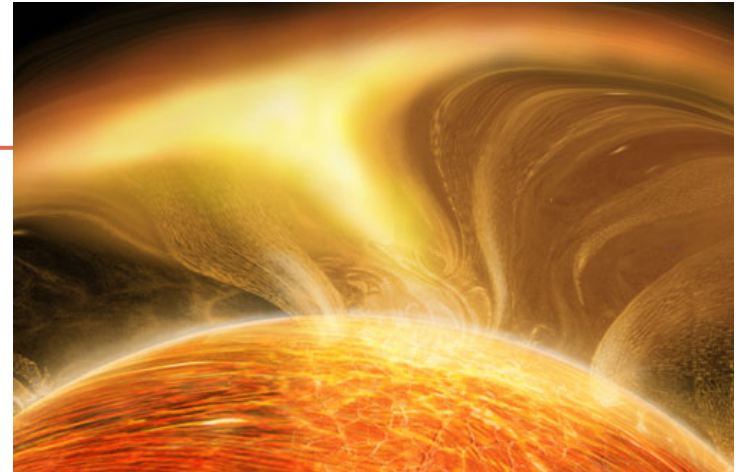


# MAGNETARS: NEUTRON STARS IN THE EXTREME

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with lot of help from P. Esposito, G.L. Israel, S. Mereghetti, L. Nobili, J. Pons, N. Rea, R. Taverna, A. Tiengo, S. Zane and many others



# Outline

- Neutron stars in a nutshell
  - The Galactic NS population
  - NS structure
  - Rotation and magnetism
  - Different classes of (isolated) NSs
- Soft Gamma Repeaters (SGRs) and Anomalous X-ray pulsars (AXPs)
  - General properties
  - SGRs & AXPs as magnetars (aka ultra-magnetized NSs)
- The persistent X-ray emission
  - X-ray spectra
  - The twisted magnetosphere model



# Outline - II

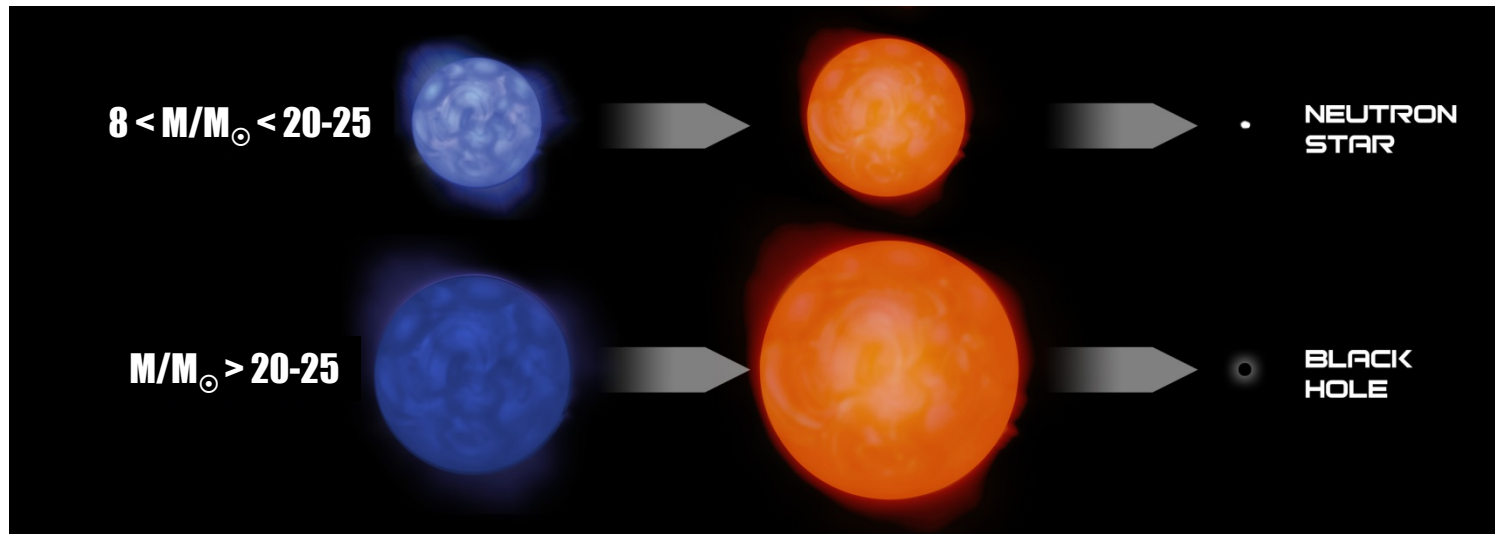
- Bursts & Flares
  - Burst phenomenology
  - Burst triggers and emission
- What makes a magnetar ?
  - SGRs/AXPs vs. high-field pulsars
  - “Low-field” magnetars (?)
  - NSs magneto-thermal evolution
  - The case of SGR 0418+5729
  - A magnetar orbiting the CG SMBH



# Neutron Star Basics

- Neutron stars (and black holes) are born in the core collapse following a supernova explosion
- Present rate of SN events in the Galaxy:  $\approx 0.01/\text{yr}$  (possibly higher in the past)
- Galactic population of compact objects:  $\approx 10^8 - 10^9$  ( $\approx 1\%$  of stars)
- Nature of compact remnant depends on progenitor mass



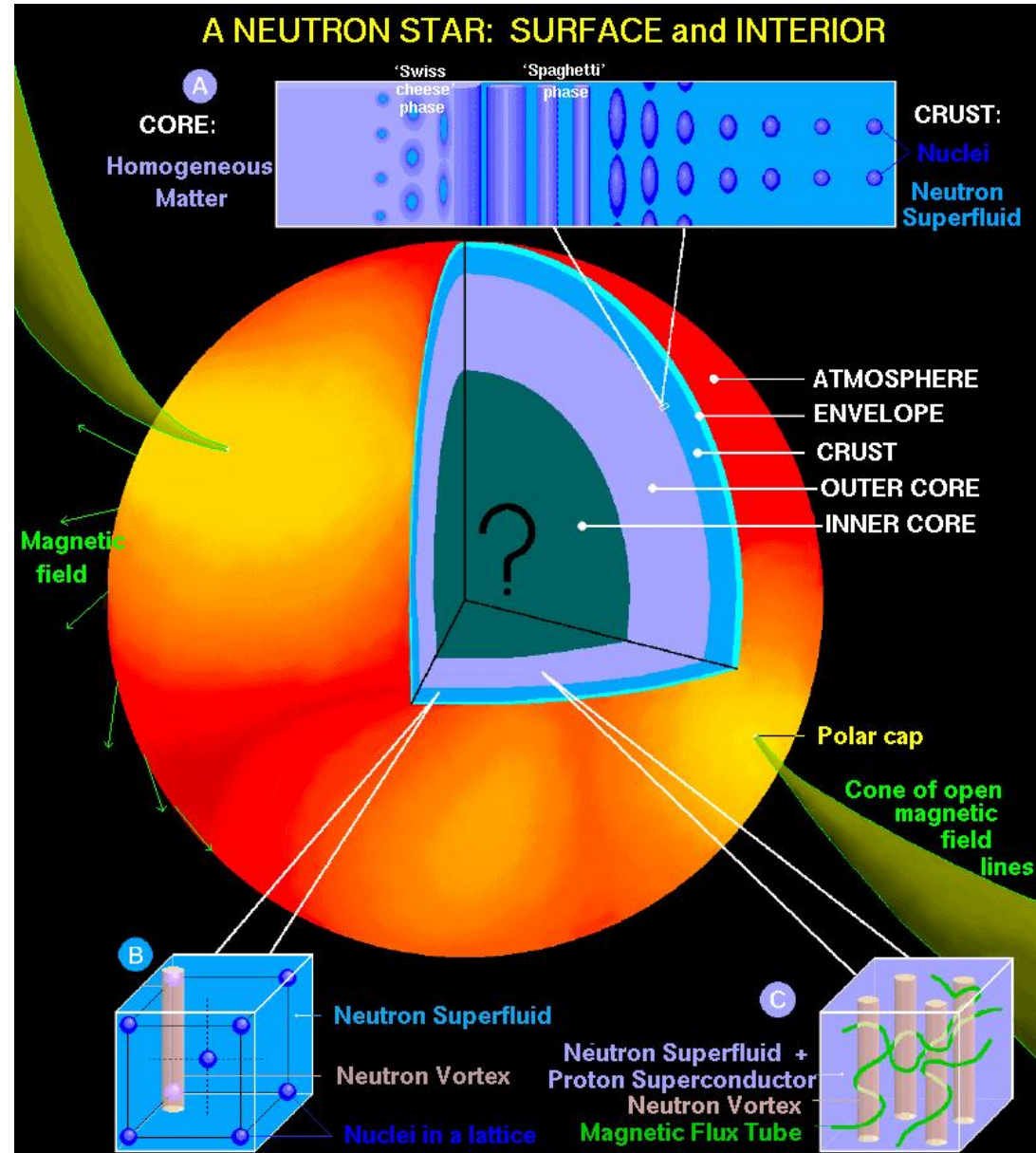


- If  $N(>M) \sim M^{-1.3}$  (Salpeter) only  $\sim 20\%$  of stars with  $M > 8 M_{\odot}$  are more massive than  $25 M_{\odot}$
- Very massive stars may form “magnetars” (Muno et al 2006, but see also Clark et al. 2014), black holes about 10% of the total
- Galactic compact objects are mostly neutron stars ( $\approx 10^8$ )



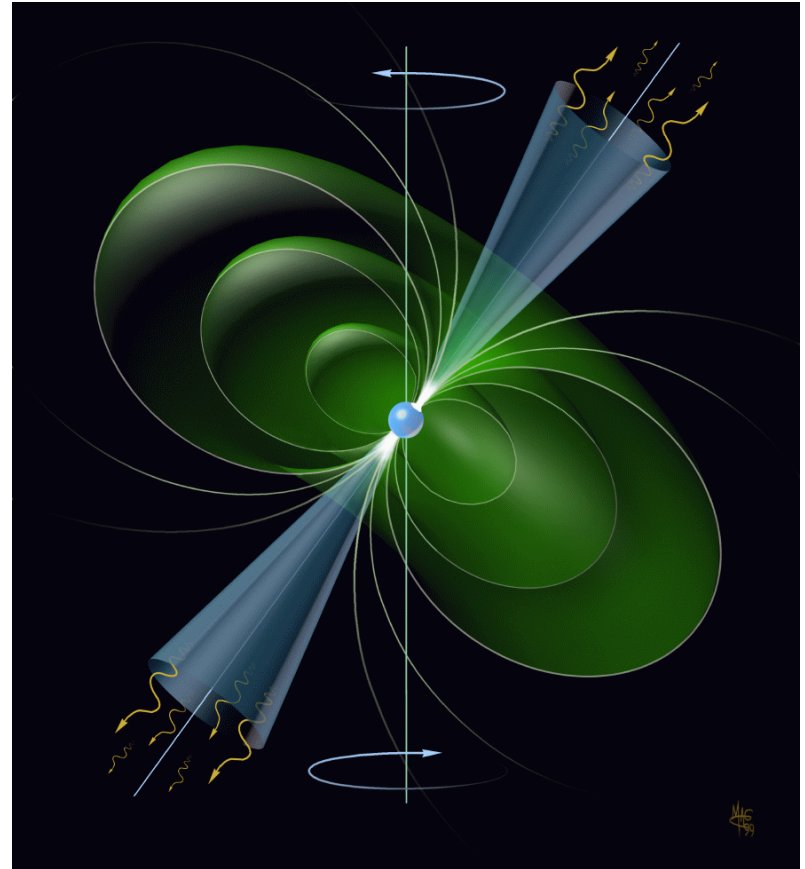
Masses  $\sim 1 - 2 M_{\odot}$   
 Radii  $\sim 10 - 15$  km ( $\sim 3$   
 Schwarzschild radii)  
 Central density  $\sim 10^{15}$  g/cm<sup>3</sup>  
 Surface gravity  $\sim 10^{14-15}$  cm/  
 s<sup>2</sup>

A 10 km sphere made of  
 neutron superfluid covered  
 by few hundred meters of  
 ordinary matter with a 10  
 cm gaseous atmosphere on  
 top (with a big question  
 mark near the centre)



# Rotation and Magnetism

- NSs rotate fast (fastest known: PSR J1748-2446ad,  $\nu=716$  Hz,  $P=1.4$  ms)
- Large dipole magnetic field ( $B \approx 10^{12}$  G in “normal” PSRs)
- Spin-down by magneto-dipolar losses



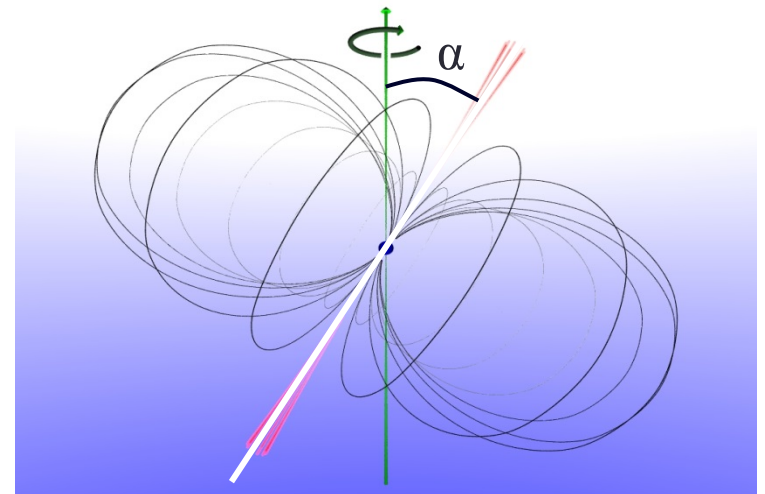
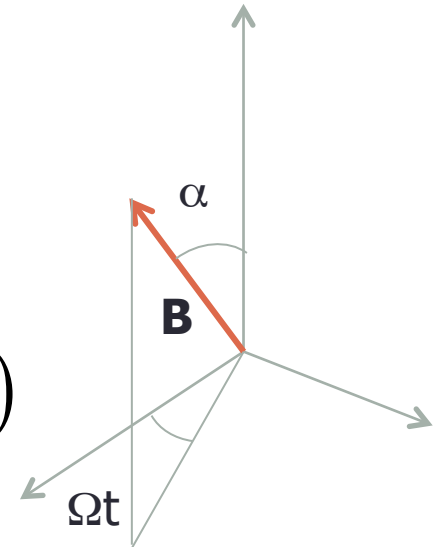
- A rotating magnetic dipole radiates at a rate

$$\dot{E} = -\frac{2}{3c^3} |\ddot{\vec{m}}|^2$$

$$\vec{m} = \frac{1}{2} B_p r^3 (\sin \alpha \cos \phi, \sin \alpha \sin \phi, \cos \alpha)$$

$$\phi = \Omega t$$

$$|\ddot{\vec{m}}| = \Omega^2 \sin \alpha |\vec{m}|$$





Total NS rotational energy  $E = \frac{1}{2} I \Omega^2$

$$\dot{E} = I \Omega \dot{\Omega} = - \frac{B_p^2 R^6 \Omega^4 \sin^2 \alpha}{6c^3}$$

$$\frac{1}{2} \left( \frac{1}{\Omega^2} - \frac{1}{\Omega_0^2} \right) = \frac{t - t_0}{\tau \Omega^2}, \quad \tau = \frac{6Ic^3}{B_p^2 R^6 \Omega^2 \sin^2 \alpha}$$

$$\Omega \ll \Omega_0 \Rightarrow t - t_0 = \tau / 2 \quad \frac{P}{\dot{P}} = \tau$$

$$B_p \sin \alpha = 6.4 \times 10^{19} R_6^{-3} I_{45}^{1/2} (P\dot{P})^{1/2} \text{ G}$$



