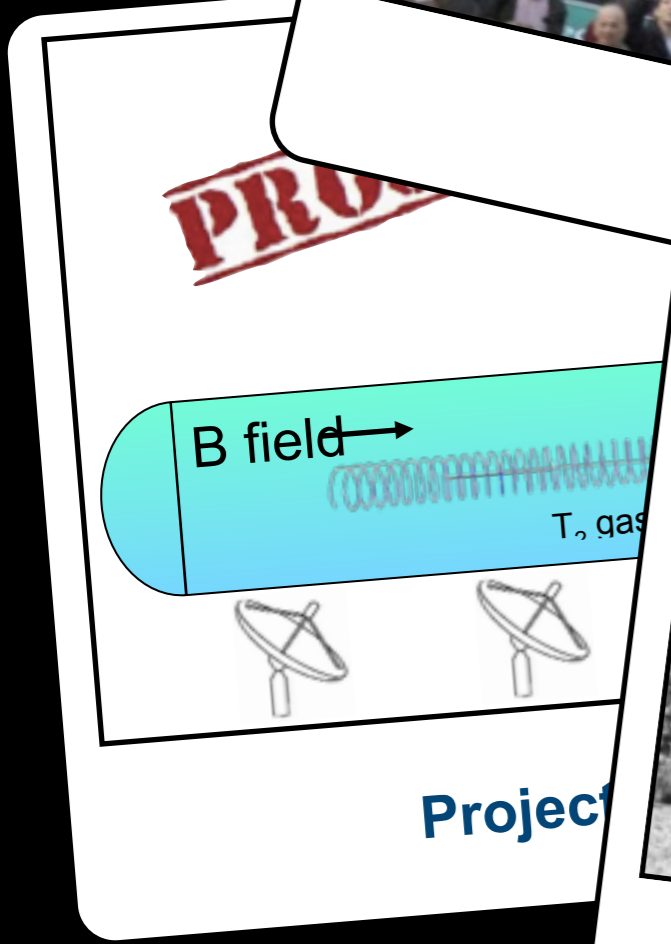




KATRIN



Project



Wilson & Penzias

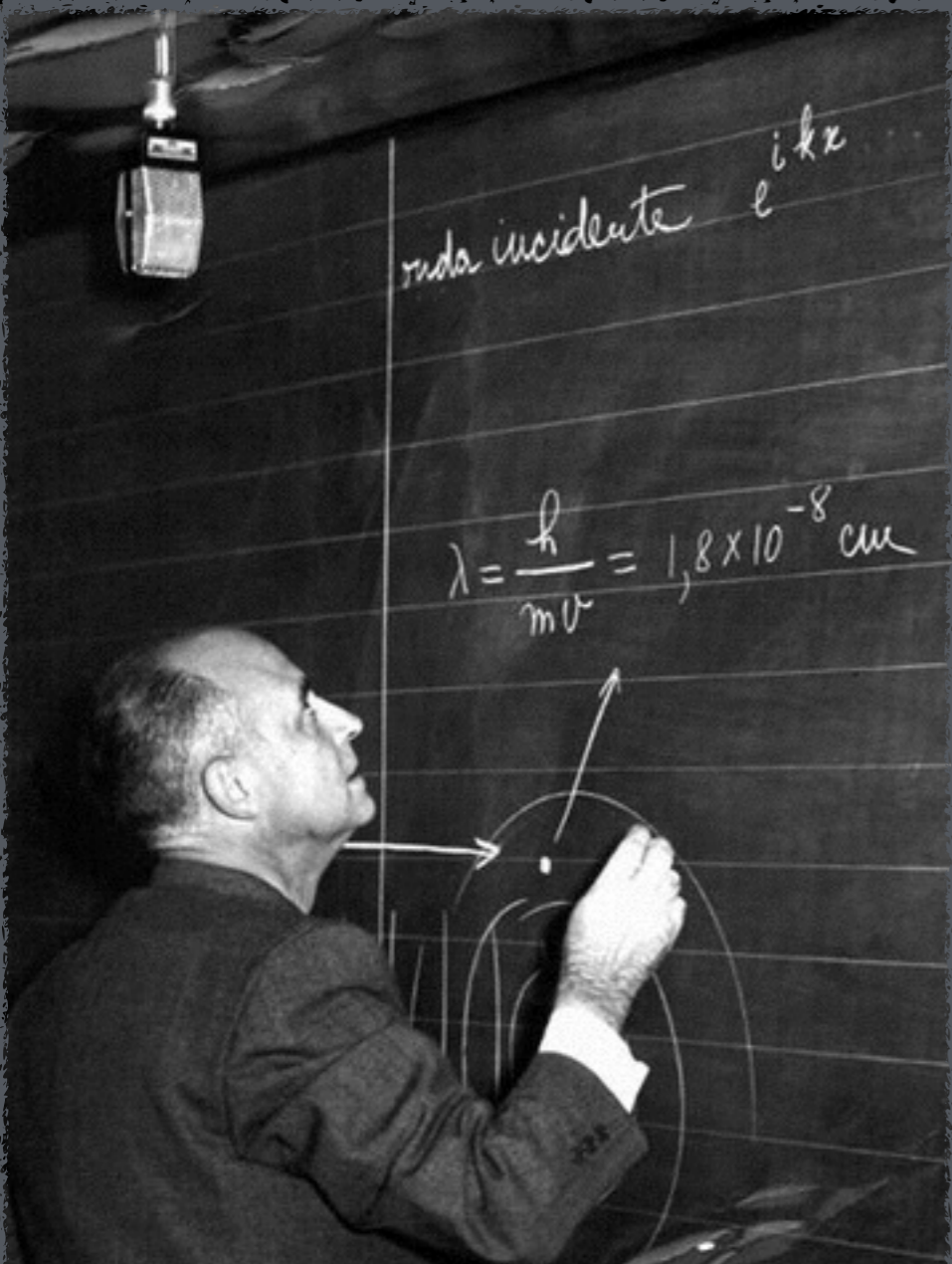
Project 8

&

Alternative Paths
in Beta Decay
Experiments

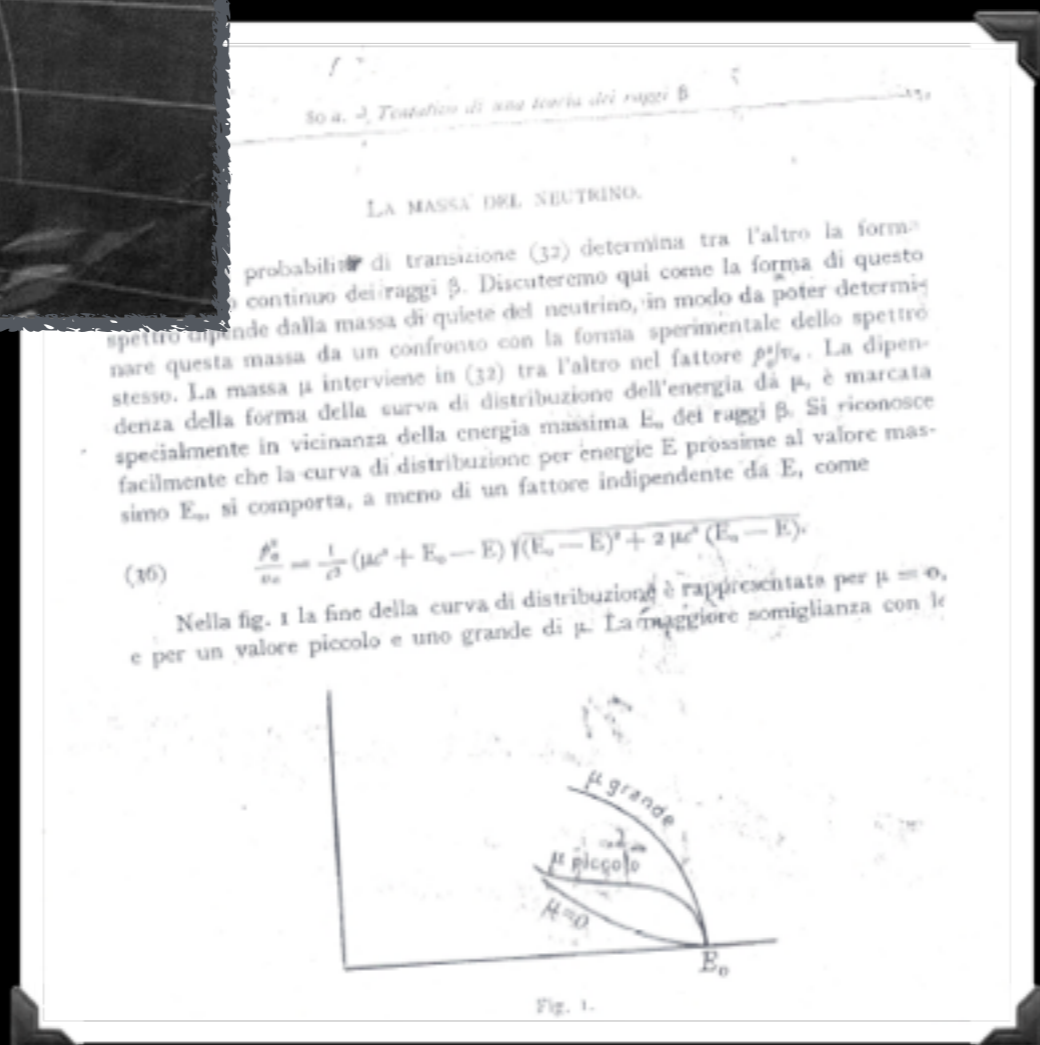
Gran Sasso National Laboratory
May 14th 2014

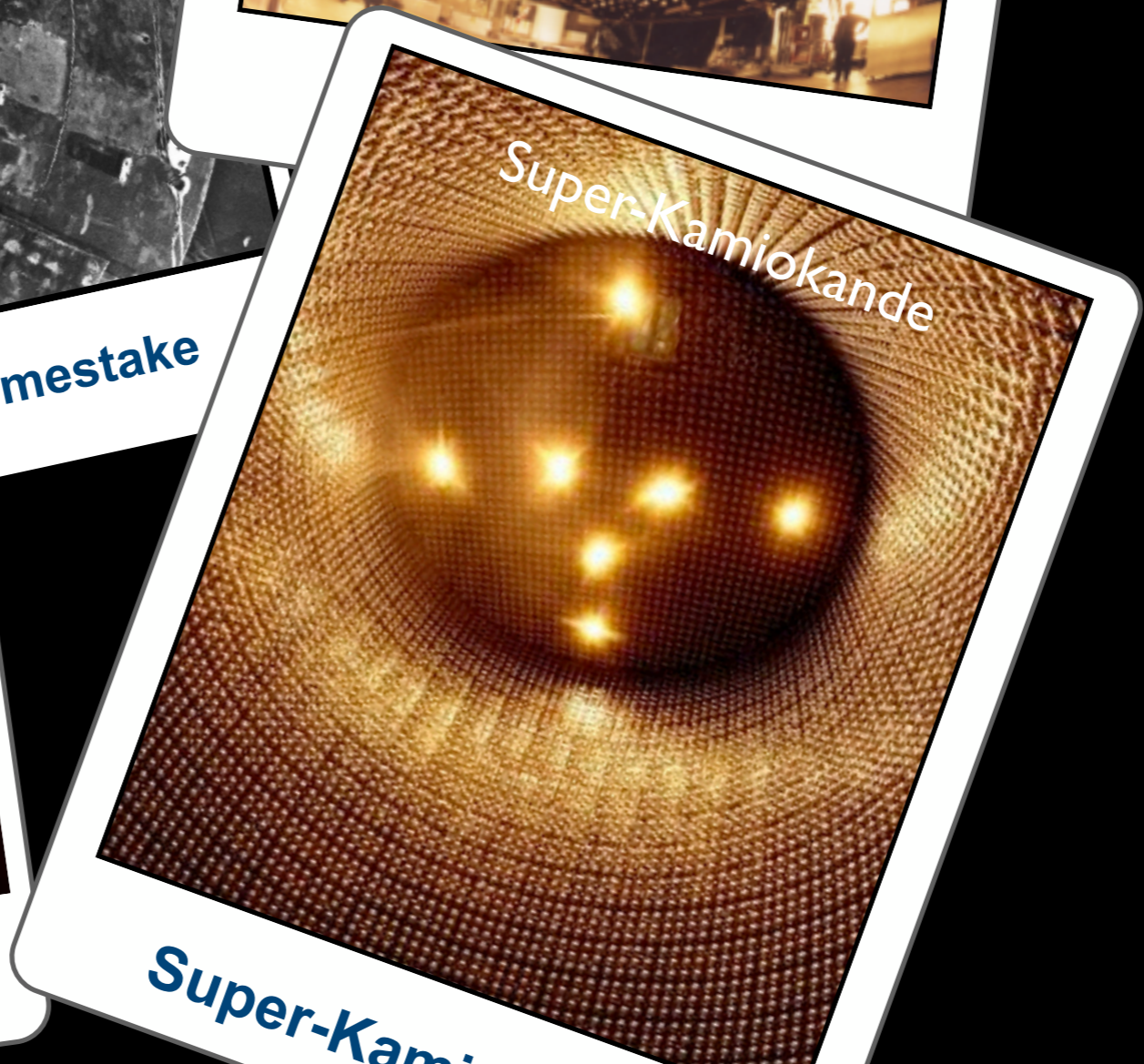
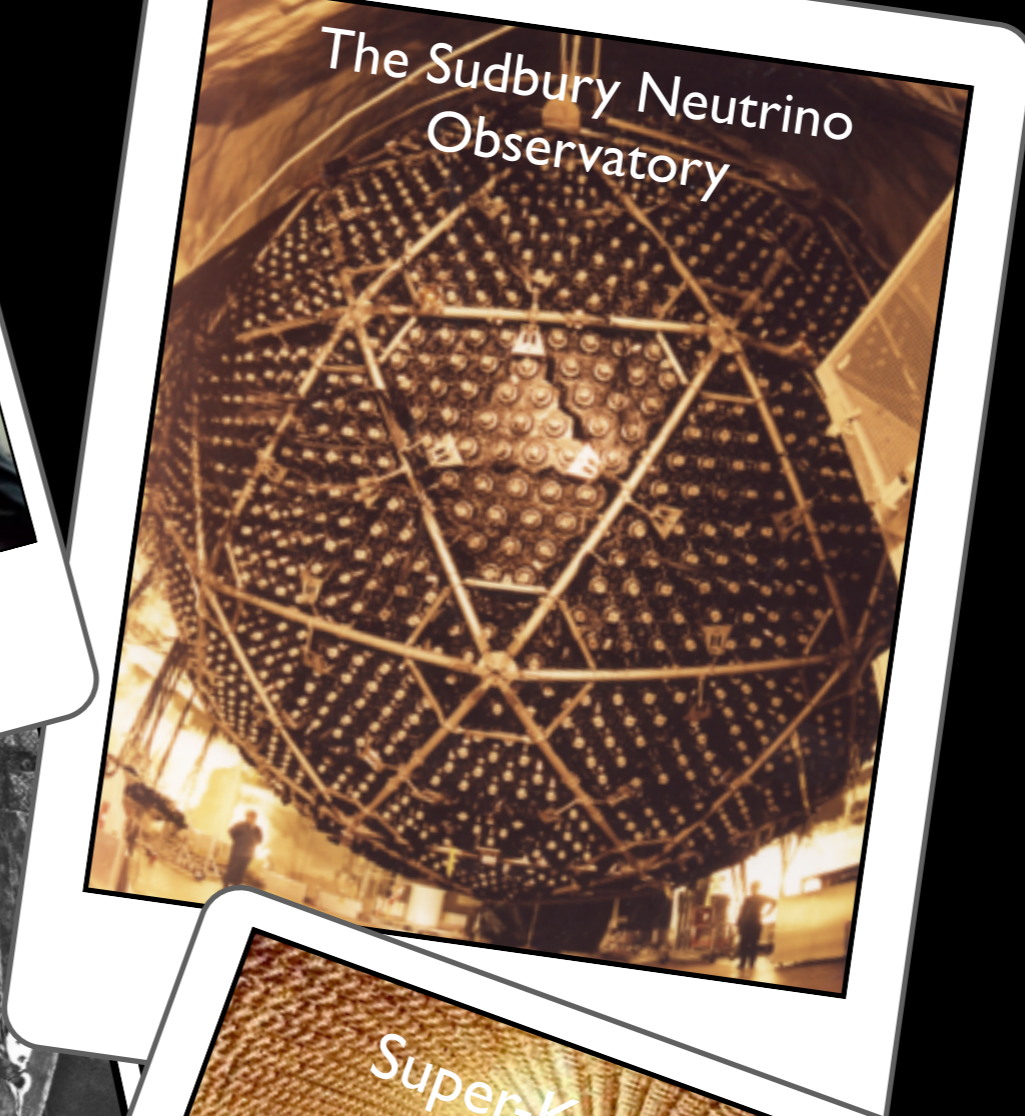
Joseph A. Formaggio
MIT



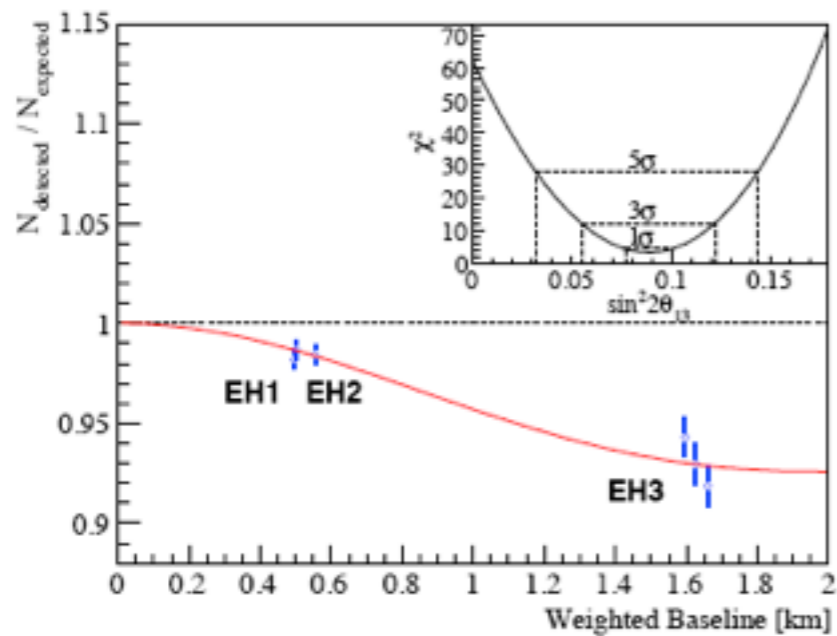
Neutrino mass measurements have a long history in physics, predating the Standard Model itself.

It should therefore be no surprise that our quest still continues to understand this fundamental property, both in its own right as well as its theoretical implications.





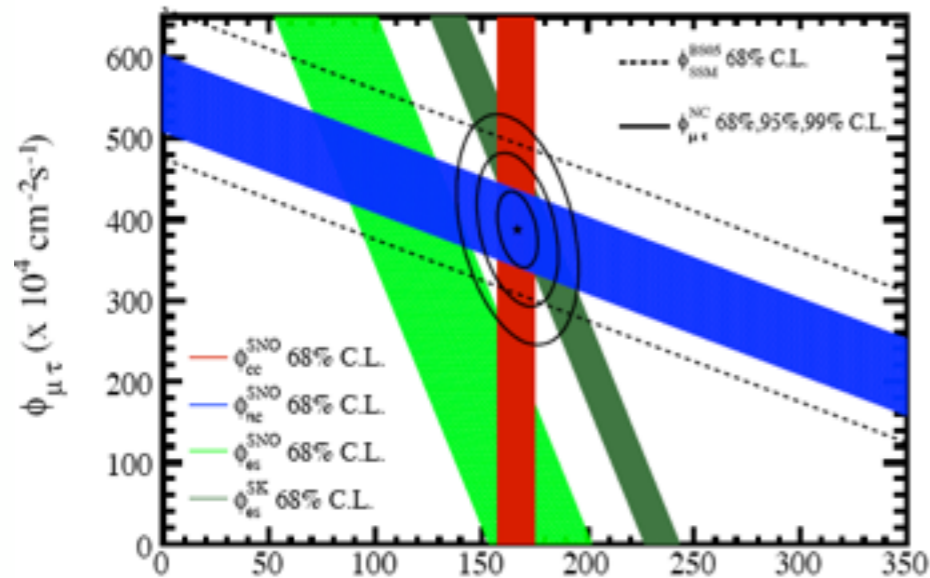
The phenomena of neutrino oscillations is now firmly established.



$$\sin^2(\theta_{13}) = 0.0241 \pm 0.0025$$

Reactor & Long Baseline

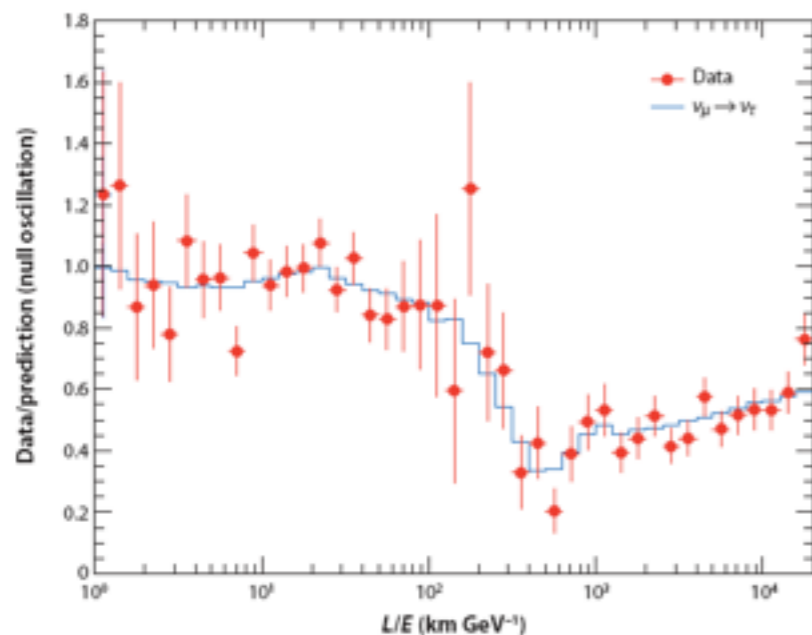
Precision measurements now exist on all three mixing angles to date.



$$\sin^2(\theta_{12}) = 0.307 \pm 0.016$$

$$\Delta m_{12}^2 = (7.54 \pm 0.26) \times 10^{-5} \text{ eV}^2$$

Solar



$$\sin^2(\theta_{23}) = 0.386 \pm 0.022$$

$$\Delta m_{23}^2 = (2.43 \pm 0.09) \times 10^{-3} \text{ eV}^2$$

Atmospheric

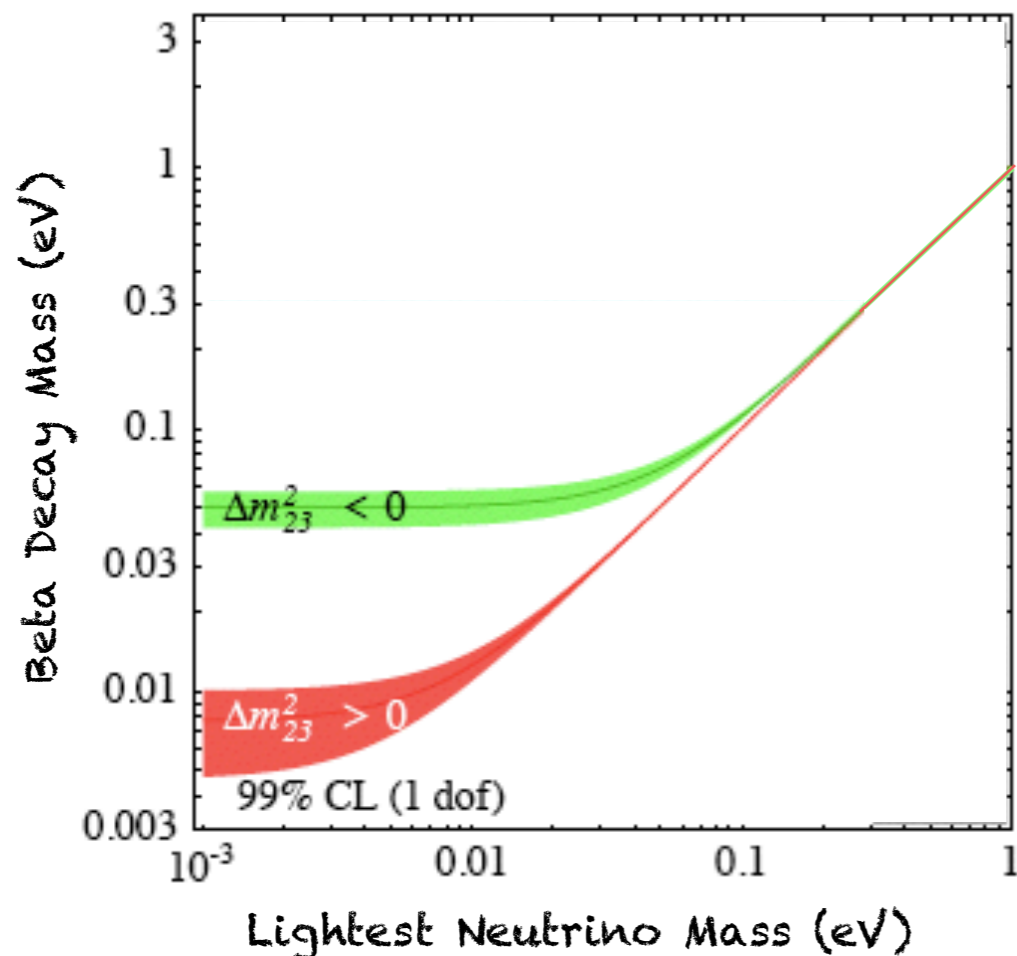
As such, oscillation measurements place a lower limit on the neutrino mass scale.

Measuring

Neutrino Masses

Neutrino oscillations have placed a lower bound on neutrino masses that can be experimentally accessed.

Lower bound depends on hierarchy of neutrinos (inverted or normal)



$$M = \sum_i^{n_\nu} m_{\nu,i}$$

Cosmological Measurements

$$\langle m_{\beta\beta}^2 \rangle = \left| \sum_i^{n_\nu} U_{ei}^2 m_{\nu,i} \right|^2$$

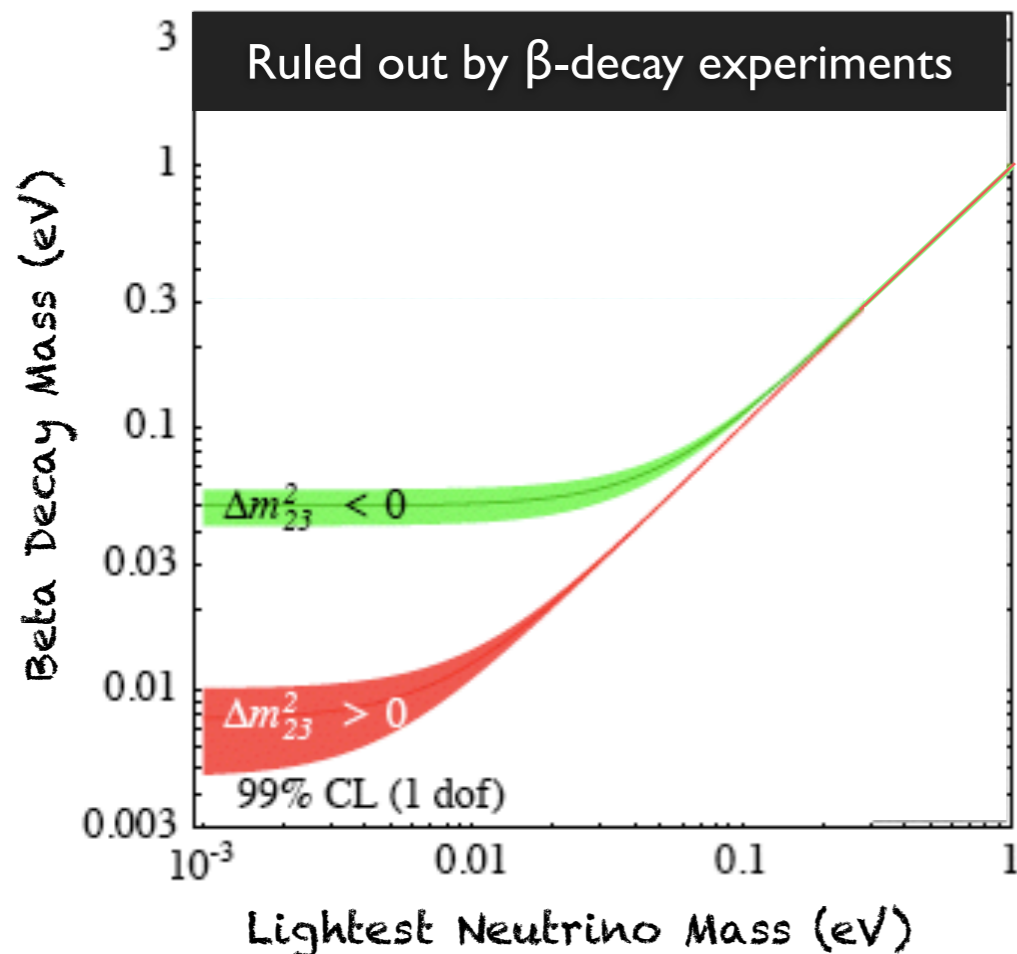
$0\nu\beta\beta$ Measurements

$$\langle m_\beta \rangle^2 = \sum_i^{n_\nu} |U_{ei}|^2 m_{\nu,i}^2$$

Beta Decay Measurements

The Neutrino Mass Scale

- The neutrino mass scale remains one of the essential "unknowns" of the Standard Model.
- Knowledge of neutrino masses can have a significant impact on many different arenas, including cosmology, the mass hierarchy, sterile neutrinos, and even relic neutrino detection.



