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# Search of Deeply Bound Kaon States @ SuperB

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4<sup>th</sup> SuperB Collaboration Meeting  
La Biodola, June 1, 2012

- **Deeply Bound Kaon States**
- **Physics Motivations for DBKS**
- **Direct decays of  $\Upsilon$  into Baryons**
  - **Production of Baryons**
  - **Production of Bound States**
- **Search of DBKS @ SuperB**

# DEEPLY BOUND KAON STATES

Existence of states with a  $\bar{K}$  deeply bound to a few nucleons  
predicted by Akaishi-Yamazaki PLB535(2002)

phenomenological  $\bar{K}$ -N potential  
constructed by exp results

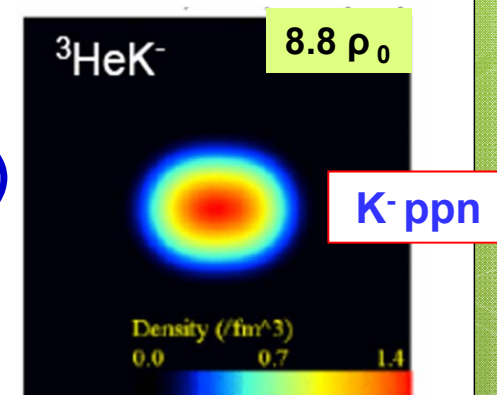


- $\bar{K}$  N scattering length
- Kaonic Hydrogen X-Ray
- Energy and width of  $\Lambda(1405)$

➤  $\bar{K}N$  strongly attractive interaction ( $I = 0$ )  
⇒ Nuclear states strongly bound ( $\sim 100$  MeV)

