

QCD@WORK 2012

INTERNATIONAL WORKSHOP ON QCD - THEORY AND EXPERIMENT

LECCE - JUNE 18-21, 2012

FLAVOUR VIOLATION PATTERNS
IN THE LIGHT OF RECENT LHCb DATA



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- New measures and old tensions in flavour data

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- Minimal Flavour Violation at large
 - * Two-Higgs-Doublet Model with MFV and Flavour-Blind Phases

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- Beyond Minimal Flavour Violation
 - * Gauged Flavour Symmetries

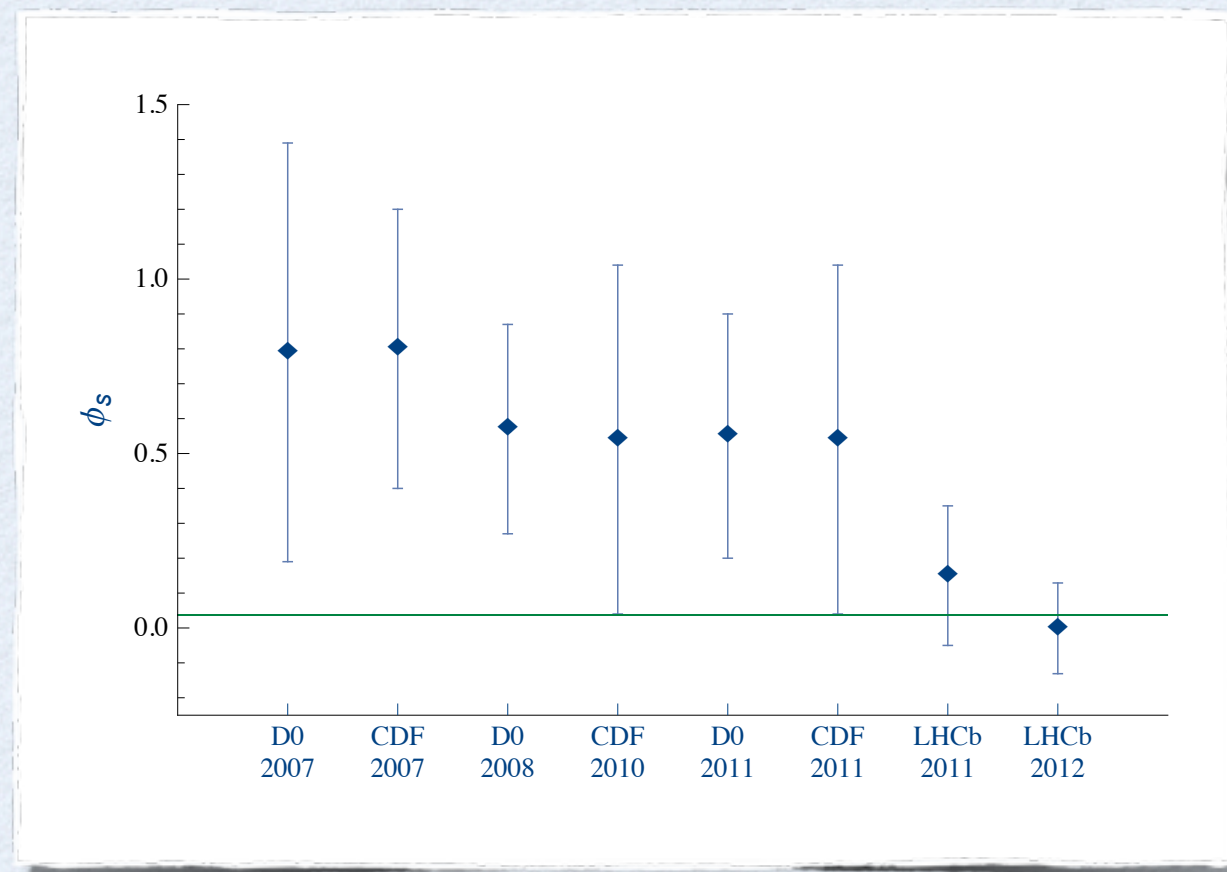
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 - * Two-Higgs-Doublet Model with MFV and Flavour-Blind Phases
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- Conclusions and Outlook

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$S_{\psi\phi}$: mixing-induced CP asymmetry in the B_s system

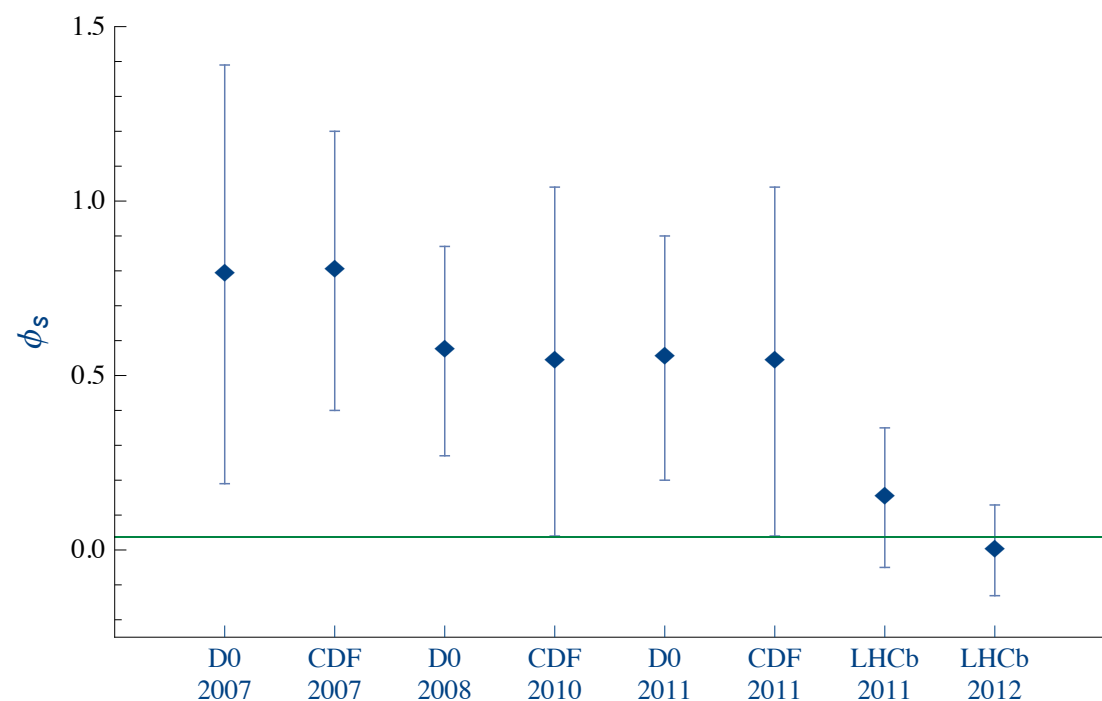


LHCb Collab., PRL **108** (2012)

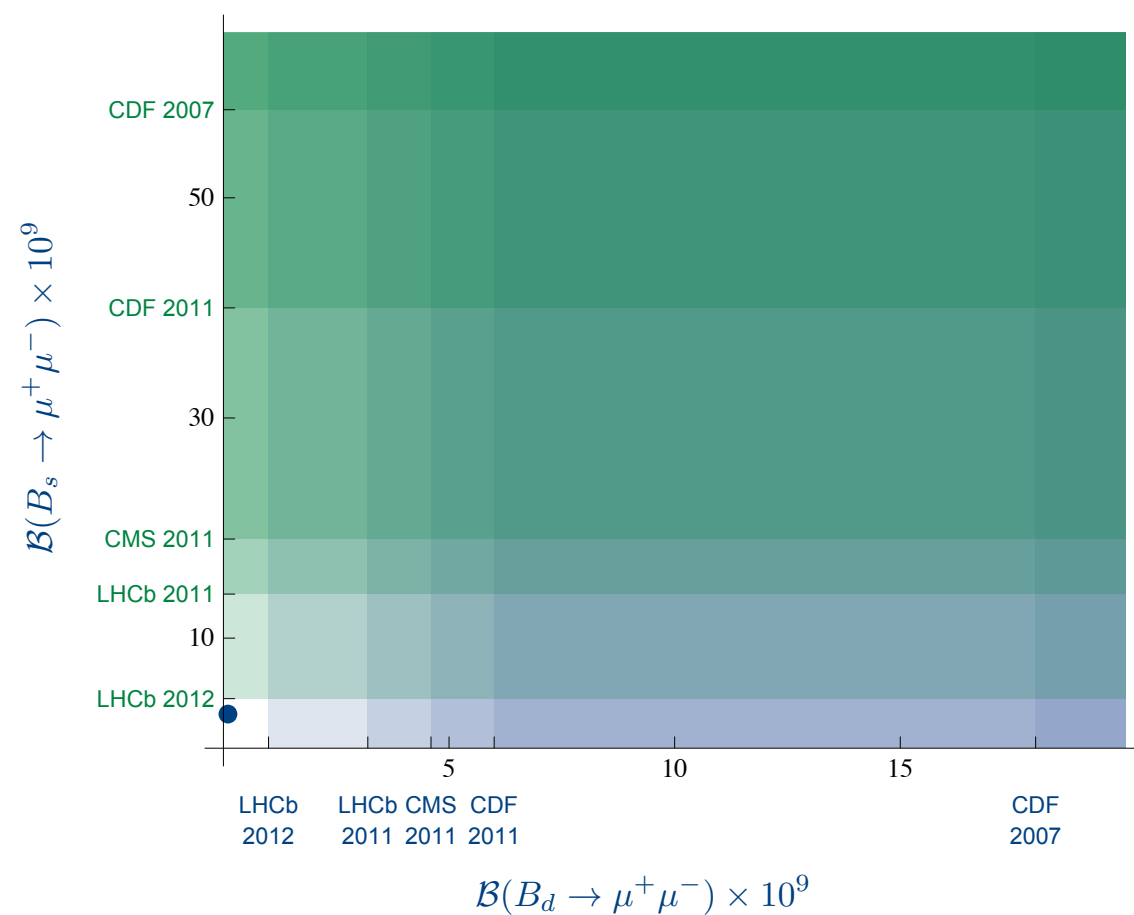
LHCb Collab., LHCb-CONF-2012-002

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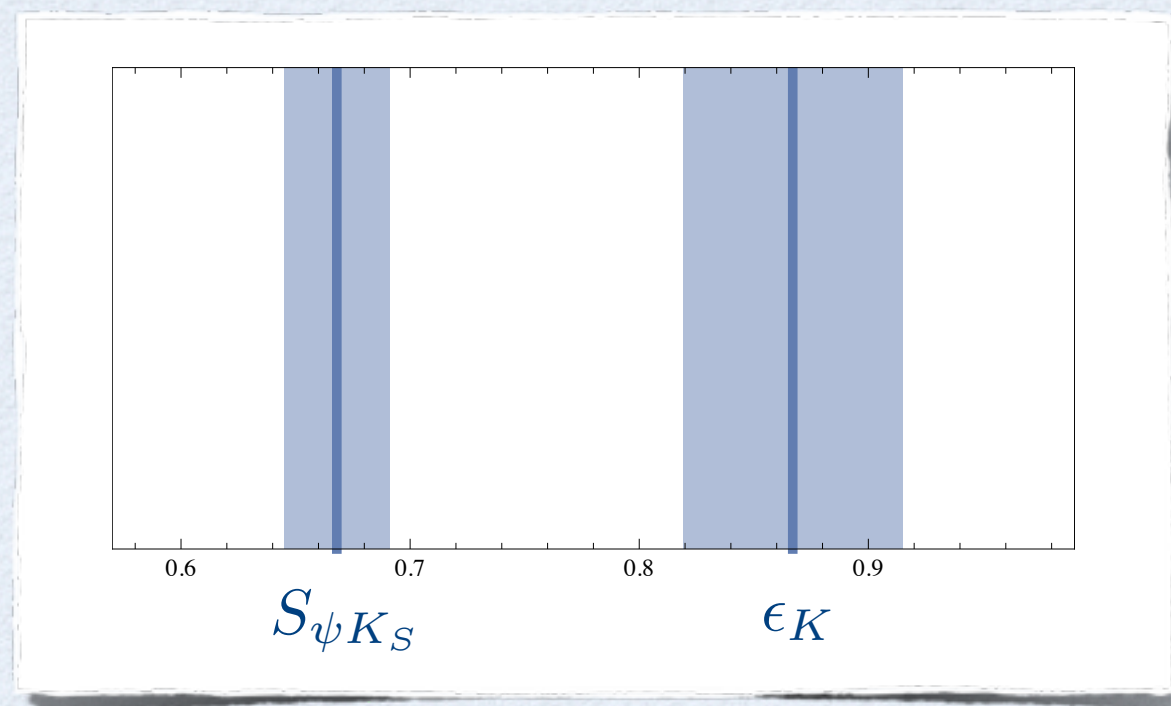
$B_{d,s} \rightarrow \mu^+\mu^-$: very sensible to contributions of new heavy particles

