

Isospin Breaking in Hadron Physics



La Sapienza Roma

September 21st 2011

Guido Martinelli

$$= \frac{i(-g^{\mu\nu})}{q^2}$$

$$= \frac{i(-g^{\mu\nu})}{q^2}$$

$$= \frac{-ig^{\mu\nu}}{q^2}$$

$$= \frac{-ig^{\mu\nu}}{q^2}$$

$$\left(\frac{-ig^{\beta\nu}}{q^2}\right)$$

$$(q^2)g^{\mu\nu} + \dots$$

$$\text{ref } (q^2)g^{\mu\nu}$$

MAIANI 70
September 21-22, 2011

Dipartimento di Fisica di Roma, Università degli Studi 'Sapienza'
<http://www.roma1.infn.it/maiani70/>

Speakers

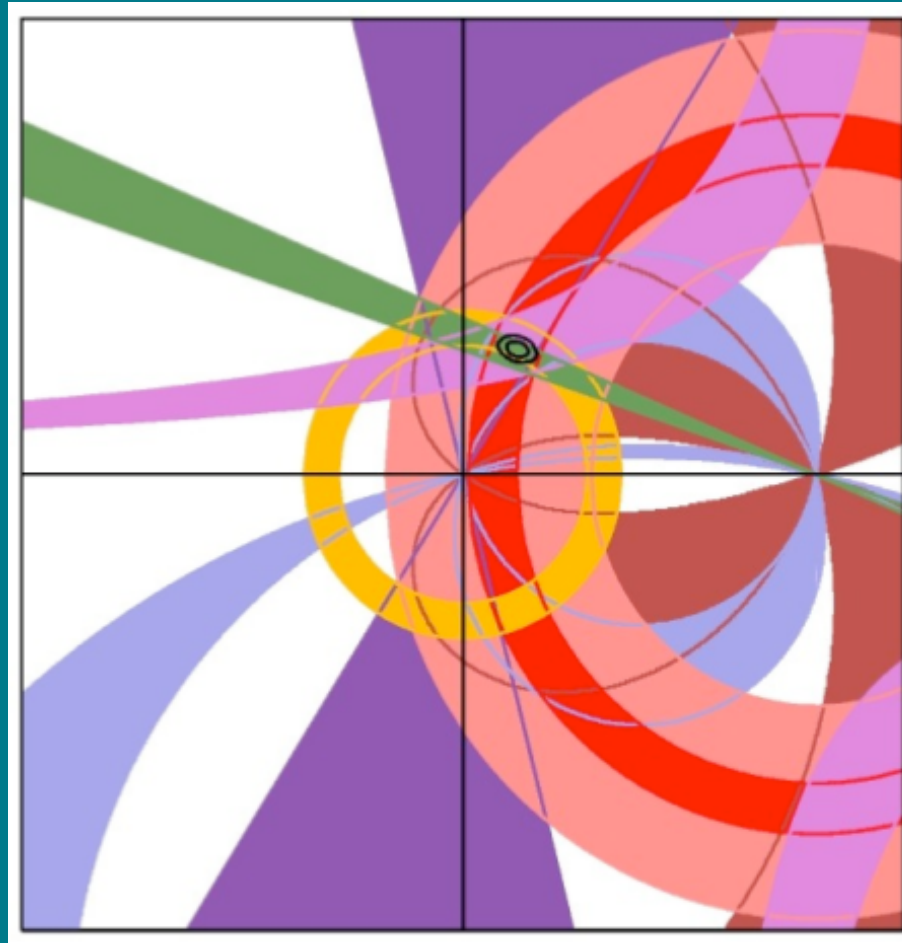
U. Amaldi
G. Altarelli
R. Barbieri
F. Buccella

J. Iliopoulos
G. Isidori
G. Martinelli
A. Masiero
G. Parisi
R. Petronzio



From the Standard Model to Dark Matter and beyond: Symmetries, Masses and Misteries

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NIS, Serbia 29/08/2011



Mixing and CP violation in K- and B-mesons:
a lattice QCD point of view

Maurizio Lusignoli, **Luciano Maiani**, Guido Martinelli, Laura Reina,
Received 11 March 1991; revised 4 July 1991;

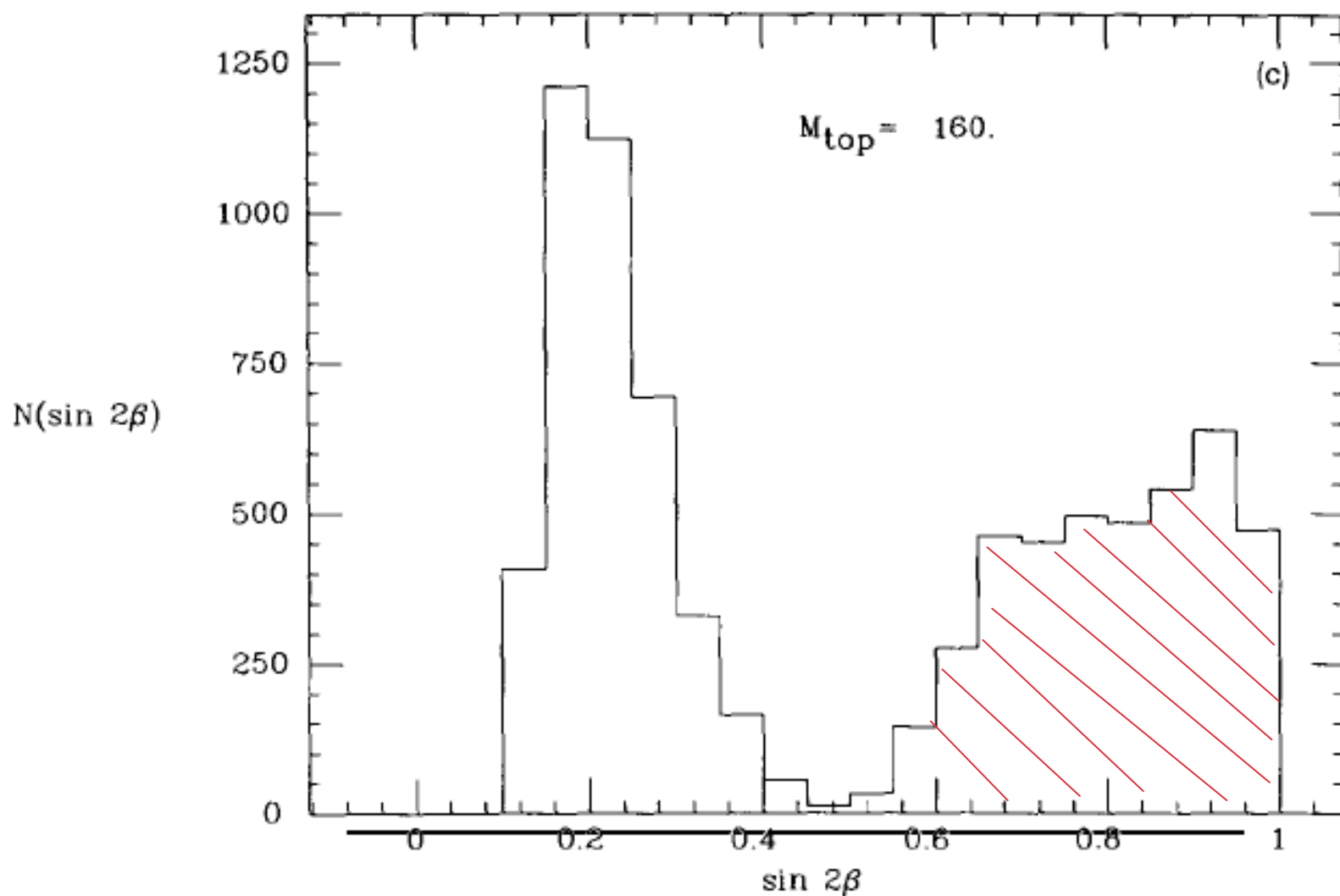
Accepted 5 August 1991. Available online 18 October 2002.

Abstract

We have analyzed $K^0 - \underline{K}^0$ and $B^0 - \underline{B}^0$ mixing on the basis of the most recent experimental results and theoretical lattice determination of the relevant matrix elements. We find that a top mass larger than 140 GeV and $f_B \geq 200$ MeV, as suggested by recent lattice calculations, imply a CP-violating asymmetry in $B \rightarrow J/\Psi + K_s$ decays much larger than previously estimated. We also report the corresponding theoretical prediction for ε'/ε , for which rather large experimental and theoretical uncertainties still exist.

