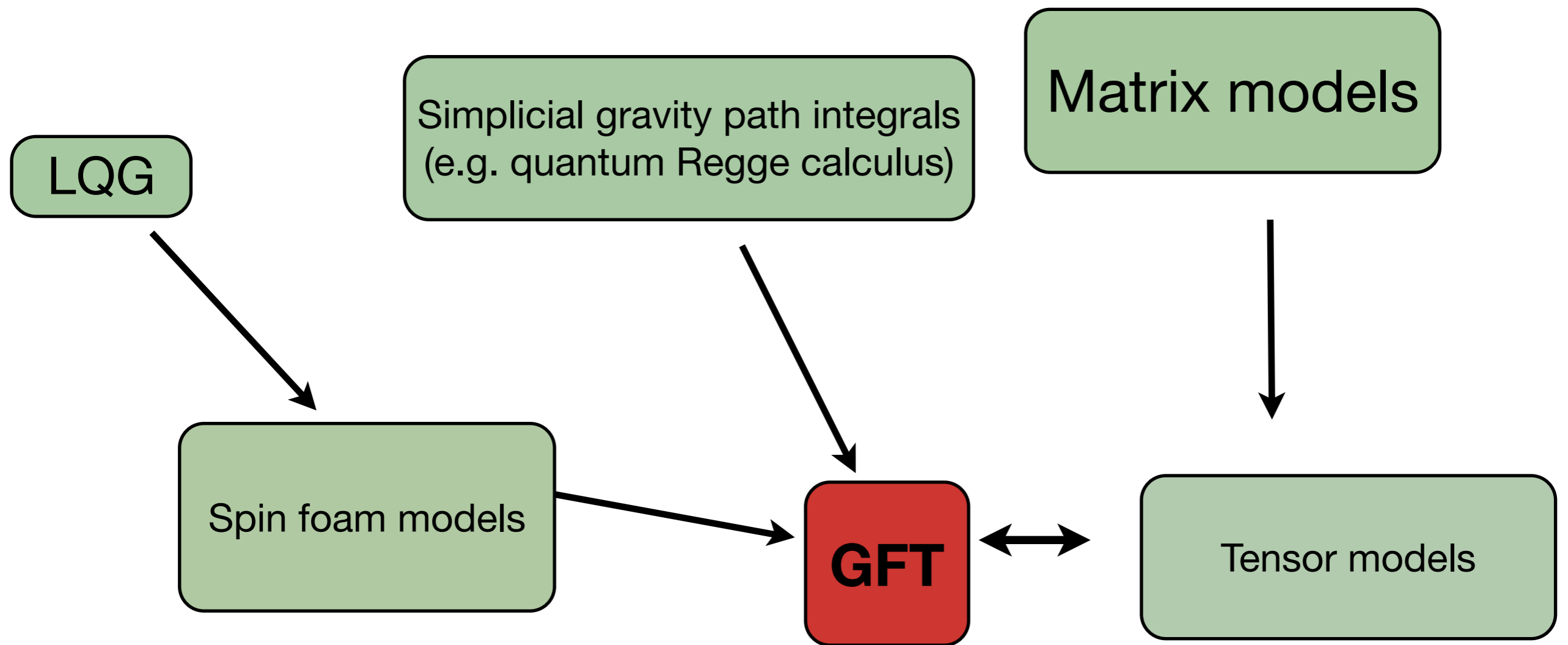
Group Field Theory: convergence of approaches



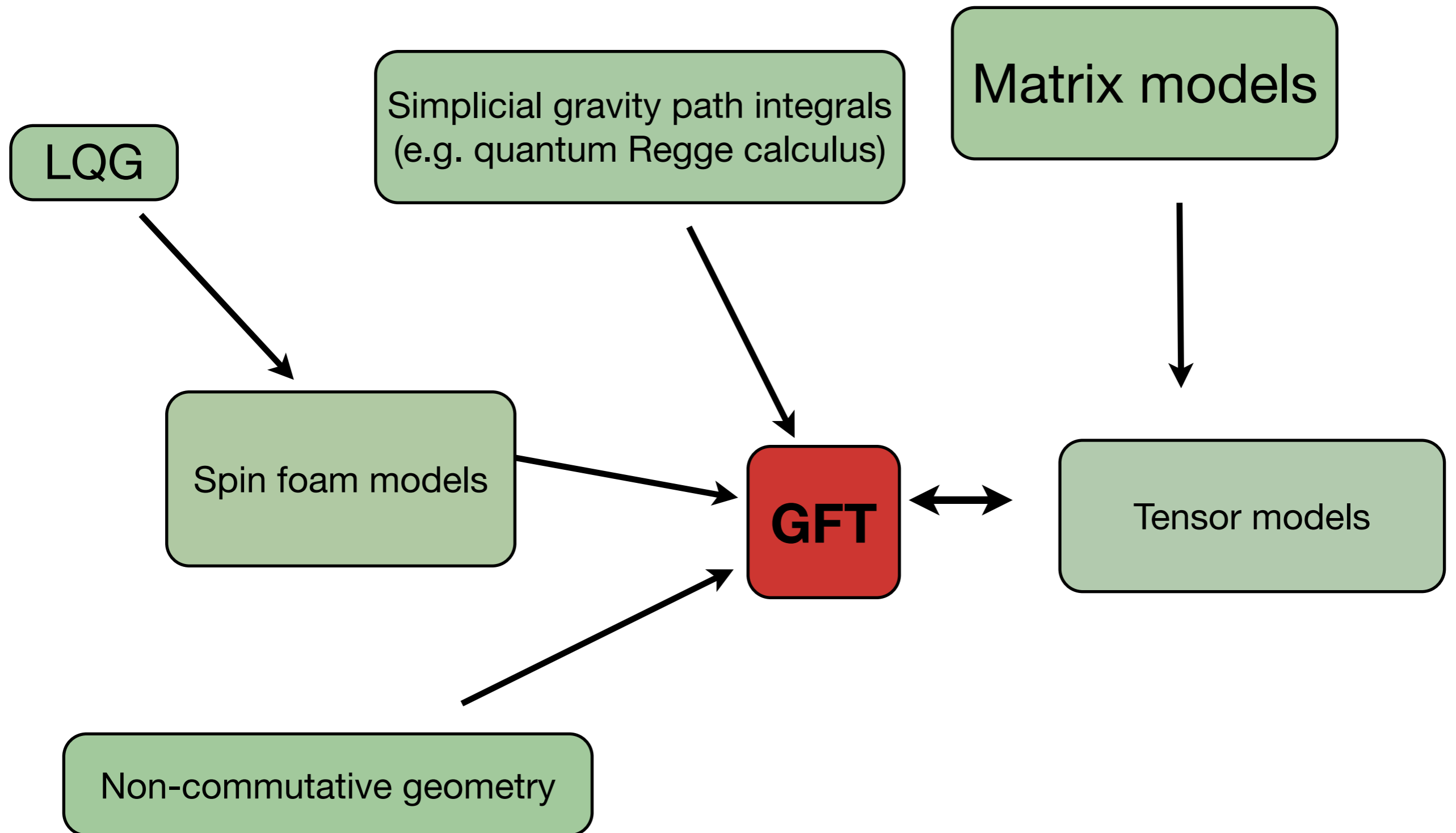
Group Field Theory: convergence of approaches



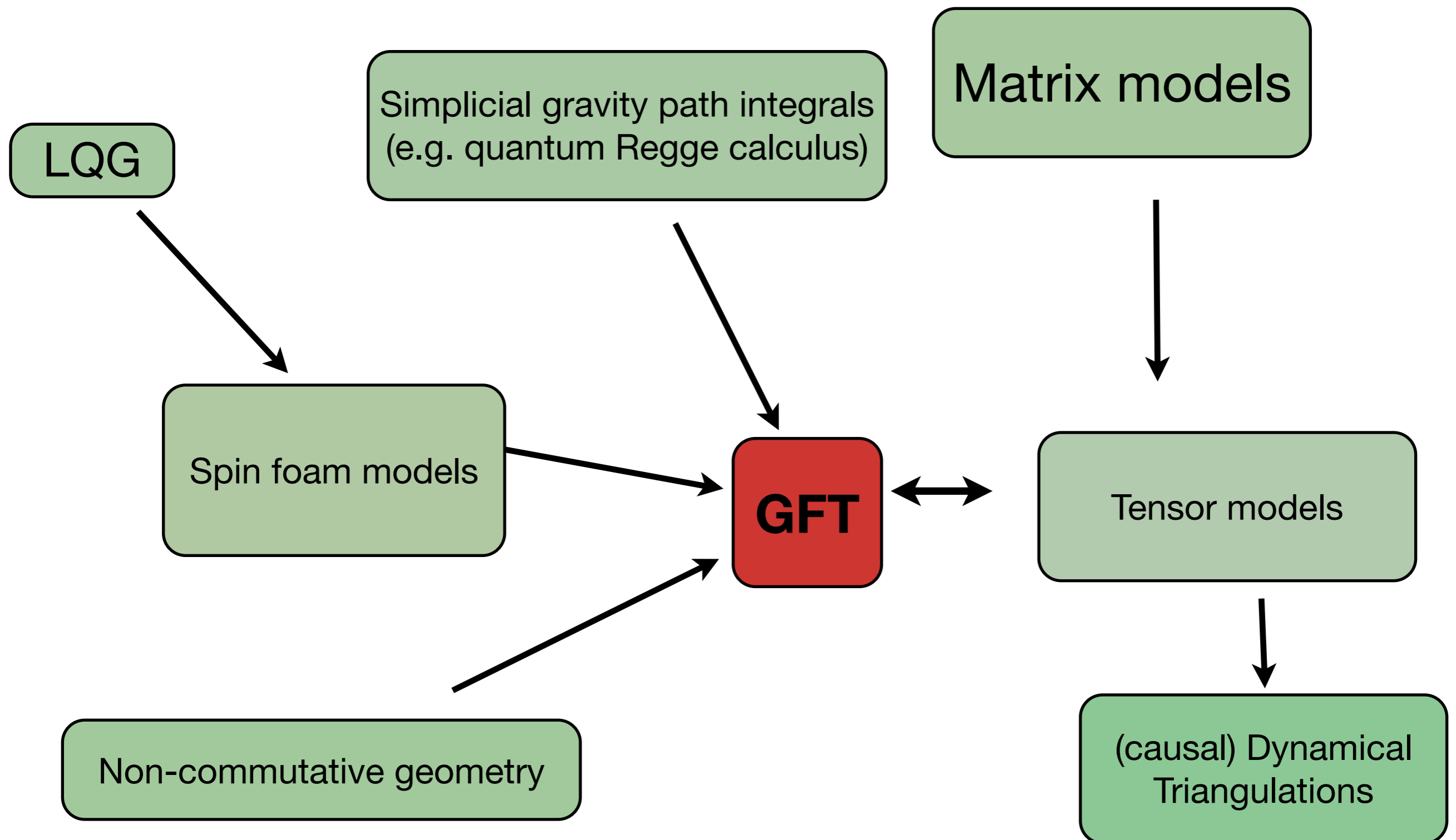
Group Field Theory: convergence of approaches



Group Field Theory: convergence of approaches



Group Field Theory: convergence of approaches



Group field theories

Quantum field theories over group manifold G (or corresponding Lie algebra) $\varphi : G^{\times d} \rightarrow \mathbb{C}$

relevant classical phase space for “GFT quanta”: $(\mathcal{T}^*G)^{\times d} \simeq (\mathfrak{g} \times G)^{\times d}$

can reduce to subspaces in specific models depending on conditions on the field

d is dimension of “spacetime-to-be”

example: $d=4$ $\varphi(g_1, g_2, g_3, g_4) \leftrightarrow \varphi(B_1, B_2, B_3, B_4) \rightarrow \mathbb{C}$

can be defined for any (Lie) group and dimension d , any signature,

very general framework; interest rests on specific models/use

Group field theories

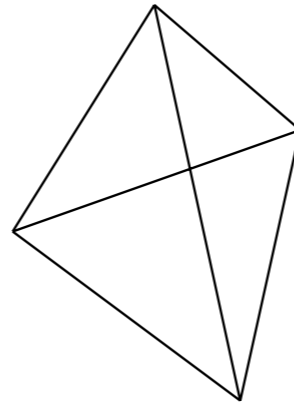
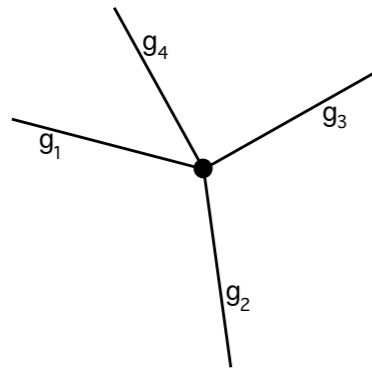
Group field theories

Fock vacuum: “no-space” (“emptiest”) state $|0\rangle$

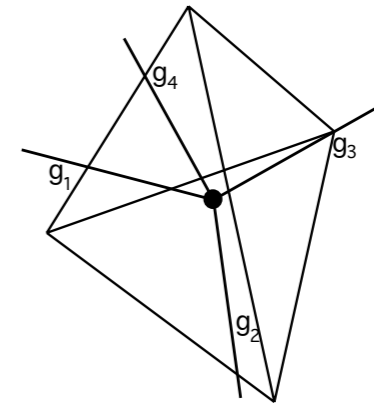
Group field theories

Fock vacuum: “no-space” (“emptiest”) state $|0\rangle$

single field “quantum”: spin network vertex or tetrahedron
 (“building block of space”)



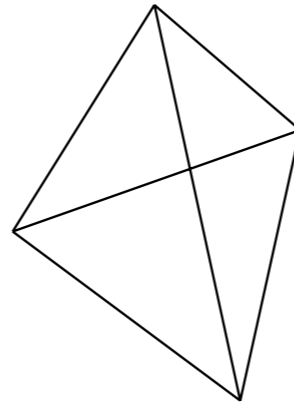
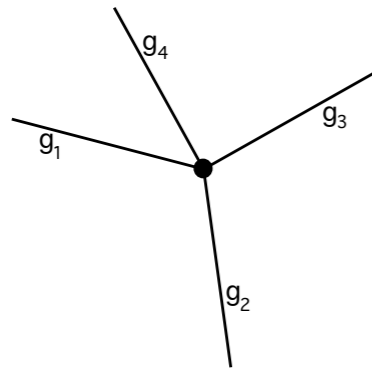
$$\varphi(g_1, g_2, g_3, g_4) \leftrightarrow \varphi(B_1, B_2, B_3, B_4) \rightarrow \mathbb{C}$$



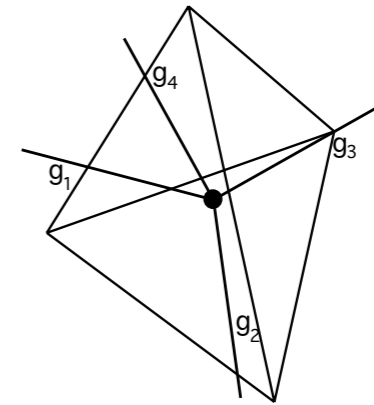
Group field theories

Fock vacuum: “no-space” (“emptiest”) state $|0\rangle$

single field “quantum”: spin network vertex or tetrahedron
 (“building block of space”)



$$\varphi(g_1, g_2, g_3, g_4) \leftrightarrow \varphi(B_1, B_2, B_3, B_4) \rightarrow \mathbb{C}$$



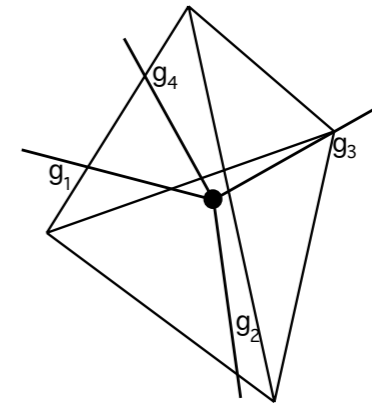
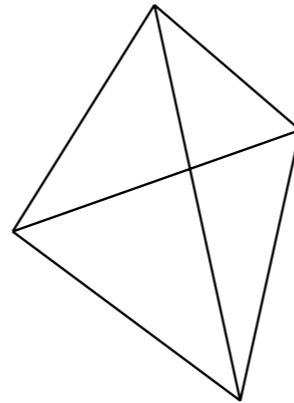
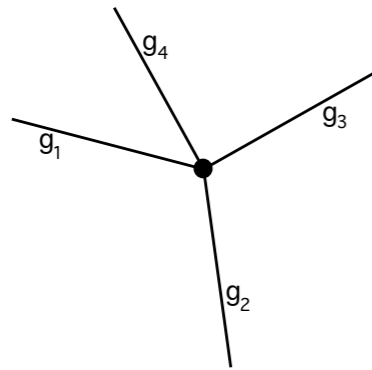
generic quantum state: arbitrary collection of spin network vertices (including glued ones) or tetrahedra (including glued ones)

Group field theories

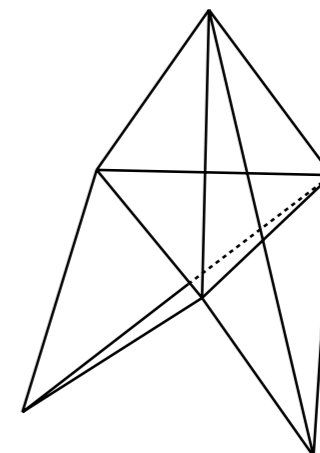
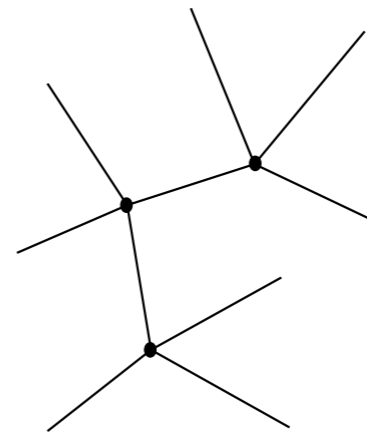
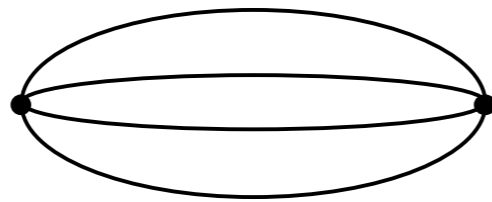
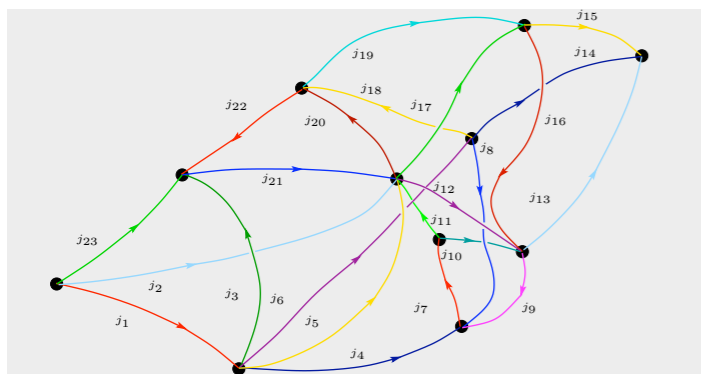
Fock vacuum: “no-space” (“emptiest”) state $|0\rangle$

single field “quantum”: spin network vertex or tetrahedron
 (“building block of space”)

$$\varphi(g_1, g_2, g_3, g_4) \leftrightarrow \varphi(B_1, B_2, B_3, B_4) \rightarrow \mathbb{C}$$



generic quantum state: arbitrary collection of spin network vertices (including glued ones) or tetrahedra (including glued ones)



Group field theories

classical action: kinetic (quadratic) term + (higher order) interaction (convolution of GFT fields)

$$S(\varphi, \bar{\varphi}) = \frac{1}{2} \int [dg_i] \overline{\varphi(g_i)} \mathcal{K}(g_i) \varphi(g_i) + \frac{\lambda}{D!} \int [dg_{ia}] \varphi(g_{i1}) \dots \varphi(\bar{g}_{iD}) \mathcal{V}(g_{ia}, \bar{g}_{iD}) + c.c.$$

Group field theories

classical action: kinetic (quadratic) term + (higher order) interaction (convolution of GFT fields)

$$S(\varphi, \bar{\varphi}) = \frac{1}{2} \int [dg_i] \overline{\varphi(g_i)} \mathcal{K}(g_i) \varphi(g_i) + \frac{\lambda}{D!} \int [dg_{ia}] \varphi(g_{i1}) \dots \varphi(\bar{g}_{iD}) \mathcal{V}(g_{ia}, \bar{g}_{iD}) + c.c.$$

“combinatorial non-locality”
in pairing of field arguments



Group field theories

classical action: kinetic (quadratic) term + (higher order) interaction (convolution of GFT fields)

$$S(\varphi, \bar{\varphi}) = \frac{1}{2} \int [dg_i] \overline{\varphi(g_i)} \mathcal{K}(g_i) \varphi(g_i) + \frac{\lambda}{D!} \int [dg_{ia}] \varphi(g_{i1}) \dots \varphi(\bar{g}_{iD}) \mathcal{V}(g_{ia}, \bar{g}_{iD}) + c.c.$$

“combinatorial non-locality”
in pairing of field arguments



simplest example (case d=4): simplicial setting

Group field theories

classical action: kinetic (quadratic) term + (higher order) interaction (convolution of GFT fields)

$$S(\varphi, \bar{\varphi}) = \frac{1}{2} \int [dg_i] \overline{\varphi(g_i)} \mathcal{K}(g_i) \varphi(g_i) + \frac{\lambda}{D!} \int [dg_{ia}] \varphi(g_{i1}) \dots \varphi(\bar{g}_{iD}) \mathcal{V}(g_{ia}, \bar{g}_{iD}) + c.c.$$

“combinatorial non-locality”
in pairing of field arguments



simplest example (case d=4): simplicial setting

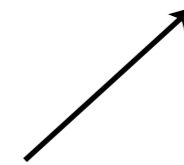
combinatorics of field arguments in interaction: gluing of 5 tetrahedra across common triangles, to form 4-simplex (“building block of spacetime”)