# Measuring $P$ and $\dot{P}$ is crucial

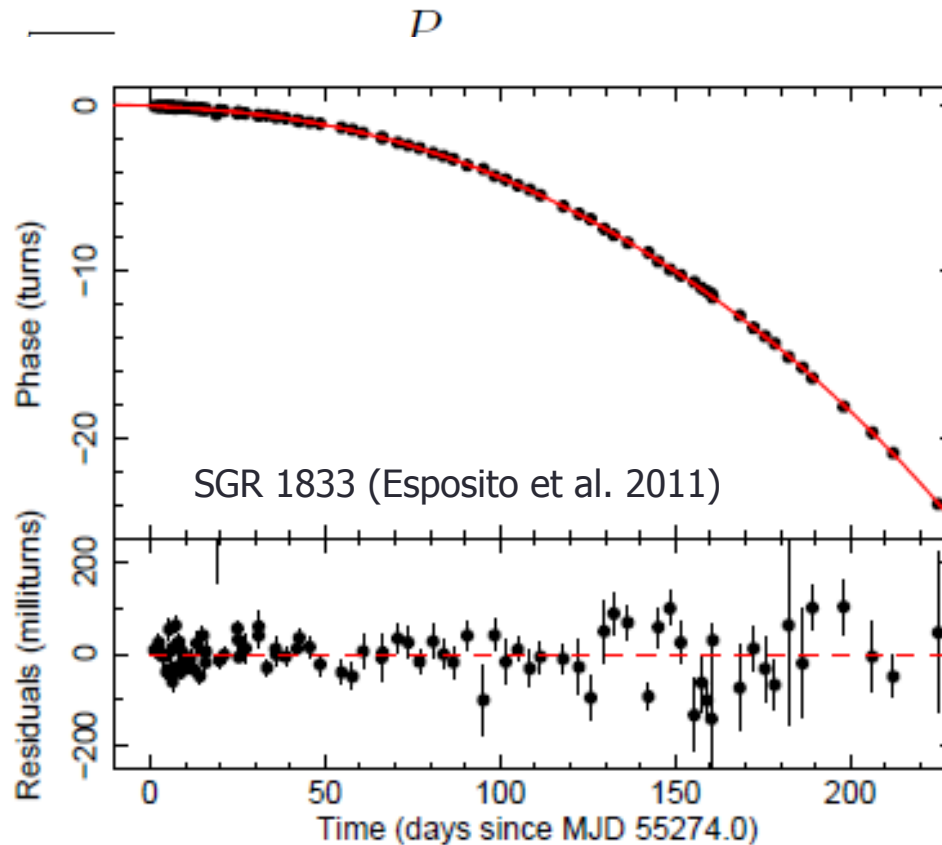
$$B \sim 3 \times 10^{12} \text{ G}$$

$$\dot{E} \sim 1.3 \times 10^{38} \text{ erg s}^{-1}$$

Phase- $\alpha$

$$\phi(t) = \int_t \dot{\phi} dt$$

$$= \phi$$




$P_0 \ll P!$

periodic

$\alpha^2 + \dots$

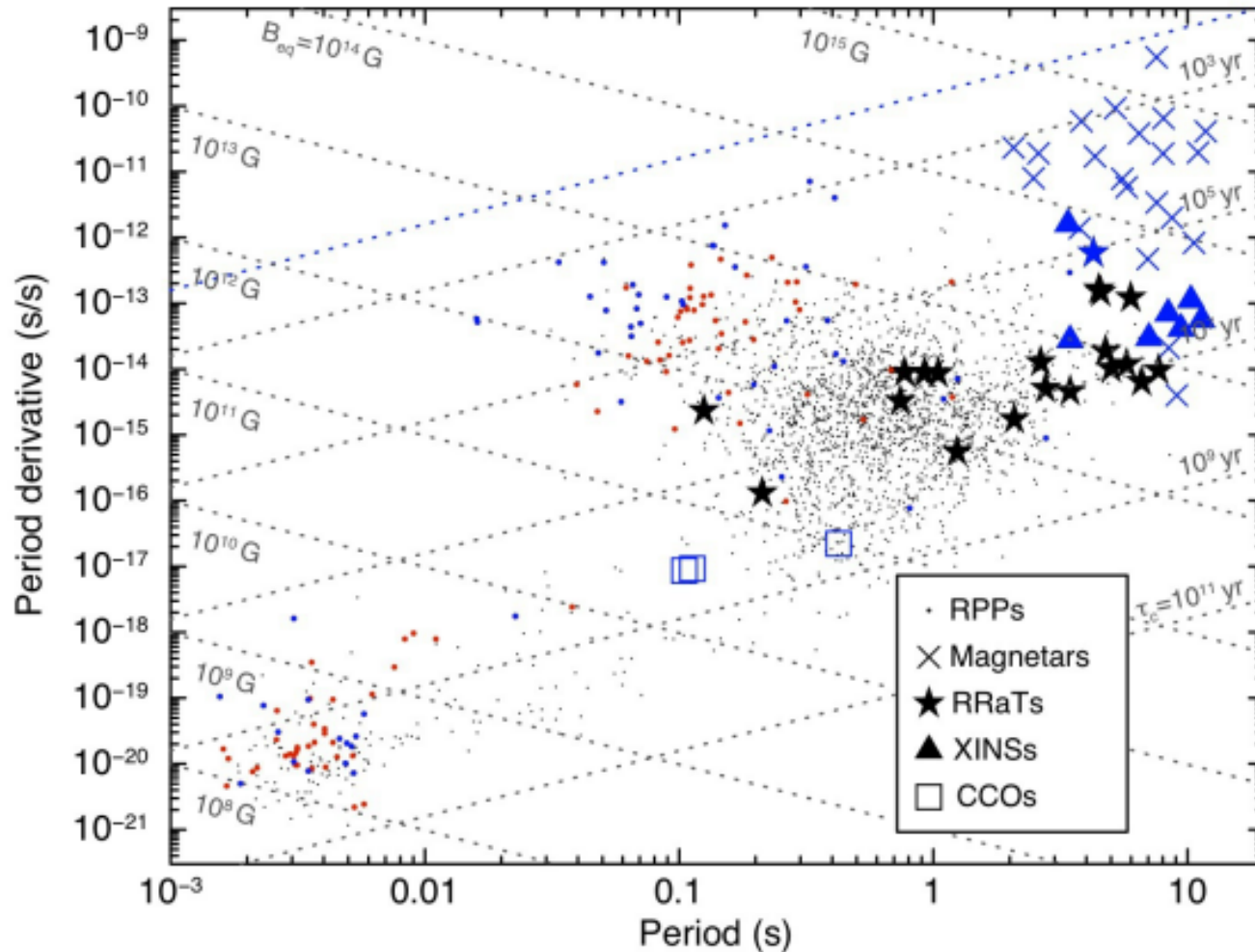


# Pulsars and...

- Most neutron stars are known through their pulsed radio-emission
  - Galactic pulsar population  $\approx 10^5$  (Vranisevic et al. 2004; > 2000 detected, ATNF catalogue)
  - The majority of neutron stars are old, dead objects
  - Observations in the X- and  $\gamma$ -rays revealed the existence of different classes of isolated neutron stars
    - Central compact objects in SNRs (CCOs)
    - X-ray dim isolated neutron stars (XDINSs)
    - Rotating Radio Transients (RRaTs)
    - Soft  $\gamma$ -repeaters (SGRs)
    - Anomalous X-ray pulsars (AXPs)
-  **Magnetars**

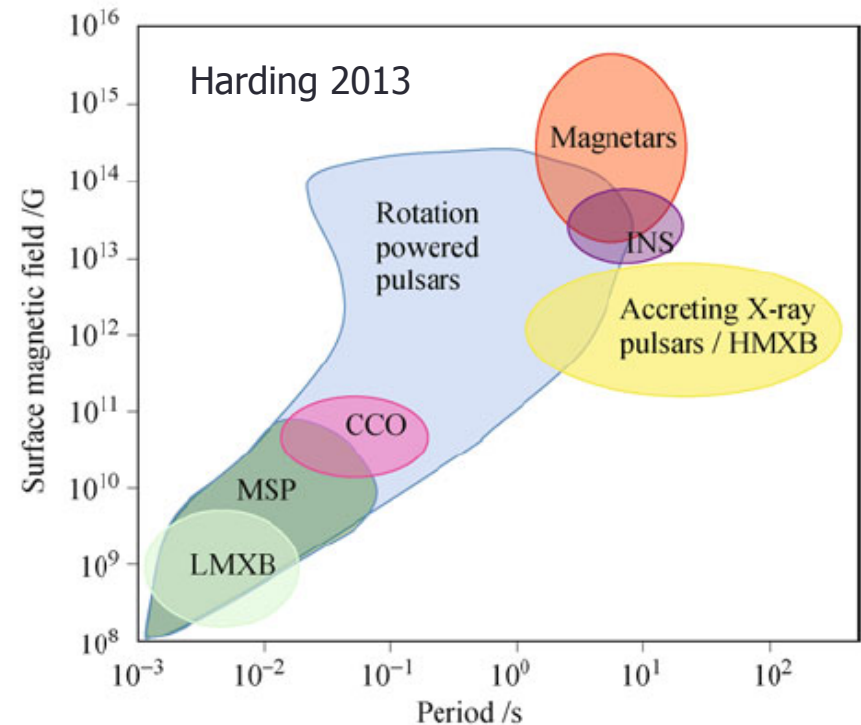


# The P - $\dot{P}$ Diagram



# The Powerhouse

- Rotation ( $\dot{E} > L_{\text{bol}}$ , RPPs)
  - PSRs, RRaTs
- Residual heat
  - XDINSs, CCOs
- Magnetic energy
  - SGRs, AXP
- Accretion
  - HMXBs, LMXBs



# Are There Enough SNe ?

- Core-collapse supernova rate,  $\beta_{\text{CCSN}} = 1.9 \pm 1.1 \text{ century}^{-1}$
- PSRs birth rate,  $\beta_{\text{PSR}} \sim 1.4 \text{ century}^{-1}$
- RRaTs birth rate,  $\beta_{\text{RRAT}} > 4 \text{ century}^{-1}$
- Magnetars birth rate,  $\beta_{\text{Mag}} \sim 0.1 \text{ century}^{-1}$
- CCOs birth rate,  $\beta_{\text{CCO}} \sim 0.04 \text{ century}^{-1}$
- XDINSs birth rate,  $\beta_{\text{XDINS}} \sim 2 \text{ century}^{-1}$

Total NS birthrate

$$\beta_{\text{TOT}} \sim 1.4 + 4 + 0.1 + 0.04 + 2 = 6.5 \text{ century}^{-1} > \beta_{\text{CCSN}}$$

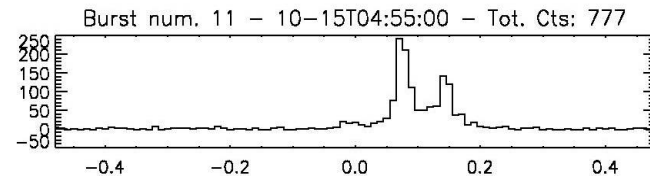
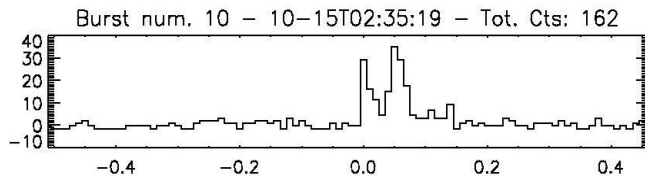
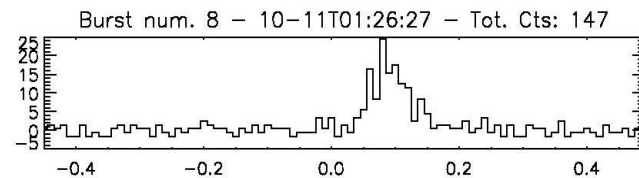
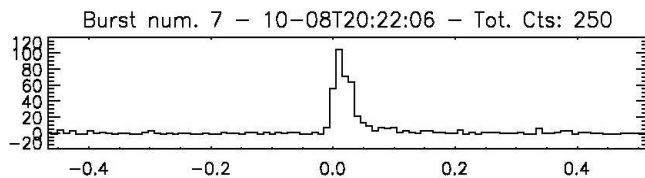
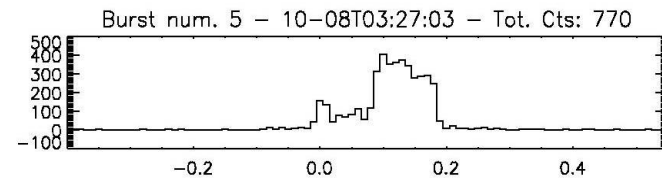
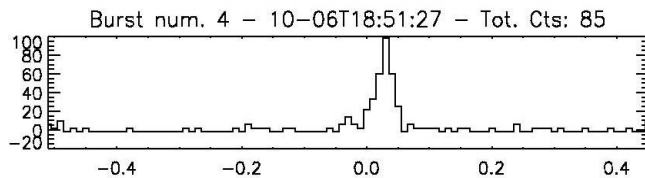
EVOLUTION !



# Soft Gamma Repeaters - I

Rare class of sources, discovered through the emission of strong, recurrent (whence the name) bursts of soft  $\gamma$ -/hard X-rays:

$L \approx 10^{39}-10^{41}$  erg/s  $\gg L_{\text{Edd}}$ , duration 0.1 - 1 s



# Soft Gamma Repeaters - II

- Much more energetic “Giant Flares” (GFs,  $L \approx 10^{44} - 10^{47}$  erg/s,  $t_{\text{peak}} \sim 1$  s,  $t_{\text{tail}} \sim 300$  s) detected from 3 sources
- No evidence for a binary companion, association with a SNR in 1 (?) case
- Persistent X-ray emitters,  $L \approx 10^{33} - 10^{35}$  erg/s
- Pulsations discovered both in GFs tails and persistent emission,  $P \approx 2 - 10$  s
- Huge spin-down rates, as compared to PSRs,  $\dot{P} \approx 10^{-11} - 10^{-10}$  ss<sup>-1</sup>





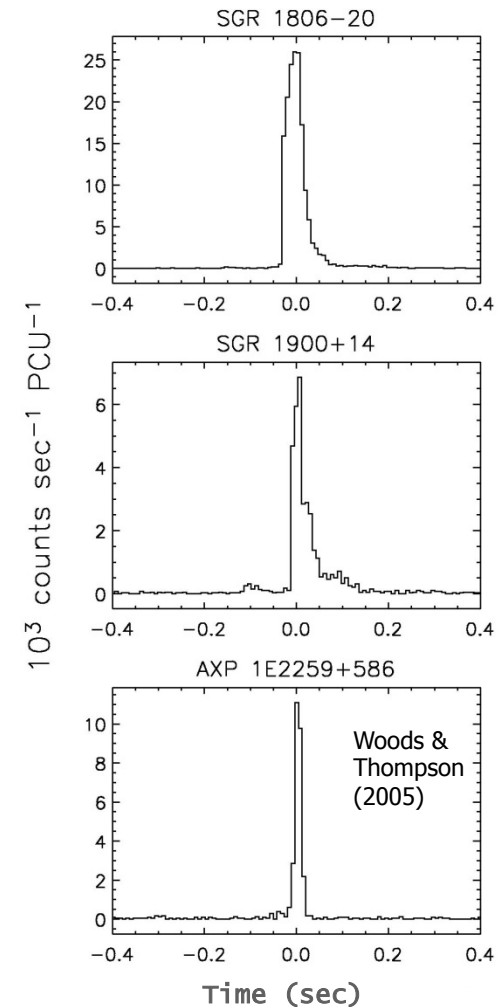
# Anomalous X-ray Pulsars - I

- Peculiar class of persistent X-ray sources,  $L \approx 10^{33} - 10^{35}$  erg/s
- Spin-down luminosity  $\dot{E} < L_X$  (not powered by rotation, hence “anomalous”)
- Pulsations with  $P \approx 2 - 10$  s
- Large spin-down rates,  $\dot{P} \approx 10^{-11}$  ss<sup>-1</sup>
- No evidence for a binary companion, association with a SNR in six (?) cases