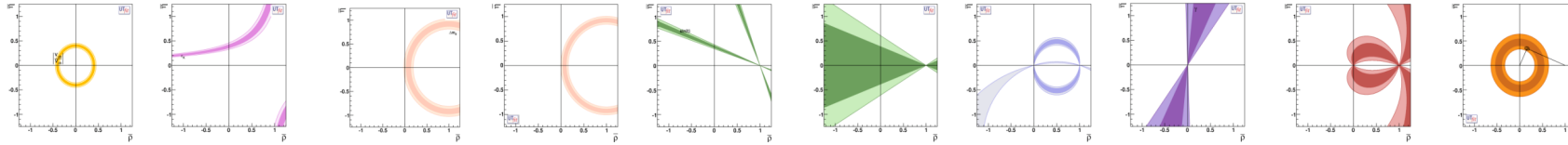
The average values of $\sin(2\beta)$ and $\sin(2\alpha)$ in correspondence to the different m_t and f_B ranges are listed in table 7 in sect. 8. Absolute lower bounds can be read from figs. 3 and 4. For large f_B , eq. (2.3a), we find

$$\sin(2\beta) \geq 0.55(0.45, 0.40), \quad \text{for } m_t = 140(160, 200) \text{ GeV}. \quad (2.14)$$

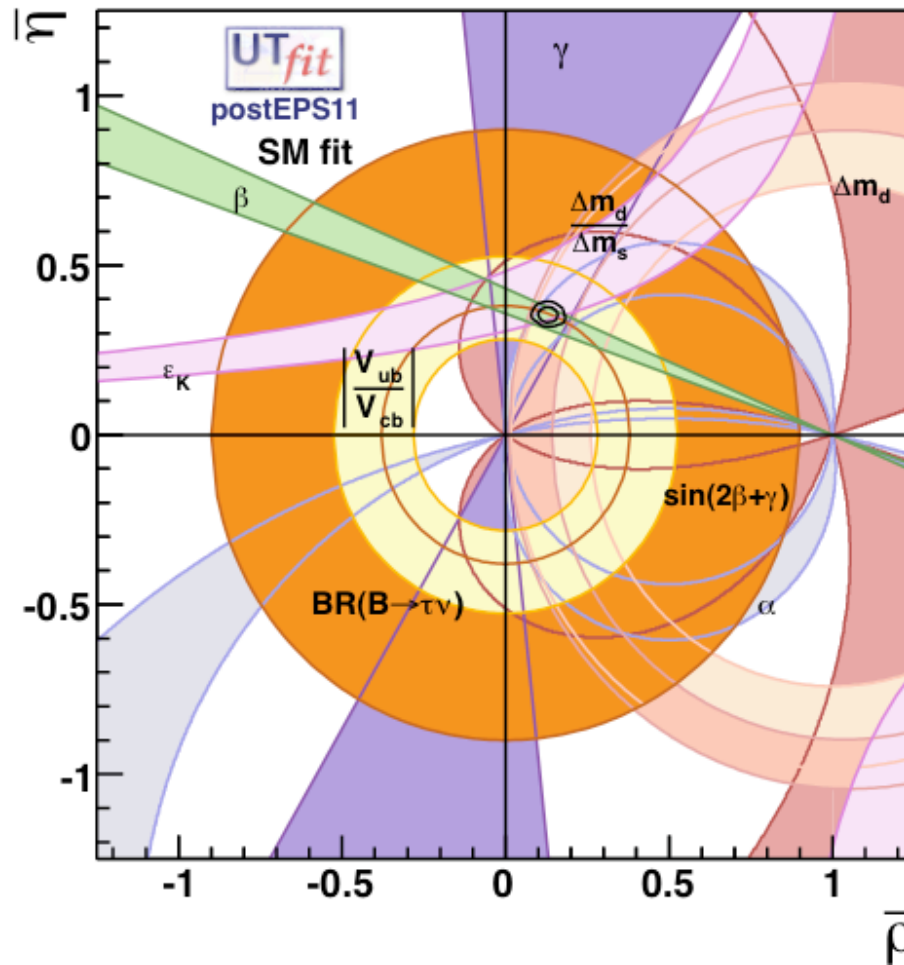


Global Fit within the SM

SM Fit



In the hadronic sector, the SM CKM pattern represents the principal part of the flavour structure and of CP violation



Consistence on an over constrained fit of the CKM parameters

$$\rho = 0.132 \pm 0.020$$

$$\eta = 0.353 \pm 0.014$$

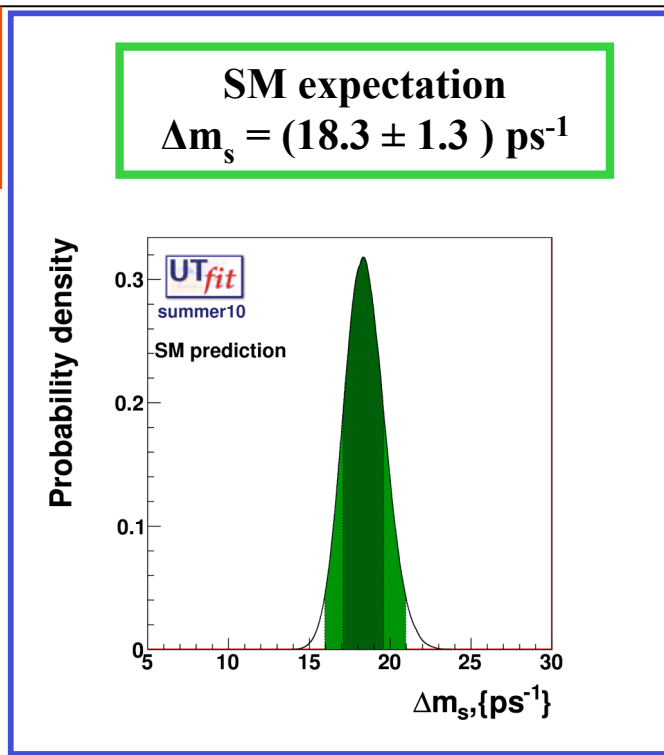
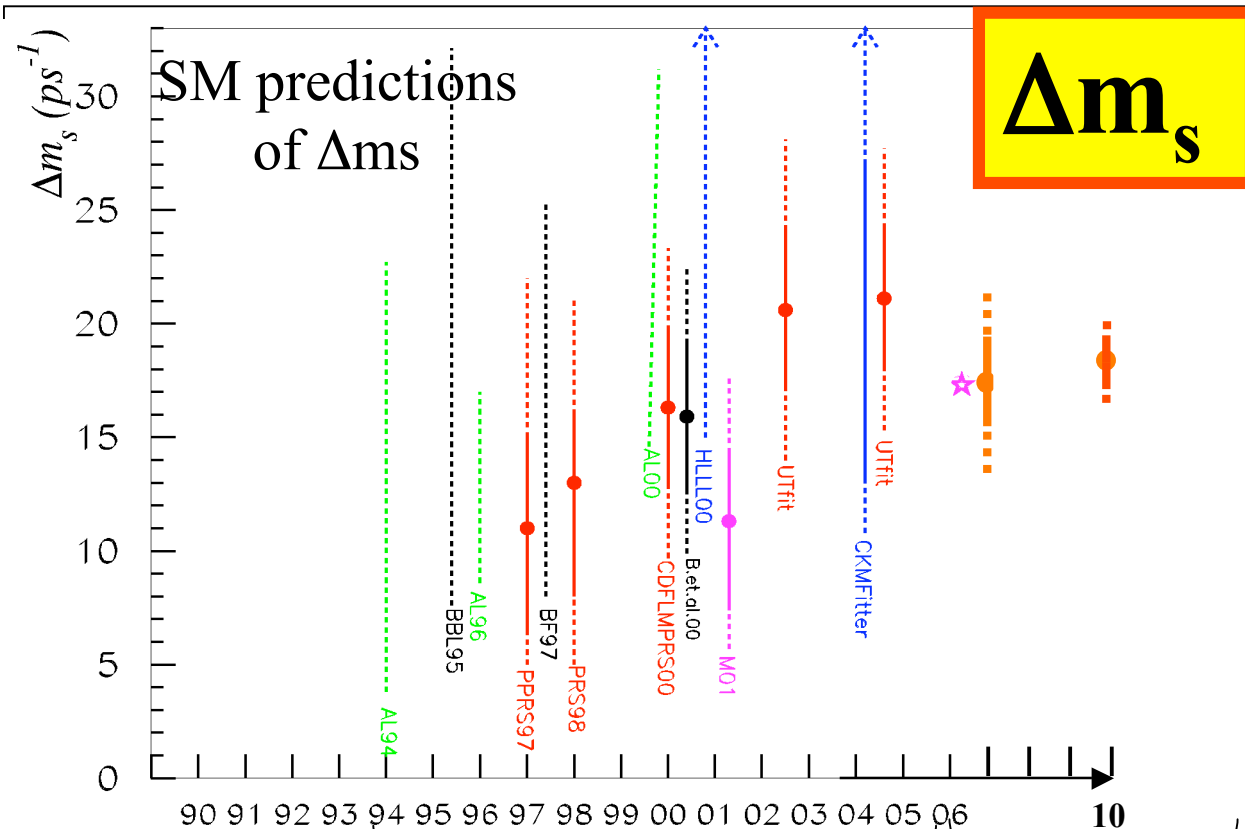
$$\alpha = (88 \pm 3)^\circ$$