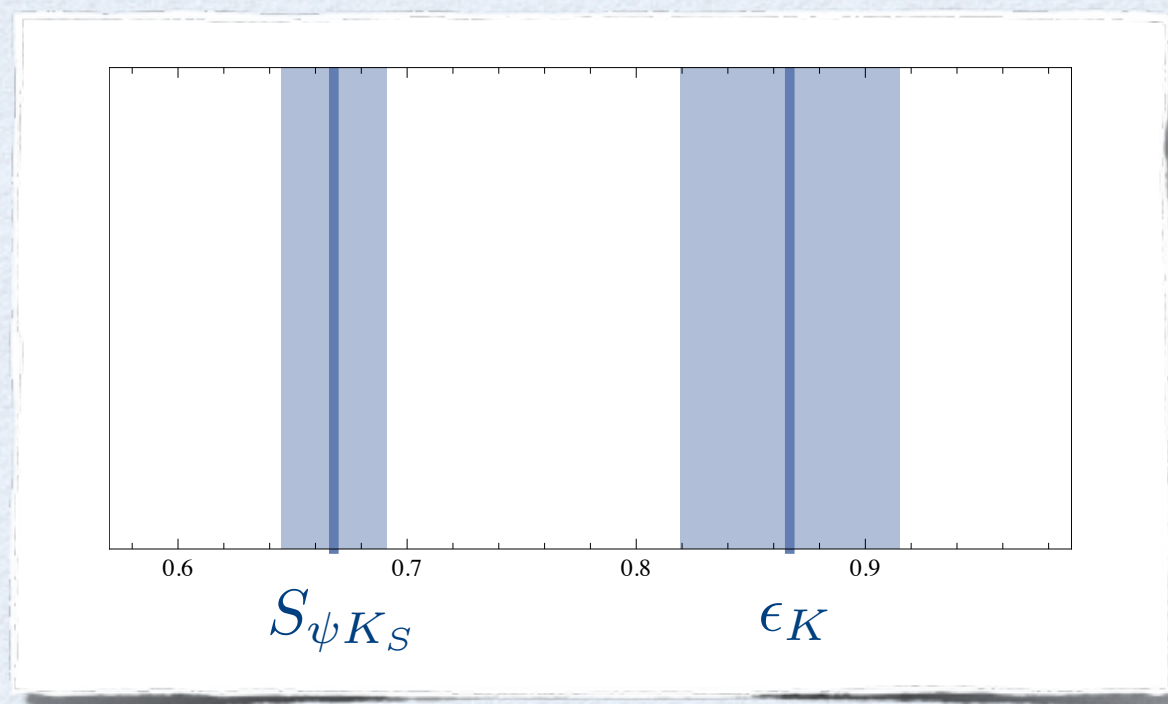


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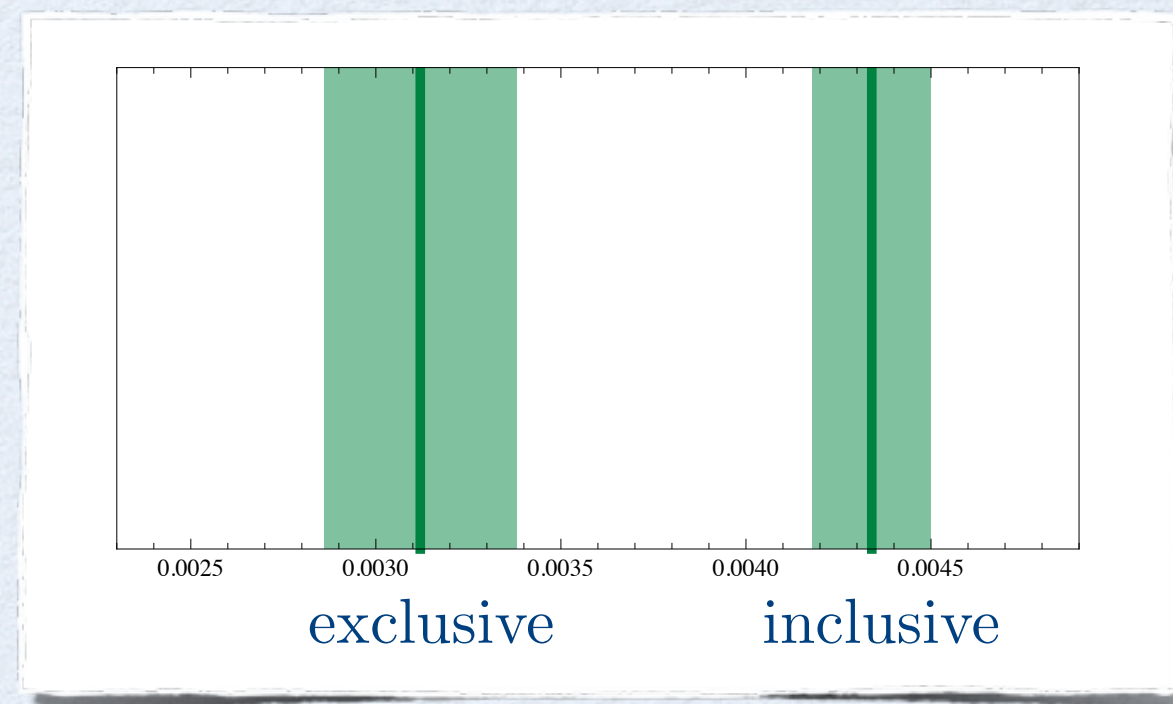
LHCb Collab., Phys. Lett. B **708** (2012)
LHCb Collab., 1203.4493

TENSIONS IN THE UNITARITY TRIANGLE: $\varepsilon_K - S_{\psi K_S}$ VS. V_{ub}

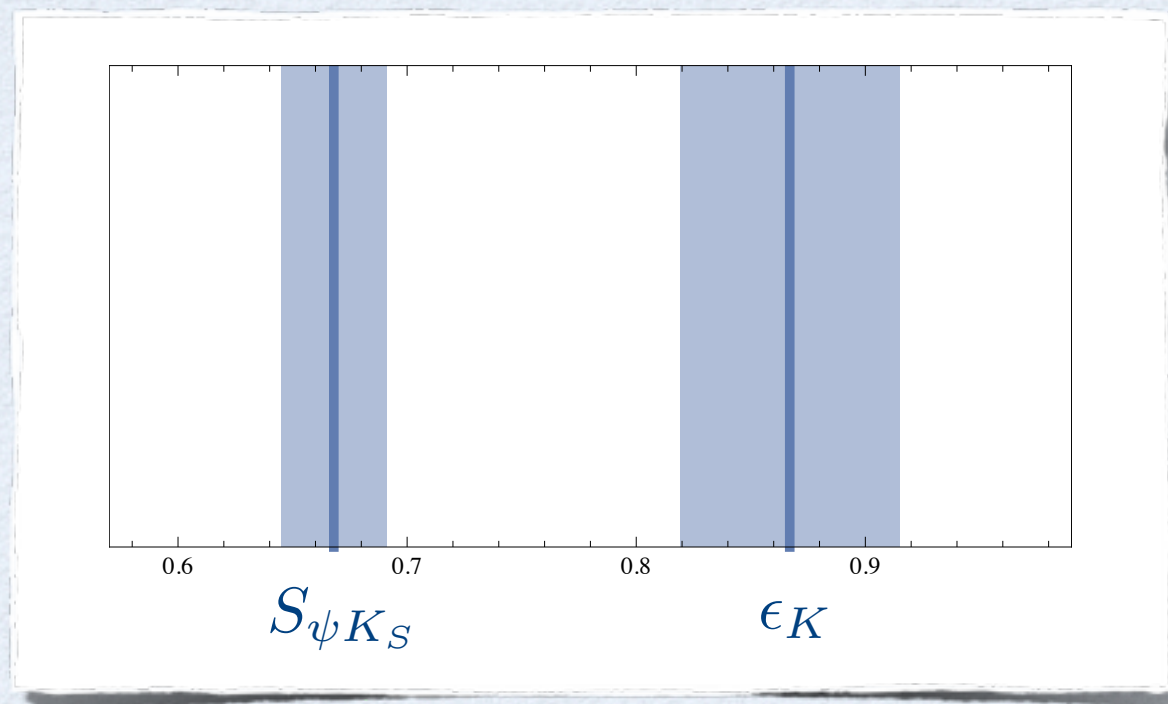
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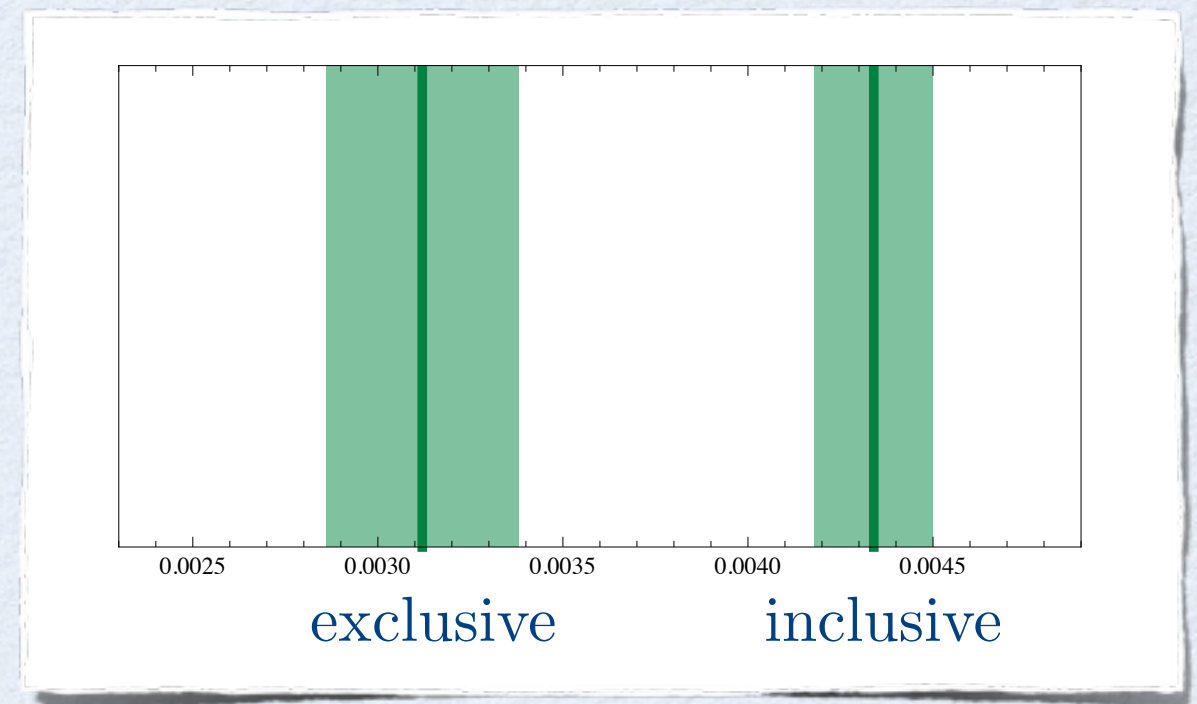
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Scenario 1

