



NUMERICAL SOLUTION OF
QUANTUM PROBLEMS
VIA BOHMIAN TRAJECTORIES

Outline of the talk

- Time Dependent Schroedinger Equation

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- Numerical Methods via Bohmian Trajectories
 - Single Particle
 - Quantum Trajectory Method
 - Single Trajectory
 - Many Particle

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 - Single Particle
 - Quantum Trajectory Method
 - Single Trajectory
 - Many Particle
- Measurement issue
 - Theoretical Model
 - Numerical Simulation

Time Dependent Schroedinger Equation

$$i\hbar \frac{\partial \psi}{\partial t} = \mathbf{H}\psi$$

Different numerical approaches for different problems

Single Particle

$$i\hbar \frac{\partial \psi}{\partial t} = \left(-\frac{\hbar^2}{2m} \nabla^2 + V\right)\psi$$

$$\psi = \psi(x, t)$$

Reference: C. L. Loprore and R. E. Wyatt: *Quantum Wavepacket Dynamics with Trajectories*, Phys. Rev. Lett. 82, 5190-5193 (1999).

Many Particle

$$i\hbar \frac{\partial \psi}{\partial t} = \left(-\sum_{i=1}^N \frac{\hbar^2}{2m_i} \nabla_i^2 + V\right)\psi$$

$$\psi = \psi(x_1, x_2, \dots, x_N, t)$$

Reference: X. Oriols: *Quantum-trajectory approach to time-dependent transport in mesoscopic systems with electron-electron interactions*, Phys. Rev. Lett. 98, 066803 (2007).

We can use the trajectories X_t of Bohmian Mechanics

$$\dot{X}_t = v \quad \text{where} \quad v = \frac{\hbar}{m} \operatorname{Im} \frac{\nabla \psi}{\psi}$$

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$$\text{Polar form} \quad \psi = R e^{i\frac{S}{\hbar}} \quad \rightarrow \quad v = \frac{\nabla S}{m}$$

Schroedinger equation

$$\frac{\partial S}{\partial t} = -\frac{\nabla^2 S}{2m} - V + \frac{\hbar^2}{2m} \frac{\nabla^2 R}{R}$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \left(\rho \frac{\nabla S}{m} \right) = 0$$

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Equation of motion

$$m \frac{d^2 X_t}{dt^2} = F_{class} + F_{quant}$$

$$F_{class} = -\nabla V$$

$$F_{quant} = -\nabla \cdot \left(-\frac{\hbar^2}{2m} \frac{\nabla^2 R}{R} \right)$$

Quantum Trajectory Method

Solve the hydrodynamic equations in the lagrangian frame for a time interval δt .



Reconstruct the wave function in the update positions of the particles by means of an interpolation method.
(eg. MWLS, RBF, ...)

Quantum Trajectory Method

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Reconstruct the wave function in the update positions of the particles by means of an interpolation method. (eg. MWLS, RBF, ...)

Single Trajectory

We solve the Schrodinger equation along a single trajectory.



Sampling we can obtain the entire wave function.

Reference : J. O. Taylor: *Connection with Bohmian Mechanics*, Ph.D. Thesis (2003).

Quantum Trajectory Method

The method uses the Lagrangian formulation of the Hamilton-Jacobi and Continuity Equation

$$m \frac{d\vec{v}(\vec{x}, t)}{dt} = -\vec{\nabla}[V(\vec{x}, t) + U(\vec{x}, t)]$$

$$\frac{d\rho(\vec{x}, t)}{dt} = -\rho(\vec{x}, t) \frac{\vec{\nabla} \cdot \vec{v}(\vec{x}, t)}{m}$$

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The Propagation at discret time is

$$\rho(\vec{x}_i(t_{n+1}), t_{n+1}) = \rho(\vec{x}_i(t_n), t_n) e^{-\vec{\nabla} \cdot \vec{v}(\vec{x}_i(t_n), t_n) \delta t}$$

$$\vec{v}(\vec{x}_i(t_{n+1}), t_{n+1}) = \vec{v}(\vec{x}_i(t_n), t_n) - \frac{\delta t}{m} \vec{\nabla}[V(\vec{x}_i(t_n), t_n) + U(\vec{x}_i(t_n), t_n)]$$

$$\vec{x}_i(t_{n+1}) = \vec{x}_i(t_n) + \vec{v}(\vec{x}_i(t_n), t_n) \delta t.$$

The TDSE is solved **along** the trajectories and **step by step**.