$m_\nu > 2$ eV (eV scale, current)

Neutrinos ruled out as dark matter

$m_\nu > 0.2$ eV (degeneracy scale)

Impact on cosmology and $0\nu\beta\beta$ reach

$m_\nu > 0.05$ eV (inverted hierarchy)

Resolve hierarchy if null result

$m_\nu > 0.01$ eV (normal hierarchy)

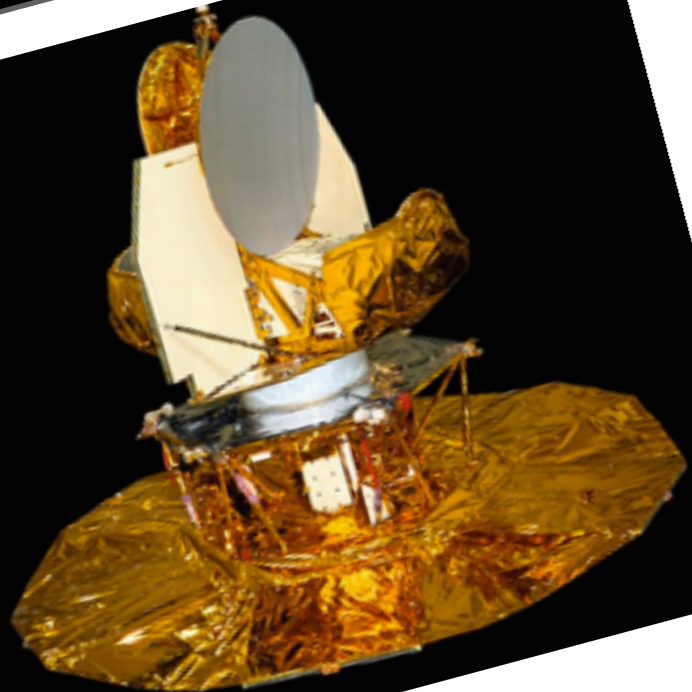
Oscillation limit; possible $C\nu B$ detection

The Era of Precision Cosmology



Cosmology has had a similar trajectory as neutrino physics, from inception to present day

The Era of Precision Cosmology



WMAP



Wil



Atacama
Cosmology Telescope

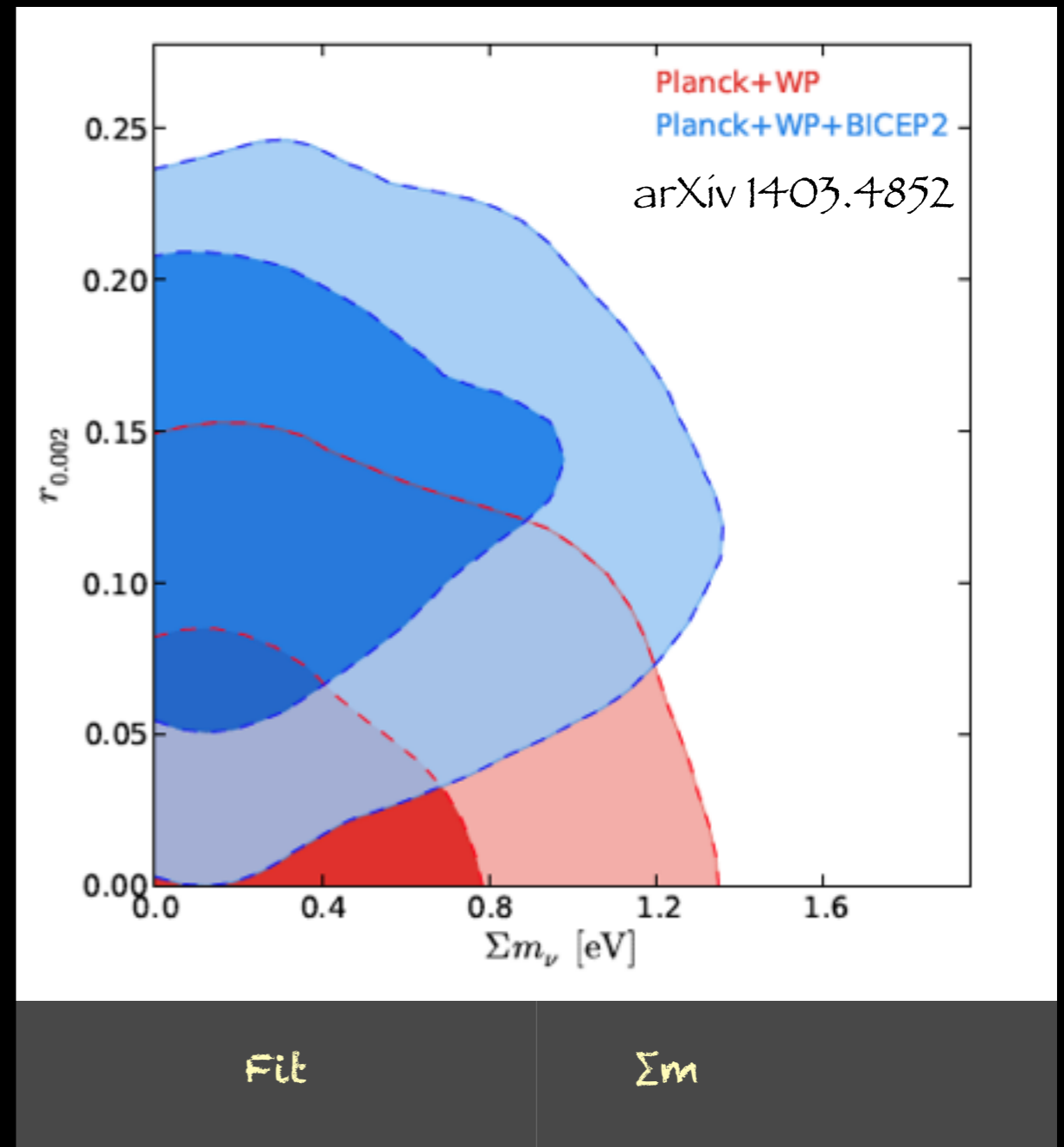
Cosmology has had a similar trajectory as neutrino physics, from inception to present day



Sloan Digital Sky Survey

PLANCK Results

- The basic PLANCK analysis looks at 6 main cosmological parameters. Neutrino masses are added as extensions to that model.
- Most conservative data combinations see no evidence for neutrino masses.
- Certainly tension exists with certain parameters (SZ clusters, Hubble constant, BICEP2) that alter the fits or in some cases favor finite masses.



Planck + WP + HighL

$< 0.66 \text{ eV}$

Planck + WP + HighL
+ BAO

$< 0.23 \text{ eV}$

Planck + SZ Clusters

$0.37 \pm 0.20 \text{ eV}$

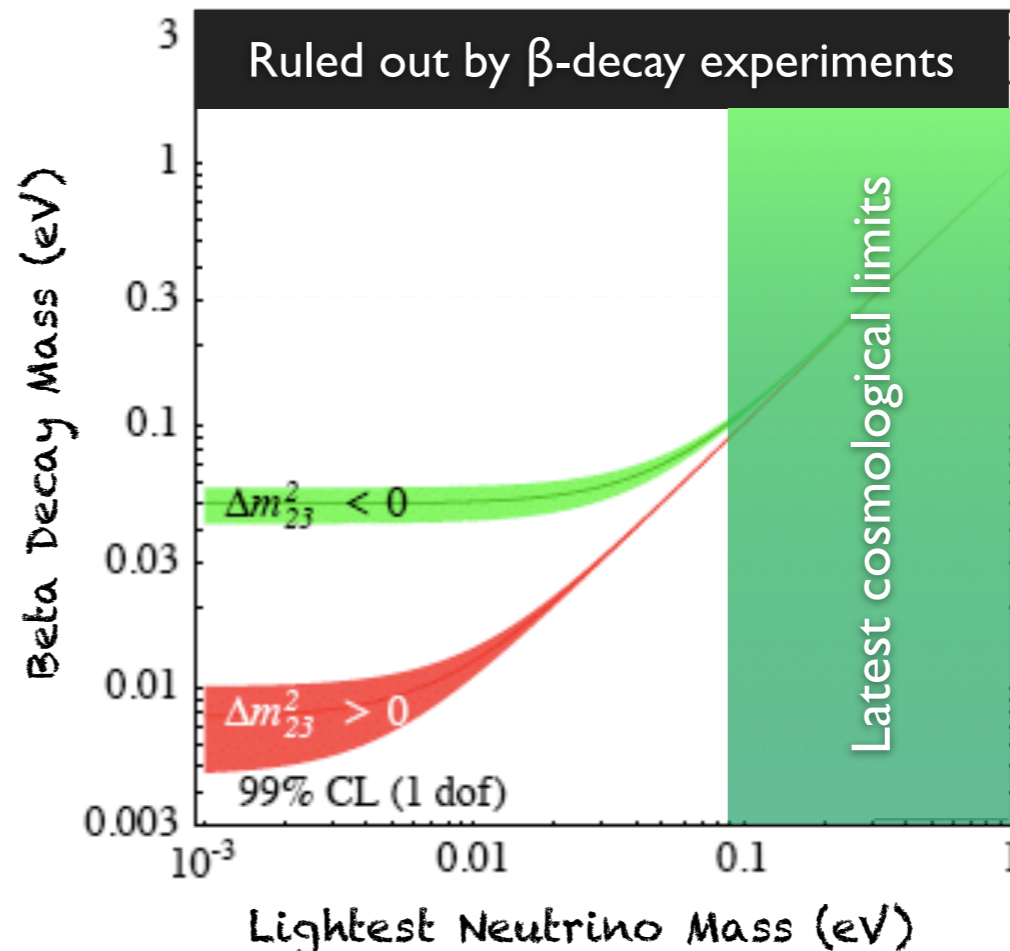
Moving Forward...

- Current cosmological limits are starting to push at the degeneracy-inverted scale.
- Future experiments (CMB-IV) could push all the way down to the normal scale.

- Model dependencies and degeneracies will still persist.

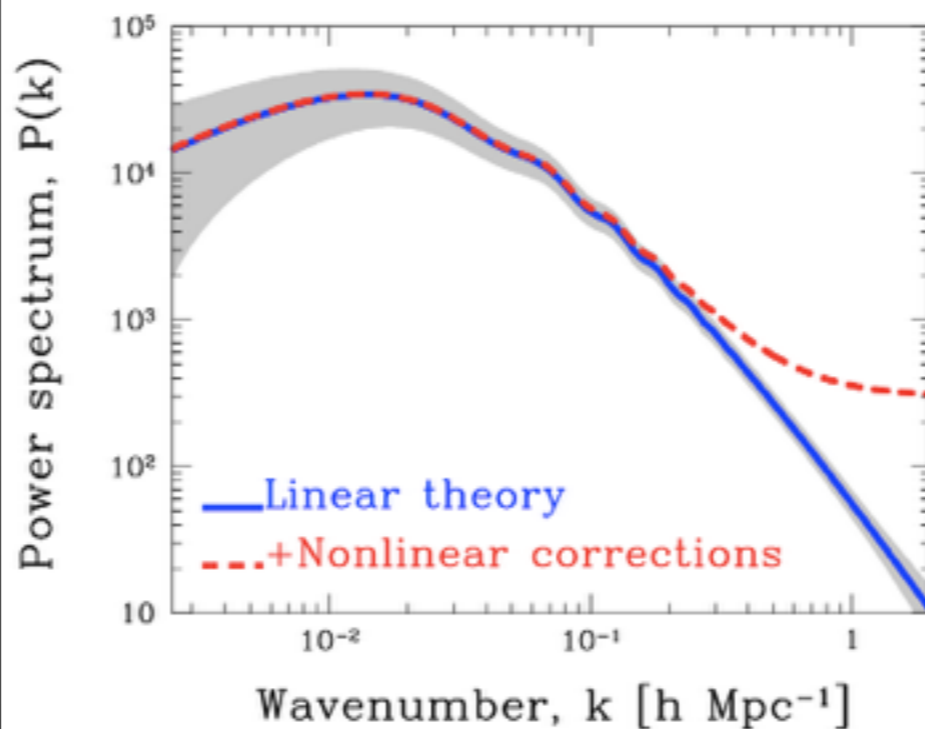
$$\frac{\Delta P}{P} \simeq -12 \frac{\Omega_\nu}{\Omega_m} \simeq 1\%$$

- Moving to the normal hierarchy scale now requires 1% precision on the power spectrum.

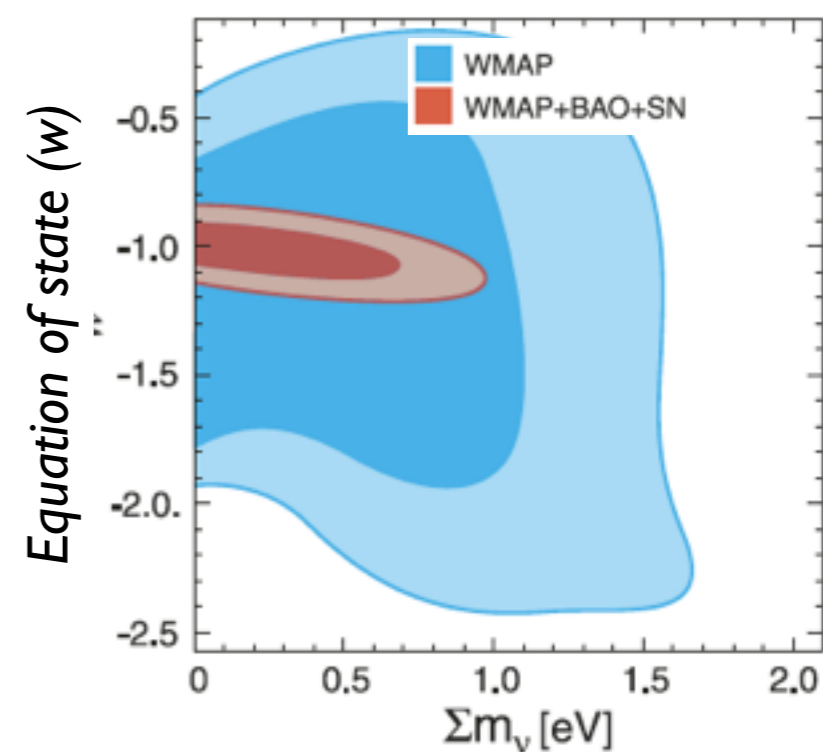


Y. Y. Y. Wong, 2010

S. Hannestad
Phys. Rev. Lett 95 221301



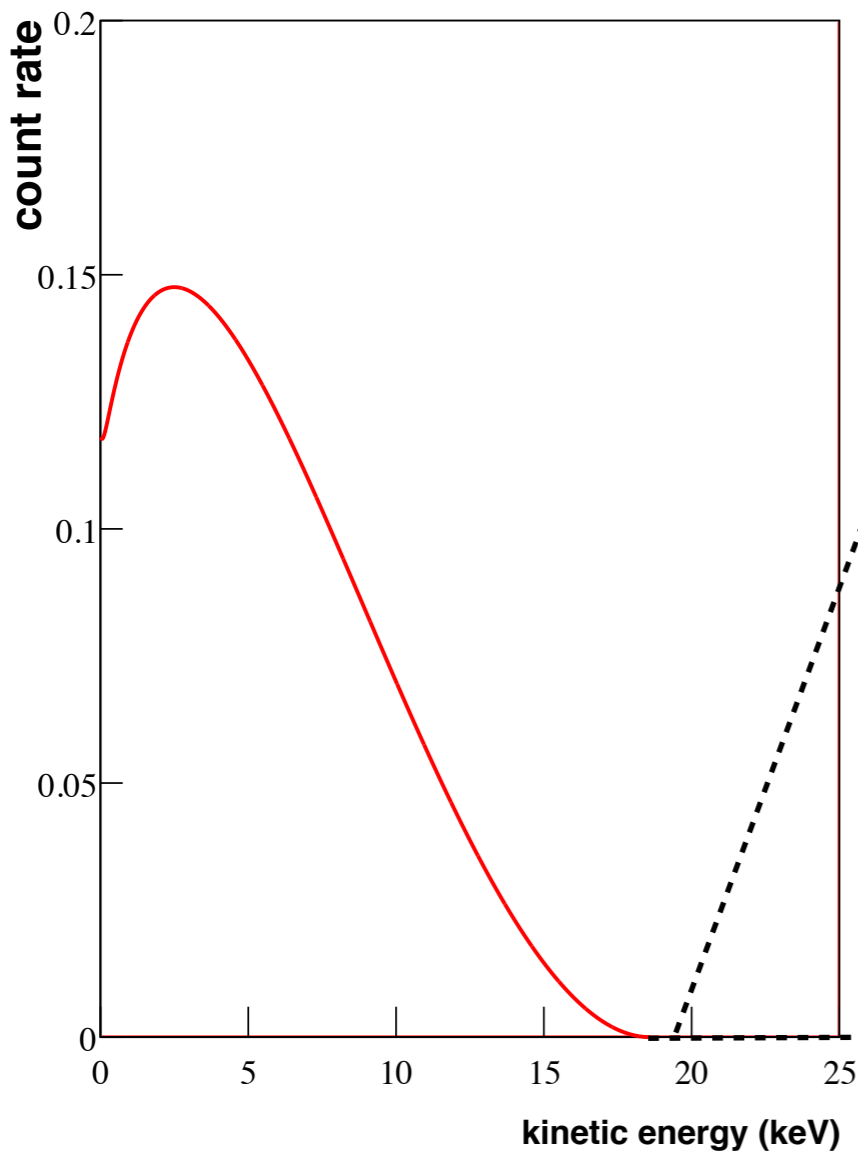
Nonlinearities



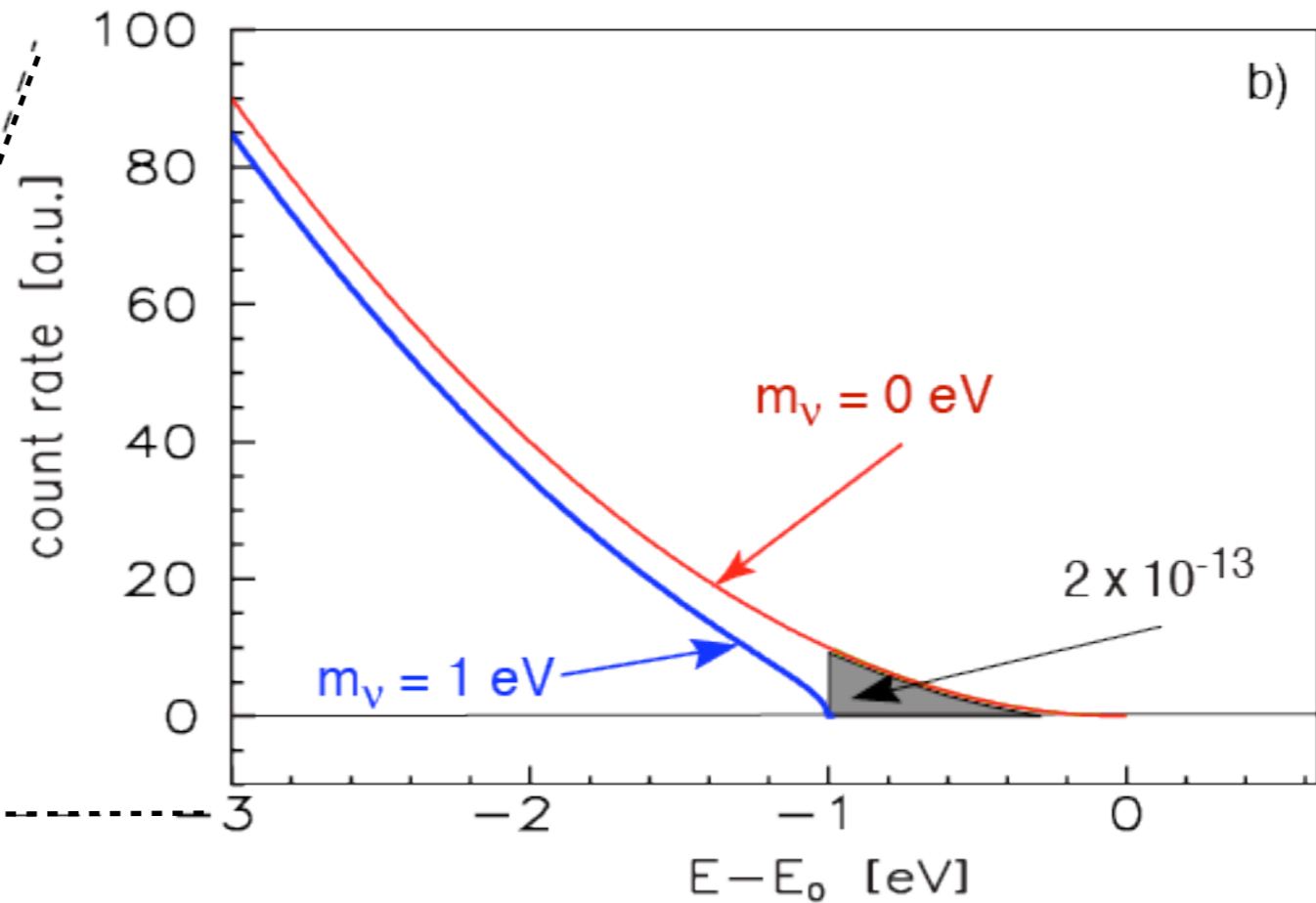
Degeneracies

Direct Probes

Electron Energy



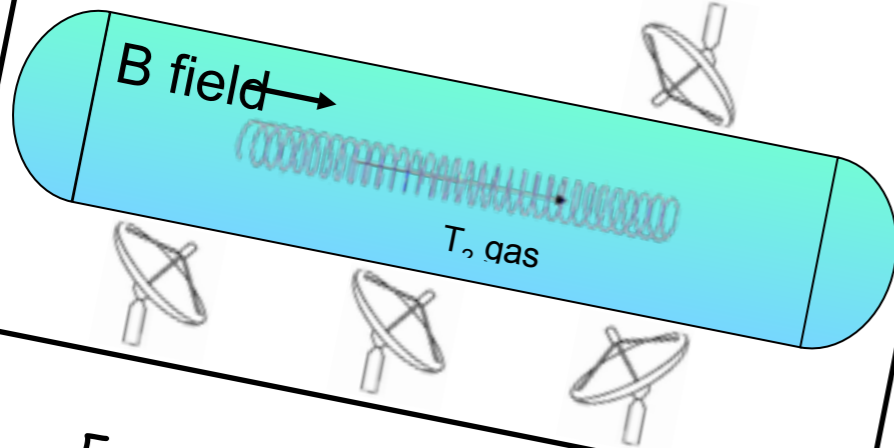
$$\dot{N} \sim p_e (K_e + m_e) \sum_i |U_{ei}|^2 \sqrt{E_0^2 - m_{\nu i}^2}$$



Beta Decay

A kinematic determination of the neutrino mass
No model dependence on cosmology or nature of mass

PROJECT 8



Frequency Techniques



MARE-HOLMES & ECHO
Calorimetry



KATRIN

Electromagnetic Spectroscopy

KATRIN is currently the prominent experiment for beta decay measurements.

New techniques being explored in the future:

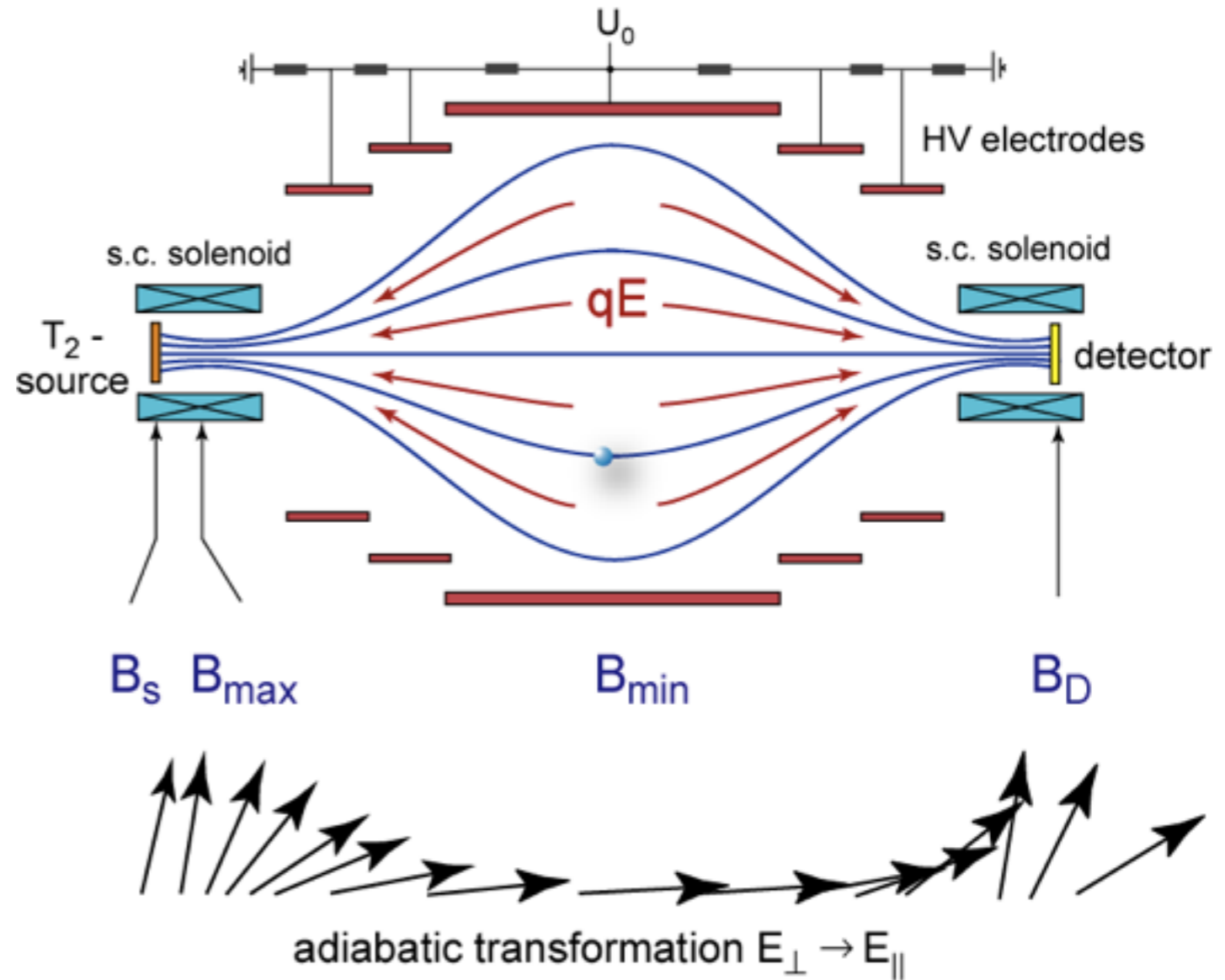
ECHO, HOLMES and Project 8

MAC-E Filter Technique

KATRIN



Spectroscopic: MAC-E Filter

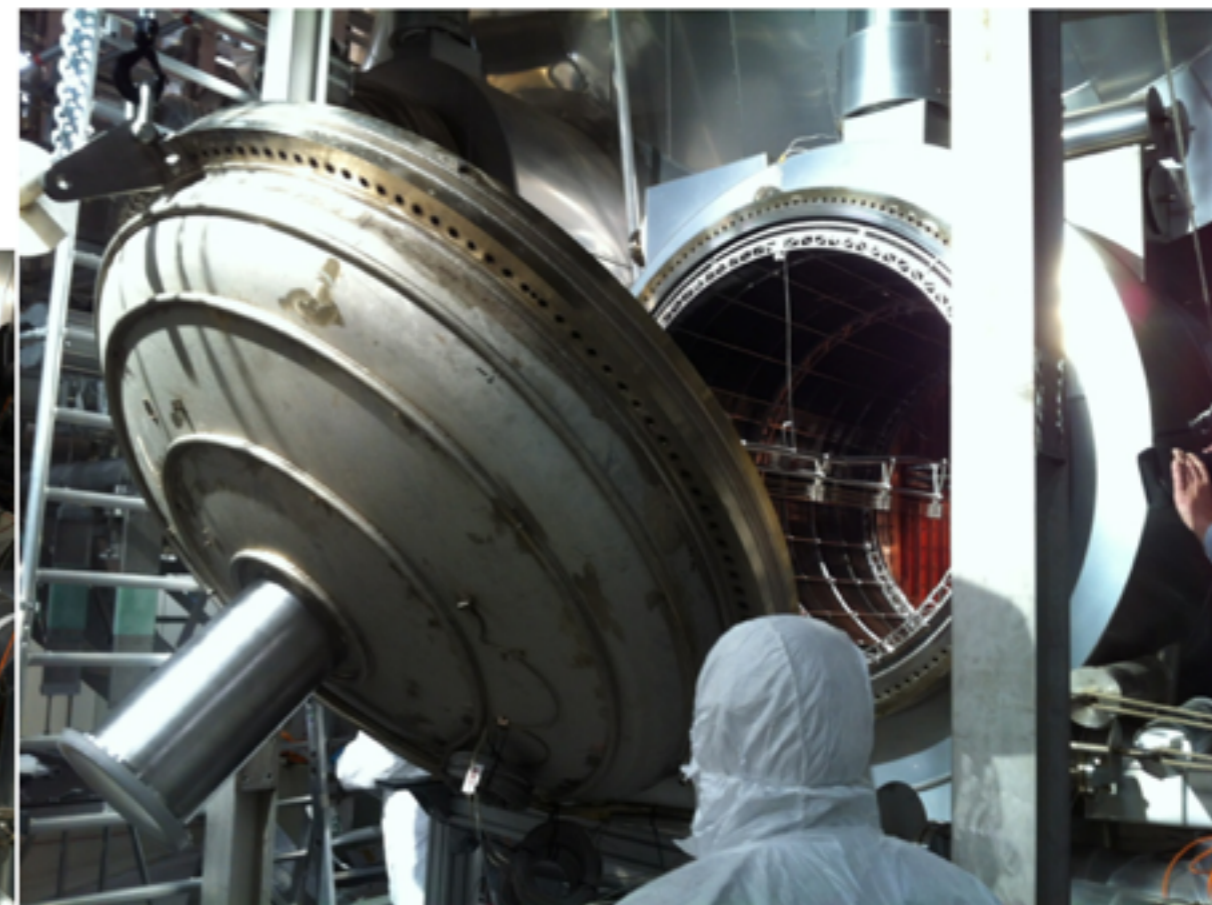


Inhomogeneous magnetic guiding field.
Retarding potential acts as high-pass filter