exclusive V_{ub}
 ϵ_K below the data
 $S_{\psi K_S}$ in agreement with data

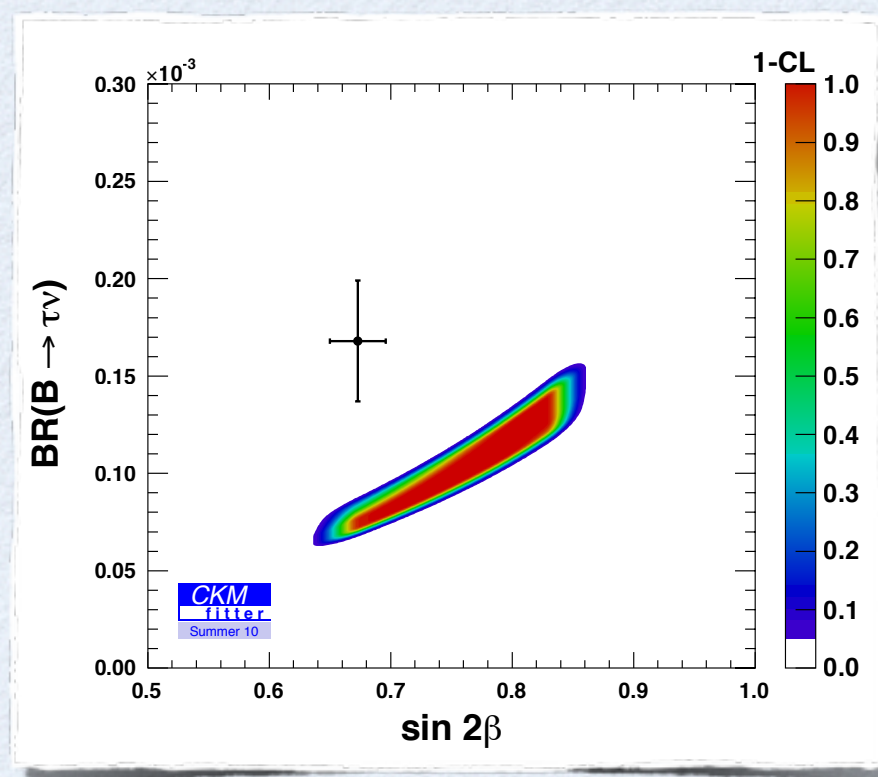
Scenario 2

inclusive V_{ub}
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MORE TENSIONS

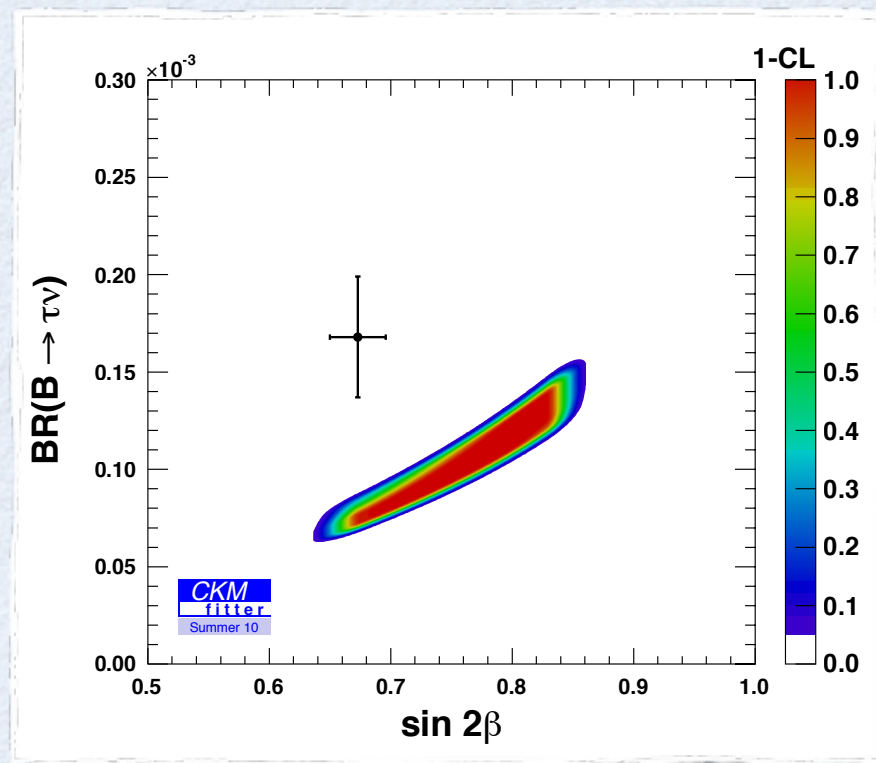
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2.9 σ discrepancy
between $\text{BR}(B \rightarrow \tau \nu)$ and
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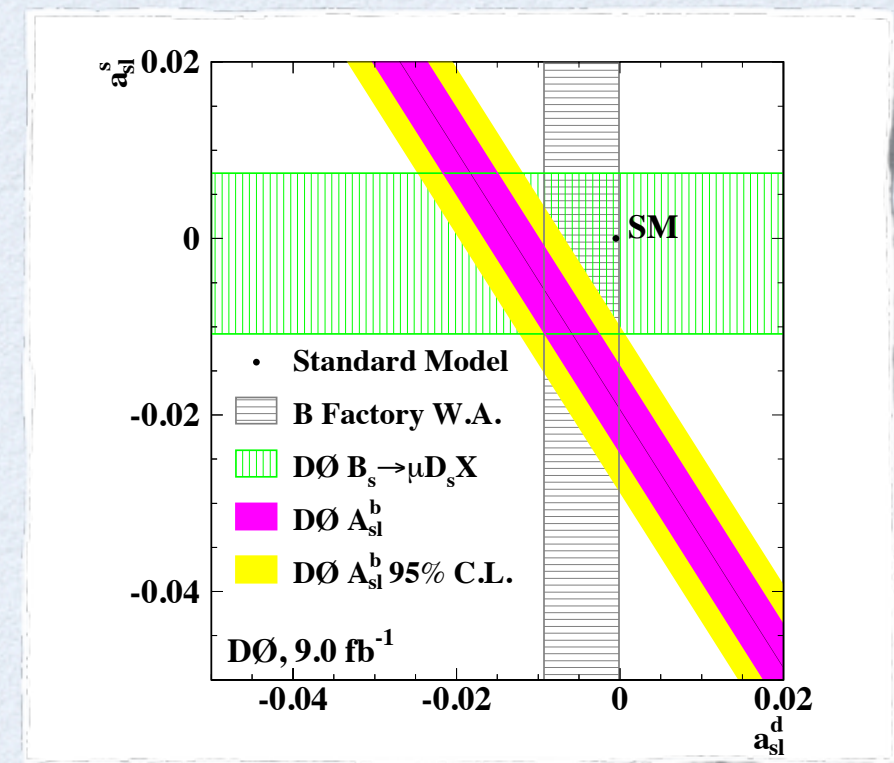


Lenz, Nierste and CKMfitter, 1203.0238

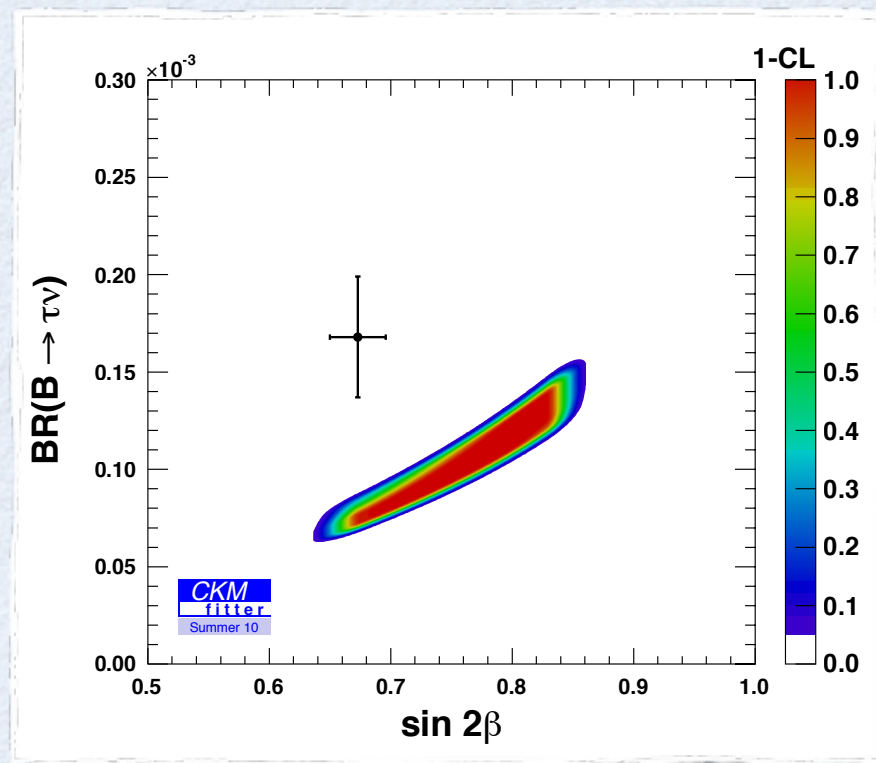
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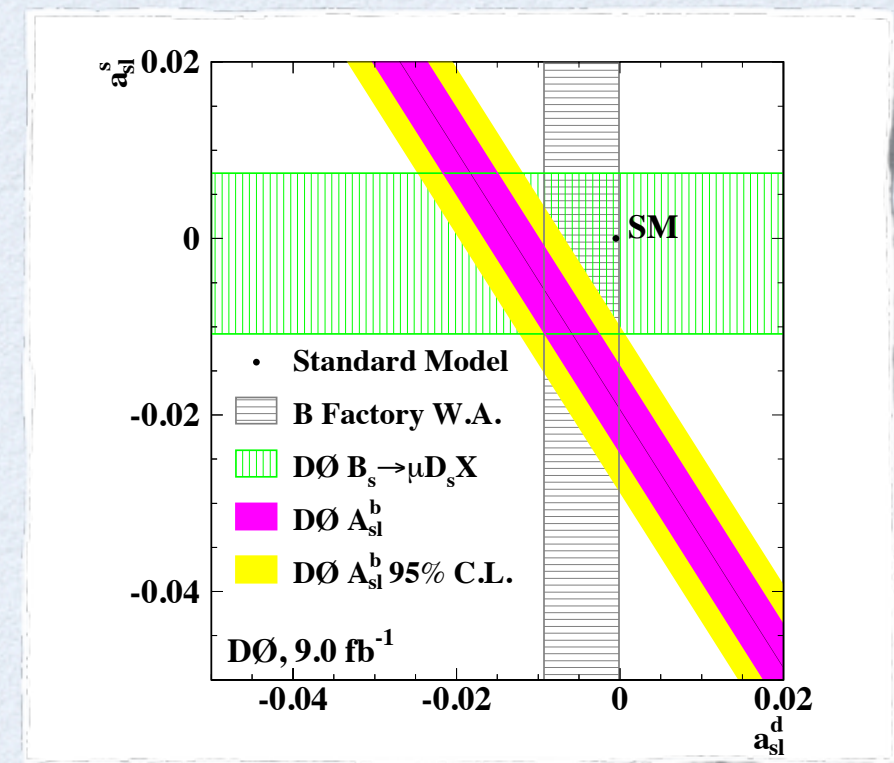
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Buras and Girschbach, 1204.5064

CONSTRAINED MINIMAL FLAVOUR VIOLATION

DEFINITION AND IMPLICATIONS

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All flavour changing transitions are governed by the **CKM matrix** with the CKM phase being the only source of CP violation.