The idea of the algorithm

$$Q_0 \rightarrow \psi_0 \rightarrow Q_1 \rightarrow \psi_1 \rightarrow \dots \rightarrow Q_n \rightarrow \psi_n$$

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Single trajectory: x_0 on $|\psi_0|^2$

Known R_{n-1} and S_{n-1} , we calculate

- $\dot{Q}_n = \frac{\nabla S_{n-1}}{m}$
- $R_n = R_{n-1}(x_0, 0) |J_{n-1}|^{\frac{1}{2}}$
- $S_n = S_{n-1}(x_0, 0) + \int_0^{t_1} m \dot{Q}_{n-1}^2 - V_{n-1} - U_{n-1}$

Single Trajectory

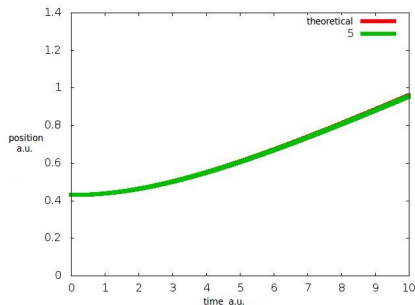
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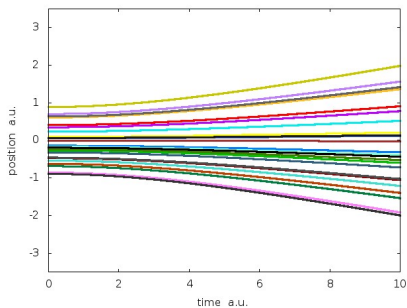
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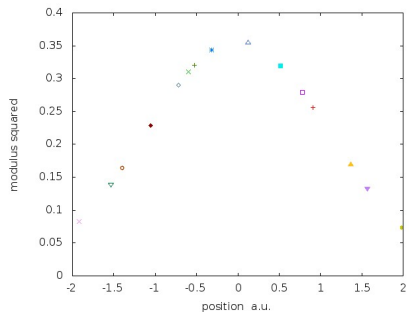
Sampling new initial position

Single Trajectory

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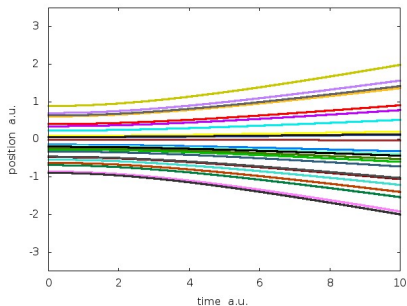


We obtain the wave function

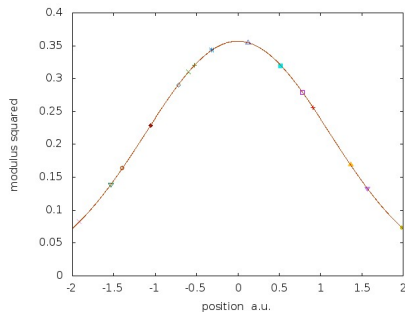


Single Trajectory

Sampling new initial position



We obtain the wave function



Interpolation gives ψ_t at every x

Many Particle - TDSE

$$i\hbar \frac{\partial \Psi(x_1, x_2, \dots, x_N, t)}{\partial t} = \left[- \sum_{i=1}^N \frac{\hbar^2}{2m_i} \nabla_i^2 + V(x_1, x_2, \dots, x_N, t) \right] \Psi(x_1, x_2, \dots, x_N, t)$$

Impossibility to solve numerically the above equation with standard methods!!!

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Using the same Bohmian method for single particle equation
we obtain a system of N coupled equations

$$\left\{ \begin{array}{l} \dots \\ \frac{d^2 x_i}{dt^2} = -\nabla_i [V(x_1, x_2, \dots, x_N, t) + U(x_1, x_2, \dots, x_N, t)]|_{x_1=x_1(t), \dots, x_N=x_N(t)} \\ \dots \end{array} \right.$$

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At this point we have to know Ψ

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We need a new approach!

Many Particle - *Trajectory*

Theorem: Many-particle Bohm trajectory

$x_a(t)$ solution of many particle Schroedinger equation $\Psi(x_a, \vec{x}, t)$

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$x_a(t)$ solution of many particle Schroedinger equation $\Psi(x_a, \vec{x}, t)$

can be calculated from a single particle wave function

$\psi_a(x_a, t) = \Psi(x_a, \vec{x}(t), t)$ solution of the pseudo-Schroedinger equation

$$i\hbar \frac{\partial \psi_a(x_a, t)}{\partial t} = \left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x_a^2} + V_a(x_a, \vec{x}(t), t) + \right. \\ \left. + G_a(x_a, \vec{x}(t), t) + iJ_a(x_a, \vec{x}(t), t) \right] \psi_a(x_a, t)$$

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$$i\hbar \frac{\partial \psi_a(x_a, t)}{\partial t} = \left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x_a^2} + V_a(x_a, \vec{x}(t), t) + G_a(x_a, \vec{x}(t), t) + iJ_a(x_a, \vec{x}(t), t) \right] \psi_a(x_a, t)$$

with a suitable approximation for

$$G_a(x_a, \vec{x}(t), t) \text{ and } iJ_a(x_a, \vec{x}(t), t)$$

we can solve many particle problems.