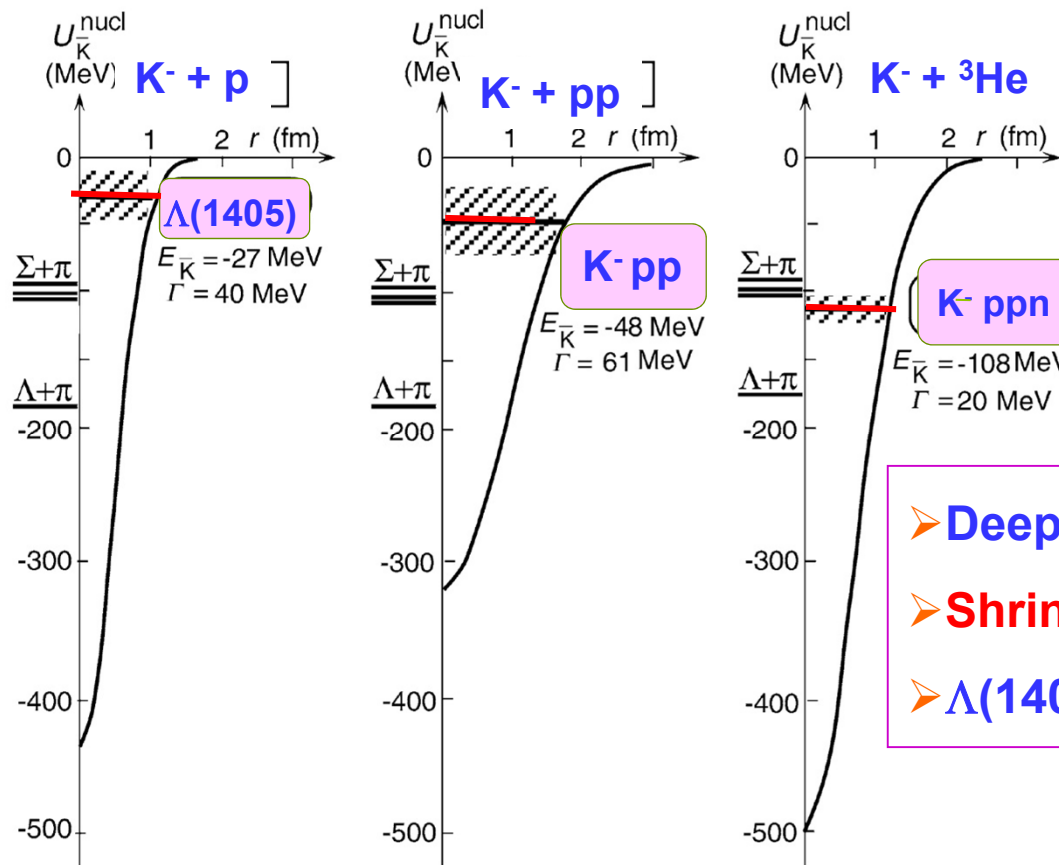
➤ Very dense systems:  $\rho > 3\rho_0$

➤ Clustering-structure → Strange Nuclear structure



# DEEPLY BOUND KAON STATES

The  $\Lambda(1405)$  is the simplest  $\bar{K}N$  bound state in the model



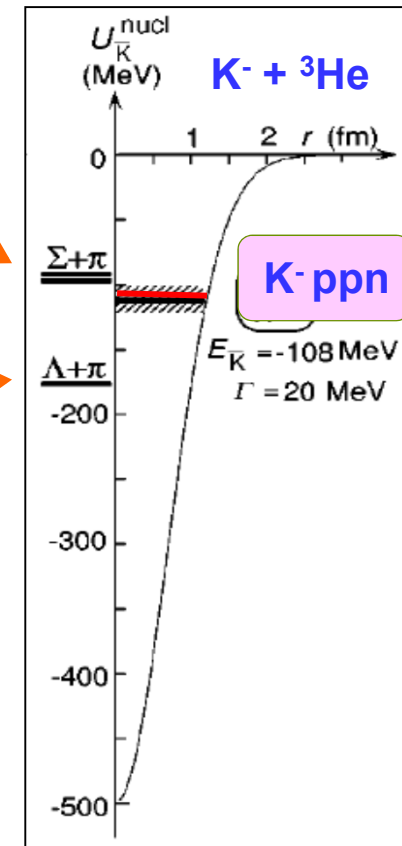
- Deeply Bound ( $> 100 \text{ MeV}$ )
- Shrinkage of the core nucleus
- $\Lambda(1405)$  as a  $\bar{K}N$  bound system

Formation of  $\Lambda(1405)$  in a nucleus as a **seed** leads to  $\bar{K}$  bound states

# DEEPLY BOUND KAON STATES

Narrow widths of **20-40 MeV** because:

- below  $\Sigma\pi$  emission threshold:  ~~$\bar{K}N \rightarrow \Sigma\pi$~~
- above  $\Lambda\pi$  emission threshold  
**BUT  $\bar{K}N \rightarrow \Lambda\pi$  suppressed**  
 due to  **$I = 0$  attraction**
- only  **$Y+nN$**  decay channels are allowed

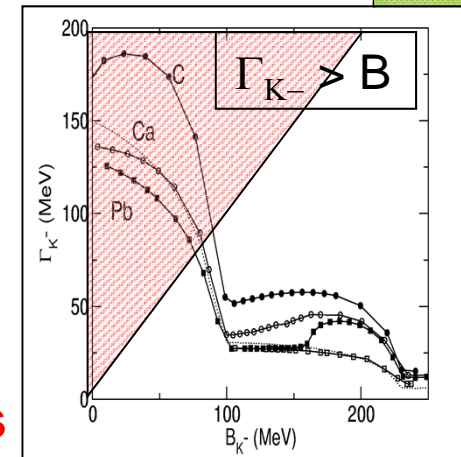


# DEEPLY BOUND KAON STATES

## Different Theoretical Approaches

Deep potential ? (150-200 MeV)  
Shallow potential ? (50-75 MeV)

- Gal, Weise, Schaffner-Bielich, Wychech
  - Possible existence of  $\bar{K}$ -nuclear aggregates,  
but **potential is shallow** and expected **widths are large**
  - Observation is difficult
- Oset, Ramos, Toki  
observed signals due to **Final State Interactions**
- Mares et al.  
high binding energy **only for heavy nuclear targets**



Their existence is not denied

But the possibility to observe them as narrow states is unlikely

Theoretical approaches: **several attitudes but limited predictive power**

# DEEPLY BOUND KAON STATES

**FINUDA @DAΦNE**

**first observation**

**of the lightest DBKS in  $p\Lambda$  spectrum**

**$K^- pp$**

**$K^-$  stopped in  
Light targets  ${}^6\text{Li}$ ,  ${}^7\text{Li}$ ,  ${}^{12}\text{C}$**

$$M = (2255 \pm 9) \text{ MeV}/c^2$$

$$B = 115^{+6}_{-5} \text{ (stat)} + {}^{+3}_{-4} \text{ (sys)} \text{ MeV}$$

$$\Gamma = 67^{+14}_{-11} \text{ (stat)} + {}^{+2}_{-3} \text{ (sys)} \text{ MeV}$$