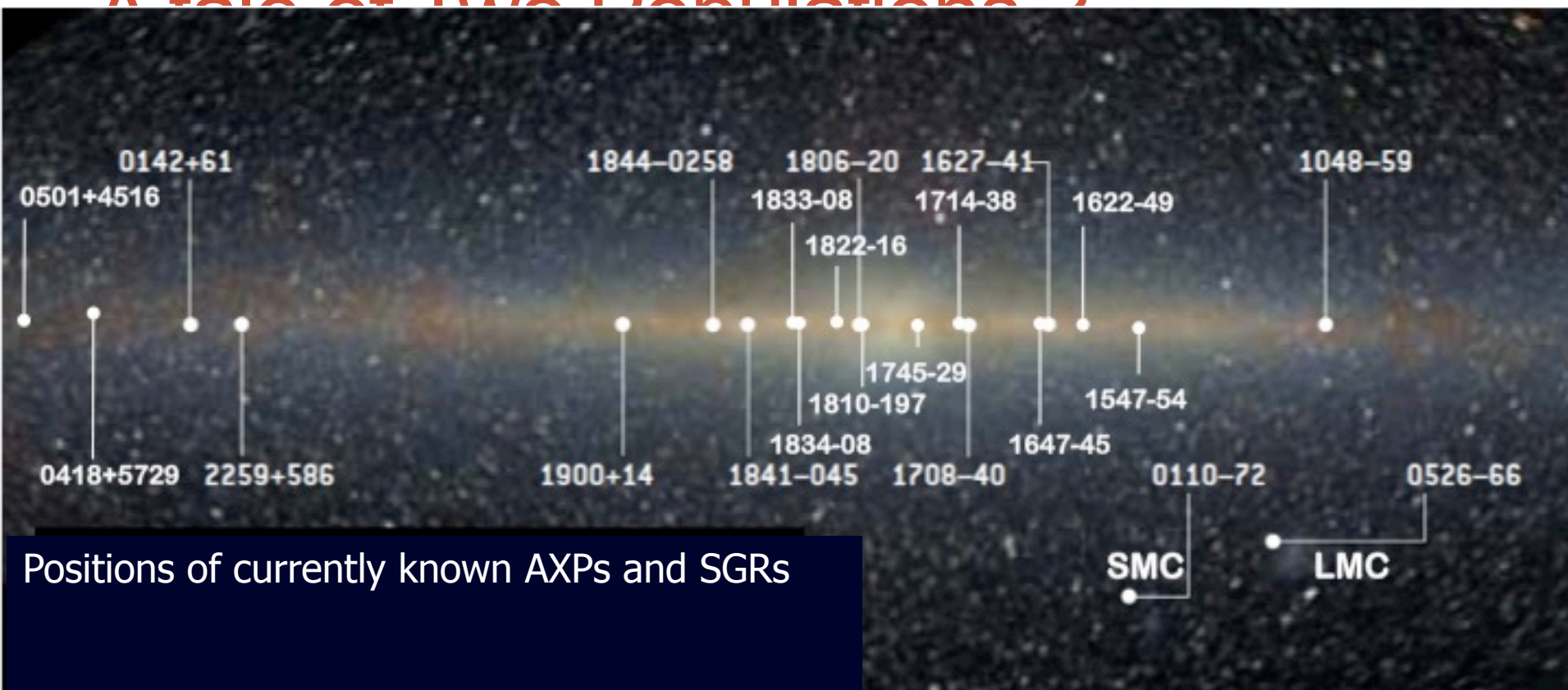


# Anomalous X-ray Pulsars - II

- Bursts of soft  $\gamma$ -/hard X-rays quite similar to those of SGRs detected first in AXP 1E 2259+586 (Gavriil et al. 2002; Kaspi et al. 2003) and then in all (but one) AXPs



## A tale of Two Populations?



- Several clues indicate that these are (isolated) neutron stars
  - $R < ct_{\text{rise}} \approx 100 \text{ km}$
  - pulsations



# Magnetars

- Strong convection in a rapidly rotating ( $P \sim 1$  ms) newborn neutron star generates a very strong magnetic field via dynamo action
- Magnetars: neutron stars powered by their own magnetic energy (surface field  $B >$  a few  $B_{\text{QED}} \sim 10^{14}$  G; Duncan & Thomson 1992; Thomson & Duncan 1993)
- Rapid spin-down due to magneto-dipolar losses,  

$$\dot{P} = -1.1 \times 10^{-11} (B/10^{14} \text{ G})^2 P^{-1} \text{ s/s}$$



# Why magnetars ?

- $L_x > \dot{E}$  → not powered by rotation
- No evidence for a companion star → not powered by accretion
- Quite young objects ( $\approx 10^3 - 10^4$  yrs): spin down to present periods (a few seconds) requires  $B > 10^{14}$  G
- Large measured spin-down rates
- Giant flares energetics requires  $B > 10^{15}$  G
- Opacity suppression for the X polarization mode in a strong B-field explains the large, super-Eddington flux in bursts

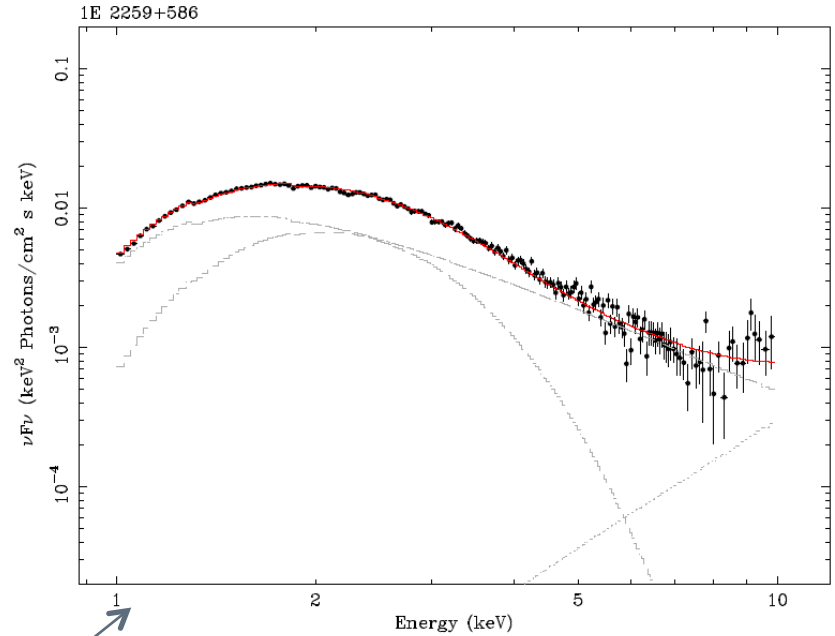
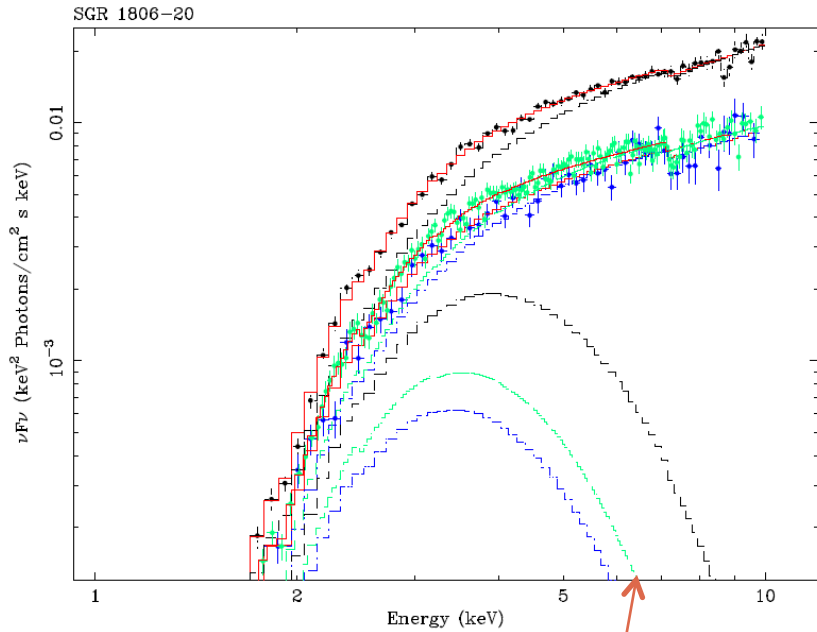
No direct measure of a super-strong field until recently



# SGRs and AXPs X-ray Spectra

- 0.5 – 10 keV emission well represented by a blackbody plus a power law
- $kT_{\text{BB}} \sim 0.5$  keV, does not change much in different sources
- Photon index  $\Gamma \approx 1 - 4$ , AXPs tend to be softer
- SGRs and AXPs persistent emission is variable (months/years)
- Variability mostly associated with the non-thermal component
- Transient spectra can be BB+BB,  $T_{\text{BB}}$  and  $R_{\text{BB}}$  decrease in time

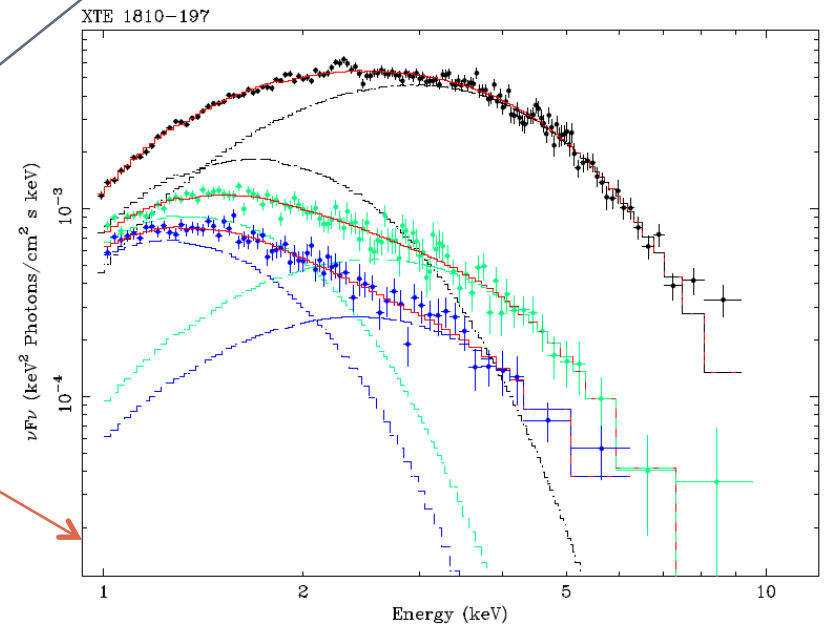




SGR 1806-20 at different epochs (BB+PL)

AXP 1E 2259-586 (BB+PL)

Transient AXP XTE 1810-197 at different epochs (BB+BB)



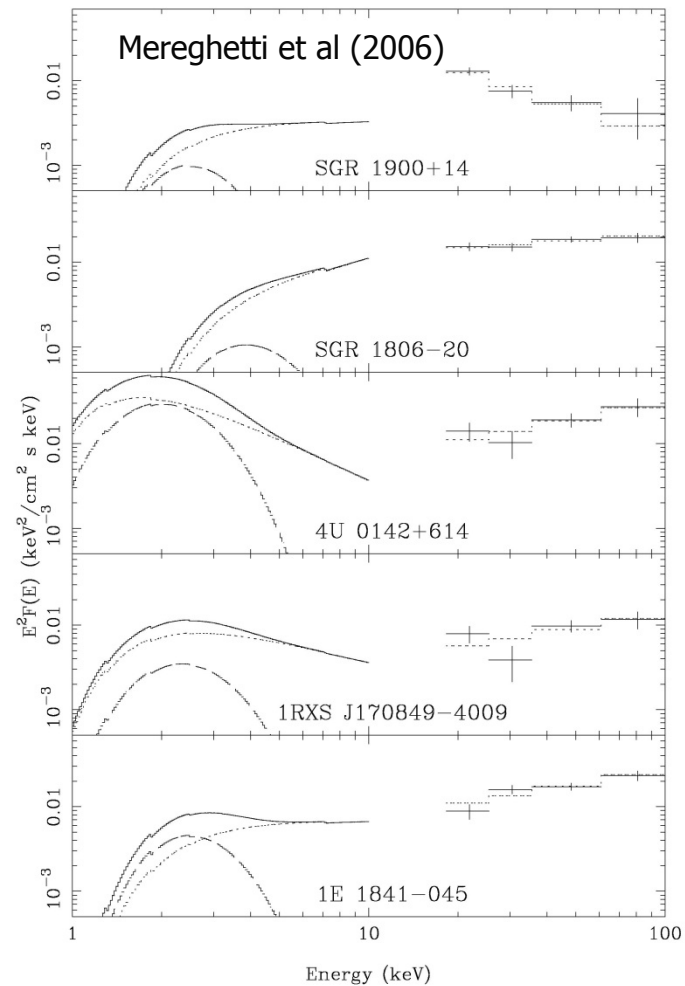
XMM Epic-pn data (Rea et al. 2008)

# Hard X-ray Emission

INTEGRAL revealed substantial emission in the 20 -100 keV band from SGRs and AXPs

Hard power law tails with  $\Gamma \approx 1-3$ , hardening wrt soft X-ray emission required in AXPs

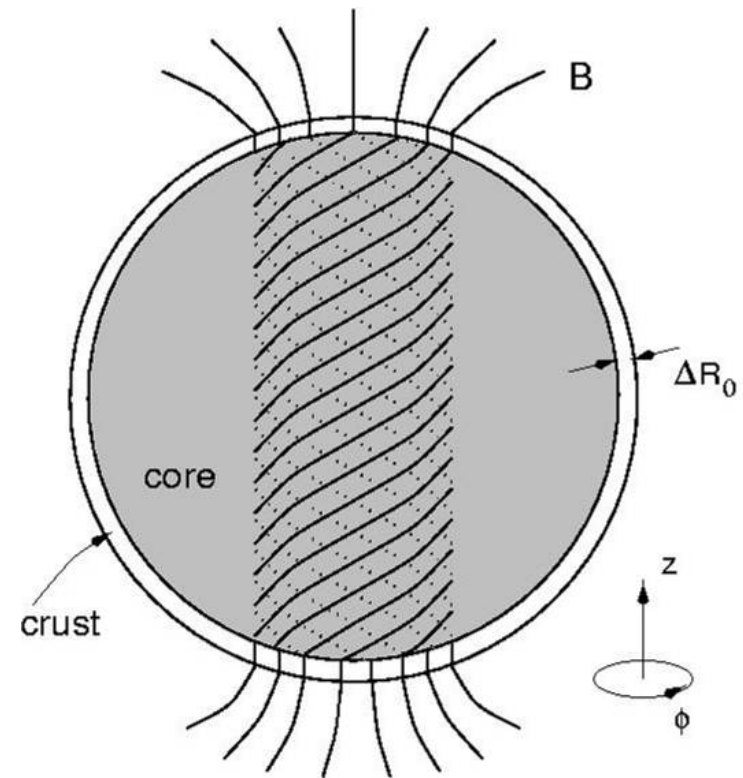
Hard emission highly pulsed





# Twisted Magnetospheres – I

- The magnetic field inside a magnetar is “wound up”
- The presence of a toroidal component induces a rotation of the surface layers
- The crust tensile strength resists
- A gradual (quasi-plastic ?) deformation of the crust
- The external field twists up  
(Thompson, Lyutikov & Kulkarni 2002)