X. Oriols: *Quantum-trajectory approach to time-dependent transport in mesoscopic systems with electron-electron interactions*, Phys. Rev. Lett. 98, 066803 (2007).

- Quantum Trajectory Method

Kinetic chemistry, dissociation problems.

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- Many Particle
Problems of transport in nanoelectric devices.
eg. G. Albareda et al. : *Many-particle Hamiltonian for open systems with full Coulomb interaction: Application to classical and quantum time-dependent simulations of nanoscale electron devices*, Phys. Rev. B **79**, 075315 (2009)

Von Neumann Model

Interaction between system and apparatus

$x \in X \rightarrow$ system $y \in Y \rightarrow$ apparatus

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$$H_{VN} = -\lambda \hat{A} \otimes \hat{P}_y \quad \text{where} \quad \hat{P}_y \equiv i\hbar \partial / \partial y$$

$$\text{Only } H_{VN} \rightarrow \psi_\alpha(x) \phi_0(y - \lambda a_\alpha t)$$

$$\text{where } A\psi_\alpha = a_\alpha \psi_\alpha$$

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$$H_{meas} = H_0^{syst} + H_0^{app} + H_{VN} \quad \text{where} \quad A(x) = \chi_{[0,+\infty)} - \chi_{(-\infty,0]}$$

We assume that $M \gg m$

Hydrodynamic equations

$$\begin{aligned}\frac{\partial S}{\partial t} &= -\frac{1}{2m} \left(\frac{\partial S}{\partial x} \right)^2 - \frac{1}{2M} \left(\frac{\partial S}{\partial y} \right)^2 + \\ &+ \lambda A(x) \frac{\partial S}{\partial y} + \frac{\hbar^2}{2mR} \frac{\partial^2 R}{\partial x^2} + \frac{\hbar^2}{2MR} \frac{\partial^2 R}{\partial y^2} \\ \frac{\partial \rho}{\partial t} &= -\frac{\partial}{\partial x} \left(\rho \frac{1}{m} \frac{\partial S}{\partial x} \right) - \frac{\partial}{\partial y} \left[\rho \left(\frac{1}{M} \frac{\partial S}{\partial y} + \lambda A(x) \right) \right]\end{aligned}$$

Guidance equations

$$v_x \equiv \frac{1}{m} \frac{\partial S}{\partial x} \quad v_y \equiv \frac{1}{M} \frac{\partial S}{\partial y} + \lambda A(x)$$

Measurement issue

Evolution of the complete wave function *system + apparatus* (Ψ)

$$\Psi_0 = \psi_0(x)\phi_0(y)$$

where

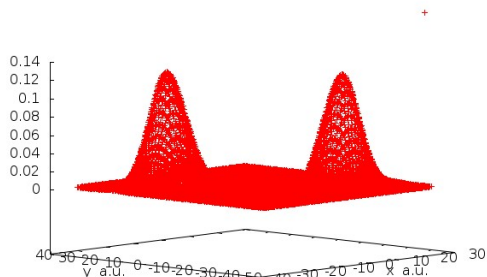
$$\psi_0(x) = \psi_+(x) + \psi_-(x)$$

$$A\psi_+ = +1 \cdot \psi_+$$

$$A\psi_- = -1 \cdot \psi_-$$

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After the interaction we have

$$\Psi_t = \psi_+(x)\phi_+(y) + \psi_-(x)\phi_-(y)$$

Wave Function of the Measured System $\psi_{\text{cond}} = \frac{\Psi_t(x, Y_t)}{\|\Psi_t(x, Y_t)\|}$

$$\psi_+$$

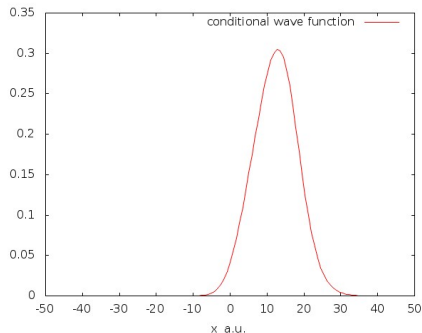
$$\psi_-$$

$$X_0 = 6.6 \text{ a.u.} \quad Y_0 = 0.6 \text{ a.u.}$$

$$X_0 = -5.2 \text{ a.u.} \quad Y_0 = 0.6 \text{ a.u.}$$

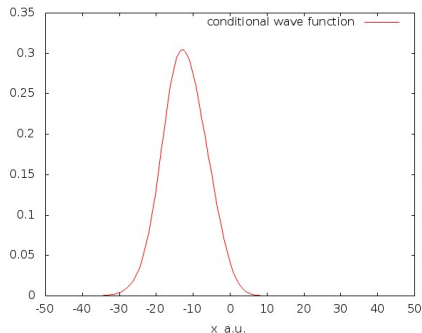
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ψ_+



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Thank you for your attention!