**Yield  $\approx 0.1\%$ /stopped  $K^-$**

**Akaishi-Yamazaki, PLB535(2002)**

**$B=48\text{MeV}$ ,  $\Gamma=61\text{MeV}$**

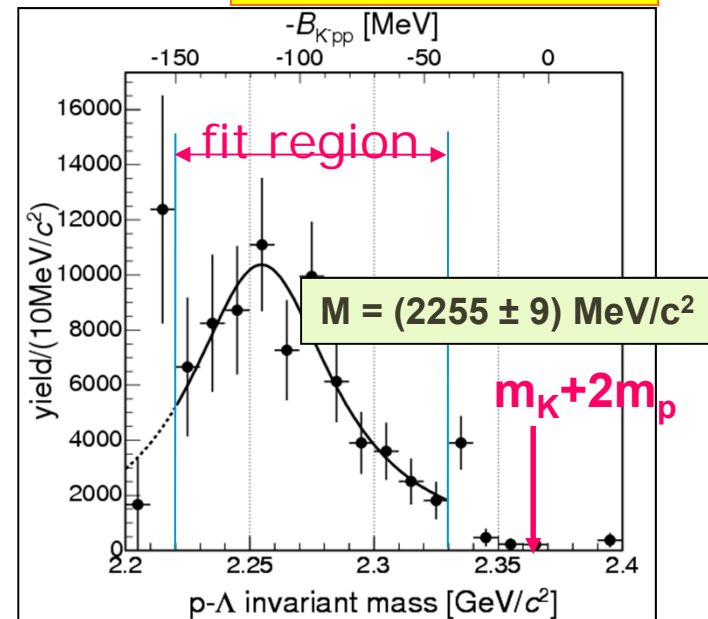
**Shevchenko, PRL98(2007)**

**$B=50-70\text{MeV}$ ,  $\Gamma\sim 100\text{MeV}$**

**Ivanov, nucl-th/0512037**

**$B=118 \text{ MeV}$ ,  $\Gamma\sim 58\text{MeV}$**

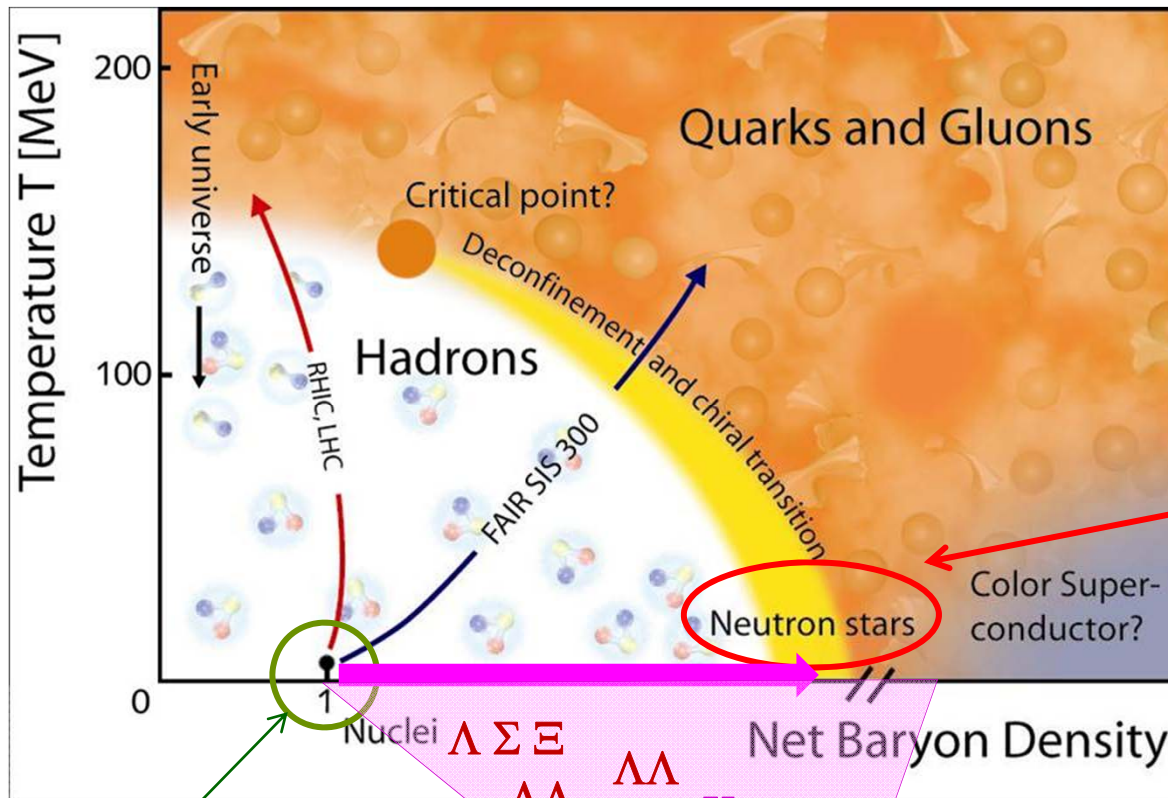