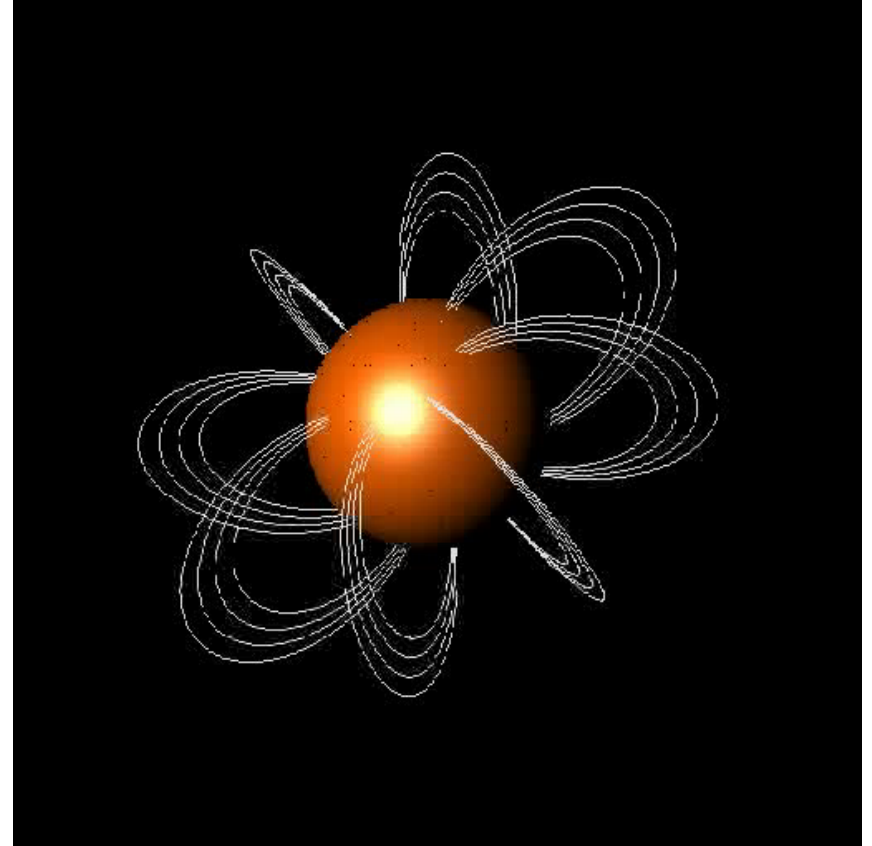
Thompson & Duncan (2001)



# Twisted Magnetospheres - II

- Twisted fields are non-potential,  $\nabla \times \vec{B} \neq 0$
- Globally twisted dipole (Thompson, Lyutikov & Kulkarni 2002, Pavan et al. 2009)
- A sequence of models labeled by the twist angle

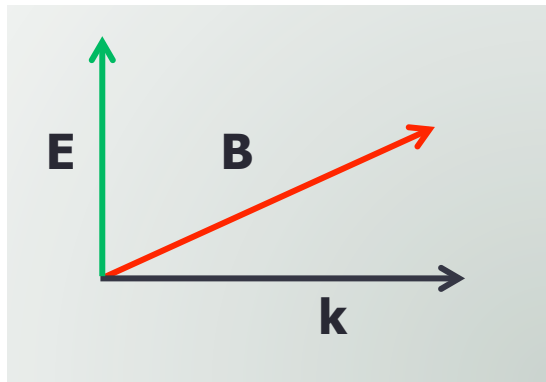
$$\Delta\phi_{N-S} = 2 \int_0^{\pi} \frac{B_{\phi}}{B_{\theta} \sin \theta} d\theta$$



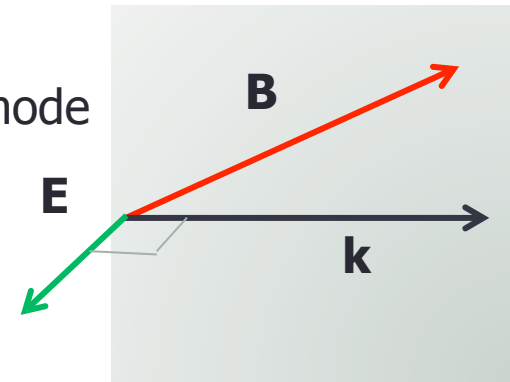
# Photons in a Magnetized Medium

- A magnetized plasma is anisotropic and birefringent, radiative processes sensitive to polarization state
- Two normal, linearly polarized modes in the magnetized vacuum: the extraordinary (X) and ordinary (O) mode
- Opacities greatly reduced for X-mode photons

O mode



X mode



The electron scattering cross section in the ERF is resonant at the cyclotron frequency and its harmonics

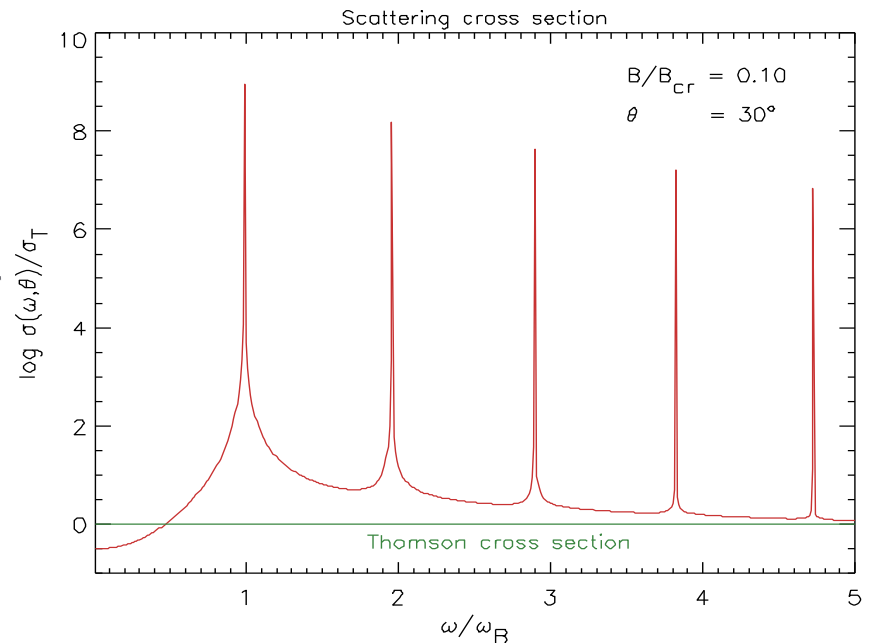
$$E_n = n\hbar\omega_B = n \frac{eB}{mc}$$

For an electron moving with velocity  $\beta$  the (first) resonance is at a frequency

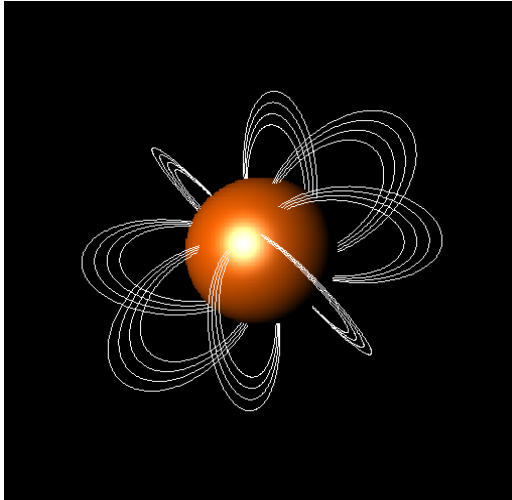
$$\omega_D = \frac{\omega_B}{\gamma(1 - \beta \cos \theta)}$$

$$\sigma_{1-1} = \frac{1}{3}\sigma_{1-2} = \frac{\pi^2 r_0 c}{2} \delta(\omega - \omega_D) \cos^2 \theta_{\text{ERF}}$$

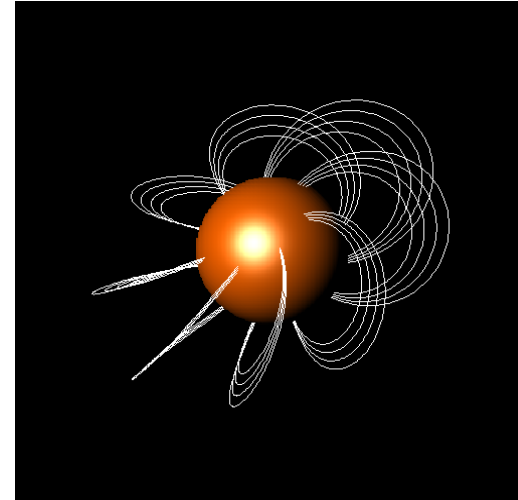
$$\sigma_{2-2} = 3\sigma_{2-1} = \frac{3\pi^2 r_0 c}{2} \delta(\omega - \omega_D),$$



# Magnetospheric Currents - I



$$\nabla \times B = 0$$



$$j = \frac{c}{4\pi} \nabla \times B$$

Contrary to PSRs, currents flow (also) along the closed field lines and  $j \gg j_{GJ}$



# Magnetospheric Currents - II

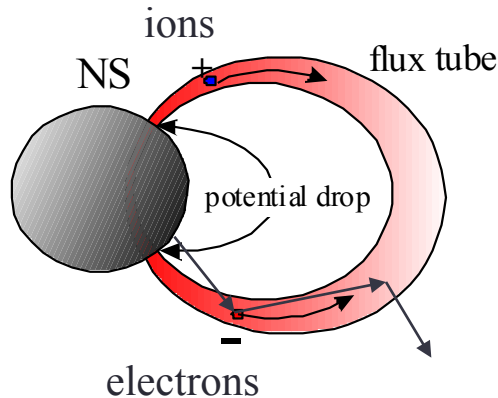
The twist must decay to support its own currents. A parallel electric field develops which accelerates the charges along the flux tube (Beloborodov & Thompson 2007; Beloborodov 2009)

$$\frac{\partial(B_{\phi}^2 / 8\pi)}{\partial t} = -E_{\parallel} j \qquad \frac{\partial E_{\parallel}}{\partial t} = j - j_B$$

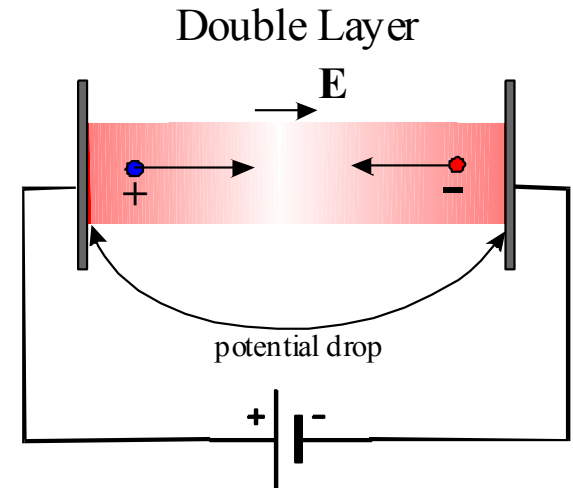
The electric field is self-regulated to ensure that the required current flows in the circuit

A potential drop  $\Phi$  is maintained between the footpoints  $j = j(\Phi)$  depends on the nature of the discharge and this fixes the duration of the twist





A simple  
analogue



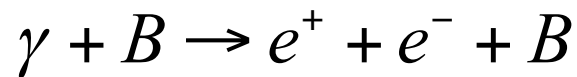
In a relativistic double layer  $j \propto \frac{m_e}{m_i} \Phi^2$ ,  $\gamma_e \approx 1 + \frac{e\Phi}{m_e c^2}$

$\Phi$  (and  $E_{\parallel}$ ) must be huge ( $\approx 10^{12}$  GeV) in order to produce  $j_B$ :  
 $\gamma_e \approx 10^9$  and the twist decays immediately



Where  $B > 2B_Q$ , 1 keV photons scatter onto  $\gamma > 1000$  electrons  
 Scattered photons have energy  $\varepsilon'$  in the MeV range and initially propagate along  $B$

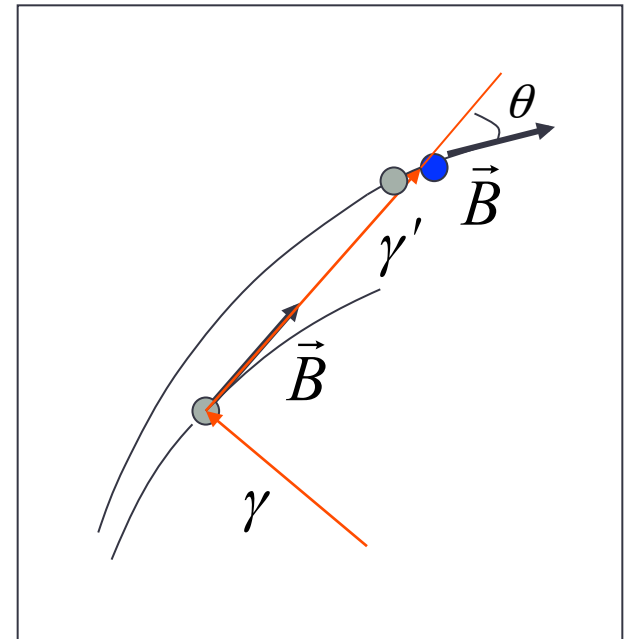
They quickly convert into pairs via



as soon as  $\varepsilon' > \frac{2m_e c^2}{\sin \theta}$

Pair production along the entire circuit screens the potential:  $j_B$  can be conducted with  $\Phi \ll \Phi_{DL}$

A quasi-stationary state in which the particle energy is just that required to ignite the pair cascade (Beloborodov & Thompson 2007)



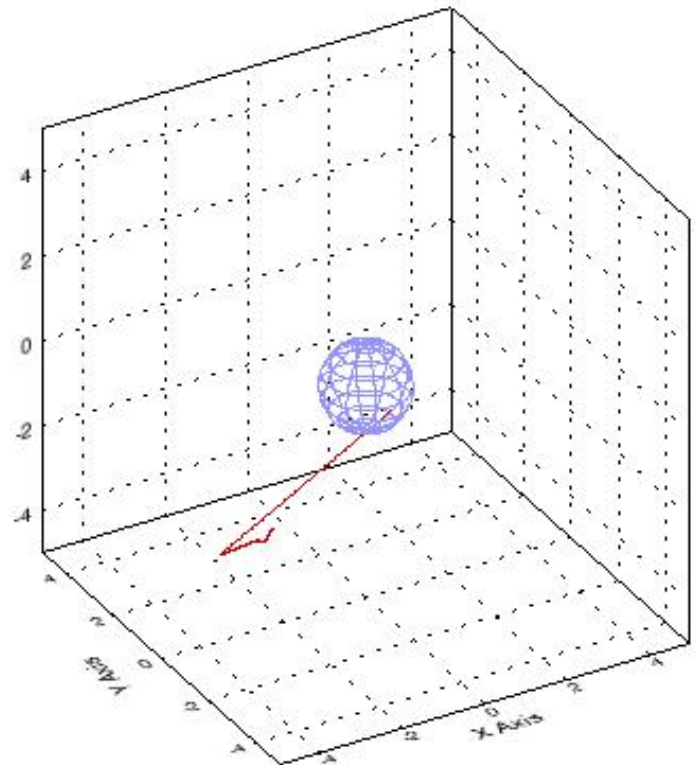


# Resonant Compton Scattering

- The current flowing along the closed field lines is

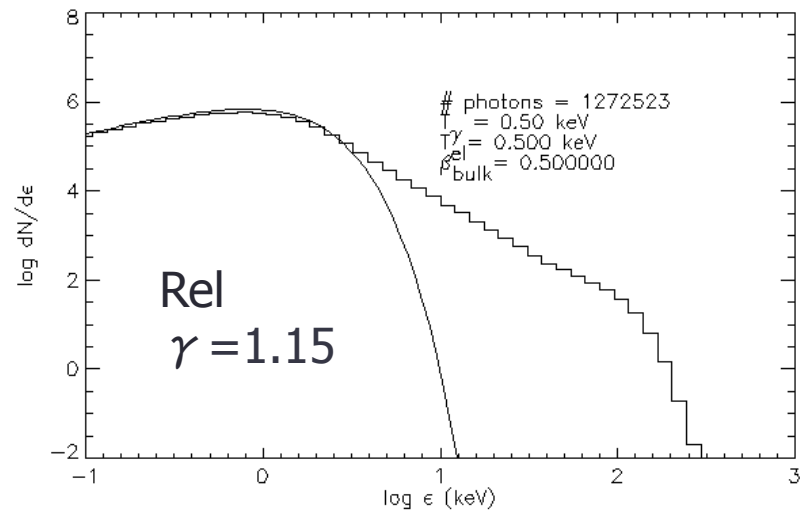
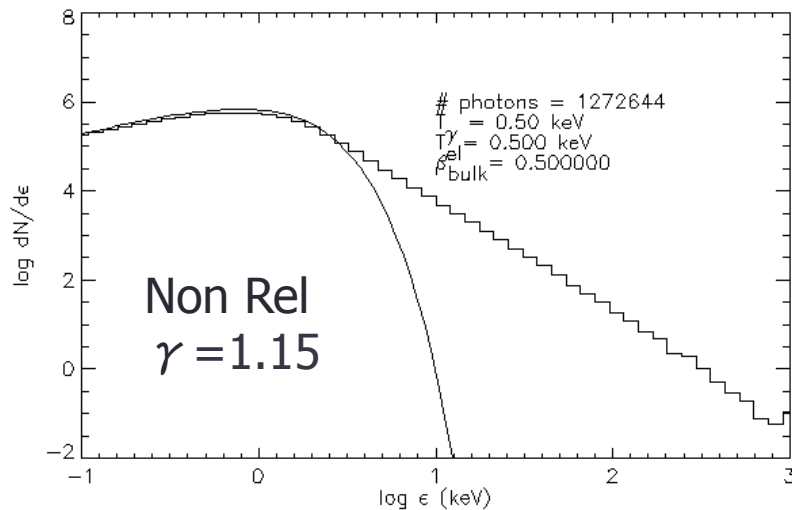
$$j = (c/4\pi) \nabla \times B \Rightarrow n_e e = p + 1/4\pi e B \downarrow \varphi / B,$$

