

# Status of the CKM matrix

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# Why precision CKM studies?

- The SM accomodates flavour & CP violation, but **we have no theory of flavour**
- We expect New Physics at the EW scale, and most models predict additional flavour and CP violation.
- The CKM mechanism is very successful  $\Rightarrow$  **flavour and CP problem** (NP must preserve agreement with data)
- To uncover small signals of physics beyond CKM, we need precision tests, in many ways a challenge for our QCD understanding

# The CKM matrix

Weak and mass eigenstates

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix} = \hat{V}_{\text{CKM}} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

$$\hat{V} = \begin{pmatrix} 1 - \frac{\lambda^2}{2} & \lambda & A\lambda^3(\varrho - i\eta) \\ -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1 - \varrho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$

Wolfenstein parameterization  $\lambda \sim 0.22$ ,  $A$ ,  $\varrho$ ,  $\eta$  are  $\mathcal{O}(1)$

To improve the accuracy, define to all orders in  $\lambda$

# The Cabibbo angle

$$\hat{V} = \begin{pmatrix} 1 - \frac{\lambda^2}{2} & \lambda & A\lambda^3(\varrho - i\eta) \\ -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1 - \varrho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$

$$\lambda = \sin(\theta_{\text{Cabibbo}}) = V_{us}$$

*Universality of charged currents  $\Leftrightarrow$  CKM unitarity*

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$$

$\swarrow O(10^{-5})$

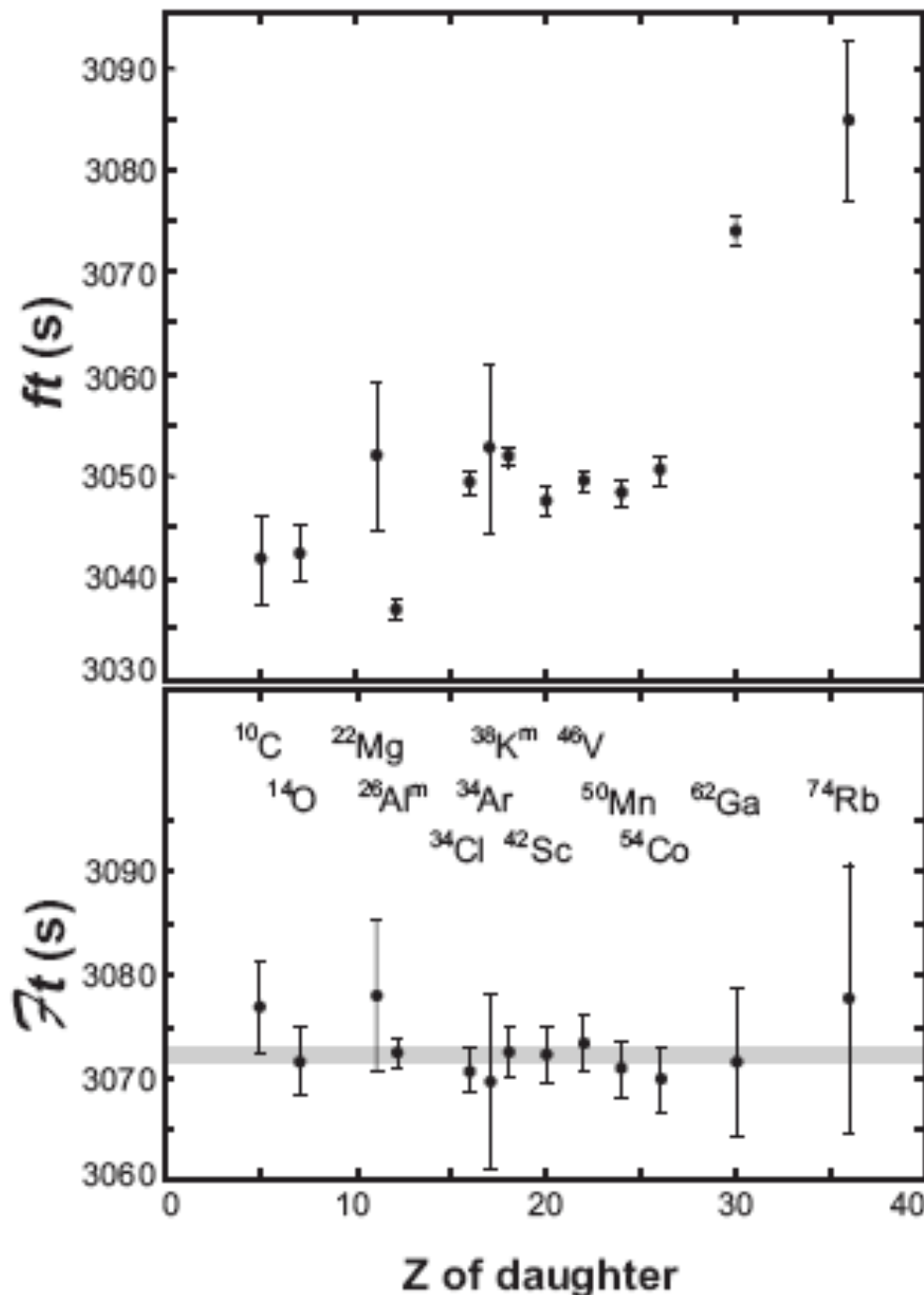
*Comparison between  $V_{ud}, V_{us}$  determinations of  $\lambda$  tests unitarity of the first line of  $V_{\text{CKM}}$*

*$\lambda$  could also be measured from 2nd line,  $V_{cd}$  (DIS) at 10%,  
 $W$  decays at LEP constrains  $\sum_{ij}|V_{ij}|^2$  at 1.3%  $V_{cs}$  at 1.3%*

# $\lambda$ from $V_{ud}$ : Fermi transitions

Superallowed Fermi transitions ( $0^+ \rightarrow 0^+$   $\beta$  decay)

$$\langle p_f; 0^+ | \bar{u} \gamma_\mu d | p_i; 0^+ \rangle = \sqrt{2} (p_i + p_f)_\mu$$



$$\mathcal{F}t \equiv ft(1 + \delta'_R)(1 + \delta_{NS} - \delta_C) \quad \text{isospin breaking, nuclear structure}$$

$$V_{ud}^2 = \frac{K}{2G_F^2(1 + \Delta_R^V)\overline{\mathcal{F}t}} \quad \Delta_R^V = (2.361 \pm 0.038)\%$$

Marciano Sirlin 2006

$$|V_{ud}| = 0.97425 \pm 0.00022 \quad \text{Hardy Towner 2008}$$

$$\Rightarrow \lambda = 0.2254(10) \text{ using unitarity}$$

Dominant error from structure indep  
RC, next structure dep ones.  
Great exp advances (Penning traps etc)

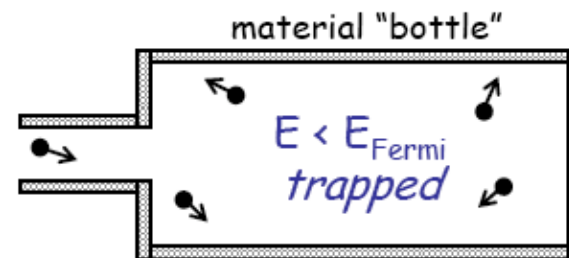
# Other $V_{ud}$ determinations

**neutron  $\beta$  decay** not pure vector, needs  $g_A/g_V$  but no nuclear structure.  $\delta V_{ud} \sim 0.002$ , will be improved through asymmetry measurements at PERKEO, Heidelberg and UCNA, LANL. 2005 measurement of  $n$  lifetime ( $6\sigma$  away) serious problem!

$$V_{ud} = 0.9746(4)_{\tau_n}(18)_{g_A}(2)_{RC}$$



Ultracold  
neutrons

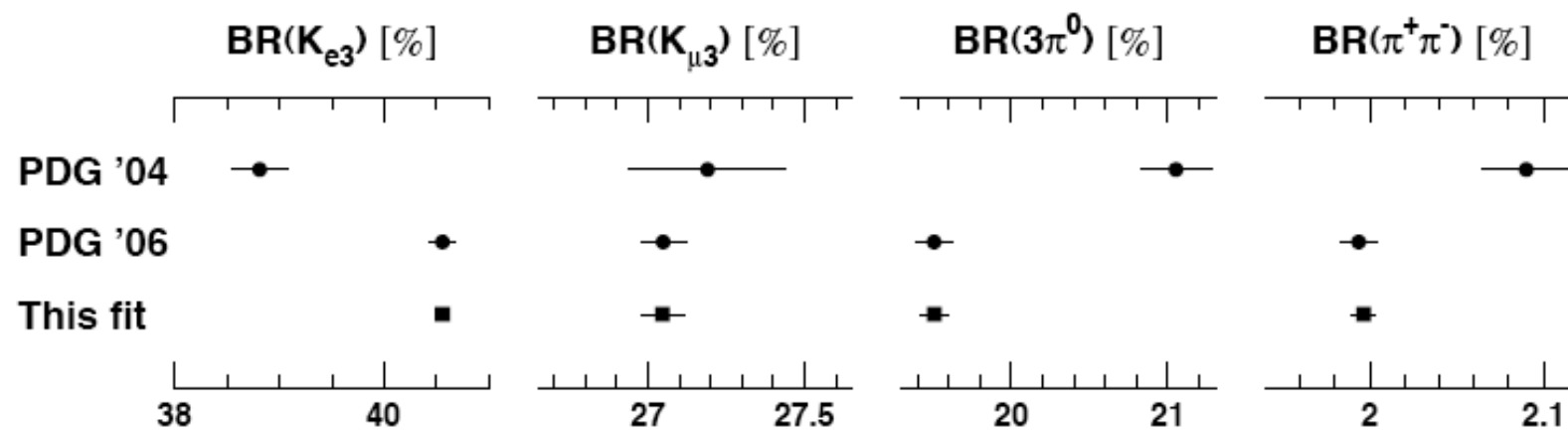


Long interaction  
times in apparatus  
↓  
Need relatively small  
number of neutrons

**$\pi^+$  decay to  $\pi^0 e \nu$**  th cleanest, promising in long term but  $BR \sim 10^{-8}$  PIBETA at PSI has  $\delta V_{ud} \sim 0.003$

$$V_{ud} = 0.9749(26) \left[ \frac{BR(\pi^+ \rightarrow e^+ \nu_e (\gamma))}{1.2352 \times 10^{-4}} \right]^{\frac{1}{2}}$$

# $\lambda$ from $K_{l3}$ - Experimental progress

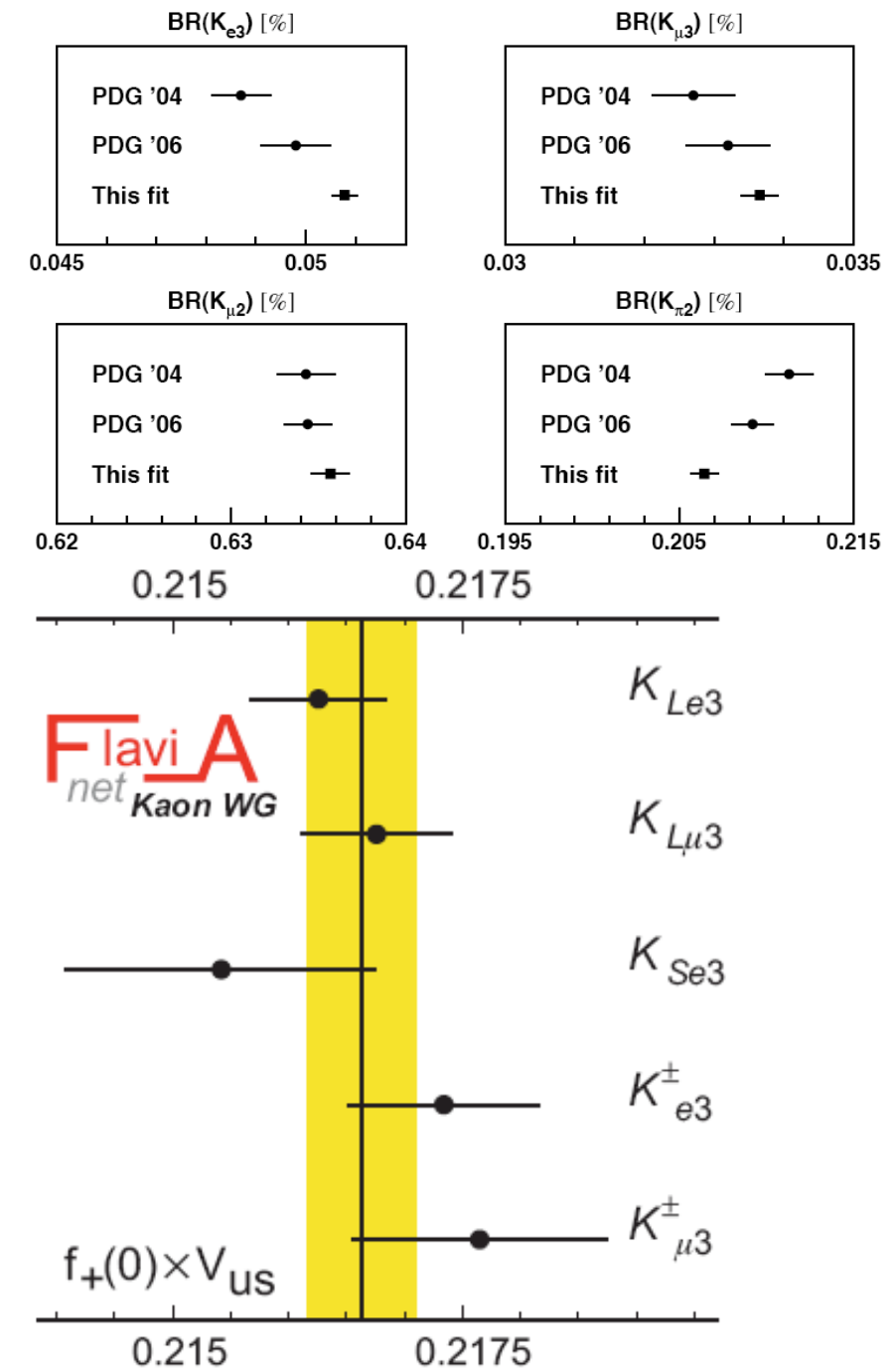


FlaviAnet Kaon WG 2008

$$\Gamma_{K\ell 3} = \frac{G_F^2 M_K^5}{192\pi^3} S_{EW} (1 + \delta_K^\ell + \delta_{SU2}) C^2 |V_{us}|^2 f_+^2(0) I_K^\ell.$$