**PRL 94(2005)212303**



**$K^- (pp) \rightarrow X \rightarrow \Lambda p$   
 $\rightarrow \Sigma^0 p \rightarrow \Lambda \gamma p$**

**$p \Lambda$  events strongly  
back-to-back correlated**

## Phase Diagram of Nuclear Matter



Nuclear density  
5-10  $\times \rho_0$

Our WORLD  
Nuclear density  $\rho_0 \sim 0.17 \text{ fm}^{-3}$   
kT = 0

Strangeness Degrees of Freedom

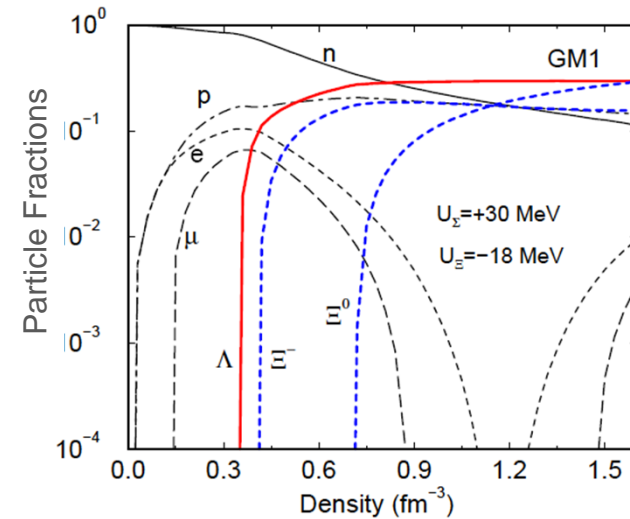
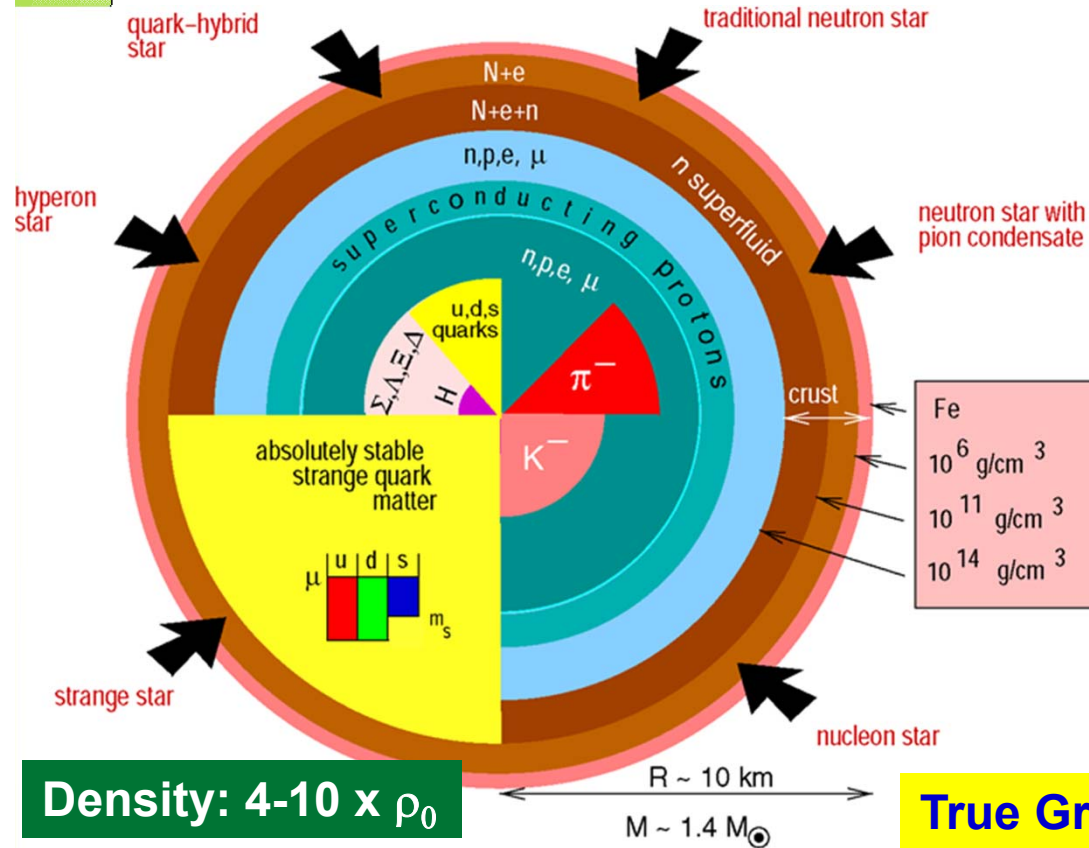