Group field theories

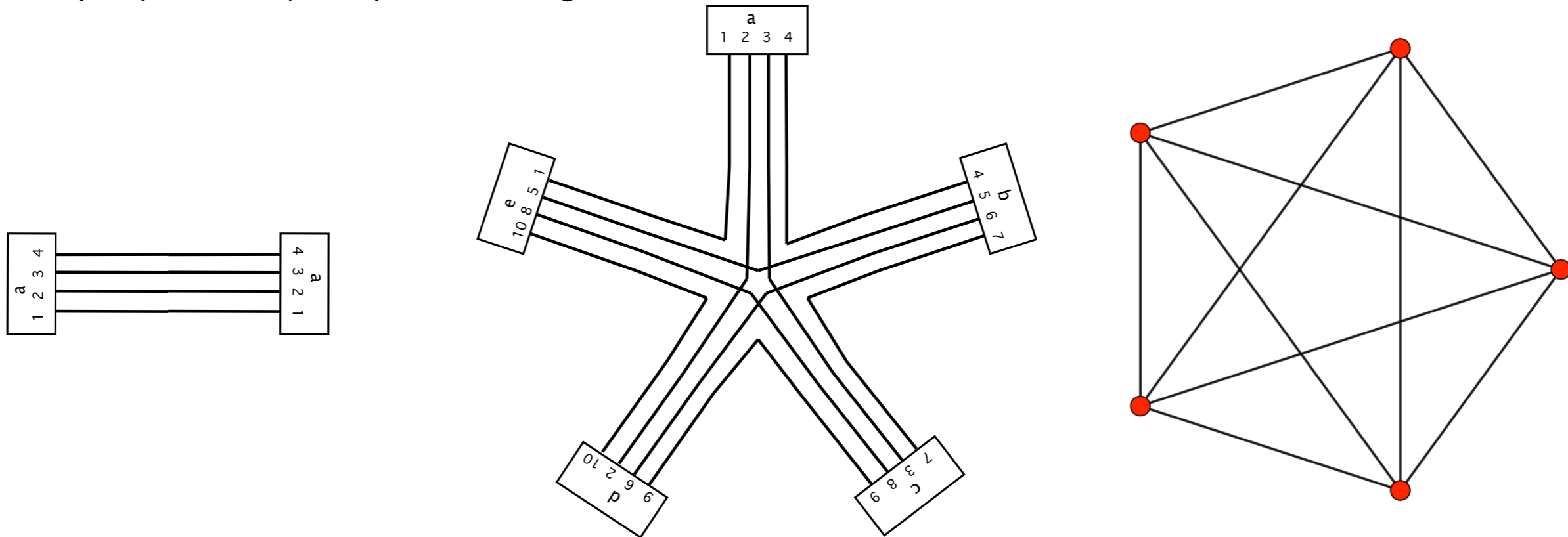
classical action: kinetic (quadratic) term + (higher order) interaction (convolution of GFT fields)

$$S(\varphi, \bar{\varphi}) = \frac{1}{2} \int [dg_i] \overline{\varphi(g_i)} \mathcal{K}(g_i) \varphi(g_i) + \frac{\lambda}{D!} \int [dg_{ia}] \varphi(g_{i1}) \dots \varphi(\bar{g}_{iD}) \mathcal{V}(g_{ia}, \bar{g}_{iD}) + c.c.$$

“combinatorial non-locality”
in pairing of field arguments



simplest example (case d=4): simplicial setting



Group field theories

Feynman perturbative expansion around trivial vacuum

$$\mathcal{Z} = \int \mathcal{D}\varphi \mathcal{D}\bar{\varphi} e^{i S_\lambda(\varphi, \bar{\varphi})} = \sum_{\Gamma} \frac{\lambda^{N_\Gamma}}{\text{sym}(\Gamma)} \mathcal{A}_\Gamma$$

Group field theories

Feynman perturbative expansion around trivial vacuum

$$\mathcal{Z} = \int \mathcal{D}\varphi \mathcal{D}\bar{\varphi} e^{i S_\lambda(\varphi, \bar{\varphi})} = \sum_{\Gamma} \frac{\lambda^{N_\Gamma}}{\text{sym}(\Gamma)} \mathcal{A}_\Gamma$$

Feynman diagrams (obtained by convoluting propagators with interaction kernels) =

= stranded diagrams dual to cellular complexes of arbitrary topology

(simplicial case: simplicial complexes obtained by gluing d-simplices in arbitrary ways)

Group field theories

Feynman perturbative expansion around trivial vacuum

$$\mathcal{Z} = \int \mathcal{D}\varphi \mathcal{D}\bar{\varphi} e^{i S_\lambda(\varphi, \bar{\varphi})} = \sum_{\Gamma} \frac{\lambda^{N_\Gamma}}{\text{sym}(\Gamma)} \mathcal{A}_\Gamma$$

Feynman diagrams (obtained by convoluting propagators with interaction kernels) =

= stranded diagrams dual to cellular complexes of arbitrary topology

(simplicial case: simplicial complexes obtained by gluing d-simplices in arbitrary ways)

Feynman amplitudes (model-dependent):

equivalently:

- spin foam models (sum-over-histories of spin networks)

Reisenberger, Rovelli, '00

- lattice path integrals
(with group+Lie algebra variables)

A. Baratin, DO, '11

Group field theories

Feynman perturbative expansion around trivial vacuum

$$\mathcal{Z} = \int \mathcal{D}\varphi \mathcal{D}\bar{\varphi} e^{i S_\lambda(\varphi, \bar{\varphi})} = \sum_{\Gamma} \frac{\lambda^{N_\Gamma}}{\text{sym}(\Gamma)} \mathcal{A}_\Gamma$$

Feynman diagrams (obtained by convoluting propagators with interaction kernels) =

= stranded diagrams dual to cellular complexes of arbitrary topology

(simplicial case: simplicial complexes obtained by gluing d-simplices in arbitrary ways)

Feynman amplitudes (model-dependent):

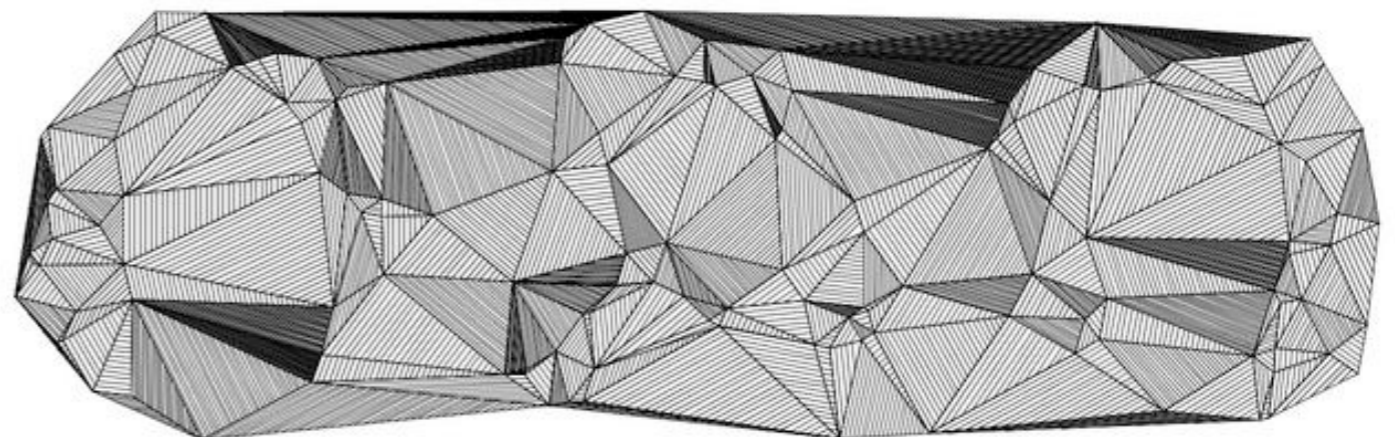
equivalently:

- spin foam models (sum-over-histories of spin networks)

Reisenberger, Rovelli, '00

- lattice path integrals (with group+Lie algebra variables)

A. Baratin, DO, '11



Group field theories

Feynman perturbative expansion around trivial vacuum

$$\mathcal{Z} = \int \mathcal{D}\varphi \mathcal{D}\bar{\varphi} e^{i S_\lambda(\varphi, \bar{\varphi})} = \sum_{\Gamma} \frac{\lambda^{N_\Gamma}}{\text{sym}(\Gamma)} \mathcal{A}_\Gamma$$

Feynman diagrams (obtained by convoluting propagators with interaction kernels) =

= stranded diagrams dual to cellular complexes of arbitrary topology

(simplicial case: simplicial complexes obtained by gluing d-simplices in arbitrary ways)

Feynman amplitudes (model-dependent):

equivalently:

- spin foam models (sum-over-histories of spin networks)

Reisenberger, Rovelli, '00

- lattice path integrals
(with group+Lie algebra variables)

A. Baratin, DO, '11

Group field theories

Feynman perturbative expansion around trivial vacuum

$$\mathcal{Z} = \int \mathcal{D}\varphi \mathcal{D}\bar{\varphi} e^{i S_\lambda(\varphi, \bar{\varphi})} = \sum_{\Gamma} \frac{\lambda^{N_\Gamma}}{\text{sym}(\Gamma)} \mathcal{A}_\Gamma$$

Feynman diagrams (obtained by convoluting propagators with interaction kernels) =

= stranded diagrams dual to cellular complexes of arbitrary topology

(simplicial case: simplicial complexes obtained by gluing d-simplices in arbitrary ways)

Feynman amplitudes (model-dependent):

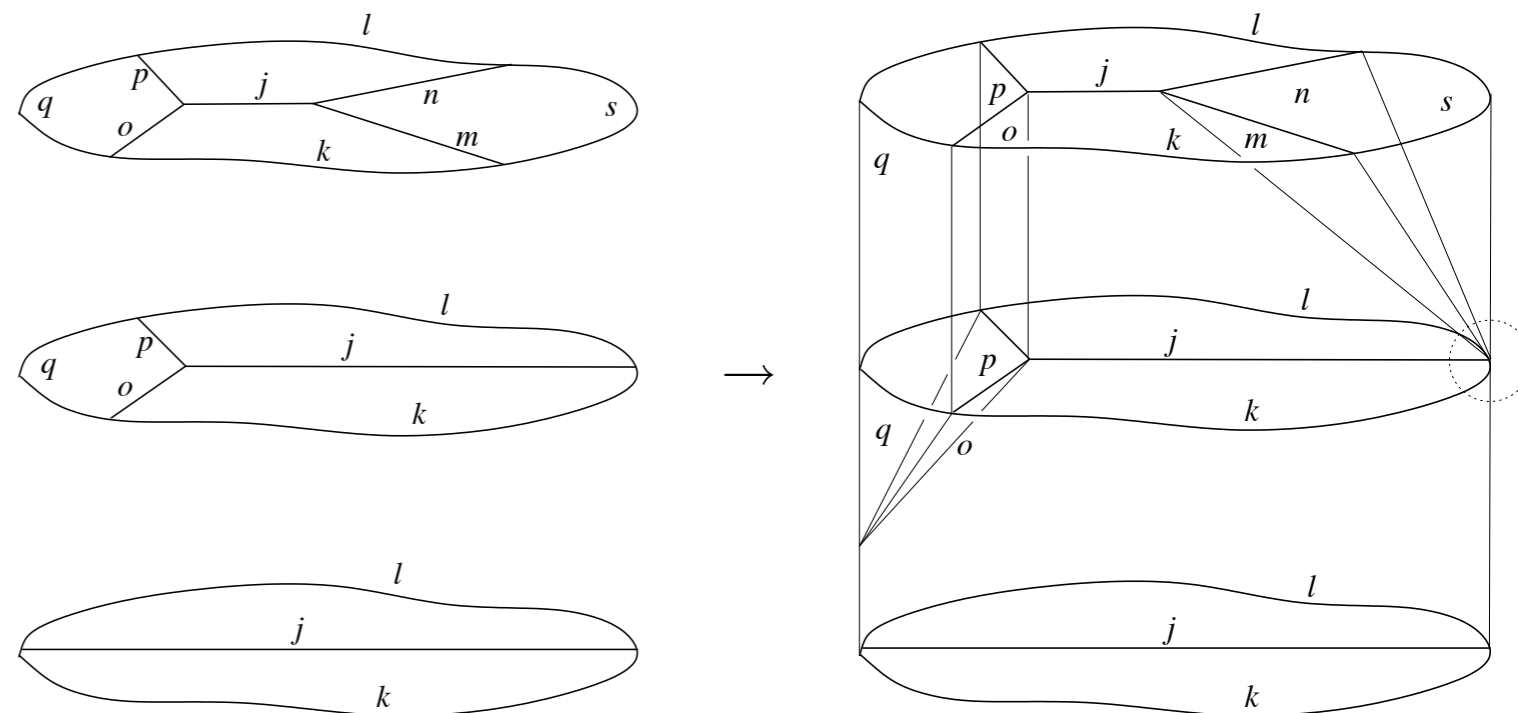
equivalently:

- spin foam models (sum-over-histories of spin networks)

Reisenberger, Rovelli, '00

- lattice path integrals (with group+Lie algebra variables)

A. Baratin, DO, '11



Group field theories

Feynman perturbative expansion around trivial vacuum

$$\mathcal{Z} = \int \mathcal{D}\varphi \mathcal{D}\bar{\varphi} e^{i S_\lambda(\varphi, \bar{\varphi})} = \sum_{\Gamma} \frac{\lambda^{N_\Gamma}}{\text{sym}(\Gamma)} \mathcal{A}_\Gamma$$

Feynman diagrams (obtained by convoluting propagators with interaction kernels) =

= stranded diagrams dual to cellular complexes of arbitrary topology

(simplicial case: simplicial complexes obtained by gluing d-simplices in arbitrary ways)

Feynman amplitudes (model-dependent):

equivalently:

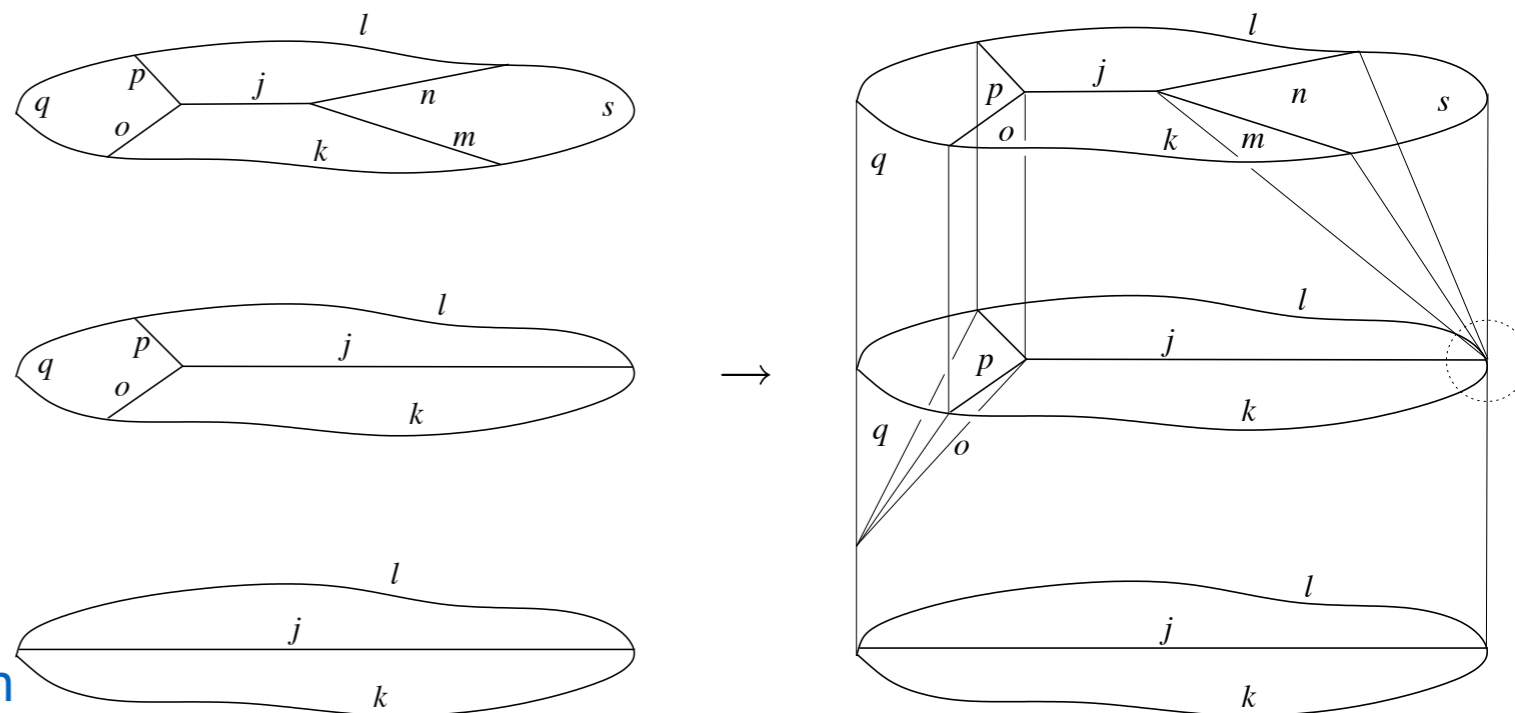
- spin foam models (sum-over-histories of spin networks)

Reisenberger, Rovelli, '00

- lattice path integrals (with group+Lie algebra variables)

A. Baratin, DO, '11

all current spin foam models have GFT formulation



GFTs, loop quantum gravity, spin foam models

appropriate conditions on GFT fields or GFT dynamics (and choice of data) turn GFT Feynman amplitudes into lattice gauge theories/discrete gravity path integrals/spin foam models

e.g. gauge invariance of GFT fields under diagonal action of group G

GFTs, loop quantum gravity, spin foam models

appropriate conditions on GFT fields or GFT dynamics (and choice of data) turn GFT Feynman amplitudes into lattice gauge theories/discrete gravity path integrals/spin foam models

e.g. gauge invariance of GFT fields under diagonal action of group G

example: $d=3$ $\varphi_\ell : SO(3)^3 / SO(3) \rightarrow \mathbb{R}$ + simplicial interaction

$\forall h \in SO(3), \quad \varphi_\ell(hg_1, hg_2, hg_3) = \varphi_\ell(g_1, g_2, g_3)$ with only delta functions

can be computed in different (equivalent) representations (group, spin, Lie algebra)

GFTs, loop quantum gravity, spin foam models

appropriate conditions on GFT fields or GFT dynamics (and choice of data) turn GFT Feynman amplitudes into lattice gauge theories/discrete gravity path integrals/spin foam models

e.g. gauge invariance of GFT fields under diagonal action of group G

example: $d=3$ $\varphi_\ell : SO(3)^3 / SO(3) \rightarrow \mathbb{R}$ + simplicial interaction

$\forall h \in SO(3), \quad \varphi_\ell(hg_1, hg_2, hg_3) = \varphi_\ell(g_1, g_2, g_3)$ with only delta functions

can be computed in different (equivalent) representations (group, spin, Lie algebra)

$$S_{kin}[\varphi_\ell] = \int [dg_i]^3 \sum_{\ell=1}^4 \varphi_\ell(g_1, g_2, g_3) \overline{\varphi}_\ell(g_1, g_2 \cdot g_3),$$

$$S_{int}[\varphi_\ell] = \lambda \int [dg_i]^6 \varphi_1(g_1, g_2, g_3) \varphi_2(g_3, g_4, g_5) \varphi_3(g_5, g_2, g_6) \varphi_4(g_6, g_4, g_1) \\ + \lambda \int [dg_i]^6 \overline{\varphi}_4(g_1, g_4, g_6) \overline{\varphi}_3(g_6, g_2, g_5) \overline{\varphi}_2(g_5, g_4, g_3) \overline{\varphi}_1(g_3, g_2, g_1)$$

GFTs, loop quantum gravity, spin foam models

appropriate conditions on GFT fields or GFT dynamics (and choice of data) turn GFT Feynman amplitudes into lattice gauge theories/discrete gravity path integrals/spin foam models

e.g. gauge invariance of GFT fields under diagonal action of group G

example: $d=3$ $\varphi_\ell : SO(3)^3 / SO(3) \rightarrow \mathbb{R}$ + simplicial interaction

$\forall h \in SO(3), \quad \varphi_\ell(hg_1, hg_2, hg_3) = \varphi_\ell(g_1, g_2, g_3)$ with only delta functions

can be computed in different (equivalent) representations (group, spin, Lie algebra)

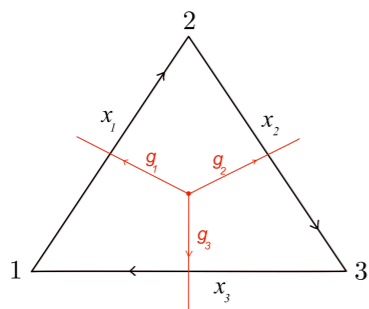
GFTs, loop quantum gravity, spin foam models

appropriate conditions on GFT fields or GFT dynamics (and choice of data) turn GFT Feynman amplitudes into lattice gauge theories/discrete gravity path integrals/spin foam models

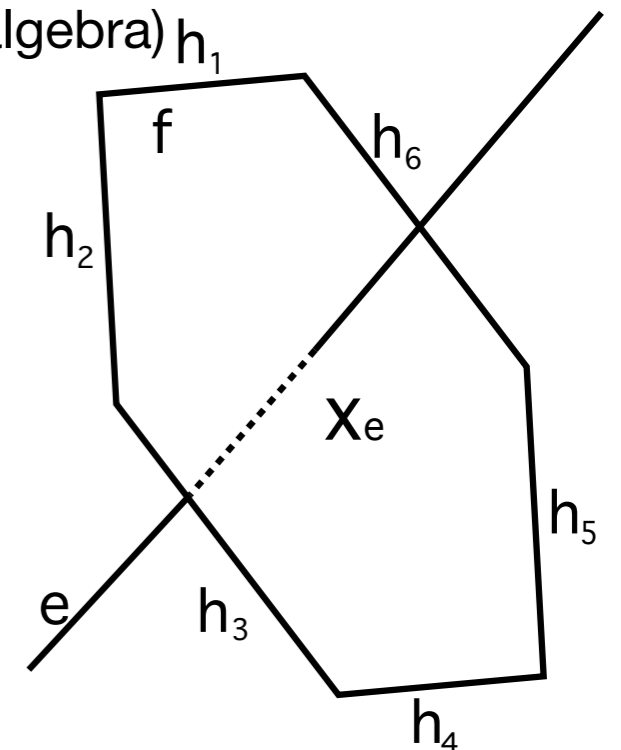
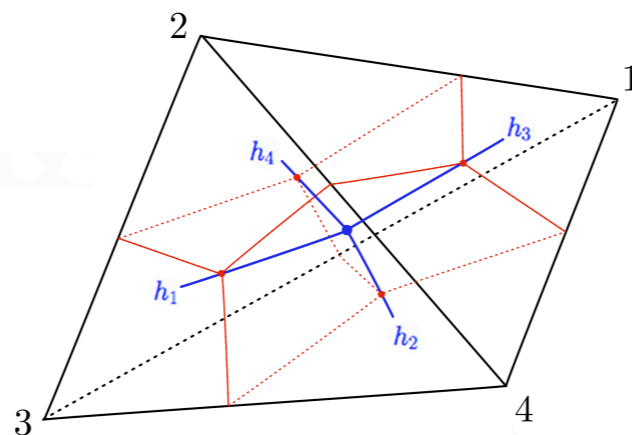
e.g. gauge invariance of GFT fields under diagonal action of group G

example: $d=3$ $\varphi_\ell : SO(3)^3 / SO(3) \rightarrow \mathbb{R}$ + simplicial interaction
 $\forall h \in SO(3), \quad \varphi_\ell(hg_1, hg_2, hg_3) = \varphi_\ell(g_1, g_2, g_3)$ with only delta functions

can be computed in different (equivalent) representations (group, spin, Lie algebra)



$$\varphi(g_1, g_2, g_3) \leftrightarrow \varphi(x_1, x_2, x_3)$$



discretization of: $S(e, \omega) = \int Tr(e \wedge F(\omega))$

GFTs, loop quantum gravity, spin foam models

appropriate conditions on GFT fields or GFT dynamics (and choice of data) turn GFT Feynman amplitudes into lattice gauge theories/discrete gravity path integrals/spin foam models

e.g. gauge invariance of GFT fields under diagonal action of group G

example: $d=3$ $\varphi_\ell : SO(3)^3 / SO(3) \rightarrow \mathbb{R}$ + simplicial interaction

$\forall h \in SO(3), \quad \varphi_\ell(hg_1, hg_2, hg_3) = \varphi_\ell(g_1, g_2, g_3)$ with only delta functions

can be computed in different (equivalent) representations (group, spin, Lie algebra)