$$f_+(0)|V_{us}| = 0.21673 \pm 0.00046$$

*0.25% accuracy!*  
*muon channels perfectly consistent*



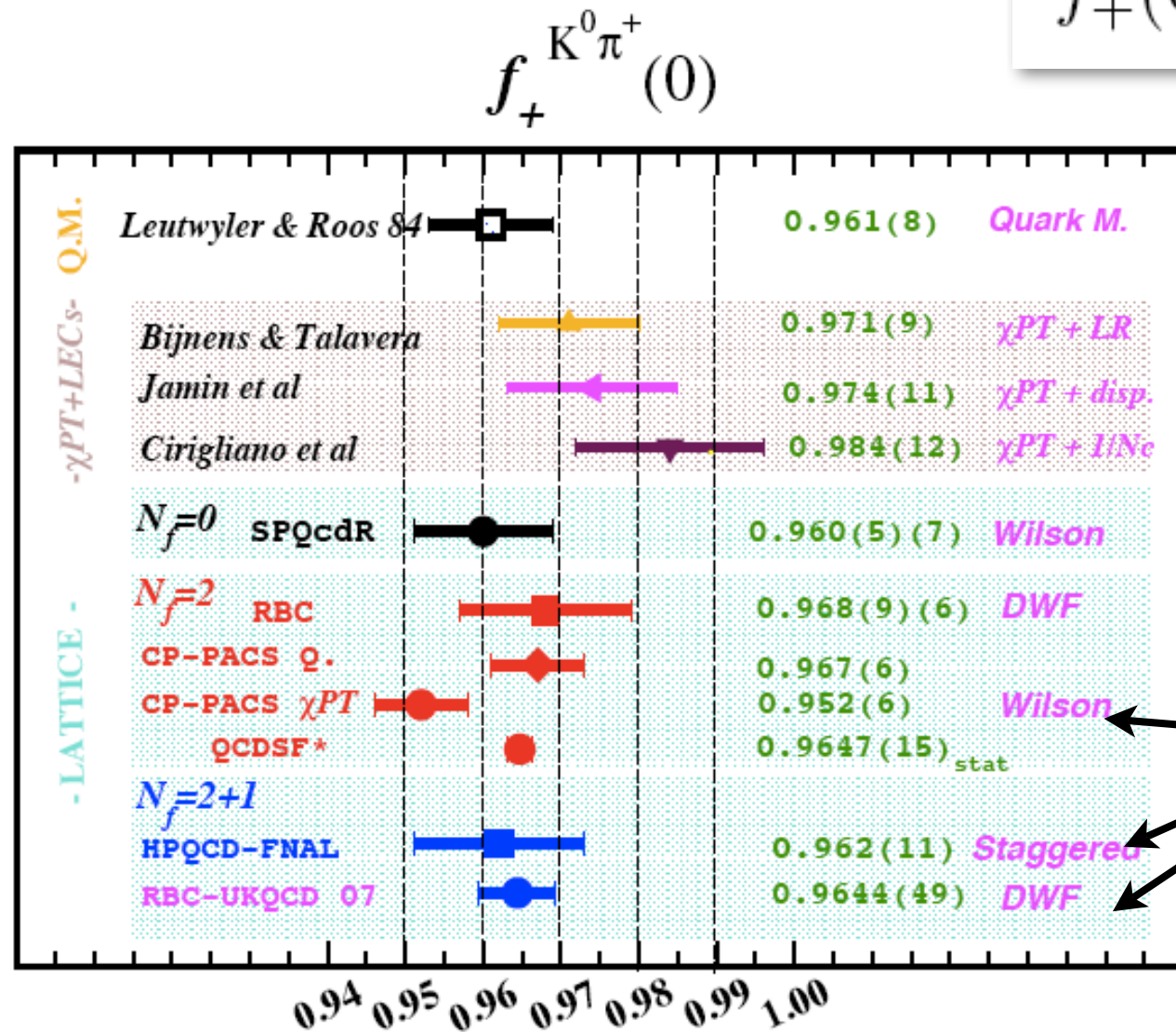


# $\lambda$ from $K_{l3}$ - Theoretical progress

$$f_+(0) = 1 + f_2 + f_4 + \dots$$

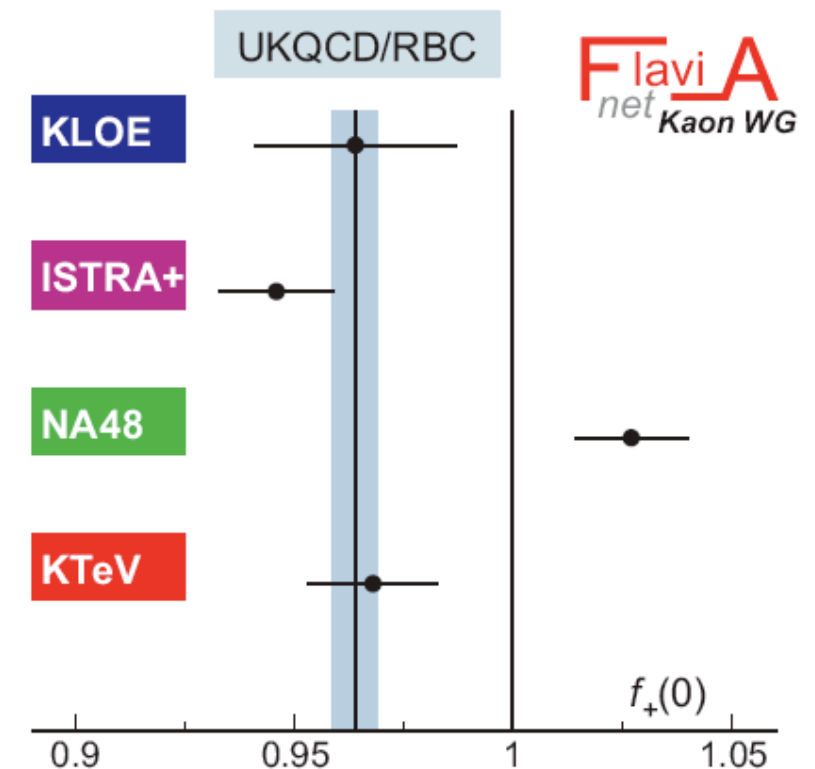
$SU(3)$  symmetry,  
Ademollo Gatto

$$= 0.964(5) \quad \text{RBC-UKQCD 07}$$



Tests of lattice are now possible from measurements of the shapes, eg from Callan-Treiman and  $f_K/f_\pi$

Various lattice actions



$$|V_{us}| = 0.2246 \pm 0.0012 \quad [K_{\ell 3} \text{ only}]$$



# $\lambda$ from $K_{l2}$

$$\frac{\Gamma(K_{\ell 2(\gamma)}^{\pm})}{\Gamma(\pi_{\ell 2(\gamma)}^{\pm})} = \left| \frac{V_{us}}{V_{ud}} \right|^2 \frac{f_K^2 m_K}{f_\pi^2 m_\pi} \left( \frac{1 - m_\ell^2/m_K^2}{1 - m_\ell^2/m_\pi^2} \right)^2 \times (1 + \delta_{\text{em}})$$

Marciano 2004

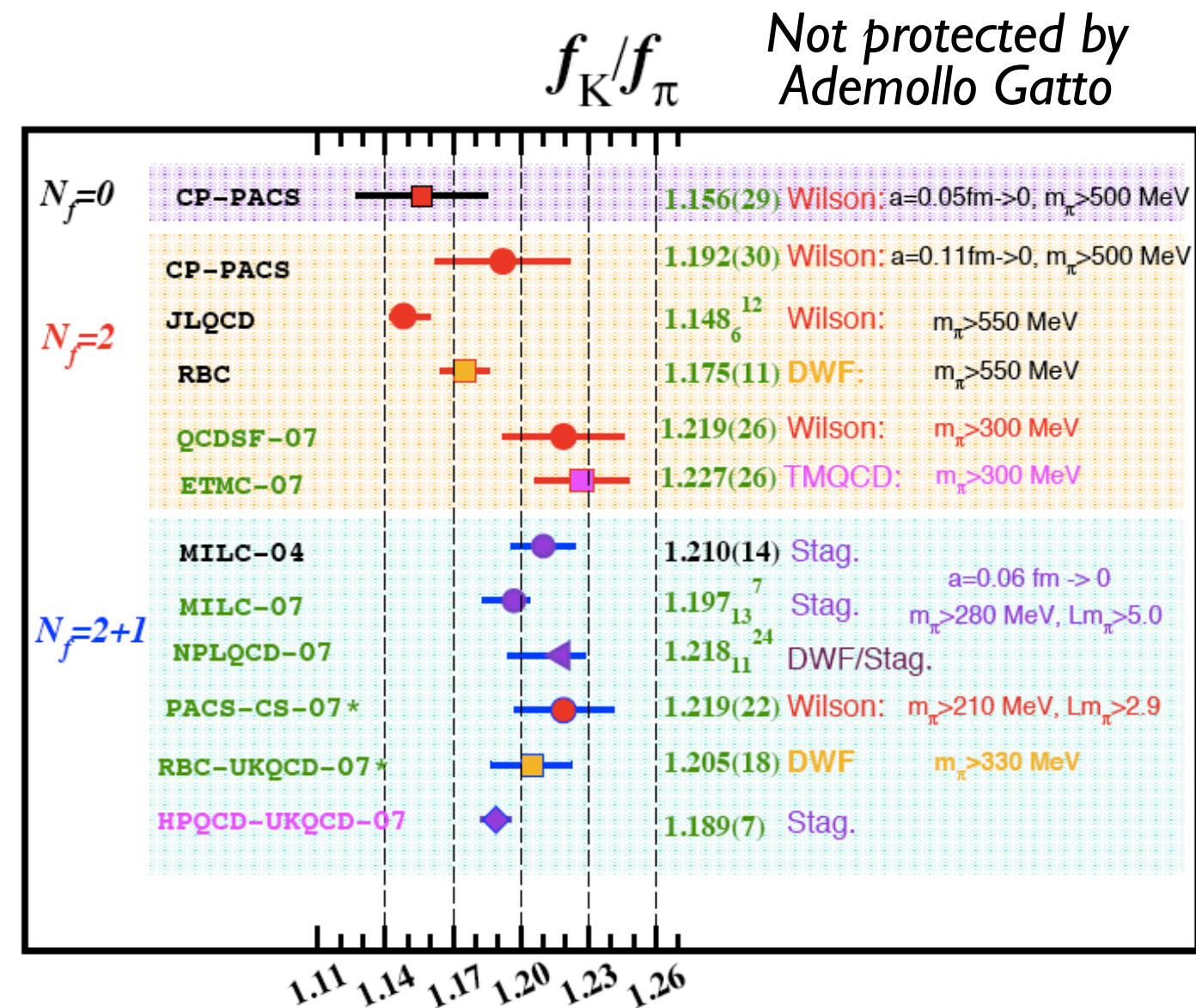
Kloe, NA48/2 find

$$R_K = (2.457 \pm 0.032) \times 10^{-5},$$

$$\left| \frac{V_{us}}{V_{ud}} \right| = 0.2321 \pm 0.0015 \quad \text{Only } K_{l2}$$

Very similar from  $\Gamma(\tau^- \rightarrow K^- \nu)/\Gamma(\tau^- \rightarrow \pi^- \nu)$   
measured at Babar:  $\left| \frac{V_{us}}{V_{ud}} \right| = 0.2315 \pm 0.0024$

Cross-check of  $K_{l2}$ !