# PHYSICS

## MOTIVATIONS

### Neutron Stars or Strange Stars ?

Strangeness/Baryon ratio  $\sim 1$   
 Strangeness expected both confined (hyperons and kaons) and deconfined (strange quark matter)



J. Schaffner-Bielich, astro-ph/0703113

Density:  $4-10 \times \rho_0$

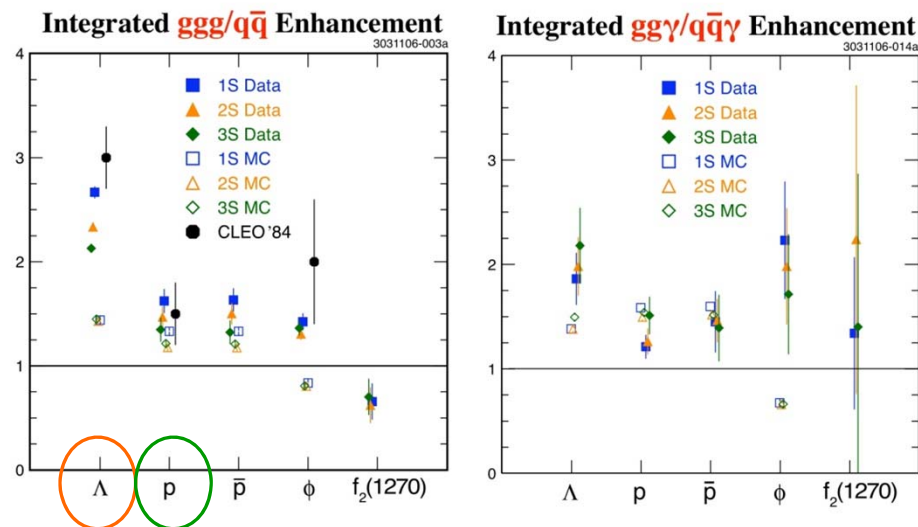
True Ground State of the Strong Interaction

Strongly attractive  $K^-$ -nuclear potential of the order of 150-200 MeV could affect  $K^-$  condensation and evolution of strangeness in high-density nuclear matter in the core of Neutron Stars, where  $Y$  could be mixed with  $N$  in a stable system and could play crucial roles

# FROM THE MEDIUM TO THE VACUUM

## Why to search for DBKS in $\Upsilon$ Decays ?

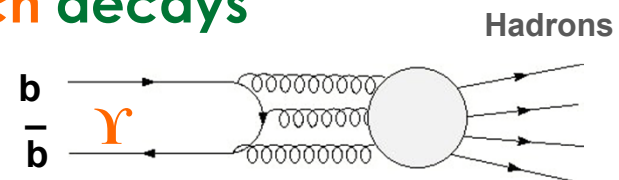
- **Baryon production** in  $\Upsilon(nS)$  decays is known to be **enhanced** of about a factor of 2 relative to the continuum hadronization



CLEO, Mitchell QCD2007

**Baryon enhancements decrease in  $gg\gamma/q\bar{q}\gamma$**

- $\Upsilon(1S)$  decays  $\sim 80\%$  via  $ggg \Rightarrow$  **glue-rich decays**  
 $\Rightarrow$  **exotic multiquark states**



# BARYON-ANTIBARYON PRODUCTION

$$e^+ e^- \rightarrow q \bar{q}$$

Different experiments observe an Enhancement of  
Correlated production

- $p-\bar{p}$
- $\Lambda-\bar{\Lambda}$
- $\Lambda_c-\bar{\Lambda}_c$

In contrast to inclusive single-particle production, compensation of Baryon and charm number is a more subtle aspect of quark fragmentation modelling

- **Local** compensation (small  $|\Delta y|$ , small opening angle)
  - $e^+e^-$  annihilation in  $qq\bar{q}$
- **Global/Primary** compensation (large  $|\Delta y|$ , large opening angle)
  - $e^+e^-$  annihilation in **diquark-antidiquark**, which hadronize in jets containing a leading Baryon and a leading antiBaryon, respectively

# BARYON-ANTI-BARYON PRODUCTION



Charmed baryons **3x** more likely to appear opposite an anti-charmed baryon than an anti-charmed meson

$e^+e^- \rightarrow c\bar{c} \rightarrow \Lambda_c \bar{\Lambda}_c$  X events @  $\Upsilon(4s)$  and below

CLEO

PRD63 (2001) 112003

**(3.35 ± 0.45 ± 0.42) times**

BaBar

PRD82 (2010) 091102(R)

**~ 4 times**

TABLE III. Production rates; statistical errors only.