High energy resolution

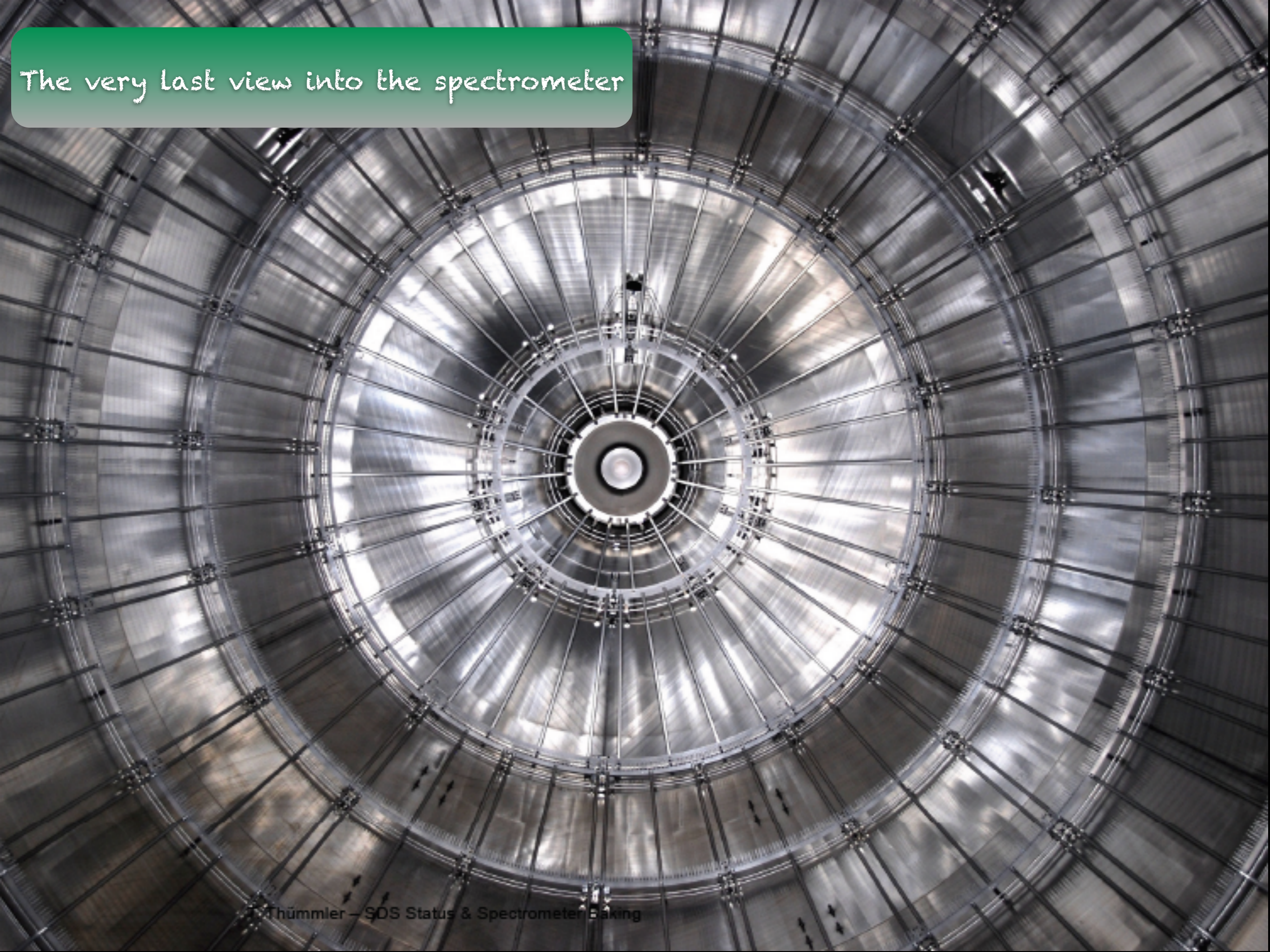
$$(\Delta E/E = B_{\min}/B_{\max} = 0.93 \text{ eV})$$

The Main Spectrometer



Recent
milestone:
Final pump
port closed
and sealed.

Tuesday,
May 8, 2012
14:11 CEST



The very last view into the spectrometer

SDS Commissioning

"First Light"

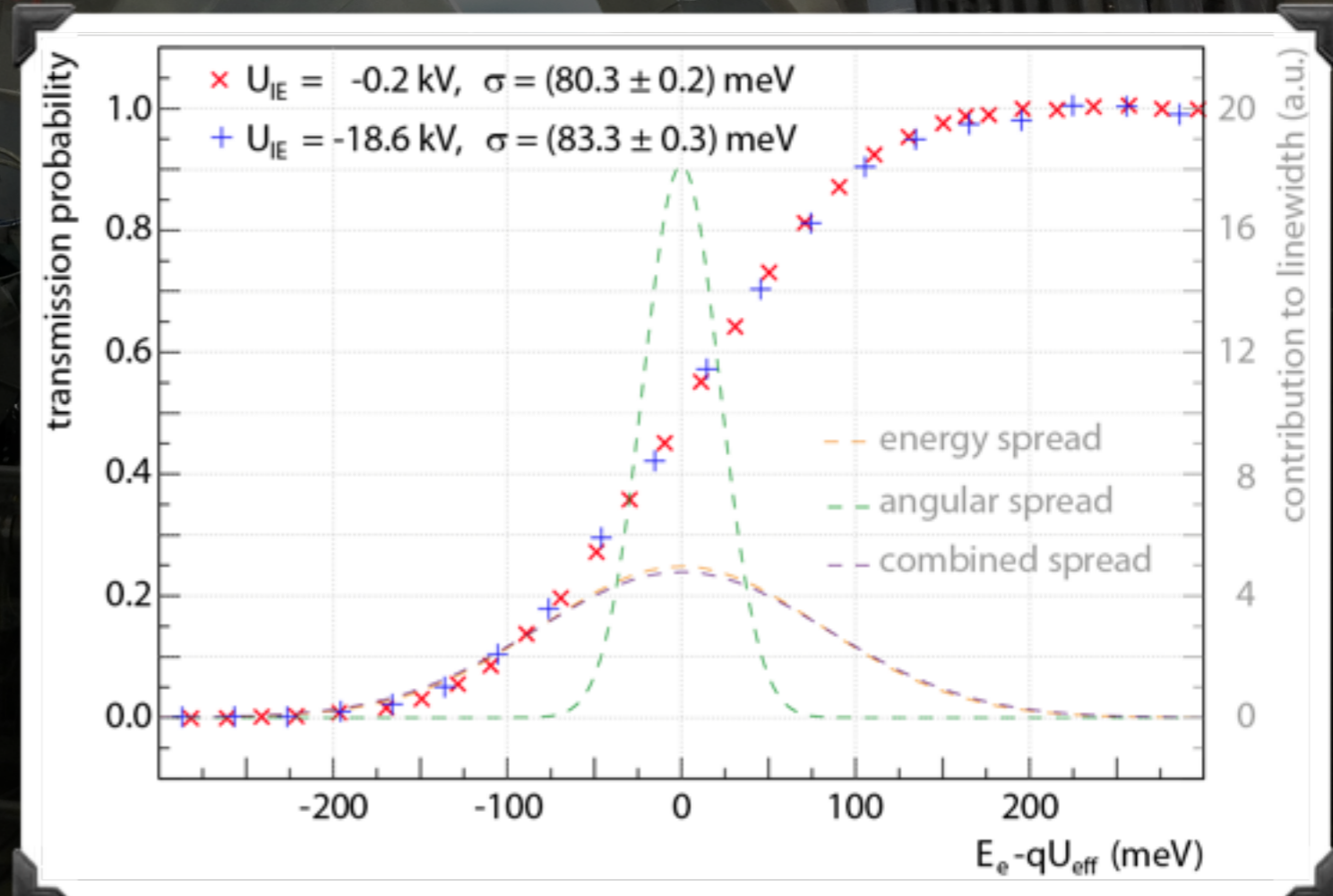
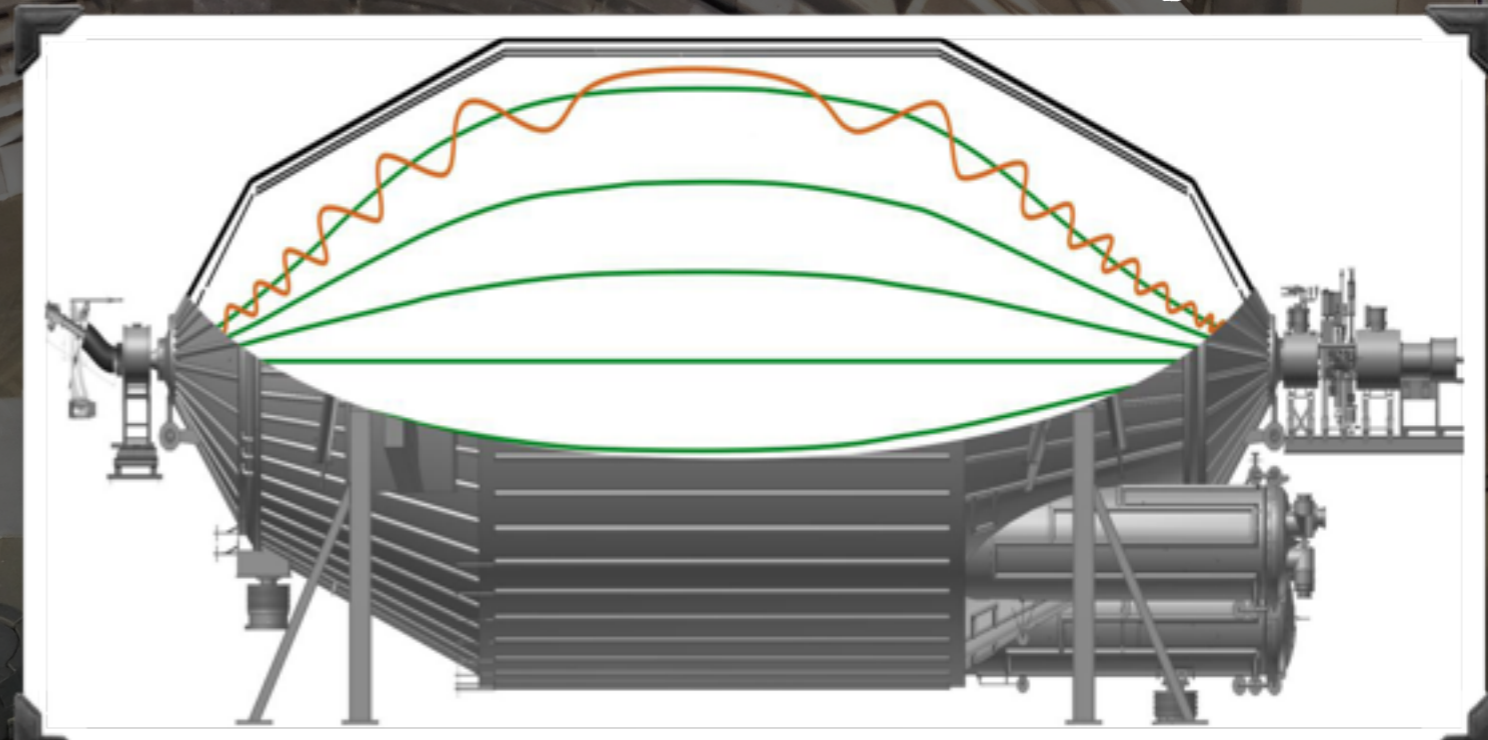
First time pre-spectrometer, main spectrometer, and detector are all connected.

First electrons in this combined system now recorded.

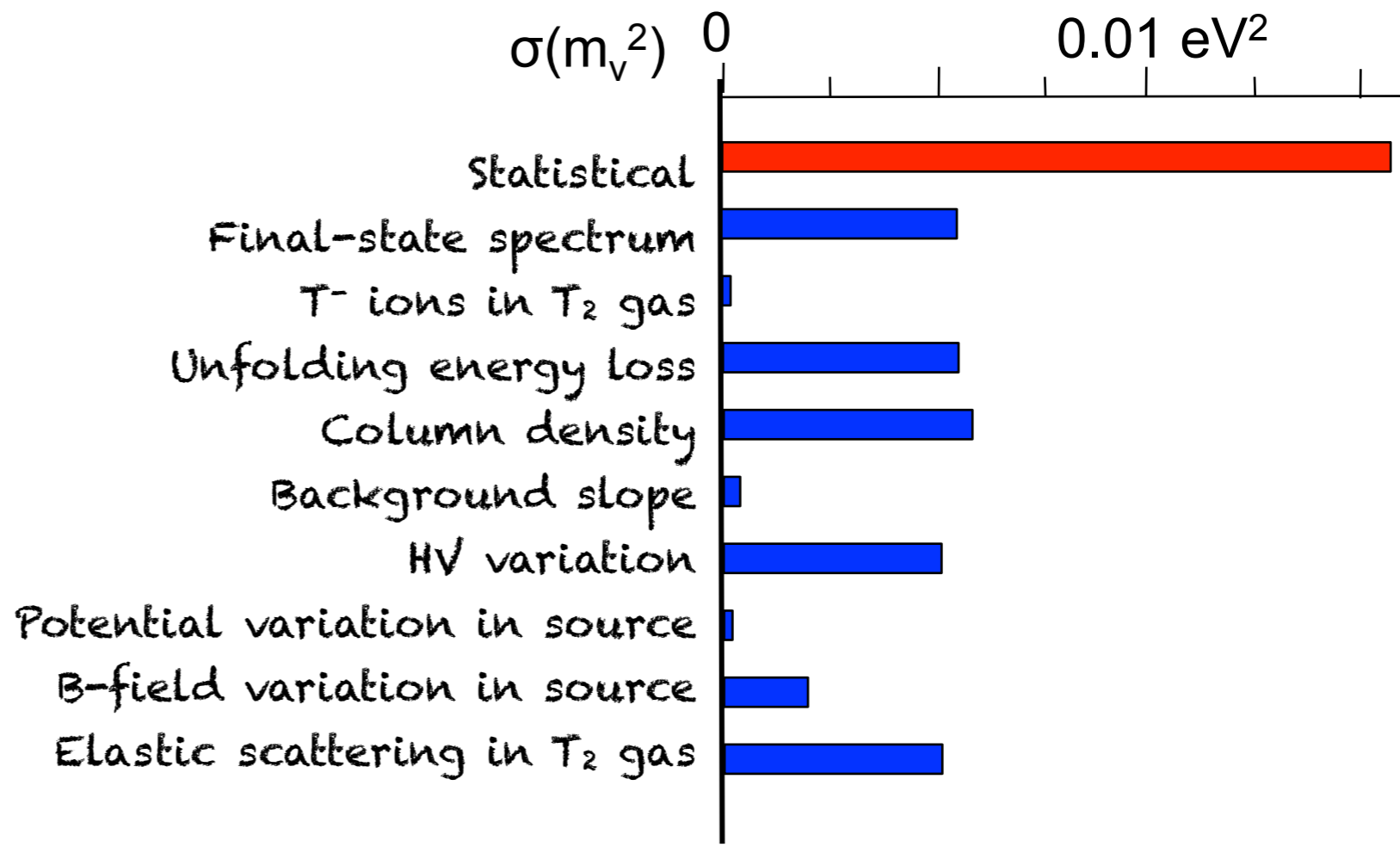
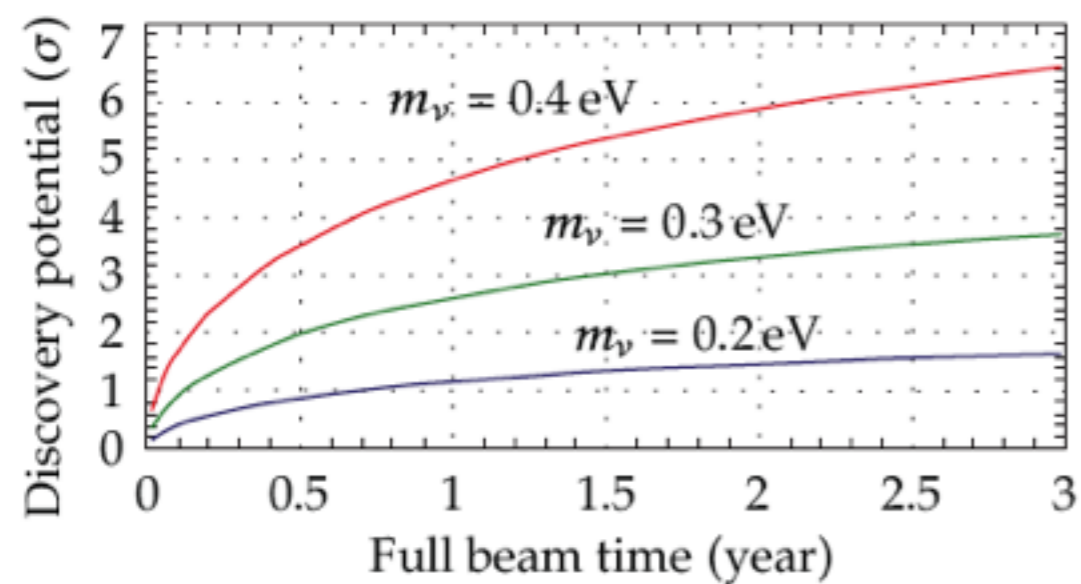
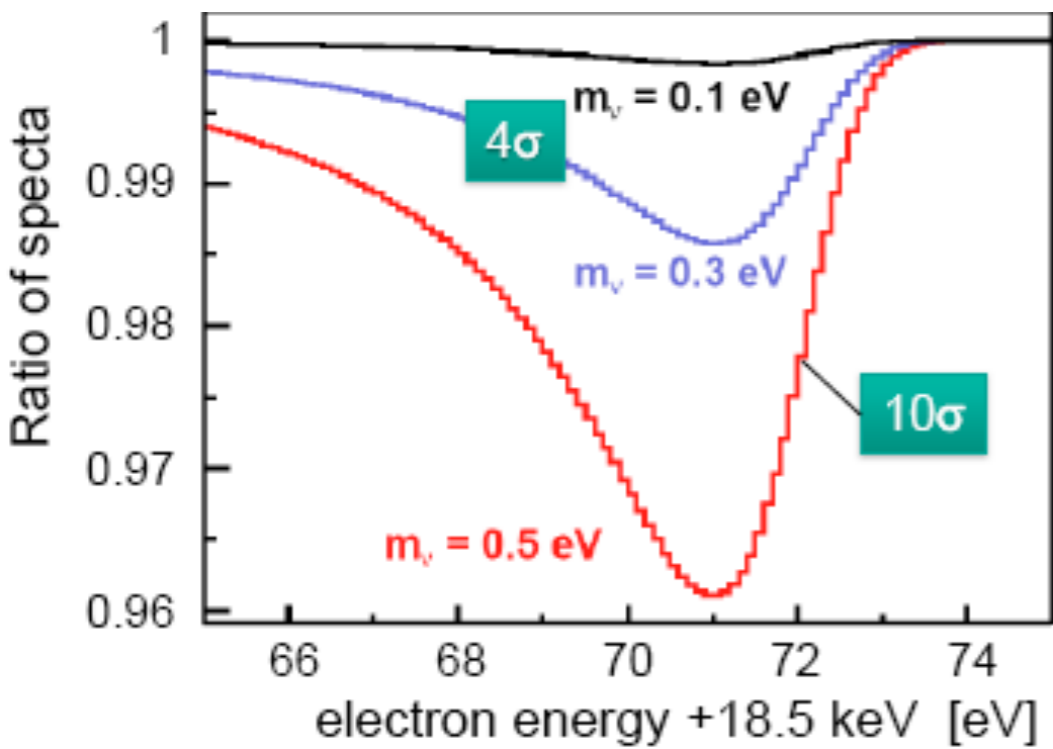
Adiabatic guidance and electric retarding work as expected.

Backgrounds measured and several background reduction techniques demonstrated.

Second phase commissioning program to commence in summer 2014.



Projected Sensitivity



Neutrino Mass Goals

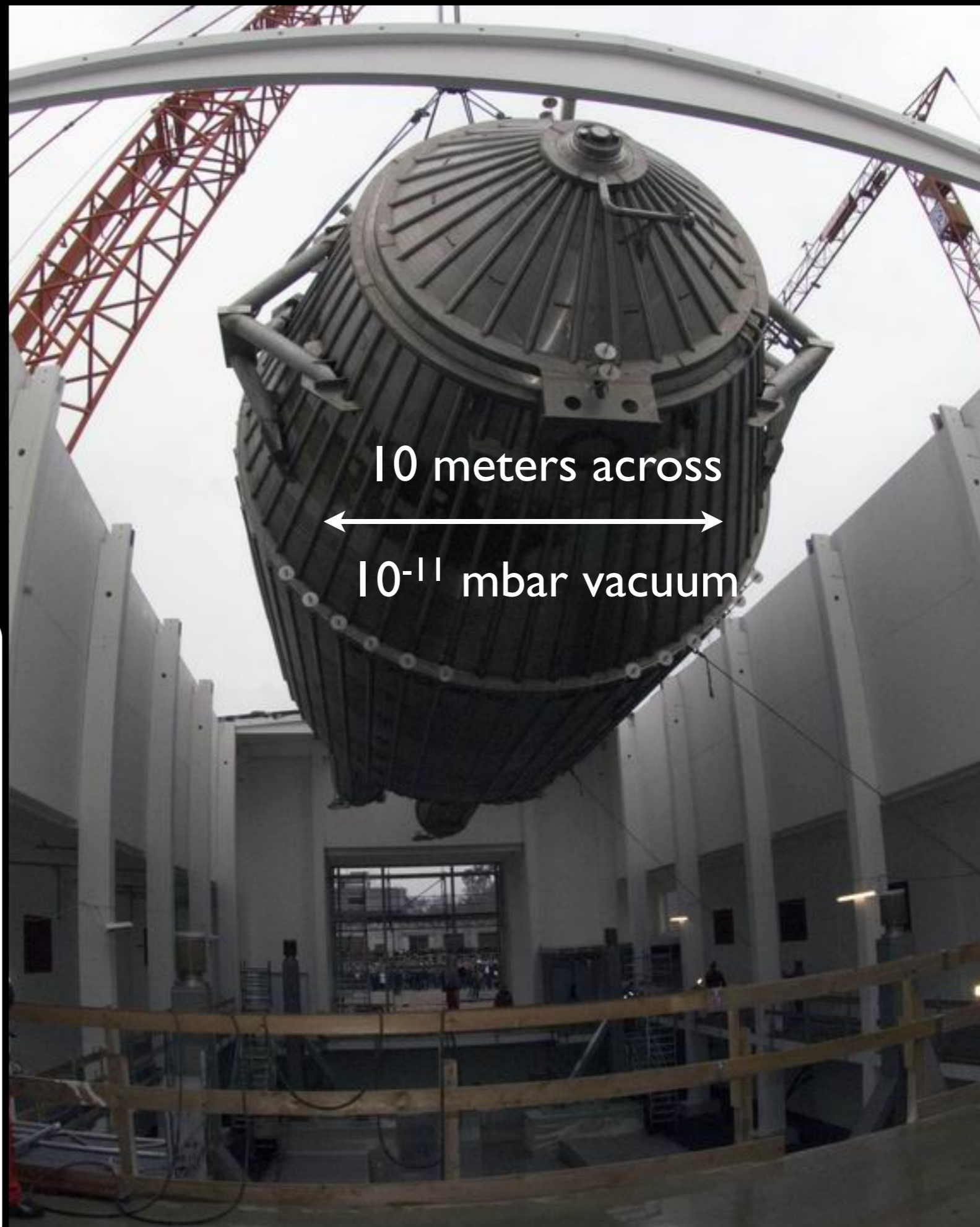
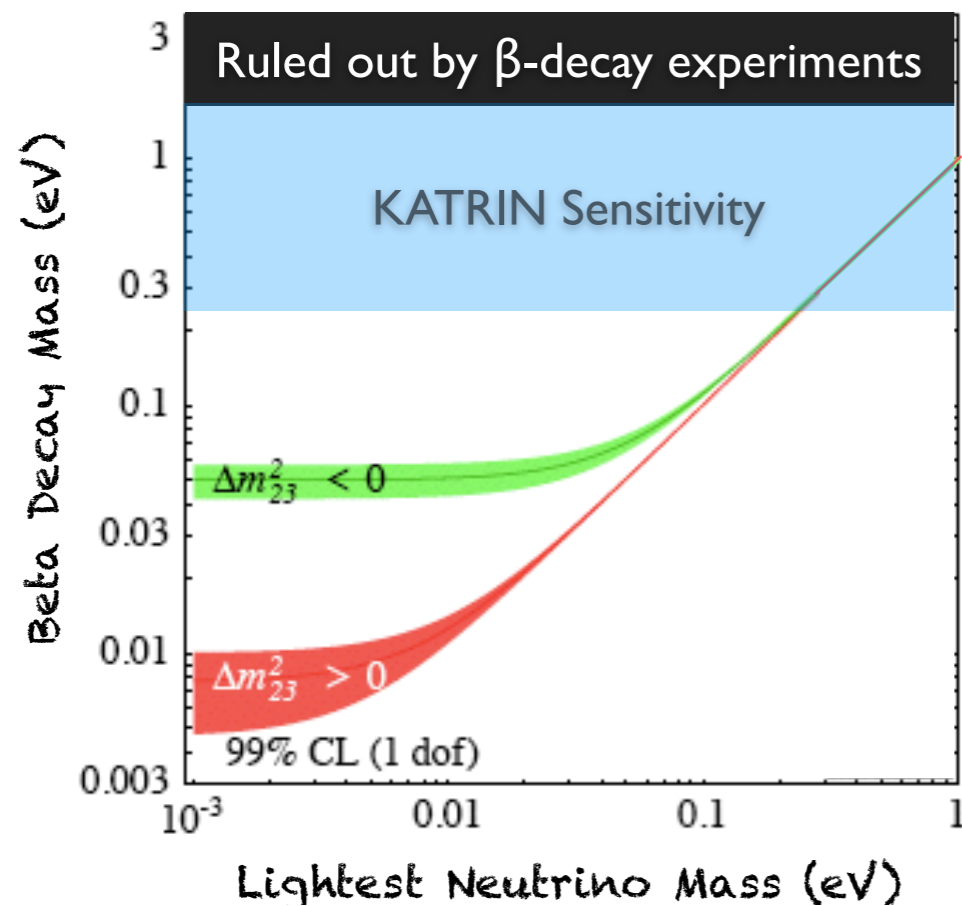
Discovery: 350 meV (at 5 σ)

Sensitivity: 200 meV (at 90% C.L.)