The only relevant **operators** in the effective Hamiltonian below the weak scale are those that are also relevant in the SM.

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All flavour changing transitions are governed by the **CKM matrix** with the CKM phase being the only source of CP violation.

The only relevant **operators** in the effective Hamiltonian below the weak scale are those that are also relevant in the SM.



- $S_{\psi K_S}$ and $S_{\psi\phi}$ are as in the SM.
- For fixed CKM parameters determined in tree-level decays, $|\varepsilon_K|$, ΔM_s and ΔM_d **can only be enhanced** relative to SM predictions, and this happens in a correlated manner.
- There are **correlations** between various observables:

$$\frac{\mathcal{B}(B_s \rightarrow \mu^+ \mu^-)}{\mathcal{B}(B_d \rightarrow \mu^+ \mu^-)} = \frac{\tau(B_d)}{\tau(B_s)} \frac{m_{B_d}}{m_{B_s}} \frac{F_{B_d}^2}{F_{B_s}^2} \left| \frac{V_{td}}{V_{ts}} \right|^2 \boxed{r(\mu^+ \mu^-)} = \frac{\hat{B}_d}{\hat{B}_s} \frac{\tau(B_d)}{\tau(B_s)} \frac{\Delta M_s}{\Delta M_d} \frac{r(\Delta M)}{r(\mu^+ \mu^-)}$$

$$\frac{\Delta M_d}{\Delta M_s} = \frac{m_{B_d}}{M_{B_s}} \frac{\hat{B}_d}{\hat{B}_s} \frac{F_{B_d}^2}{F_{B_s}^2} \left| \frac{V_{td}}{V_{ts}} \right|^2 \boxed{r(\Delta M)}$$

CONSTRAINED MINIMAL FLAVOUR VIOLATION

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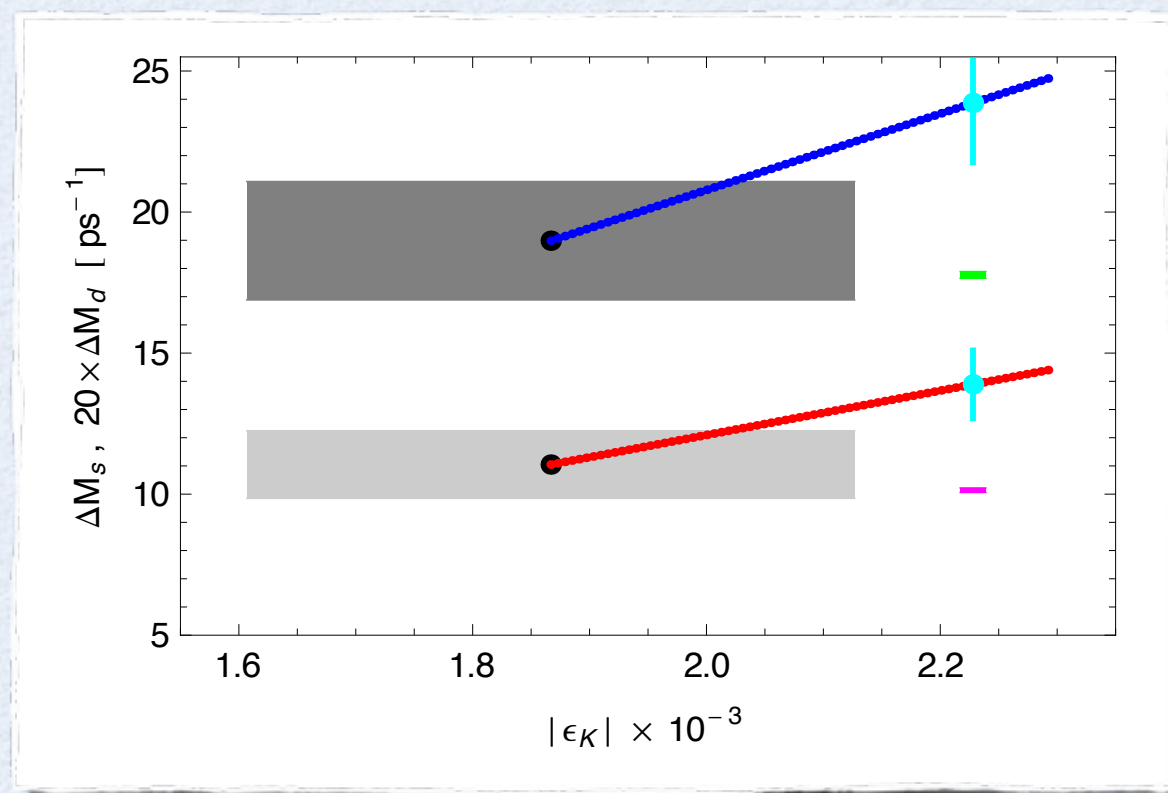
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- $S_{\psi K_S}$ cannot be modified \longrightarrow Scenario 1 for V_{ub}
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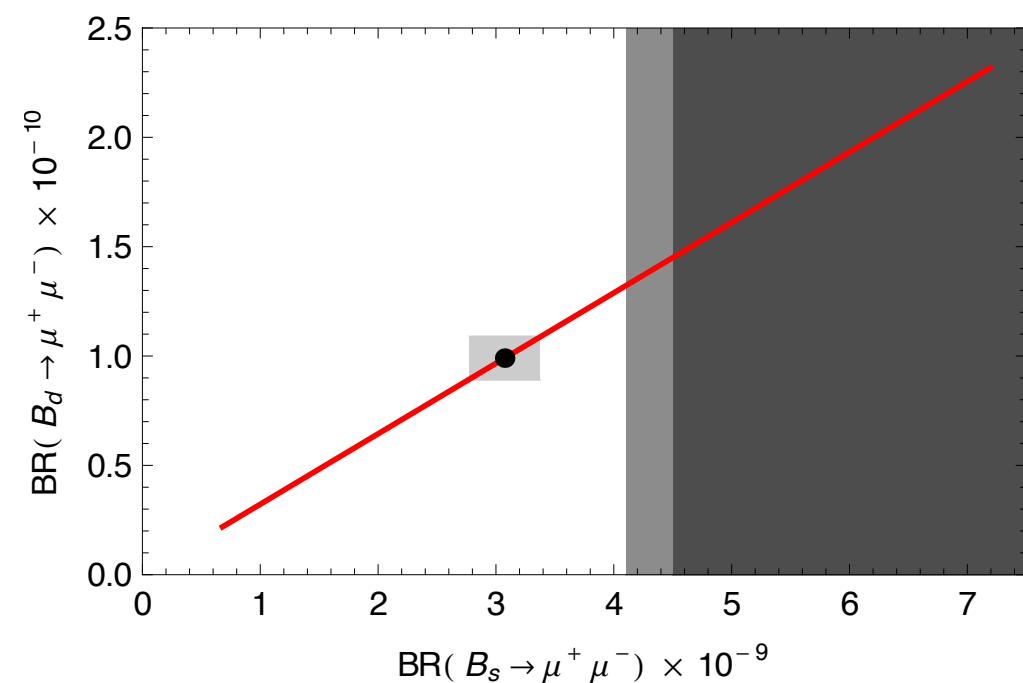
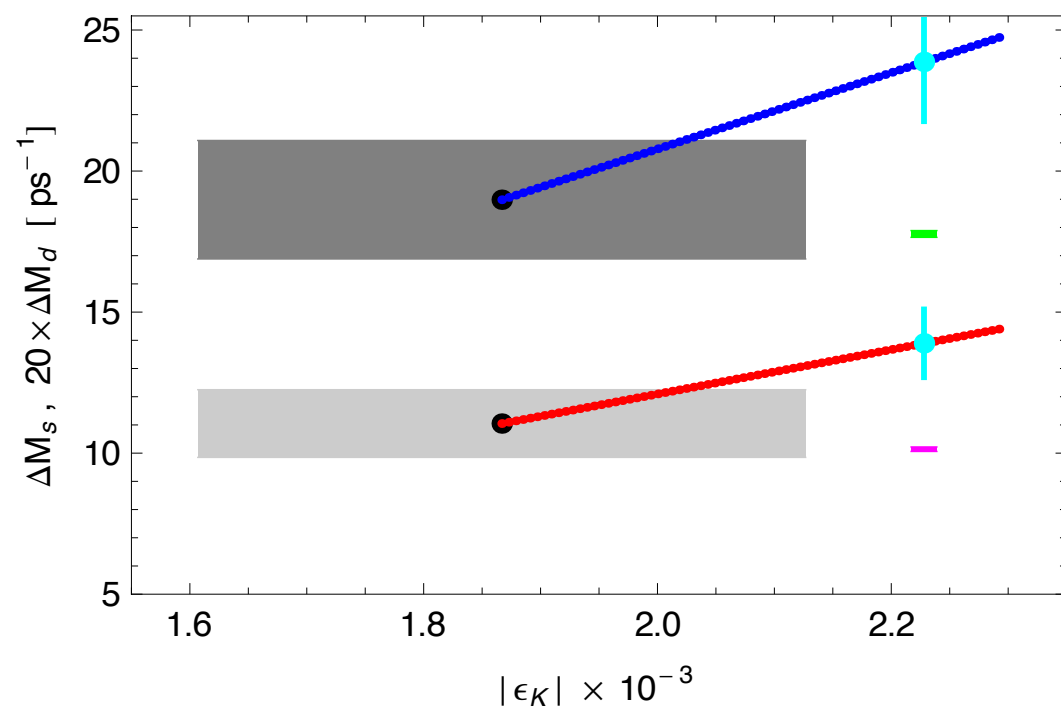


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Correlations:
an independent test
of CMFV



MFV AT LARGE: $2\text{HDM}_{\overline{\text{MFV}}}$

DEFINITION AND APPLICATION TO 2HDM

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Successful symmetry + symmetry-breaking pattern of the SM:

large global flavour symmetry in the gauge sector

$$\mathcal{G}_q = (SU(3) \otimes U(1))^3$$

broken only by the Yukawa couplings

$$\mathcal{L}_Y = -\bar{Q}_L Y_d D_R H - \bar{Q}_L Y_u U_R H^c$$

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Flavour symmetry is formally recovered by promoting the Yukawa couplings to **spurions**