$$\sin 2\beta = 0.695 \pm 0.025??$$

$$\beta = (22 \pm 1)^\circ$$

$$\gamma = (69 \pm 3)^\circ$$

CKM matrix is the dominant source of flavour mixing and CP violation



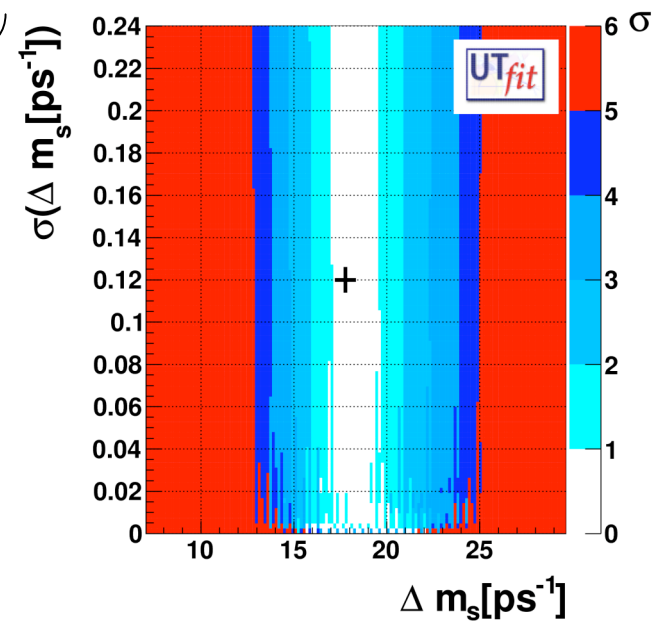
90 91 92 93 94 95 96 97 98 99 00 01 02 03 04 05 06 10

Prediction "era" Monitoring "era"

Exp
 $\Delta m_s = (17.77 \pm 0.12) \text{ ps}^{-1}$

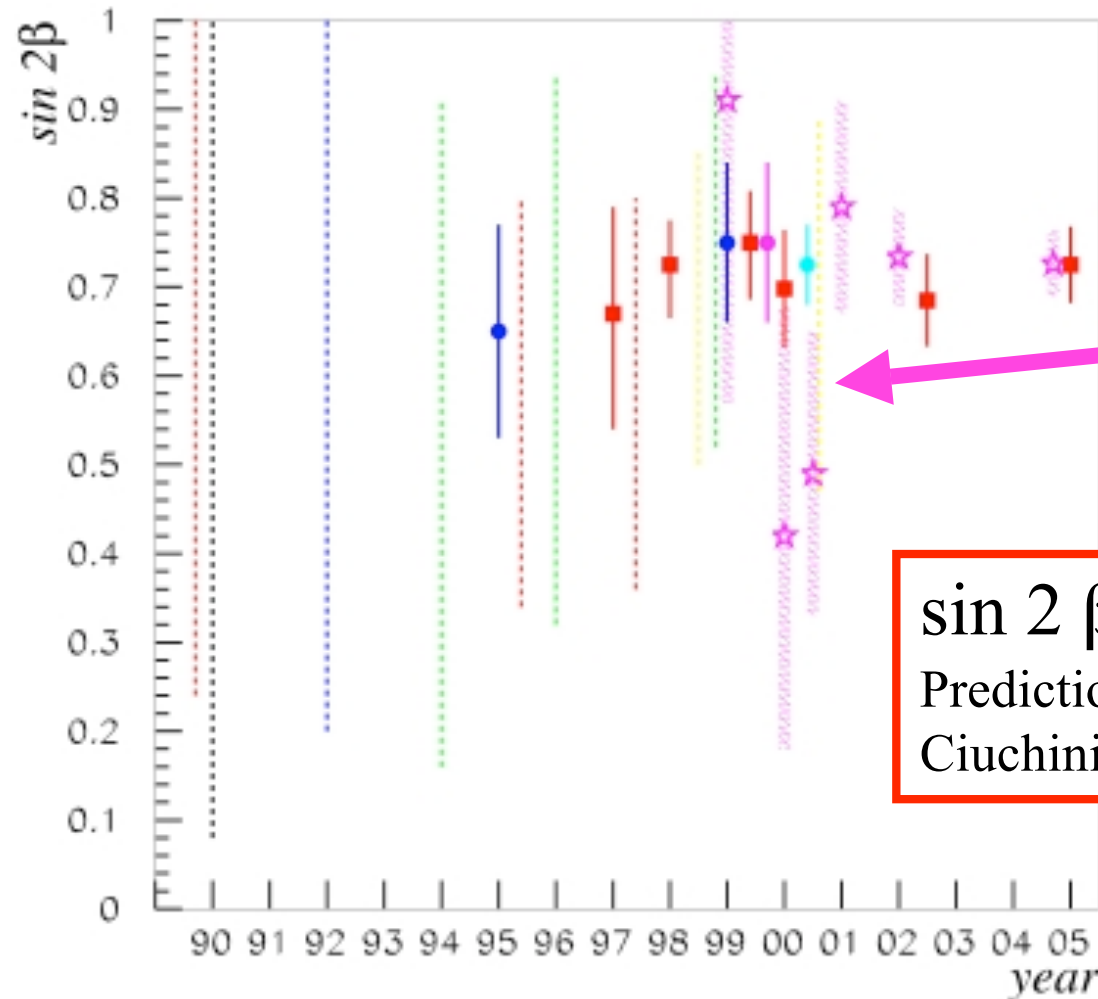
Legenda
 agreement between the predicted values and the measurements at better than :

□ 1σ	■ 3σ	■ 5σ
■ 2σ	■ 4σ	■ 6σ



Theoretical predictions of $\sin 2\beta$ in the years

predictions
exist since '95



experiments

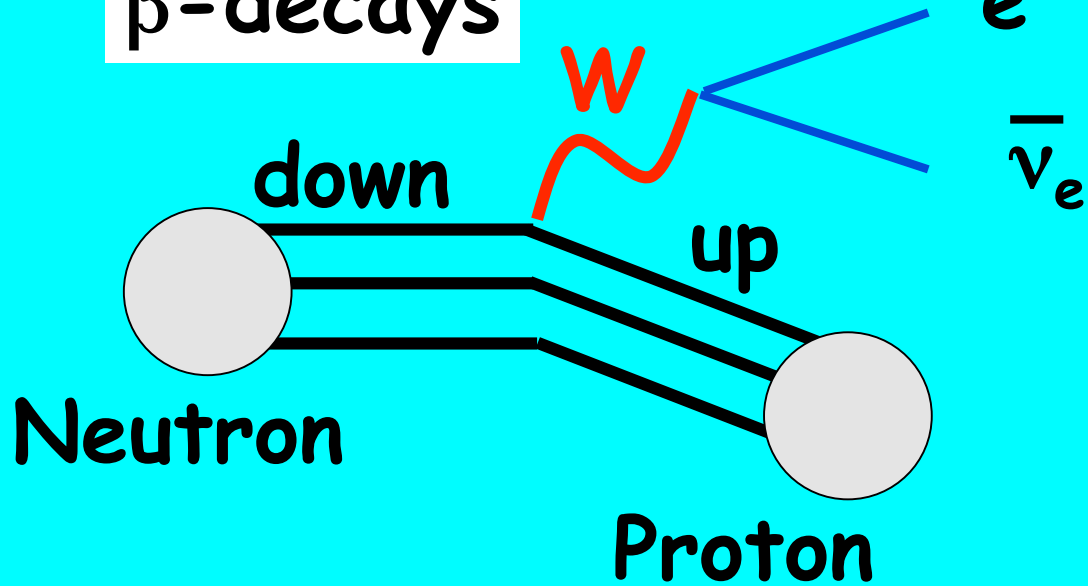
$\sin 2\beta_{\text{UTA}} = 0.65 \pm 0.12$
Prediction 1995 from
Ciuchini, Franco, G.M., Reina, Silvestrini

Quark masses & Generation Mixing

V_{ud}	V_{us}	V_{ub}
V_{cd}	V_{cs}	V_{cb}
V_{td}	V_{ts}	V_{tb}

GIM

β -decays



$|V_{ud}|$

- $|V_{ud}| = 0.9735(8)$
- $|V_{us}| = 0.2196(23)$
- $|V_{cd}| = 0.224(16)$
- $|V_{cs}| = 0.970(9)(70)$
- $|V_{cb}| = 0.0406(8)$
- $|V_{ub}| = 0.00409(25)$
- $|V_{tb}| = 0.99(29)$
(0.999)

Form factor, decay constants and unitarity

- unitarity:

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$$

- experiment:
via non pert. quantities

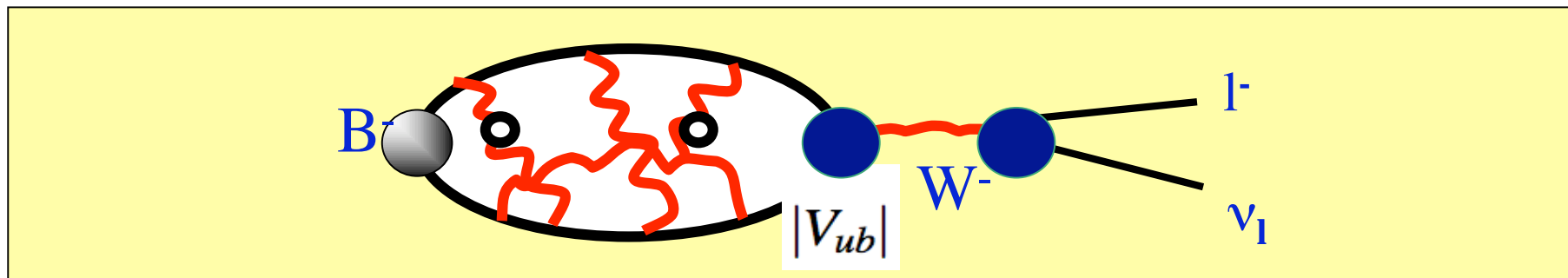
$$|V_{ub}| = 3.93 (36) \cdot 10^{-3}$$

- Kaon decays:

$$|V_{us}| f_+(0) = \boxed{0.2163(5)}$$

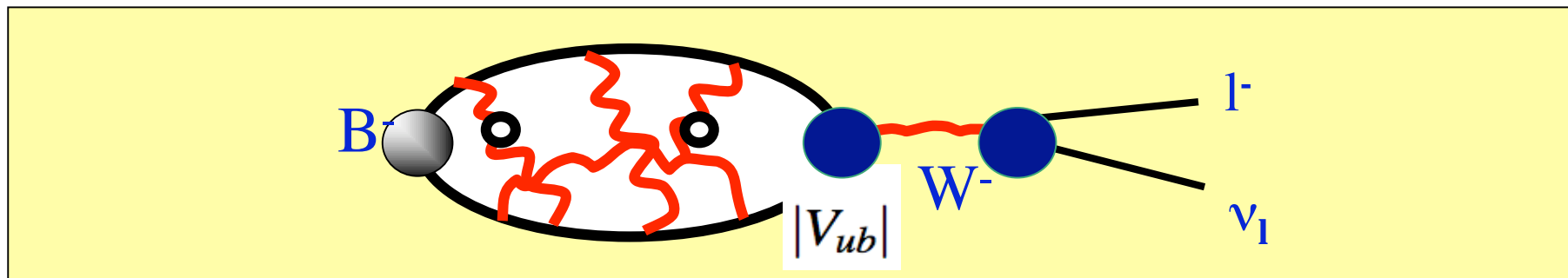
$$\left| \frac{V_{us} f_K}{V_{ud} f_\pi} \right| = \boxed{0.2758(5)}$$

- 3 expressions, 4 unknowns; need one more input



Form factor, decay constants and unitarity

- unitarity: $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$
- experiment:
via non pert. quantities $|V_{ub}| = 3.93 (36) \cdot 10^{-3}$
- Kaon decays: $|V_{us}| f_+(0) = \mathbf{0.2163(5)}$
- $\left| \frac{V_{us} f_K}{V_{ud} f_\pi} \right| = \mathbf{0.2758(5)}$
- 3 expressions, 4 unknowns; need one more input



Form factor, decay constants and unitarity

$$f_+(0) = 0.964 (3) (4) \quad (N_f = 2 + 1)$$

$$f_+(0) = 0.956 (6) (6) \quad (N_f = 2)$$

2010

publication status
 chiral extrapolation
 finite volume errors
 continuum extrapolation

Precision at the per mille level !!

The K → π vector form-factor at zero momentum transfer on the lattice.

D. Becirevic, G. Isidori, V. Lubicz, G. Martinelli,

F. Mescia, S. Simula, C. Tarantino, G. Villadoro Published in Nucl.Phys. B705 (2005) 339-362

$$f_K/f_\pi = 1.190 (2) (10) \quad (N_f = 2 + 1)$$

$$f_K/f_\pi = 1.210 (6) (17) \quad (N_f = 2)$$

Table 1: Colour code for the data on $f_+(0)$.

2010

publication status
 chiral extrapolation
 finite volume errors
 continuum extrapolation

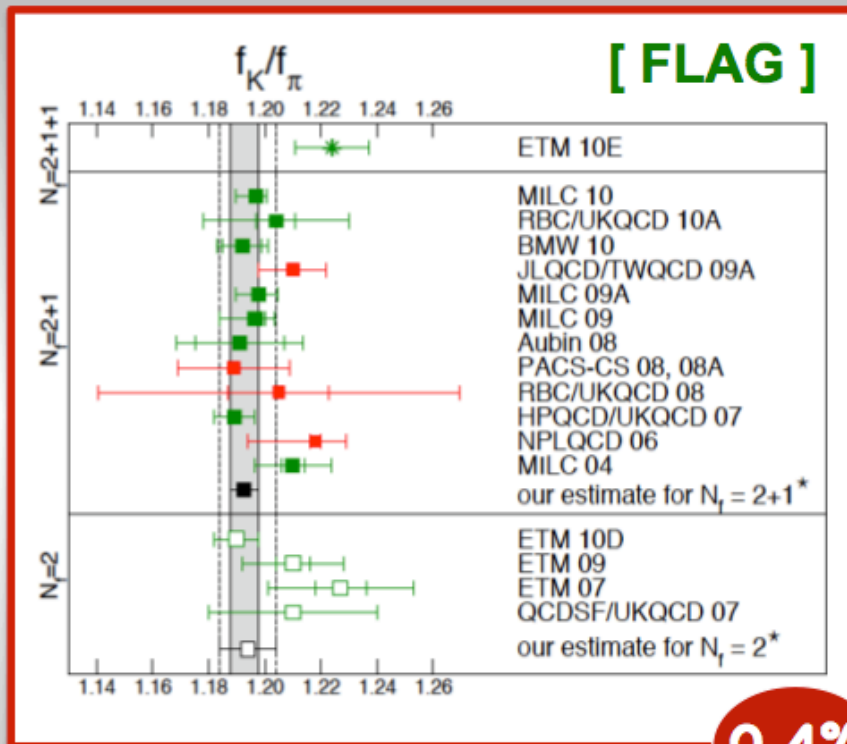
most systematics
OK

Collaboration	N_f	publication status	chiral extrapolation	finite volume errors	continuum extrapolation	f_K/f_π
MILC 09A	2+1	C	★	★	★	1.198(2)($^{+6}_{-8}$)
MILC 09	2+1	P	★	★	★	1.197(3)($^{+6}_{-13}$)
ALVdW 08	2+1	C	★	●	●	1.191(16)(17)
PACS-CS 08, 08B	2+1	A	★	■	■	1.189(20)
BMW 08	2+1	C	★	★	★	1.18(1)(1)
HPQCD/UKQCD 08	2+1	A	★	●	★	1.189(2)(7)
RBC/UKQCD 08	2+1	A	●	★	■	1.205(18)(62)
NPLQCD 06	2+1	A	●	■	■	1.218(2)($^{+11}_{-24}$)
ETM 09	2	A	●	●	★	1.210(6)(15)(9)
QCDSF/UKQCD 07	2	C	●	★	●	1.21(3)

Table 1: Colour code for the data on f_K/f_π .