GFTs, loop quantum gravity, spin foam models

appropriate conditions on GFT fields or GFT dynamics (and choice of data) turn GFT Feynman amplitudes into lattice gauge theories/discrete gravity path integrals/spin foam models

e.g. gauge invariance of GFT fields under diagonal action of group G

example: d=3 $\varphi_\ell : SO(3)^3 / SO(3) \rightarrow \mathbb{R}$ + simplicial interaction

$\forall h \in SO(3), \quad \varphi_\ell(hg_1, hg_2, hg_3) = \varphi_\ell(g_1, g_2, g_3)$ with only delta functions

can be computed in different (equivalent) representations (group, spin, Lie algebra)

$$\begin{aligned} \mathcal{A}_\Gamma &= \int \prod_l dh_l \prod_f \delta(H_f(h_l)) = \int \prod_l dh_l \prod_f \delta\left(\overrightarrow{\prod}_{l \in \partial f} h_l\right) = \\ &= \sum_{\{j_e\}} \prod_e d_{j_e} \prod_\tau \left\{ \begin{array}{ccc} j_1^\tau & j_2^\tau & j_3^\tau \\ j_4^\tau & j_5^\tau & j_6^\tau \end{array} \right\} = \int \prod_l [dh_l] \prod_e [d^3 x_e] e^{i \sum_e \text{Tr } x_e H_e} \end{aligned}$$

GFTs, loop quantum gravity, spin foam models

appropriate conditions on GFT fields or GFT dynamics (and choice of data) turn GFT Feynman amplitudes into lattice gauge theories/discrete gravity path integrals/spin foam models

e.g. gauge invariance of GFT fields under diagonal action of group G

example: d=3 $\varphi_\ell : SO(3)^3 / SO(3) \rightarrow \mathbb{R}$ + simplicial interaction
 $\forall h \in SO(3), \quad \varphi_\ell(hg_1, hg_2, hg_3) = \varphi_\ell(g_1, g_2, g_3)$ with only delta functions

can be computed in different (equivalent) representations (group, spin, Lie algebra)

$$\begin{aligned} \mathcal{A}_\Gamma &= \int \prod_l dh_l \prod_f \delta(H_f(h_l)) = \int \prod_l dh_l \prod_f \delta\left(\overrightarrow{\prod}_{l \in \partial f} h_l\right) = \text{← lattice gauge theory formulation of 3d gravity/BF theory} \\ &= \sum_{\{j_e\}} \prod_e d_{j_e} \prod_\tau \left\{ \begin{matrix} j_1^\tau & j_2^\tau & j_3^\tau \\ j_4^\tau & j_5^\tau & j_6^\tau \end{matrix} \right\} = \int \prod_l [dh_l] \prod_e [d^3 x_e] e^{i \sum_e \text{Tr } x_e H_e} \end{aligned}$$

GFTs, loop quantum gravity, spin foam models

appropriate conditions on GFT fields or GFT dynamics (and choice of data) turn GFT Feynman amplitudes into lattice gauge theories/discrete gravity path integrals/spin foam models

e.g. gauge invariance of GFT fields under diagonal action of group G

example: $d=3$ $\varphi_\ell : SO(3)^3 / SO(3) \rightarrow \mathbb{R}$ + simplicial interaction

$\forall h \in SO(3), \quad \varphi_\ell(hg_1, hg_2, hg_3) = \varphi_\ell(g_1, g_2, g_3)$ with only delta functions

can be computed in different (equivalent) representations (group, spin, Lie algebra)

$$\begin{aligned} \mathcal{A}_\Gamma &= \int \prod_l dh_l \prod_f \delta(H_f(h_l)) = \int \prod_l dh_l \prod_f \delta\left(\overrightarrow{\prod_{l \in \partial f}} h_l\right) = \text{--- lattice gauge theory formulation of 3d gravity/BF theory} \\ &= \sum_{\{j_e\}} \prod_e d_{j_e} \prod_\tau \left\{ \begin{array}{ccc} j_1^\tau & j_2^\tau & j_3^\tau \\ j_4^\tau & j_5^\tau & j_6^\tau \end{array} \right\} = \int \prod_l [dh_l] \prod_e [d^3 x_e] e^{i \sum_e \text{Tr } x_e H_e} \end{aligned}$$

spin foam formulation of 3d gravity/BF theory

GFTs, loop quantum gravity, spin foam models

appropriate conditions on GFT fields or GFT dynamics (and choice of data) turn GFT Feynman amplitudes into lattice gauge theories/discrete gravity path integrals/spin foam models

e.g. gauge invariance of GFT fields under diagonal action of group G

example: $d=3$ $\varphi_\ell : SO(3)^3 / SO(3) \rightarrow \mathbb{R}$ + simplicial interaction

$\forall h \in SO(3), \quad \varphi_\ell(hg_1, hg_2, hg_3) = \varphi_\ell(g_1, g_2, g_3)$ with only delta functions

can be computed in different (equivalent) representations (group, spin, Lie algebra)

$$\begin{aligned} \mathcal{A}_\Gamma &= \int \prod_l dh_l \prod_f \delta(H_f(h_l)) = \int \prod_l dh_l \prod_f \delta\left(\overrightarrow{\prod}_{l \in \partial f} h_l\right) = \\ &= \sum_{\{j_e\}} \prod_e d_{j_e} \prod_\tau \left\{ \begin{matrix} j_1^\tau & j_2^\tau & j_3^\tau \\ j_4^\tau & j_5^\tau & j_6^\tau \end{matrix} \right\} = \int \prod_l [dh_l] \prod_e [d^3 x_e] e^{i \sum_e \text{Tr } x_e H_e} \end{aligned}$$

lattice gauge theory formulation of 3d gravity/BF theory

discrete 1st order path integral for 3d gravity/BF theory on simplicial complex dual to GFT Feynman diagram

spin foam formulation of 3d gravity/BF theory

GFTs, loop quantum gravity, spin foam models

GFT models of 4d gravity:

based on classical (Plebanski) formulation of GR as BF theory + (simplicity) constraints

start from GFT formulation of 4d BF theory
+ impose simplicity constraints (geometricity of simplicial structures)

(Barbieri, Baez, Barrett, Crane, Reisenberger, Perez, De Pietri, Engle, Pereira, Freidel, Krasnov, Rovelli, Livine, Speziale, Baratin, DO,)

GFTs, loop quantum gravity, spin foam models

GFT models of 4d gravity:

based on classical (Plebanski) formulation of GR as BF theory + (simplicity) constraints

start from GFT formulation of 4d BF theory
+ impose simplicity constraints (geometricity of simplicial structures)

(Barbieri, Baez, Barrett, Crane, Reisenberger, Perez, De Pietri, Engle, Pereira, Freidel, Krasnov, Rovelli, Livine, Speziale, Baratin, DO,)

“geometricity operator” = simplicity constraints + gauge invariance:

$$G \triangleright \phi \equiv C \triangleright S^\beta \triangleright \phi = S^\beta \triangleright C \triangleright \phi \equiv \Psi$$

GFTs, loop quantum gravity, spin foam models

GFT models of 4d gravity:

based on classical (Plebanski) formulation of GR as BF theory + (simplicity) constraints

start from GFT formulation of 4d BF theory
+ impose simplicity constraints (geometricity of simplicial structures)

(Barbieri, Baez, Barrett, Crane, Reisenberger, Perez, De Pietri, Engle, Pereira, Freidel, Krasnov, Rovelli, Livine, Speziale, Baratin, DO,

“geometricity operator” = simplicity constraints + gauge invariance:

$$G \triangleright \phi \equiv C \triangleright S^\beta \triangleright \phi = S^\beta \triangleright C \triangleright \phi \equiv \Psi$$

concrete, well-defined GFT (spin foam) model(s) for 4d QG dynamics - nice discrete geometry, lots of results

GFTs, loop quantum gravity, spin foam models

GFT models of 4d gravity:

based on classical (Plebanski) formulation of GR as BF theory + (simplicity) constraints

start from GFT formulation of 4d BF theory
+ impose simplicity constraints (geometricity of simplicial structures)

(Barbieri, Baez, Barrett, Crane, Reisenberger, Perez, De Pietri, Engle, Pereira, Freidel, Krasnov, Rovelli, Livine, Speziale, Baratin, DO,

“geometricity operator” = simplicity constraints + gauge invariance:

$$G \triangleright \phi \equiv C \triangleright S^\beta \triangleright \phi = S^\beta \triangleright C \triangleright \phi \equiv \Psi$$

concrete, well-defined GFT (spin foam) model(s) for 4d QG dynamics - nice discrete geometry, lots of results

all current spin foam models have a GFT formulation

GFTs, loop quantum gravity, spin foam models

GFT models of 4d gravity:

based on classical (Plebanski) formulation of GR as BF theory + (simplicity) constraints

start from GFT formulation of 4d BF theory
+ impose simplicity constraints (geometricity of simplicial structures)

(Barbieri, Baez, Barrett, Crane, Reisenberger, Perez, De Pietri, Engle, Pereira, Freidel, Krasnov, Rovelli, Livine, Speziale, Baratin, DO,)

“geometricity operator” = simplicity constraints + gauge invariance:

$$G \triangleright \phi \equiv C \triangleright S^\beta \triangleright \phi = S^\beta \triangleright C \triangleright \phi \equiv \Psi$$

concrete, well-defined GFT (spin foam) model(s) for 4d QG dynamics - nice discrete geometry, lots of results

all current spin foam models have a GFT formulation

decompose GFT field in SU(2) data +
geometricity conditions



GFT dynamics to LQG quantum states

GFTs, loop quantum gravity, spin foam models

second quantized version of Loop Quantum Gravity
but dynamics not derived from canonical quantization of GR

DO, 1310.7786 [gr-qc]
DO, J. Ryan, J. Thuerigen, '14

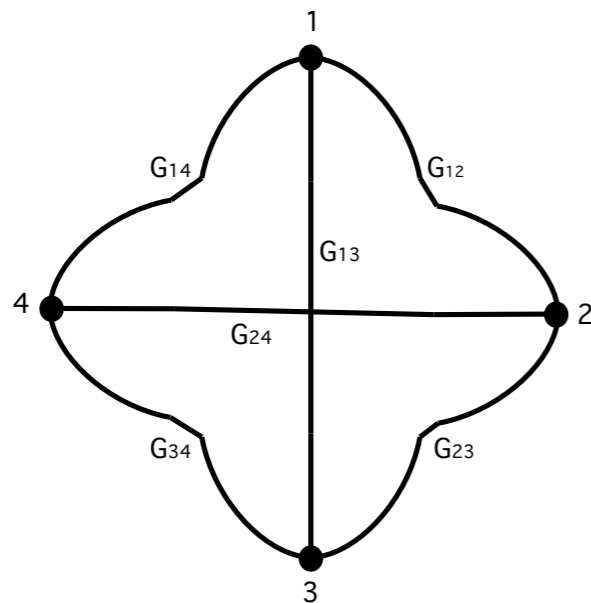
(LQG spin network states \sim many-particles states, “particle” \sim spin network vertex)

GFTs, loop quantum gravity, spin foam models

second quantized version of Loop Quantum Gravity
but dynamics not derived from canonical quantization of GR

DO, 1310.7786 [gr-qc]
DO, J. Ryan, J. Thuerigen, '14

(LQG spin network states \sim many-particles states, "particle" \sim spin network vertex)

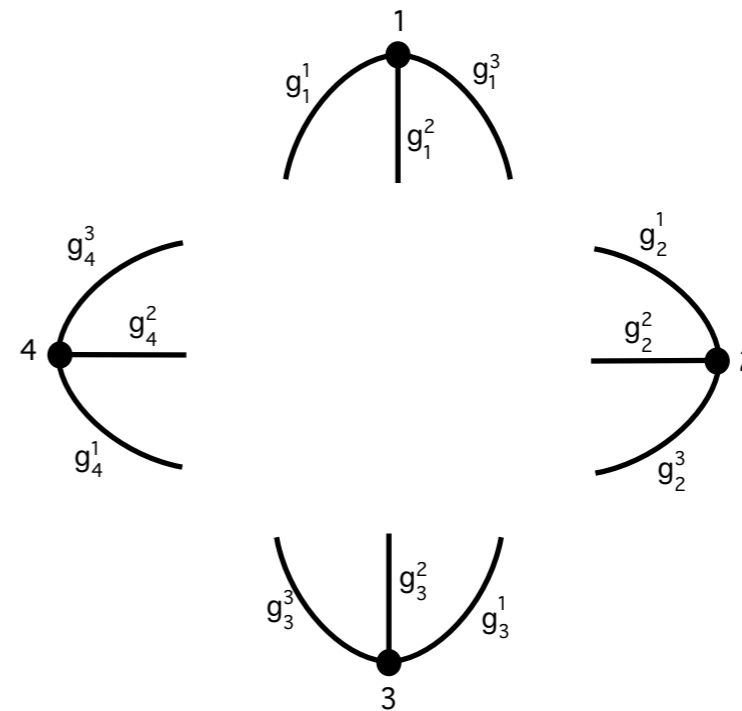
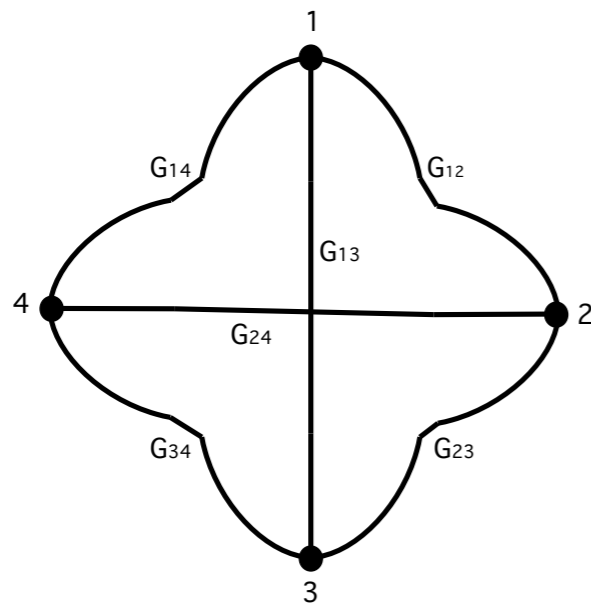


GFTs, loop quantum gravity, spin foam models

second quantized version of Loop Quantum Gravity
but dynamics not derived from canonical quantization of GR

DO, 1310.7786 [gr-qc]
DO, J. Ryan, J. Thuerigen, '14

(LQG spin network states \sim many-particles states, "particle" \sim spin network vertex)

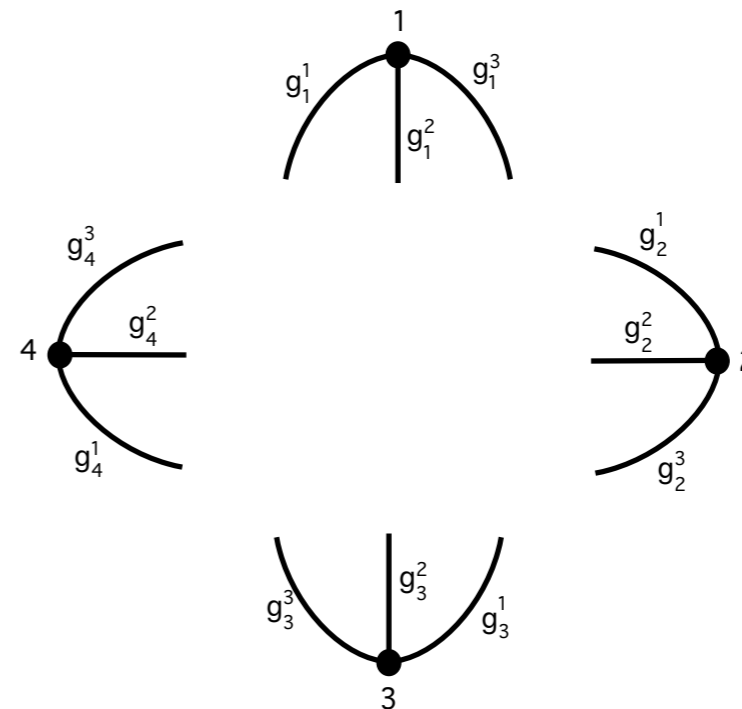
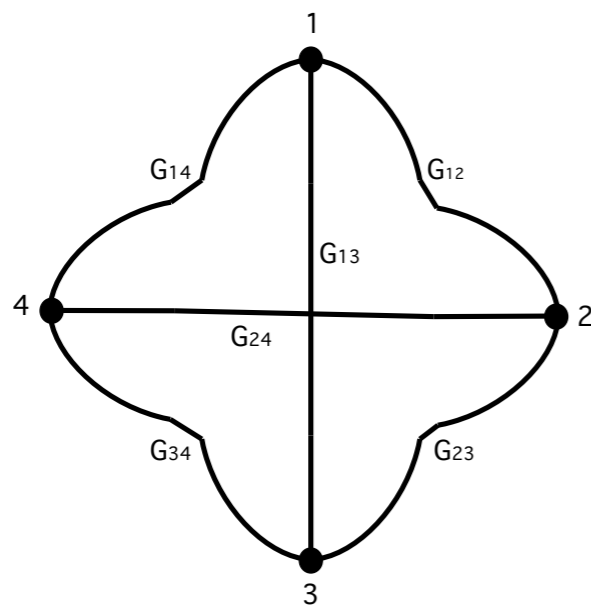


GFTs, loop quantum gravity, spin foam models

second quantized version of Loop Quantum Gravity
but dynamics not derived from canonical quantization of GR

DO, 1310.7786 [gr-qc]
DO, J. Ryan, J. Thuerigen, '14

(LQG spin network states \sim many-particles states, "particle" \sim spin network vertex)



GFT Hilbert space = Fock space of open spin network vertices - contains any LQG state (all spin networks)

any LQG observable has a 2nd quantised, GFT counterpart

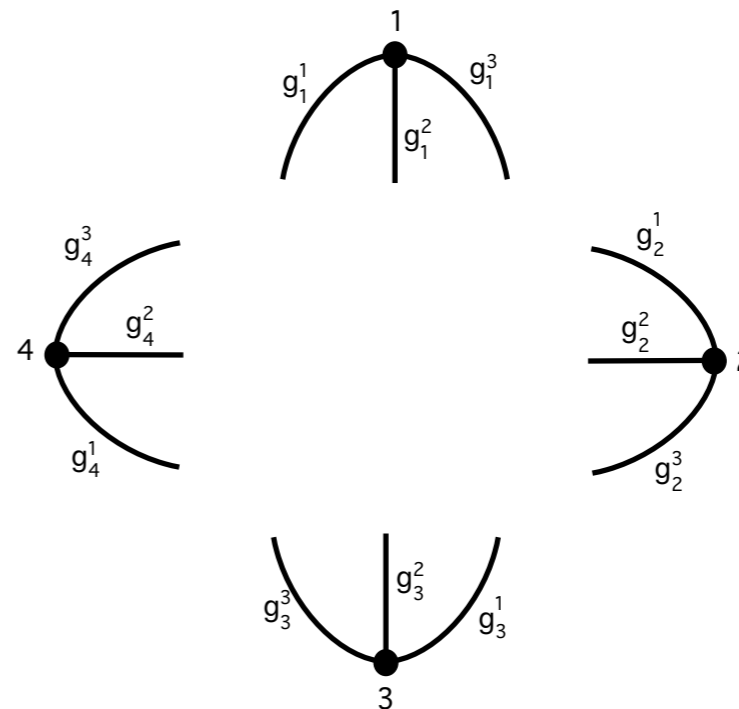
choice of LQG dynamics (Hamiltonian constraint operator) translates into choice of GFT action

GFTs, loop quantum gravity, spin foam models

second quantized version of Loop Quantum Gravity
but dynamics not derived from canonical quantization of GR

DO, 1310.7786 [gr-qc]
DO, J. Ryan, J. Thuerigen, '14

(LQG spin network states \sim many-particles states, "particle" \sim spin network vertex)



GFT Hilbert space = Fock space of open spin network vertices - contains any LQG state (all spin networks)

any LQG observable has a 2nd quantised, GFT counterpart

choice of LQG dynamics (Hamiltonian constraint operator) translates into choice of GFT action

GFTs, loop quantum gravity, spin foam models

second quantized version of Loop Quantum Gravity
but dynamics not derived from canonical quantization of GR

DO, 1310.7786 [gr-qc]
DO, J. Ryan, J. Thuerigen, '14

(LQG spin network states \sim many-particles states, "particle" \sim spin network vertex)

GFT Hilbert space = Fock space of open spin network vertices - contains any LQG state (all spin networks)

any LQG observable has a 2nd quantised, GFT counterpart

choice of LQG dynamics (Hamiltonian constraint operator) translates into choice of GFT action

GFTs, loop quantum gravity, spin foam models

second quantized version of Loop Quantum Gravity
but dynamics not derived from canonical quantization of GR

DO, 1310.7786 [gr-qc]
DO, J. Ryan, J. Thuerigen, '14

(LQG spin network states \sim many-particles states, “particle” \sim spin network vertex)

GFTs, loop quantum gravity, spin foam models

second quantized version of Loop Quantum Gravity
but dynamics not derived from canonical quantization of GR

DO, 1310.7786 [gr-qc]
DO, J. Ryan, J. Thuerigen, '14

(LQG spin network states \sim many-particles states, “particle” \sim spin network vertex)

QFT methods (i.e. GFT reformulation of LQG and spin foam models) useful to address physics of large numbers of LQG d.o.f.s, i.e. many and refined graphs (continuum limit)

(superpositions of “many-vertices” states, refinement as creation of new vertices, etc)

1. making sense of quantum dynamics and LQG partition function (correlations)
2. understanding LQG phase structure
3. extracting effective continuum dynamics

1st message

we have a solid candidate formalism for a theory of quantum gravity (a QFT for the “atoms of quantum space”)

grounded in LQG (and discrete gravity, tensor models)

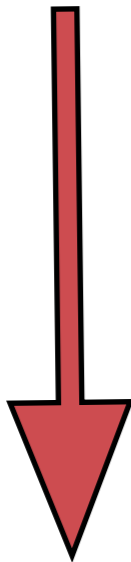
rigorous mathematics, clear pre-geometric meaning

promising fundamental dynamical models

lots of results

Quantum spacetime: the difficult path from microstructure to cosmology

Quantum Gravity problem:
identify microscopic d.o.f. of quantum spacetime and their fundamental dynamics



derive effective (QG-inspired) models for fundamental (quantum) cosmology:
explain features of early Universe, obtain testable QG predictions

various models: loop quantum cosmology,

task is daunting (imagine analogue problem in condensed matter theory)

The problem of the continuum limit in QG

new (non-geometric, non-spatio-temporal) physical degrees of freedom (“building blocks”) for space-time

The problem of the continuum limit in QG

new (non-geometric, non-spatio-temporal) physical degrees of freedom (“building blocks”) for space-time



new direction to explore: number of fundamental degrees of freedom

The problem of the continuum limit in QG

new (non-geometric, non-spatio-temporal) physical degrees of freedom (“building blocks”) for space-time



new direction to explore: number of fundamental degrees of freedom

main point:

physics of few d.o.f.s is different from physics of (very) many d.o.f.s

The problem of the continuum limit in QG

new (non-geometric, non-spatio-temporal) physical degrees of freedom (“building blocks”) for space-time



new direction to explore: number of fundamental degrees of freedom

main point:

physics of few d.o.f.s is different from physics of (very) many d.o.f.s

(quantum) continuum, geometric space-time should be recovered in the regime of large number N of non-spatio-temporal d.o.f.s