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Application to 2HDM:

$$\mathcal{L}_{\text{MFV}}^{\text{FCNC}} = \frac{1}{\sin \beta} \bar{d}_L^i \left[\left(a_0 V^\dagger \lambda_u^2 V + a_1 V^\dagger \lambda_u^2 V \Delta + a_0 \Delta V^\dagger \lambda_u^2 V \right) \lambda_d \right]_{ij} d_R^j \frac{S_2 + iS_3}{\sqrt{2}} + \text{h.c.}$$

double CKM **suppression** + down-type Yukawa **suppression**

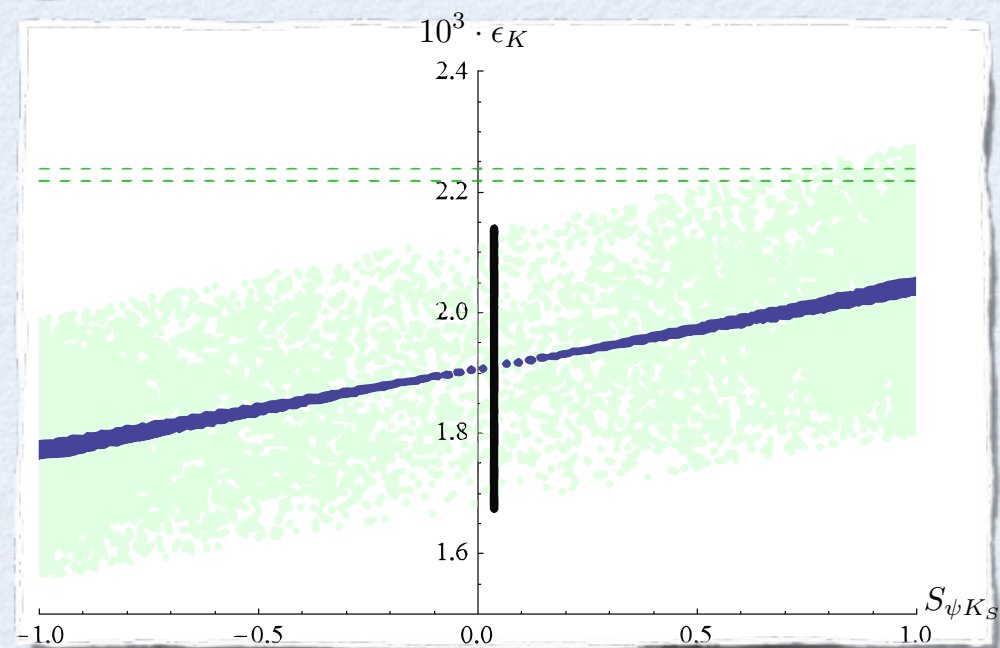
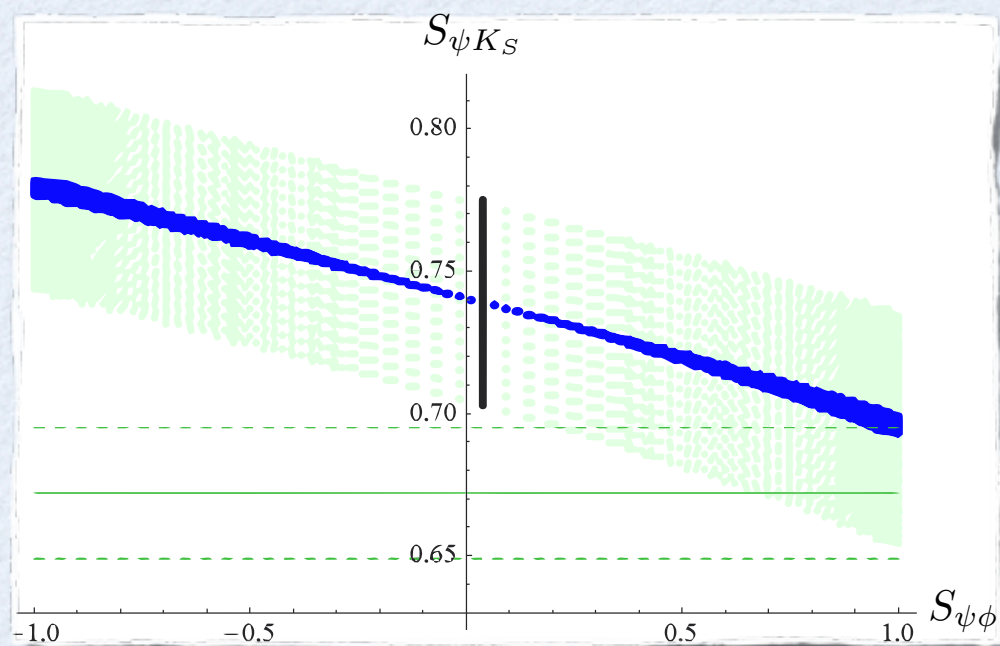
In general \mathbf{a}_i can be complex \longrightarrow **Flavour Blind Phases**

MFV AT LARGE: $2\text{HDM}_{\overline{\text{MFV}}}$

PHENOMENOLOGICAL ANALYSIS

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Room for a large $S_{\psi\phi}$ that automatically solves the $S_{\psi K_S} - \epsilon_K$ tension!

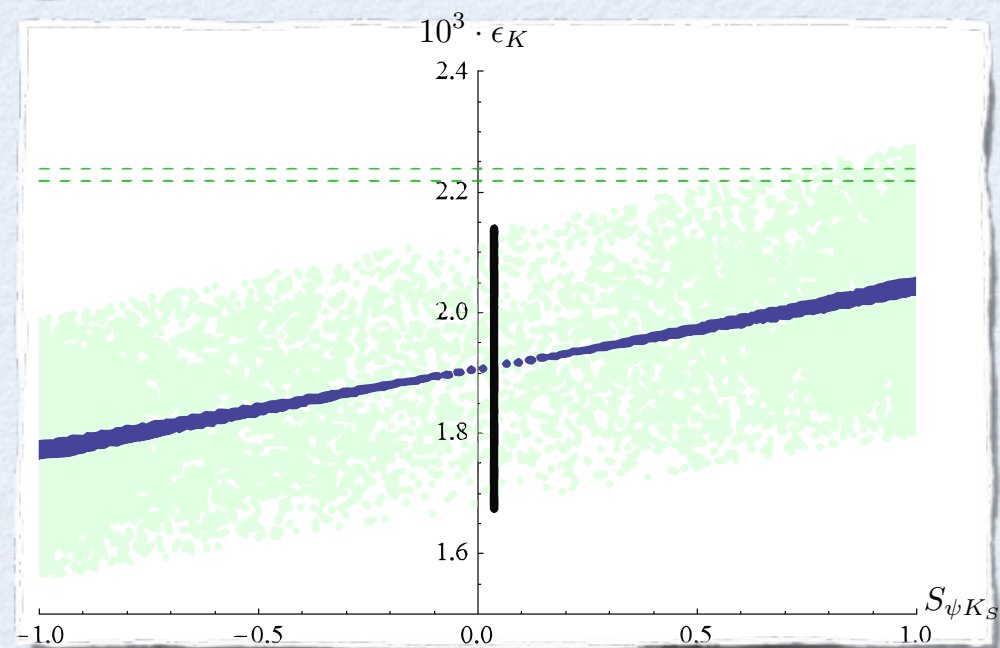
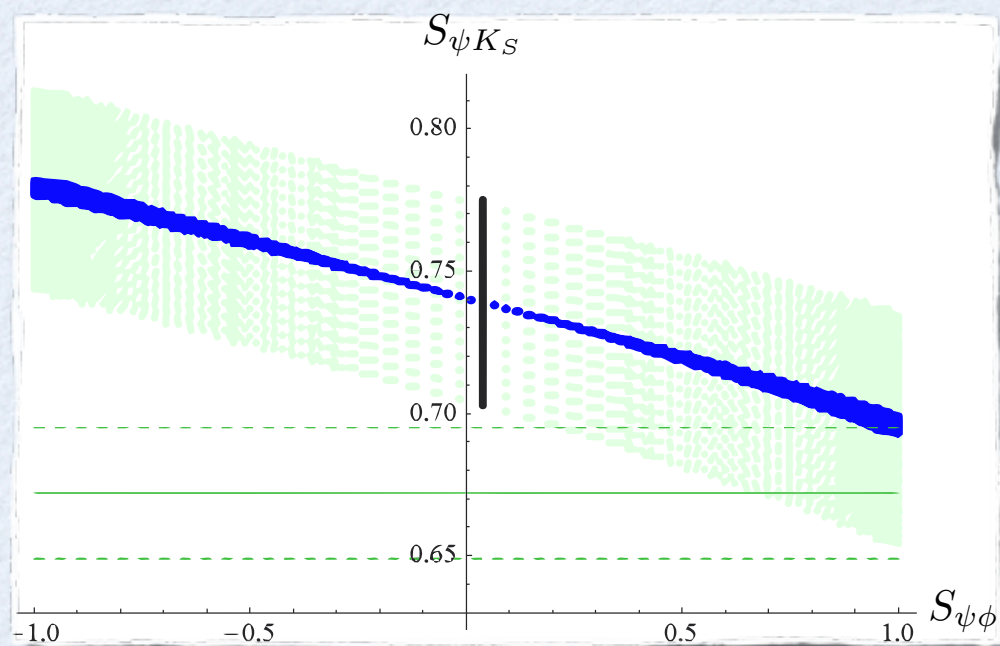


Buras, MVC, Gori and Isidori, JHEP 1010

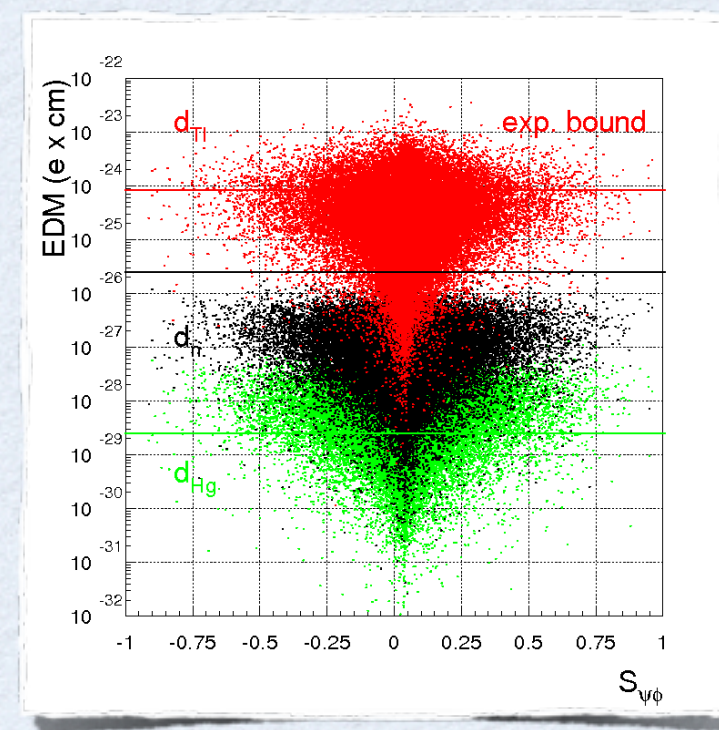
Buras, Isidori and Paradisi, Phys. Lett. B **694** (2011)