LATTICE RESULTS FOR f_K/f_π & $f_+(0)$

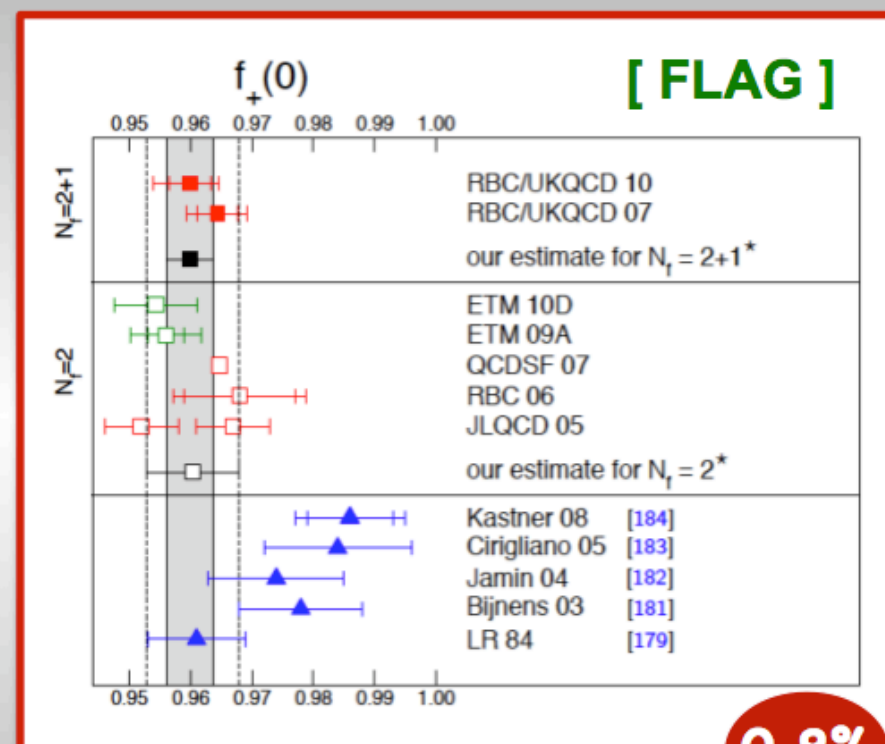
Lubicz LP2011



$$f_K/f_\pi = 1.193(5) \quad (N_f=2+1)$$

$$f_K/f_\pi = 1.210(18) \quad (N_f=2)$$

First result with $N_f=2+1+1$ available

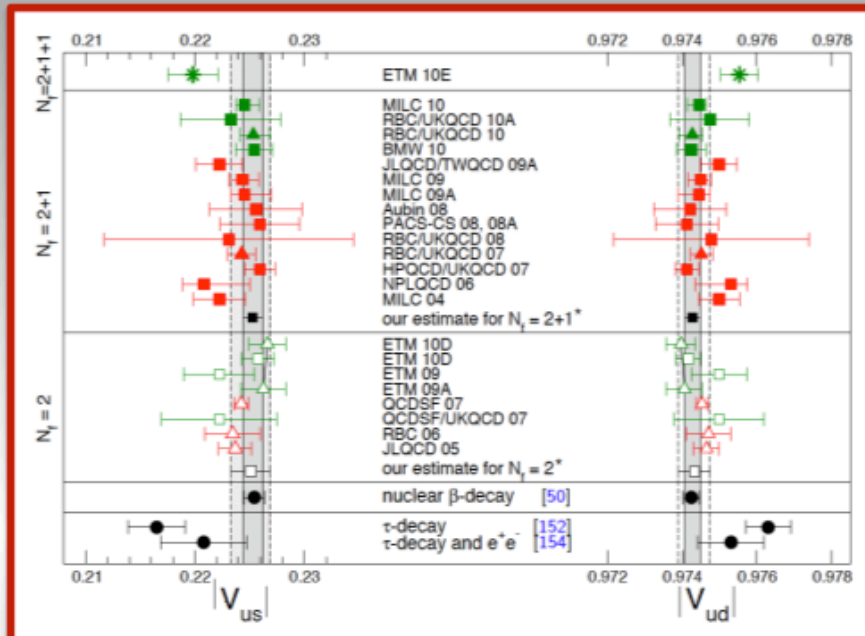


$$f_+(0) = 0.956(8) \quad (N_f=2 \text{ \& } 2+1)$$

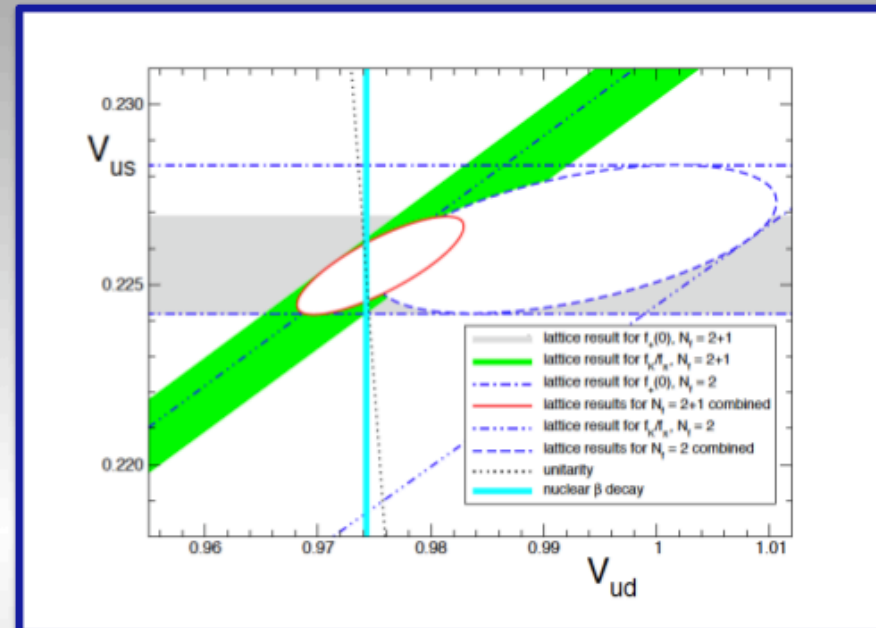
Predictions of analytical model tends to be larger than lattice results¹⁰

First row V_{CKM} Unitarity Test

Lubicz LP2011



Results from f_K/f_π [squares] and $f_+(0)$ [triangles]



The 1st row unitarity plot [FLAG]

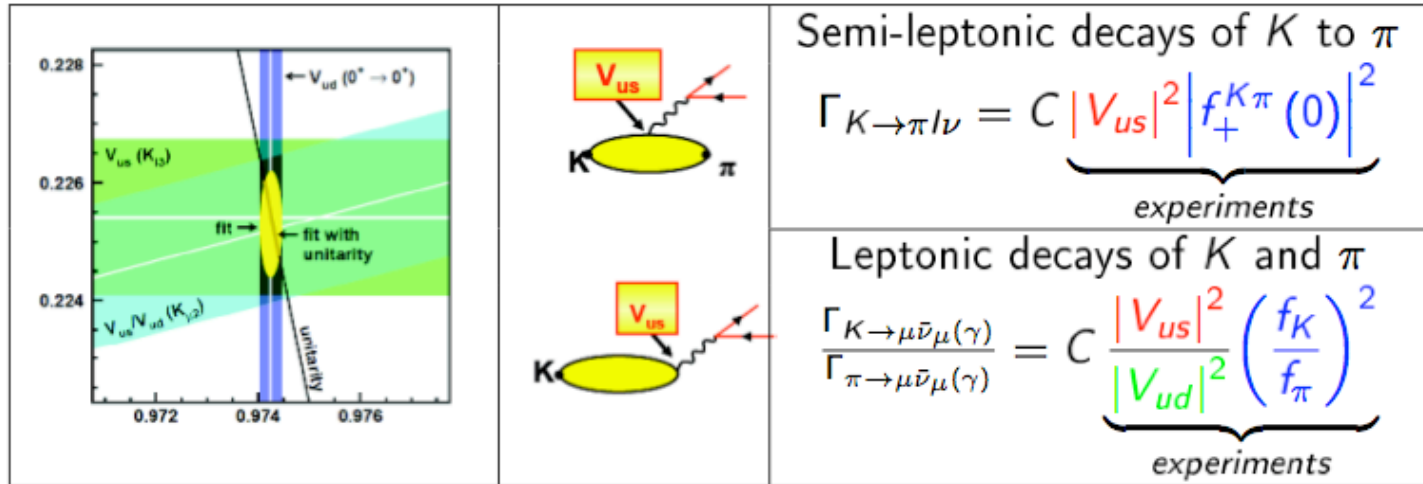
From K12 and
K13 decays

$$\bullet |V_{ud}| = 0.9743(2) \quad \bullet |V_{us}| = 0.2254(9)$$

Combining with $|V_{ud}|$ from nuclear β decays:

$$\Delta_{CKM} = |V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 - 1 = (0 \pm 7) \cdot 10^{-4}$$

Examples: IB effects in Unitary triangle analysis - V_{us}



Lattice: $f_+^{K\pi}(0) = 0.956(8)$
 $f_K/f_\pi = 1.193(5)$ in the isospin symmetric limit.

→ At current precision (0.5–1%), IB corrections **not negligible** ←

Indeed $Ch - PT$ estimates of effects are:

$\left(\frac{f_+^{K^+\pi^0}}{f_+^{K^-\pi^+}} - 1\right)^{QCD} = 2.9(4)\%$ A. Kastner, H. Neufeld (EPJ C57 2008)	$\left(\frac{f_{K^+}/f_{\pi^+}}{f_K/f_\pi} - 1\right)^{QCD} = -0.22(6)\%$ V. Cirigliano, H. Neufeld (arXiv:1102.0563)
--	--

RM123 collaboration:

- F. Sanfilippo: La Sapienza (ROMA 1)
- G.deDivitiis, P.Dimopoulos, R.Frezzotti, Tor Vergata (ROMA 2)
- R.Petronzio, G.Rossi, N.Tantalo:
- V. Lubicz, S. Simula, C. Tarantino: Roma Tre (ROMA 3)
- G. Martinelli: Renegade @ SISSA

Ahhh the
modern PhD
students !!!