The problem of the continuum limit in QG

new (non-geometric, non-spatio-temporal) physical degrees of freedom (“building blocks”) for space-time



new direction to explore: number of fundamental degrees of freedom

main point:

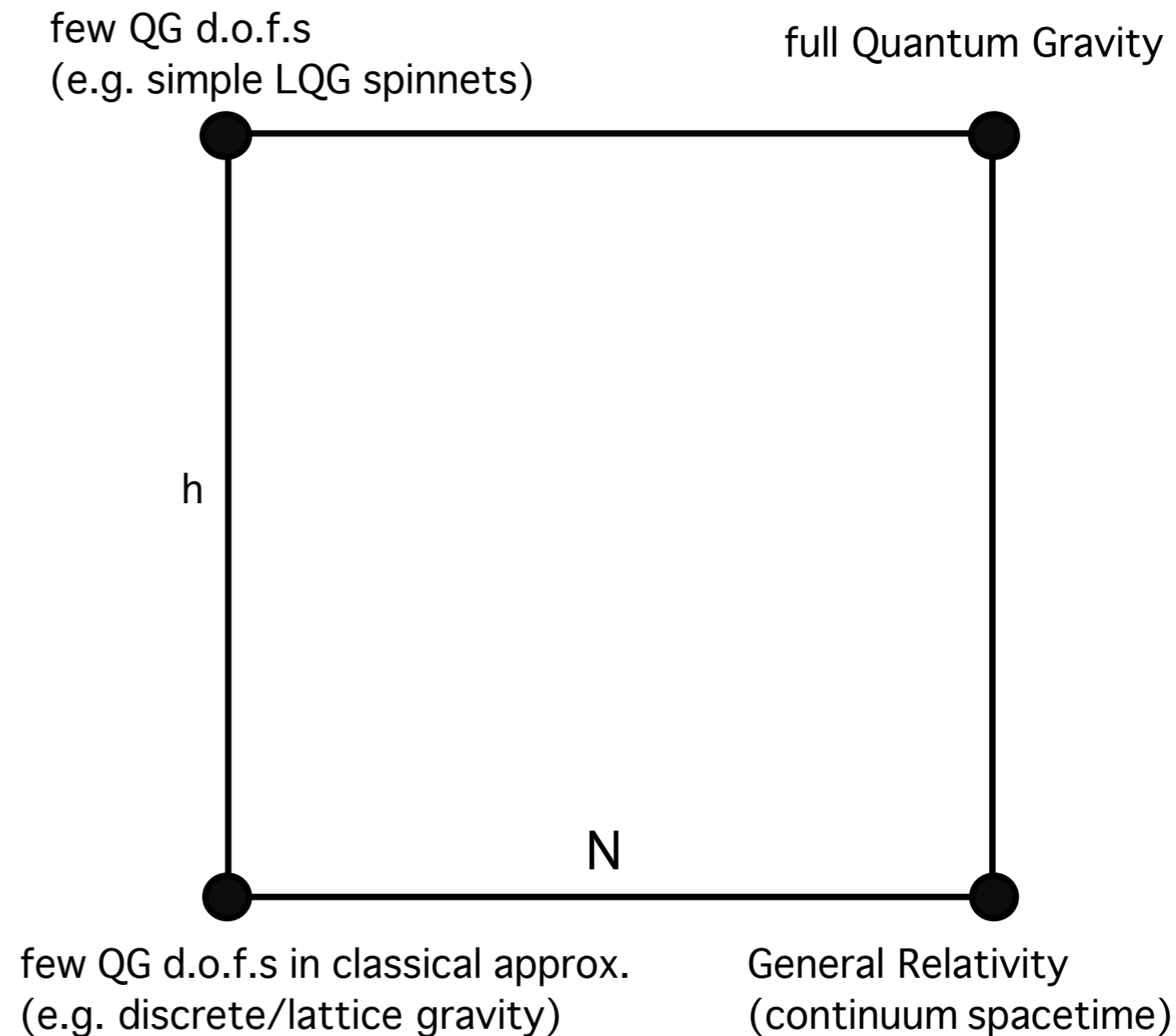
physics of few d.o.f.s is different from physics of (very) many d.o.f.s

(quantum) continuum, geometric space-time should be recovered in the regime of large number N of non-spatio-temporal d.o.f.s

GFTs are a formulation of LQG/spin foams that is most suited to tackle this problem, thanks to QFT tools

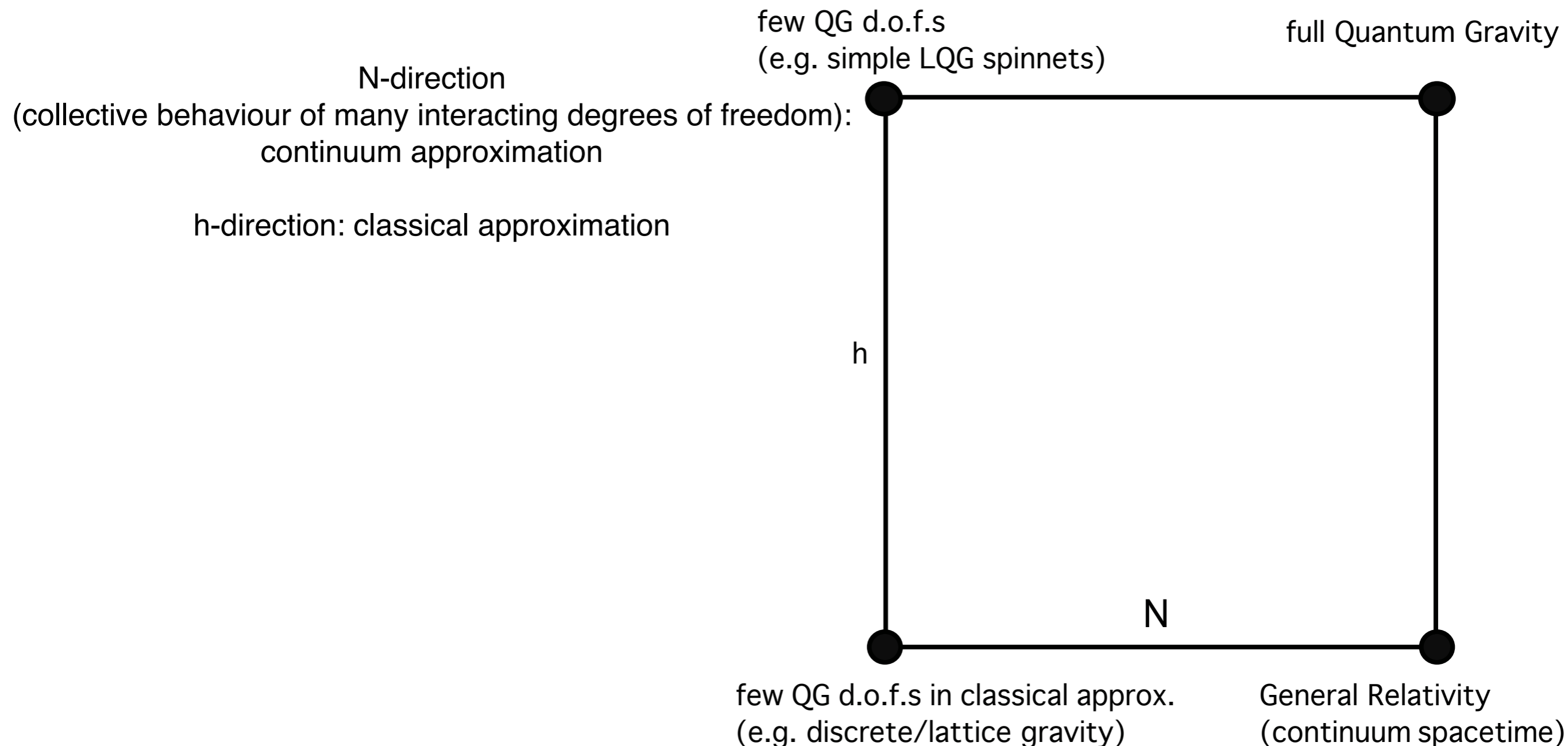
Moving along the N-direction: the case of QG

continuum approximation very different (conceptually and technically) from classical approximation



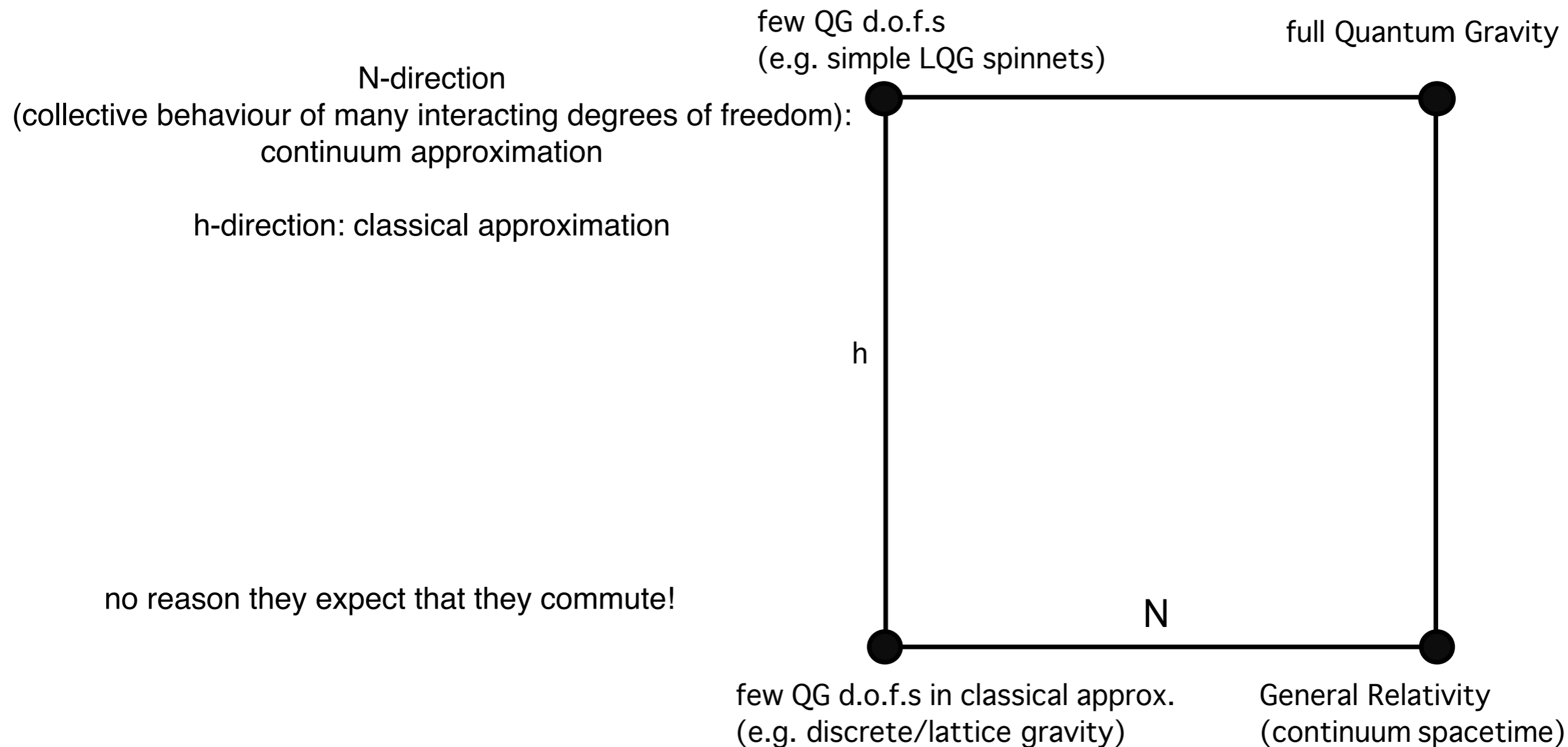
Moving along the N-direction: the case of QG

continuum approximation very different (conceptually and technically) from classical approximation



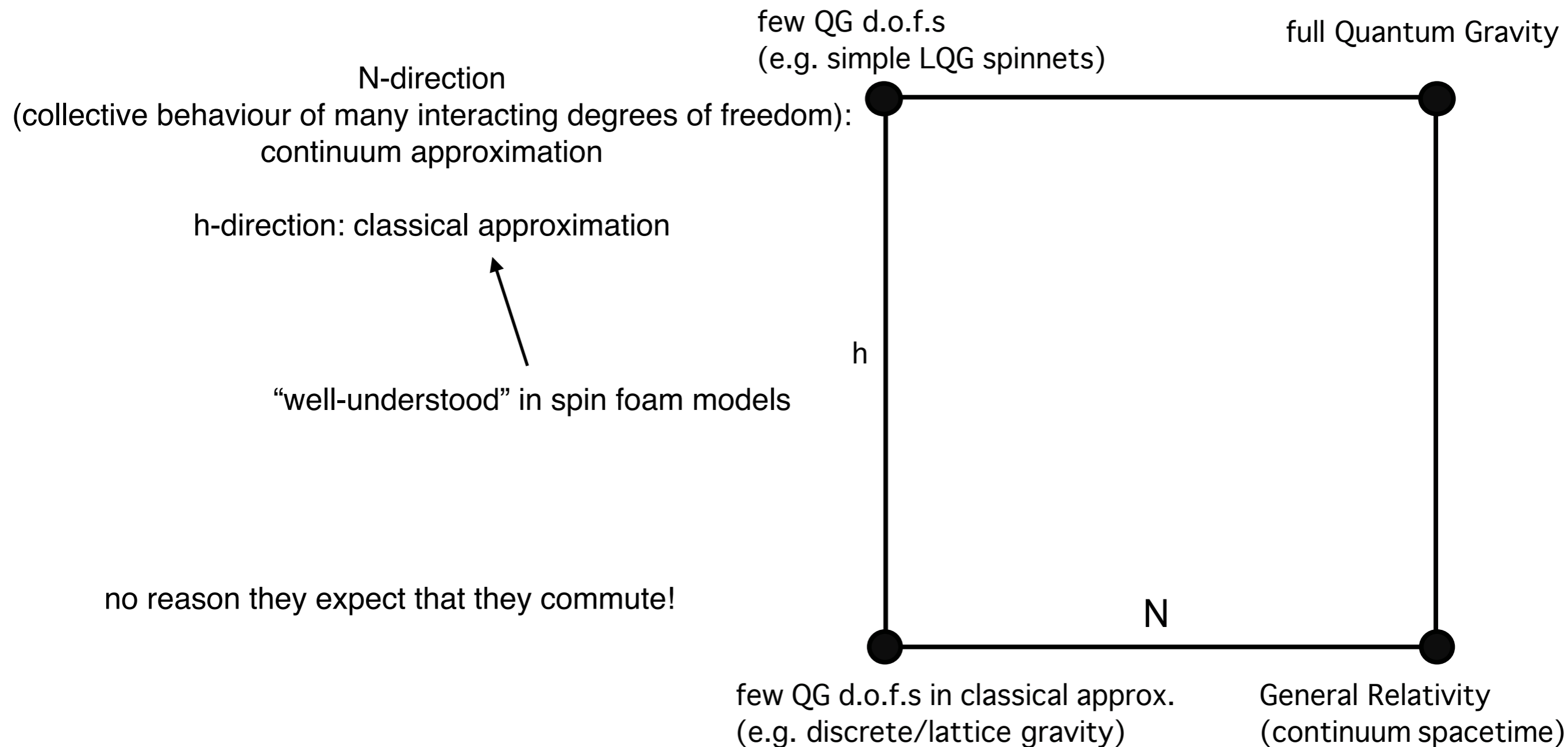
Moving along the N-direction: the case of QG

continuum approximation very different (conceptually and technically) from classical approximation



Moving along the N-direction: the case of QG

continuum approximation very different (conceptually and technically) from classical approximation



Problem of the continuum in QG: role of RG

Renormalization Group is crucial tool (mathematical, conceptual, physical)

renormalization is not about “curing or hiding divergences”, but
taking into account the physics of more and more d.o.f.s

Problem of the continuum in QG: role of RG

Renormalization Group is crucial tool (mathematical, conceptual, physical)

renormalization is not about “curing or hiding divergences”, but
taking into account the physics of more and more d.o.f.s

- for our QG models (LQG/spin foams), do not expect to have a unique continuum limit

Problem of the continuum in QG: role of RG

Renormalization Group is crucial tool (mathematical, conceptual, physical)

renormalization is not about “curing or hiding divergences”, but
taking into account the physics of more and more d.o.f.s

- for our QG models (LQG/spin foams), do not expect to have a unique continuum limit

Problem of the continuum in QG: role of RG

Renormalization Group is crucial tool (mathematical, conceptual, physical)

renormalization is not about “curing or hiding divergences”, but
taking into account the physics of more and more d.o.f.s

- for our QG models (LQG/spin foams), do not expect to have a unique continuum limit
collective behaviour of (interacting) fundamental d.o.f.s should lead to different macroscopic phases,
separated by phase transitions

Problem of the continuum in QG: role of RG

Renormalization Group is crucial tool (mathematical, conceptual, physical)

renormalization is not about “curing or hiding divergences”, but
taking into account the physics of more and more d.o.f.s

- for our QG models (LQG/spin foams), do not expect to have a unique continuum limit
collective behaviour of (interacting) fundamental d.o.f.s should lead to different macroscopic phases,
separated by phase transitions

Problem of the continuum in QG: role of RG

Renormalization Group is crucial tool (mathematical, conceptual, physical)

renormalization is not about “curing or hiding divergences”, but
taking into account the physics of more and more d.o.f.s

- for our QG models (LQG/spin foams), do not expect to have a unique continuum limit
collective behaviour of (interacting) fundamental d.o.f.s should lead to different macroscopic phases,
separated by phase transitions
- for a non-spatio-temporal QG system (e.g. LQG in GFT formulation),

Problem of the continuum in QG: role of RG

Renormalization Group is crucial tool (mathematical, conceptual, physical)

renormalization is not about “curing or hiding divergences”, but
taking into account the physics of more and more d.o.f.s

- for our QG models (LQG/spin foams), do not expect to have a unique continuum limit
collective behaviour of (interacting) fundamental d.o.f.s should lead to different macroscopic phases,
separated by phase transitions
- for a non-spatio-temporal QG system (e.g. LQG in GFT formulation),
which of the macroscopic phases is described by a smooth geometry with matter fields?

Problem of the continuum in QG: role of RG

Renormalization Group is crucial tool (mathematical, conceptual, physical)

renormalization is not about “curing or hiding divergences”, but taking into account the physics of more and more d.o.f.s

- for our QG models (LQG/spin foams), do not expect to have a unique continuum limit
- collective behaviour of (interacting) fundamental d.o.f.s should lead to different macroscopic phases, separated by phase transitions
- for a non-spatio-temporal QG system (e.g. LQG in GFT formulation), which of the macroscopic phases is described by a smooth geometry with matter fields?
- in specific GFT case:
- fundamental formulation of QG = QFT, defined perturbatively around “no-space” (degenerate) vacuum
- need to prove consistency of the theory: **perturbative GFT renormalizability**
- need to understand effective dynamics at different “GFT scales”:
RG flow of effective actions & **phase structure & phase transitions**

Geometrogenesis in LQG/GFT

idea of “geometrogenesis” in LQG/GFT :
starting from degenerate phase, continuum geometric physics in new phase
new phase can be “condensate” phase of QG “atoms of space”

in canonical LQG context:
T. Koslowski, 0709.3465 [gr-qc]

in covariant SF/GFT context:
DO, 0710.3276 [gr-qc]

also in tensor models
V. Rivasseau, ‘13

Geometrogenesis in LQG/GFT

idea of “geometrogenesis” in LQG/GFT :
starting from degenerate phase, continuum geometric physics in new phase
new phase can be “condensate” phase of QG “atoms of space”

in canonical LQG context:
T. Koslowski, 0709.3465 [gr-qc]

in covariant SF/GFT context:
DO, 0710.3276 [gr-qc]

also in tensor models
V. Rivasseau, '13

need to prove, in the full quantum dynamics, a phase transition to non-degenerate (e.g. condensate) phase

some experience and results in tensor models and GFTs

V. Bonzom, R. Gurau, A. Riello, V. Rivasseau, '11;
A. Baratin, S. Carrozza, DO, J. Ryan, M. Smerlak, '13

Geometrogenesis in LQG/GFT

idea of “geometrogenesis” in LQG/GFT :
starting from degenerate phase, continuum geometric physics in new phase
new phase can be “condensate” phase of QG “atoms of space”

in canonical LQG context:
T. Koslowski, 0709.3465 [gr-qc]

in covariant SF/GFT context:
DO, 0710.3276 [gr-qc]

also in tensor models
V. Rivasseau, '13

need to prove, in the full quantum dynamics, a phase transition to non-degenerate (e.g. condensate) phase

some experience and results in tensor models and GFTs

V. Bonzom, R. Gurau, A. Riello, V. Rivasseau, '11;
A. Baratin, S. Carrozza, DO, J. Ryan, M. Smerlak, '13

first possible interpretation:

other phases and phase transition -not- physical, just formal: theory makes sense only in geometric phase

this is point of view in CDT (J. Ambjorn, R. Loll, ...) , but see J. Mielczarek, '14

Geometrogenesis in LQG/GFT

idea of “geometrogenesis” in LQG/GFT :
starting from degenerate phase, continuum geometric physics in new phase
new phase can be “condensate” phase of QG “atoms of space”

in canonical LQG context:
T. Koslowski, 0709.3465 [gr-qc]

in covariant SF/GFT context:
DO, 0710.3276 [gr-qc]

also in tensor models
V. Rivasseau, '13

need to prove, in the full quantum dynamics, a phase transition to non-degenerate (e.g. condensate) phase

some experience and results in tensor models and GFTs

V. Bonzom, R. Gurau, A. Riello, V. Rivasseau, '11;
A. Baratin, S. Carrozza, DO, J. Ryan, M. Smerlak, '13

first possible interpretation:

other phases and phase transition -not- physical, just formal: theory makes sense only in geometric phase

this is point of view in CDT (J. Ambjorn, R. Loll, ...) , but see J. Mielczarek, '14

second possible interpretation:

other phases are physical; phase transitions are physical; we live in the geometric phase

Geometrogenesis in LQG/GFT

idea of “geometrogenesis” in LQG/GFT :
starting from degenerate phase, continuum geometric physics in new phase
new phase can be “condensate” phase of QG “atoms of space”

in canonical LQG context:
T. Koslowski, 0709.3465 [gr-qc]

in covariant SF/GFT context:
DO, 0710.3276 [gr-qc]

also in tensor models
V. Rivasseau, '13

need to prove, in the full quantum dynamics, a phase transition to non-degenerate (e.g. condensate) phase

some experience and results in tensor models and GFTs

V. Bonzom, R. Gurau, A. Riello, V. Rivasseau, '11;
A. Baratin, S. Carrozza, DO, J. Ryan, M. Smerlak, '13

first possible interpretation:

other phases and phase transition -not- physical, just formal: theory makes sense only in geometric phase

this is point of view in CDT (J. Ambjorn, R. Loll, ...) , but see J. Mielczarek, '14

second possible interpretation:

other phases are physical; phase transitions are physical; we live in the geometric phase

if geometric phase transition is physical, which physics does it describe?