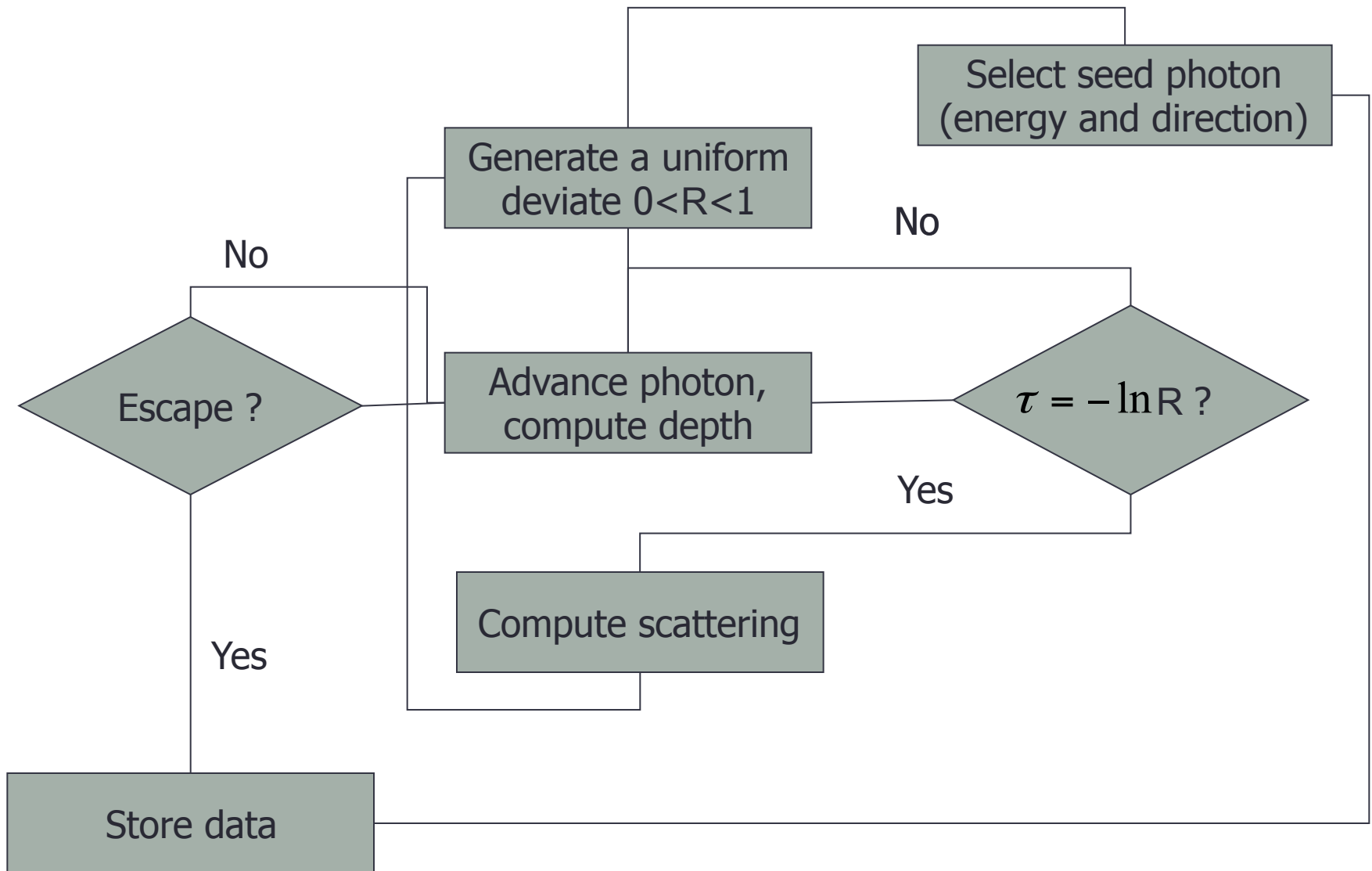
- The optical depth for Thomson  $n_e \sigma_T r \approx 10^{-4}$
- At resonance  $\sigma \approx 10^5 \sigma_T \rightarrow$  resonant cyclotron scattering
- Up-scattering of thermal photons from cooling surface onto mildly relativistic electrons

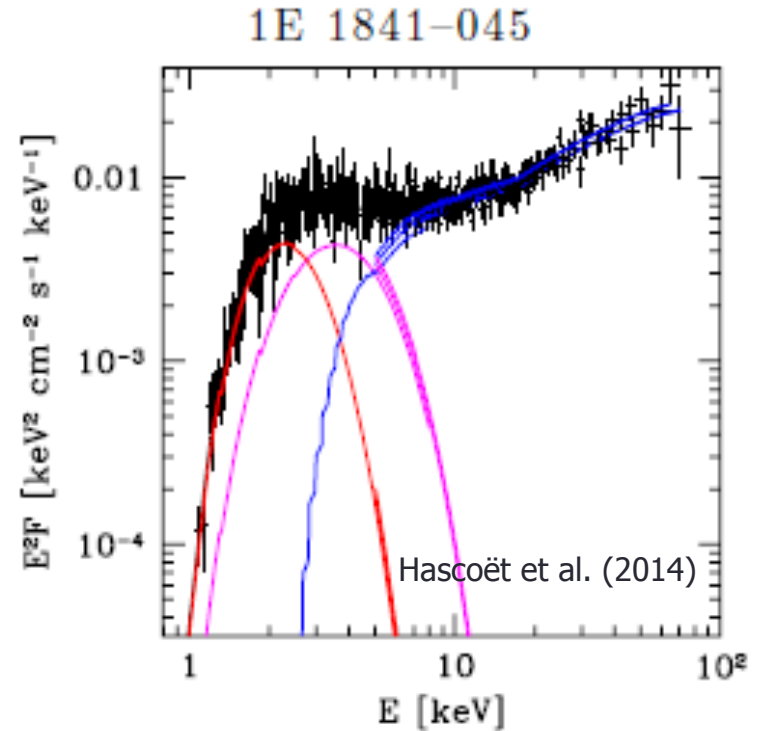
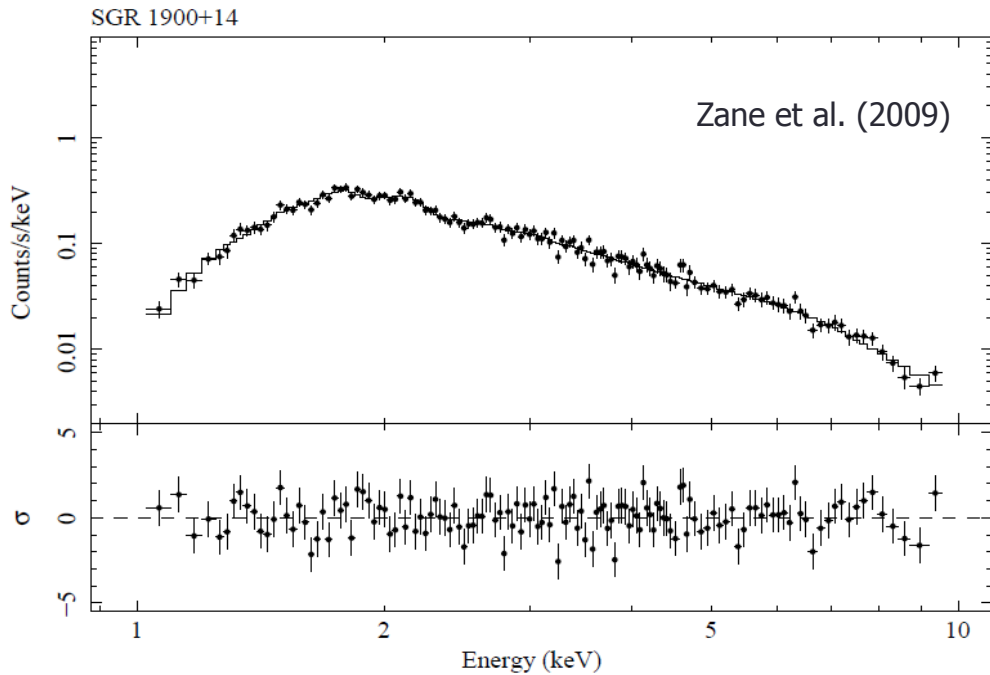


Repeated scatterings lead to the formation of a power-law tail because  $\omega_D = \omega_D(r, \theta)$  and  $r_{\text{current}} > R_{\text{NS}}$

Spectral formation in twisted magnetospheres investigated quite in detail using Montecarlo methods (Lyutikov & Gavriil 2006; Fernandez & Thompson 2007; Nobili, Turolla & Zane 2008a, b)







RCS models quite successful in explaining magnetars soft X-ray spectra ( $\sim 0.5 - 10$  keV) and also high-energy tails

Spectral fits provide information on the physical state of the star/magnetosphere (twist angle, charge velocity, surface temperature, etc)

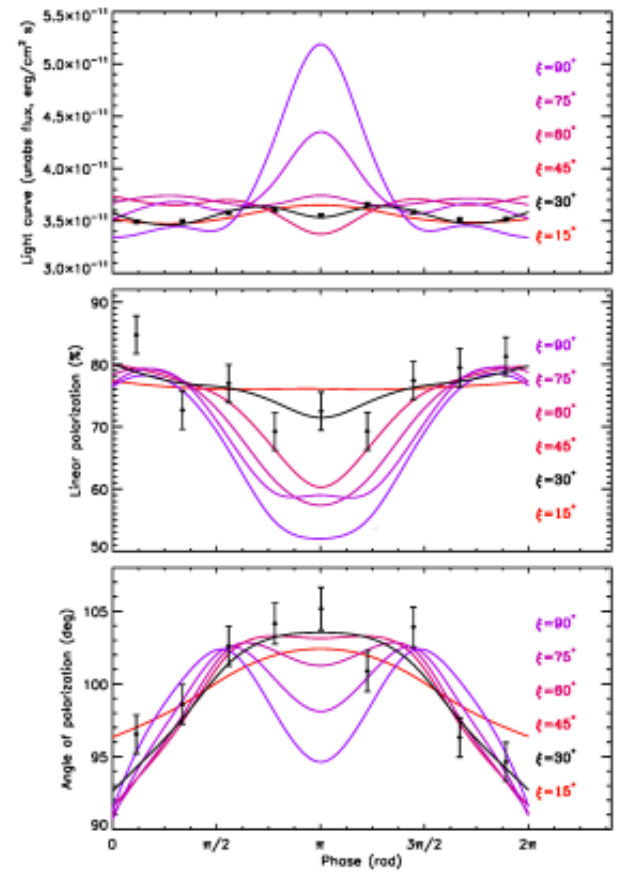
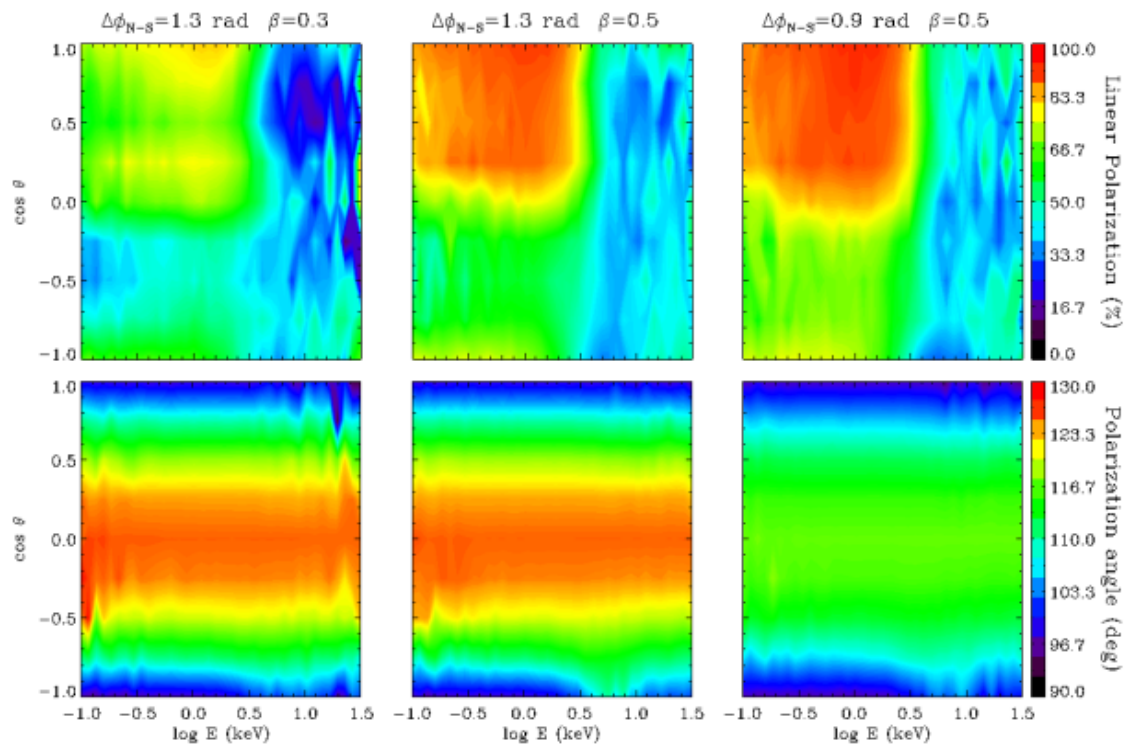
No single spectral model can consistently explain observations in the 0.5-100 keV band though



# X-ray Polarization

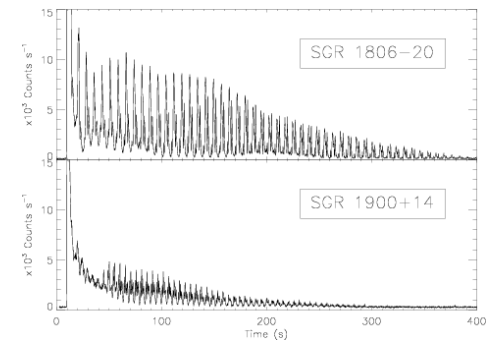
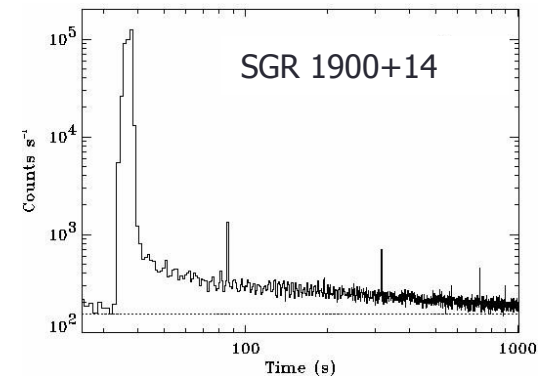
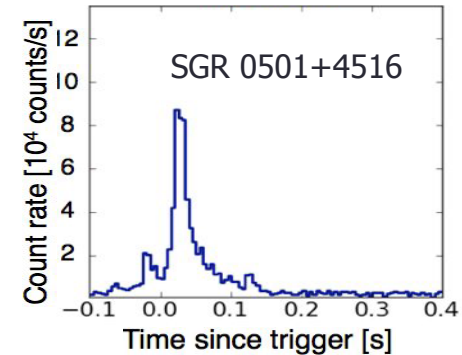
- Thermal surface emission highly polarized in the X-mode
- Scatterings can change the photon polarization state
- The observed polarization fraction and polarization angle depend on QED effects (“vacuum polarization”) and on magnetic field geometry (Stokes parameters rotation)
- $\Pi_L$  and  $\chi_p$  very sensitive to the source geometry (inclination of the LOS and of the magnetic axis wrt the rotation axis)
- X-ray polarimetry will provide an entirely new tool in magnetar studies
- XIPE and IXPE proposed for ESA M4 and NASA Smex programmes