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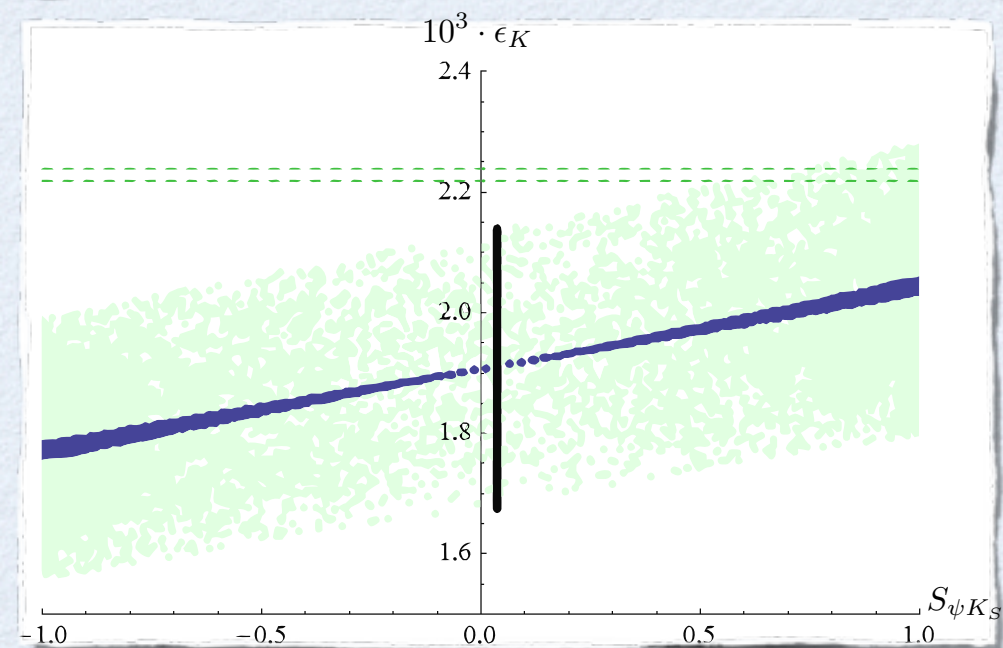
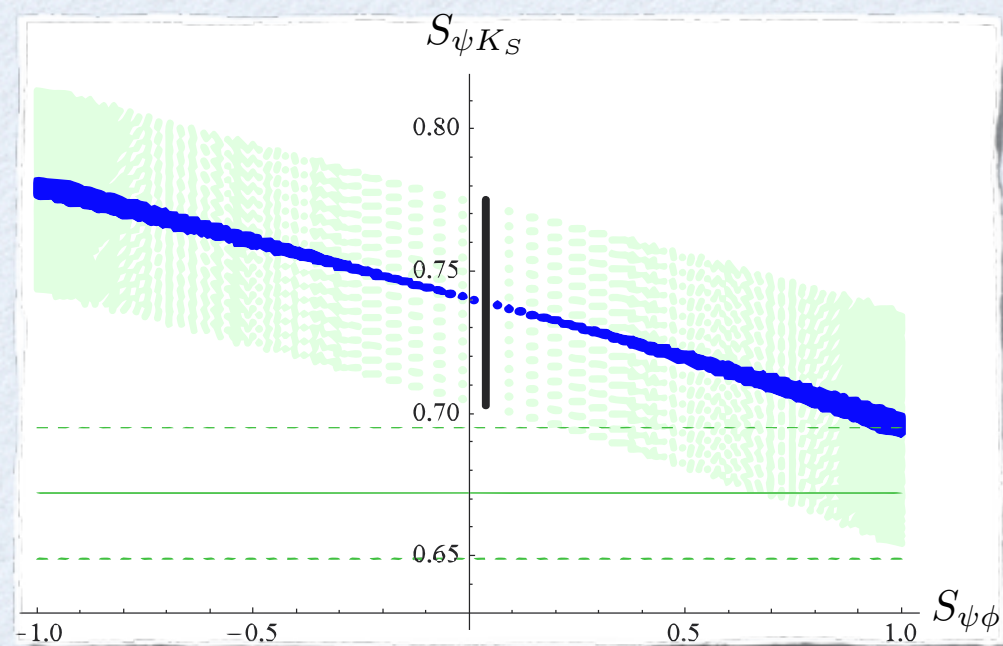


A large $S_{\psi\phi}$ does not spoil the EDMs

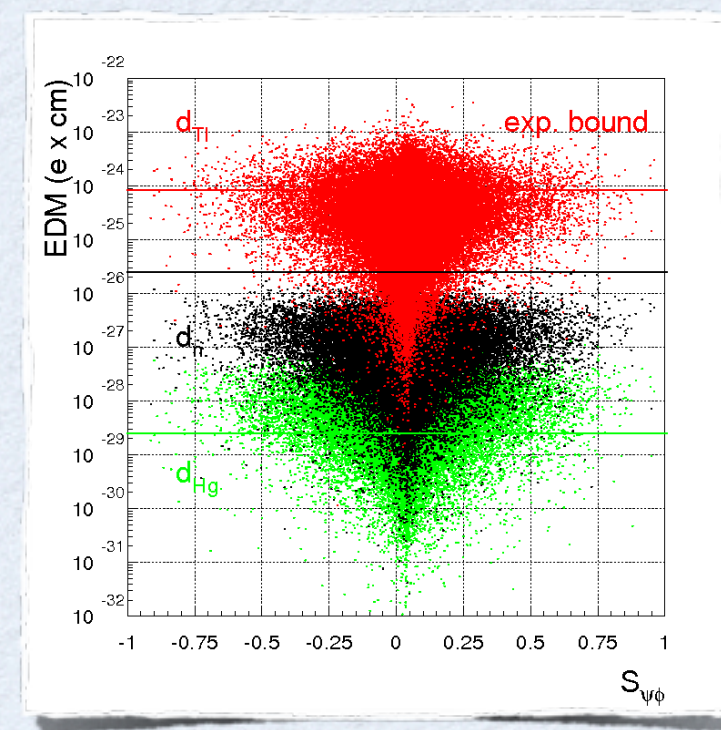


PHENOMENOLOGICAL ANALYSIS

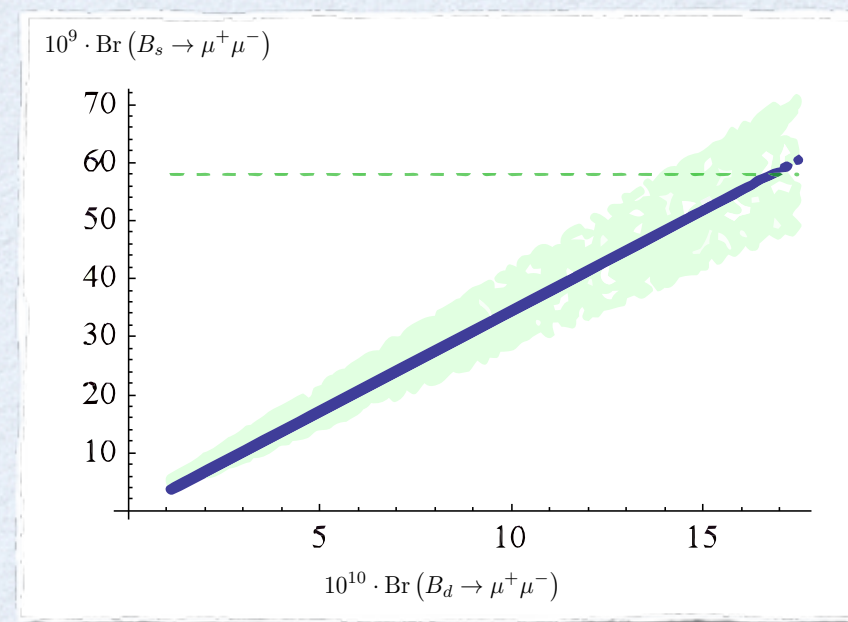
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$B_{d,s} \rightarrow \mu^+ \mu^-$ correlation



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UPDATES AFTER LHCb

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Negligible contributions to ε_K

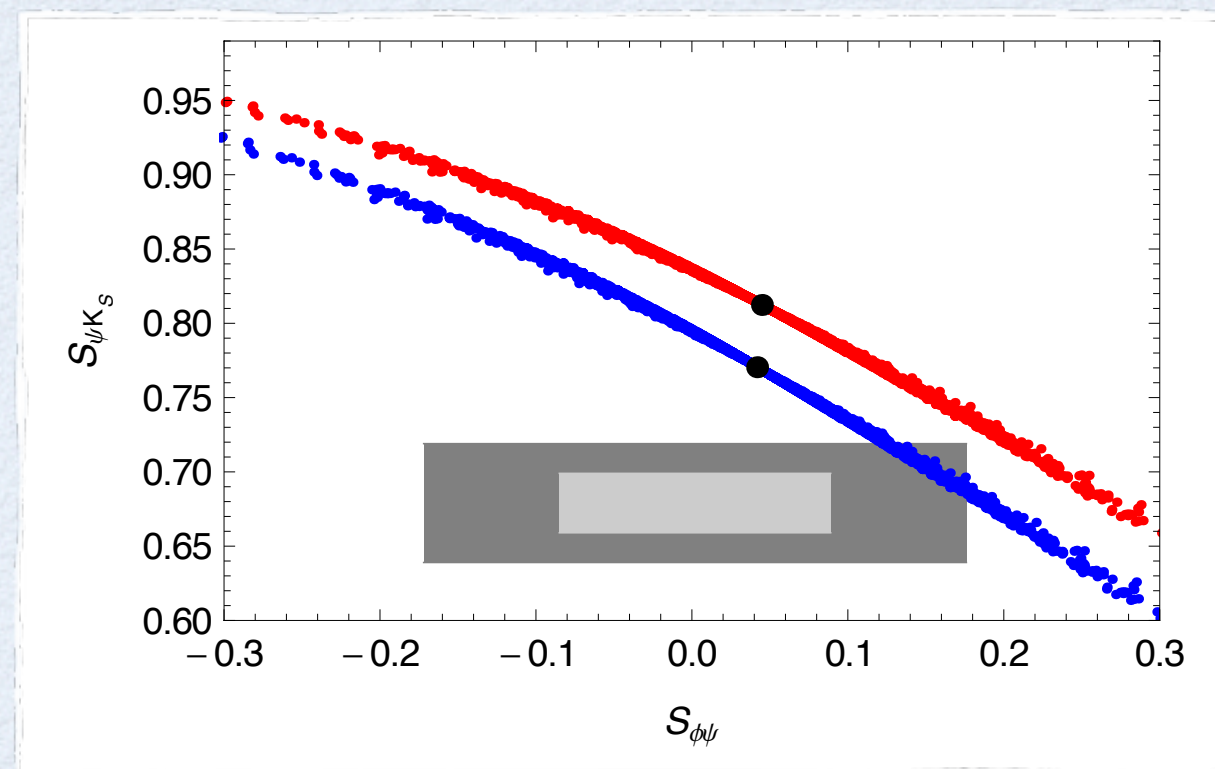


Scenario 2 for V_{ub}



softening of the $\text{BR}(B \rightarrow \tau \nu)$ tension

UPDATES AFTER LHCb

Negligible contributions to ε_K Scenario 2 for V_{ub} softening of the $\text{BR}(B \rightarrow \tau \nu)$ tensionNo large $S_{\psi\phi}$ the $S_{\psi K_s} - \varepsilon_K$ tension cannot be improved better than to 2σ 

THE MODEL

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Structure

The flavour group G_q is a gauge symmetry of Nature.