Geometrogenesis in LQG/GFT

idea of “geometrogenesis” in LQG/GFT :
starting from degenerate phase, continuum geometric physics in new phase
new phase can be “condensate” phase of QG “atoms of space”

in canonical LQG context:
T. Koslowski, 0709.3465 [gr-qc]

in covariant SF/GFT context:
DO, 0710.3276 [gr-qc]

also in tensor models
V. Rivasseau, '13

need to prove, in the full quantum dynamics, a phase transition to non-degenerate (e.g. condensate) phase

some experience and results in tensor models and GFTs

V. Bonzom, R. Gurau, A. Riello, V. Rivasseau, '11;
A. Baratin, S. Carrozza, DO, J. Ryan, M. Smerlak, '13

first possible interpretation:

other phases and phase transition -not- physical, just formal: theory makes sense only in geometric phase

this is point of view in CDT (J. Ambjorn, R. Loll, ...) , but see J. Mielczarek, '14

second possible interpretation:

other phases are physical; phase transitions are physical; we live in the geometric phase

if geometric phase transition is physical, which physics does it describe?

natural hypothesis: very early Universe - big bang as QG phase transition

Geometrogenesis in LQG/GFT

Geometrogenesis in LQG/GFT

- GFT is QG analogue of QFT for atoms in condensed matter system

Geometrogenesis in LQG/GFT

- GFT is QG analogue of QFT for atoms in condensed matter system

Geometrogenesis in LQG/GFT

- GFT is QG analogue of QFT for atoms in condensed matter system

Geometrogenesis in LQG/GFT

- GFT is QG analogue of QFT for atoms in condensed matter system
- continuum spacetime (with GR-like dynamics) emerges from collective behaviour of large numbers of GFT building blocks (e.g. spin nets, simplices), possibly only in one phase of microscopic system

Geometrogenesis in LQG/GFT

- GFT is QG analogue of QFT for atoms in condensed matter system
- continuum spacetime (with GR-like dynamics) emerges from collective behaviour of large numbers of GFT building blocks (e.g. spin nets, simplices), possibly only in one phase of microscopic system

Geometrogenesis in LQG/GFT

- GFT is QG analogue of QFT for atoms in condensed matter system
- continuum spacetime (with GR-like dynamics) emerges from collective behaviour of large numbers of GFT building blocks (e.g. spin nets, simplices), possibly only in one phase of microscopic system
- continuum spacetime as a peculiar quantum fluid

Geometrogenesis in LQG/GFT

- GFT is QG analogue of QFT for atoms in condensed matter system
- continuum spacetime (with GR-like dynamics) emerges from collective behaviour of large numbers of GFT building blocks (e.g. spin nets, simplices), possibly only in one phase of microscopic system
- continuum spacetime as a peculiar quantum fluid

Geometrogenesis in LQG/GFT

- GFT is QG analogue of QFT for atoms in condensed matter system
- continuum spacetime (with GR-like dynamics) emerges from collective behaviour of large numbers of GFT building blocks (e.g. spin nets, simplices), possibly only in one phase of microscopic system
- continuum spacetime as a peculiar quantum fluid

Geometrogenesis in LQG/GFT

- GFT is QG analogue of QFT for atoms in condensed matter system
- continuum spacetime (with GR-like dynamics) emerges from collective behaviour of large numbers of GFT building blocks (e.g. spin nets, simplices), possibly only in one phase of microscopic system
- continuum spacetime as a peculiar quantum fluid
- more specific hypothesis: continuum spacetime is GFT condensate

Geometrogenesis in LQG/GFT

- GFT is QG analogue of QFT for atoms in condensed matter system
- continuum spacetime (with GR-like dynamics) emerges from collective behaviour of large numbers of GFT building blocks (e.g. spin nets, simplices), possibly only in one phase of microscopic system
- continuum spacetime as a peculiar quantum fluid
- more specific hypothesis: continuum spacetime is GFT condensate
- GR-like dynamics from GFT condensate hydrodynamics

Geometrogenesis in LQG/GFT

- GFT is QG analogue of QFT for atoms in condensed matter system
- continuum spacetime (with GR-like dynamics) emerges from collective behaviour of large numbers of GFT building blocks (e.g. spin nets, simplices), possibly only in one phase of microscopic system
- continuum spacetime as a peculiar quantum fluid
- more specific hypothesis: continuum spacetime is GFT condensate
- GR-like dynamics from GFT condensate hydrodynamics

simple candidates for physical (geometric) vacuum: GFT condensates

what is their definition? do they have a continuum geometric interpretation?
what is their effective quantum dynamics? does it relate to GR?

DO, L. Sindoni, '10; S. Gielen, DO, L. Sindoni, 1303.3576 [gr-qc], 1311.1238 [gr-qc];
S. Gielen, '14; G. Calcagni, '14; L. Sindoni, '14; S. Gielen, DO, '14

Continuum Phases & phase transitions in LQG

canonical LQG:

purely kinematical, inequivalent representations (phases) of quantum algebra of observables

Continuum Phases & phase transitions in LQG

canonical LQG:

purely kinematical, inequivalent representations (phases) of quantum algebra of observables

AL vacuum $|0\rangle_{AL}$



totally degenerate geometry (emptiest state)
connection highly fluctuating
diffeo invariant

$${}_{AL}\langle 0|E_S|0\rangle_{AL} = 0 \quad \forall S$$

$$\delta_{AL}E_S \ll 1 \quad \delta_{AL}A_S \gg 1$$

J. Lewandowski, A. Okolow, H. Sahlmann T. Thiemann '06
C. Fleischack, '06

KS vacuum $|0\rangle_{KS}$



non-degenerate geometry (triad condensate)
connection highly fluctuating
diffeo covariant

$${}_{KS}\langle 0|E_S|0\rangle_{KS} = E_S \quad \forall S$$

$$\delta_{KS}E_S \ll 1 \quad \delta_{KS}A_S \gg 1$$

T. Koslowski, H. Sahlmann, 1109.4688 [gr-qc]

DG vacuum
(or BF vacuum) $|0\rangle_{DG}$



non-degenerate flat connection
metric highly fluctuating
diffeo covariant
simplicial context

$${}_{DG}\langle 0|F(A)|0\rangle_{DG} = 0$$

$$\delta_{DG}A \ll 1 \quad \delta_{DG}E_S \gg 1$$

B. Dittrich, M. Geiller, 1401.6441 [gr-qc]

Continuum Phases & phase transitions in LQG

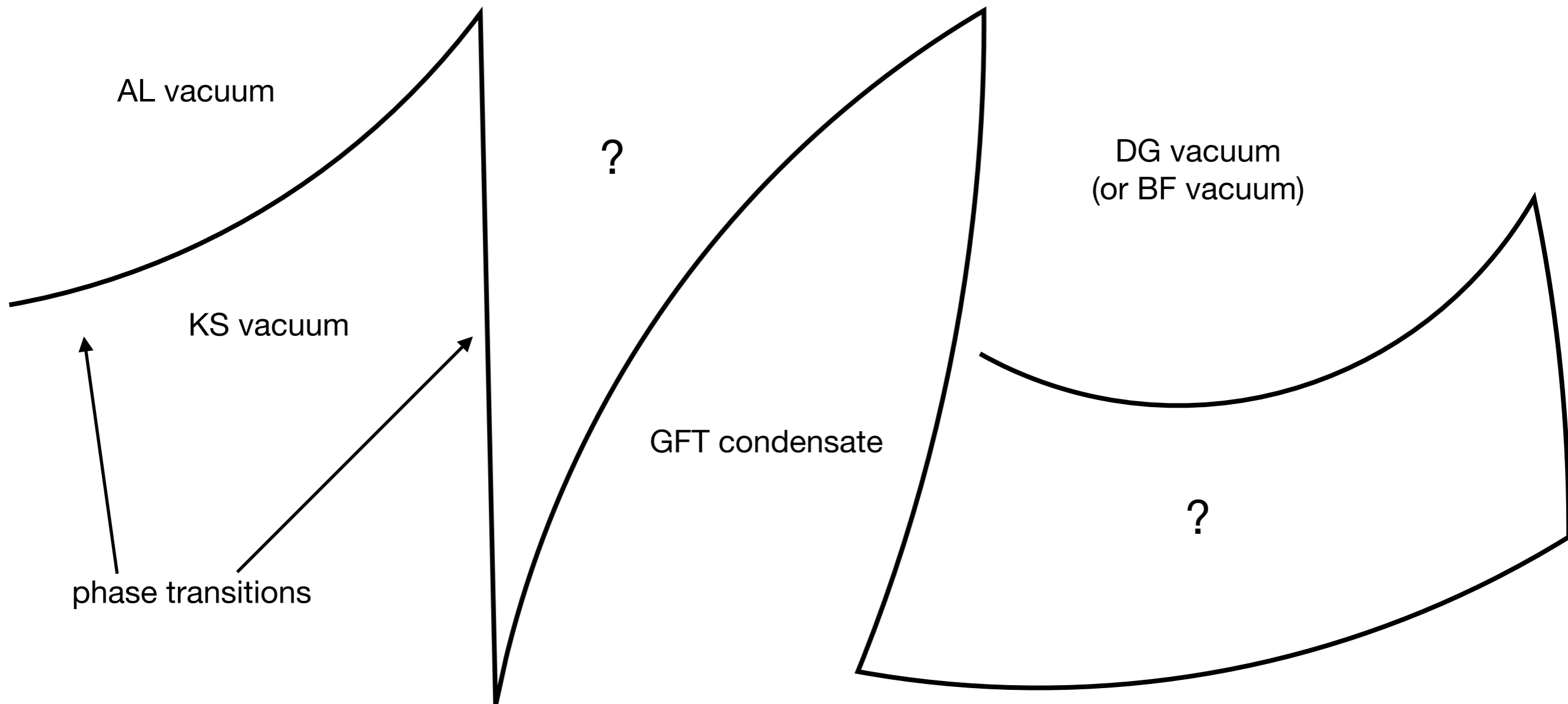
canonical LQG:

purely kinematical, inequivalent representations (phases) of quantum algebra of observables

Continuum Phases & phase transitions in LQG

canonical LQG:

purely kinematical, inequivalent representations (phases) of quantum algebra of observables



Continuum Phases & phase transitions in LQG

spin foam models (without GFT framework)

- rewrite them as lattice gauge theory path integrals
- define (background independent) coarse graining procedure
- look for flow of effective actions and fixed points

technically (numerically) very challenging

many results, mainly in simplified models (simpler algebraic data, dimensionally reduced models)

work by:

B. Bahr, B. Dittrich, F. Eckert, F. Hellmann, W. Kaminski, M. Martin-Benito, S. Steinhaus, - '09-'15

Renormalization of GFTs: where are we?

- power counting and radiative corrections in GFT models
(cut-off of fields in representation space)

- topological simplicial GFT models (BF theory):

- partial power counting and scaling theorems - large-N scaling

L. Freidel, R. Gurau, DO, '09; J. Magnen et al., '09; J. Ben Geloun, J. Magnen, V. Rivasseau, '10 ; R. Gurau, '11; S. Carrozza, DO, '11,'12

- radiative corrections of 2-point function: need for Laplacian kinetic term

J. Ben Geloun, V. Bonzom, '11

- super-renormalizability in abelian case (3d, with Laplacian)

J. Ben Geloun, '13

- 4d gravity models

- super-renormalizability of some versions of BC model

A. Perez, C. Rovelli, '00, '01

- radiative corrections of 2-point function in EPRL-FK model

T. Kraiewski, J. Magnen, V. Rivasseau, A. Tanasa, P. Vitale, '10; A. Riello, '13

Renormalization of GFTs: where are we?

- systematic renormalisation group analysis of tensorial GFT models:
 - requires subtle analysis of combinatorics of diagrams (dual to cellular complexes)

Renormalization of GFTs: where are we?

- **systematic renormalisation group analysis of tensorial GFT models:**
requires subtle analysis of combinatorics of diagrams (dual to cellular complexes)

many results: **perturbative renormalizability (around free theory) and renormalisation group flow**

J. Ben Geloun, D. Ousmane-Samary, V. Rivasseau, S. Carrozza, DO, E. Livine, F. Vignes-Tourneret, A. Tanasa, M. Raasakka,

Renormalization of GFTs: where are we?

- systematic renormalisation group analysis of tensorial GFT models:

requires subtle analysis of combinatorics of diagrams (dual to cellular complexes)

many results: perturbative renormalizability (around free theory) and renormalisation group flow

J. Ben Geloun, D. Ousmane-Samary, V. Rivasseau, S. Carrozza, DO, E. Livine, F. Vignes-Tourneret, A. Tanasa, M. Raasakka,

- several renormalizable abelian TGFT models (different groups and dimension, with/without gauge invariance)

J. Ben Geloun, V. Rivasseau, '11; J. Ben Geloun, D. Ousmane-Samary, '11 S. Carrozza, DO, V. Rivasseau, '12

- first renormalizable non-abelian TGFT model in 3d with gauge invariance (3d BF + laplacian)

S. Carrozza, DO, V. Rivasseau, '13

- proof of asymptotic freedom for abelian TGFT models without gauge invariance

J. Ben Geloun, D. Ousmane-Samary, '11; J. Ben Geloun, '12

- study of asymptotic freedom/safety for non-abelian TGFT models with gauge invariance

S. Carrozza, '14

Renormalization of GFTs: where are we?

- **systematic renormalisation group analysis of tensorial GFT models:**

requires subtle analysis of combinatorics of diagrams (dual to cellular complexes)

many results: **perturbative renormalizability (around free theory) and renormalisation group flow**

J. Ben Geloun, D. Ousmane-Samary, V. Rivasseau, S. Carrozza, DO, E. Livine, F. Vignes-Tourneret, A. Tanasa, M. Raasakka,

- **several renormalizable abelian TGFT models (different groups and dimension, with/without gauge invariance)**

J. Ben Geloun, V. Rivasseau, '11; J. Ben Geloun, D. Ousmane-Samary, '11 S. Carrozza, DO, V. Rivasseau, '12

- **first renormalizable non-abelian TGFT model in 3d with gauge invariance (3d BF + laplacian)**

S. Carrozza, DO, V. Rivasseau, '13

- **proof of asymptotic freedom for abelian TGFT models without gauge invariance**

J. Ben Geloun, D. Ousmane-Samary, '11; J. Ben Geloun, '12

- **study of asymptotic freedom/safety for non-abelian TGFT models with gauge invariance**

S. Carrozza, '14

- **Functional Renormalization Group for TGFTs**

D. Benedetti, J. Ben Geloun, DO, '14; J. Ben Geloun, R.Martini, DO, '15

Renormalization group flow of GFTs from FRG

D. Benedetti, J. Ben Geloun, DO, '14

Renormalization group flow of GFTs from FRG

D. Benedetti, J. Ben Geloun, DO, '14

an example: $\varphi(g_1, g_2, g_3) = \sum_{p_1, p_2, p_3} \varphi_{p_1 p_2 p_3} e^{ip_1 \theta_1} e^{ip_2 \theta_2} e^{ip_3 \theta_3} \in \mathbb{R}$ $g_i = e^{i\theta_i} \in U(1)$ $\theta_i \in [-\pi, \pi)$ $p_i \in \mathbb{Z}$

$$\Gamma_N(\varphi) = \frac{Z_N}{2} \text{Tr}_2(\varphi \cdot K \cdot \varphi) + \frac{m_N}{2} \text{Tr}_2(\varphi^2) + S^{\text{int}}$$

$$\Delta S_N(\phi) = \frac{1}{2} \text{Tr}(\phi \cdot R_N \cdot \phi)$$

$$S^{\text{int}} = \frac{\lambda_N}{4} \left(\text{Tr}_{4;1}(\varphi^4) + \text{Sym}(1 \rightarrow 2 \rightarrow 3) \right)$$

$$R_N(\{p_i\}; \{p'_i\}) = \delta_{p_i, p'_i} Z_N \left(N - \frac{1}{3} \sum_{i=1}^3 p_i \right) \Theta \left(N - \frac{1}{3} \sum_{i=1}^3 p_i \right)$$

$$\text{Tr}_2(\varphi \cdot K \cdot \varphi) = \sum_{p_i \in \mathbb{N}} \varphi_{123} \left(\frac{1}{3} \sum_i p_i \right) \varphi_{123}$$

$$\text{Tr}_2(\varphi^2) = \sum_{p_i \in \mathbb{N}} \varphi_{123}^2$$

$$\text{Tr}_{4;1}(\varphi^4) = \sum_{p_i, p'_i \in \mathbb{N}} \varphi_{123} \varphi_{1'23} \varphi_{1'2'3'} \varphi_{12'3'}$$

Renormalization group flow of GFTs from FRG

D. Benedetti, J. Ben Geloun, DO, '14

an example: $\varphi(g_1, g_2, g_3) = \sum_{p_1, p_2, p_3} \varphi_{p_1 p_2 p_3} e^{ip_1 \theta_1} e^{ip_2 \theta_2} e^{ip_3 \theta_3} \in \mathbb{R}$ $g_i = e^{i\theta_i} \in U(1)$ $\theta_i \in [-\pi, \pi)$ $p_i \in \mathbb{Z}$

$$\Gamma_N(\varphi) = \frac{Z_N}{2} \text{Tr}_2(\varphi \cdot K \cdot \varphi) + \frac{m_N}{2} \text{Tr}_2(\varphi^2) + S^{\text{int}}$$

$$\Delta S_N(\phi) = \frac{1}{2} \text{Tr}(\phi \cdot R_N \cdot \phi)$$

$$S^{\text{int}} = \frac{\lambda_N}{4} \left(\text{Tr}_{4;1}(\varphi^4) + \text{Sym}(1 \rightarrow 2 \rightarrow 3) \right)$$

$$R_N(\{p_i\}; \{p'_i\}) = \delta_{p_i, p'_i} Z_N \left(N - \frac{1}{3} \sum_{i=1}^3 p_i \right) \Theta \left(N - \frac{1}{3} \sum_{i=1}^3 p_i \right)$$

$$\text{Tr}_2(\varphi \cdot K \cdot \varphi) = \sum_{p_i \in \mathbb{N}} \varphi_{123} \left(\frac{1}{3} \sum_i p_i \right) \varphi_{123}$$

$$\text{Tr}_2(\varphi^2) = \sum_{p_i \in \mathbb{N}} \varphi_{123}^2$$

$$\text{Tr}_{4;1}(\varphi^4) = \sum_{p_i, p'_i \in \mathbb{N}} \varphi_{123} \varphi_{1'23} \varphi_{1'2'3'} \varphi_{12'3'}$$

flow equations:

$$\partial_t \mu_N = -\mu_N \eta - \frac{\lambda_N N}{(N + \mu_N)^2} \left\{ \frac{9}{2} (3N + 2)(N + 1) + \frac{\eta}{2} (11 + 18N + 9N^2) \right\}$$

$$\partial_t \lambda_N = -\frac{1}{3} \lambda_N^2 \frac{\eta}{(N + \mu_N)^3} N(3N + 1)(6N + 13) + 9 \lambda_N^2 \frac{(\eta + 1)}{(N + \mu_N)^3} N(N + 1)(N + 2) - \dots$$

non-autonomous system (due to external scale $a = \text{size of grow manifold}$)

Renormalization group flow of GFTs from FRG

D. Benedetti, J. Ben Geloun, DO, '14

an example: $\varphi(g_1, g_2, g_3) = \sum_{p_1, p_2, p_3} \varphi_{p_1 p_2 p_3} e^{ip_1 \theta_1} e^{ip_2 \theta_2} e^{ip_3 \theta_3} \in \mathbb{R}$ $g_i = e^{i\theta_i} \in U(1)$ $\theta_i \in [-\pi, \pi)$ $p_i \in \mathbb{Z}$

$$\Gamma_N(\varphi) = \frac{Z_N}{2} \text{Tr}_2(\varphi \cdot K \cdot \varphi) + \frac{m_N}{2} \text{Tr}_2(\varphi^2) + S^{\text{int}}$$

$$\Delta S_N(\phi) = \frac{1}{2} \text{Tr}(\phi \cdot R_N \cdot \phi)$$

$$S^{\text{int}} = \frac{\lambda_N}{4} \left(\text{Tr}_{4;1}(\varphi^4) + \text{Sym}(1 \rightarrow 2 \rightarrow 3) \right)$$

$$R_N(\{p_i\}; \{p'_i\}) = \delta_{p_i, p'_i} Z_N \left(N - \frac{1}{3} \sum_{i=1}^3 p_i \right) \Theta \left(N - \frac{1}{3} \sum_{i=1}^3 p_i \right)$$

$$\text{Tr}_2(\varphi \cdot K \cdot \varphi) = \sum_{p_i \in \mathbb{N}} \varphi_{123} \left(\frac{1}{3} \sum_i p_i \right) \varphi_{123}$$

$$\text{Tr}_2(\varphi^2) = \sum_{p_i \in \mathbb{N}} \varphi_{123}^2$$

$$\text{Tr}_{4;1}(\varphi^4) = \sum_{p_i, p'_i \in \mathbb{N}} \varphi_{123} \varphi_{1'23} \varphi_{1'2'3'} \varphi_{12'3'}$$

Renormalization group flow of GFTs from FRG

D. Benedetti, J. Ben Geloun, DO, '14

an example: $\varphi(g_1, g_2, g_3) = \sum_{p_1, p_2, p_3} \varphi_{p_1 p_2 p_3} e^{ip_1 \theta_1} e^{ip_2 \theta_2} e^{ip_3 \theta_3} \in \mathbb{R}$ $g_i = e^{i\theta_i} \in U(1)$ $\theta_i \in [-\pi, \pi)$ $p_i \in \mathbb{Z}$

$$\Gamma_N(\varphi) = \frac{Z_N}{2} \text{Tr}_2(\varphi \cdot K \cdot \varphi) + \frac{m_N}{2} \text{Tr}_2(\varphi^2) + S^{\text{int}}$$

$$\Delta S_N(\phi) = \frac{1}{2} \text{Tr}(\phi \cdot R_N \cdot \phi)$$

$$S^{\text{int}} = \frac{\lambda_N}{4} \left(\text{Tr}_{4;1}(\varphi^4) + \text{Sym}(1 \rightarrow 2 \rightarrow 3) \right)$$

$$R_N(\{p_i\}; \{p'_i\}) = \delta_{p_i, p'_i} Z_N \left(N - \frac{1}{3} \sum_{i=1}^3 p_i \right) \Theta \left(N - \frac{1}{3} \sum_{i=1}^3 p_i \right)$$

$$\text{Tr}_2(\varphi \cdot K \cdot \varphi) = \sum_{p_i \in \mathbb{N}} \varphi_{123} \left(\frac{1}{3} \sum_i p_i \right) \varphi_{123}$$