# Unitarity of the first row

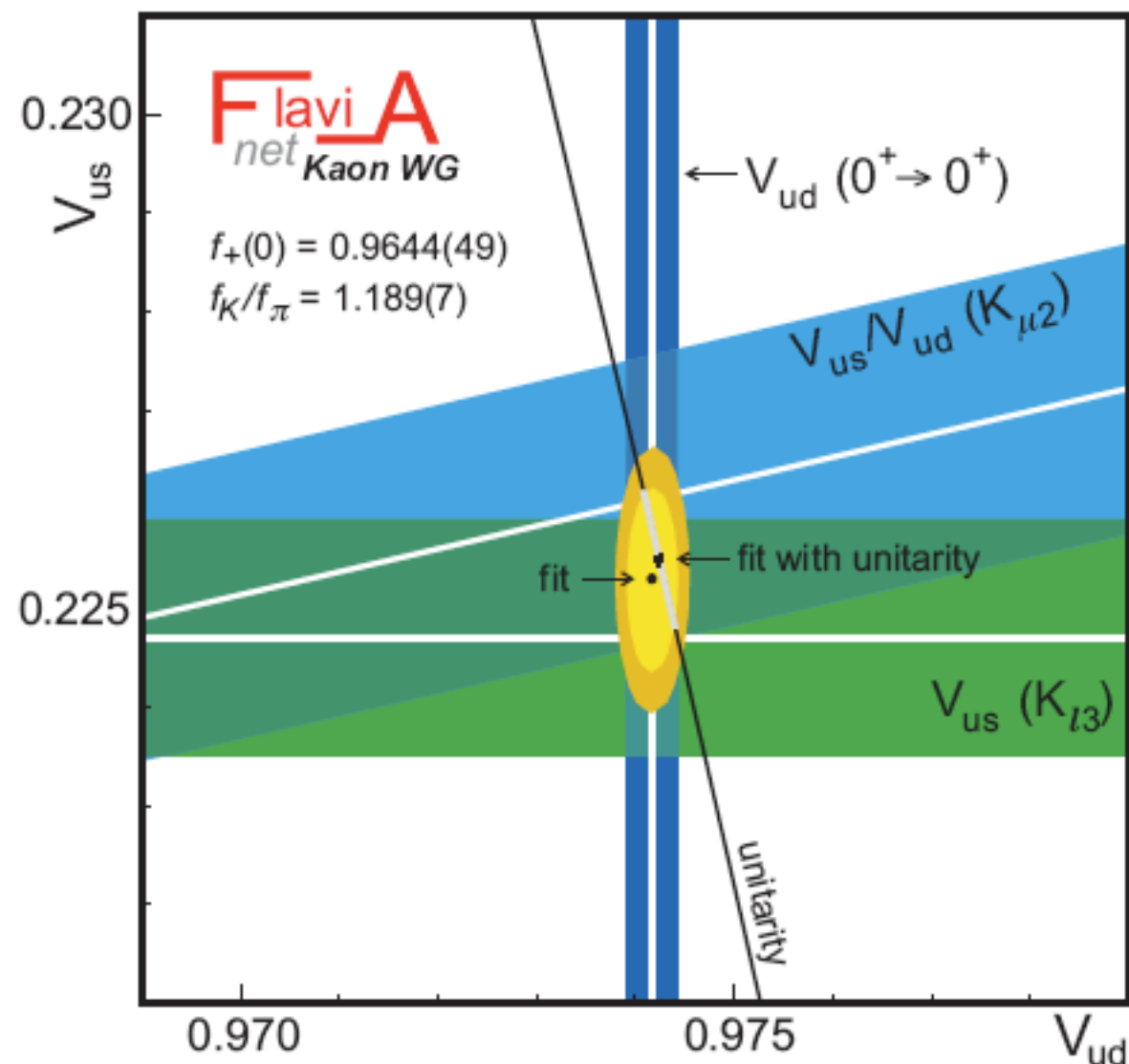
$$|V_{us}| = \sin \theta_C = \lambda = 0.2255(7) \quad [\text{with unitarity}]$$

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = \\ = 0.9999 \pm 0.0006$$

Inclusive Tau decays also give  $\lambda$  with  $\sim 1\%$  error but need  $m_s(m_\tau)$ .

Preliminary Belle and Babar data suggest  $0.2165(27)$  but there are some doubts on experimental analyses (missing modes)

Gamiz et al 2007, ...

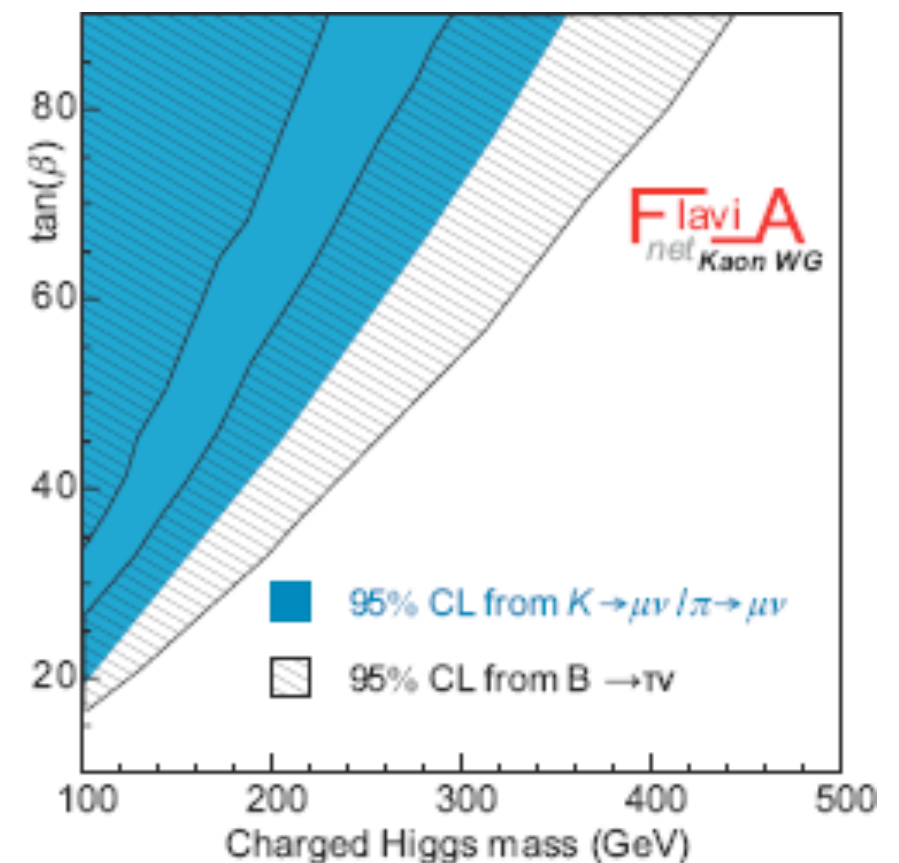


# Constraints on New Physics

- Bounds on **non-universality** :  $G_F^{CKM} = 1.1663(4) 10^{-5} \text{ GeV}^{-2}$   
 $G_\mu = 1.166371(6) 10^{-5} \text{ GeV}^{-2}$   
 $G_F^{EW} = 1.1656(11) 10^{-5} \text{ GeV}^{-2}$

*Can be tree-level (mixing with heavy quark)  
or loop induced (squarks vs sleptons)*

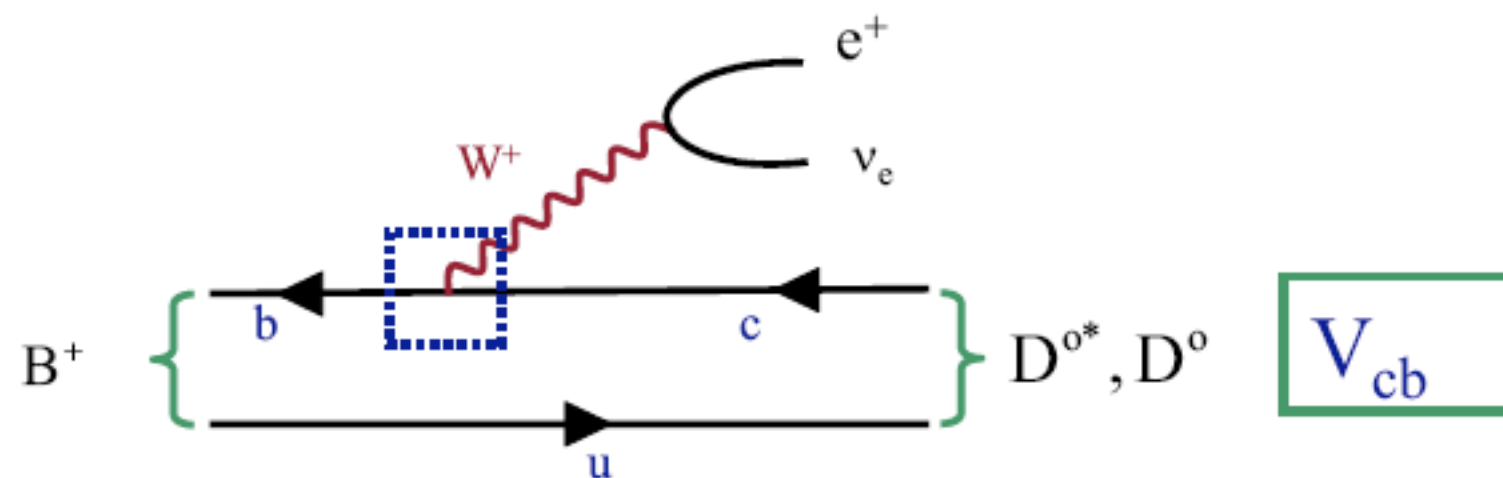
- Bounds on scalar currents, eg charged Higgs interactions in  $K_{l2}$



# Determination of A

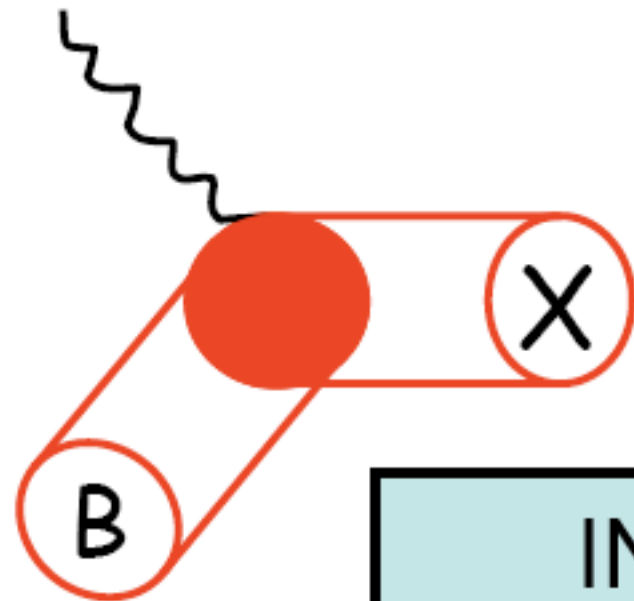
$$\hat{V} = \begin{pmatrix} 1 - \frac{\lambda^2}{2} & \lambda & A\lambda^3(\varrho - i\eta) \\ -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1 - \varrho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$

A can be determined from either  $V_{cb}$  or  $V_{ts}$



Two roads to  $V_{cb}$ : inclusive and exclusive

# Inclusive vs exclusive B decays



Simplicity: ew or em currents probe the B dynamics

INCLUSIVE	EXCLUSIVE
<b>OPE</b> : non-pert physics described by B matrix elements of local operators can be extracted by exp suppressed by $1/m_b^2$	<b>Form factors</b> : in general computed by non pert methods (lattice, sum rules,...) symmetry can provide normalization

As precision increases, simplicity evaporates...