Y_u and Y_d are physical scalar fields, whose vevs break G_q .

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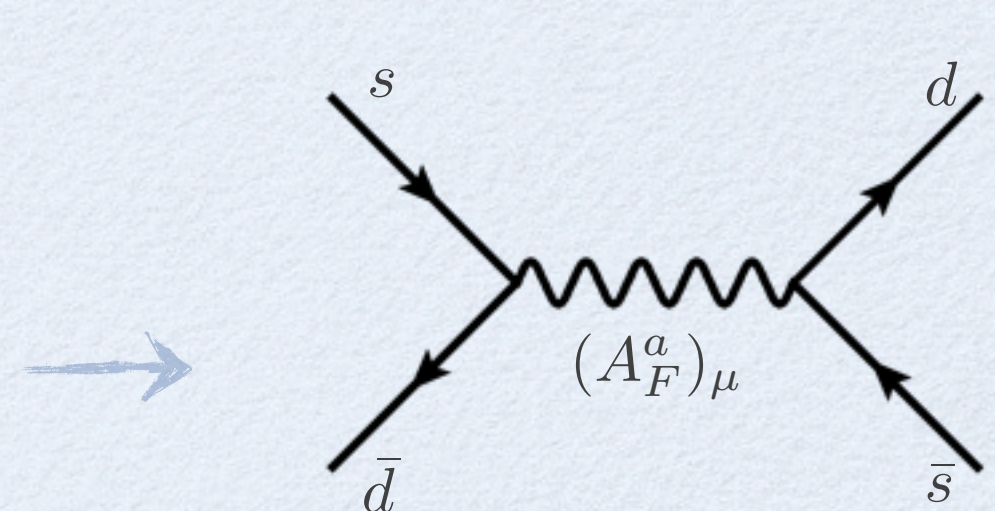
The flavour gauge bosons mediate **FCNCs**, but the new quarks generate a mechanism of **inverted hierarchy** that suppresses them effectively.

SM fermion mass terms:

$$\frac{\langle H \rangle}{\langle Y_D \rangle} \bar{Q}_L D_R \rightarrow \langle Y_F \rangle \propto \frac{1}{\text{SM Yukawas}}$$

Boson mass terms:

$$\langle H \rangle^2 \langle Y_D \rangle^2 (A_D)_\mu (A_D)^\mu \rightarrow m_{A_F} \propto \frac{1}{\text{SM Yukawas}}$$



Suppressed by (down SM Yukawa)⁴

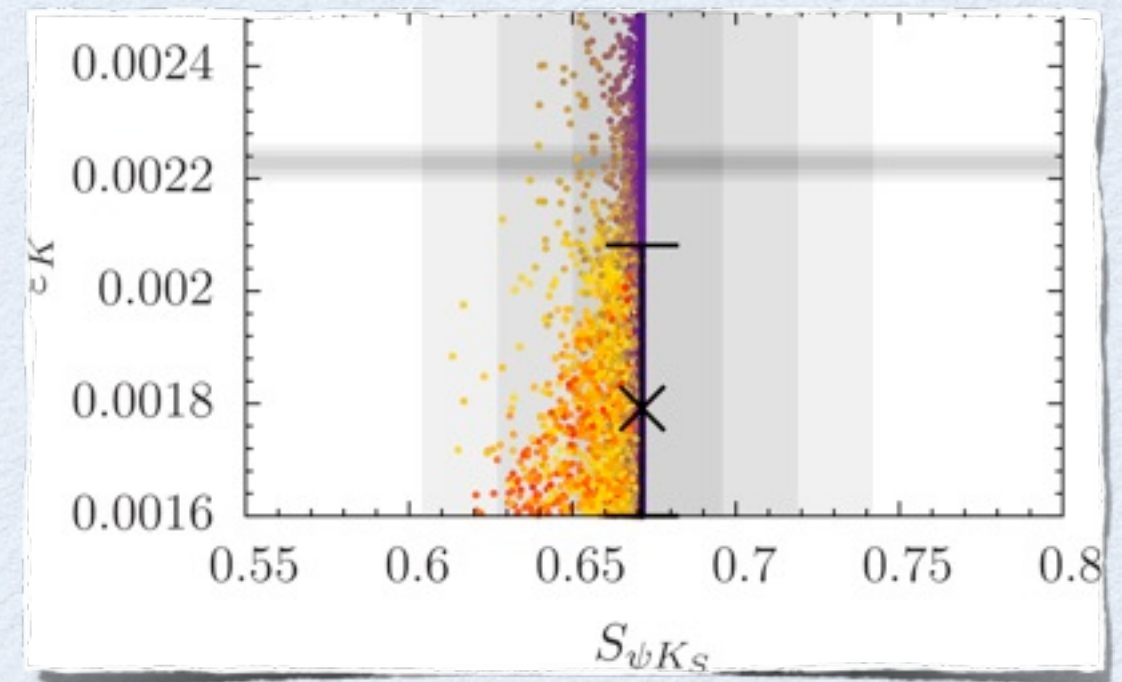
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PHENOMENOLOGICAL ANALYSIS - I

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PHENOMENOLOGICAL ANALYSIS - I

The $S_{\psi K_S} - \epsilon_K$ tension can be solved...

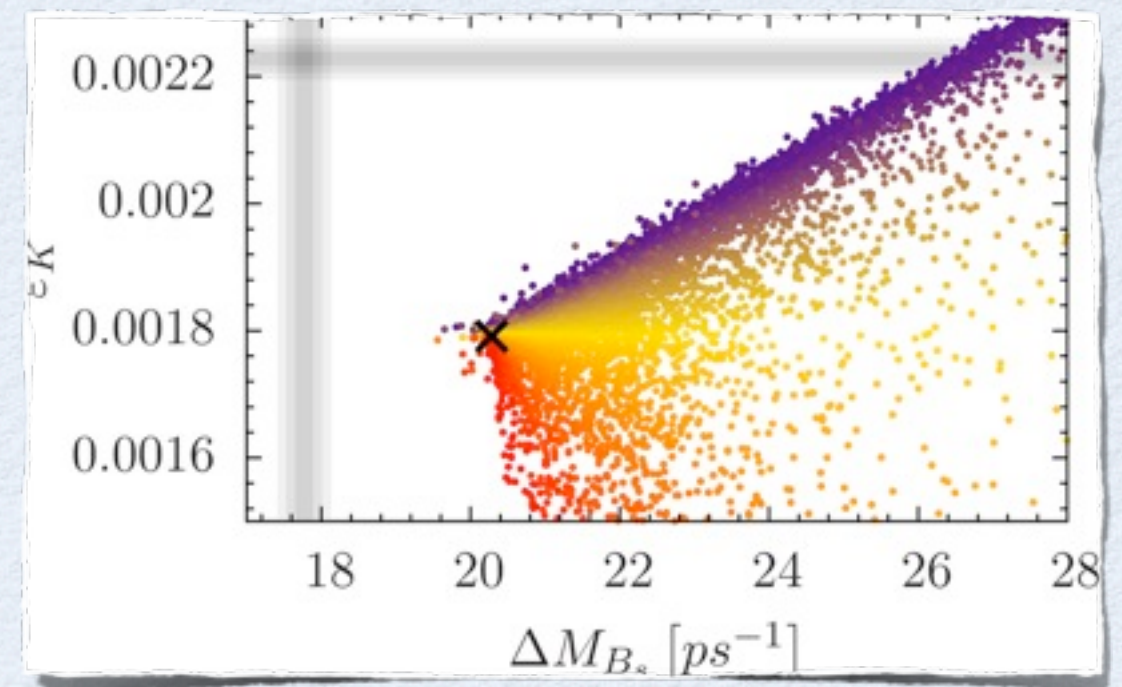
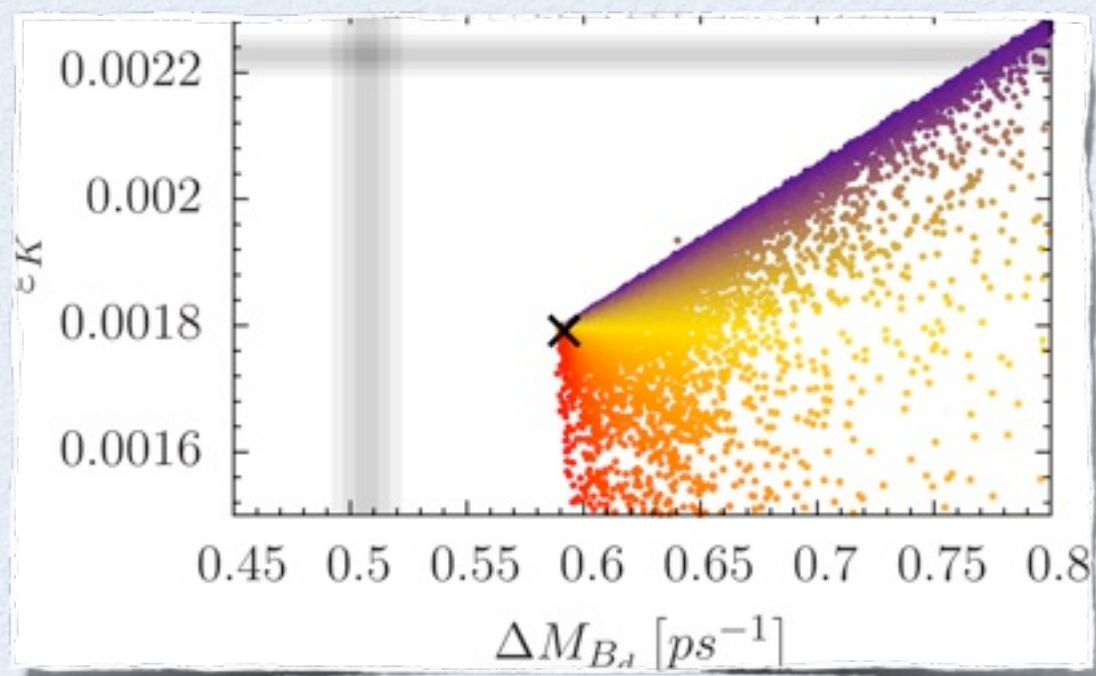
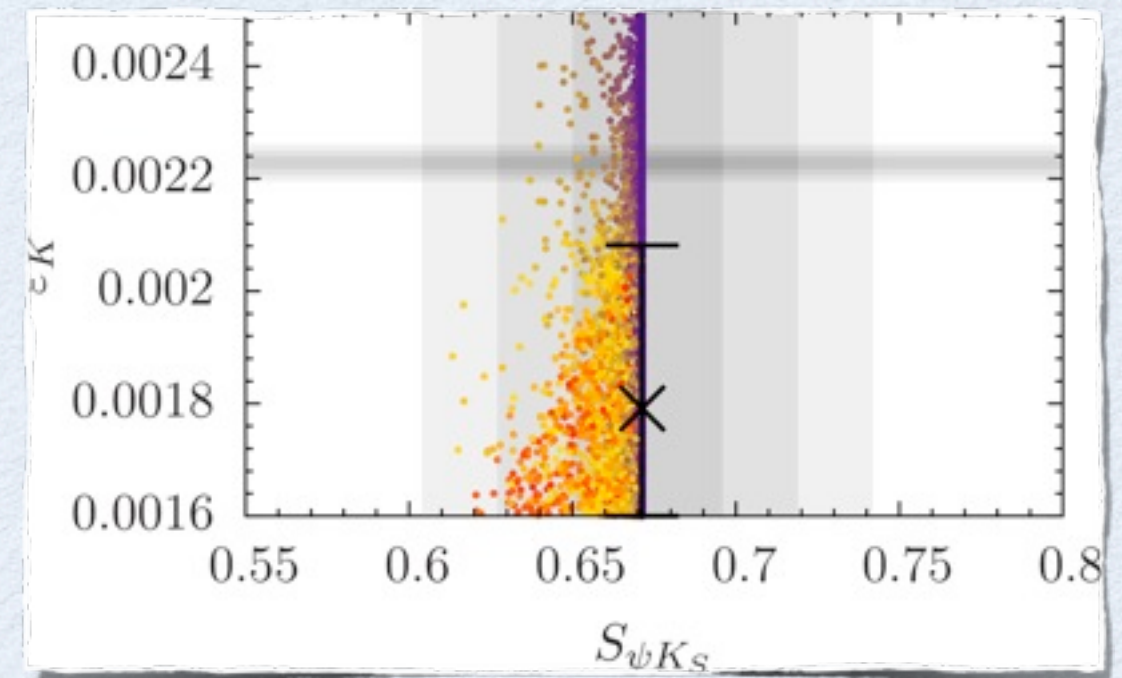


BEYOND MFV: GAUGED FLAVOUR SYMMETRIES

PHENOMENOLOGICAL ANALYSIS - I

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... but this automatically worsens the agreement of ΔM_s and ΔM_d .