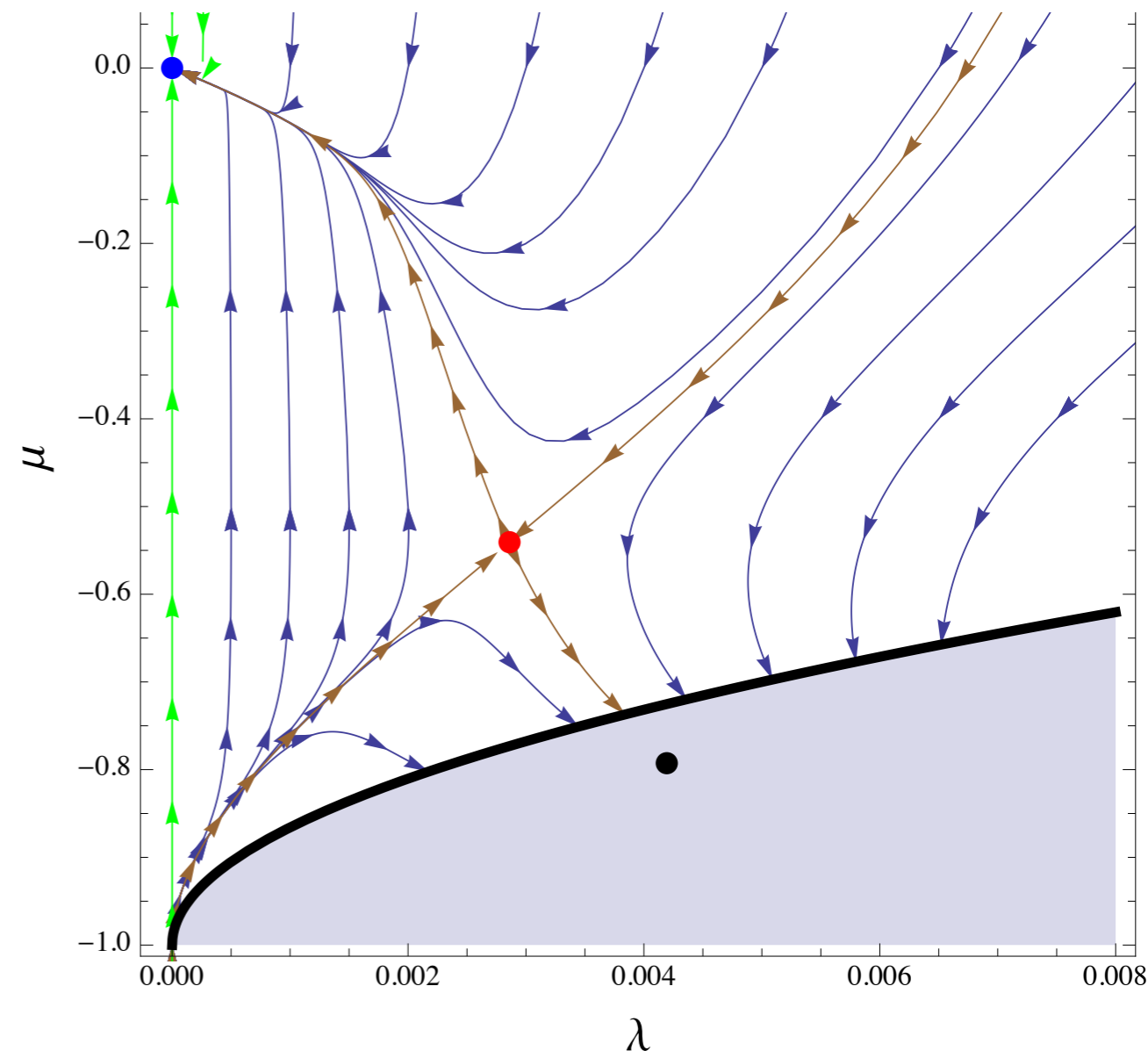
$$\text{Tr}_2(\varphi^2) = \sum_{p_i \in \mathbb{N}} \varphi_{123}^2$$

$$\text{Tr}_{4;1}(\varphi^4) = \sum_{p_i, p'_i \in \mathbb{N}} \varphi_{123} \varphi_{1'2'3} \varphi_{1'2'3'} \varphi_{12'3'}$$

large-N regime (“formal UV”):

asymptotic freedom

(Gaussian fixed point with two relev. directions)



Renormalization group flow of GFTs from FRG

D. Benedetti, J. Ben Geloun, DO, '14

an example: $\varphi(g_1, g_2, g_3) = \sum_{p_1, p_2, p_3} \varphi_{p_1 p_2 p_3} e^{ip_1 \theta_1} e^{ip_2 \theta_2} e^{ip_3 \theta_3} \in \mathbb{R}$ $g_i = e^{i\theta_i} \in U(1)$ $\theta_i \in [-\pi, \pi)$ $p_i \in \mathbb{Z}$

$$\Gamma_N(\varphi) = \frac{Z_N}{2} \text{Tr}_2(\varphi \cdot K \cdot \varphi) + \frac{m_N}{2} \text{Tr}_2(\varphi^2) + S^{\text{int}}$$

$$\Delta S_N(\phi) = \frac{1}{2} \text{Tr}(\phi \cdot R_N \cdot \phi)$$

$$S^{\text{int}} = \frac{\lambda_N}{4} \left(\text{Tr}_{4;1}(\varphi^4) + \text{Sym}(1 \rightarrow 2 \rightarrow 3) \right)$$

$$R_N(\{p_i\}; \{p'_i\}) = \delta_{p_i, p'_i} Z_N \left(N - \frac{1}{3} \sum_{i=1}^3 p_i \right) \Theta \left(N - \frac{1}{3} \sum_{i=1}^3 p_i \right)$$

$$\text{Tr}_2(\varphi \cdot K \cdot \varphi) = \sum_{p_i \in \mathbb{N}} \varphi_{123} \left(\frac{1}{3} \sum_i p_i \right) \varphi_{123}$$

$$\text{Tr}_2(\varphi^2) = \sum_{p_i \in \mathbb{N}} \varphi_{123}^2$$

$$\text{Tr}_{4;1}(\varphi^4) = \sum_{p_i, p'_i \in \mathbb{N}} \varphi_{123} \varphi_{1'23} \varphi_{1'2'3'} \varphi_{12'3'}$$

Renormalization group flow of GFTs from FRG

D. Benedetti, J. Ben Geloun, DO, '14

an example: $\varphi(g_1, g_2, g_3) = \sum_{p_1, p_2, p_3} \varphi_{p_1 p_2 p_3} e^{ip_1 \theta_1} e^{ip_2 \theta_2} e^{ip_3 \theta_3} \in \mathbb{R}$ $g_i = e^{i\theta_i} \in U(1)$ $\theta_i \in [-\pi, \pi)$ $p_i \in \mathbb{Z}$

$$\Gamma_N(\varphi) = \frac{Z_N}{2} \text{Tr}_2(\varphi \cdot K \cdot \varphi) + \frac{m_N}{2} \text{Tr}_2(\varphi^2) + S^{\text{int}}$$

$$\Delta S_N(\phi) = \frac{1}{2} \text{Tr}(\phi \cdot R_N \cdot \phi)$$

$$S^{\text{int}} = \frac{\lambda_N}{4} \left(\text{Tr}_{4;1}(\varphi^4) + \text{Sym}(1 \rightarrow 2 \rightarrow 3) \right)$$

$$R_N(\{p_i\}; \{p'_i\}) = \delta_{p_i, p'_i} Z_N \left(N - \frac{1}{3} \sum_{i=1}^3 p_i \right) \Theta \left(N - \frac{1}{3} \sum_{i=1}^3 p_i \right)$$

$$\text{Tr}_2(\varphi \cdot K \cdot \varphi) = \sum_{p_i \in \mathbb{N}} \varphi_{123} \left(\frac{1}{3} \sum_i p_i \right) \varphi_{123}$$

$$\text{Tr}_2(\varphi^2) = \sum_{p_i \in \mathbb{N}} \varphi_{123}^2$$

$$\text{Tr}_{4;1}(\varphi^4) = \sum_{p_i, p'_i \in \mathbb{N}} \varphi_{123} \varphi_{1'2'3} \varphi_{1'2'3'} \varphi_{12'3'}$$

small-N regime (“formal IR”):

two IR fixed points:

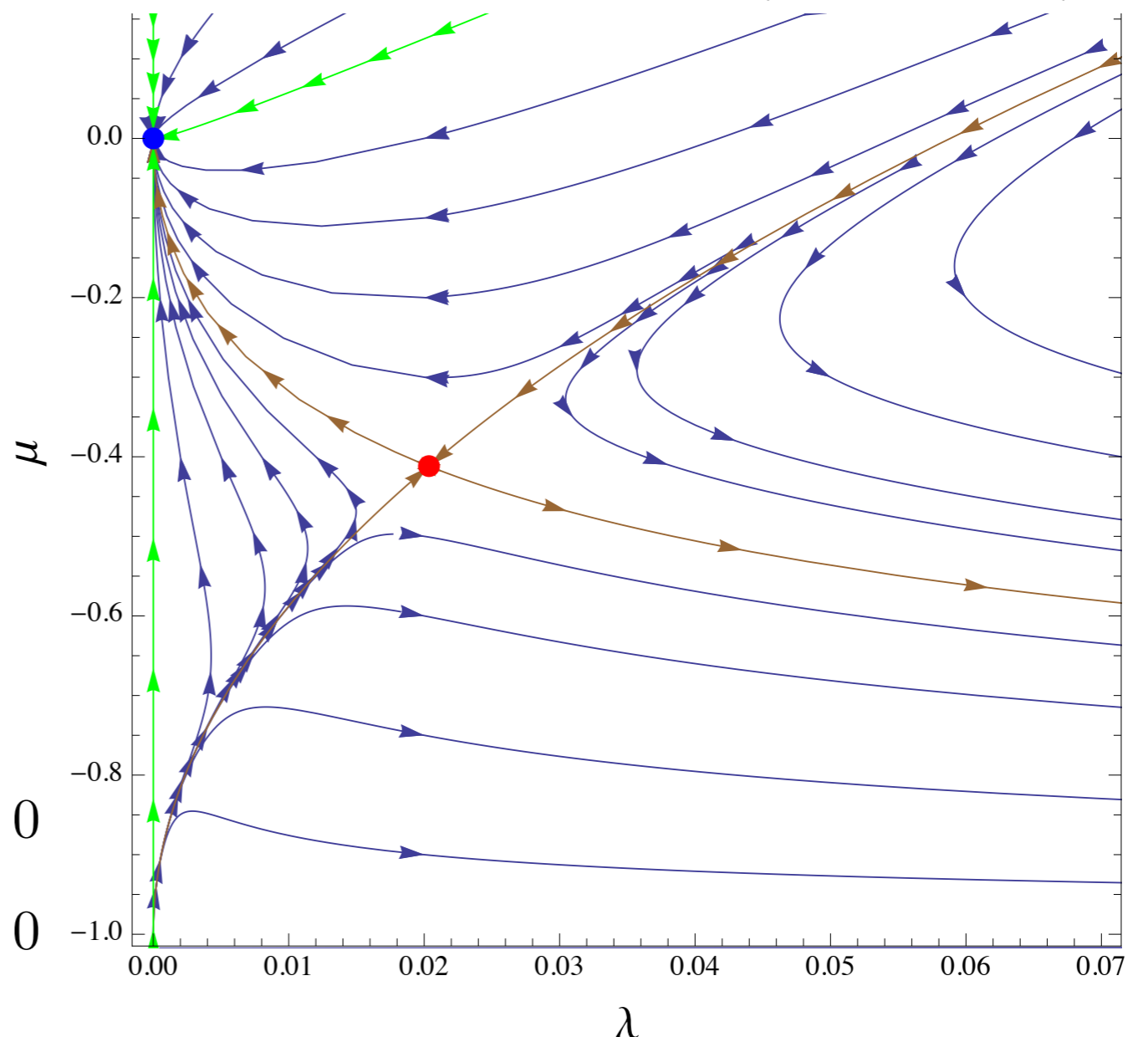
Gaussian - 2 relevant directions

non-Gaussian - 1 rel., 1 irrel. directions

two phases:

“degenerate” $\mu > 0$ $\langle \varphi \rangle = 0$

“condensate” $\mu < 0$ $\langle \varphi \rangle \neq 0$



Renormalization group flow of GFTs from FRG

D. Benedetti, J. Ben Geloun, DO, '14

an example: $\varphi(g_1, g_2, g_3) = \sum_{p_1, p_2, p_3} \varphi_{p_1 p_2 p_3} e^{ip_1 \theta_1} e^{ip_2 \theta_2} e^{ip_3 \theta_3} \in \mathbb{R}$ $g_i = e^{i\theta_i} \in U(1)$ $\theta_i \in [-\pi, \pi)$ $p_i \in \mathbb{Z}$

$$\Gamma_N(\varphi) = \frac{Z_N}{2} \text{Tr}_2(\varphi \cdot K \cdot \varphi) + \frac{m_N}{2} \text{Tr}_2(\varphi^2) + S^{\text{int}}$$

$$\Delta S_N(\phi) = \frac{1}{2} \text{Tr}(\phi \cdot R_N \cdot \phi)$$

$$S^{\text{int}} = \frac{\lambda_N}{4} \left(\text{Tr}_{4;1}(\varphi^4) + \text{Sym}(1 \rightarrow 2 \rightarrow 3) \right)$$

$$R_N(\{p_i\}; \{p'_i\}) = \delta_{p_i, p'_i} Z_N \left(N - \frac{1}{3} \sum_{i=1}^3 p_i \right) \Theta \left(N - \frac{1}{3} \sum_{i=1}^3 p_i \right)$$

$$\text{Tr}_2(\varphi \cdot K \cdot \varphi) = \sum_{p_i \in \mathbb{N}} \varphi_{123} \left(\frac{1}{3} \sum_i p_i \right) \varphi_{123}$$

$$\text{Tr}_2(\varphi^2) = \sum_{p_i \in \mathbb{N}} \varphi_{123}^2$$

$$\text{Tr}_{4;1}(\varphi^4) = \sum_{p_i, p'_i \in \mathbb{N}} \varphi_{123} \varphi_{1'23} \varphi_{1'2'3'} \varphi_{12'3'}$$

2nd message

problem of the continuum in QG (GFT,LQG): crucial to connect to macrophysics, difficult and open

we are addressing it

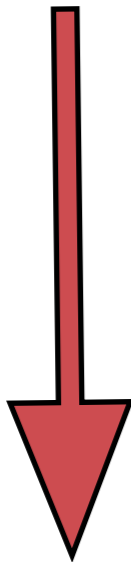
RG is crucial tool, different strategies, many results (renormalizable models, RG flows,...)

QG phase transition (condensation?) could be physical

cosmological interpretation: realization of “Emergent Spacetime” and of “Universe as a Condensate” ideas

Quantum spacetime: the difficult path from microstructure to cosmology

Quantum Gravity problem:
identify microscopic d.o.f. of quantum spacetime and their fundamental dynamics



derive effective (QG-inspired) models for fundamental (quantum) cosmology:
explain features of early Universe, obtain testable QG predictions

various models: loop quantum cosmology,

task is daunting (imagine analogue problem in condensed matter theory)

(Quantum) Cosmology from GFT

S. Gielen, DO, L. Sindoni, PRL, [arXiv:1303.3576 \[gr-qc\]](#); JHEP, [arXiv:1311.1238 \[gr-qc\]](#)

problem 1:

identify quantum states in fundamental theory with continuum spacetime interpretation

problem 2:

extract from fundamental theory an effective macroscopic dynamics for such states

(Quantum) Cosmology from GFT

S. Gielen, DO, L. Sindoni, PRL, [arXiv:1303.3576 \[gr-qc\]](#); JHEP, [arXiv:1311.1238 \[gr-qc\]](#)

problem 1:

identify quantum states in fundamental theory with continuum spacetime interpretation

many results in LQG (weaves, coherent states, statistical geometry, approximate symmetric states,...)

problem 2:

extract from fundamental theory an effective macroscopic dynamics for such states

(Quantum) Cosmology from GFT

S. Gielen, DO, L. Sindoni, PRL, [arXiv:1303.3576 \[gr-qc\]](#); JHEP, [arXiv:1311.1238 \[gr-qc\]](#)

problem 1:

identify quantum states in fundamental theory with continuum spacetime interpretation

many results in LQG (weaves, coherent states, statistical geometry, approximate symmetric states,...)

Quantum GFT condensates are continuum (homogeneous) (quantum) spaces

problem 2:

extract from fundamental theory an effective macroscopic dynamics for such states

(Quantum) Cosmology from GFT

S. Gielen, DO, L. Sindoni, PRL, [arXiv:1303.3576 \[gr-qc\]](#); JHEP, [arXiv:1311.1238 \[gr-qc\]](#)

problem 1:

identify quantum states in fundamental theory with continuum spacetime interpretation

many results in LQG (weaves, coherent states, statistical geometry, approximate symmetric states,...)

Quantum GFT condensates are continuum (homogeneous) (quantum) spaces

described by single collective wave function
(depending on homogeneous anisotropic geometric data)

problem 2:

extract from fundamental theory an effective macroscopic dynamics for such states

(Quantum) Cosmology from GFT

S. Gielen, DO, L. Sindoni, PRL, [arXiv:1303.3576 \[gr-qc\]](#); JHEP, [arXiv:1311.1238 \[gr-qc\]](#)

problem 1:

identify quantum states in fundamental theory with continuum spacetime interpretation

many results in LQG (weaves, coherent states, statistical geometry, approximate symmetric states,...)

Quantum GFT condensates are continuum (homogeneous) (quantum) spaces

described by single collective wave function
(depending on homogeneous anisotropic geometric data)

similar constructions in LQG (Alesci, Cianfrani) and LQC (Bojowald, Wilson-Ewing,

problem 2:

extract from fundamental theory an effective macroscopic dynamics for such states

(Quantum) Cosmology from GFT

S. Gielen, DO, L. Sindoni, PRL, [arXiv:1303.3576 \[gr-qc\]](#); JHEP, [arXiv:1311.1238 \[gr-qc\]](#)

problem 1:

identify quantum states in fundamental theory with continuum spacetime interpretation

many results in LQG (weaves, coherent states, statistical geometry, approximate symmetric states,...)

Quantum GFT condensates are continuum (homogeneous) (quantum) spaces

described by single collective wave function
(depending on homogeneous anisotropic geometric data)

similar constructions in LQG (Alesci, Cianfrani) and LQC (Bojowald, Wilson-Ewing,

problem 2:

extract from fundamental theory an effective macroscopic dynamics for such states

following procedures of standard BEC

(Quantum) Cosmology from GFT

S. Gielen, DO, L. Sindoni, PRL, [arXiv:1303.3576 \[gr-qc\]](#); JHEP, [arXiv:1311.1238 \[gr-qc\]](#)

problem 1:

identify quantum states in fundamental theory with continuum spacetime interpretation

many results in LQG (weaves, coherent states, statistical geometry, approximate symmetric states,....)

Quantum GFT condensates are continuum (homogeneous) (quantum) spaces

described by single collective wave function
(depending on homogeneous anisotropic geometric data)

similar constructions in LQG (Alesci, Cianfrani) and LQC (Bojowald, Wilson-Ewing,

problem 2:

extract from fundamental theory an effective macroscopic dynamics for such states

following procedures of standard BEC

QG (GFT) analogue of Gross-Pitaevskii hydrodynamic equation in BECs

is

non-linear and non-local extension of quantum cosmology equation for collective wave function

(Quantum) Cosmology from GFT

S. Gielen, DO, L. Sindoni, PRL, [arXiv:1303.3576 \[gr-qc\]](#); JHEP, [arXiv:1311.1238 \[gr-qc\]](#)

problem 1:

identify quantum states in fundamental theory with continuum spacetime interpretation

many results in LQG (weaves, coherent states, statistical geometry, approximate symmetric states,...)

Quantum GFT condensates are continuum (homogeneous) (quantum) spaces

described by single collective wave function
(depending on homogeneous anisotropic geometric data)

similar constructions in LQG (Alesci, Cianfrani) and LQC (Bojowald, Wilson-Ewing,

problem 2:

extract from fundamental theory an effective macroscopic dynamics for such states

following procedures of standard BEC

QG (GFT) analogue of Gross-Pitaevskii hydrodynamic equation in BECs
is

non-linear and non-local extension of quantum cosmology equation for collective wave function

similar equations obtained in non-linear extension of LQC (Bojowald et al. '12)

GFT states and approximate continuum geometries

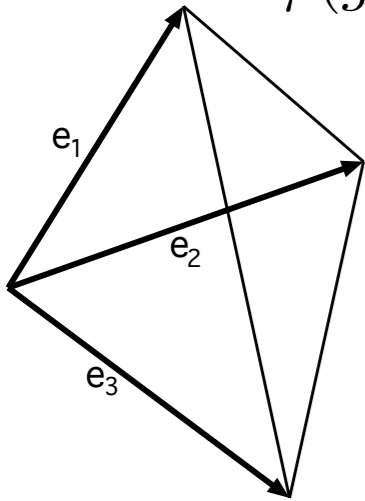
- work with **GFT with simplicial geometric interpretation** (Riemannian SO(4) case for simplicity)

$$\varphi(g_1, g_2, g_3, g_4) \leftrightarrow \varphi(B_1, B_2, B_3, B_4) \rightarrow \mathbb{C}$$

describes geometric tetrahedron

$$B_i^{AB} = \epsilon_i^{jk} e_j^A e_k^B$$

(closure + simplicity constraints)



GFT states and approximate continuum geometries

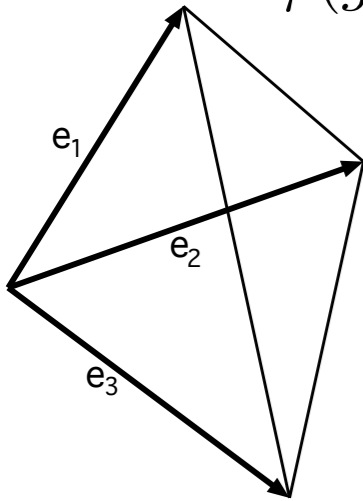
- work with **GFT with simplicial geometric interpretation** (Riemannian SO(4) case for simplicity)

$$\varphi(g_1, g_2, g_3, g_4) \leftrightarrow \varphi(B_1, B_2, B_3, B_4) \rightarrow \mathbb{C}$$

describes geometric tetrahedron

$$B_i^{AB} = \epsilon_i^{jk} e_j^A e_k^B$$

(closure + simplicity constraints)



many results in LQG,
simplicial geometry

GFT states and approximate continuum geometries

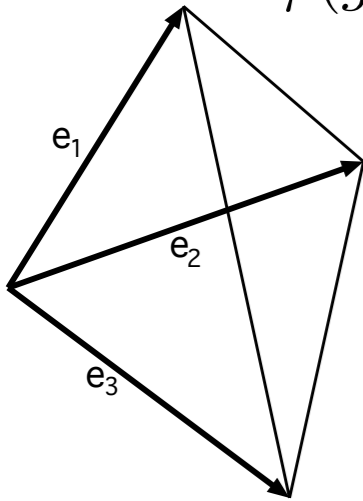
- work with **GFT with simplicial geometric interpretation** (Riemannian SO(4) case for simplicity)

$$\varphi(g_1, g_2, g_3, g_4) \leftrightarrow \varphi(B_1, B_2, B_3, B_4) \rightarrow \mathbb{C}$$

describes geometric tetrahedron

$$B_i^{AB} = \epsilon_i^{jk} e_j^A e_k^B$$

(closure + simplicity constraints)



many results in LQG,
simplicial geometry

- generic N-particle GFT state** (N geometric tetrahedra):

$$|B_{I(m)}\rangle := \prod_{m=1}^N \hat{\varphi}^\dagger(B_{1(m)}, \dots, B_{4(m)})|0\rangle$$