Motivations

Sources of Isospin Breaking (in SM)

Electro-magnetic: quarks have different charge $Q_{Up} \neq Q_{Down}$,
 $O(\alpha_{em}) \simeq 1/100$

Strong: quarks have different mass, $m_{Up} \neq m_{Down}$,
 $O((m_{Down} - m_{Up})/\Lambda_{QCD}) \simeq 1/100$

We concentrated on **strong** effects: we left EM for the (near) future.

Current situation

- Lattice typically neglect these effects
- Important corrections at current precision in flavour physics
- Typically estimated with $Ch - PT$ but
 - not always possible (e.g. nucleon mass splitting) or reliable
 - not from first principles

A new strategy: $m_d - m_u$ expansion

Separate mass lagrangian in two contributions:

$$\mathcal{L}_{mass} = \underbrace{\left(\frac{m_d + m_u}{2} \right)}_{m_{ud}} (\bar{u}u + \bar{d}d) - \underbrace{\left(\frac{m_d - m_u}{2} \right)}_{\delta m} (\bar{u}u - \bar{d}d)$$

Split action in two parts:

$$S = S_0 - \delta m \hat{S}, \quad \begin{cases} S_0 & \text{isospin symmetric action} \\ \hat{S} & \text{perturbation} = \sum_x (\bar{u}u - \bar{d}d) \end{cases}$$

Expand functional integral:

$$\langle O \rangle = \frac{\int D\psi O e^{-S_0 + \delta m \hat{S}}}{\int D\psi e^{-S_0 + \delta m \hat{S}}} \stackrel{1st}{\approx} \frac{\int D\psi O e^{-S_0} (1 + \delta m \hat{S})}{\int D\psi e^{-S_0} (1 + \delta m \hat{S})} \approx \frac{\langle O \rangle_0 + \delta m \langle O \hat{S} \rangle_0}{1 + \delta m \langle \hat{S} \rangle_0}$$

Isospin correction determination

One can determine relative correction to an observable O as:

$$\frac{\delta \langle O \rangle}{\langle O \rangle_0} \equiv \frac{\langle O \rangle - \langle O \rangle_0}{\langle O \rangle_0} \simeq \delta m \frac{\langle \hat{S} O \rangle_0}{\langle O \rangle_0}, \quad \hat{S} = \sum_x (\bar{u}u - \bar{d}d)$$

Advantages

- **No need** for new gauge configurations
- Matrix elements are **large** (correction are small due to δm)
- Advantage from **statistical correlations** between $\langle \hat{S} O \rangle$ and $\langle O \rangle$.

Diagrammatically

The diagram shows the expansion of quark propagators. The top row represents the u quark propagator, with a blue arrow labeled 'u' above it. It is equal to a sum of terms: a blue arrow, plus δm times a blue arrow with a circled cross in the middle, plus an ellipsis. The bottom row represents the d quark propagator, with a green arrow labeled 'd' below it. It is equal to a sum of terms: a blue arrow, minus δm times a blue arrow with a circled cross in the middle, plus an ellipsis.

After fixing δm this method can be applied in principle to any observable.

Numerical analysis

Focus

- Applicability of the method
- First observable chosen: f_K/f_π .

Lattice data

- $N_f = 2$ twisted mass configuration provided by ETMC
- m_{ud} , α already determined (arxiv:1010.3659)
- We fix δm from $M_{K_0} - M_{K_+}$
- Electromagnetic corrections are taken from *FLAG*.

Pion effects start at second order in δm (also for π^0)



The diagram shows the expansion of a quark loop with two external lines labeled 'u' and 'd'. The left side is a loop with two vertices. The right side is a sum of diagrams: a bare loop, a loop with a cross on the top vertex, a loop with a cross on the bottom vertex, and an ellipsis, followed by a plus sign and $O(\delta m^2)$.

$$C_{\pi^0\pi^0}(t) - C_{\pi^+\pi^-}(t) = -4 \left[\text{diagram with two crosses} - \text{diagram with two crosses} \right] + \mathcal{O}(\epsilon_{ud}^3)$$

$$C_{nn}^{\pm}(t) = - \left[\text{Diagram 1} \right] + \left[\text{Diagram 2} \right],$$

$$C_{pp}^{\pm}(t) = - \left[\text{Diagram 1} \right] + \left[\text{Diagram 2} \right],$$

$$\left[\text{Diagram 1} \right] - \left[\text{Diagram 2} \right] = 2 \left[\text{Diagram 3} \right] + \mathcal{O}(\varepsilon_{ud}^2),$$

$$\left[\text{Diagram 1} \right] - \left[\text{Diagram 2} \right] = 2 \left[\text{Diagram 3} \right] + \mathcal{O}(\varepsilon_{ud}^2).$$

$$C_{K^0 \pi^-}^{\mu}(t) = - \left[\text{Triangle Diagram} \right] = - \left[\text{Triangle Diagram 1} \right] + \left[\text{Triangle Diagram 2} \right] - \left[\text{Triangle Diagram 3} \right] + \mathcal{O}(\varepsilon_{ud}^2). \quad (20)$$

In order to derive the correction in the $K^+ \rightarrow \pi^0 l^+ \nu$ case, in accordance to our general recipe outlined in the previous section, we start from the correlation function $C_{K^+ \pi^0}^{\mu}(t)$ in the full theory, that has disconnected diagrams, and expand all the light quark propagators in powers of ε_{ud} , namely

$$\begin{aligned} C_{K^+ \pi^0}^{\mu}(t) &= - \left[\text{Triangle Diagram} \right] + \left[\text{Bubble Diagram 1} \right] - \left[\text{Bubble Diagram 2} \right] \\ &= - \left[\text{Triangle Diagram 1} \right] + \left[\text{Bubble Diagram 1} \right] - \left[\text{Bubble Diagram 2} \right] \\ &\quad - \left[\text{Triangle Diagram 2} \right] - \left[\text{Triangle Diagram 3} \right] + \left[\text{Bubble Diagram 3} \right] \\ &\quad + \left[\text{Bubble Diagram 4} \right] - \left[\text{Bubble Diagram 5} \right] + \left[\text{Bubble Diagram 6} \right] \\ &= - \left[\text{Triangle Diagram 1} \right] - \left[\text{Triangle Diagram 2} \right] - \left[\text{Triangle Diagram 3} \right] + 2 \left[\text{Bubble Diagram 3} \right] + \mathcal{O}(\varepsilon_{ud}^2). \end{aligned} \quad (21)$$



Electromagnetic and “strong” QCD IB corrections

$$\begin{aligned}
 C_{K^+K^-}(t) &= - \text{[quark loop]} - \text{[gluon loop]} - e_s^2 \text{[gluon loop]} - e_u^2 \text{[gluon loop]} - e_s e_u \text{[photon loop]} + \mathcal{O}(\alpha_{em} \epsilon_{ud}), \\
 C_{K^0\bar{K}^0}(t) &= - \text{[quark loop]} + \text{[gluon loop]} - e_s^2 \text{[gluon loop]} - e_d^2 \text{[gluon loop]} - e_s e_d \text{[photon loop]} + \mathcal{O}(\alpha_{em} \epsilon_{ud}).
 \end{aligned}$$

The electromagnetic corrections are logarithmically divergent, corresponding to the renormalization of the up and down quark masses (or $\bar{q}q$ or $\bar{q}\tau_3q$)

Using Dashen theorem:

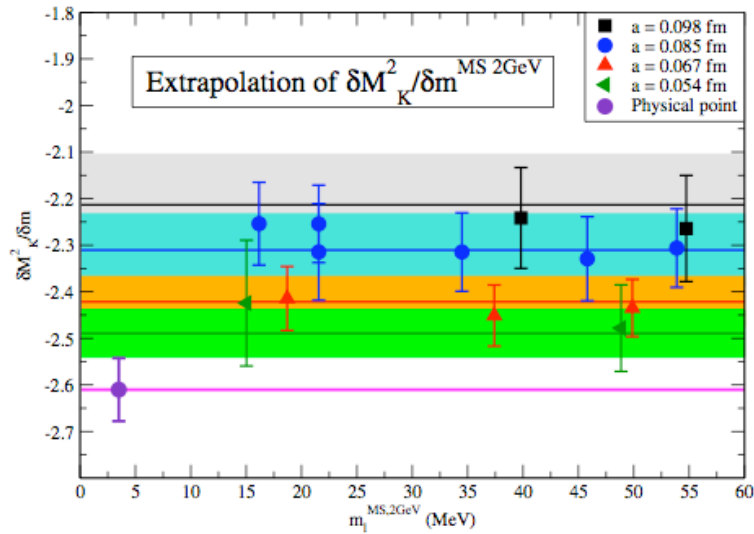
$$[M_{K^0}^2 - M_{K^+}^2]^{QCD} = [M_{K^0}^2 - M_{K^+}^2]^{exp} - (1 + \epsilon_\gamma) [M_{\pi^0}^2 - M_{\pi^+}^2]^{exp}$$