# Exclusive decays: $B \rightarrow D^* l \nu$

At zero recoil, where rate vanishes, the  $ff$  is

$$\mathcal{F}(1) = \eta_A(1 + \delta_{1/m^2})$$

Recent progress in the measurement of slopes and shape parameters *Despite extrapolation, exp error ~2%*

Main problem is normalization  $F(1)$ : **non-pert quantities relevant for excl decays cannot be experimentally determined**

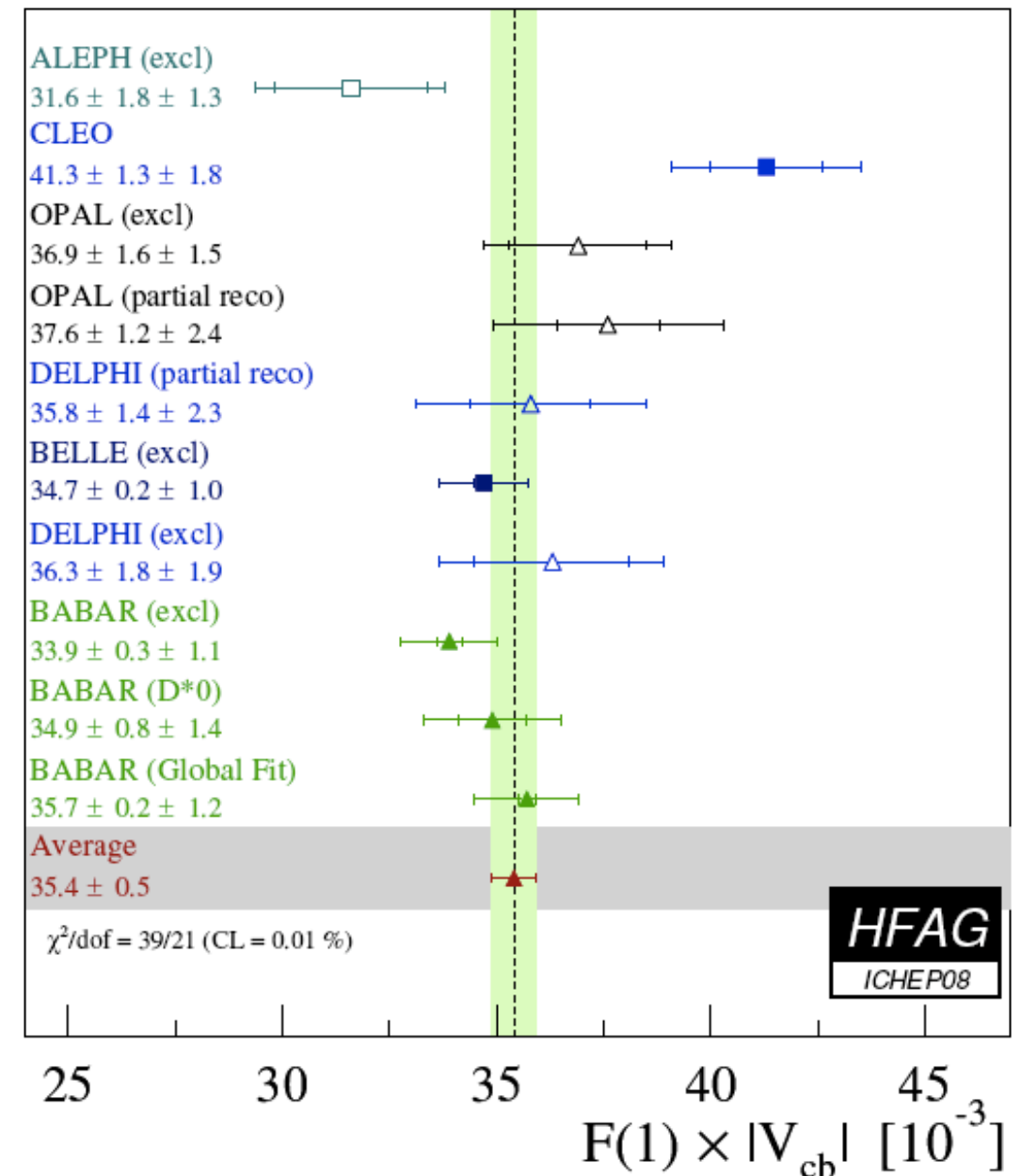
New and **only** unquenched Lattice QCD:

$F(1) = 0.921(24)$  Laiho et al 2008, HQET, double ratio

$$|V_{cb}| = 38.2(0.5)(1.1) \times 10^{-3}$$

$\sim 2.4\sigma$  from inclusive determination

**NB** Heavy Quark Sum rules give higher  $|V_{cb}|$ :  
 $F(1) = 0.87(4)$  Uraltsev in agreement with inclusive work in progress



Lattice promising alternative: step scaling, w dependence, only quenched de Divitiis et al

$B \rightarrow D l \nu$  gives consistent but much less precise results

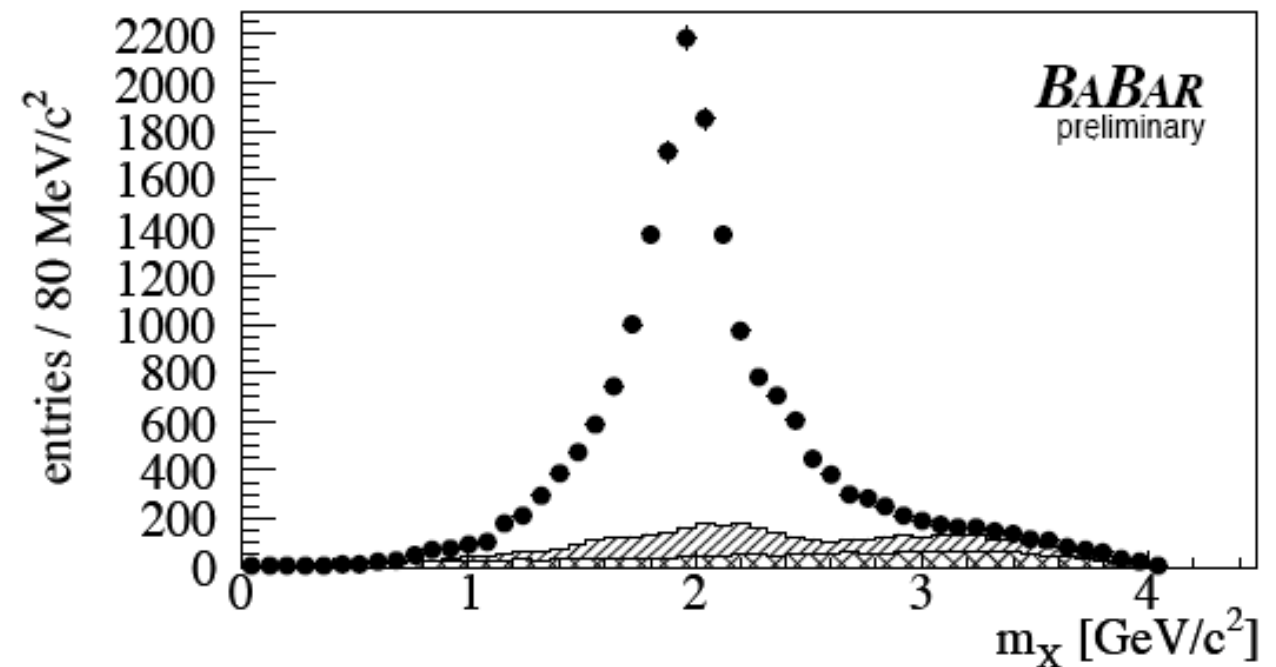
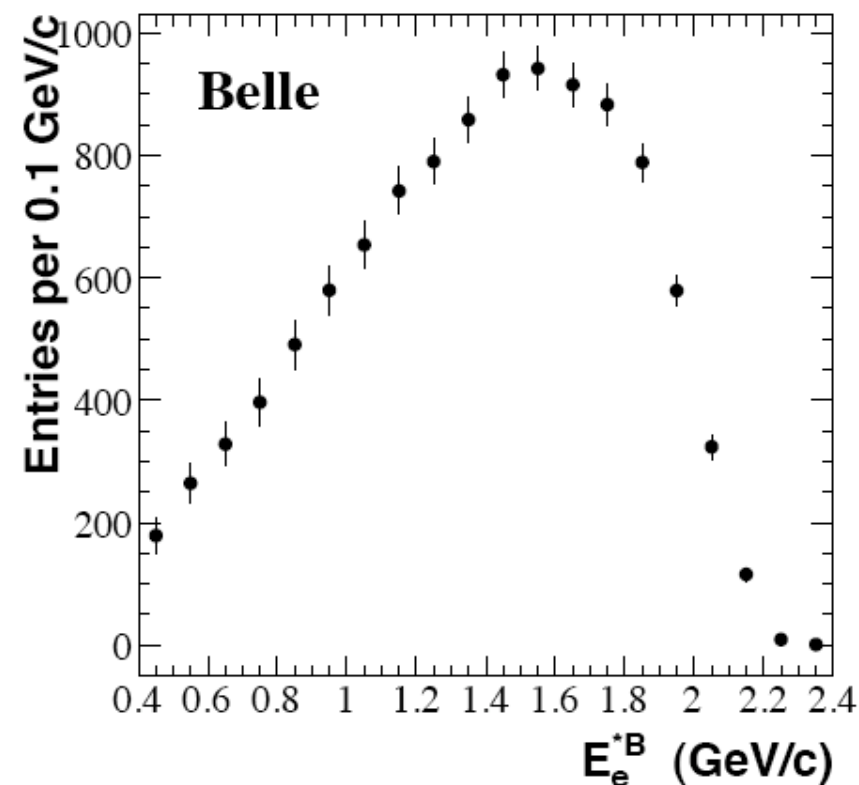


# Inclusive $|V_{cb}|$ : basic features

- **Simple idea:** inclusive decay do not depend on final state, factorize long distance dynamics of the meson. OPE allows to express it in terms of matrix elements of local operators
- The Wilson coefficients are perturbative, matrix elements of local ops parameterize non-pert physics: **double series in  $\alpha_s, \Lambda/m_b$**
- Lowest order: decay of a free  $b$ , linear  $\Lambda/m_b$  absent. Depends on  $m_{b,c}$ , 2 parameters at  $O(1/m_b^2)$ , 2 more at  $O(1/m_b^3)$ ...

$$\mu_\pi^2(\mu) = \frac{1}{2M_B} \left\langle B \left| \bar{b} (i\vec{D})^2 b \right| B \right\rangle_\mu \quad \mu_G^2(\mu) = \frac{1}{2M_B} \left\langle B \left| \bar{b} \frac{i}{2} \sigma_{\mu\nu} G^{\mu\nu} b \right| B \right\rangle_\mu$$

# Fitting OPE parameters to the moments



Total **rate** gives  $|V_{cb}|$ , global **shape** parameters (moments of the distributions) tell us about  $B$  structure

*OPE parameters describe universal properties of the  $B$  meson and of the quarks*