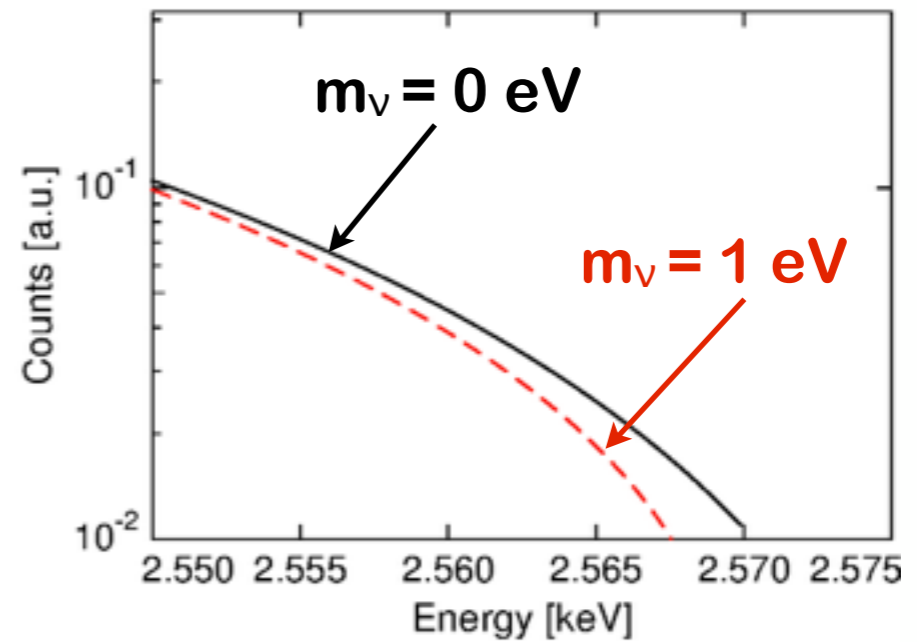
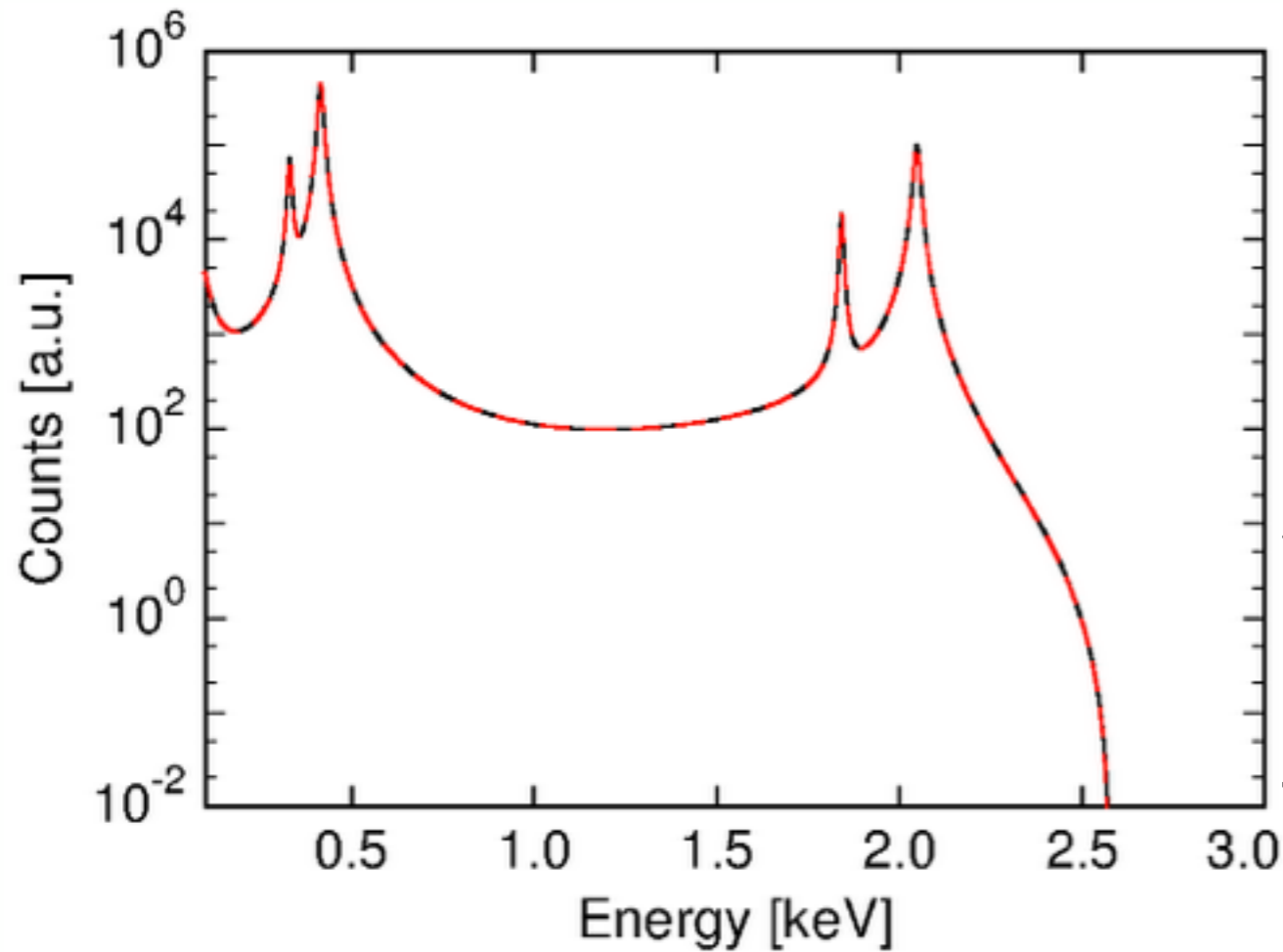
Partial loading in 2015.
Full Tritium Running in 2016.

Can we push further?

- Can direct measurements push to the inverted hierarchy scale?
- To do so, they must have better scaling law.



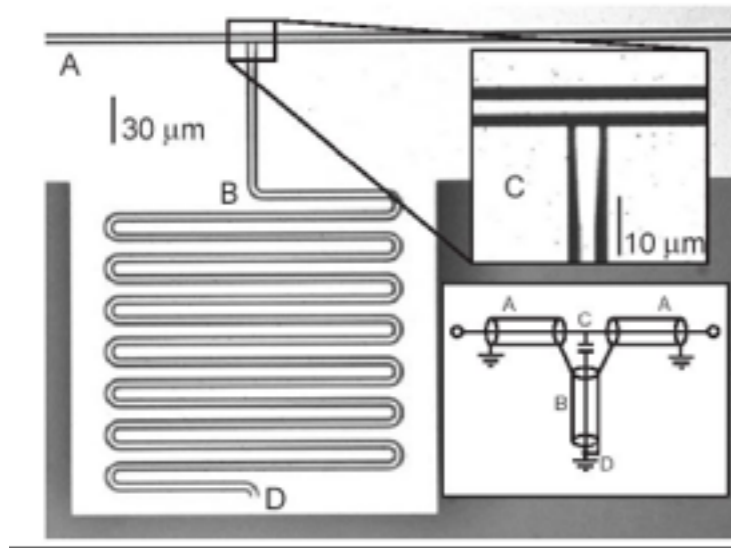
$$\dot{N} \sim (Q_{EC} - E_C)^2 \sum_i |U_{ei}|^2 \sqrt{1 - \frac{m_{\nu i}^2}{(Q_{EC} - E_C)^2}} \sum_H B_H \psi_H^2(0) \frac{\frac{\Gamma_H}{2\pi}}{(E_{EC} - E_H)^2 + \frac{\Gamma_H^2}{4}}$$



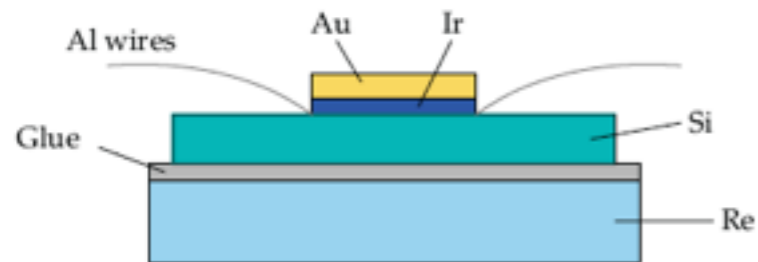
isotope
 New ~~kid~~ on the block:
 Electron Capture

The HOLMES Experiment

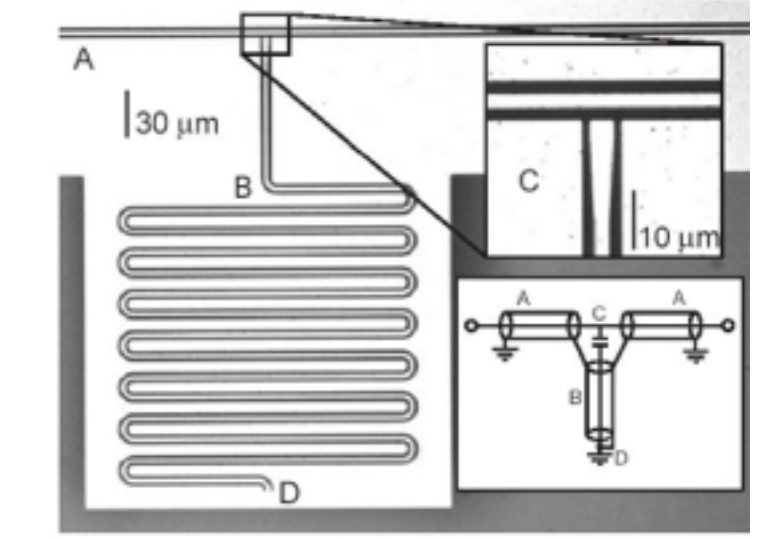
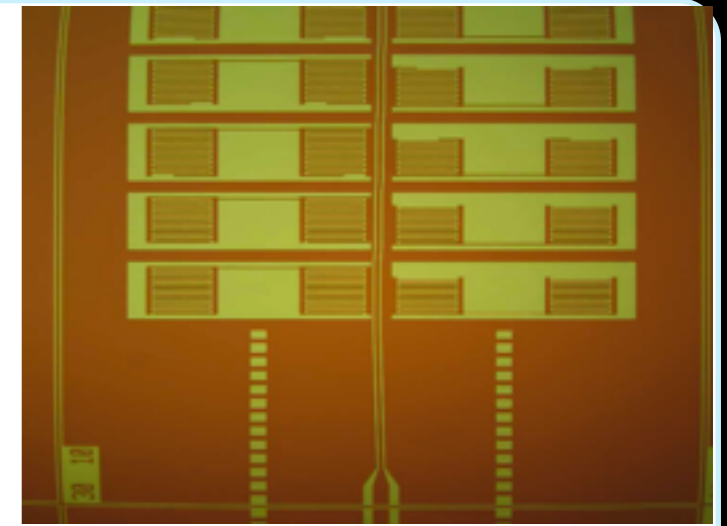
Technologies:



Superconducting Resonators



Transition Edge Sensors

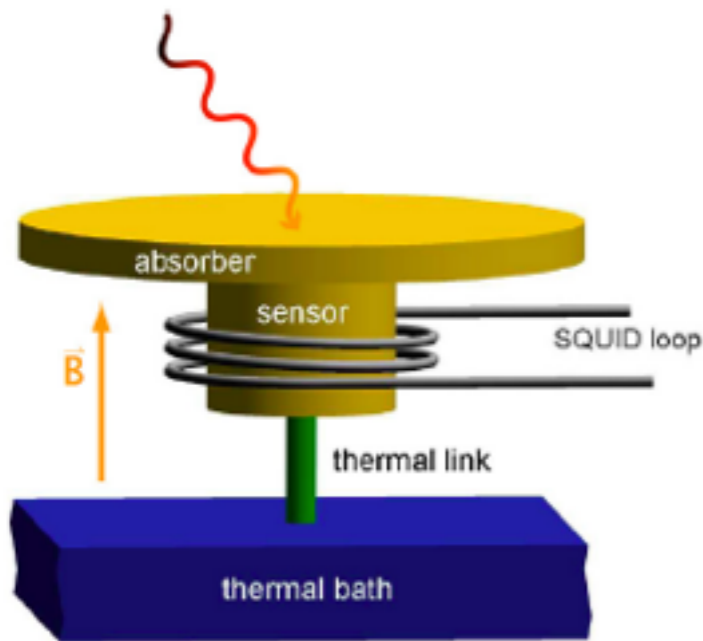


HOLMES

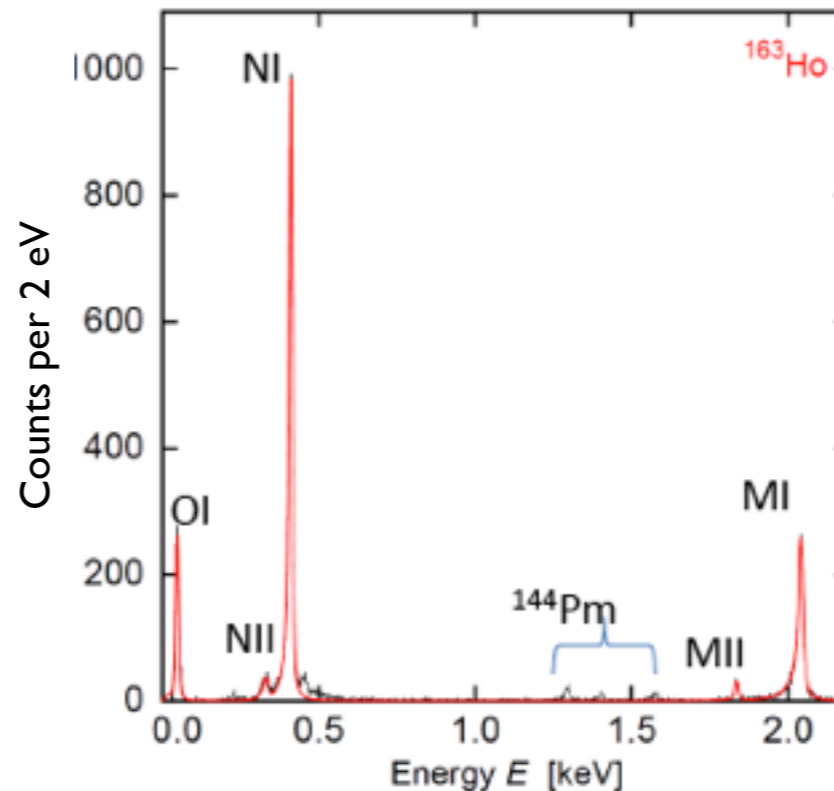
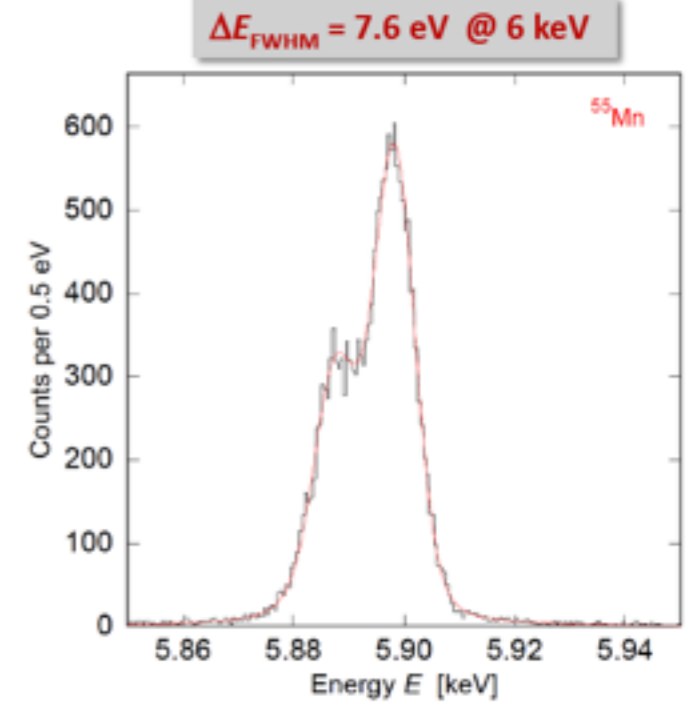
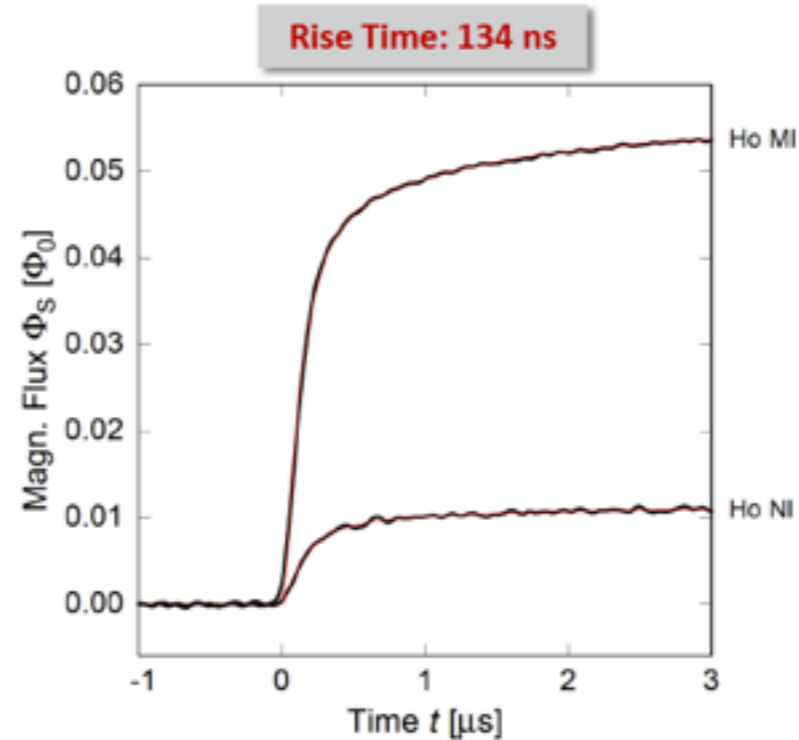
- MARE (Phase I) explored various technology approaches, such as Transition-Edge Sensors (TES) and Microwave Kinetic Inductance Detectors (MKIDs).
- Successful extraction of Ho^+ ions for metal production and implantation onto detectors.
- Successful funding received for one thousand channel Ho detector experiment (the HOLMES experiment).

The ECHO Experiment

Technology:



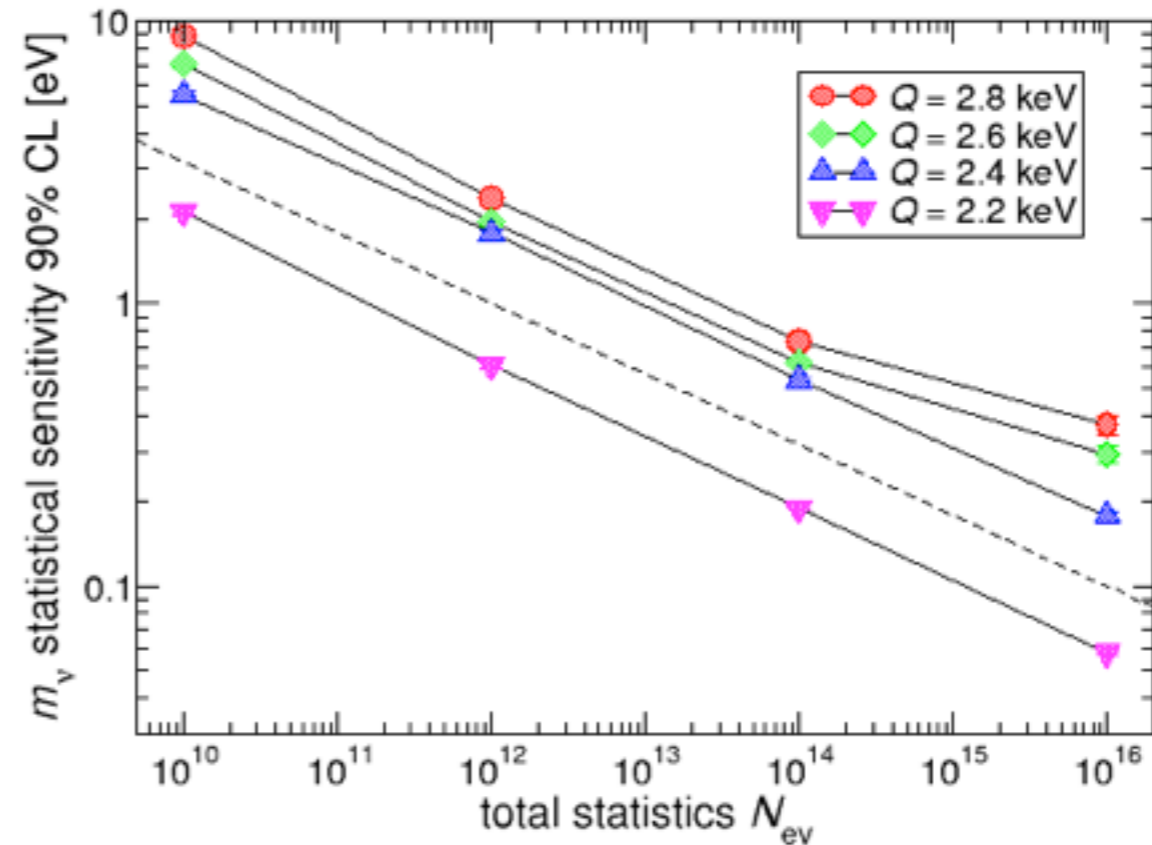
Metallic Magnetic Calorimeters



- The ECHO experiment uses metallic magnetic calorimeters to achieve goals.
- Fast rise times and good energy resolutions and linearity demonstrated.
- Endpoint measured at $2.80 \pm 0.08 \text{ keV}$.

Advantages & Challenges

Challenges:



Source Activity

$N_{ev} > 10^{14}$ to reach sub-eV level

Advantages:

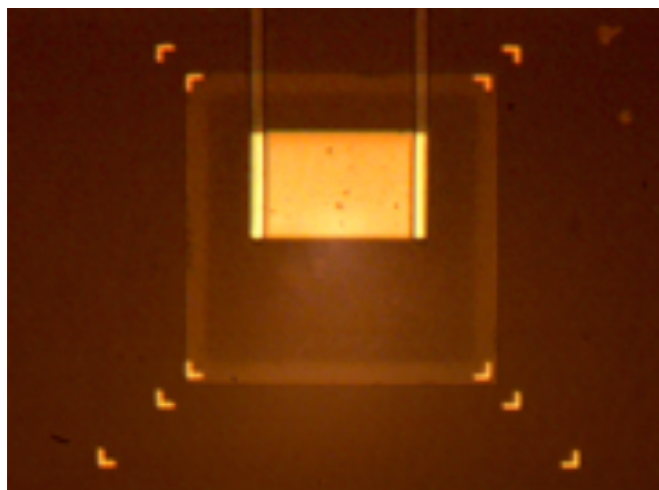
- Source = detector
- No backscattering
- No molecular final state effects.
- Self-calibrating

Detector Response

$\Delta E_{FWHM} < 10$ eV
Trisetime < 1 μ s

Experimental Challenges:

- Fast rise times to avoid pile-up effects.
- Good energy resolution & linearity
- Abundant isotope production



Project 8

Coherent radiation emitted can be collected and used to measure the energy of the electron in non-destructively.

PROJECT 8

Frequency Approach



I. I. Rabi



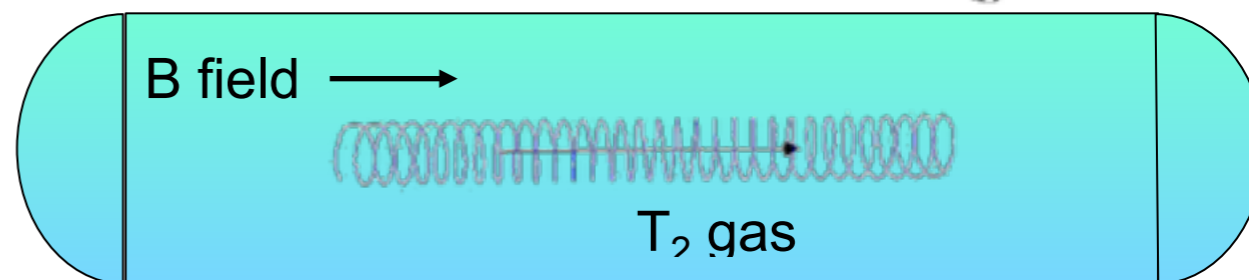
A. L. Schawlow

“Never measure anything but frequency.”

- Use cyclotron frequency to extract electron energy.

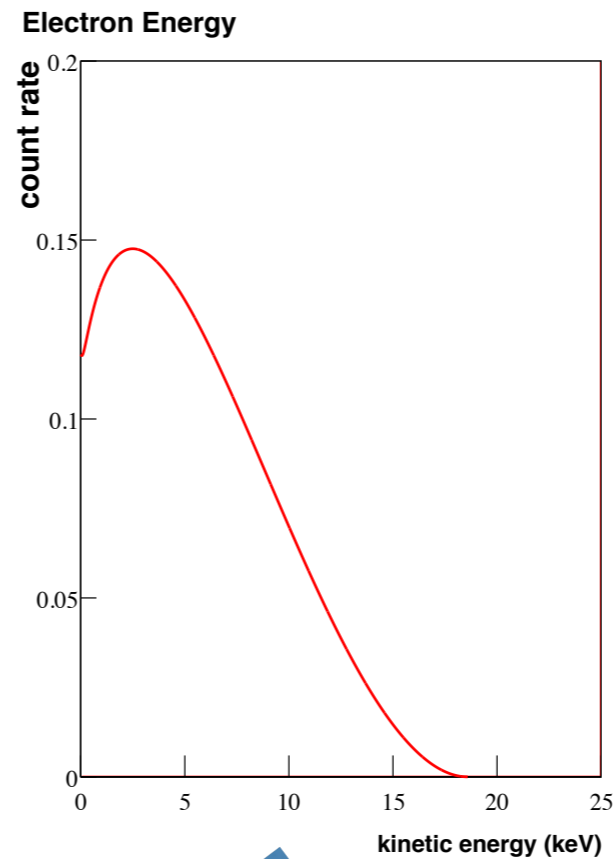
$$\omega(\gamma) = \frac{\omega_0}{\gamma} = \frac{eB}{K + m_e}$$

- Non-destructive measurement of electron energy.



Unique Advantages

- **Source = Detector**
(no need to separate the electrons from the tritium)
- **Frequency Measurement**
(can pin electron energies to well-known frequency standards)
- **Full Spectrum Sampling**
(full differential spectrum measured at once, large leverage for stability and statistics)

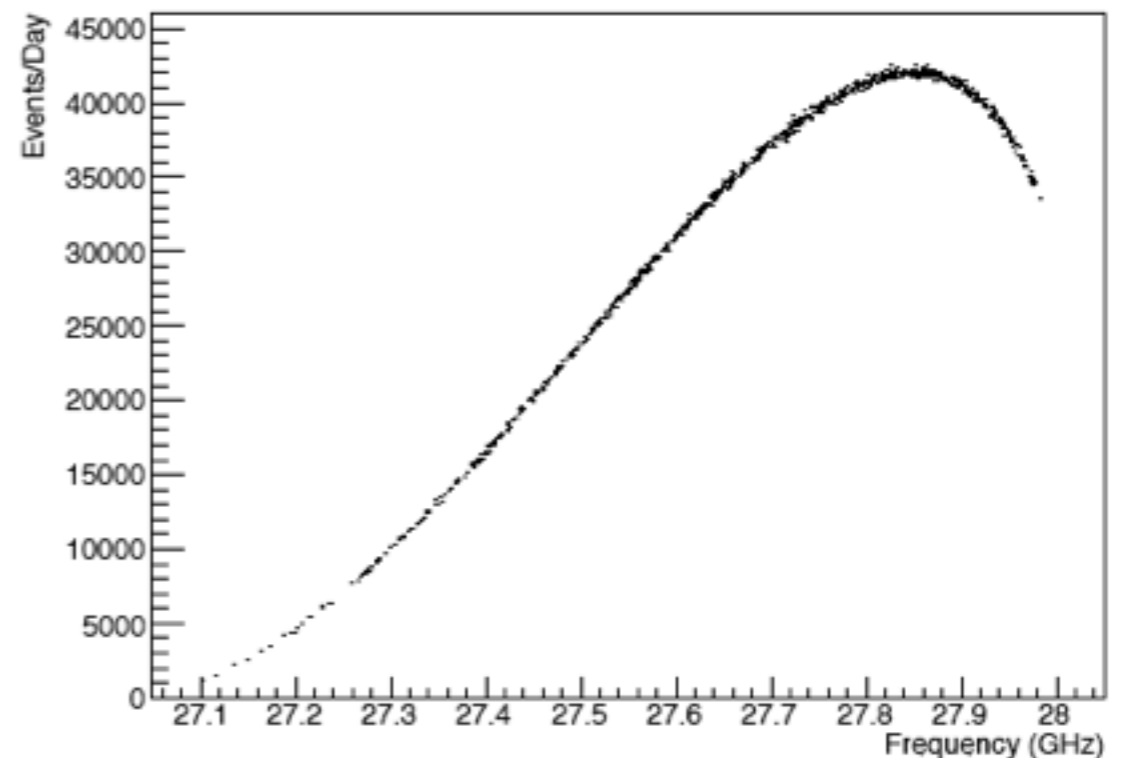


Beta spectrum



Cyclotron Frequency

$$\omega(\gamma) = \frac{\omega_0}{\gamma} = \frac{eB}{K + m_e}$$



Beta (frequency) spectrum

...and Challenges

⦿ Power Emitted

Less than 1 fW of power radiated (depends on antenna geometry) is challenging.

