# A quantum cellular automaton 

## extension of quantum field theory

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## INFORMATION-THEORETICAL PRINCIPLES

## Principles for Quantum Theory

Pi. Causality
P2. Local Readability


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\text { PRA } 84 \text { OI23II (2OI2) }
$$

$\mathrm{P}_{3}$. Conservation of Information (Purification)
P4. Indivisibility of Composition
$\mathrm{P}_{5}$. Discriminability of Specific Ini
P6. Lossless Compressibility


Algorithm

## INFORMATIO

Principles for Qua
Pi. Causality
P2. Local Readability
$\mathrm{P}_{3}$. Conservation of Inform
Topological Homogeneity


Physical law

Deutsch-Cburch-Turing principle Every finite physical protocol must be perfectly simulated by a quantum computer made with a finite number of qubits and a finite number of gates

- 2uantum-information density is bounded Locality of interactions


Algorithm

## FINITE INFORMATION

Localized states over a locally quiescient state (vacuum)


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## PROGRAM: <br> QCA-EXTENSION OF QFT



In this talk


THE DIRAC QU-AUTOMATON

$\boldsymbol{\psi}_{n}:=\left[\begin{array}{l}\psi_{n}^{+} \\ \psi_{n}^{n}\end{array}\right]$ $\psi_{n-1}^{t} \psi_{n-1}^{-1} \psi_{t}^{4} \psi_{n} \psi_{n+1}^{+1} \psi_{n+1} \cdots$

$$
\boldsymbol{\psi}(t+\mathfrak{t})=U^{\dagger} \boldsymbol{\psi}(t) U=\mathbf{U} \boldsymbol{\psi}(t)
$$

$$
c^{2}+s^{2}=1
$$

$$
\begin{aligned}
& A_{n}=\exp \left[-i \theta\left(\sigma_{2 n-1}^{-} \sigma_{2 n}^{+}+\sigma_{2 n-1}^{+} \sigma_{2 n}^{-}\right)\right] \\
& B_{n}=\exp \left[-i \frac{\pi}{2}\left(\sigma_{2 n}^{+} \sigma_{2 n+1}^{-}+\sigma_{2 n}^{-} \sigma_{2 n+1}^{+}\right)\right]
\end{aligned}
$$

$$
\mathbf{U}=\left[\begin{array}{cc}
\mathrm{S} \hat{\partial}_{-} & -i \mathrm{c} \\
-i \mathrm{c} & \mathrm{~s} \hat{\partial}_{+}
\end{array}\right]
$$

$c=\cos \theta=\omega t=\frac{\mathfrak{a}}{\lambda}=\frac{m}{\mathfrak{m}}, \quad s=\sin \theta=\zeta=\sqrt{1-\left(\frac{m}{m}\right)^{2}}$

$$
\mathfrak{m}:=\frac{\hbar}{\mathfrak{a} \mathfrak{c}}
$$



## THE DIRAC QU-AUTOMATON



$$
c^{2}+s^{2}=1
$$

$$
A_{n}=\exp \left[-i \theta\left(\sigma_{2 n-1}^{-} \sigma_{2 n}^{+}+\sigma_{2 n-1}^{+} \sigma_{2 n}^{-}\right)\right]
$$

$$
B_{n}=\exp \left[-i \frac{\pi}{2}\left(\sigma_{2 n}^{+} \sigma_{2 n+1}^{-}+\sigma_{2 n}^{-} \sigma_{2}^{+}, \quad\right)\right]
$$

$$
\mathbf{U}=\left[\begin{array}{ll}
\mathrm{s} \widehat{\partial}_{-} & -i \mathrm{c} \\
-i \mathrm{c} & \mathrm{~s} \widehat{\partial}_{+}
\end{array}\right]
$$

Planck length
$\mathrm{c}=\cos \theta=\omega \mathfrak{t}=\frac{\mathfrak{a}}{\lambda}=\frac{m}{\mathfrak{m}}, \quad \mathrm{~s}=$
$\hbar:=\mathfrak{m a c}$