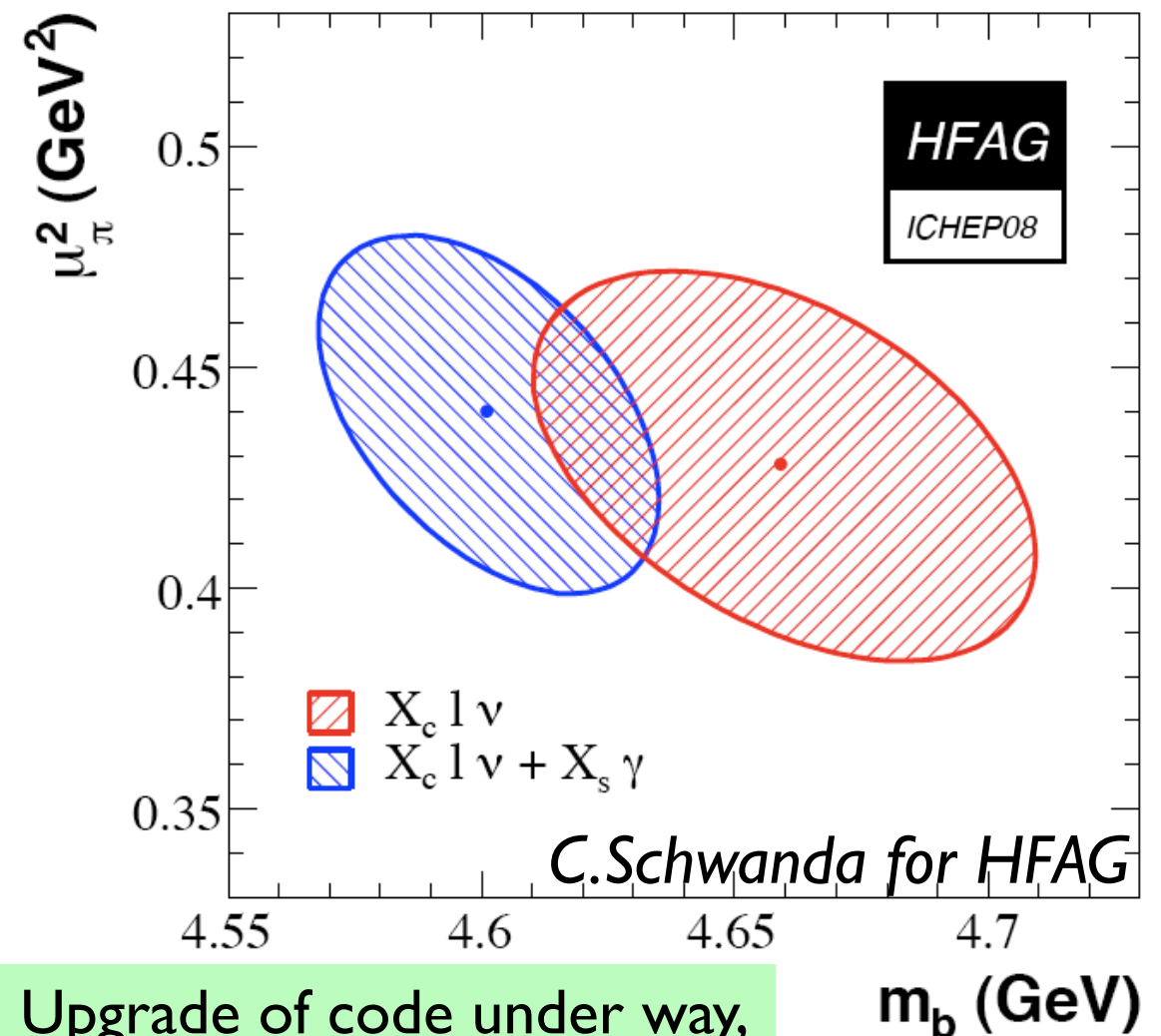
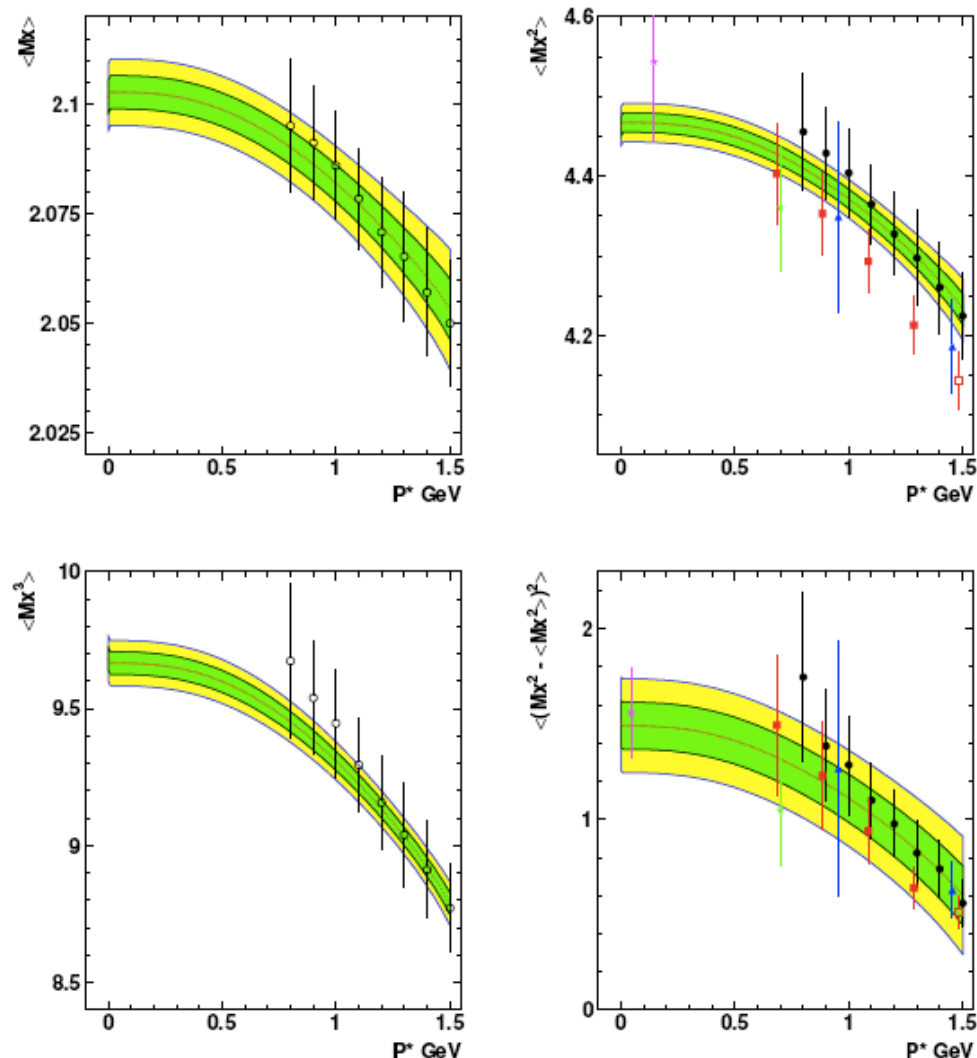
# Global fit (kinetic scheme)

Inputs	$ V_{cb}  \cdot 10^3$	$m_b^{\text{kin}}$	$\chi^2/\text{ndf}$
$b \rightarrow c$ & $b \rightarrow s\gamma$	41.67(44)(58)	4.601(34)	29.7/57
$b \rightarrow c$ only	41.48(48)(58)	4.659(49)	24.1/46

Based on PG, Uraltsev & Benson et al

**In the kinetic scheme** the contributions of gluons with energy below  $\mu \approx 1 \text{ GeV}$  are absorbed in the OPE parameters

Here scheme means also a number of different assumptions and a recipe for theory errors



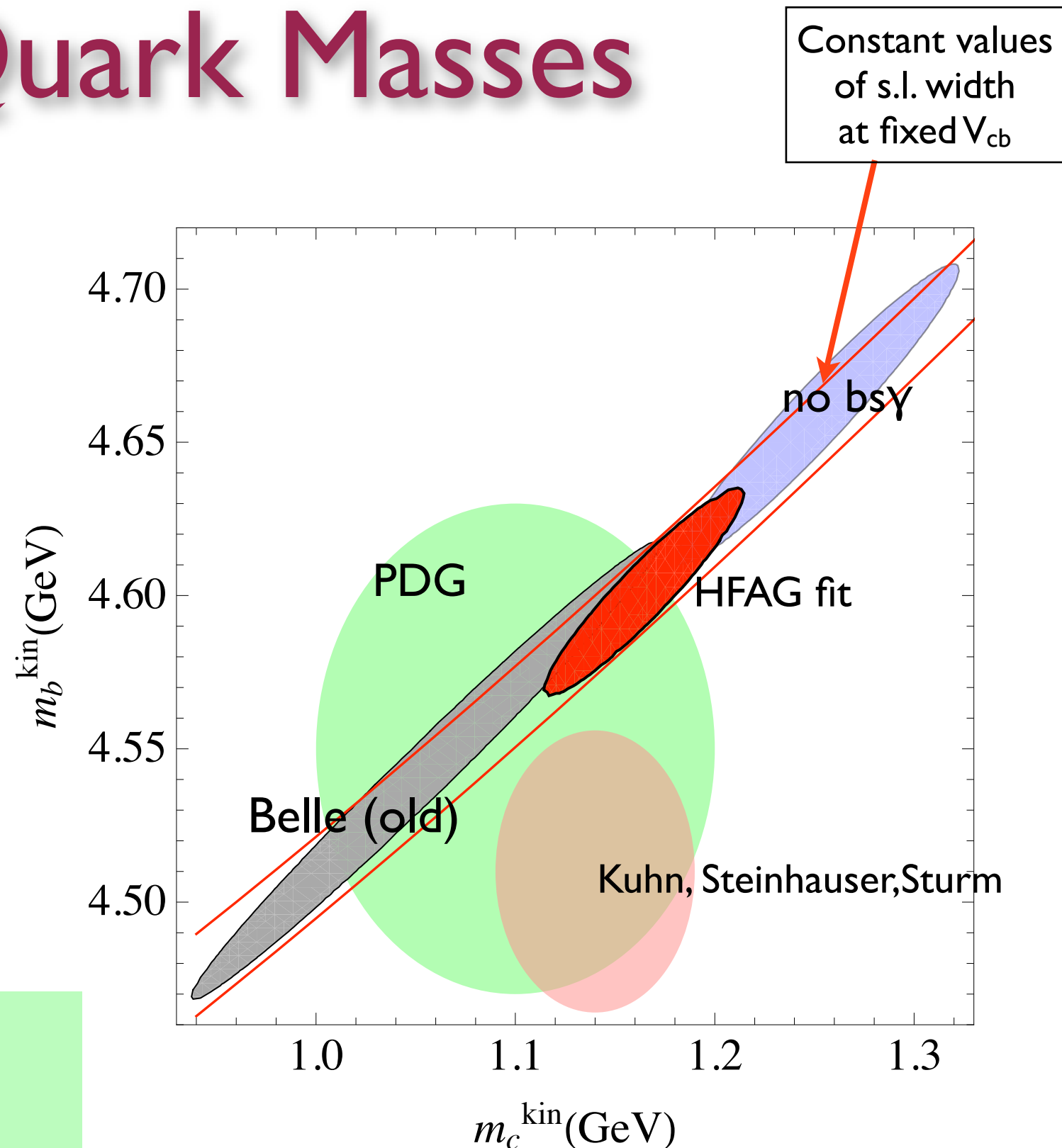
# Fits & Quark Masses

Assume quark-hadron duality but self-consistently check it

Semileptonic moments identify a strip in  $(m_b, m_c)$  plane along which the minimum is **shallow**.

Inclusion of radiative moments controversial as OPE fails at  $O(\alpha_s)$ . At present the role of radiative moments in the fits is similar to using PDG bound on  $m_b$ .

Incl.  $|V_{cb}|$  looks OK  
Heavy quark masses?



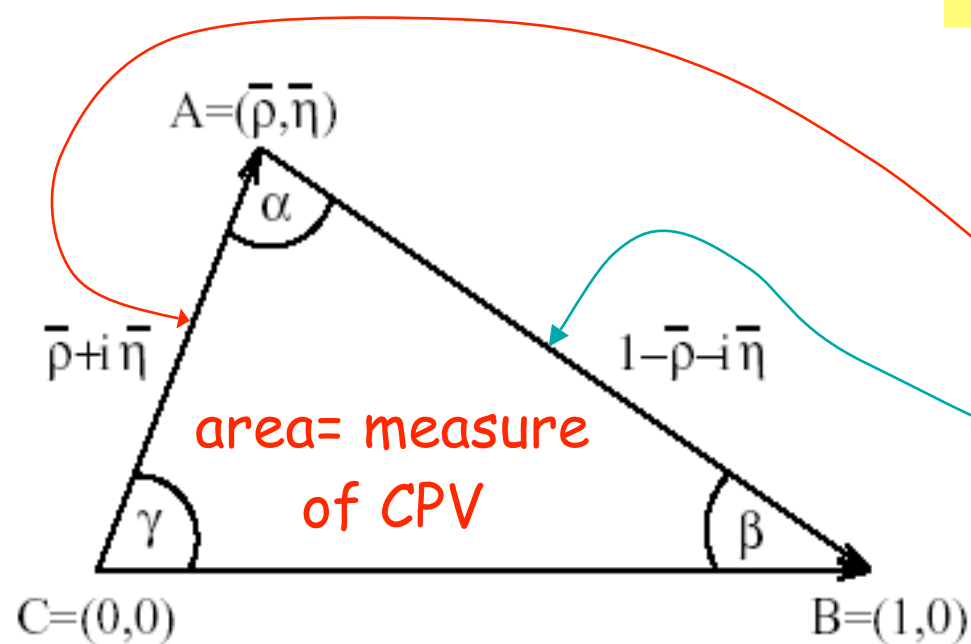
# The Unitarity Triangle

$$V_{ji} V_{jk}^* = \delta_{ik}$$

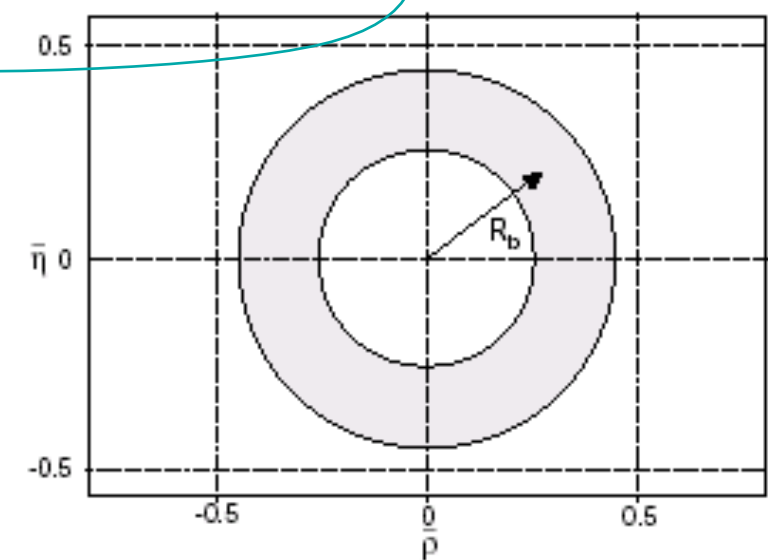
Unitarity determines several triangles in complex plane

$$V_{ud} V_{ub}^* + V_{cd} V_{cb}^* + V_{td} V_{tb}^* = 0 \quad \mathcal{O}(\lambda^3)$$

$$1 + \frac{V_{ud} V_{ub}^*}{V_{cd} V_{cb}^*} + \frac{V_{td} V_{tb}^*}{V_{cd} V_{cb}^*} = 0$$