⦿ Confinement Period

One needs time to make sufficiently accurate measurement ($> 10 \mu\text{s}$).

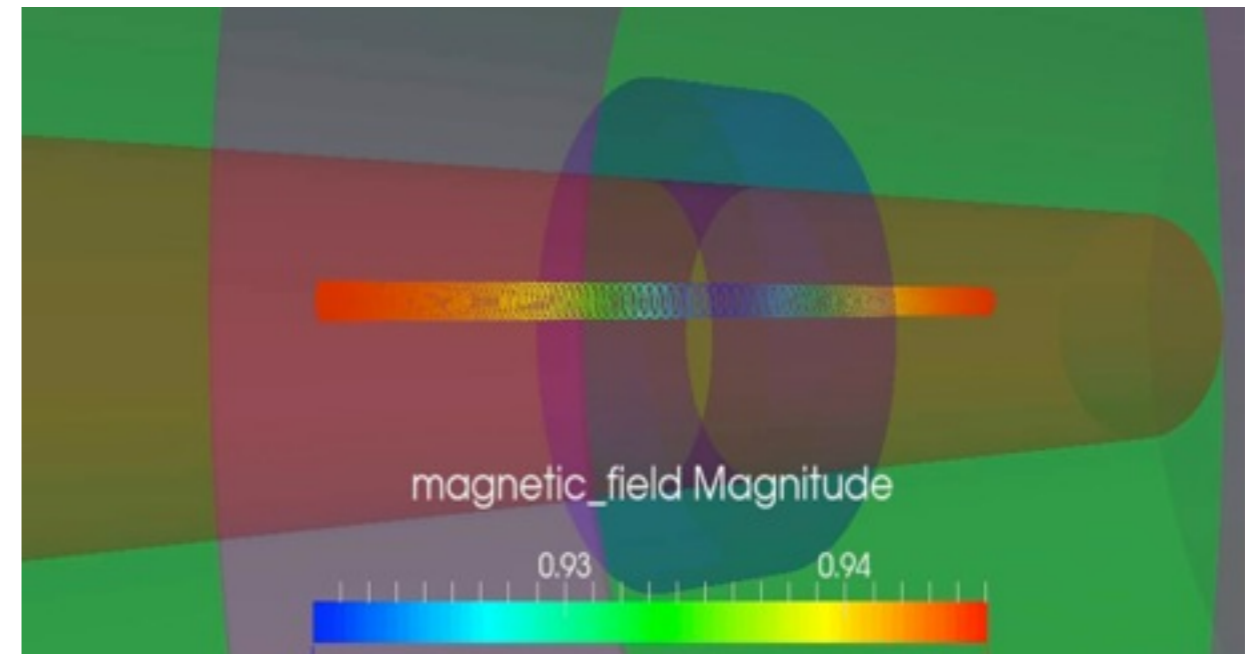
Employ magnetic bottle for trapping.

⦿ Full Spectrum

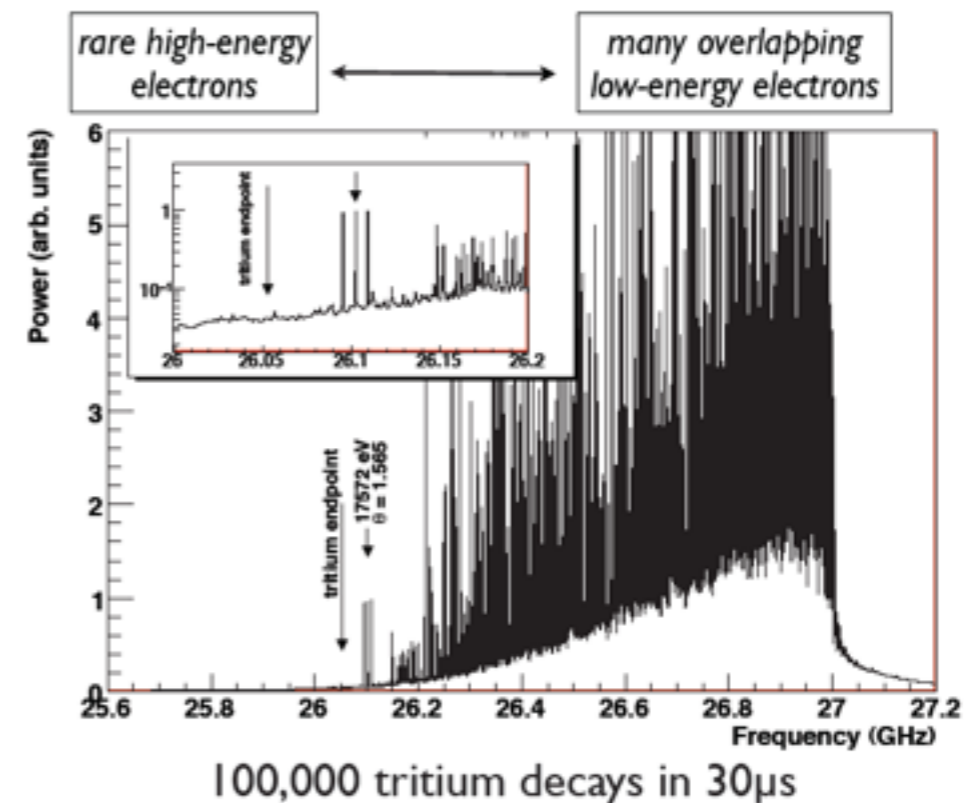
The full spectrum is available. Fortunately, linearity of frequency space helps separate regions of interest.

$$P_{\text{tot}}(\beta_{\parallel}, \beta) = \frac{1}{4\pi\epsilon_0} \frac{2e^2\omega_0^2}{3c} \frac{\beta_{\parallel}^2}{1 - \beta^2}$$

(Free) Radiative Power Emitted



Simulation of electron motion in magnetic bottle



Simulation of beta (frequency) spectrum

Project 8 Collaboration



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P. Monhanmurthy*
Massachusetts Institute of Technology

R. Bradley
National Radio Astronomy Observatory

T. Thuemmler
Karlsruhe Institute of Technology

R. Patterson
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Shepard Doelman, Alan Rogers
Haystack Observatories

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University of California, Santa Barbara

J. Kofron*, L. McBride*, R.G.H. Robertson, Leslie Rosenberg, Gray Rybka
University of Washington, Center for Experimental Nuclear Physics and Astrophysics

D.M. Asner, J. Fernandes, A.M. Jones, J.F. Kelly, B.A. VanDevender
Pacific Northwest National Lab

* indicates graduate student

A Phased Approach

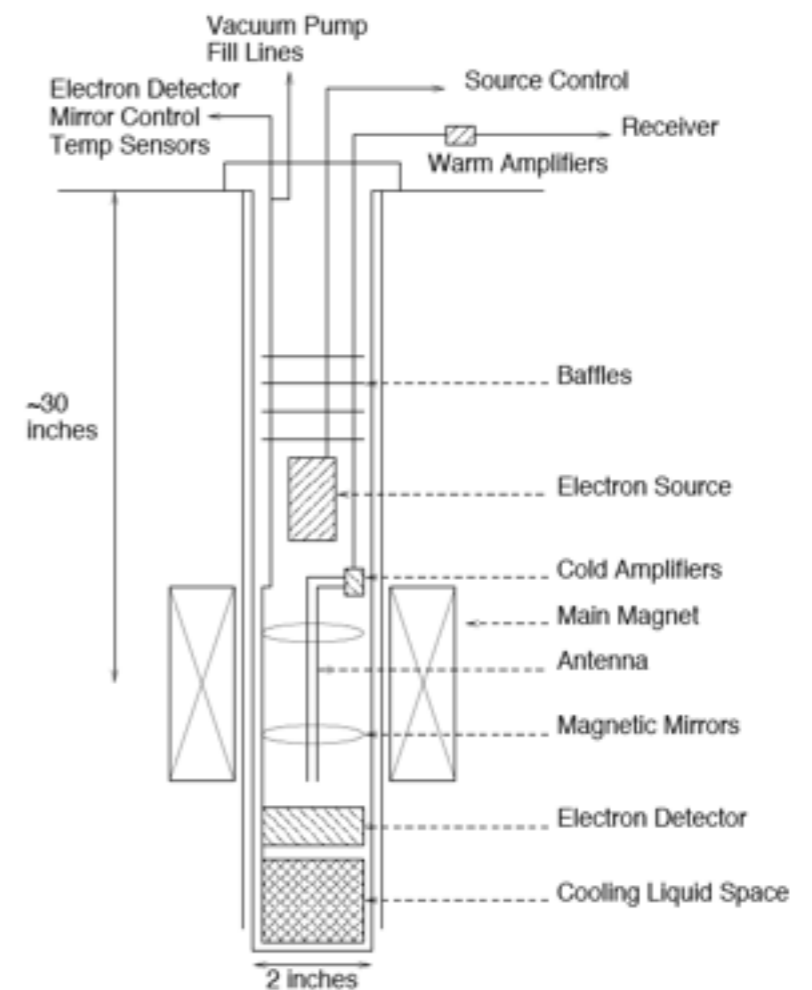
Given the novelty of the project, we are pursuing a phased approach toward neutrino mass measurements:

	Timeline	Scientific Goal	Source	R&D Milestone
Phase I	2010-2014	Proof of principle; Kr spectrum	^{83}m	Single electron detection
Phase II	2014-2016	T-He mass difference	T	Tritium spectrum; calibration and error studies
Phase III	2016-2018	0.2 eV scale	T	
Phase IV	2018+	0.05 eV scale	T	High rate sensitivity

We have commenced Phase I, we are designing Phase II

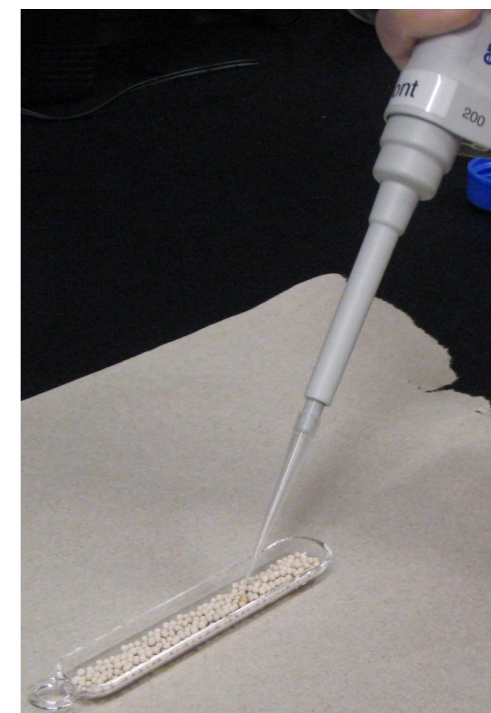
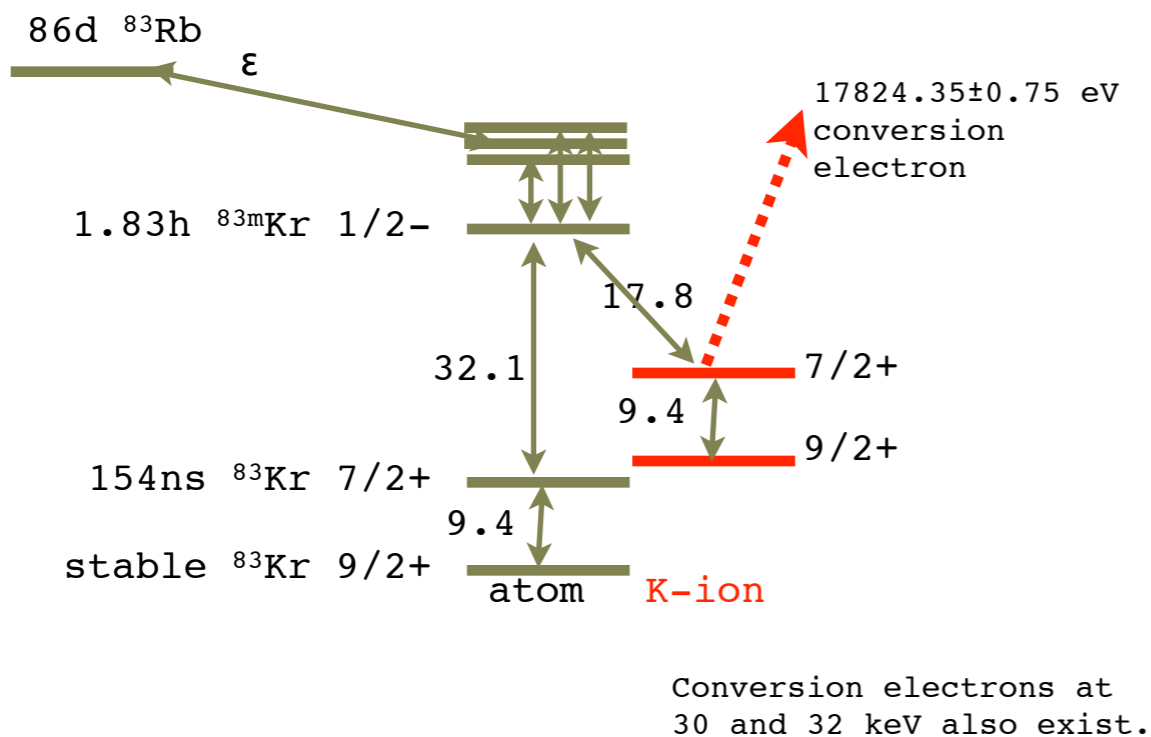
Basic Layout of Phase I

- ① **Gas/Electron System**
Provides mono-energetic electrons for signal detection.
- ② **Magnet System**
Provides magnetic field and trapping of electrons.
- ③ **RF Detection/Calibration System**
Detection of microwave signal.

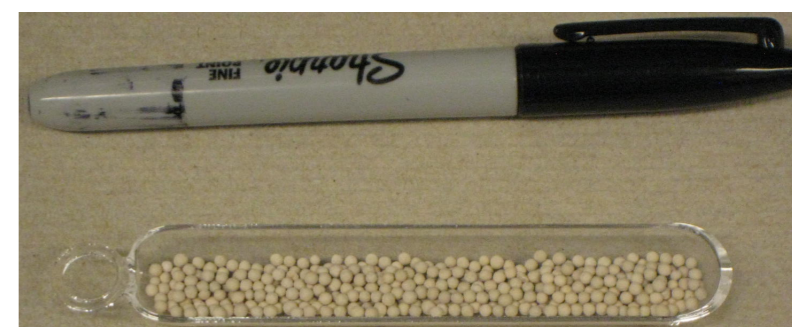


The Electron Source

Initial Demonstration Source: ^{83m}Kr



Zeolite
loading



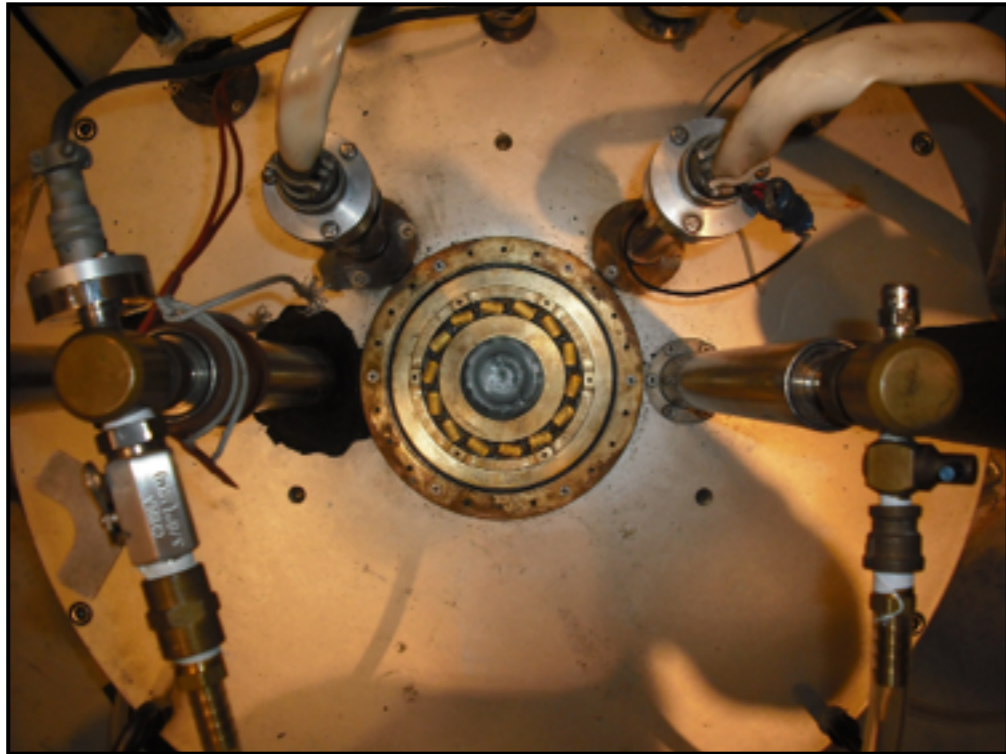
Mono-energetic gaseous electron source

Collaboration taking a phased approach to understand the scaling and systematics of the experiment.

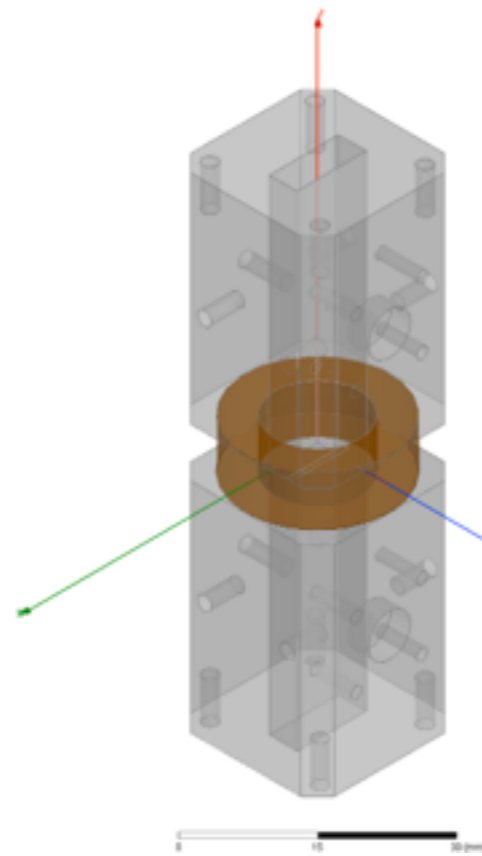
First phase (single electron detection) requires single electron detection.

Using ^{83m}Kr (^{83}Rb implanted in zeolite beads) as source

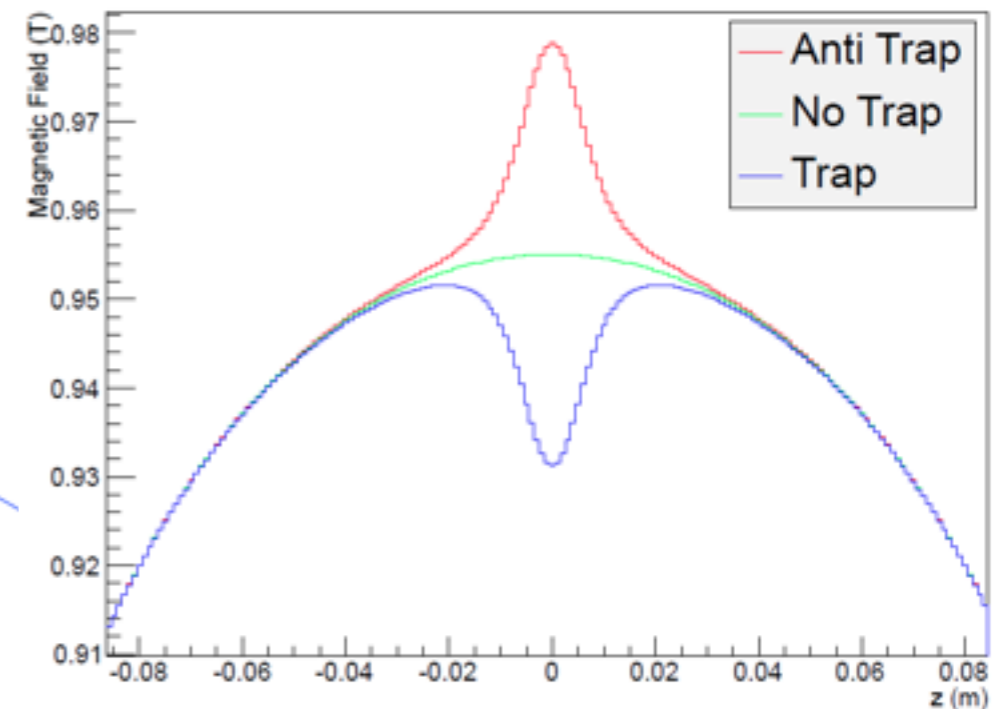
The Magnet System



Main Superconducting Magnet
1 T field (27 GHz)



~100 G Trapping coil



Trapping / Anti-trapping
configuration

Using 0.94 Tesla field, where signal occurs at ~26 GHz.

Electrons are trapped in a small "magnetic bottle" provided by a small trapping coil.

Trapping coil provides means to observe electron for sufficient time to extract signal.

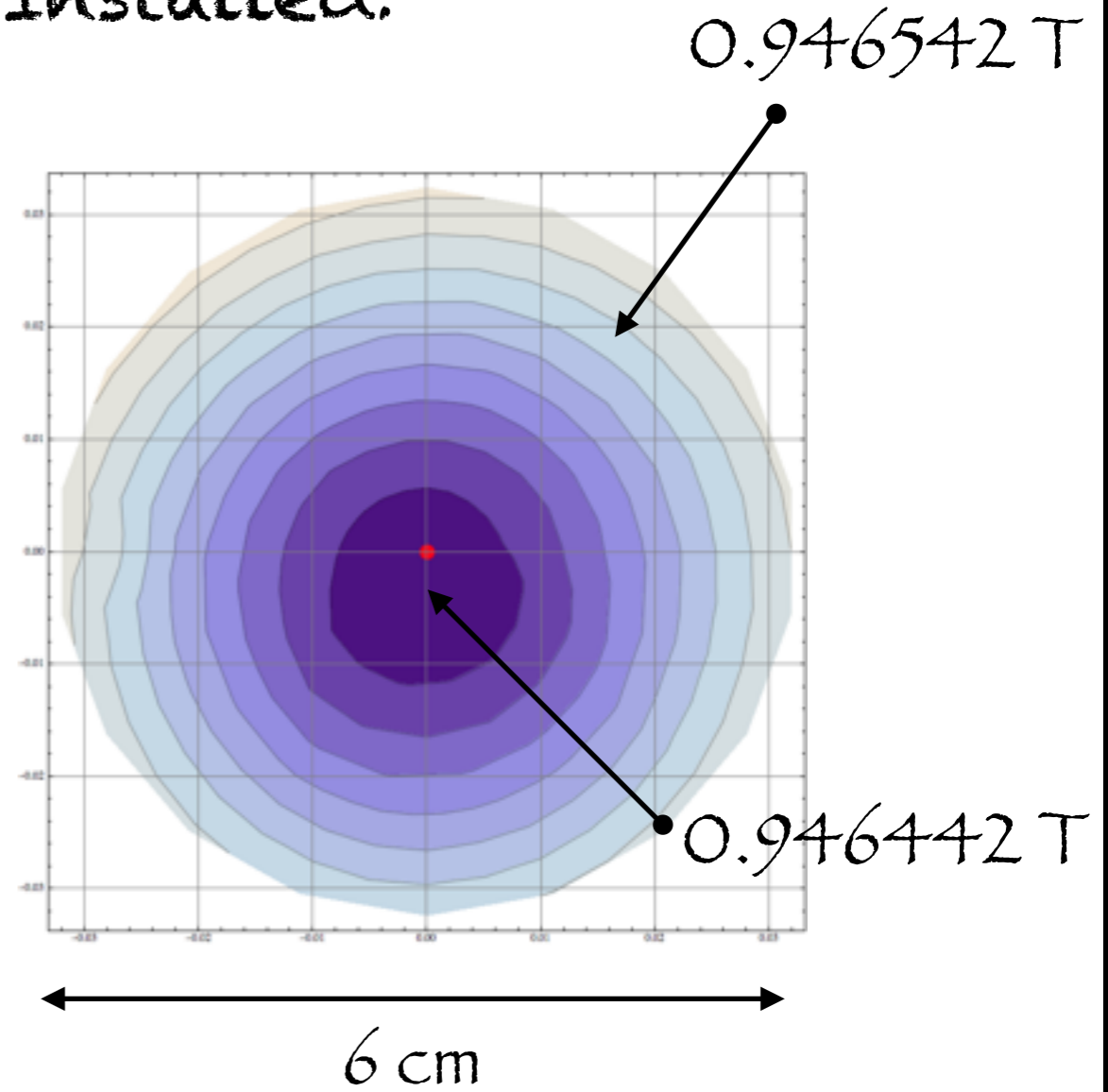
New NMR Magnet Installed!



Old
"Green Giant"



New Bruker
NMR Magnet

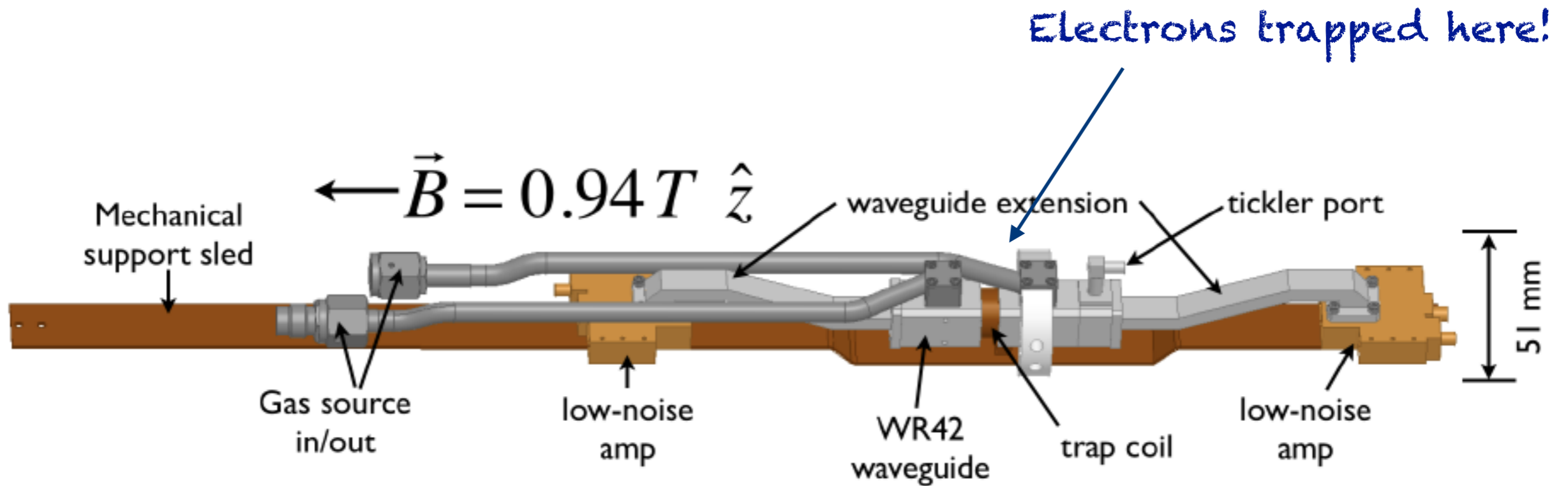


As of 2014, new NMR magnet installed for prototype.

Magnet much easier to use and maneuver for testing.

Improved field homogeneity and stability.