BEYOND MFV: GAUGED FLAVOUR SYMMETRIES

PHENOMENOLOGICAL ANALYSIS - II

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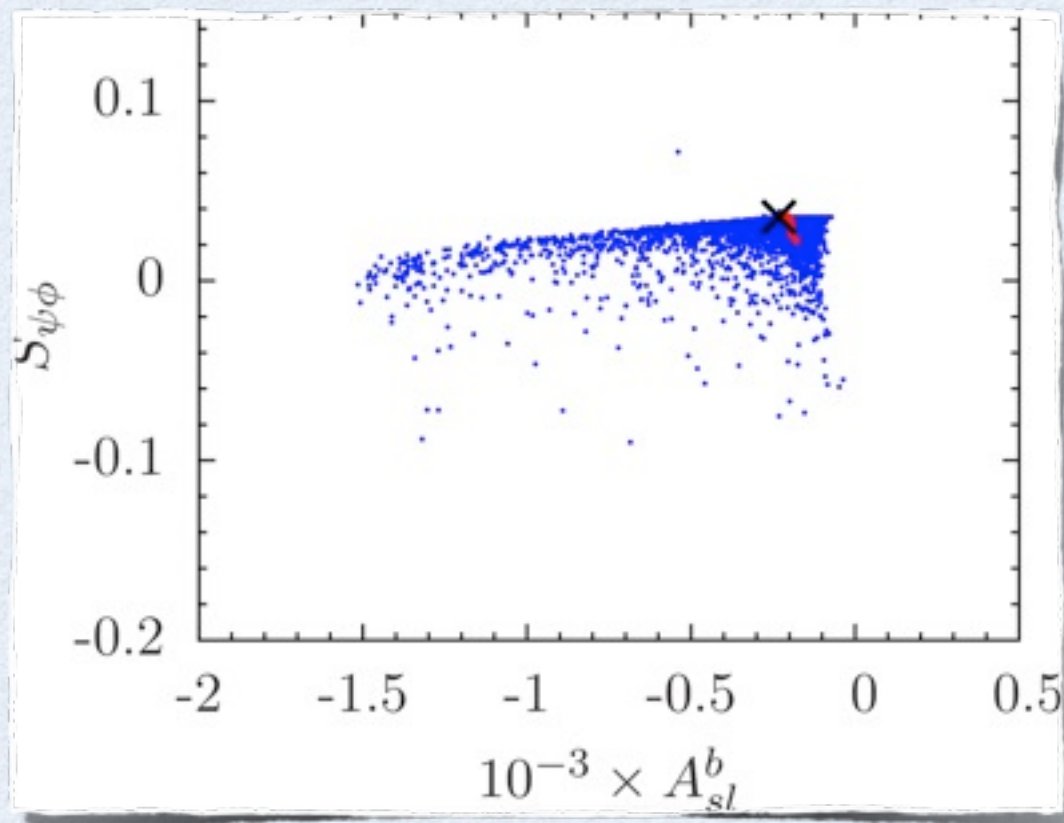
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The discrepancy in the dimuon asymmetry can be **improved**...

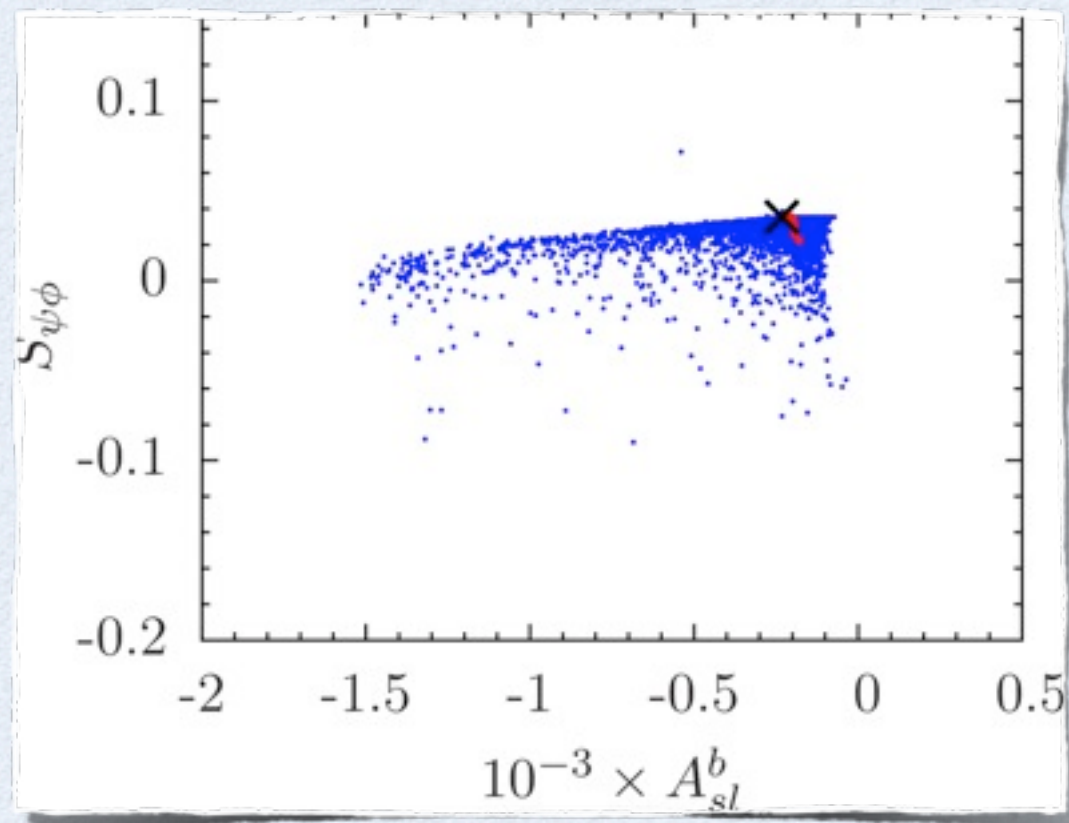


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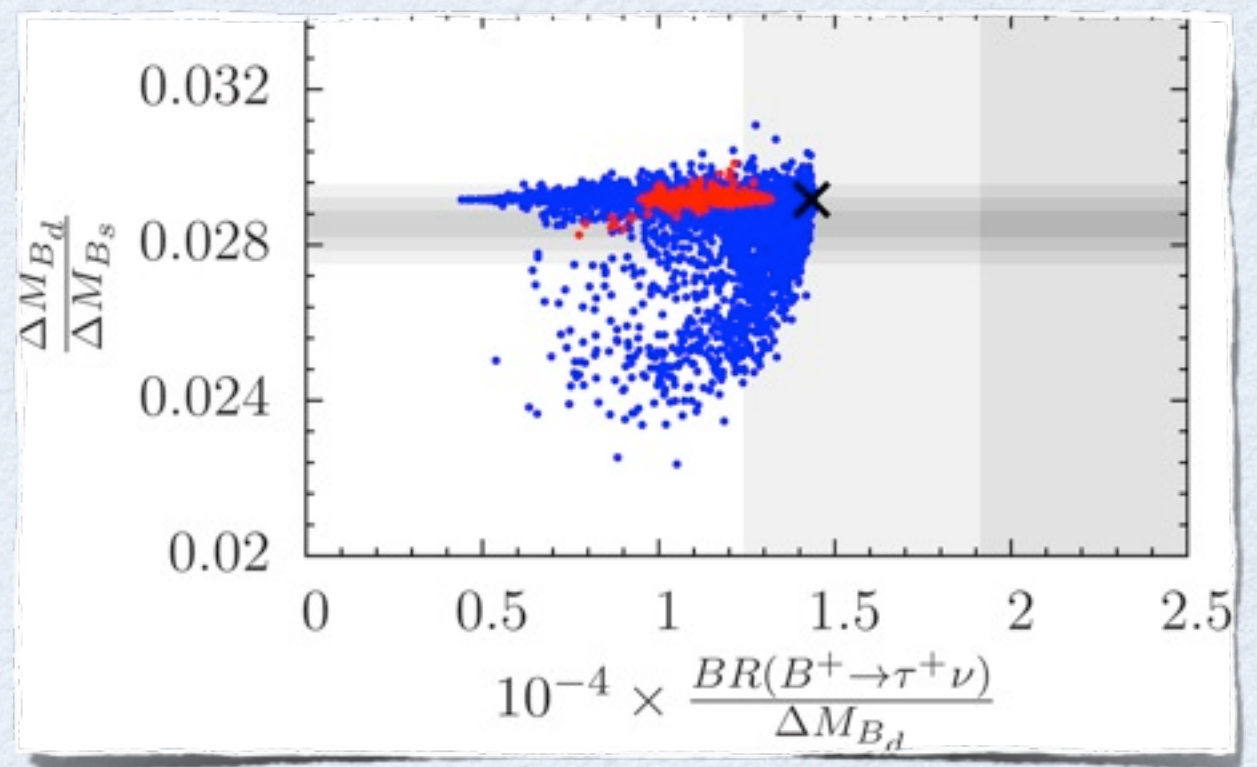
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... but the one in $BR(B \rightarrow \tau \nu)$ is even **worsened**.



CONCLUSIONS AND OUTLOOK

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- Last data from LHCb indicates that the SM has passed also to the tests of $S_{\psi\phi}$ and $B_{d,s} \rightarrow \mu^+\mu^-$. Nevertheless, the 3σ tensions in $\varepsilon_K - S_{\psi K_S}$, V_{ub} , $BR(B \rightarrow \tau \nu)$, A_{SL} , together with a few more $1-2\sigma$ tensions, continue to seriously afflict the SM.

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- More than 20 different BSM scenarios have been studied in Munich during the last decade; this work has produced a database of deviations from SM and correlations for each of them. Hence, we have already the tools to test the models, and we are starting to receive the data to make more conclusive statements on them.

THANKS!