Double tags Single tags	Data fraction
$2 \times \frac{\Lambda_c^+ \bar{\Lambda}_c^-}{\Lambda_c^+ D^0}$	<b>(7.19 ± 1.08) × 10<sup>-3</sup></b>
$\frac{\Lambda_c^+ D^0}{D^0}$	(2.05 ± 0.16) × 10 <sup>-3</sup>
$\frac{\Lambda_c^+ D^-}{D^-}$	(2.03 ± 0.26) × 10 <sup>-3</sup>
$\frac{\Lambda_c^+ \bar{\Lambda}_c^-}{\bar{\Lambda}_c^-}$	(19.3 ± 1.1) × 10 <sup>-3</sup>

Particles in two **opposite** emispheres



**Correlated**  
production of  $\Lambda_c^+ \bar{\Lambda}_c^-$

**Baryon # conserved by a leading Baryon and AntiBaryon, rather than locally along the hadronization chain**



**Primary diquark-antidiquark production and/or production of multiple intermediate mesons between a Baryon and antiBaryon**

# BARYON-ANTIBARYON PRODUCTION

$\Lambda \bar{\Lambda}$

**3  $\sigma$  excess has been found in the estimation of the primary correlated  $\Lambda \bar{\Lambda}$  yield**

**CLEO**

PRD636(2002) 052002

$e^+e^- \rightarrow q\bar{q}$  events @  $\Upsilon(4s)$  and below

**872  $\pm$  288 excess events**

**$\sim 3\sigma$  excess of particles produced in two **opposite** emispheres**



**Correlated production of  $\Lambda \bar{\Lambda}$**

**Here the result is **model dependent** since **JETSET** event generator has been heavily used to estimate non primary contribution**

**NO significant statistical signal for primary ( $\Lambda \bar{\Lambda}$ ) production was found in previous experiments at **LEP****

# ANTI-DEUTERON PRODUCTION

$d \bar{d}$

CLEO

- $B^*(\Upsilon(1S) \rightarrow \bar{d} + X) = (3.36 \pm 0.23 \pm 0.25) \times 10^{-5}$  direct decays to  $ggg$  and  $\gamma gg$
- $B(\Upsilon(4S) \rightarrow \bar{d} + X) < 1.3 \times 10^{-5}$
- $B(\gamma^* \rightarrow q\bar{q} \rightarrow \bar{d} + X) < 1.3 \times 10^{-5}$  determined off resonance

➤  $ggg + \gamma gg$  is at least about **3x more likely** than  $\gamma^* \rightarrow q\bar{q}$  to produce deuterons

➤ How often is an anti-d **compensated** by a d as compared to (n, p) combinations ?

Roughly equal compensation by **nn**, **pp**, **np** relative to each other

**3  $\bar{d}d$  events**

**$R(e^+e^- \rightarrow \bar{d}dX) / R(e^+e^- \rightarrow \bar{d}NNX) \sim 1\%$**

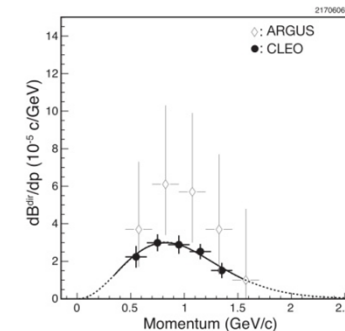
# ANTI-DEUTERON PRODUCTION

## Theoretical descriptions of antideuteron formation based on **Coalescence Model**

PLB98(1981)153 - PRep131(1986)223



$\bar{p}$  and  $\bar{n}$  nearby to each other  
in phase space bind together



Fit to Maxwell distribution as used in fireball model

**Coalescence disfavoured** if anti-deuterons always compensated by deuterons, instead of n and p

Yield =  $(5.9 \pm 1.8 \pm 0.5) \times 10^{-6}$   $Z^0$  into  $\bar{d}$  @ Aleph

Yield <  $(0.8 \pm 0.5) \times 10^{-6}$   $Z^0$  into  $\bar{d}$  @ OPAL

$\bar{d} / \bar{p} \rightarrow (5.0 \pm 1.0 \pm 0.5) \times 10^{-4}$ ,  $\sigma = (2.7 \pm 0.5 \pm 0.2)$  nb , in  $\gamma p$  @ H1

High yields of anti-deuterons or anti-nuclei indirect signature of New Physics, such as **Dark Matter**

## AntiNuclei & AntiHypernuclei

TABLE I: Particle ratios from Au+Au collisions at  $\sqrt{s_{NN}} = 200$  GeV/c. The  ${}^3\text{He}$  ( ${}^3\bar{\text{He}}$ ) yield have been corrected for  ${}^3_{\Lambda}\text{H}$  ( ${}^3_{\Lambda}\bar{\text{H}}$ ) feed-down contribution.