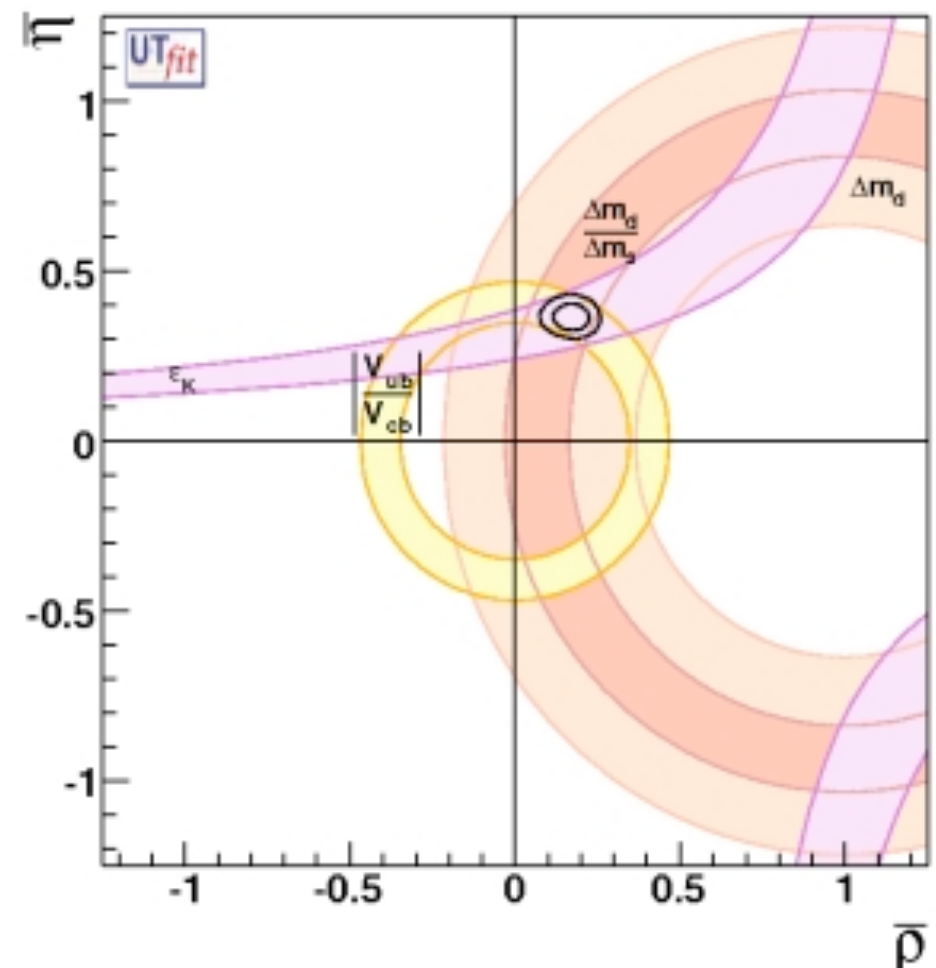
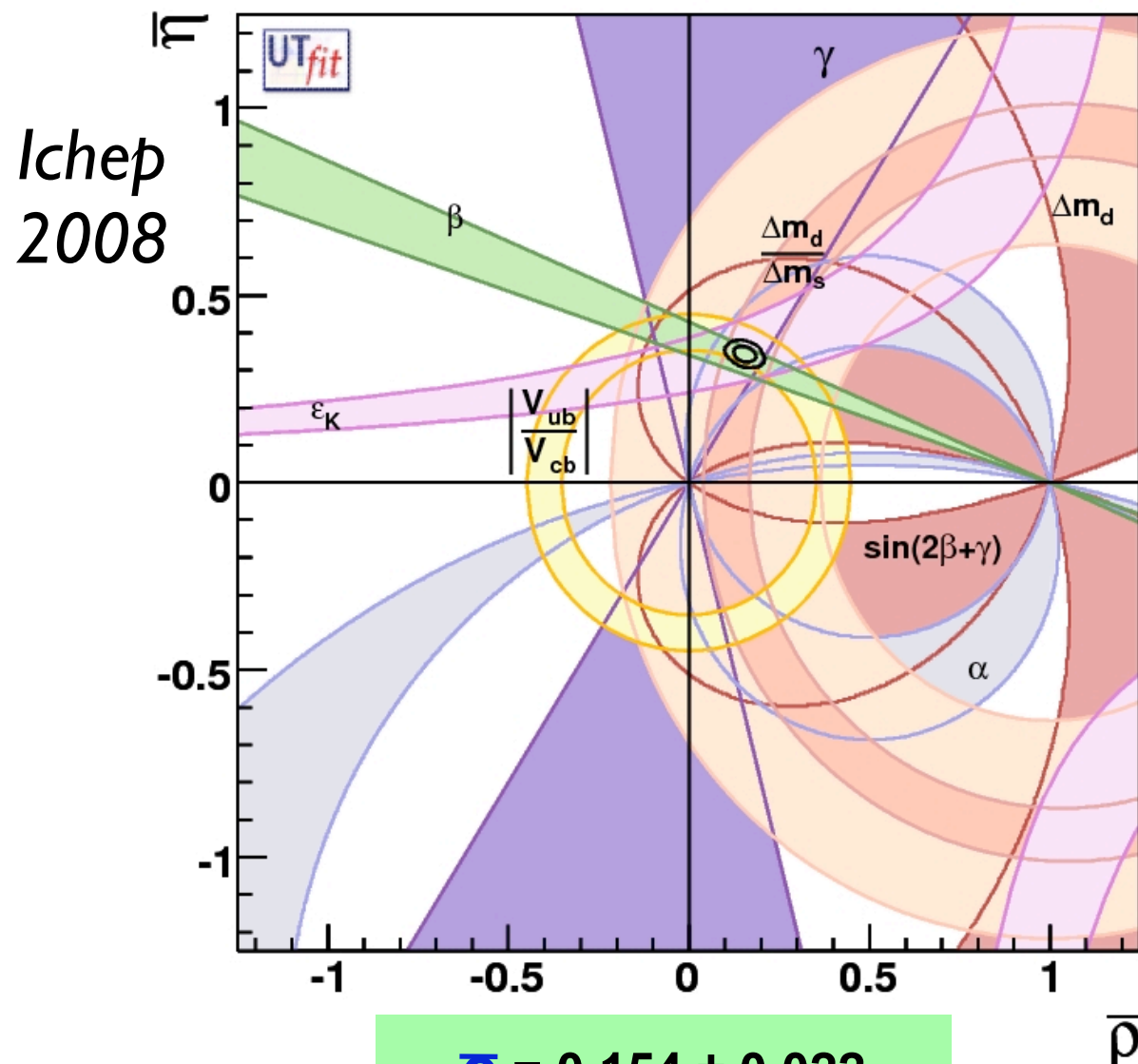


$V_{td}$  cannot be accessed directly:  
need FCNC loops sensitive to  
new physics eg  $B_d$ ,  $B_s$  mixing





# The Unitarity Triangle



Almost identical results by CKMfitter @ ICHEP 2008

$B_K = 0.725(50)$  Lellouch LAT08

$$\sin 2\beta_{\text{charmionium}} = 0.672 \pm 0.024 \text{ HFAG}$$

getting closer to 5% accuracy?



# Exclusive determination of $|V_{ub}|$

High exp accuracy. Various parameterizations based on analyticity etc + experimental data on the  $q^2$  spectrum: model independently

$$|V_{ub}|f_+(0) = (9.1 \pm 0.6(\text{shape}) \pm 0.3(\text{BR})) \times 10^{-4}$$

*ff on lattice or with LC sum rules, no symmetry helps. LCSR cannot be much improved, while lattice can*

LCSR:  $|V_{ub}| = (3.5 \pm 0.4 \pm 0.1) \times 10^{-4}$   
*Ball-Zwicky*

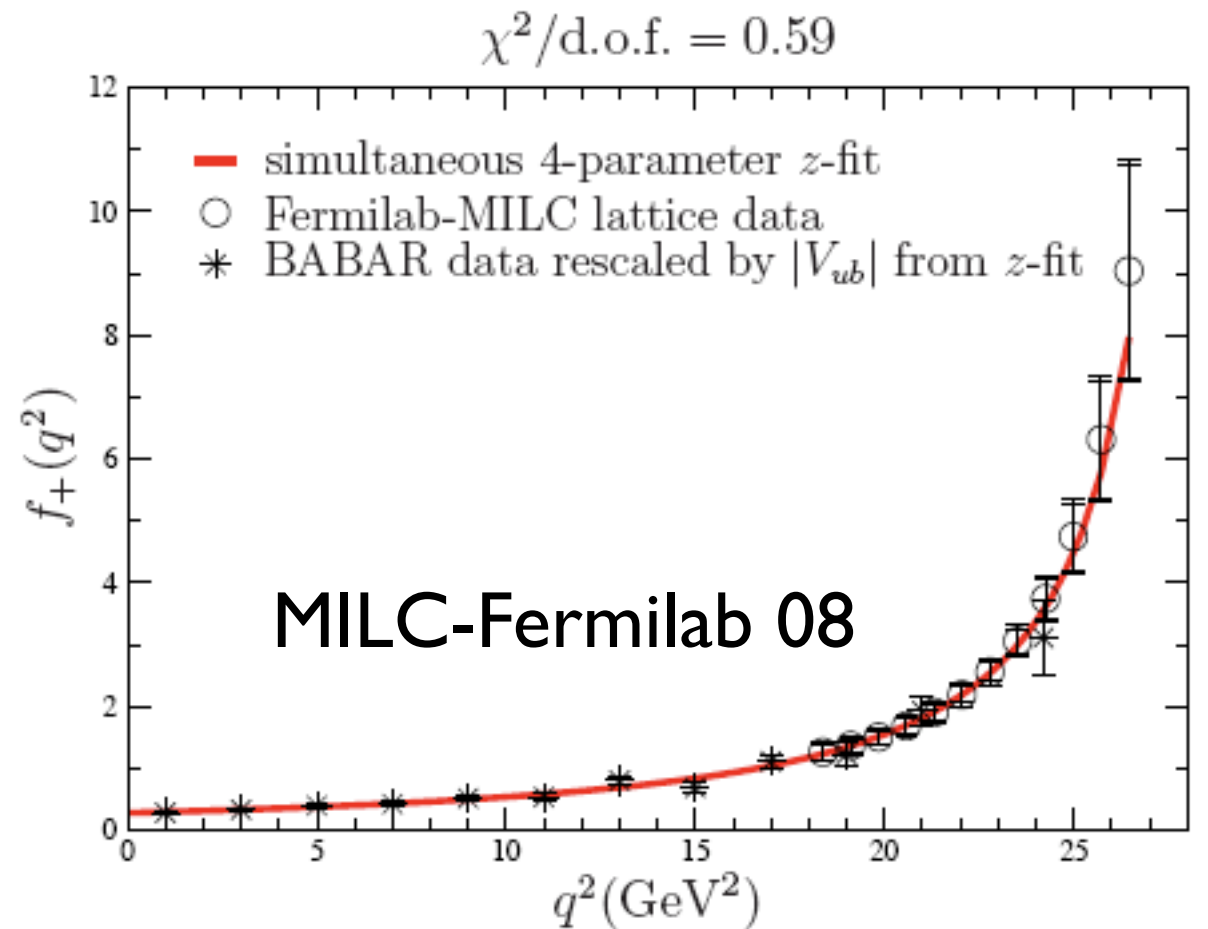
$$|V_{ub}| = \left( 3.5 \pm 0.4|_{th} \pm 0.2|_{shape} \pm 0.1|_{BR} \right) \times 10^{-3}$$