The RF Detection System

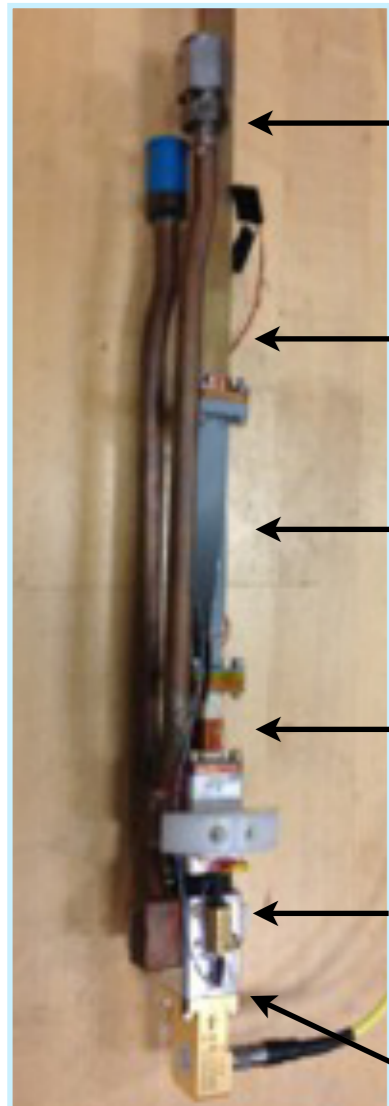


Basic Layout of "Detector"

Fundamental frequency is carried via WR-42 waveguide to low-noise amplifiers.

Signal is mixed down to MHz range, further amplified, and then digitized to data for analysis.

The RF Detection System



Gas source
in/out

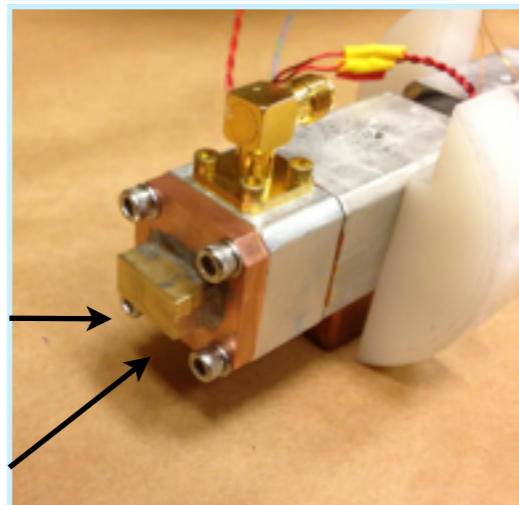
WR-42
waveguide

WR-42 twist

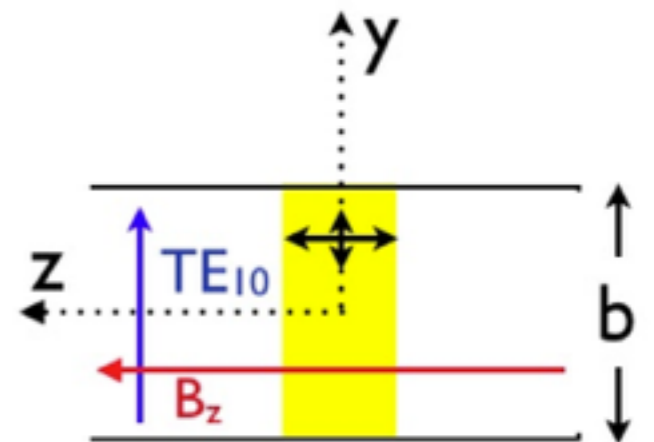
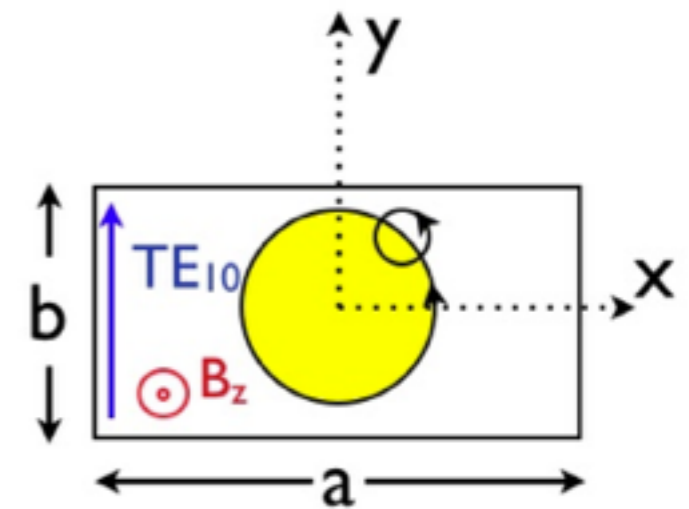
DPPH source

Trapping section

Tickler port



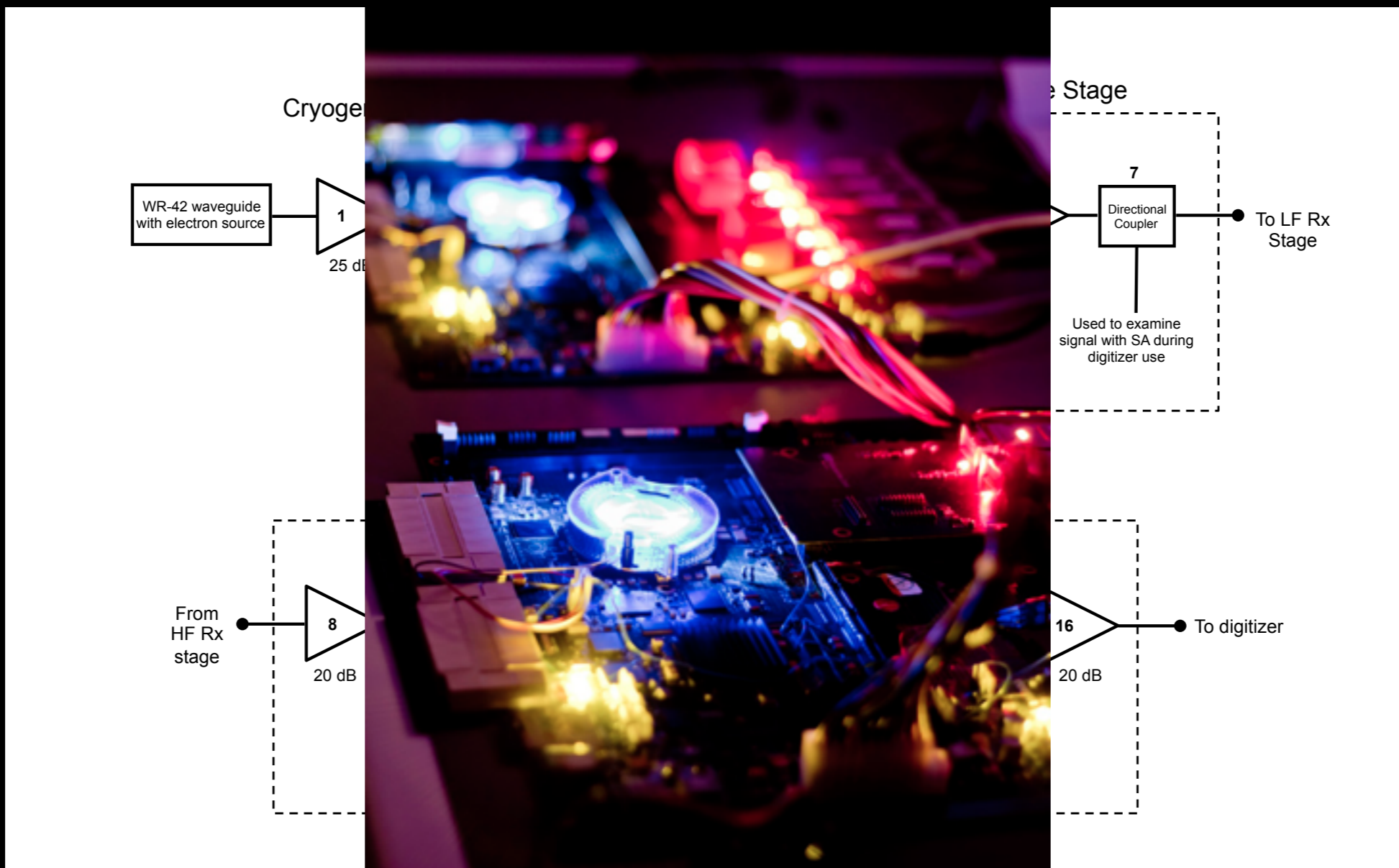
RF
waveguide
used as
main
chamber



Fundamental frequency is carried via WR-42 waveguide to low-noise amplifiers.

Signal is mixed down to MHz range, further amplified, and then digitized to data for analysis.

The Signal Chain & Digitization

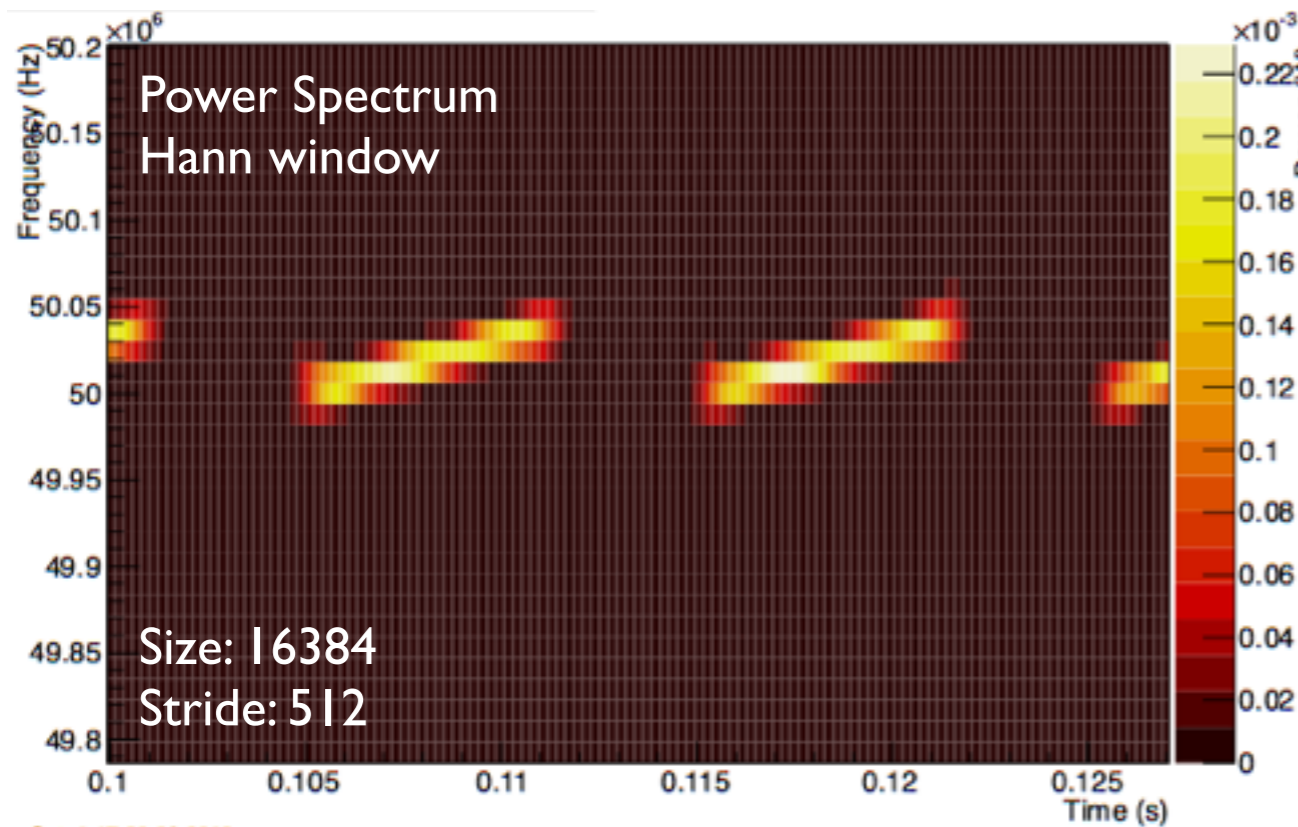


Fundamental frequency is carried via WR-42 waveguide to low-noise amplifiers.

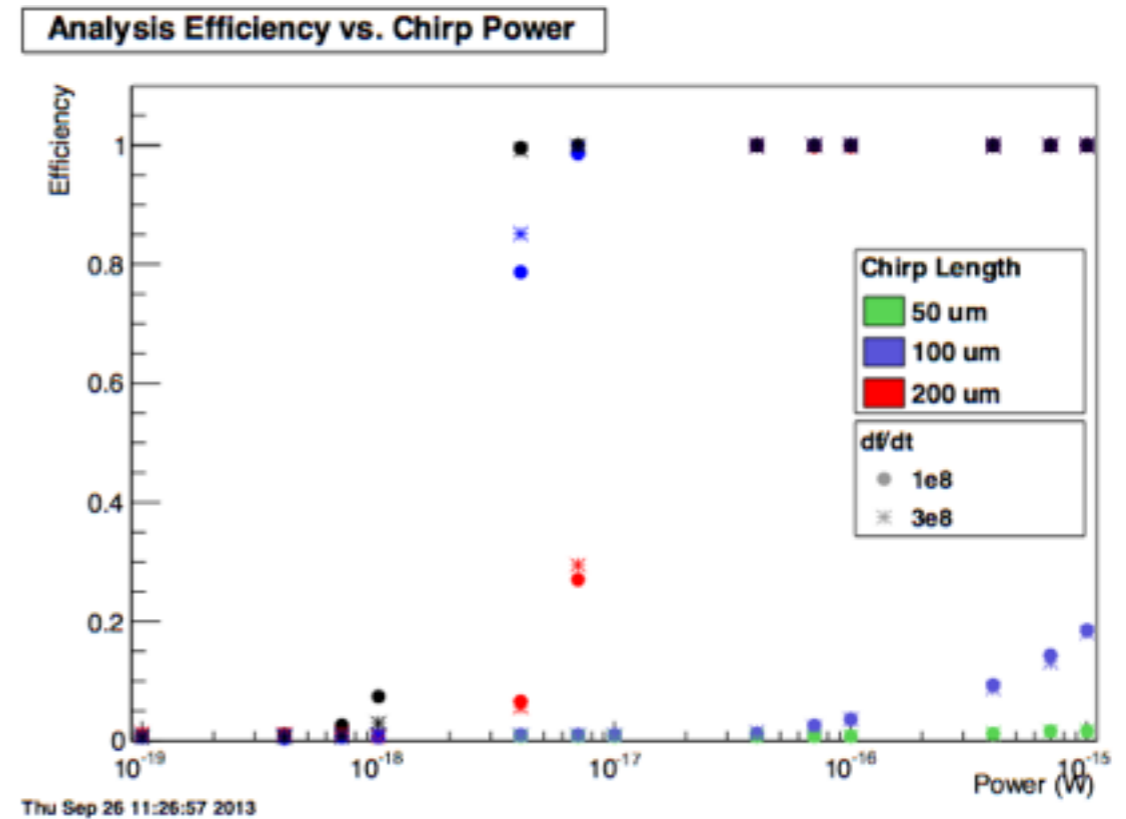
Signal is mixed down to MHz range, further amplified, and then digitized to data for analysis.

Developing CASPER-ROACH system for future triggering and data processing.

What Does a Signal Look Like?



Simulated "Chirp" Signal



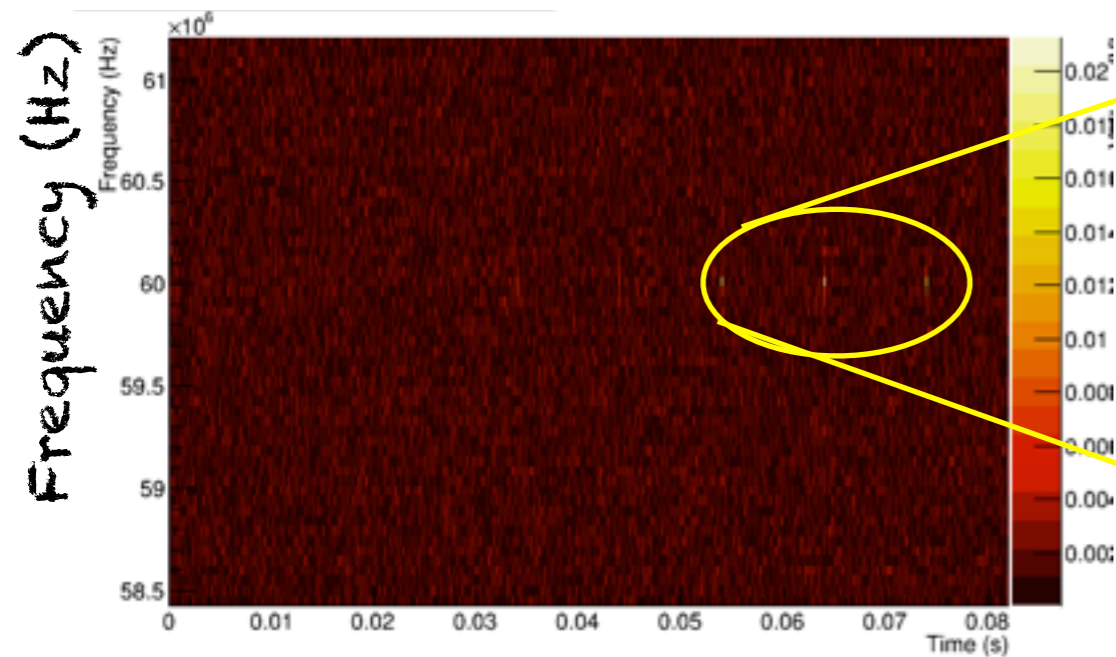
Efficiency for finding candidates

Signal exhibits itself as a short duration pulse ("chirp") in frequency.

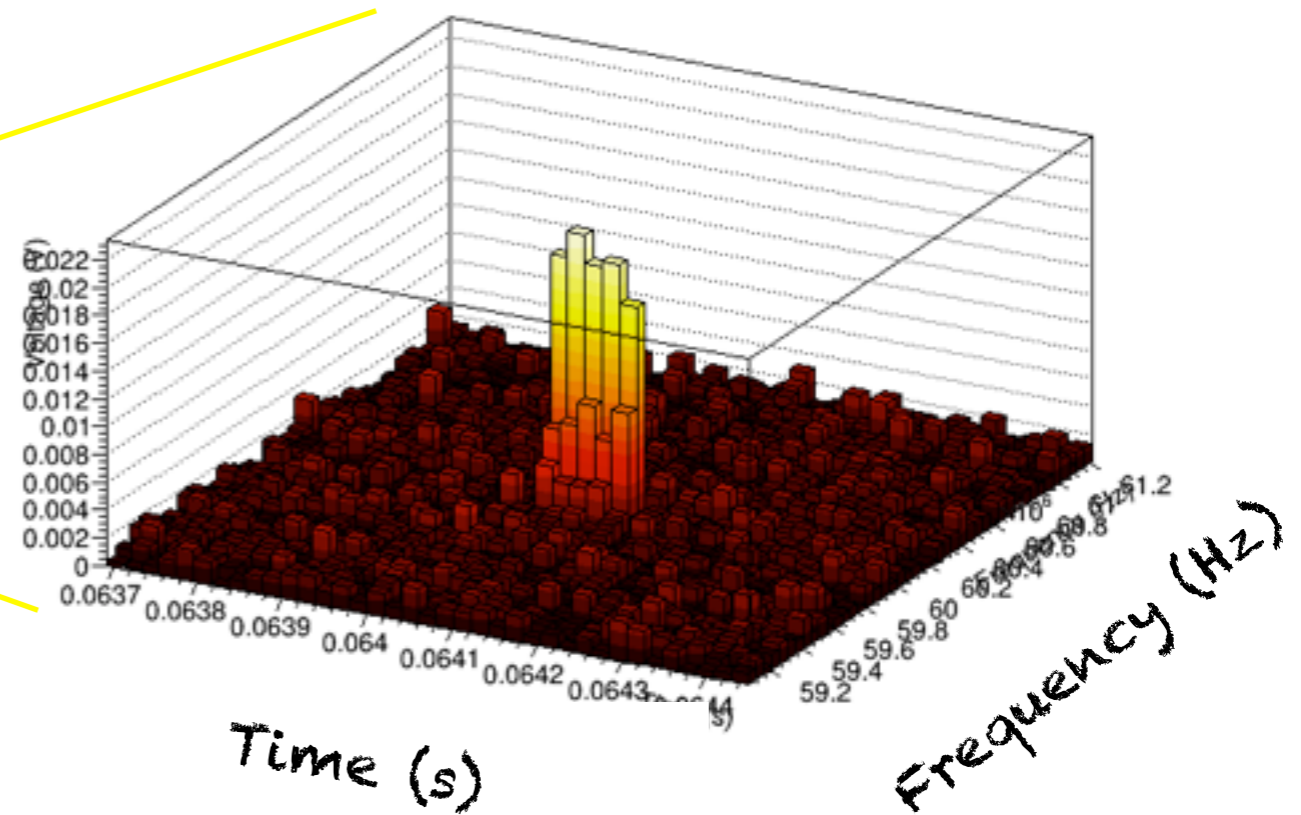
Signal is digitized in frequency/time space and scanned for clusters.

Sensitivity studies demonstrate excellent efficiency down to 10^{-17} fW for 50 μs signals.

Calibrating the Signal



Time (s)



Time (s)

Frequency (Hz)

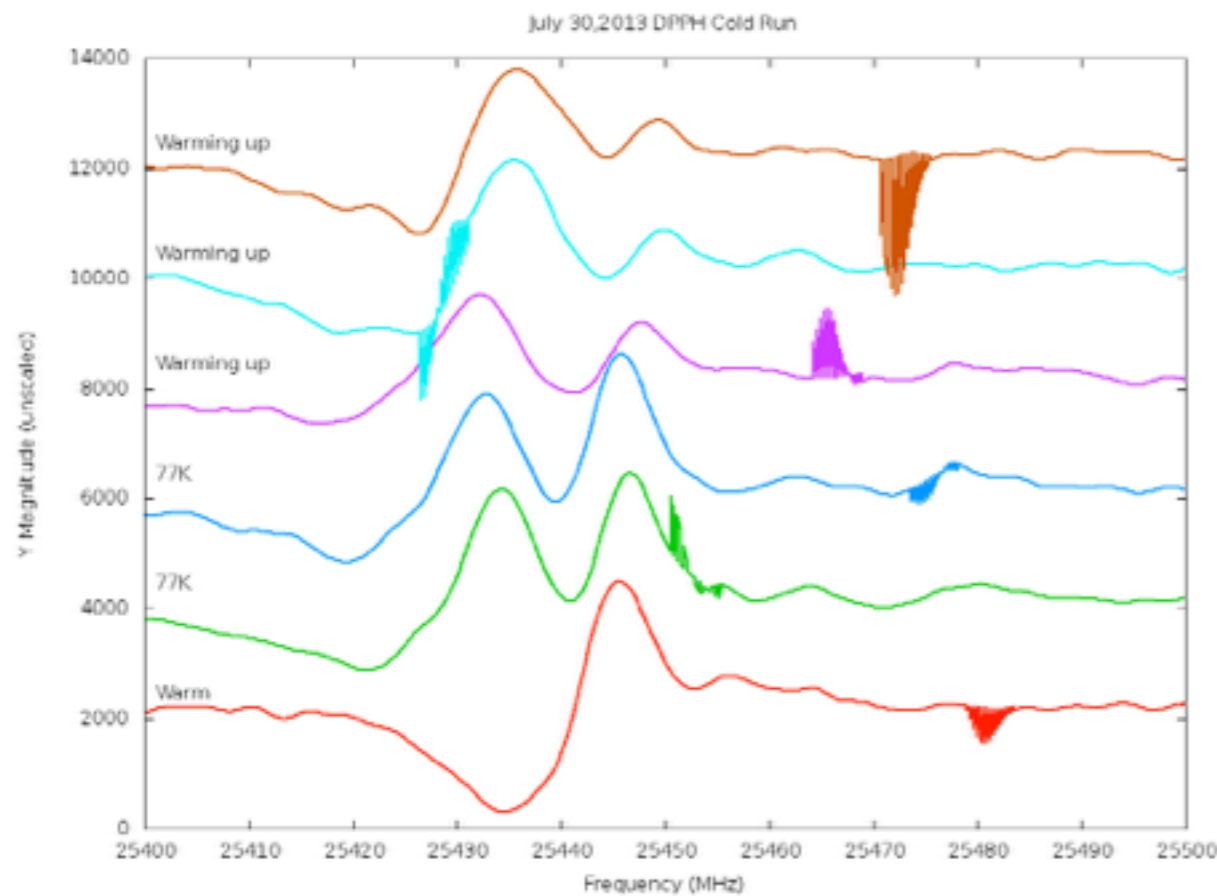
Injected Fake Signal

We have the ability to insert artificial chirps and tones into our cavity.

Using an RF switch, the fake-signal can be "pulsed".

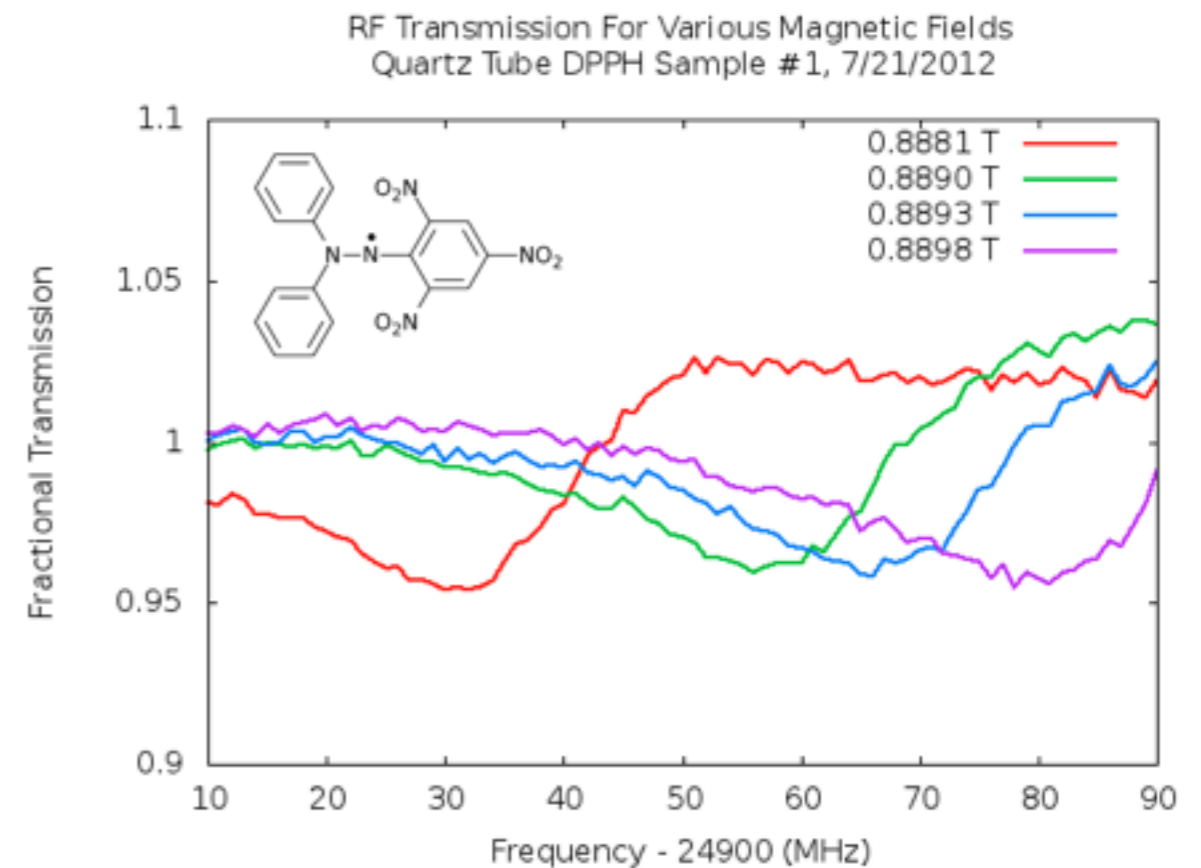
Tests system's sensitivity to amplitude and duration.