# Bursts & Flares

- Short bursts
  - $t \sim 0.1 - 1$  s,  $L \sim 10^{39} - 10^{41}$  erg/s ,  
thermal spectrum ( $kT \sim 10$  keV), seen  
in both SGRs and AXPs
- Intermediate bursts
  - $t \sim 1 - 40$  s,  $L \sim 10^{41} - 10^{43}$  erg/s ,  
thermal spectrum, seen in both SGRs  
and AXPs
- Giant flares
  - only three observed, each from a  
different SRG,  $L \sim 10^{44} - 10^{47}$  erg/s,  
initial spike ( $\sim 0.1$ s) + pulsating tail ( $\sim$   
100 s)



# Burst Trigger Mechanism(s)

Rapid magnetic field reconfiguration is a key ingredient, but no precise model as yet

Secular magnetic evolution builds stresses that are released catastrophically in the bursts

Alvén speed  $v_A = 10^8 \text{cm/s} \left( \frac{B}{10^{16} \text{G}} \right) \left( \frac{10^{15} \text{g/cm}^3}{\rho} \right)^{1/2}$

Shear velocity

$$v_s = 1.1 \times 10^8 \text{cm/s} \left( \frac{\rho}{10^{14} \text{g/cm}^3} \right)^{1/6} \left( \frac{Z}{38} \right) \left( \frac{302}{A} \right)^{2/3} \left( \frac{1 - X_n}{0.25} \right)^{2/3}$$





- Magnetic evolution leads to an unstable configuration in the core  $\Rightarrow$  MHD instabilities with growth time  $\approx R/v_A \approx 0.1$  s
- Magnetic stresses rupture the crust  $\Rightarrow$  release of elastic energy over a timescale  $\approx \pi R/v_s \approx 0.3$  s
- Core and crust evolve smoothly, stresses are released in the magnetosphere via plasma instabilities/magnetic reconnection  $\Rightarrow$  very short timescale,  $< 0.01$  s ( $v_A \sim c$ )

All three scenarios provide timescale in rough agreement with burst duration/rise time

No serious problem with energetics (including giant flares)



# Burst Emission

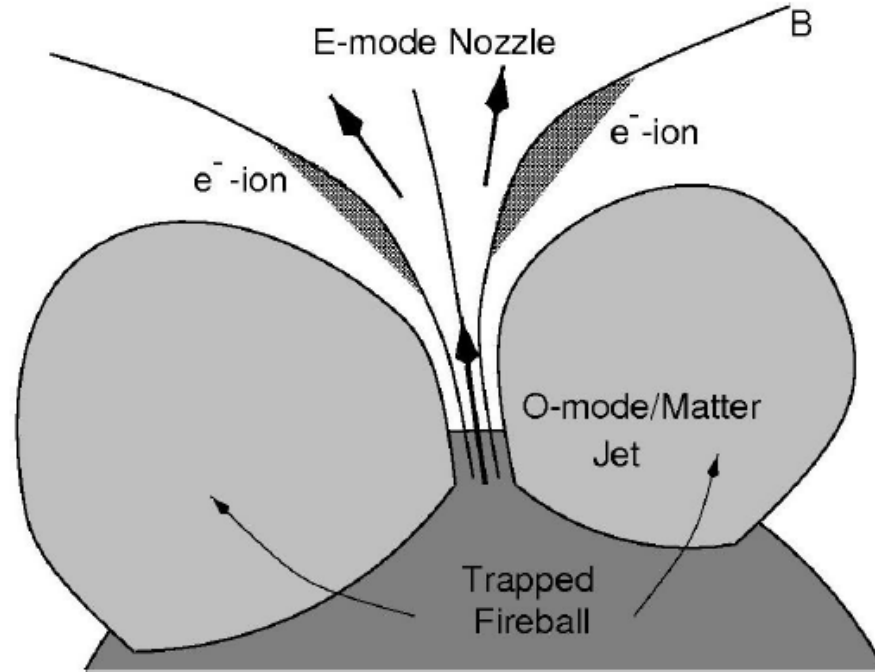
- Magnetic reconfiguration produces particle acceleration
- Electrons moving along the curved field lines emit  $\gamma$ -rays which drive a pair cascade

- The pair plasma is confined by the magnetic field if

$$B_{\text{dipole}} > 2 \times 10^{14} \left( \frac{E_{\text{fireball}}}{10^{44} \text{ erg}} \right)^{1/2} \left( \frac{\Delta R}{10 \text{ km}} \right)^{-3/2} \left( \frac{1 + \Delta R/R}{2} \right)^3 \text{ G}$$

- Confinement leads to an optically thick “fireball”
- Radiation escapes preferentially in the X-mode, due to the much reduced opacity





No detailed model for burst emission available as yet

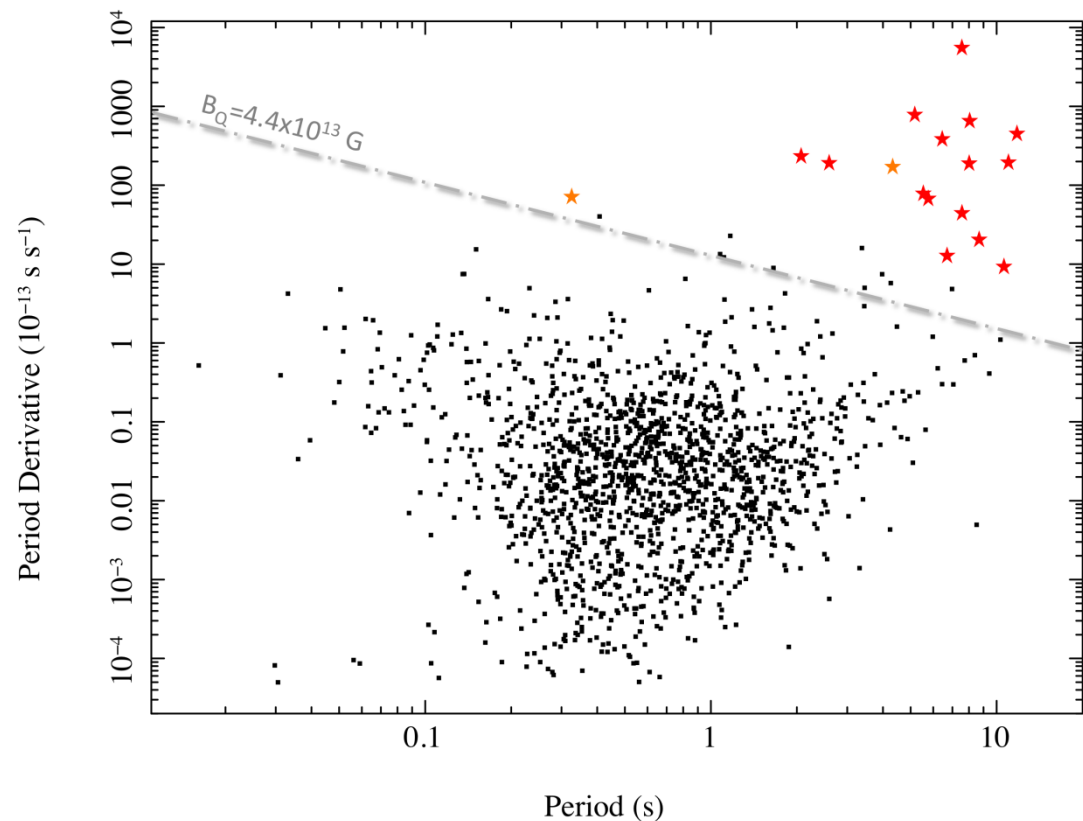


# SGRs/AXPs vs. High-B Pulsars

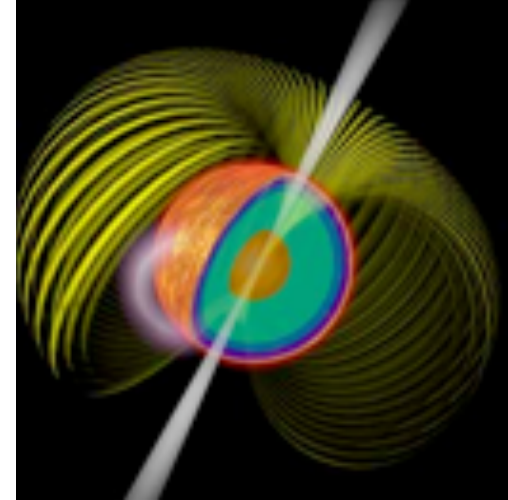
“Magnetar activity” (bursts, outbursts, ...) for a long time detected only in high-B sources ( $B_p > 5 \times 10^{13}$  G) : AXP+SGRs (★) and PSR J1846-0258, PSR J1622-4950 (★)

The ATNF Catalogue lists 20 PSRs with  $B_p > 5 \times 10^{13}$  G (HBPSRs)

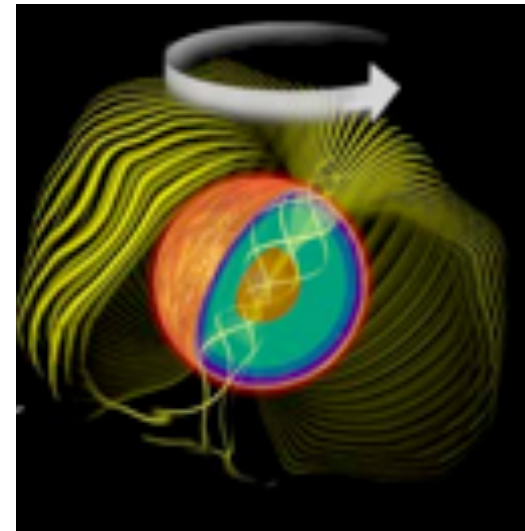
A high dipole field does not always make a magnetar, but a magnetar has necessary a high dipole field



- What really matters is the internal toroidal field  $B_\phi$
- A large  $B_\phi$  induces a rotation of the surface layers
- Deformation of the crust  $\Rightarrow$  fractures  $\Rightarrow$  bursts/twist of the external field



High-B PSR

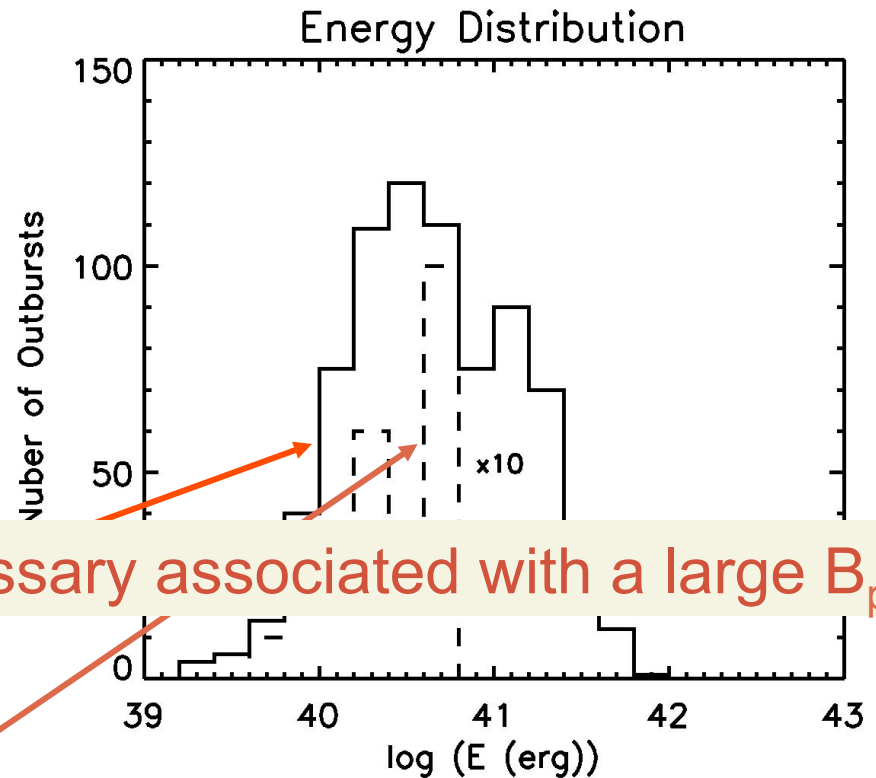


SGR/AXP



Calculation of magnetic stresses acting on the NS crust at different ages  
(Perna & Pons 2011; Pons & Perna 2011)

Activity strongly enhanced when  $B_{\text{tor},0} > B_{\text{p},0}$



A large  $B_{\text{tor}}$  is necessary associated with a large  $B_{\text{p}}$

$$B_{\text{tor},0} = 2.5 \times 10^{14} \text{ G}$$

$$B_{\text{tor},0} = 8 \times 10^{14} \text{ G}$$

$$B_{\text{p},0} = 1.6 \times 10^{14} \text{ G}$$



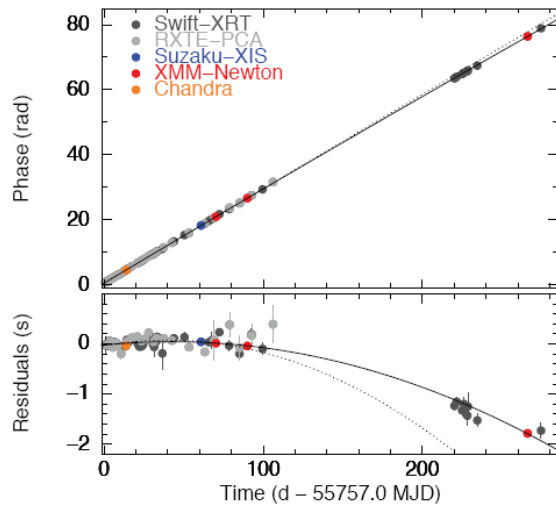
# The “Low-Field” Magnetars

- Three peculiar magnetar candidates discovered since 2009: **SGR 0418+5729** (van der Horst et al. 2010, Esposito et al. 2010, Rea et al. 2010), **Swift J1822.3–1606** (Rea et al. 2012, Scholz et al. 2012) and **3XMM J1852+0033** (Rea et al. 2014)
- All the features of a (transient) magnetar
  - Rapid, large flux increase and decay
  - Emission of bursts
  - Periods in the range  $\approx 8-11$  s



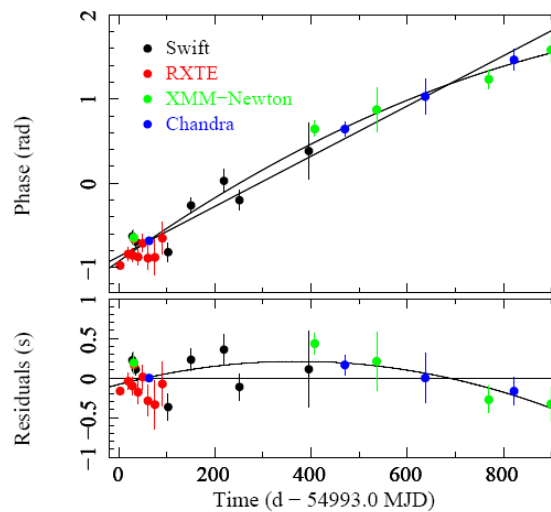
# Hunting for $\dot{P}$

Swift J1822 (Rea et al. 2012)