Particle type	Ratio
${}^3_{\Lambda}\bar{\text{H}}/{}^3_{\Lambda}\text{H}$	$0.49 \pm 0.18$ (stat.) $\pm 0.07$ (sys.)
${}^3\bar{\text{He}}/{}^3\text{He}$	$0.45 \pm 0.02$ (stat.) $\pm 0.04$ (sys.)
${}^3_{\Lambda}\bar{\text{H}}/{}^3\bar{\text{He}}$	$0.89 \pm 0.28$ (stat.) $\pm 0.13$ (sys.)
${}^3_{\Lambda}\text{H}/{}^3\text{He}$	$0.82 \pm 0.16$ (stat.) $\pm 0.12$ (sys.)

Coalescence



$${}^3_{\Lambda}\bar{\text{H}}/{}^3_{\Lambda}\text{H} \propto (\bar{p}/p)(\bar{n}/n)(\bar{\Lambda}/\Lambda)$$

$${}^3\bar{\text{He}}/{}^3\text{He} \propto (\bar{p}/p)^2 (\bar{n}/n)$$

$0.45 \sim 0.77 \cdot 0.77 \cdot 0.77$

**Ratios favour Coalescence**

AntiHypernuclei produced with similar yields of AntiNuclei,  
unlike what happens at lower energies





# SEARCH OF DBKS

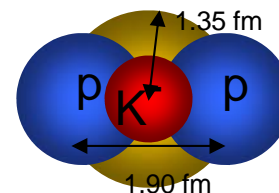
@ SUPERB

$$e^+e^- \rightarrow \Upsilon(1S) \rightarrow (K^- pp) + X$$

$$K^- pp \rightarrow \Lambda p$$

$$\rightarrow \Sigma^0 p \rightarrow \Lambda \gamma p$$

$\Upsilon(1S)$  Strong Decays  $\sim 80\%$   
through  $ggg$  hadronization



$I = 0$  Strong attractive potential  
Very compact system

➤  $\Lambda p$  Invariant Mass spectrum

❖ Search for Mass  $m \sim 2.25 \text{ GeV}/c^2$  and narrow width  $\Gamma < 100 \text{ MeV}$

➤ Check  $\Lambda p$  Angular correlations

We rely in the high performance of the SuperB detector

- At least as frequent as Hypernucleus formation ( ${}^3_{\Lambda}\text{H}$ ,  ${}^4_{\Lambda}\text{H}$ ,  ${}^4_{\Lambda}\text{He}$ )
- No medium  $\rightarrow$  different FSI  $\rightarrow$  Identification easier than in HI exp

## SUMMARY

- Kaonic Nuclear States have been predicted to be **deeply bound** due to the  $\bar{K}N$  strong interaction in the  **$l=0$  term**
- **$K^-pp$** , the lightest DBKS, has been first observed by **FINUDA** at DAFNE
- Study of DBKS is important to understand **high density nuclear matter** inside the **neutron stars** (kaon condensation)
- Nuclear bound systems, such as **antideuterons**, have been already **observed in  $\Upsilon$  decays** at B-factories
- Search of DBKS in **direct decays of  $\Upsilon(1S)$**  at SuperB can be pursued exploiting the high performance of the apparatus
- Measurements of the production yields could give important clues about both **DBKS** and **hadronization processes** in  **$\Upsilon(1S)$  decays**

# BACKUP SLIDES

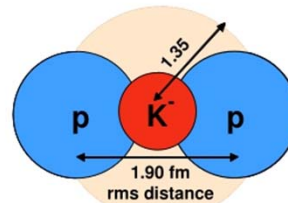
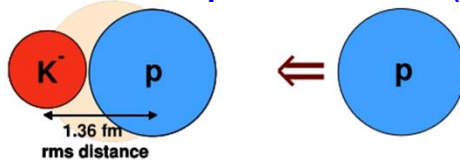


# DEEPLY BOUND KAON STATES

**K<sup>-</sup> pp**

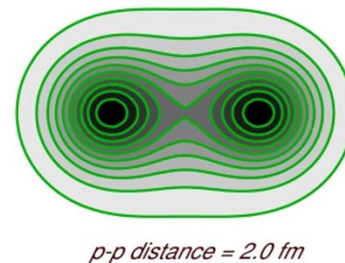
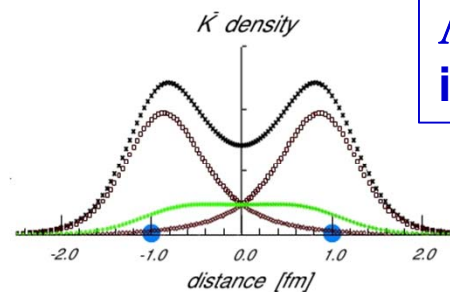
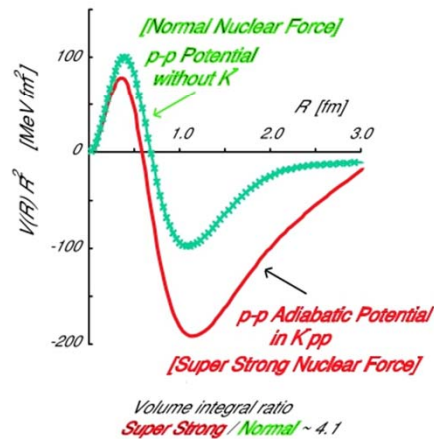
**K<sup>-</sup> combines two protons into a strongly bound system**