*Duplancic et al*

$$|V_{ub}| \times 10^3 = 3.36 \pm 0.23$$

*Bourelly et al 08*

*Combines older lattice results +LCSR with larger (!) errors and a new parameterization*



$$|V_{ub}| \times 10^3 = 3.38 \pm 0.36. \quad \text{MILC-Fermilab 08}$$

# $|V_{ub}|$ inclusive

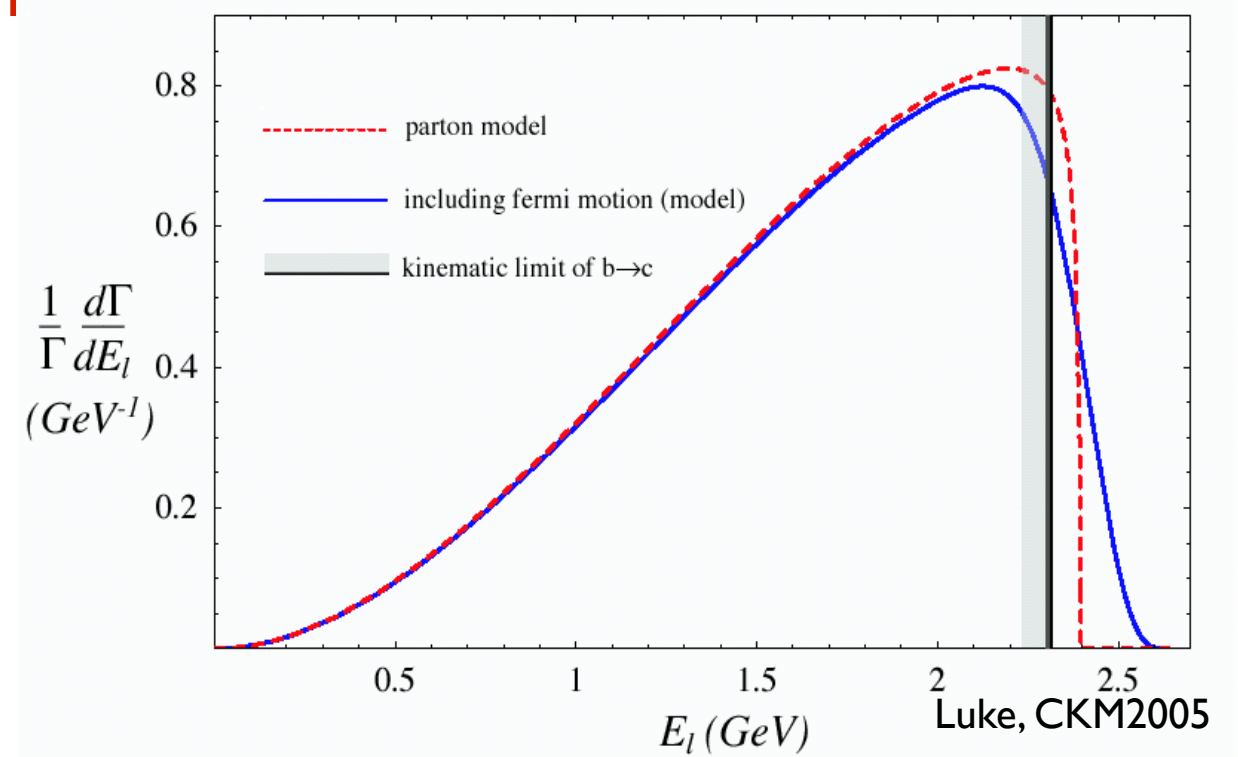
$|V_{ub}|$  from total  $\text{BR}(b \rightarrow u\ell\nu)$  like incl  $|V_{cb}|$  but we need kinematic cuts to avoid the  $\sim 100\times$  larger  $b \rightarrow c\ell\nu$  background:

$$m_X < M_D \quad E_\ell > (M_B^2 - M_D^2)/2M_B \quad q^2 > (M_B - M_D)^2 \dots$$

or combined  $(m_X, q^2)$  cuts

The cuts destroy convergence of the OPE that work so well in  $b \rightarrow c$ .  
OPE expected to work only away from  
pert singularities

Rate becomes sensitive to “local”  
b-quark wave function properties  
like Fermi motion Dominant non-  
pert contributions can be resummed  
into a **SHAPE FUNCTION**  $f(k_+)$



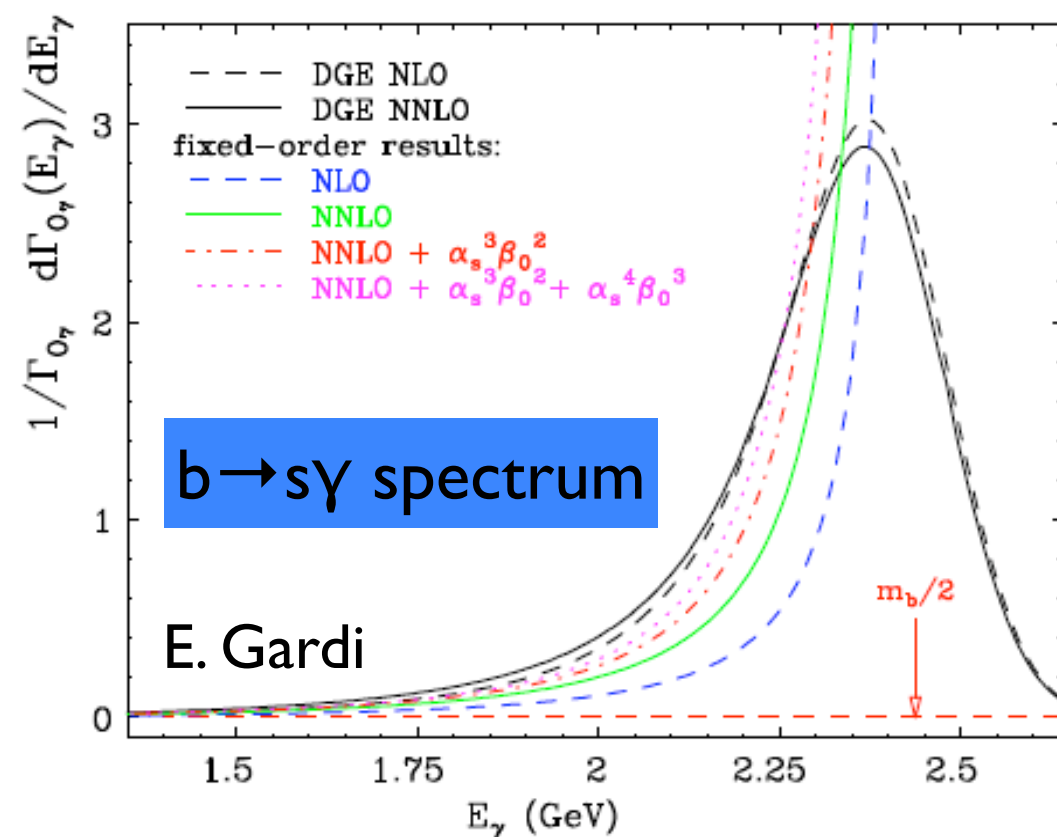
# SF from perturbation theory

Resummed perturbation theory is qualitatively different: **Support properties**; **stability!** (E. Gardi)

**b quark SF emerges from soft-gluon resummed pQCD but needs resummation of running coupling corrections and power corrections for  $b \rightarrow B$**

Dress Gluon Exponentiation (DGE) by Gardi et al employs renormalon resummation to define Fermi motion. Power corrections can be partly accommodated.

Aglietti et al (ADFR) use Analytic Coupling in the IR



# The SF in the OPE

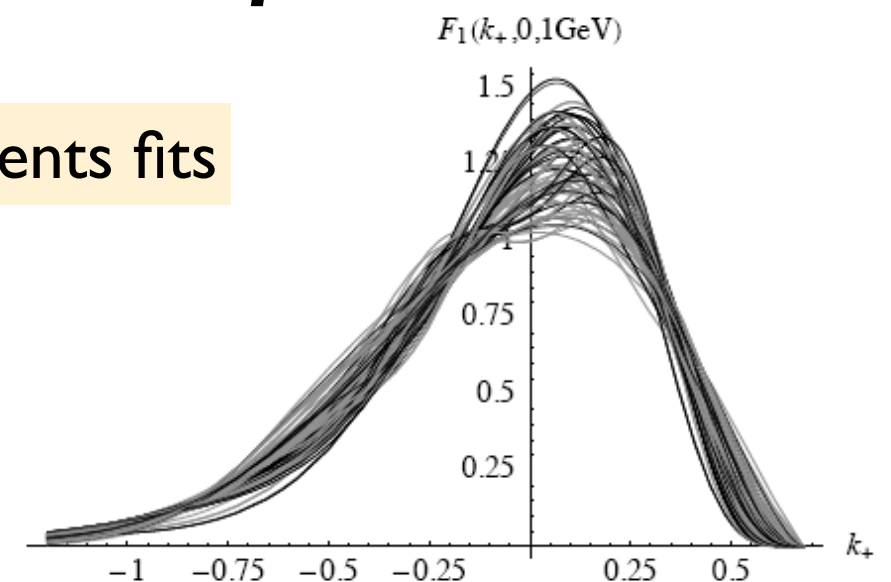
Local OPE has also threshold singularities and SF can be equivalently introduced resumming dominant singularities *Bigi et al, Neubert*

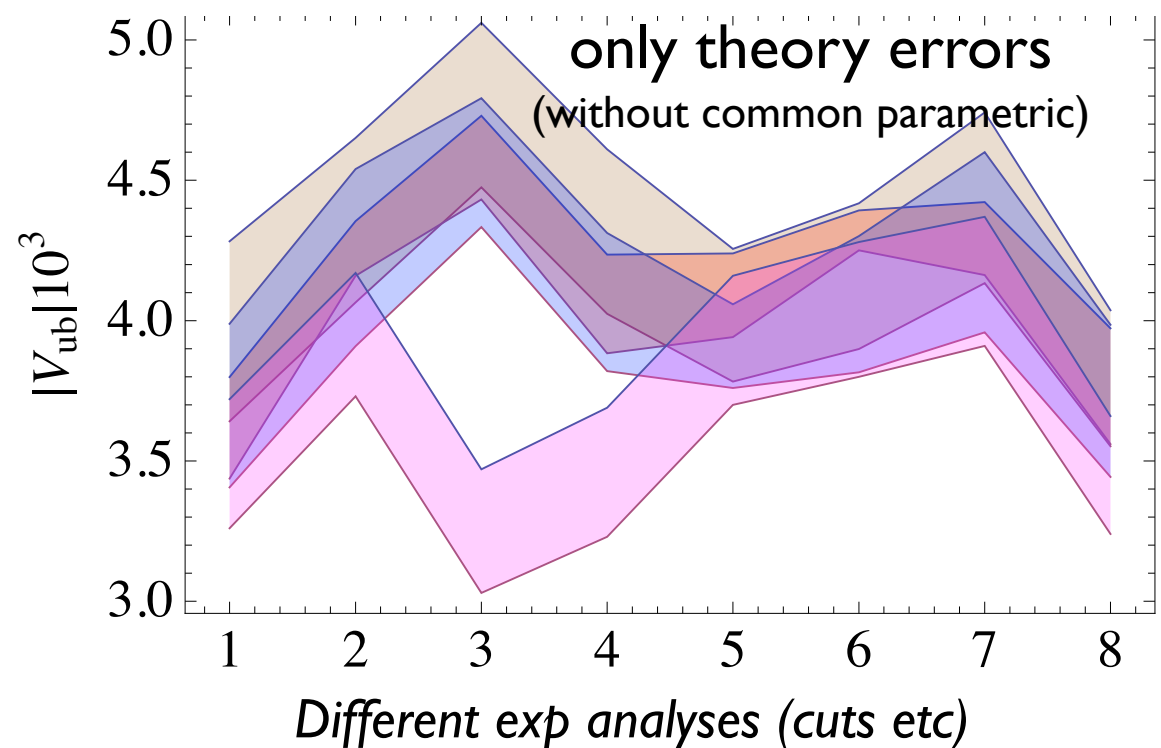
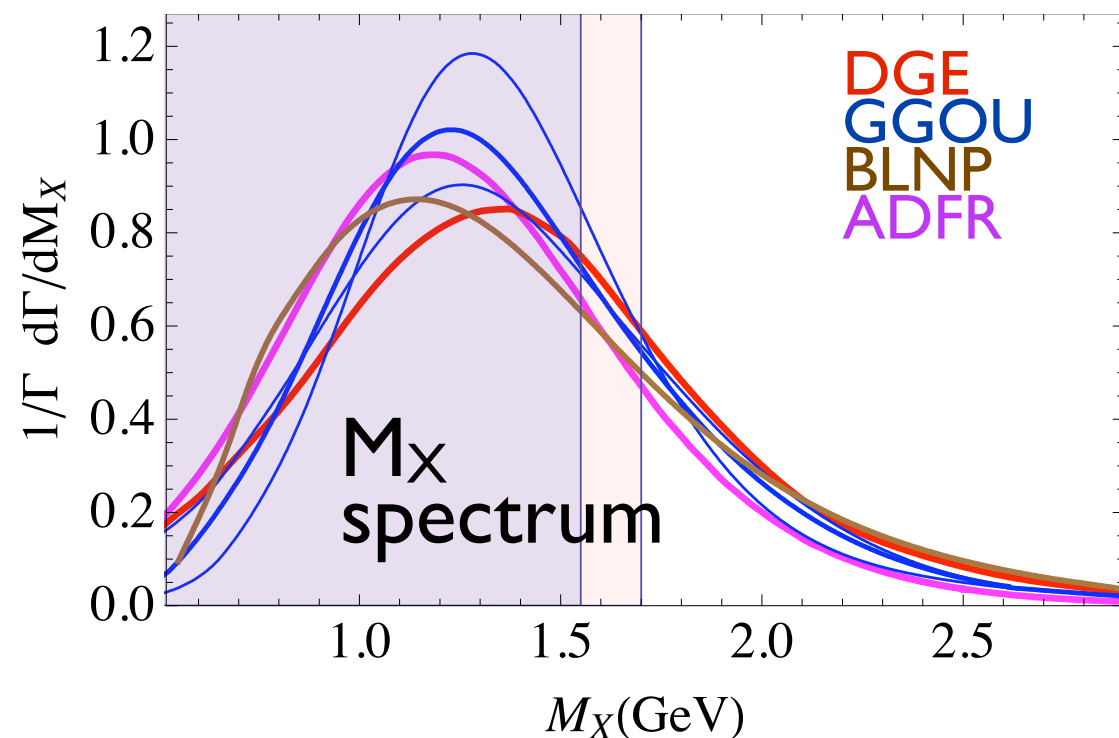
Fermi motion can be parameterized within the OPE like PDFs in DIS. At leading order in  $m_b$  only a single universal function of one parameter enters (SF).