FLAG

Eur. Phys. J. C71 (2011) 1695

$$\epsilon_\gamma = 0.7(5),$$

$$[M_{K^0}^2 - M_{K^+}^2]^{QCD} = 6.05(63) \times 10^3 \text{ MeV}^2.$$



No details on the lattice analysis only a bunch of (preliminary) results

$$[m_d - m_u]^{QCD}(\overline{MS}, 2\text{GeV}) = 2.29(5)(24) \text{ MeV} \times \frac{[M_{K^0}^2 - M_{K^+}^2]^{QCD}}{6.05 \times 10^3 \text{ MeV}},$$

**Cirigliano
Neufeld
-0.0022(6)**

$$\left[\frac{F_{K^+}/F_{\pi^+}}{F_K/F_\pi} - 1 \right]^{QCD} = -0.00376(29)(4) \times \frac{[M_{K^0}^2 - M_{K^+}^2]^{QCD}}{6.05 \times 10^3 \text{ MeV}},$$

$$[M_n - M_p]^{QCD} = 2.8(8)(3) \text{ MeV} \times \frac{[M_{K^0}^2 - M_{K^+}^2]^{QCD}}{6.05 \times 10^3 \text{ MeV}},$$

$$\left[\frac{f_+^{K^0\pi^-}(0) - f_+^{K\pi}(0)}{f_+^{K\pi}(0)} \right]^{QCD} = 1.9(4)(2) \times 10^{-4} \times \frac{[M_{K^0}^2 - M_{K^+}^2]^{QCD}}{6.05 \times 10^3 \text{ MeV}}.$$



WHY RARE DECAYS ?

Rare decays are a manifestation of broken (accidental) symmetries e.g. of physics beyond the Standard Model

Proton decay

**baryon and lepton
number conservation**

$$\mu \rightarrow e + \gamma$$

lepton flavor number

$$\nu_i \rightarrow \nu_k$$

RARE DECAYS WHICH ARE ALLOWED IN THE STANDARD MODEL

FCNC :

$$q_i \rightarrow q_k + \nu \bar{\nu}$$

$$q_i \rightarrow q_k + l^+ l^-$$

$$q_i \rightarrow q_k + \gamma$$

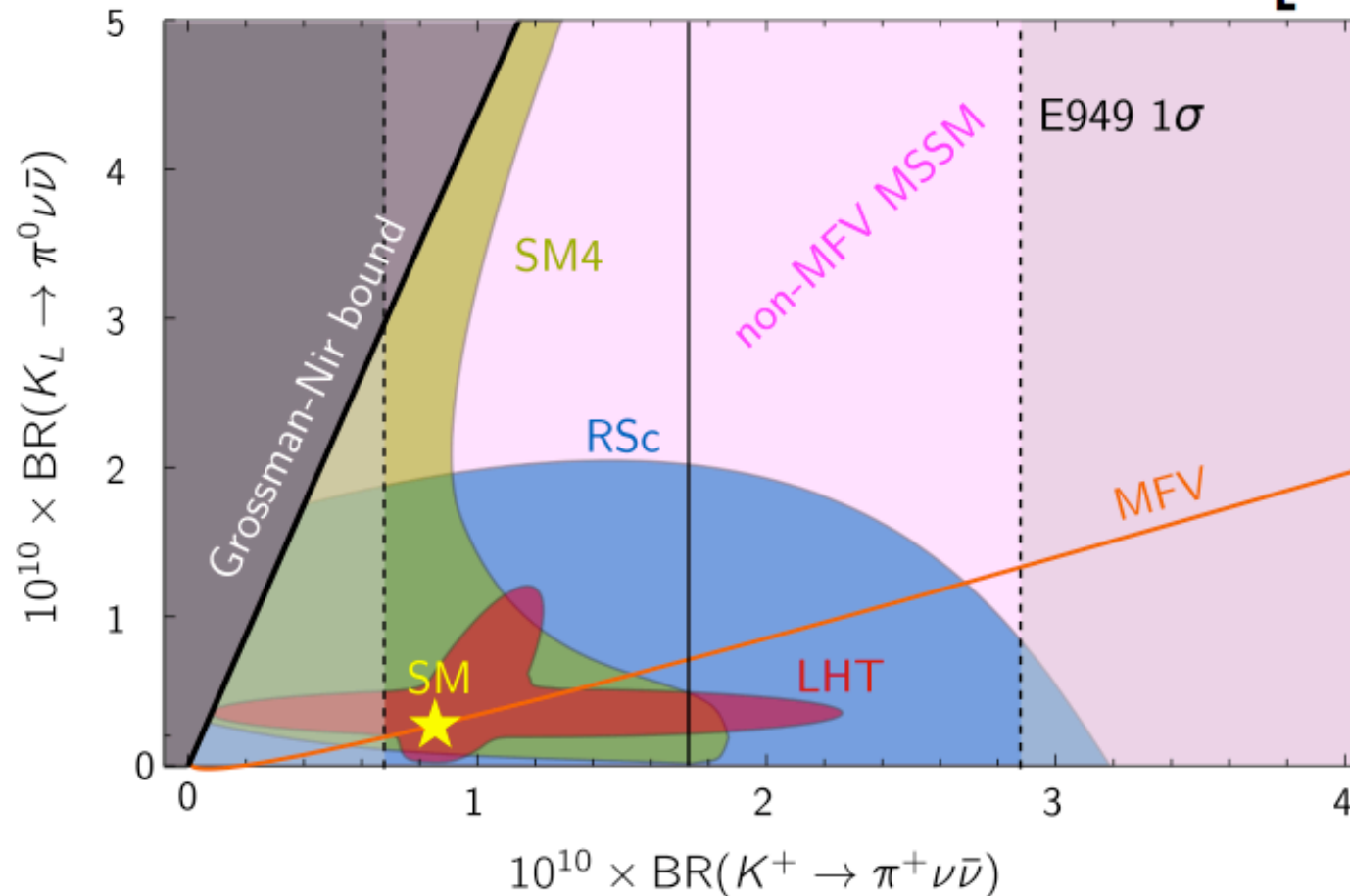
these decays occur only via loops because of *GIM* and are suppressed by *CKM*

THUS THEY ARE SENSITIVE TO
NEW PHYSICS

Rare Decays as Probes of NP models

FCNC in rare K decays

[Straub]



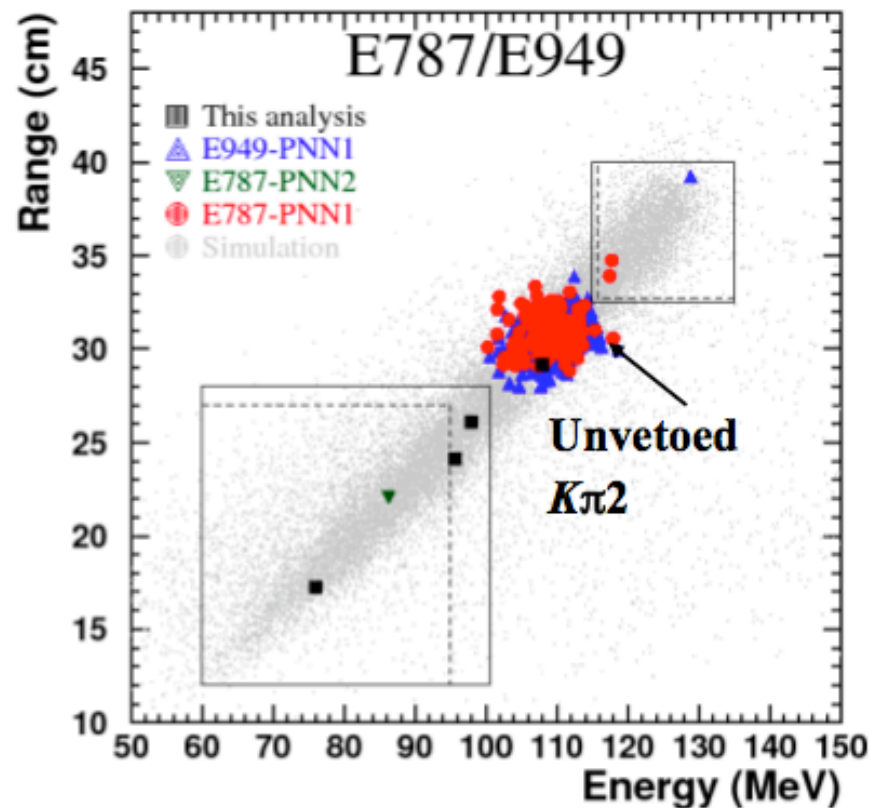
- While only a schematic picture :
 - Correlation between different measurements a powerful probe of NP models
 - Large number of potential channels ... will talk only about a few
 - RD have a bright future: final data sets from B-factories, LHCb, Super Flavour Factories, Kaon experiments...