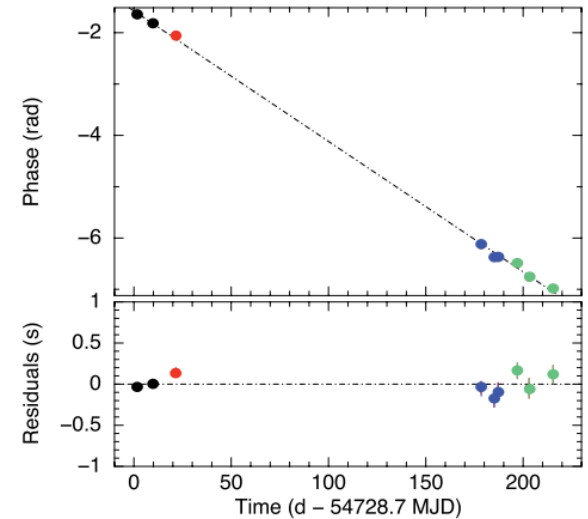
$$\begin{aligned}\dot{P} &= 8.3 \times 10^{-14} \text{ s/s} \\ B_p &= 2.7 \times 10^{13} \text{ G} \\ \tau_c &= 1.6 \text{ Myr}\end{aligned}$$

SGR 0418 (Rea et al. 2013)



$$\begin{aligned}\dot{P} &= 5.14 \times 10^{-15} \text{ s/s} \\ B_p &= 6.9 \times 10^{12} \text{ G} \\ \tau_c &= 29.5 \text{ Myr}\end{aligned}$$

3XMM J1852 (Rea et al. 2014)

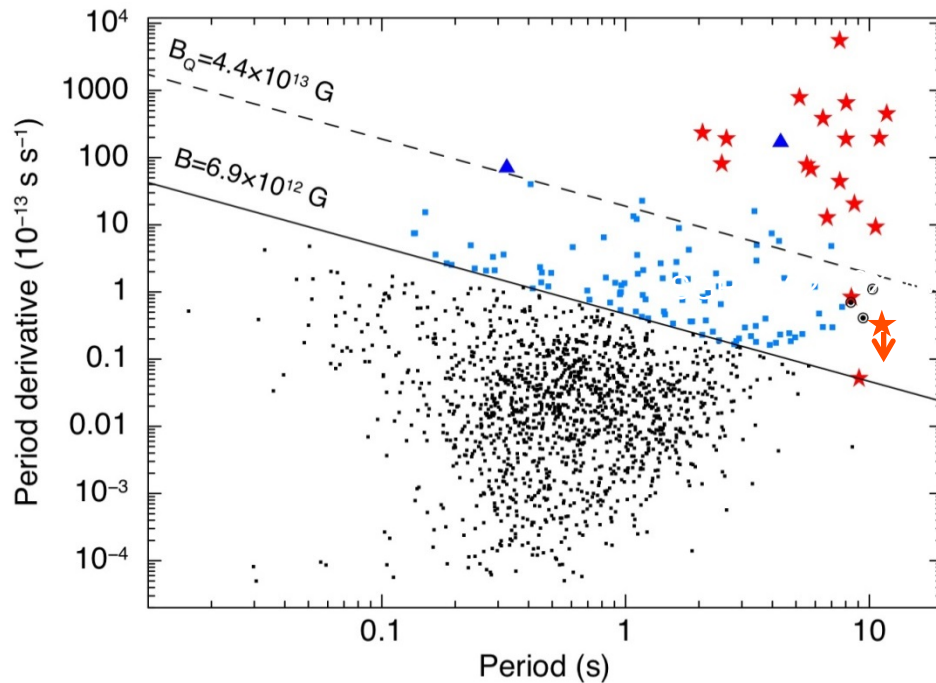


$$\begin{aligned}\dot{P} &< 1.4 \times 10^{-13} \text{ s/s} \\ B_p &< 4.1 \times 10^{13} \text{ G} \\ \tau_c &> 0.1 \text{ Myr}\end{aligned}$$





# Dr Pulsar and Mr Magnetar



Three “active magnetars”  
with B-field well within  
PSR range

More than 20% of known  
PSRs have  $B_p$  stronger  
than SGR 0418

A continuum of  
magnetar-like activity  
across the P- $\dot{P}$  diagram

**A supercritical  $B_p$  not required to make a magnetar !**



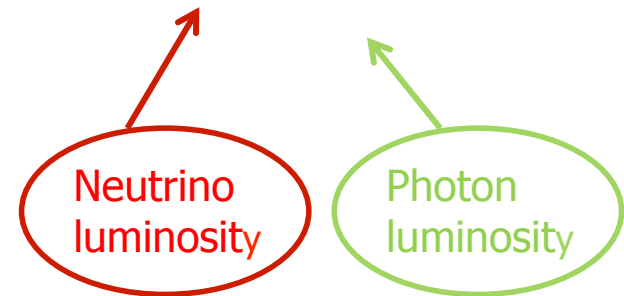
# Neutron Star Evolution

- Rotational evolution

$$I\Omega\dot{\Omega} = -\frac{B_p^2 R^6 \Omega^4 \sin^2 \alpha}{6c^3}$$

- Thermal evolution

$$dE_{\text{th}}/dt = c\nu dT/dt = -L_{\nu} - L_{\gamma}$$



- Magnetic evolution

$$\partial B / \partial t = -\nabla \times \{ \eta \nabla \times B + c/4\pi en \downarrow e (\nabla \times B) \times B \}$$

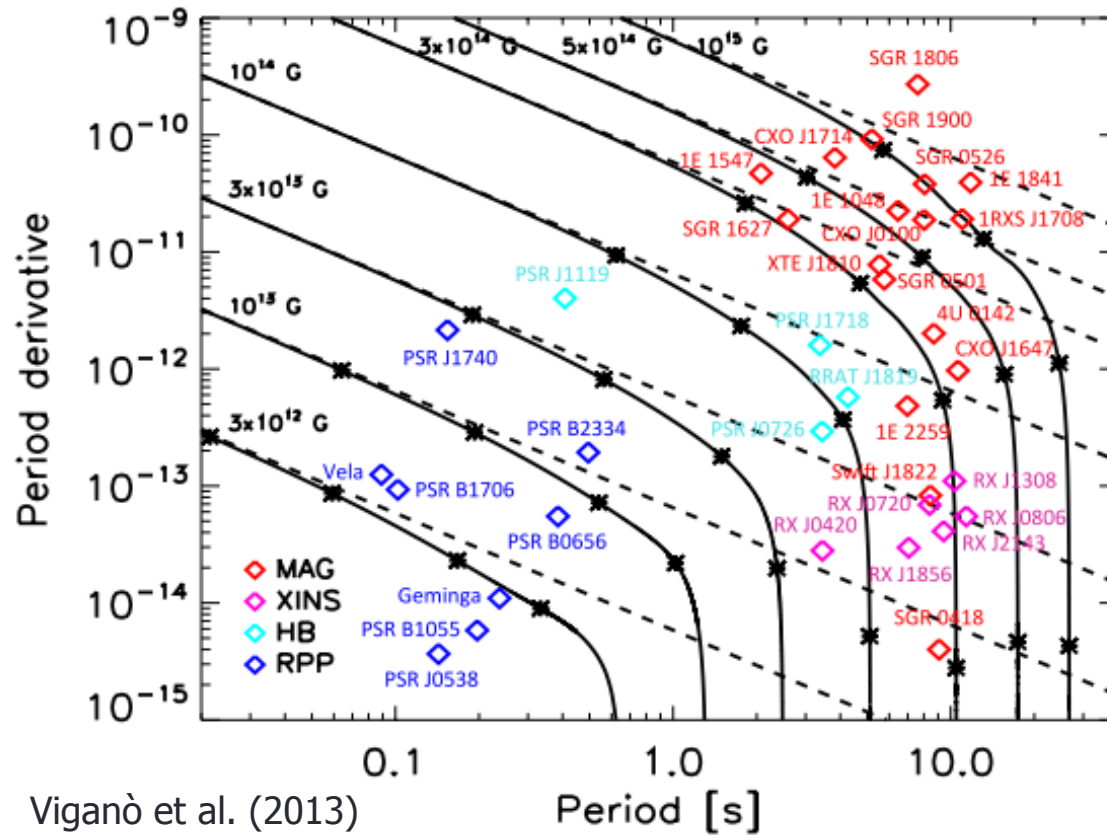


## Faraday induction equation

$\eta$  is the magnetic diffusivity and strongly depends on T

C

N  
(F)

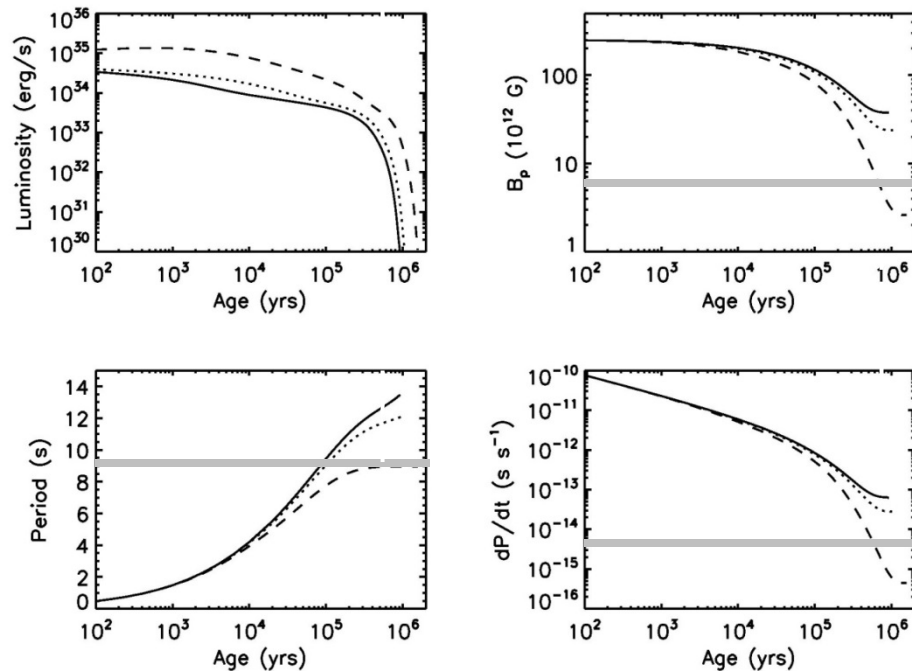


Viganò et al. (2013)

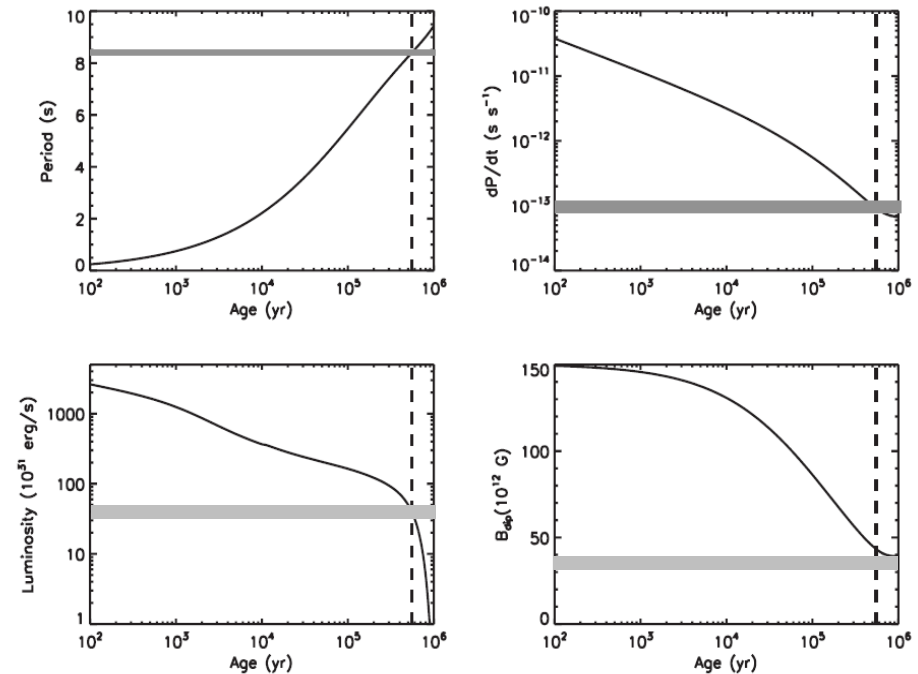


# Are low-field sources old ?

SGR 0418 (Turolla et al. 2011)



SGR 1822 (Rea et al. 2012)



“Low-field” sources look indeed oldish ( $\approx 10^6$  yr) magnetars in which the surface magnetic field substantially decayed



# Wear and Tear

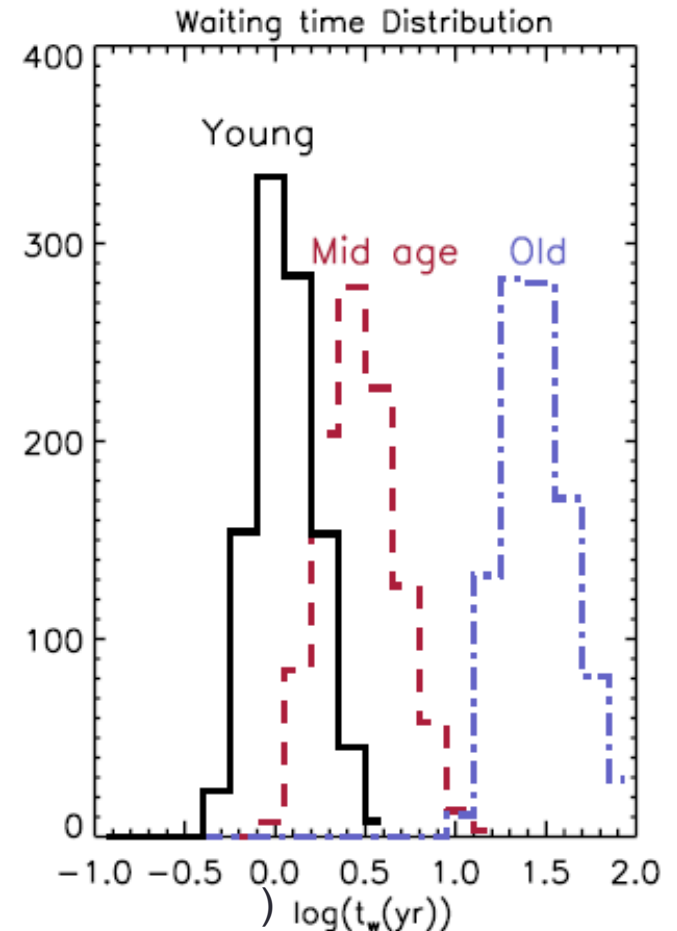
Crustal fractures possible also at late evolutionary phases ( $\approx 10^5 - 10^6$  yr; Perna & Pons 2011)

Burst energetics decreases and recurrence time increases as the NS ages

For  $B_{p,0} = 2 \times 10^{14}$  G and  $B_{tor,0} = 10^{15}$  G,  
 $\Delta t \approx 10 - 100$  yr

Very close to what required for SGR 1822

Fiducial model for SGR 0418 has similar  $B_{p,0}$  and larger  $B_{tor,0}$   $\Rightarrow$  comparable (at least) bursting properties