Unlike resummed pQCD, **the OPE does not predict the SF**, only its first few moments. One then **needs an ansatz for its functional form**.

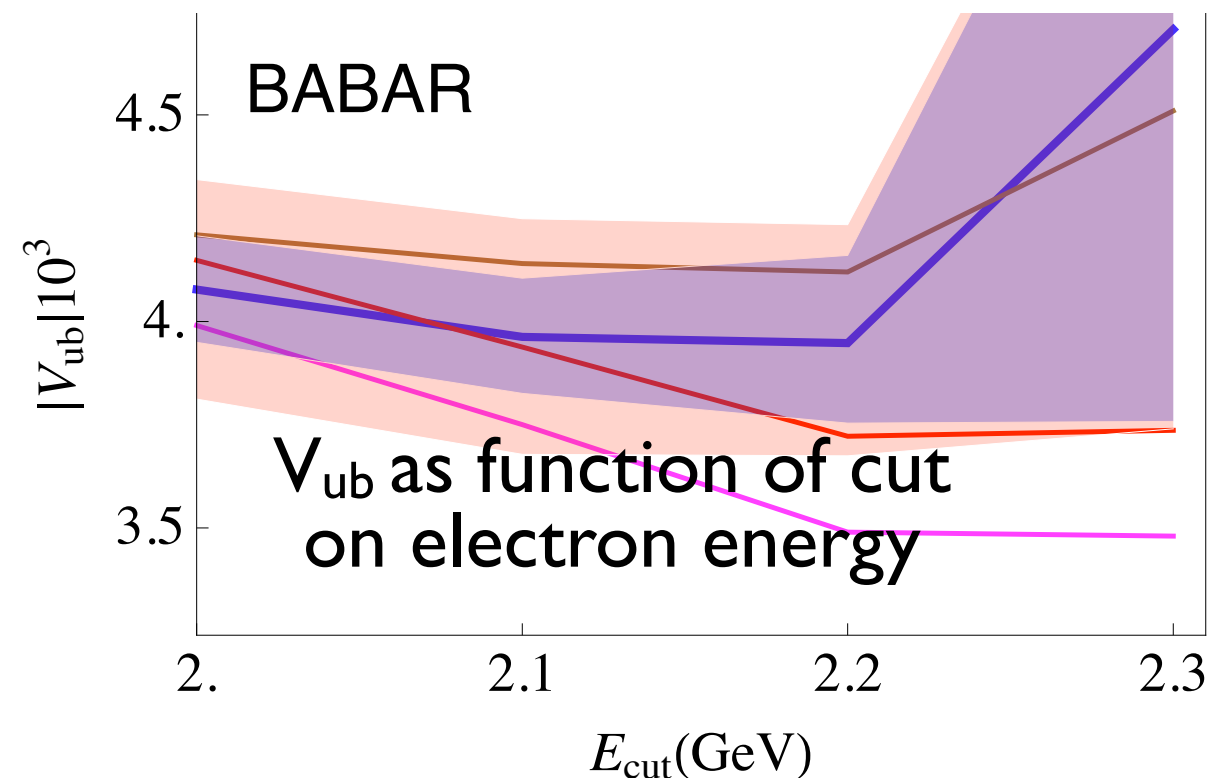
$$\int dk_+ k_+^n F_i(k_+, q^2) = \text{local OPE prediction} \Leftarrow \text{moments fits}$$

Two very different implementations:  
PG, Giordano, Ossola, Uraltsev (GGOU)  
Bosch, Lampe, Neubert, Paz (BLNP)





- \* Comparison at common inputs
- \* Overall good agreement with one exception **SPREAD WITHIN TH ERRORS**
- \* Precise measurement of spectra may allow to discriminate between them, but difficult for exp
- \* Strong  $m_b$  dependence  $V_{ub} \propto m_b^{4.5}$



- *Not all observables are equally clean. eg high  $q^2$  tail is sensitive to WA*
- Need spectra: only way to test frameworks (see  $E_l$  spectrum).
- *More inclusive measurements, less dependence on  $m_b$*
- Theory errors are partly parametric ( $m_b$ )

	Average $ V_{ub}  \times 10^3$
DGE	$4.26(14)_{\text{ex}}^{+19}_{-13}$
BLNP	$4.31(16)_{\text{ex}}^{+32}_{-27}$
GGOU	$3.96(15)_{\text{ex}}^{+20}_{-23}$

2.1, 1.9, 1.3 $\sigma$  from  $B \rightarrow \pi l \nu$   
(MILC-FNAL)

3.1, 2.4, 1.5 $\sigma$  from UFit  
(because of  $\sin 2\beta$ )

**NEW** preliminary Belle Multivariate analysis only  $E_l > 1 \text{ GeV}$

$$|V_{ub}| = (4.42 \pm 0.26^{+0.14}_{-0.22}) \times 10^{-3}_{\text{GGOU}}$$

2.1 $\sigma$  from excl, 2.5 $\sigma$  from UFit

*This includes about 90% of the rate really inclusive measurement, no need for SF. Only crucial input  $m_b$  needs to be confirmed!*

**NEW PHYSICS?**  
eg LR models Chen, Nam





# CP violation in B vs K decays

*Recent lattice results for  $B_K$  and previously neglected contributions lead to 15% smaller  $\epsilon_K$ , in conflict with  $\sin 2\beta$*

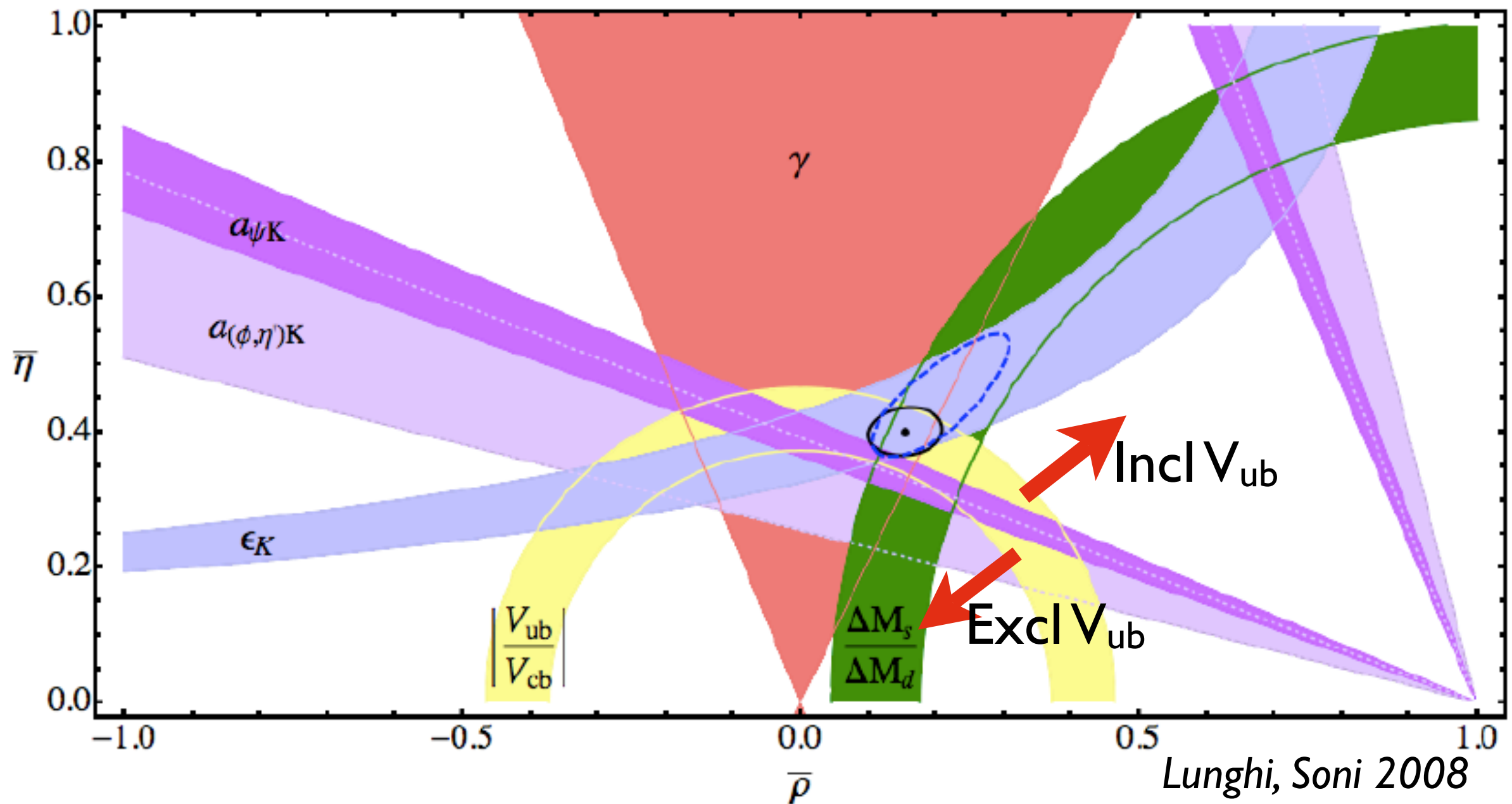
*Assuming SM, use  $\sin 2\beta$ ,  $|V_{cb}|$ ,  $\lambda$ ,  $\Delta M_s/\Delta M_d$ ,  $\xi$ ,  $B_K=0.720(39)$  RBC-UKQCD*

$$|\epsilon_K|^{SM} = (1.78 \pm 0.25) 10^{-3} \quad \text{vs} \quad |\epsilon_K|^{exp} = (2.229 \pm 0.012) 10^{-3}$$

**1.8-2.1  $\sigma$**  *depending on assigned errors* Buras, Guadagnoli, Lunghi, Soni

Easy to find new physics explanations, even in CMFV

*UTfit without  $\sin 2\beta_{exp}=0.672(24)$  gives  $\sin 2\beta=0.732 \pm 0.034 \Rightarrow 1.5\sigma$   
CKMfitter finds 2.5 tension between  $\sin 2\beta_{exp}$  and  $B \rightarrow \tau \nu$  (depends strongly on  $B_B$ )*



*Perhaps  $\sin 2\beta$  is simply too low...*  
*... or  $\text{incl } V_{ub}$  and  $B_K$  both wrong*

# Conclusions

CKM is overall in a good shape

Great progress in lattice calculations

A few  $\sim 2\sigma$  problems: inclusive and exclusive  $|V_{xb}|$  tend to clash,  $\sin 2\beta$  seems a bit low, too early to invoke new physics

We need better  $m_{b,c}$  *determinations* and to exploit data to check theory calculations (shapes, distributions etc)

Important constraints on new physics but no time to discuss them