GFT states and approximate continuum geometries

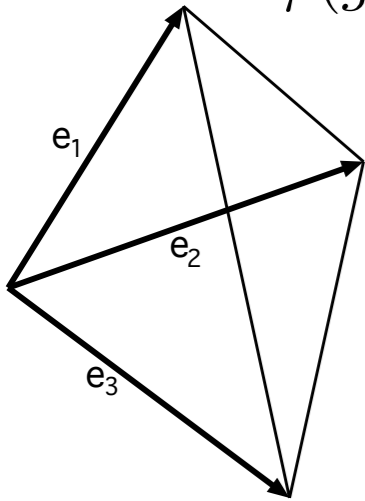
- work with **GFT with simplicial geometric interpretation** (Riemannian SO(4) case for simplicity)

$$\varphi(g_1, g_2, g_3, g_4) \leftrightarrow \varphi(B_1, B_2, B_3, B_4) \rightarrow \mathbb{C}$$

describes geometric tetrahedron

$$B_i^{AB} = \epsilon_i^{jkl} e_j^A e_k^B$$

(closure + simplicity constraints)



many results in LQG,
simplicial geometry

- generic N-particle GFT state** (N geometric tetrahedra):

$$|B_{I(m)}\rangle := \prod_{m=1}^N \hat{\varphi}^\dagger(B_{1(m)}, \dots, B_{4(m)})|0\rangle$$

- from B's of each GFT quantum,
construct:

$$g_{ij} = \frac{1}{8 \operatorname{tr}(B_1 B_2 B_3)} \epsilon_i^{kl} \epsilon_j^{mn} \tilde{B}_{km} \tilde{B}_{ln} \quad \tilde{B}_{ij} := B_i^{AB} B_j{}_{AB}$$

interpretation: spatial **metric coefficients (and conjugate variables)** “sampled” at N points

$$B_{I(m)} \leftrightarrow g_{ij}(x_m) \leftrightarrow a_i(x_m) \quad g_{I(m)} \leftrightarrow K_{ij}(x_m) \leftrightarrow p_{a_i}(x_m)$$

GFT states and approximate continuum geometries

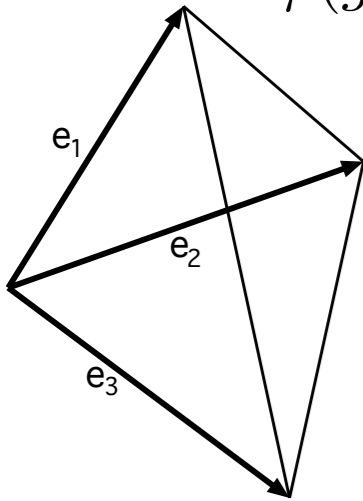
- work with **GFT with simplicial geometric interpretation** (Riemannian SO(4) case for simplicity)

$$\varphi(g_1, g_2, g_3, g_4) \leftrightarrow \varphi(B_1, B_2, B_3, B_4) \rightarrow \mathbb{C}$$

describes geometric tetrahedron

$$B_i^{AB} = \epsilon_i^{jk} e_j^A e_k^B$$

(closure + simplicity constraints)



many results in LQG,
simplicial geometry

- generic N-particle GFT state** (N geometric tetrahedra):

$$|B_{I(m)}\rangle := \prod_{m=1}^N \hat{\varphi}^\dagger(B_{1(m)}, \dots, B_{4(m)})|0\rangle$$

- from B's of each GFT quantum,
construct:

$$g_{ij} = \frac{1}{8 \text{tr}(B_1 B_2 B_3)} \epsilon_i^{kl} \epsilon_j^{mn} \tilde{B}_{km} \tilde{B}_{ln} \quad \tilde{B}_{ij} := B_i^{AB} B_{j AB}$$

interpretation: spatial **metric coefficients (and conjugate variables)** “sampled” at N points

$$B_{I(m)} \leftrightarrow g_{ij}(x_m) \leftrightarrow a_i(x_m) \quad g_{I(m)} \leftrightarrow K_{ij}(x_m) \leftrightarrow p_{a_i}(x_m)$$

- classical criterion for homogeneity (for GFT data):** $g_{ij(m)} = g_{ij(k)} \quad \forall k, m = 1, \dots, N$

i.e. all GFT quanta are labelled by the same (gauge invariant) data

Homogeneous geometries & GFT condensates

Homogeneous geometries & GFT condensates

- lift homogeneity criterion to quantum level (and include conjugate information):

Homogeneous geometries & GFT condensates

- lift homogeneity criterion to quantum level (and include conjugate information):

all GFT quanta have the same (gauge invariant) “wave function”, i.e. are in the same quantum state

$$\Psi (B_{i(1)}, \dots, B_{i(N)}) = \frac{1}{N!} \prod_{m=1}^N \Phi(B_i(m))$$

Homogeneous geometries & GFT condensates

- lift homogeneity criterion to quantum level (and include conjugate information):

all GFT quanta have the same (gauge invariant) “wave function”, i.e. are in the same quantum state

$$\Psi (B_{i(1)}, \dots, B_{i(N)}) = \frac{1}{N!} \prod_{m=1}^N \Phi(B_i(m))$$

- in GFT: such states can be expressed in 2nd quantized language and one can consider superpositions of states of arbitrary N

Homogeneous geometries & GFT condensates

- lift homogeneity criterion to quantum level (and include conjugate information):

all GFT quanta have the same (gauge invariant) “wave function”, i.e. are in the same quantum state

$$\Psi (B_{i(1)}, \dots, B_{i(N)}) = \frac{1}{N!} \prod_{m=1}^N \Phi(B_i(m))$$

- in GFT: such states can be expressed in 2nd quantized language and one can consider superpositions of states of arbitrary N
- sending N to infinity means improving arbitrarily the accuracy of the sampling

Homogeneous geometries & GFT condensates

- lift homogeneity criterion to quantum level (and include conjugate information):

all GFT quanta have the same (gauge invariant) “wave function”, i.e. are in the same quantum state

$$\Psi (B_{i(1)}, \dots, B_{i(N)}) = \frac{1}{N!} \prod_{m=1}^N \Phi(B_i(m))$$

- in GFT: such states can be expressed in 2nd quantized language and one can consider superpositions of states of arbitrary N
- sending N to infinity means improving arbitrarily the accuracy of the sampling



quantum GFT condensates are continuum (homogeneous) (quantum) spaces

Homogeneous geometries & GFT condensates

- lift homogeneity criterion to quantum level (and include conjugate information):

all GFT quanta have the same (gauge invariant) “wave function”, i.e. are in the same quantum state

$$\Psi (B_{i(1)}, \dots, B_{i(N)}) = \frac{1}{N!} \prod_{m=1}^N \Phi(B_i(m))$$

- in GFT: such states can be expressed in 2nd quantized language and one can consider superpositions of states of arbitrary N
- sending N to infinity means improving arbitrarily the accuracy of the sampling



quantum GFT condensates are continuum (homogeneous) (quantum) spaces

similar constructions in LQG (Alesci, Cianfrani) and LQC (Bojowald, Wilson-Ewing,

Quantum GFT condensates

a simple choice of **quantum GFT condensate**
(**homogeneous continuum quantum space**)

other constructions possible, depending on how much information the condensate state has to encode (in a coarse-grained form)

S. Gielen, DO, L. Sindoni, '13; DO, D. Pranzetti, J. Ryan, L. Sindoni, '15

various procedures for estimating validity of chosen ansatz for vacuum state,
e.g. L. Sindoni, arXiv:1408.3095 [gr-qc]

Quantum GFT condensates

a simple choice of **quantum GFT condensate**
(**homogeneous continuum quantum space**)

single-particle condensate
(Gross-Pitaevskii approximation)

other constructions possible, depending on how much information the condensate state has to encode (in a coarse-grained form)

S. Gielen, DO, L. Sindoni, '13; DO, D. Pranzetti, J. Ryan, L. Sindoni, '15

various procedures for estimating validity of chosen ansatz for vacuum state,
e.g. L. Sindoni, arXiv:1408.3095 [gr-qc]

$$|\sigma\rangle := \exp(\hat{\sigma}) |0\rangle$$

$$\hat{\sigma} := \int d^4g \sigma(g_I) \hat{\varphi}^\dagger(g_I) \quad \sigma(g_I k) = \sigma(g_I)$$

Quantum GFT condensates

a simple choice of **quantum GFT condensate**
(**homogeneous continuum quantum space**)

single-particle condensate
(Gross-Pitaevskii approximation)

other constructions possible, depending on how much information the condensate state has to encode (in a coarse-grained form)

S. Gielen, DO, L. Sindoni, '13; DO, D. Pranzetti, J. Ryan, L. Sindoni, '15

various procedures for estimating validity of chosen ansatz for vacuum state,
e.g. L. Sindoni, arXiv:1408.3095 [gr-qc]

$$|\sigma\rangle := \exp(\hat{\sigma}) |0\rangle$$

$$\hat{\sigma} := \int d^4g \sigma(g_I) \hat{\varphi}^\dagger(g_I) \quad \sigma(g_I k) = \sigma(g_I)$$

- data for homogeneous anisotropic geometries
- truly non-perturbative quantum states (infinite QG dofs, superposition of graphs)
- **support perturbations at any sampling scale N**
- 2nd quantized coherent states $\hat{\varphi}(g_I)|\sigma\rangle = \sigma(g_I)|\sigma\rangle$
- can be studied using BEC techniques

Effective cosmological dynamics from GFT

follow closely procedure used in real BECs

single-particle GFT condensate:

$$|\sigma\rangle := \exp(\hat{\sigma}) |0\rangle \quad \hat{\sigma} := \int d^4g \sigma(g_I) \hat{\varphi}^\dagger(g_I) \quad \sigma(g_I k) = \sigma(g_I)$$

S. Gielen, DO, L. Sindoni,

PRL, [arXiv:1303.3576 \[gr-qc\]](#);

JHEP, [arXiv:1311.1238 \[gr-qc\]](#)

Effective cosmological dynamics from GFT

follow closely procedure used in real BECs

S. Gielen, DO, L. Sindoni,

PRL, [arXiv:1303.3576 \[gr-qc\]](#);

JHEP, [arXiv:1311.1238 \[gr-qc\]](#)

single-particle GFT condensate:

$$|\sigma\rangle := \exp(\hat{\sigma}) |0\rangle \quad \hat{\sigma} := \int d^4g \sigma(g_I) \hat{\varphi}^\dagger(g_I) \quad \sigma(g_I k) = \sigma(g_I)$$

from truncation of SD equations for GFT model

applied to (coherent) GFT condensate state,
gives equation for “wave function”:

$$\int [dg'_i] \tilde{\mathcal{K}}(g_i, g'_i) \sigma(g'_i) + \lambda \frac{\delta \tilde{\mathcal{V}}}{\delta \varphi(g_i)} \Big|_{\varphi \equiv \sigma} = 0$$

basically (up to some approximations), the “classical GFT eqns”

Effective cosmological dynamics from GFT

follow closely procedure used in real BECs

S. Gielen, DO, L. Sindoni,

PRL, [arXiv:1303.3576 \[gr-qc\]](#);

JHEP, [arXiv:1311.1238 \[gr-qc\]](#)

single-particle GFT condensate:

$$|\sigma\rangle := \exp(\hat{\sigma}) |0\rangle \quad \hat{\sigma} := \int d^4g \sigma(g_I) \hat{\varphi}^\dagger(g_I) \quad \sigma(g_I k) = \sigma(g_I)$$

from truncation of SD equations for GFT model

applied to (coherent) GFT condensate state,
gives equation for “wave function”:

$$\int [dg'_i] \tilde{\mathcal{K}}(g_i, g'_i) \sigma(g'_i) + \lambda \frac{\delta \tilde{\mathcal{V}}}{\delta \varphi(g_i)} \Big|_{\varphi \equiv \sigma} = 0$$

basically (up to some approximations), the “classical GFT eqns”

non-linear and non-local extension of quantum cosmology-like equation for “collective wave function”

QG (GFT) analogue of Gross-Pitaevskii hydrodynamic equation in BECs

Effective cosmological dynamics from GFT

follow closely procedure used in real BECs

S. Gielen, DO, L. Sindoni,

PRL, [arXiv:1303.3576 \[gr-qc\]](#);

JHEP, [arXiv:1311.1238 \[gr-qc\]](#)

single-particle GFT condensate:

$$|\sigma\rangle := \exp(\hat{\sigma}) |0\rangle \quad \hat{\sigma} := \int d^4g \sigma(g_I) \hat{\varphi}^\dagger(g_I) \quad \sigma(g_I k) = \sigma(g_I)$$

from truncation of SD equations for GFT model

applied to (coherent) GFT condensate state,
gives equation for “wave function”:

$$\int [dg'_i] \tilde{\mathcal{K}}(g_i, g'_i) \sigma(g'_i) + \lambda \frac{\delta \tilde{\mathcal{V}}}{\delta \varphi(g_i)} \Big|_{\varphi \equiv \sigma} = 0$$

basically (up to some approximations), the “classical GFT eqns”

non-linear and non-local extension of quantum cosmology-like equation for “collective wave function”

QG (GFT) analogue of Gross-Pitaevskii hydrodynamic equation in BECs

similar equations obtained in non-linear extension of LQC (Bojowald et al. '12)

Effective cosmological dynamics from GFT

follow closely procedure used in real BECs

S. Gielen, DO, L. Sindoni,

PRL, [arXiv:1303.3576 \[gr-qc\]](#);

JHEP, [arXiv:1311.1238 \[gr-qc\]](#)

single-particle GFT condensate:

$$|\sigma\rangle := \exp(\hat{\sigma}) |0\rangle \quad \hat{\sigma} := \int d^4g \sigma(g_I) \hat{\varphi}^\dagger(g_I) \quad \sigma(g_I k) = \sigma(g_I)$$

from truncation of SD equations for GFT model

applied to (coherent) GFT condensate state,
gives equation for “wave function”:

$$\int [dg'_i] \tilde{\mathcal{K}}(g_i, g'_i) \sigma(g'_i) + \lambda \frac{\delta \tilde{\mathcal{V}}}{\delta \varphi(g_i)} \Big|_{\varphi \equiv \sigma} = 0$$

basically (up to some approximations), the “classical GFT eqns”

non-linear and non-local extension of quantum cosmology-like equation for “collective wave function”

QG (GFT) analogue of Gross-Pitaevskii hydrodynamic equation in BECs

similar equations obtained in non-linear extension of LQC (Bojowald et al. '12)

toy examples of effective dynamics have been studied (S. Gielen, DO, L. Sindoni '13; S. Gielen, '14; G. Calcagni '14; S. Gielen, DO, '14), with approximate Friedmann eq emerging in isotropic case

Effective cosmological dynamics from GFT

follow closely procedure used in real BECs

S. Gielen, DO, L. Sindoni,

PRL, [arXiv:1303.3576 \[gr-qc\]](#);

JHEP, [arXiv:1311.1238 \[gr-qc\]](#)

single-particle GFT condensate:

$$|\sigma\rangle := \exp(\hat{\sigma}) |0\rangle \quad \hat{\sigma} := \int d^4g \sigma(g_I) \hat{\varphi}^\dagger(g_I) \quad \sigma(g_I k) = \sigma(g_I)$$

from truncation of SD equations for GFT model

applied to (coherent) GFT condensate state,
gives equation for “wave function”:

$$\int [dg'_i] \tilde{\mathcal{K}}(g_i, g'_i) \sigma(g'_i) + \lambda \frac{\delta \tilde{\mathcal{V}}}{\delta \varphi(g_i)} \Big|_{\varphi \equiv \sigma} = 0$$

basically (up to some approximations), the “classical GFT eqns”

non-linear and non-local extension of quantum cosmology-like equation for “collective wave function”

QG (GFT) analogue of Gross-Pitaevskii hydrodynamic equation in BECs

similar equations obtained in non-linear extension of LQC (Bojowald et al. '12)

toy examples of effective dynamics have been studied (S. Gielen, DO, L. Sindoni '13; S. Gielen, '14; G. Calcagni '14; S. Gielen, DO, '14), with approximate Friedmann eq emerging in isotropic case

effective cosmological dynamics from full-blown 4d gravity models? under way (DO, L. Sindoni, E. Wilson-Ewing, '15)

Cosmological variables and LQC lattice refinement from GFT condensate cosmology

S. Gielen, DO, arXiv:1407.8167 [gr-qc]

key new element in 2nd quantised framework: number operator N

crucial in identifying macroscopic (cosmological) variables + enters effective cosmological dynamics

Cosmological variables and LQC lattice refinement from GFT condensate cosmology

S. Gielen, DO, arXiv:1407.8167 [gr-qc]

key new element in 2nd quantised framework: number operator N

crucial in identifying macroscopic (cosmological) variables + enters effective cosmological dynamics

microscopic (single vertex, 1st quantized) variables:

$$[\hat{g}, \hat{B}^i] = -i\kappa\tau^i \hat{g}, \quad [\hat{B}^i, \hat{B}^j] = -i\kappa \epsilon^{ij}_k \hat{B}^k$$
$$g = \sqrt{1 - \vec{\pi}[g]^2} \mathbf{1} - i\vec{\sigma} \cdot \vec{\pi}[g], \quad |\vec{\pi}[g]| \leq 1$$

Cosmological variables and LQC lattice refinement from GFT condensate cosmology

S. Gielen, DO, arXiv:1407.8167 [gr-qc]

key new element in 2nd quantised framework: number operator N

crucial in identifying macroscopic (cosmological) variables + enters effective cosmological dynamics

microscopic (single vertex, 1st quantized) variables:

$$[\hat{g}, \hat{B}^i] = -i\kappa\tau^i \hat{g}, \quad [\hat{B}^i, \hat{B}^j] = -i\kappa \epsilon^{ij}_k \hat{B}^k$$
$$g = \sqrt{1 - \vec{\pi}[g]^2} \mathbf{1} - i\vec{\sigma} \cdot \vec{\pi}[g], \quad |\vec{\pi}[g]| \leq 1$$

macroscopic (2nd quantized) variables:

$$\hat{b}_a^i = i\kappa \int (dg)^4 \hat{\varphi}^\dagger(g_I) \frac{d}{dt} \hat{\varphi}(\exp(\tau_a^i t) g_I) \Big|_{t=0}$$

total (non-commutative) flux

$$\hat{\Pi}[g_a] = \int (dg)^4 \vec{\pi}[g_a] \hat{\varphi}^\dagger(g_I) \hat{\varphi}(g_I)$$

“total holonomy”

satisfying: $\left[\hat{b}_a^i, \widehat{\vec{\Pi}[g_a]} \right] \propto \hat{N}$

Cosmological variables and LQC lattice refinement from GFT condensate cosmology

S. Gielen, DO, arXiv:1407.8167 [gr-qc]

key new element in 2nd quantised framework: number operator \hat{N}

crucial in identifying macroscopic (cosmological) variables + enters effective cosmological dynamics

microscopic (single vertex, 1st quantized) variables:

$$[\hat{g}, \hat{B}^i] = -i\kappa\tau^i \hat{g}, \quad [\hat{B}^i, \hat{B}^j] = -i\kappa \epsilon^{ij}_k \hat{B}^k$$

$$g = \sqrt{1 - \vec{\pi}[g]^2} \mathbf{1} - i\vec{\sigma} \cdot \vec{\pi}[g], \quad |\vec{\pi}[g]| \leq 1$$

macroscopic (2nd quantized) variables:

$$\hat{b}_a^i = i\kappa \int (dg)^4 \hat{\varphi}^\dagger(g_I) \frac{d}{dt} \hat{\varphi}(\exp(\tau_a^i t) g_I) \Big|_{t=0}$$

total (non-commutative) flux

$$\hat{\Pi}[g_a] = \int (dg)^4 \vec{\pi}[g_a] \hat{\varphi}^\dagger(g_I) \hat{\varphi}(g_I)$$

“total holonomy”

satisfying: $\left[\hat{b}_a^i, \widehat{\vec{\Pi}[g_a]} \right] \propto \hat{N}$

entering effective (semiclassical) cosmological equations via expectation values:

$$\langle \hat{b}_a^i \rangle \quad \langle \widehat{\vec{\Pi}[g_a]} \rangle$$

Cosmological variables and LQC lattice refinement from GFT condensate cosmology

S. Gielen, DO, arXiv:1407.8167 [gr-qc]

key new element in 2nd quantised framework: number operator \hat{N}

crucial in identifying macroscopic (cosmological) variables + enters effective cosmological dynamics

microscopic (single vertex, 1st quantized) variables:

$$[\hat{g}, \hat{B}^i] = -i\kappa\tau^i \hat{g}, \quad [\hat{B}^i, \hat{B}^j] = -i\kappa \epsilon^{ij}_k \hat{B}^k$$

$$g = \sqrt{1 - \vec{\pi}[g]^2} \mathbf{1} - i\vec{\sigma} \cdot \vec{\pi}[g], \quad |\vec{\pi}[g]| \leq 1$$

macroscopic (2nd quantized) variables:

$$\hat{b}_a^i = i\kappa \int (dg)^4 \hat{\varphi}^\dagger(g_I) \frac{d}{dt} \hat{\varphi}(\exp(\tau_a^i t) g_I) \Big|_{t=0}$$

total (non-commutative) flux

$$\hat{\Pi}[g_a] = \int (dg)^4 \vec{\pi}[g_a] \hat{\varphi}^\dagger(g_I) \hat{\varphi}(g_I)$$

“total holonomy”

satisfying: $\left[\hat{b}_a^i, \widehat{\vec{\Pi}[g_a]} \right] \propto \hat{N}$

entering effective (semiclassical) cosmological equations via expectation values: $\langle \hat{b}_a^i \rangle$ $\langle \widehat{\vec{\Pi}[g_a]} \rangle$

macroscopic geometric conjugate variables are instead:

$$(\vec{\Pi}[g_a]^{\text{av.}}) = \langle \hat{\Pi}[g_a] \rangle / \langle \hat{N} \rangle$$

“average holonomy”

one extensive, other intensive

$$B_a^i = \langle \hat{b}_a^i \rangle$$

Cosmological variables and LQC lattice refinement from GFT condensate cosmology

S. Gielen, DO, arXiv:1407.8167 [gr-qc]

Cosmological variables and LQC lattice refinement from GFT condensate cosmology

S. Gielen, DO, arXiv:1407.8167 [gr-qc]

natural definition of “cosmological flux variable”:

$$\hat{f}_I^i = i\kappa \int (dg)^4 \hat{\varphi}^\dagger(\pi[g_J]) \frac{\partial}{\partial \pi_i^I} \hat{\varphi}(\pi[g_J])$$

commutative limit of “total flux” B

Cosmological variables and LQC lattice refinement from GFT condensate cosmology

S. Gielen, DO, arXiv:1407.8167 [gr-qc]

natural definition of “cosmological flux variable”:

$$\hat{f}_I^i = i\kappa \int (dg)^4 \hat{\varphi}^\dagger(\pi[g_J]) \frac{\partial}{\partial \pi_i^I} \hat{\varphi}(\pi[g_J])$$

commutative limit of “total flux” B

natural definition of “cosmological connection”:

$$\mu \vec{\omega} := - \frac{\langle \vec{\Pi} \rangle}{|\langle \vec{\Pi} \rangle|} \arcsin \frac{|\langle \vec{\Pi} \rangle|}{N}$$

Cosmological variables and LQC lattice refinement from GFT condensate cosmology

S. Gielen, DO, arXiv:1407.8167 [gr-qc]

natural definition of “cosmological flux variable”:

$$\hat{f}_I^i = i\kappa \int (dg)^4 \hat{\varphi}^\dagger(\pi[g_J]) \frac{\partial}{\partial \pi_i^I} \hat{\varphi}(\pi[g_J])$$

commutative limit of “total flux” B

natural definition of “cosmological connection”:

$$\mu \vec{\omega} := -\frac{\langle \vec{\Pi} \rangle}{|\langle \vec{\Pi} \rangle|} \arcsin \frac{|\langle \vec{\Pi} \rangle|}{N}$$