Radio-frequency (RF) Calibration



In-situ measurement of magnetic field
via DPPH absorption

ESR Measurement



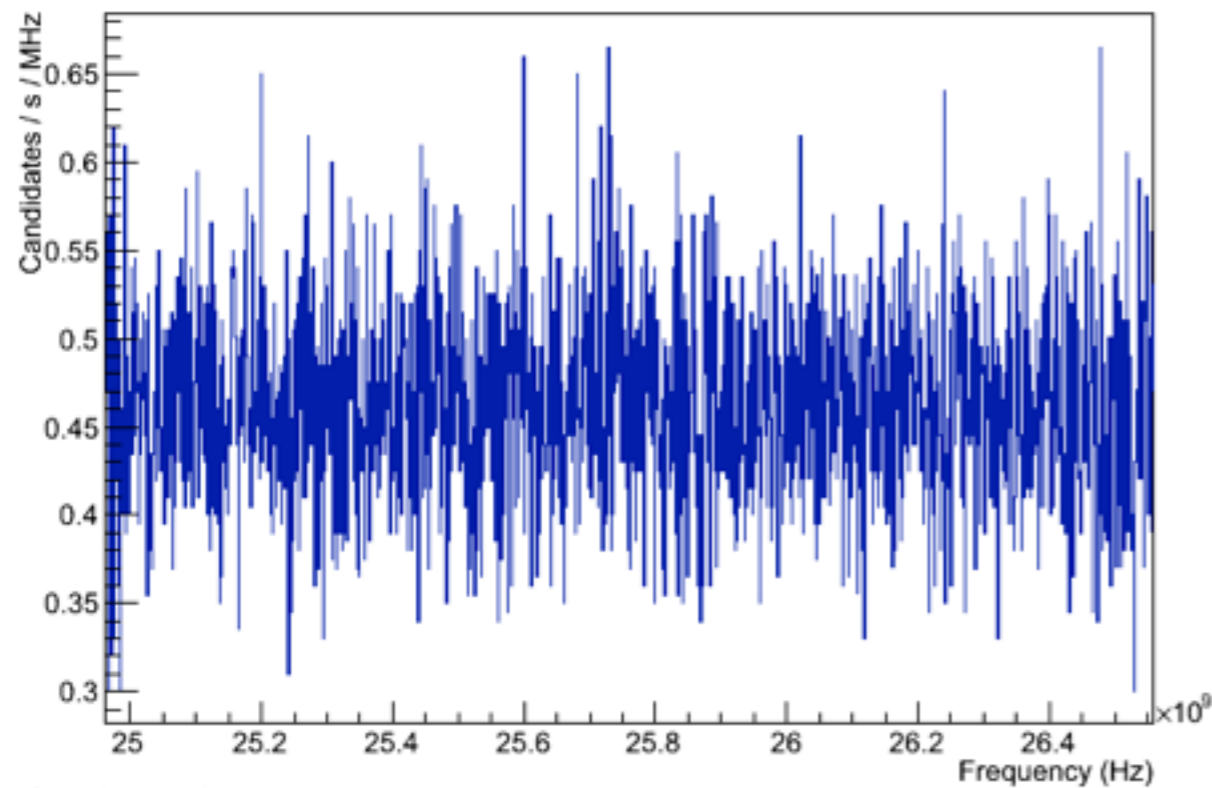
Magnetic field monitoring

Can use absorption of RF line from ESR to monitor field strength in-situ.

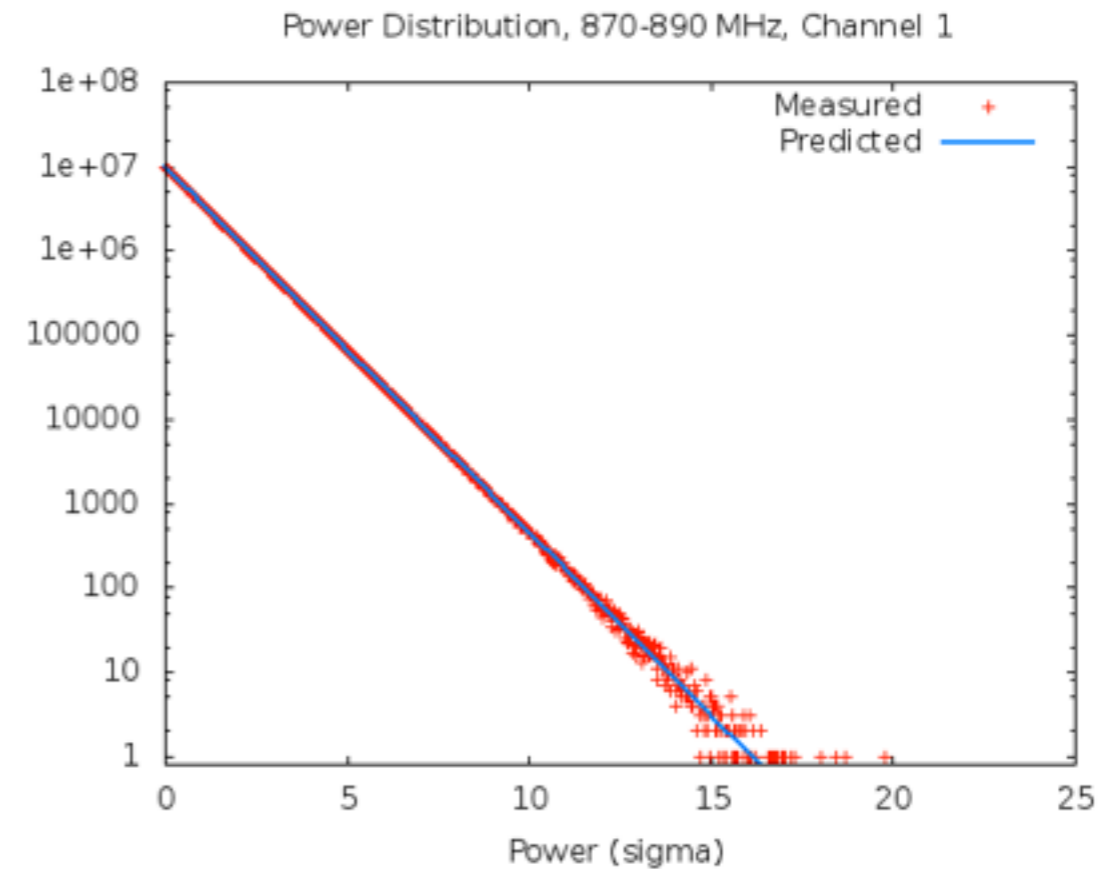
Using DPPH source for testing of RF system + magnetic field monitoring.

System works well in new NMR magnet.

Testing System on the January Data



Candidate background spectrum

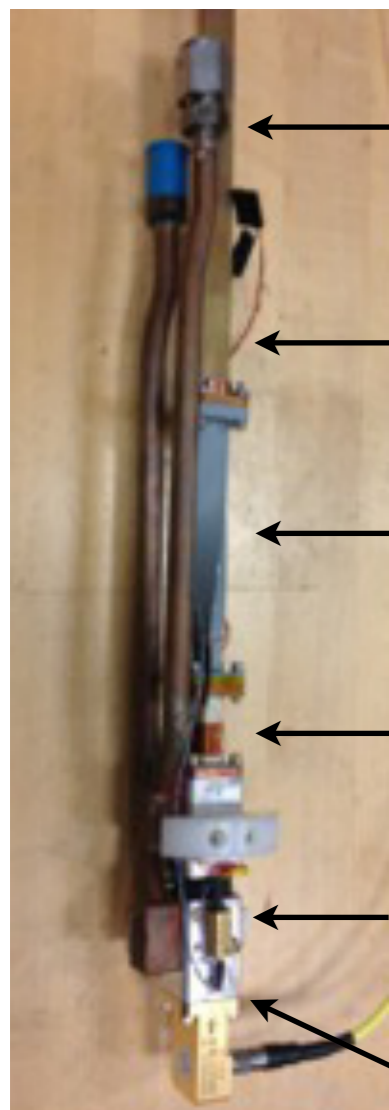


Thermal noise spectrum, January run
(Not a fit)

Short (8 hour) test run taking in January, 2013.

January data with rather large (~150 K) temperatures as initial test of system.

Preparing for Run II



Gas source
in/out

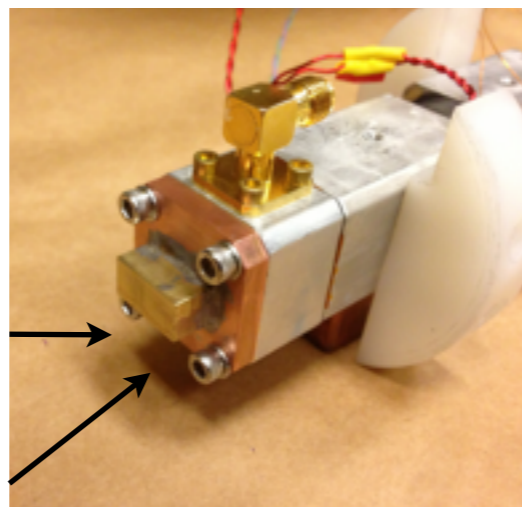
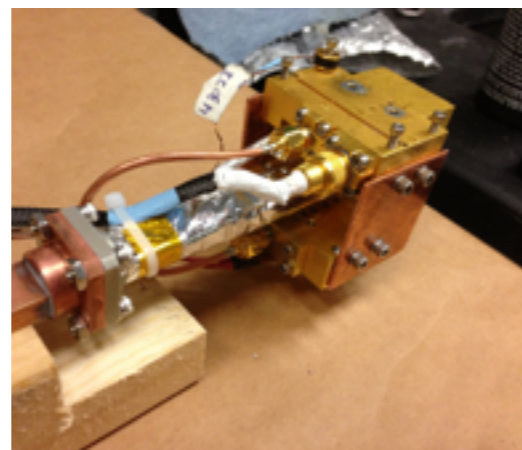
WR-42
waveguide

WR-42 twist

DPPH source

Trapping section

Tickler port



System

Specification

Achieved

Magnetic
Field

1 T Field
10

< few 10
DPPH
Monitoring

Gas System

10

< 10
PIPS detector

Noise
Temperature

T

T

Sensitivity
Analysis

SNR > 10 for
0.5 fW signal

SNR > 10 dB
for 0.1 fW

A new run is planned with a target noise temperature down to ~ 50 K.
System undergoing commissioning now.

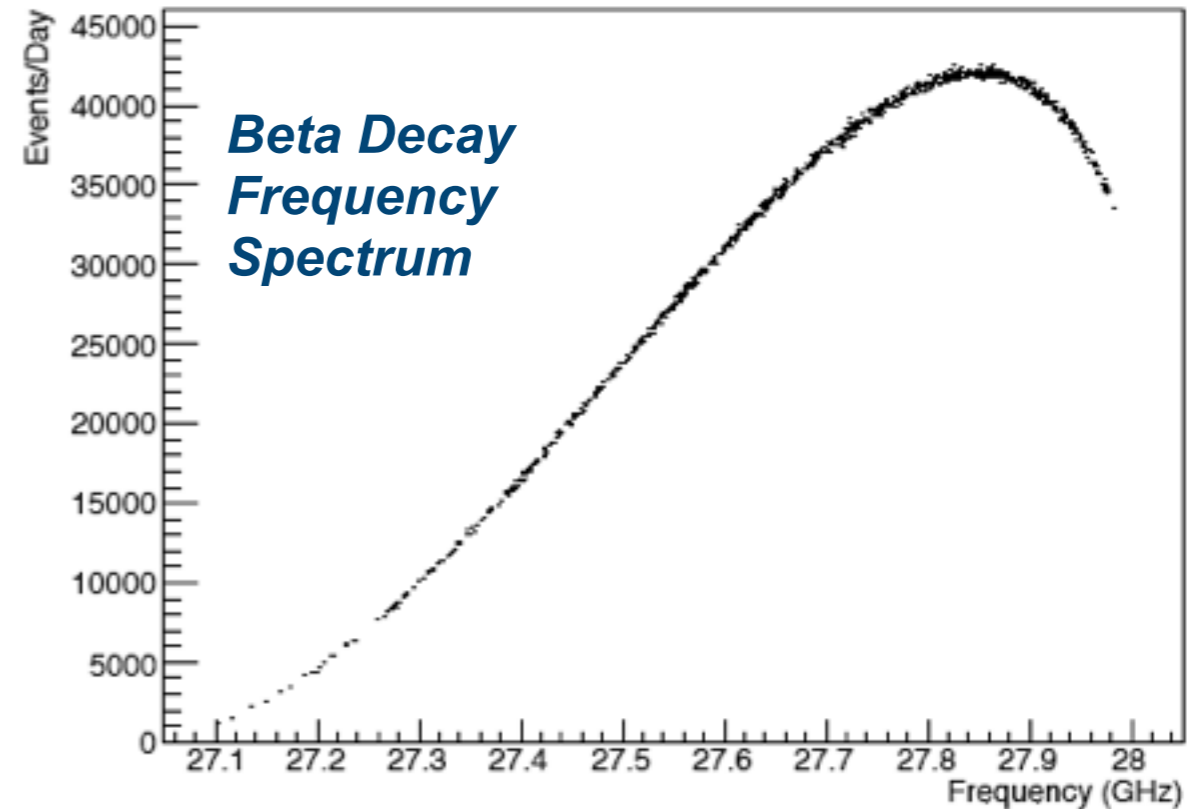
Analysis & simulation show signal efficiency of $> 90\%$ for electrons down
Sensitivity of < 0.1 fW and > 50 μ s trapping time.



Attempting to
see the
inverted
scale...

Sensitivity to Neutrino Masses

- There are distinct advantages that are specific to frequency-based measurements:
- You get the **entire spectrum** (and background) at once.
- The background is **extremely small**:
 - There is no detector.
 - There might not even be any surfaces.
 - Cosmic ray interactions and radioactive backgrounds are interacting with a gas, very little target material.



$$\frac{dN}{dE_e} = 3rt(E_0 - E_e)[(E_0 - E_e)^2 - m_\nu^2]^{1/2}$$

Mass sensitivity depends on:

Target activity (volume x density)

Background

Field homogeneity

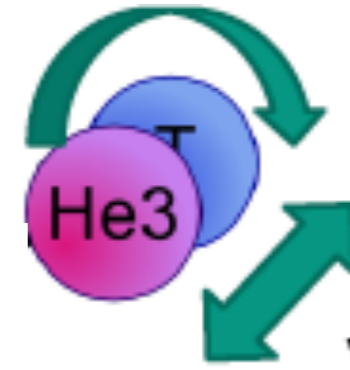
Lifetime of electron in trap (density)

Final states, doppler shifts, temperature

Moving Beyond the Degeneracy Scale

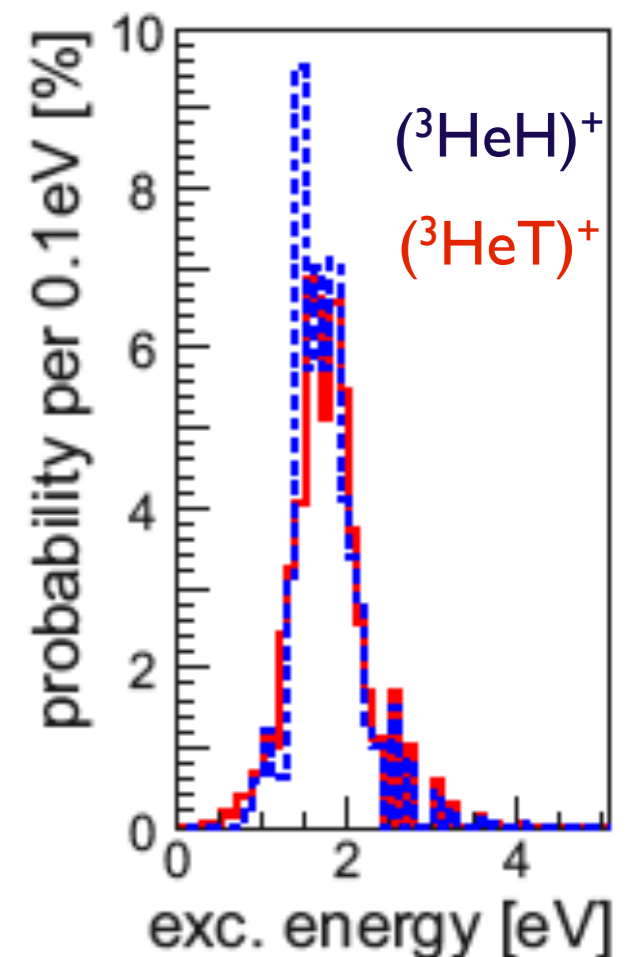
- Most effective tritium source achieved so far involves the use of gaseous molecular tritium.
- Method will eventually hit a resolution "wall" which is dictated by the rotational-vibrational states of T_2 . This places a resolution limit of 0.36 eV.
- One needs to either switch to (extremely pure) atomic tritium or other isotope with equivalent yield.
- The trapping conditions necessary for electrons also lends itself for atomic trapping of atomic tritium (R. G. H. Robertson)

rotational

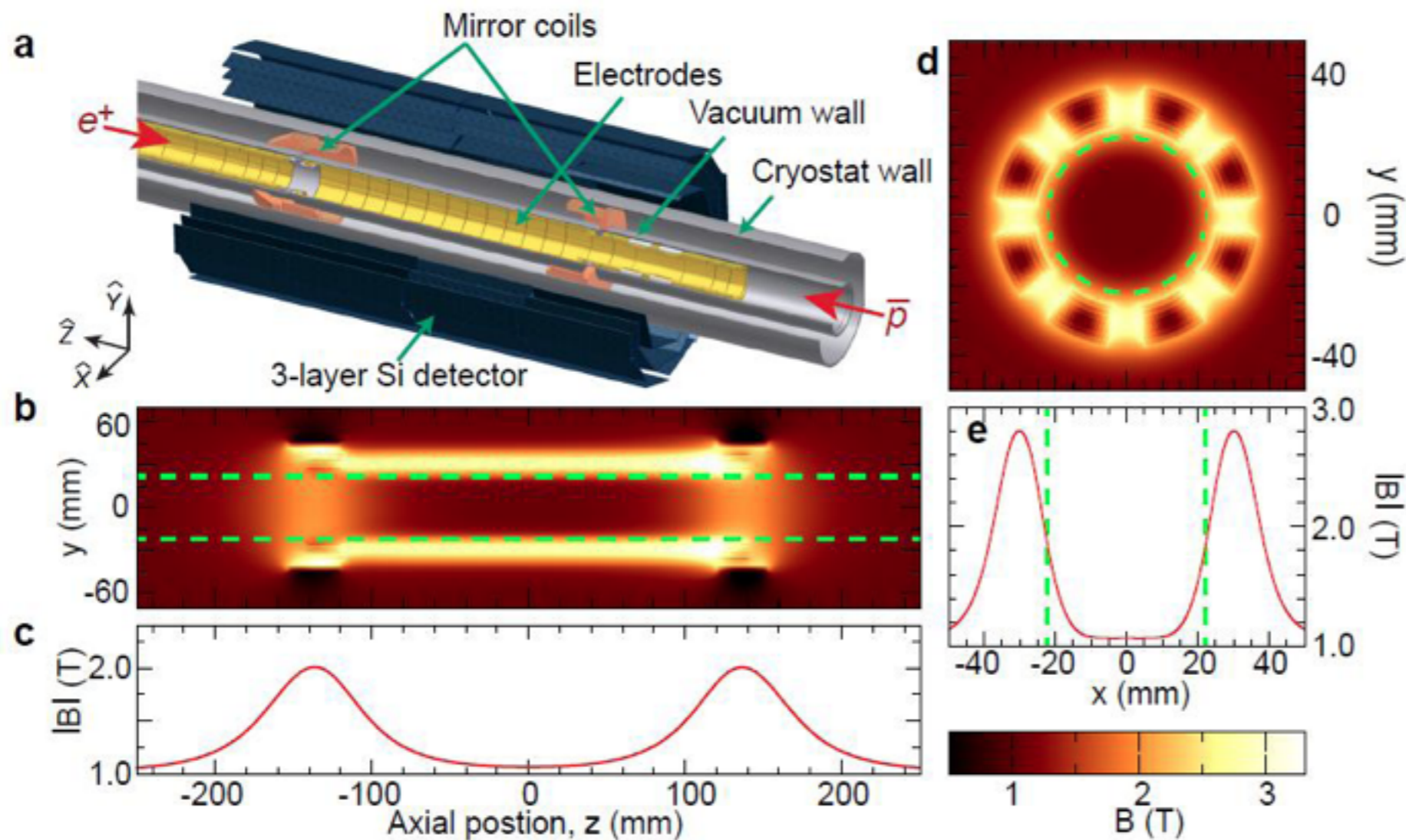


vibrational

Inherent
0.36 eV
final state
smearing



Trapping of Atomic Tritium



Similar design to anti-hydrogen trapping:

Solenoidal field for uniformity

Pinch coils for axial confinement

Ioffe multipoles for radial confinement

Cooling polarized tritium down to $\sim 1\text{K}$ is necessary (and the main challenge)

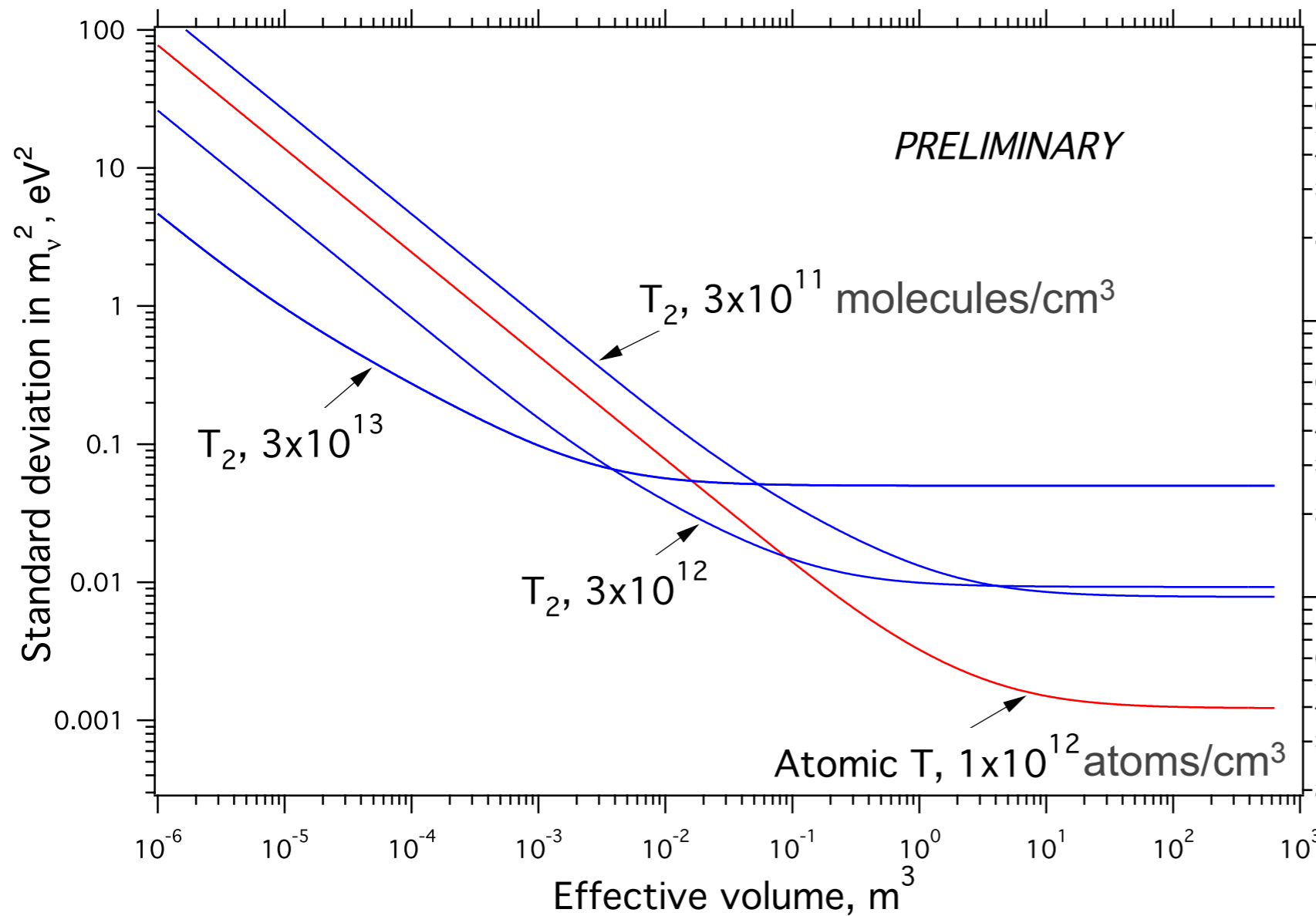
ALPHA Collaboration: Nature Phys.7:558-564,2011; arXiv 1104.4982

In order to achieve atomic tritium purity, it is necessary to cool and trap polarized atomic tritium in both a radial and axial magnetic trap (Ioffe-Pritchard traps).

Technique quite similar to hydrogen BEC (MIT) and anti-hydrogen trapping (ALPHA).

Densities low, so recombination is highly suppressed.

Projected Sensitivity (Molecular & Atomic)



Systematics include:

Degeneracy scale

Statistical uncertainties
(1 year run)

Final state interactions

Volume $\approx 0.05 \text{ m}^3$
Thermal broadening
($\approx 70 \text{ mCi}$)

Scattering

Inverted

Volume $\approx 5 \text{ m}^3$
(0.25 Ci)

Field inhomogeneity

1% uncertainty in resolution
distribution

Sensitivity for both molecular and atomic tritium are shown.

Systematics include final state interactions, thermal broadening, statistical uncertainties, and scattering.

Can calibrate against frequency standards.

Degeneracy and Beyond..

Spectroscopy (KATRIN)

Technique PROVEN. State-of-the-art.

Experiment soon to commence with 0.2 eV reach.

Integral measurement with TOF possibility.

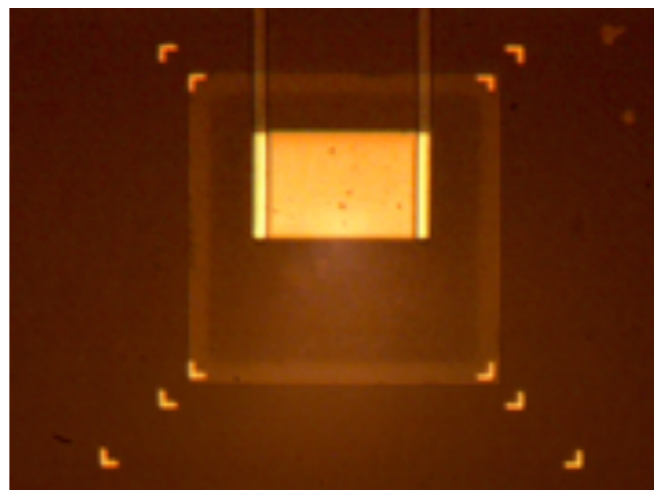


Calorimetry (HOLMES & ECHO)

Technique highly advanced.

New experiment(s) planned to reach 0.1 eV scale.

Statistics & systematics next hurdle.

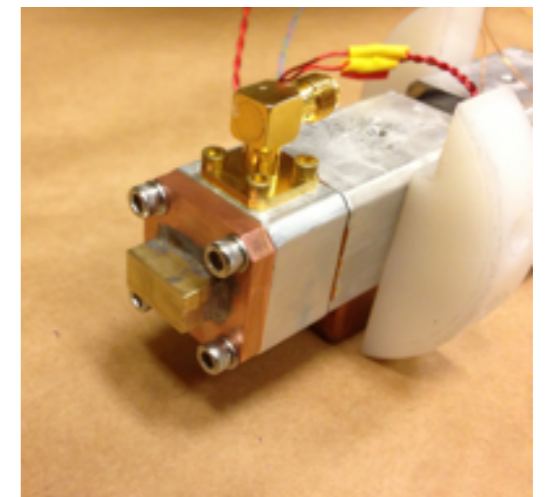


Frequency (Project 8)

Technique in R&D phase, with results soon.

Potential of scalability and exploring atomic sources to inverted scale.

Need to establish the limits of the technique.

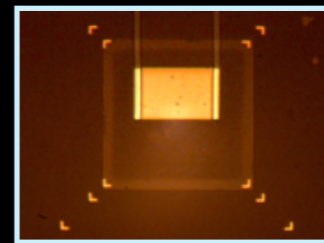


... but the road is long.

Spectroscopy
(KATRIN)



Calorimetry
(HOLMES & ECHO)



Frequency
(Project 8)!

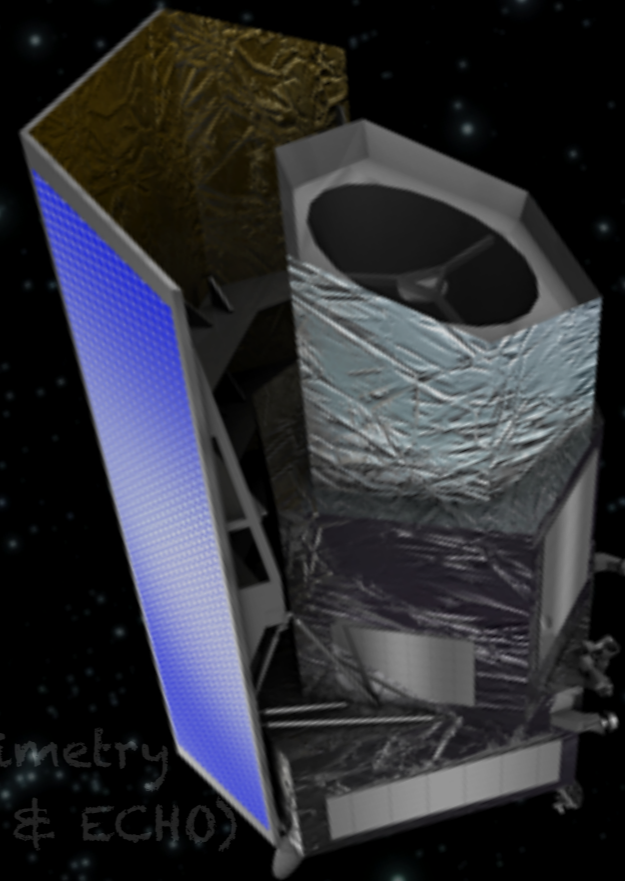


...or will cosmology just settle the problem?