proton approaches  
a bound K-p “atom”,  $\Lambda(1405)$



**Kaonic hydrogen  
molecular ion**

$\Lambda(1405)$  proton **doorway**  
in the formation process



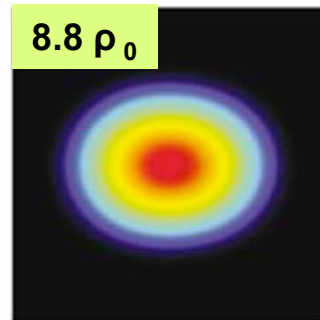
Yamazaki-Akaishi  
nucl-th0706.3651v2(2007)

Despite the drastic dynamical change caused by the strong  $\bar{K}N$  interaction, the **identity of the constituent «atom»,  $\Lambda(1405)$ , is preserved** due to the presence of a **short-range repulsion** between the two protons

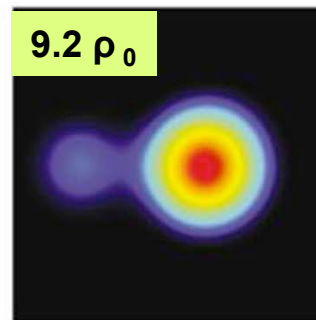
# DEEPLY BOUND KAON STATES

AMD Calculations by Dote et al.  
(Antisymmetrized Molecular Dynamics)

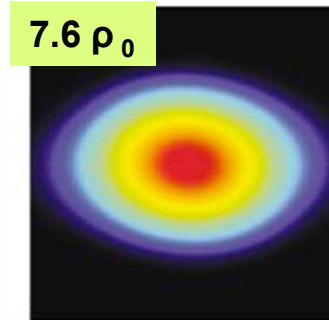
NPA738 (2004) 372



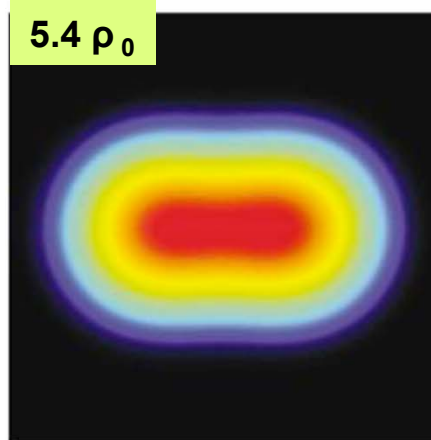
ppnK<sup>-</sup>



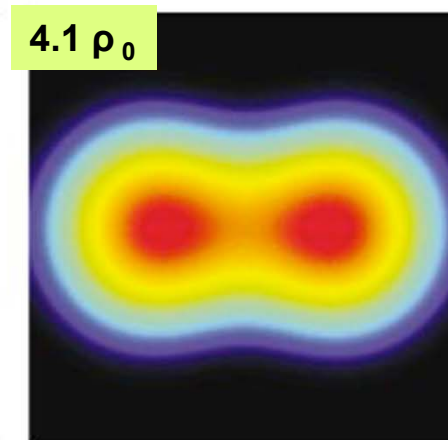
pppK<sup>-</sup>



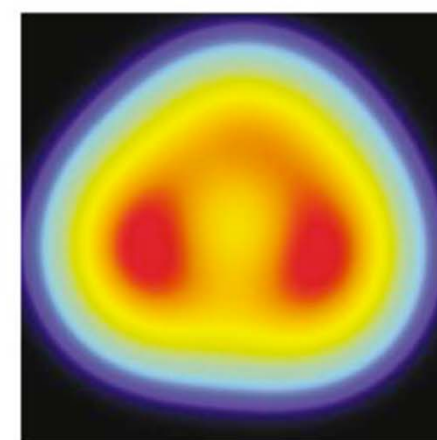
pppnK<sup>-</sup>



<sup>6</sup>BeK<sup>-</sup>



<sup>9</sup>BK<sup>-</sup>



<sup>11</sup>CK<sup>-</sup>

Dense Nuclear States