μ encodes choice of coordinate system in defining the cosmological connection
can be chosen as in “improved” LQC: $\mu = N^{-1/3}$

Cosmological variables and LQC lattice refinement from GFT condensate cosmology

S. Gielen, DO, arXiv:1407.8167 [gr-qc]

natural definition of “cosmological flux variable”: $\hat{f}_I^i = i\kappa \int (dg)^4 \hat{\varphi}^\dagger(\pi[g_J]) \frac{\partial}{\partial \pi_i^I} \hat{\varphi}(\pi[g_J])$

commutative limit of “total flux” B

natural definition of “cosmological connection”: $\mu \vec{\omega} := -\frac{\langle \vec{\Pi} \rangle}{|\langle \vec{\Pi} \rangle|} \arcsin \frac{|\langle \vec{\Pi} \rangle|}{N}$

μ encodes choice of coordinate system in defining the cosmological connection
can be chosen as in “improved” LQC: $\mu = N^{-1/3}$

intrinsic dependence of cosmological holonomies on N = average number of microscopic building blocks

GFT condensate counterpart of the “lattice refinement” in LQC

Cosmological variables and LQC lattice refinement from GFT condensate cosmology

S. Gielen, DO, arXiv:1407.8167 [gr-qc]

natural definition of “cosmological flux variable”:

$$\hat{f}_I^i = i\kappa \int (dg)^4 \hat{\varphi}^\dagger(\pi[g_J]) \frac{\partial}{\partial \pi_i^I} \hat{\varphi}(\pi[g_J])$$

commutative limit of “total flux” B

natural definition of “cosmological connection”:

$$\mu \vec{\omega} := -\frac{\langle \vec{\Pi} \rangle}{|\langle \vec{\Pi} \rangle|} \arcsin \frac{|\langle \vec{\Pi} \rangle|}{N}$$

μ encodes choice of coordinate system in defining the cosmological connection
can be chosen as in “improved” LQC: $\mu = N^{-1/3}$

intrinsic dependence of cosmological holonomies on N = average number of microscopic building blocks

GFT condensate counterpart of the “lattice refinement” in LQC

two immediate (generic) consequences:

- 1) GFT condensate cosmology gives quantum corrections to cosmological equations akin to LQC ones
- 2) effective cosmological equations will carry a dependence on $\langle N \rangle$ (purely quantum observable) when expressed in terms of cosmological variables

exact relation between $\langle N \rangle$ and cosmological variables depends on quantum state

3rd message(s):

GFT condensates are interesting candidates for physical, geometric vacua of QG theory

3rd message(s):

GFT condensates are interesting candidates for physical, geometric vacua of QG theory

derivation of (quantum) cosmological equations from GFT quantum dynamics **very general**
it rests on:

- continuum homogeneous quantum space (at microscopic scales) ~ GFT condensate
- good encoding of discrete geometry in GFT states
- 2nd quantized GFT formalism

3rd message(s):

GFT condensates are interesting candidates for physical, geometric vacua of QG theory

derivation of (quantum) cosmological equations from GFT quantum dynamics **very general**
it rests on:

- continuum homogeneous quantum space (at microscopic scales) ~ GFT condensate
- good encoding of discrete geometry in GFT states
- 2nd quantized GFT formalism

non-linear quantum cosmology-like equations emerging as hydrodynamics for GFT condensate

3rd message(s):

GFT condensates are interesting candidates for physical, geometric vacua of QG theory

derivation of (quantum) cosmological equations from GFT quantum dynamics **very general**
it rests on:

- continuum homogeneous quantum space (at microscopic scales) ~ GFT condensate
- good encoding of discrete geometry in GFT states
- 2nd quantized GFT formalism

non-linear quantum cosmology-like equations emerging as hydrodynamics for GFT condensate

derivation of (quantum) cosmology from fundamental QG formalism!

3rd message(s):

GFT condensates are interesting candidates for physical, geometric vacua of QG theory

derivation of (quantum) cosmological equations from GFT quantum dynamics **very general**
it rests on:

- continuum homogeneous quantum space (at microscopic scales) ~ GFT condensate
- good encoding of discrete geometry in GFT states
- 2nd quantized GFT formalism

non-linear quantum cosmology-like equations emerging as hydrodynamics for GFT condensate

derivation of (quantum) cosmology from fundamental QG formalism!

exact form of equations depends on specific model considered
now: derive effective cosmological dynamics from most promising GFT (spin foam) models

3rd message(s):

GFT condensates are interesting candidates for physical, geometric vacua of QG theory

derivation of (quantum) cosmological equations from GFT quantum dynamics **very general**
it rests on:

- continuum homogeneous quantum space (at microscopic scales) ~ GFT condensate
- good encoding of discrete geometry in GFT states
- 2nd quantized GFT formalism

non-linear quantum cosmology-like equations emerging as hydrodynamics for GFT condensate

derivation of (quantum) cosmology from fundamental QG formalism!

exact form of equations depends on specific model considered
now: derive effective cosmological dynamics from most promising GFT (spin foam) models

non-linear quantum cosmology is QG analogue of Gross-Pitaevskii hydrodynamics for BECs

consistent with “geometrogenesis” hypothesis and general “macro-from-micro” scenario

GFT condensate cosmology: phenomenology?

“standard” directions
(but calculations to be done):

phenomenology as in LQC, here derived from fundamental theory

basis for most LQC phenomenology: modified Friedmann equation
plus quantum corrections

$$\left(\frac{\dot{a}}{a}\right)^2 = (8\pi G \rho/3) \left(1 - \frac{\rho}{\rho_{\text{crit}}}\right)$$

GFT condensate cosmology: phenomenology?

“standard” directions
(but calculations to be done):

phenomenology as in LQC, here derived from fundamental theory

basis for most LQC phenomenology: modified Friedmann equation
plus quantum corrections

$$\left(\frac{\dot{a}}{a}\right)^2 = (8\pi G \rho/3) \left(1 - \frac{\rho}{\rho_{\text{crit}}}\right)$$

- in GFT condensate cosmology, modified gravity equations at effective level basically inevitable

GFT condensate cosmology: phenomenology?

“standard” directions
(but calculations to be done):

phenomenology as in LQC, here derived from fundamental theory

basis for most LQC phenomenology: modified Friedmann equation
plus quantum corrections

$$\left(\frac{\dot{a}}{a}\right)^2 = (8\pi G \rho/3) \left(1 - \frac{\rho}{\rho_{\text{crit}}}\right)$$

- in GFT condensate cosmology, modified gravity equations at effective level basically inevitable

modifications from several ingredients (in addition to involved microscopic dynamics):

GFT condensate cosmology: phenomenology?

“standard” directions
(but calculations to be done):

phenomenology as in LQC, here derived from fundamental theory

basis for most LQC phenomenology: modified Friedmann equation
plus quantum corrections

$$\left(\frac{\dot{a}}{a}\right)^2 = (8\pi G \rho/3) \left(1 - \frac{\rho}{\rho_{\text{crit}}}\right)$$

- in GFT condensate cosmology, modified gravity equations at effective level basically inevitable

modifications from several ingredients (in addition to involved microscopic dynamics):

“expected from LQC”:

- holonomy corrections
- inverse triad corrections
- non-commutativity of fluxes

GFT condensate cosmology: phenomenology?

“standard” directions
(but calculations to be done):

phenomenology as in LQC, here derived from fundamental theory

basis for most LQC phenomenology: modified Friedmann equation
plus quantum corrections

$$\left(\frac{\dot{a}}{a}\right)^2 = (8\pi G \rho/3) \left(1 - \frac{\rho}{\rho_{\text{crit}}}\right)$$

- in GFT condensate cosmology, modified gravity equations at effective level basically inevitable

modifications from several ingredients (in addition to involved microscopic dynamics):

“expected from LQC”:

- holonomy corrections
- inverse triad corrections
- non-commutativity of fluxes

new, due to embedding into full theory:

- new QG observable N: number of “QG atoms of space”
- non-linear terms in effective cosmological equations
- hydrodynamic character of cosmological dynamics

GFT condensate cosmology: phenomenology?

“standard” directions
(but calculations to be done):

phenomenology as in LQC, here derived from fundamental theory

basis for most LQC phenomenology: modified Friedmann equation
plus quantum corrections

$$\left(\frac{\dot{a}}{a}\right)^2 = (8\pi G \rho/3) \left(1 - \frac{\rho}{\rho_{\text{crit}}}\right)$$

- in GFT condensate cosmology, modified gravity equations at effective level basically inevitable

modifications from several ingredients (in addition to involved microscopic dynamics):

“expected from LQC”:

- holonomy corrections
- inverse triad corrections
- non-commutativity of fluxes

new, due to embedding into full theory:

- new QG observable N: number of “QG atoms of space”
- non-linear terms in effective cosmological equations
- hydrodynamic character of cosmological dynamics

other source of LQC phenomenology (Bojowald et al.): deformation of diffeomorphism algebra & signature change

GFT condensate cosmology: phenomenology?

“standard” directions
(but calculations to be done):

phenomenology as in LQC, here derived from fundamental theory

basis for most LQC phenomenology: modified Friedmann equation plus quantum corrections

$$\left(\frac{\dot{a}}{a}\right)^2 = (8\pi G \rho/3) \left(1 - \frac{\rho}{\rho_{\text{crit}}}\right)$$

- in GFT condensate cosmology, modified gravity equations at effective level basically inevitable

modifications from several ingredients (in addition to involved microscopic dynamics):

“expected from LQC”:

- holonomy corrections
- inverse triad corrections
- non-commutativity of fluxes

new, due to embedding into full theory:

- new QG observable N: number of “QG atoms of space”
- non-linear terms in effective cosmological equations
- hydrodynamic character of cosmological dynamics

other source of LQC phenomenology (Bojowald et al.): deformation of diffeomorphism algebra & signature change

- diffeos in GFT also expected to be deformed: A. Baratin, F. Girelli, DO, ‘11
 - simplicial diffeos realised as global quantum group symmetry in topological models
 - expect more surprises at effective cosmological level

GFT condensate cosmology: phenomenology?

Big Bounce?

GFT condensate cosmology: phenomenology?

Big Bounce?

given effective cosmological equations for GFT condensates,
it can be derived via the same type of calculations done in LQC

GFT condensate cosmology: phenomenology?

Big Bounce?

given effective cosmological equations for GFT condensates,
it can be derived via the same type of calculations done in LQC (Big Bounce from the full theory!)

GFT condensate cosmology: phenomenology?

Big Bounce?

given effective cosmological equations for GFT condensates,
it can be derived via the same type of calculations done in LQC (Big Bounce from the full theory!)

.... provided the GFT hydrodynamics approximation (and other assumptions) does not break down in that regime

GFT condensate cosmology: phenomenology?

Big Bounce?

given effective cosmological equations for GFT condensates,
it can be derived via the same type of calculations done in LQC (Big Bounce from the full theory!)

.... provided the GFT hydrodynamics approximation (and other assumptions) does not break down in that regime

if it does break, one has to go back to the full GFT theory, and improve the
construction (ansatz for vacuum, approximation of SD equations,)
and then try again

GFT condensate cosmology: phenomenology?

Big Bounce?

given effective cosmological equations for GFT condensates,
it can be derived via the same type of calculations done in LQC (Big Bounce from the full theory!)

.... provided the GFT hydrodynamics approximation (and other assumptions) does not break down in that regime

if it does break, one has to go back to the full GFT theory, and improve the
construction (ansatz for vacuum, approximation of SD equations,)
and then try again

novelty: it can be done!

GFT condensate cosmology: phenomenology?

Big Bounce?

given effective cosmological equations for GFT condensates,
it can be derived via the same type of calculations done in LQC (Big Bounce from the full theory!)

.... provided the GFT hydrodynamics approximation (and other assumptions) does not break down in that regime

if it does break, one has to go back to the full GFT theory, and improve the
construction (ansatz for vacuum, approximation of SD equations,)
and then try again

novelty: it can be done!

exactly as one would do in a BEC....

GFT condensate cosmology: phenomenology?

“standard” directions

(but conceptual and technical issues to be solved, first)

effective dynamics of cosmological perturbations
from first principles, i.e. from full QG formalism

needed for computation of CMB spectrum
needed for tests of fate of Lorentz invariance

GFT condensate cosmology: phenomenology?

“standard” directions

(but conceptual and technical issues to be solved, first)

effective dynamics of cosmological perturbations
from first principles, i.e. from full QG formalism

needed for computation of CMB spectrum
needed for tests of fate of Lorentz invariance

several strategies:

GFT condensate cosmology: phenomenology?

“standard” directions
(but conceptual and technical issues to be solved, first)

effective dynamics of cosmological perturbations
from first principles, i.e. from full QG formalism

needed for computation of CMB spectrum
needed for tests of fate of Lorentz invariance

several strategies:

“cheap” (similar to Agullo, Ashtekar, Nelson):

1. define modified FRW metric from expectation values for cosmological variables derived from GFT
2. use it inside standard effective QFT for fields

GFT condensate cosmology: phenomenology?

“standard” directions
(but conceptual and technical issues to be solved, first)

effective dynamics of cosmological perturbations
from first principles, i.e. from full QG formalism

needed for computation of CMB spectrum
needed for tests of fate of Lorentz invariance

several strategies:

“cheap” (similar to Agullo, Ashtekar, Nelson):

1. define modified FRW metric from expectation values for cosmological variables derived from GFT
2. use it inside standard effective QFT for fields

“ambitious”:

1. develop statistical aspects of GFT condensate hydrodynamics, in terms of “homogeneous patches” (S. Gielen, '15)
2. derive **effective dynamics for GFT fluctuations above condensate from full theory**
3. recast it in standard spacetime-based QFT form using information from background GFT condensate
(difficulty is: the formalism naturally gives it in diffeo-invariant variables, spacetime-free form)

GFT condensate cosmology: phenomenology?

“standard” directions
(but conceptual and technical issues to be solved, first)

effective dynamics of cosmological perturbations
from first principles, i.e. from full QG formalism

needed for computation of CMB spectrum
needed for tests of fate of Lorentz invariance

several strategies:

“cheap” (similar to Agullo, Ashtekar, Nelson):

1. define modified FRW metric from expectation values for cosmological variables derived from GFT
2. use it inside standard effective QFT for fields

“ambitious”:

1. develop statistical aspects of GFT condensate hydrodynamics, in terms of “homogeneous patches” (S. Gielen, '15)
2. derive **effective dynamics for GFT fluctuations above condensate from full theory**
3. recast it in standard spacetime-based QFT form using information from background GFT condensate
(difficulty is: the formalism naturally gives it in diffeo-invariant variables, spacetime-free form)

expect deformation of standard QFT:

- holonomization of the connection and non-commutativity of triad variables, both entering definition of basic variables for perturbations (momenta, positions)
- derivation of effective dynamics of perturbations around mean field in topological GFT:
non-commutative scalar field theory on non-commutative flat space

GFT condensate cosmology: phenomenology?

“non-standard” directions

GFT condensate cosmology: phenomenology?

“non-standard” directions

- “superfluidity of the metric”?

GFT condensate cosmology: phenomenology?

“non-standard” directions

- “superfluidity of the metric”?
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field

GFT condensate cosmology: phenomenology?

“non-standard” directions

- “superfluidity of the metric”?
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?

GFT condensate cosmology: phenomenology?

“non-standard” directions

- “superfluidity of the metric”?
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?

GFT condensate cosmology: phenomenology?

“non-standard” directions

- “superfluidity of the metric”?
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?
- “cosmological dissipation”?

GFT condensate cosmology: phenomenology?

“non-standard” directions

- **“superfluidity of the metric”?**
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?
- **“cosmological dissipation”?**
effective cosmological counterpart of interaction between QG atoms forming GFT condensate and GFT fluctuations over it

GFT condensate cosmology: phenomenology?

“non-standard” directions

- **“superfluidity of the metric”?**
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?
- **“cosmological dissipation”?**
effective cosmological counterpart of interaction between QG atoms forming GFT condensate and GFT fluctuations over it
S. Liberati, L. Maccione, ‘13

GFT condensate cosmology: phenomenology?

“non-standard” directions

- **“superfluidity of the metric”?**
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?
- **“cosmological dissipation”?**
effective cosmological counterpart of interaction between QG atoms forming GFT condensate and GFT fluctuations over it
S. Liberati, L. Maccione, ‘13

GFT condensate cosmology: phenomenology?

“non-standard” directions

- **“superfluidity of the metric”?**
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?
- **“cosmological dissipation”?**
effective cosmological counterpart of interaction between QG atoms forming GFT condensate and GFT fluctuations over it
S. Liberati, L. Maccione, ‘13
- **cosmological signature of “depletion factor”?**

GFT condensate cosmology: phenomenology?

“non-standard” directions

- **“superfluidity of the metric”?**
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?
- **“cosmological dissipation”?**
effective cosmological counterpart of interaction between QG atoms forming GFT condensate and GFT fluctuations over it
S. Liberati, L. Maccione, ‘13
- **cosmological signature of “depletion factor”?**
assumption of all QG atoms being in condensate state is unrealistic: new observable: depletion factor

GFT condensate cosmology: phenomenology?

“non-standard” directions

- **“superfluidity of the metric”?**
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?
- **“cosmological dissipation”?**
effective cosmological counterpart of interaction between QG atoms forming GFT condensate and GFT fluctuations over it
S. Liberati, L. Maccione, ‘13
- **cosmological signature of “depletion factor”?**
assumption of all QG atoms being in condensate state is unrealistic: new observable: depletion factor
what is its effective cosmological significance?

GFT condensate cosmology: phenomenology?

“non-standard” directions

- **“superfluidity of the metric”?**
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?
- **“cosmological dissipation”?**
effective cosmological counterpart of interaction between QG atoms forming GFT condensate and GFT fluctuations over it
S. Liberati, L. Maccione, ‘13
- **cosmological signature of “depletion factor”?**
assumption of all QG atoms being in condensate state is unrealistic: new observable: depletion factor
what is its effective cosmological significance?
F. Girelli, S. Liberati, L. Sindoni, ‘09

GFT condensate cosmology: phenomenology?

“non-standard” directions

- **“superfluidity of the metric”?**
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?
- **“cosmological dissipation”?**
effective cosmological counterpart of interaction between QG atoms forming GFT condensate and GFT fluctuations over it
S. Liberati, L. Maccione, ‘13
- **cosmological signature of “depletion factor”?**
assumption of all QG atoms being in condensate state is unrealistic: new observable: depletion factor
what is its effective cosmological significance?
F. Girelli, S. Liberati, L. Sindoni, ‘09

GFT condensate cosmology: phenomenology?

“non-standard” directions

- **“superfluidity of the metric”?**
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?
- **“cosmological dissipation”?**
effective cosmological counterpart of interaction between QG atoms forming GFT condensate and GFT fluctuations over it
S. Liberati, L. Maccione, ‘13
- **cosmological signature of “depletion factor”?**
assumption of all QG atoms being in condensate state is unrealistic: new observable: depletion factor
what is its effective cosmological significance?
F. Girelli, S. Liberati, L. Sindoni, ‘09
- **direct cosmological signatures of geometrogenesis phase transition?**

GFT condensate cosmology: phenomenology?

“non-standard” directions

- **“superfluidity of the metric”?**
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?
- **“cosmological dissipation”?**
effective cosmological counterpart of interaction between QG atoms forming GFT condensate and GFT fluctuations over it
S. Liberati, L. Maccione, ‘13
- **cosmological signature of “depletion factor”?**
assumption of all QG atoms being in condensate state is unrealistic: new observable: depletion factor
what is its effective cosmological significance?
F. Girelli, S. Liberati, L. Sindoni, ‘09
- **direct cosmological signatures of geometrogenesis phase transition?**
C. Contaldi, J. Magueijo, L. Smolin, ‘06

GFT condensate cosmology: phenomenology?

“non-standard” directions

- **“superfluidity of the metric”?**
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?
- **“cosmological dissipation”?**
effective cosmological counterpart of interaction between QG atoms forming GFT condensate and GFT fluctuations over it
S. Liberati, L. Maccione, ‘13
- **cosmological signature of “depletion factor”?**
assumption of all QG atoms being in condensate state is unrealistic: new observable: depletion factor
what is its effective cosmological significance?
F. Girelli, S. Liberati, L. Sindoni, ‘09
- **direct cosmological signatures of geometrogenesis phase transition?**
C. Contaldi, J. Magueijo, L. Smolin, ‘06

GFT condensate cosmology: phenomenology?

“non-standard” directions

- **“superfluidity of the metric”?**
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?
- **“cosmological dissipation”?**
effective cosmological counterpart of interaction between QG atoms forming GFT condensate and GFT fluctuations over it
S. Liberati, L. Maccione, ‘13
- **cosmological signature of “depletion factor”?**
assumption of all QG atoms being in condensate state is unrealistic: new observable: depletion factor
what is its effective cosmological significance?
F. Girelli, S. Liberati, L. Sindoni, ‘09
- **direct cosmological signatures of geometrogenesis phase transition?**
C. Contaldi, J. Magueijo, L. Smolin, ‘06

a new avenue for analogue gravity: analogue GFT condensate quantum cosmology in real BECs?

GFT condensate cosmology: phenomenology?

“non-standard” directions

- **“superfluidity of the metric”?**
what is the “velocity” of the GFT condensate? indications: the homogeneous metric/triad field
a GFT condensate is not an ordinary fluid, but possibly a superfluid - what emergent property of effective metric corresponds to superfluidity?
 - **“cosmological dissipation”?**
effective cosmological counterpart of interaction between QG atoms forming GFT condensate and GFT fluctuations over it
S. Liberati, L. Maccione, ‘13
 - **cosmological signature of “depletion factor”?**
assumption of all QG atoms being in condensate state is unrealistic: new observable: depletion factor
what is its effective cosmological significance?
F. Girelli, S. Liberati, L. Sindoni, ‘09
 - **direct cosmological signatures of geometrogenesis phase transition?**
C. Contaldi, J. Magueijo, L. Smolin, ‘06
- a new avenue for analogue gravity: analogue GFT condensate quantum cosmology in real BECs?**
- need to simulate mini-superspace, not spacetime
BEC hydrodynamics needs to reproduce GFT “classical” equations”, not GR equations
no problem with lack of diffeo invariance or relativistic symmetry in the lab

Foundational issues: the universe as a (quantum) fluid

- GFT condensates encoded in “collective wave function” identical to Quantum Cosmology one
- corresponding dynamical equation is non-linear and non-local (on mini-superspace)
- **cosmological dynamics is hydrodynamics of fundamental GFT** (analogue of GP equation for BEC)
- still, fundamental dynamics is (more or less) standard quantum mechanics for QG d.o.f.s (although big interpretational issues (e.g. concerning probabilities, unitarity, etc) remain)

calls for new interpretation of “quantum cosmology” (see also Bojowald, '15):

- quantum cosmology is not quantum at all; rather, “cosmological hydrodynamics”
no probability for “the whole universe”, no “Hilbert space of states of the Universe”
- can still use expectation values (average quantities) but in “hydrodynamics” (realistic/statistical) sense
- no problem of “collapse of cosmological wave function” or spontaneous collapse due to non-linearities?

related work by Pearle, Sudarsky, Perez, Peter, Martin, ...

4th message

a new promising direction to extract effective cosmological dynamics
(and associated phenomenology)

directly from full QG theory!

Thank you for your attention!