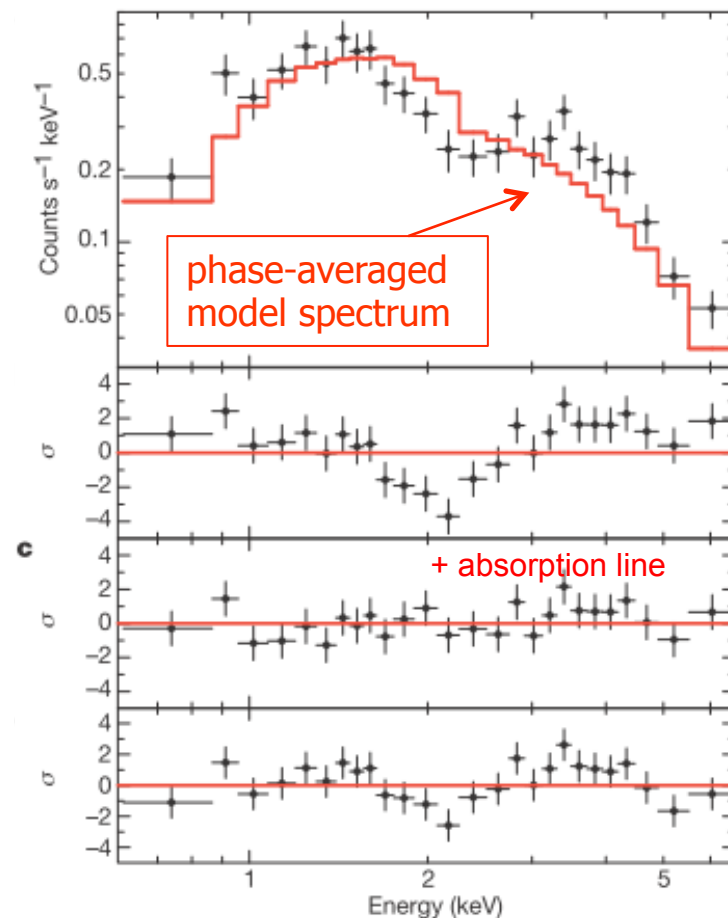
# SGR 0418+5729: The “Surprise Egg”

SGR 0418 went in outburst on June 5, 2009. Subsequent monitoring with RXTE, Swift, Chandra and XMM (Esposito et al. 2010; Rea et al. 2010)

67-ks XMM observation on August 12, 2009 when the source was still quite bright

Phase-averaged spectrum is BB+BB or BB+PL:  $kT_h \sim 0.91$  keV,  $R_h \sim 0.9$  km

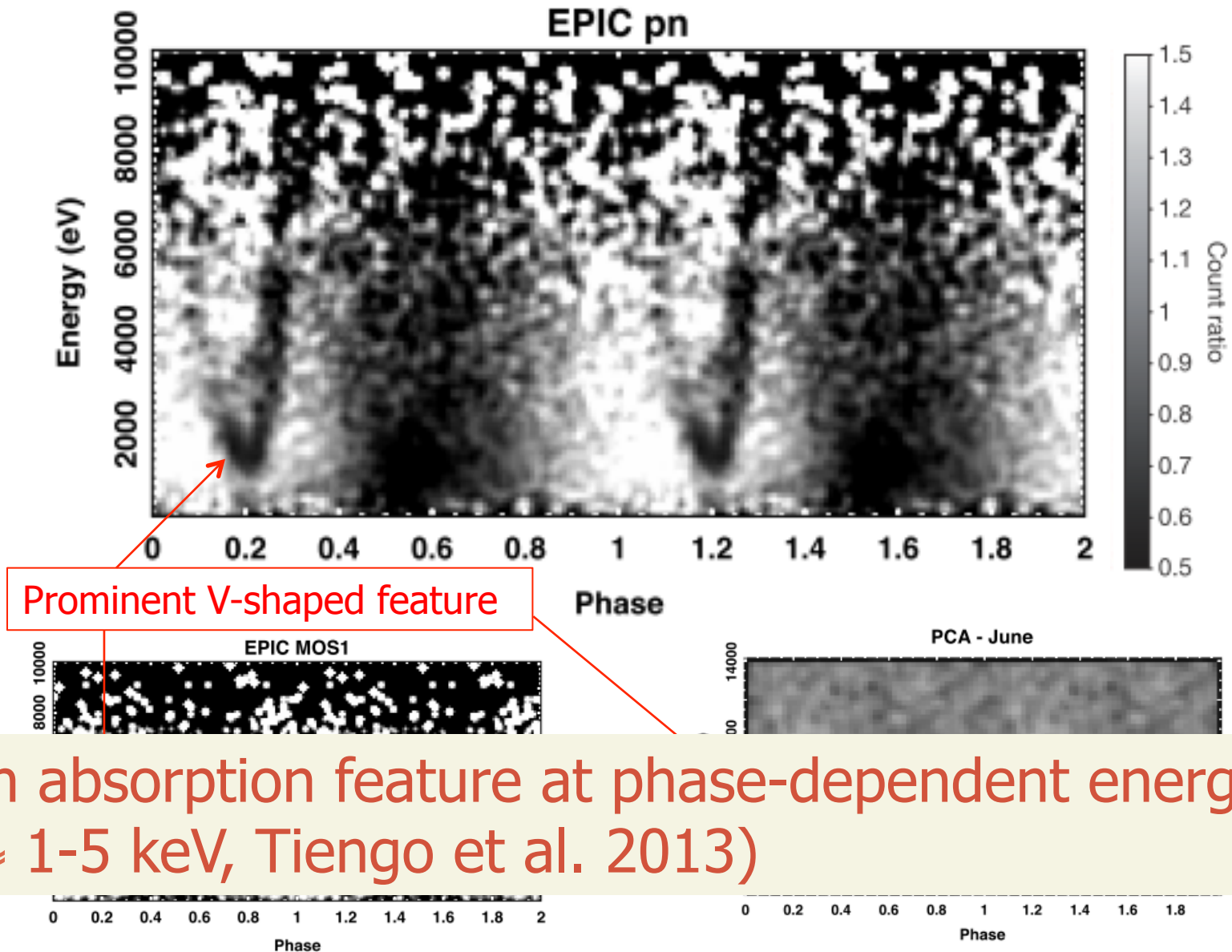
Phase-resolved spectra well fitted by rescaling the phase-averaged one, but not around phases 0.1-0.3 (and 0.5-0.6)



XMM Epic pn phase-resolved spectrum of SGR 0418 (phase interval 0.15-0.17)



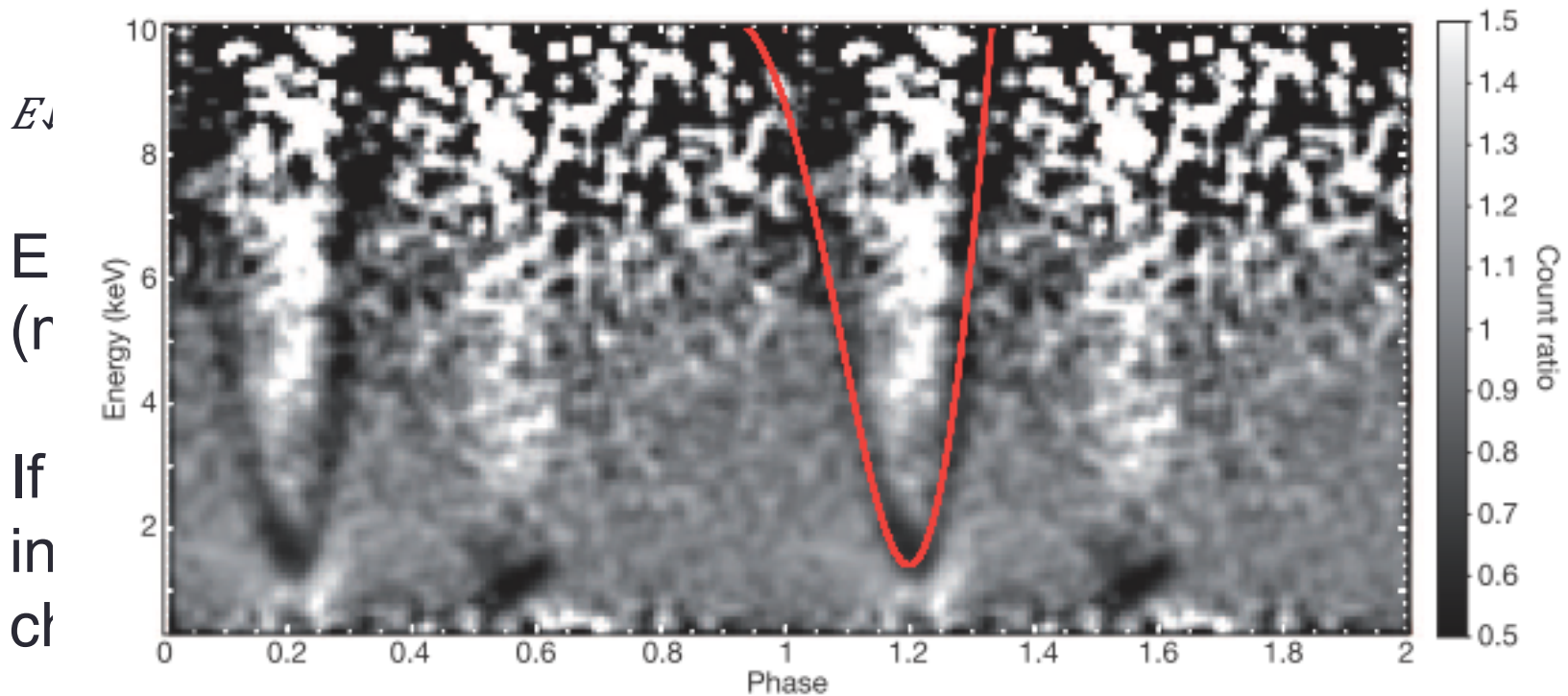
# Normalized phase-energy images (independent on spectral modelling)



An absorption feature at phase-dependent energy ( $\approx 1-5$  keV, Tiengo et al. 2013)

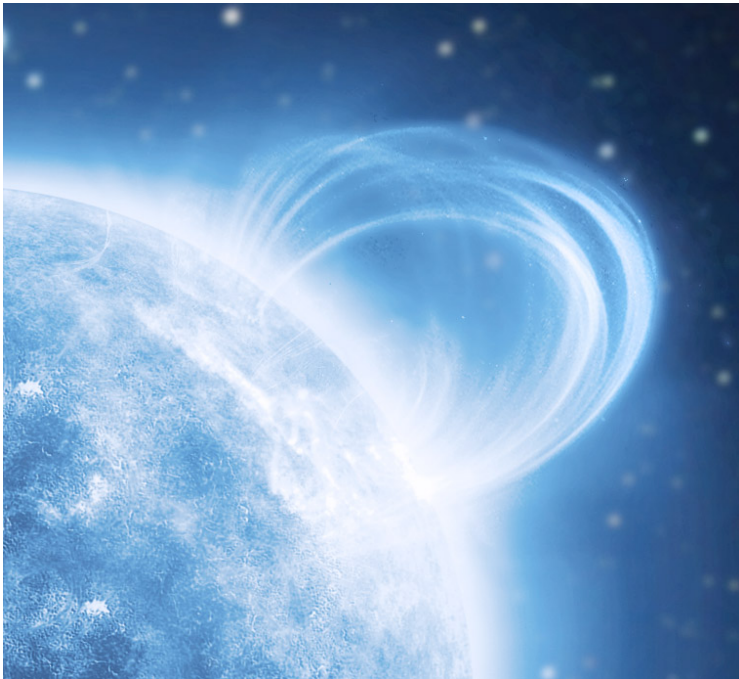
# Inferences

Most probably a cyclotron line (atomic transitions in H/He  
be'





First direct measure of the magnetic field at the surface of a magnetar,  $B \sim 2 \times 10^{14} - 10^{15} \text{ G}$



An artist impression of SGR 0418 with the ejected magnetic loop



Of course, reality is a trifle more complicated...

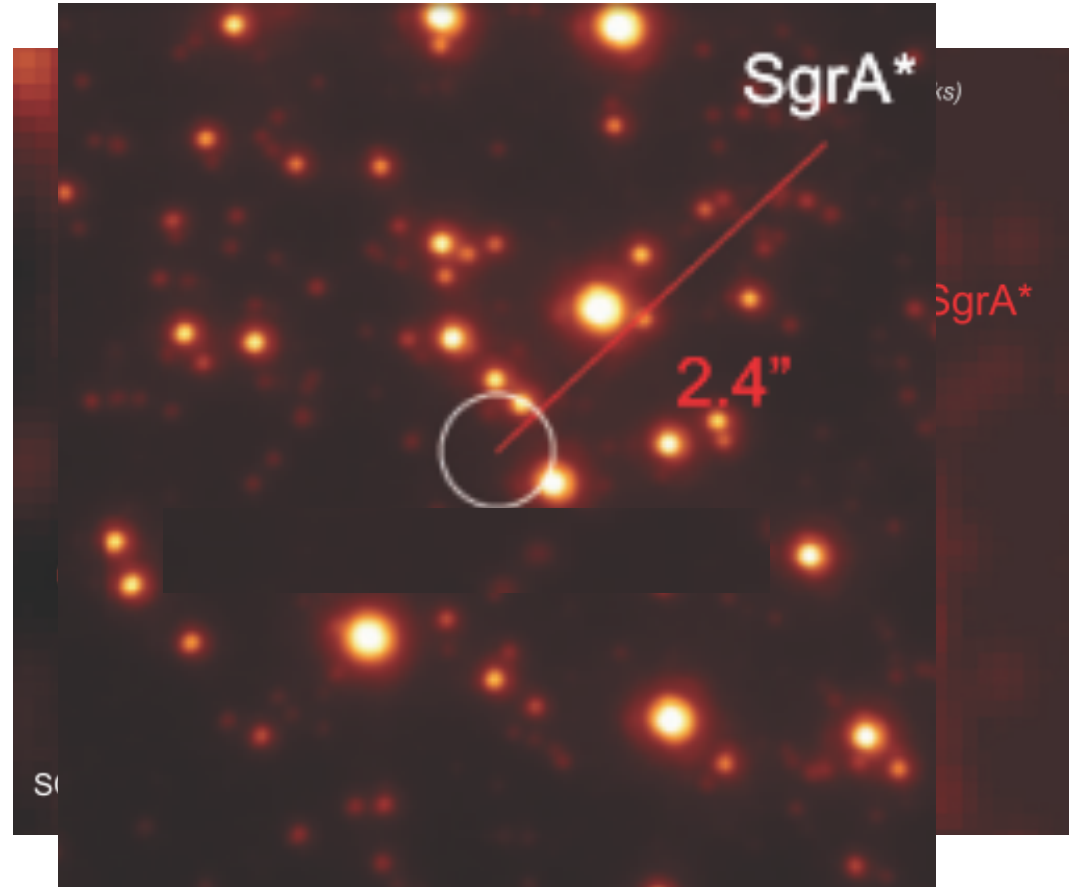


# A Space Oddity

The Galactic Center hosts the supermassive black hole Sgr A\* and a cluster of young, massive stars

Compact stellar remnants expected. Less expected was the presence of a magnetar at only 2.4" from Sgr A\*

SGR J1745-2900 was discovered on April 24, 2013 when it entered an outburst phase (Mori et al. 2013; Rea et al. 2013)



VLT/NaCo  $K_s$  image of the GC region



# SGR J1745-2900

- Full-fledged magnetar:  $P = 3.76 \text{ s}$ ,  $\dot{P} = 6.61 \times 10^{-12} \text{ s/s} \rightarrow B_p \sim 1.6 \times 10^{14} \text{ G}$ ,  $\tau_c \sim 9 \text{ kyr}$
- Probability that it is a foreground/ background object  $\sim 10^{-6}$
- Estimated distance to Sgr A\*  $\sim 0.1\text{-}2 \text{ pc}$
- Very likely ( $\sim 90\%$ ) in a bound orbit around Sgr A\*



# Future Developments

- Further support to the magnetar model: search for cyclotron line in other sources (preliminary results promising)
- Twisted magnetosphere model in general agreement with observations; hard vs soft power-laws
  - Nustar (and ASTRO-H)
  - More detailed modeling of magnetospheric currents
- Magnetar emission polarized: polarization measures key
- Only a general picture for the burst emission: need a quantitative model to compare with observations
- Extragalactic magnetars: relation with (long) GRBs ?
- The neutron star zoo: evolutionary links among different classes