Spectroscopy
(KATRIN)



Calorimetry
(HOLMES & ECHO)



Frequency
(Project 8)!



Some good
advice...

"Joe, you are wrong..."

[It matters not that cosmology may measure it first. Thermal cosmology track neutrinos over eons, from being relativistic to non-relativistic. A positive measurement of neutrino mass in cosmology AND nuclear physics would be the ultimate confirmation of that model.]

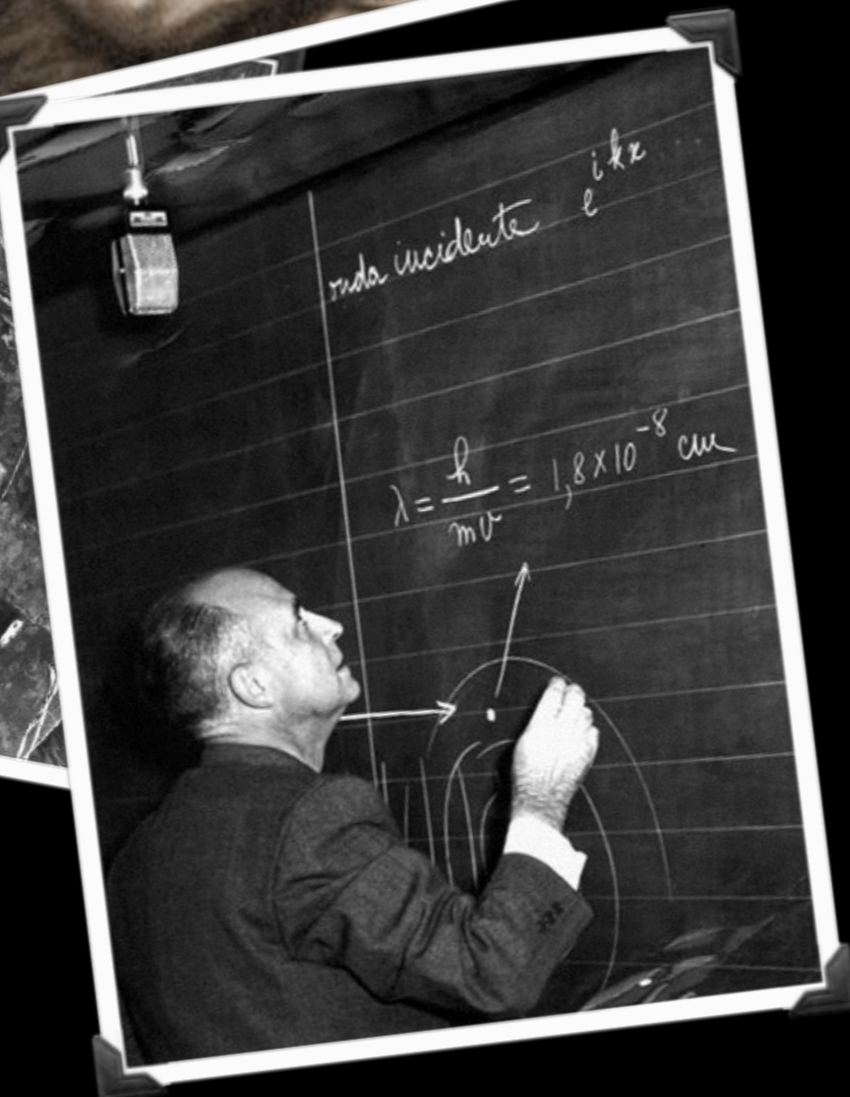
[You can't beat that kind of impact...] - G. Fuller





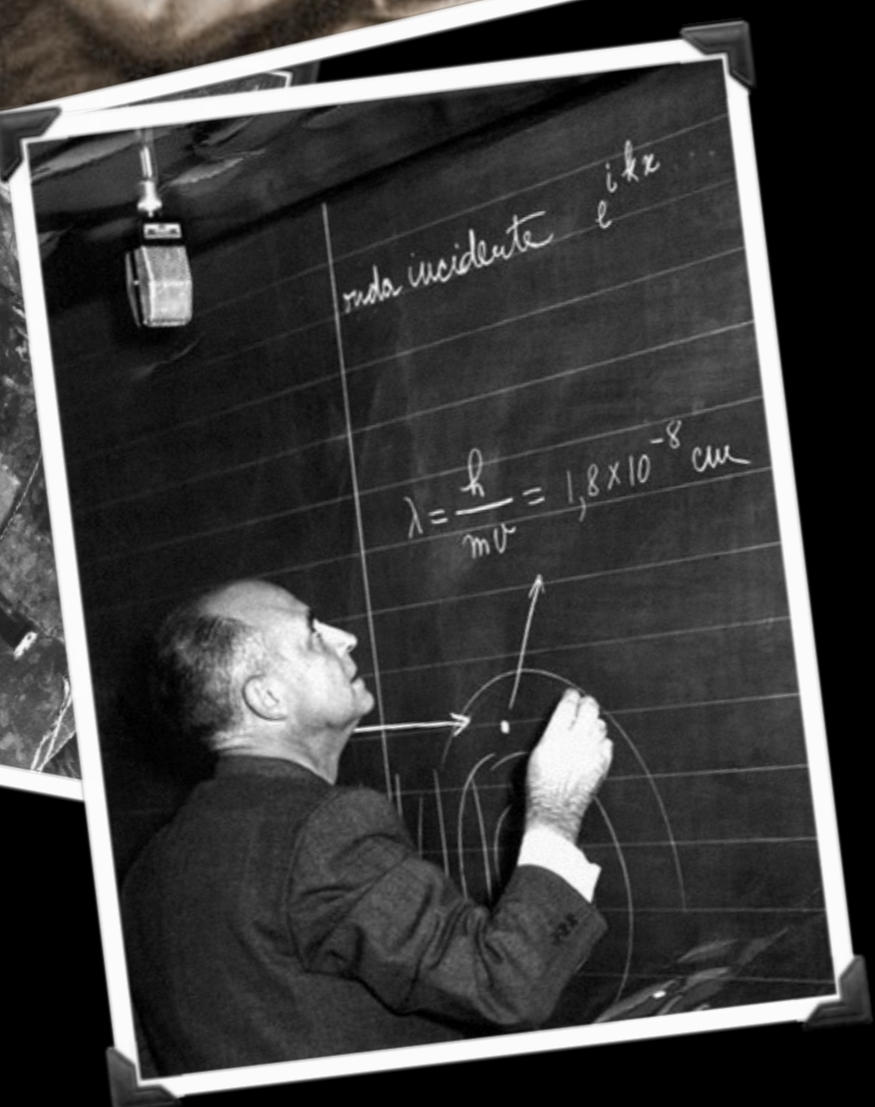
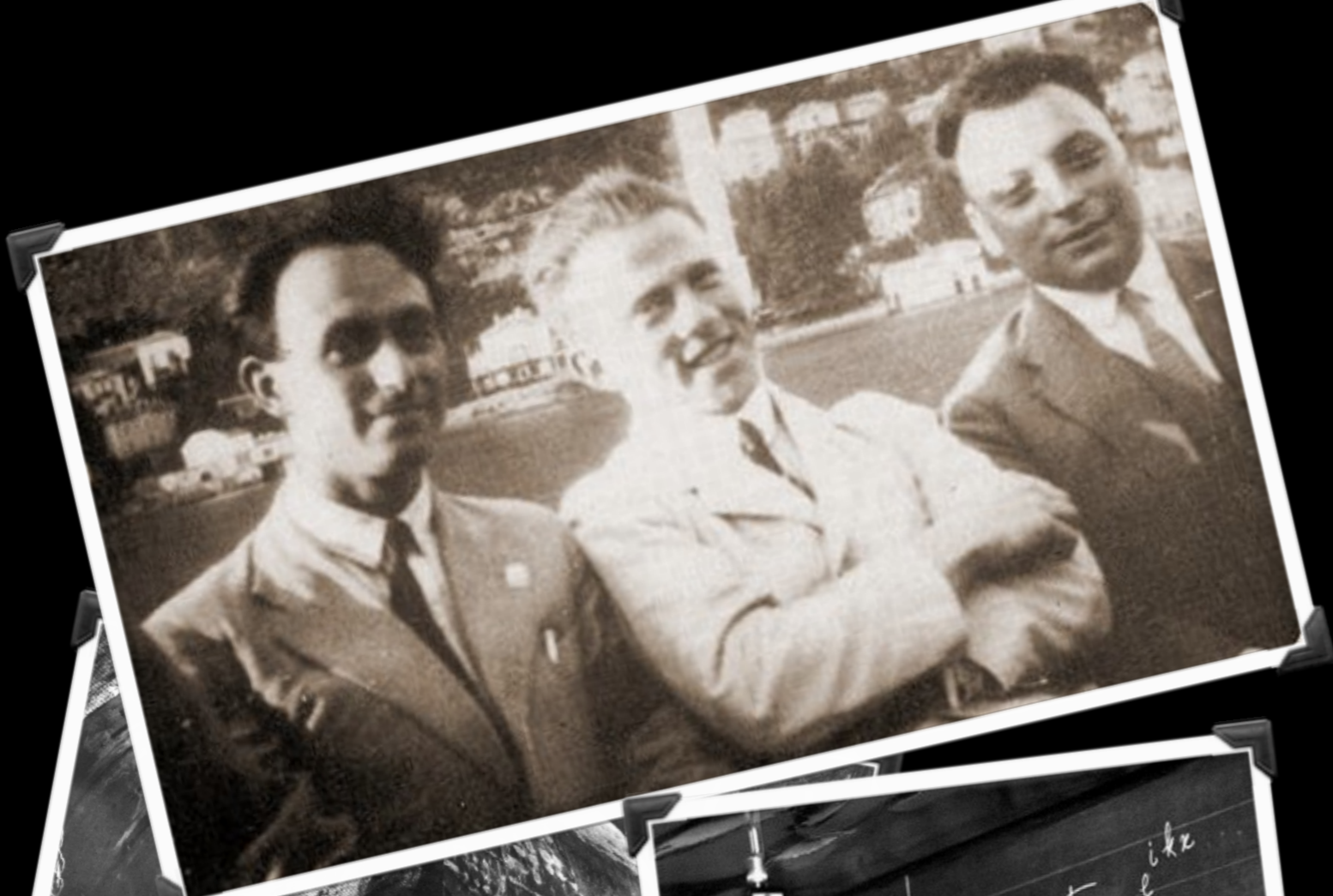
Direct probes on neutrino mass measurements may provide a robust test of cosmology (and vice-versa).

Over the next decade, KATRIN, ECHO and HOLMES can probe the degeneracy scale.



Project 8 introduces a new technique into the mix. Potentially extendable to the inverted scale.

Lots of exciting work to come!

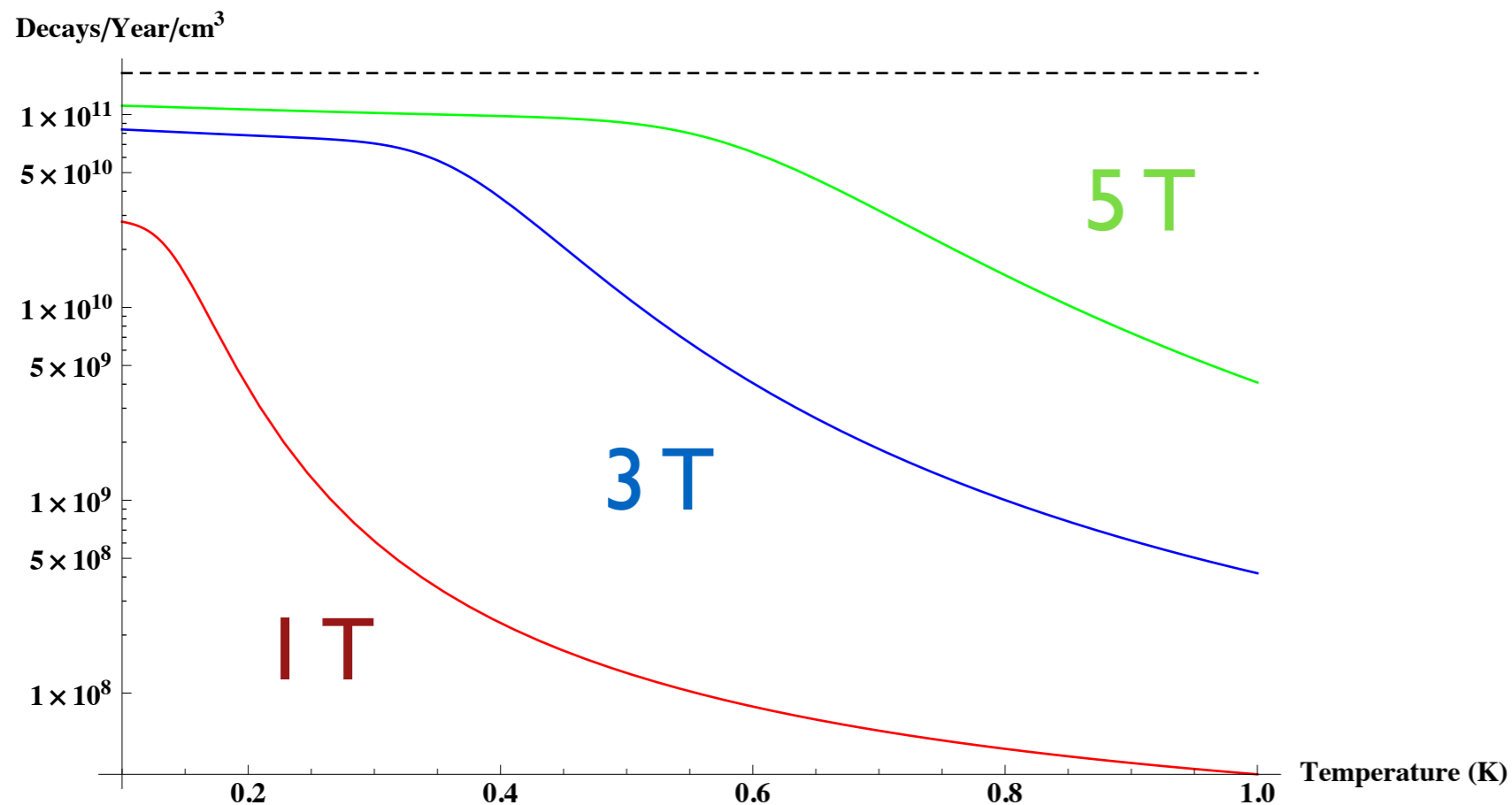


onda incidente e ikz

$$\lambda = \frac{h}{mv} = 1,8 \times 10^{-8} \text{ cm}$$

Thank you for
your attention

Trapping of Atomic Tritium



Important to cool tritium target to clear out T_2 contamination and keep source stable.

Source increases in stability with magnetic field and temperature.

Stability based on hydrogen trapping (preliminary)

In order to achieve atomic tritium purity, it is necessary to cool and trap polarized atomic tritium in both a radial and axial magnetic trap (Ioffe-Pritchard traps).

Technique quite similar to hydrogen BEC (MIT) and anti-hydrogen trapping (ALPHA).

Densities low, so recombination is highly suppressed.

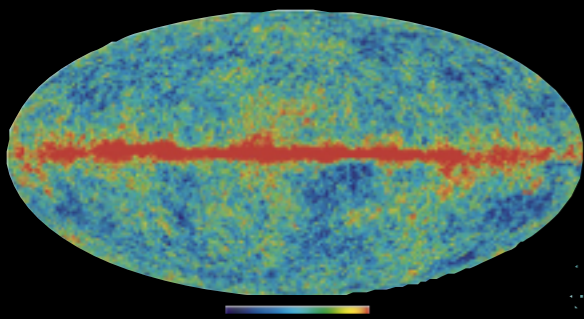
Systematics and Target Specifications

Uncertainty	Target Specification	In-situ Calibration	Ex-Situ	Monitoring
Field Homogeneity	$\pm 10^{-7}$	ESR; NMR	NMR probe	
Field Stability	$\pm 10^{-7}$	ESR; NMR		Hall probe
Target Density	$\pm 10^{-5}$	T spectrum; ESR	e-gun	P monitoring
Target Purity	$\pm 10^{-6}$	ESR; NMR	gas monitoring	
Energy Scale	$\pm 10^{-5}$	(see field calibration)	^{83m}Kr injection	
Antenna Response	Unknown			
Final States (eV)	0.36 eV (T_2 only)			
Space-charge	Unknown		e^- Kr studies	
Ion Trapping	Unknown			
Doppler (eV)	0.1 (T_2) 0.025 (T)			T monitoring
Background	$1\mu\text{Hz}$	T spectrum		
Quantum Limit	Unknown			calculation

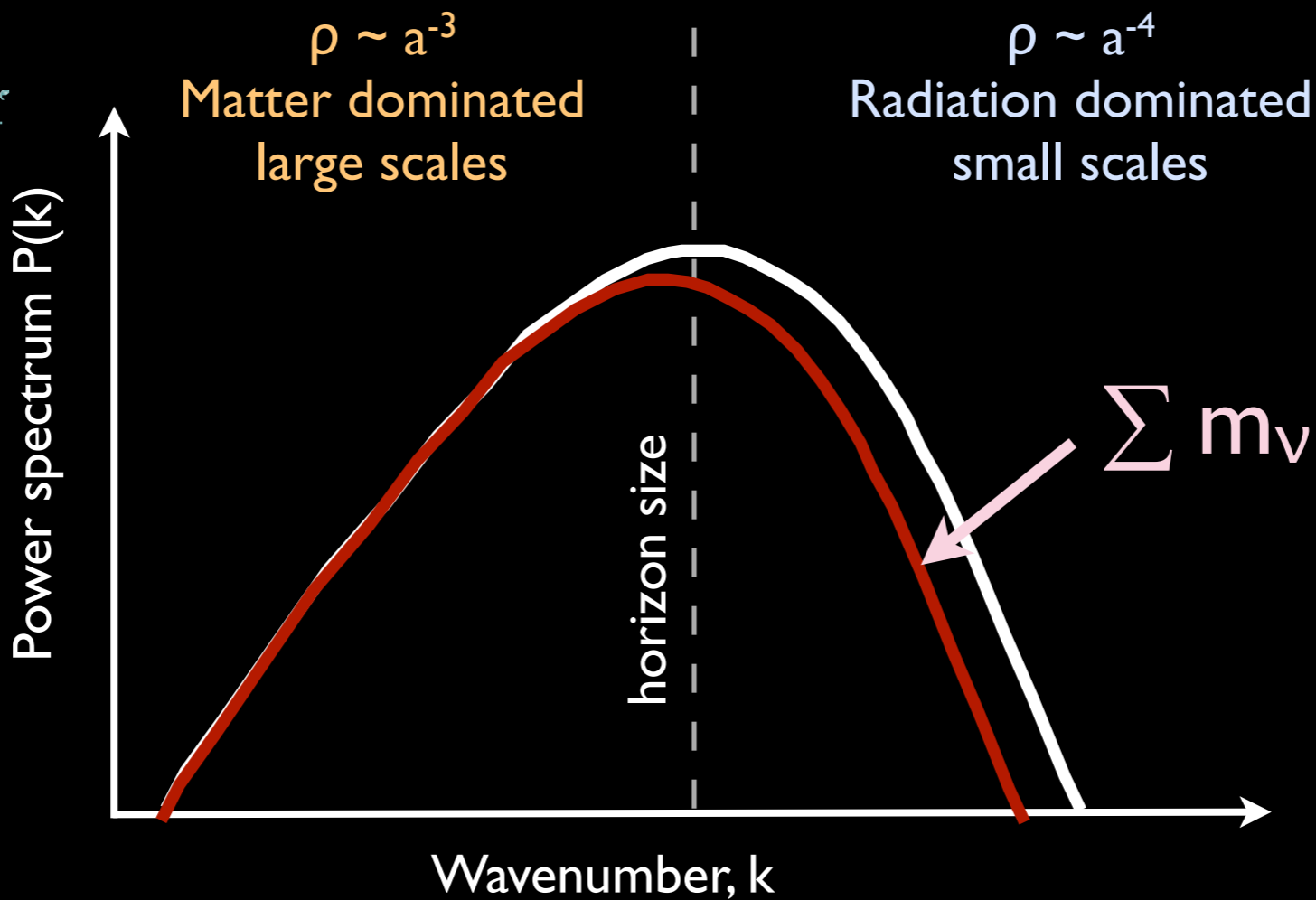
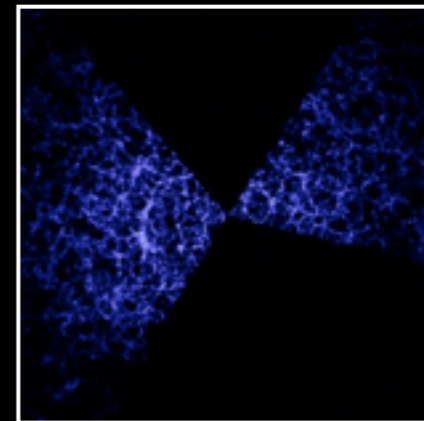
Systematic table is being built in order to guide next atomic tritium design.

The Strategy (a naive view)

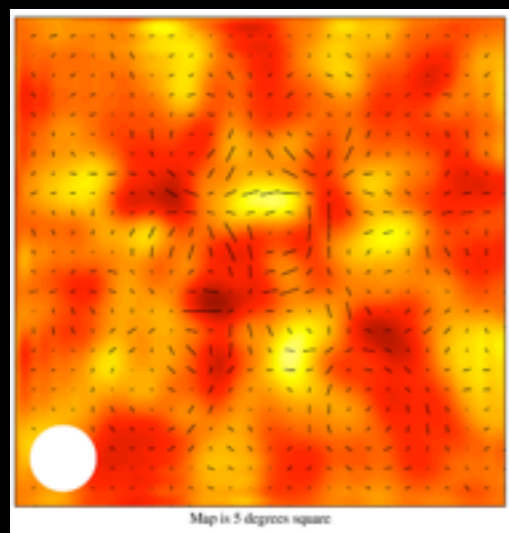
WMAP Temperature Map



Galaxy Surveys



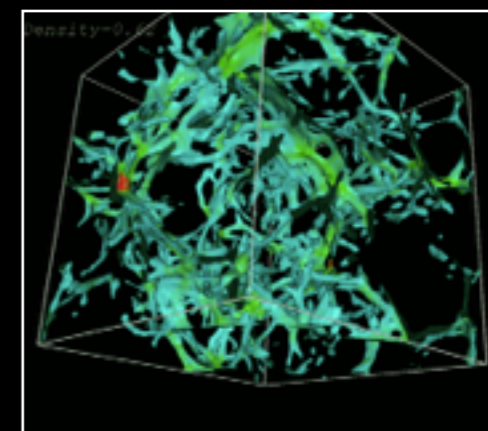
Weak lensing



CMB Polarization

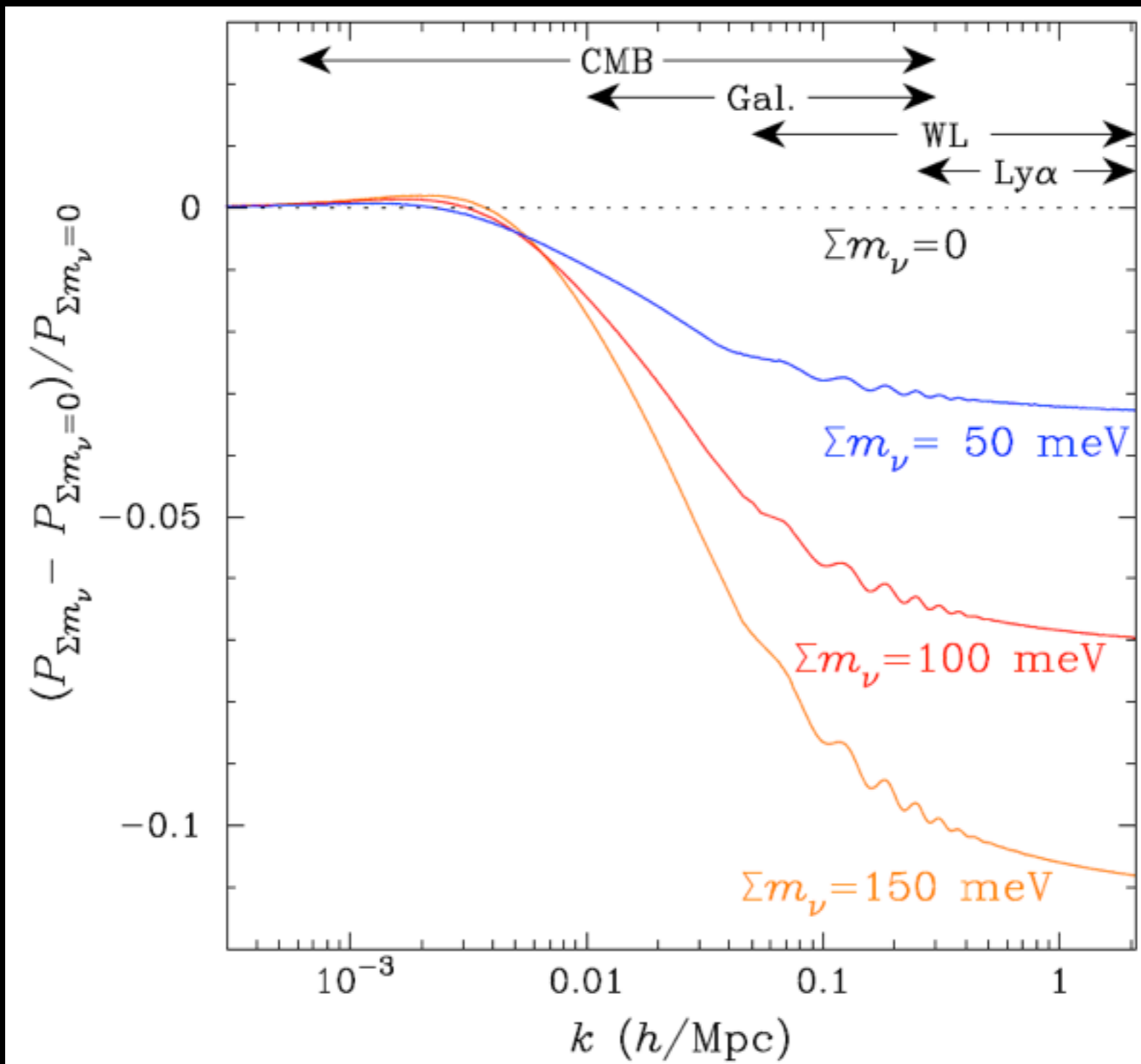
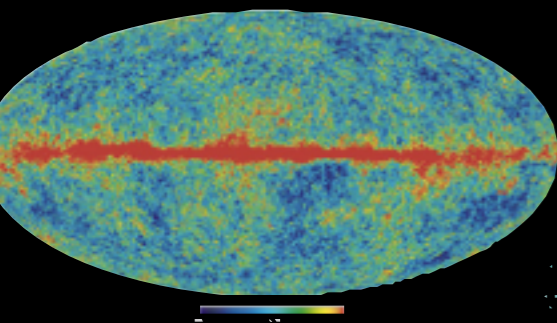
$\delta(x) = (\rho(x) - \bar{\rho}) / \bar{\rho}$
Neutrinos come to affect the power spectrum,

particularly at small distance scales
 $P(k) = \langle |\delta(k)|^2 \rangle$

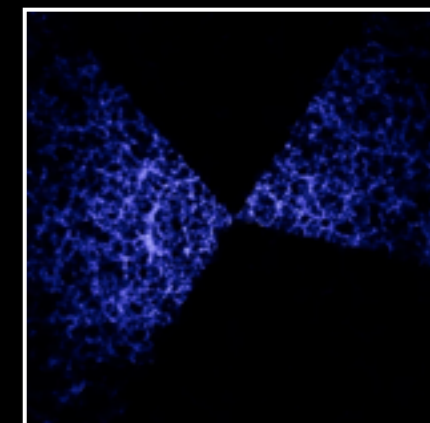


Lyman α

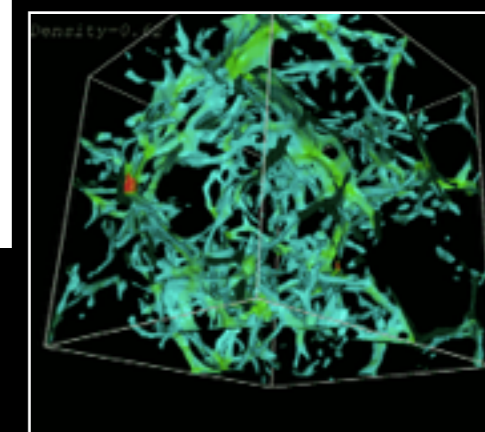
Temperature Map



Galaxy Surveys



Weak lensing



Lyman α

Large scale structure tends to weaken power spectrum at small wavelengths...

CMB Polarization