## PARTICLES

## Dirac QCA- hlist Quanlration



## Dirac QCA: First Quantization

| A | $\left[\begin{array}{ll}\mathrm{s} & \\ & -i \mathrm{c}\end{array}\right]$ |
| :---: | :---: |
| B | $\mathbf{U}=\left[\begin{array}{ll}\mathrm{s} \hat{\partial}_{-} & \\ -i \mathrm{c} & \mathrm{s} \widehat{\partial}_{+}\end{array}\right]$ |

## Planckian Particles

(Foldy-Wouthuysen) $\mathbf{U}=\left(\begin{array}{cc}\mathrm{s} \widehat{\partial}_{-} & -i \mathrm{c} \\ -i \mathrm{c} & \mathrm{s} \widehat{\partial}_{+}\end{array}\right)=\left(\begin{array}{cc}\mathrm{s} e^{i k} & -i \mathrm{c} \\ -i \mathrm{c} & \mathrm{s} e^{-i k}\end{array}\right)$
$\frac{1}{2 N^{( \pm)}(k)}\left(\left.\begin{array}{c}i \mathrm{c} \\ \left.\mathrm{s} e^{i k}-e^{ \pm i \omega(k)}\right) \\ \text { igenenvectors } k \text {-space }\end{array} \right\rvert\,\right.$

Dispersion relation
$\omega(k)=\cos ^{-1}(\mathrm{~s} \cos \mathrm{k})$

## Alessandro Bisio <br> Alessandro Tosini



## Planckian Particles

$$
i \partial_{t} A(x, t)=\left[-i d \partial_{x}-\frac{D}{2} \partial_{x}^{2}\right] A(x, t)
$$

$$
A(x, t)=\frac{1}{\sqrt[4]{2 \pi \Delta^{2}(t)}} \exp \left[-\frac{(x-x(t))_{2}^{2}}{4 \Delta^{2}(t)}\right]
$$

$$
x(t)=d \quad \Delta(t)=\Delta \sqrt{1+i \frac{D}{2 \Delta^{2}} t}
$$

$$
d^{( \pm)}(\mathrm{k})
$$

$\pi$

## MECHANICS EMERGING FROM COMPUTATION

$i \hbar \hat{\partial}_{t} \boldsymbol{\psi}=[\boldsymbol{\psi}, H]$


$$
H=\frac{i \hbar}{2 t} \boldsymbol{\psi}^{\dagger}\left(\mathbf{U}-\mathbf{U}^{\dagger}\right) \boldsymbol{\psi}
$$

## MECHANICS EMERGING FROM COMPUTATION

- PATH-INTEGRAL



## Are we able to simulate our theory (even with a quantum computer)?

## Simulating Physics with Computers Richard P. Feynman

The question is, if we wrote a Hamiltonian which involved only these operators, locally coupled to corresponding operators on the other space-time points, could we imitate every quantum mechanical system which is discrete and has a finite number of degrees of freedom? I know, almost certainly, that we could do that for any quantum mechanical system which involves Bose particles. I'm not sure whether Fermi particles could be described by such a system. So I leave that open. Well, that's an example of what I meant by a general quantum mechanical simulator. I'm not sure that it's sufficient, because I'm not sure that it takes care of Fermi particles.


$$
\text { Int. J. of Th. Phys. } 21467 \text { (1982) }
$$

## SPACE-TIME AT PLANCK SCALE

FOLIATION: TIME AS A COMPUTER CLOCK


## SPACE-TIME AT PLANCK SCALE

## THE COMPUTATIONAL TOMONAGA-SCHWINGER



## SPACE-TIME AT PLANCK SCALE TIME-DILATION AND SPACE-CONTRACTION



## Why information is quantum?

Should we consider a network-axiom for QT?

- "Direction" of information imprinted in the state using minimal informational resources.
- A थuantum-Digital World: restoration of the isotropy of information flow.



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First Quantization two particle states


First Quantization particle antiparticle



First Quantization particle antiparticle



First Quantization two particle states

## THANK YOU!