$K^+ \rightarrow \pi^+ \nu \bar{\nu}$ - Experiment

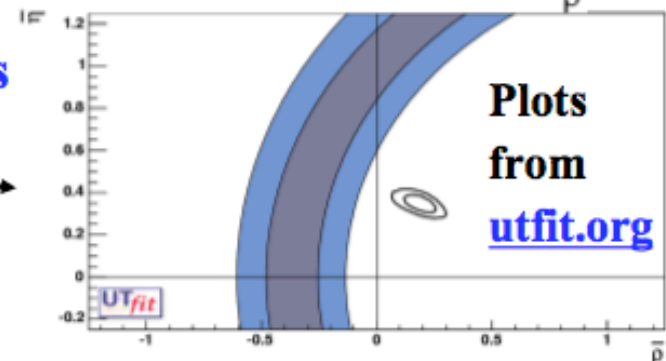
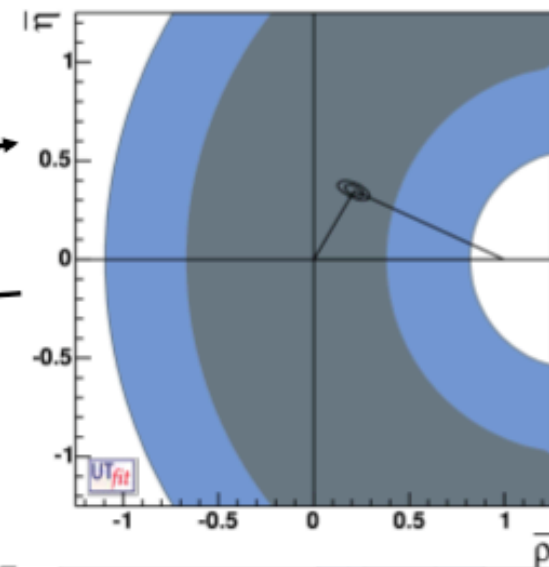
Final results from E787/E949 (2008) [Spadaro]

Combined results, from E787 (1995-8 runs) & E949 (12-weeks run in 2001)

$$BR(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = (1.73^{+1.15}_{-1.05}) \times 10^{-10}$$



Same central value, 100 evts



Prob. all 7 obs. evts are bkg is $\sim 10^{-3}$

Why we like $K \rightarrow \pi \nu \bar{\nu}$?

For the same reason as $A_{J/\psi K_S}$:

1) Dominated by short distance dynamics

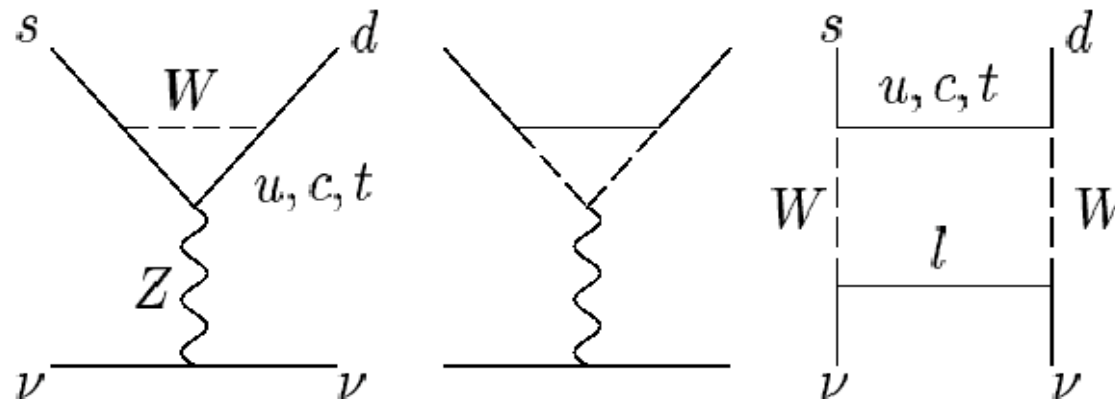
(hard GIM suppression, calculable in pert. theory)

2) Negligible hadronic uncertainties

(matrix element known)

$O(G_F^2)$ Z and W penguin/box $s \rightarrow d \nu \bar{\nu}$ diagrams

**SM
Diagrams**



$$H_{eff} = G_F^2 \alpha / (2\sqrt{2} \pi s_W^2) [V_{td} V_{ts}^* X_t + V_{cd} V_{cs}^* X_c] \times (\bar{s} \gamma_\mu (1 - \gamma_5) d) (\bar{\nu} \gamma^\mu (1 - \gamma_5) \nu)$$

- ☺ NLO QCD corrections to $X_{t,c}$ and $O(G_F^3 m_t^4)$ contributions known
- ☺ The hadronic matrix element $\langle \pi | s \gamma_\mu (1 - \gamma_5) d | K \rangle$ is known with very high accuracy from K13 decays
- ☺ Sensitive to $V_{td} V_{ts}^*$ and expected large \cancel{CP}

CP Violating

$$K_L \rightarrow \pi^0 \nu \bar{\nu}$$

*dominated by the
top quark contribution
→ short distances
(or new physics)*

$$\begin{aligned} & O(\lambda^5 m_t^2) + i O(\lambda m_t^2) \\ & O(\lambda m_c^2) + i O(\lambda^5 m_c^2) \\ & O(\lambda \Lambda_{\text{QCD}}^2) \end{aligned}$$

theoretical error ~ 2 %

$$\begin{aligned} \text{BR}(K_L)_{\text{SM}} &= 4.30 \times 10^{-10} (m_t / 170 \text{ GeV})^{2.3} \times \\ & (\text{Im}(V_{ts}^* V_{td}) / \lambda^5)^2 = (2.8 \pm 1.0) \times 10^{-11} \end{aligned}$$

Using $\Gamma(K_L \rightarrow \pi^0 \nu \bar{\nu}) < \Gamma(K^+ \rightarrow \pi^+ \nu \bar{\nu})$

One gets $\text{BR}(K_L \rightarrow \pi^0 \nu \bar{\nu}) < 1.8 \times 10^{-9}$ (90% C.L.)

2 order of magnitude larger than the SM expectations

$$A(s \rightarrow d \nu \bar{\nu})$$

$$O(\lambda^5 m_t^2) + i O(\lambda^5 m_t^2)$$

$$O(\lambda m_c^2) + i O(\lambda^5 m_c^2)$$

$$O(\lambda \Lambda_{\text{QCD}}^2)$$

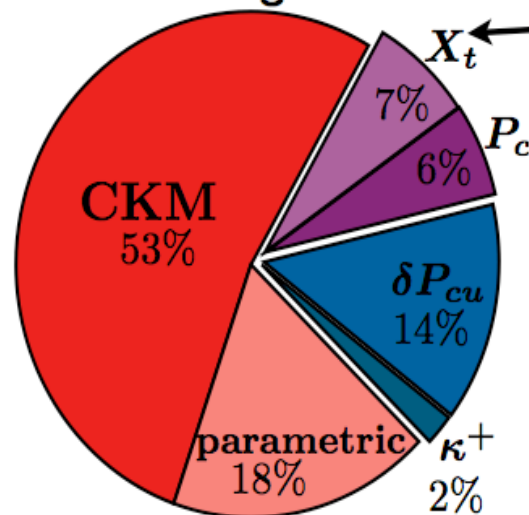
CKM suppressed

GIM

CP conserving: error of $O(10\%)$ due to NNLO corrections in the charm contribution and

Error can be halved by LQCD !!

$K^+ \rightarrow \pi^+ \nu \bar{\nu}$ - error budget



Improved theory prediction

$$Br^{exp} = 1.73_{-1.05}^{+1.15} \times 10^{-10}$$

$$Br^{the} = 8.22_{-0.65}^{+0.74} \pm 0.29 \times 10^{-10}$$

[22] L. Maiani, G. Martinelli, M. L. Paciello, B. Taglienti, Nucl. Phys. **B293** (1987) 420.

Citation summary results

	All papers	Published only
Total number of citable papers analyzed:	<u>20</u>	<u>16</u>
Total number of citations:	2,177	2,169
Average citations per paper:	108.8	135.6
Breakdown of papers by citations:		
Renowned papers (500+)	<u>1</u>	<u>1</u>
Famous papers (250-499)	<u>1</u>	<u>1</u>
Very well-known papers (100-249)	<u>4</u>	<u>4</u>
Well-known papers (50-99)	<u>4</u>	<u>4</u>
Known papers (10-49)	<u>5</u>	<u>5</u>
Less known papers (1-9)	<u>4</u>	<u>1</u>
Unknown papers (0)	<u>1</u>	<u>0</u>
Additional Citation Metrics ?		
h-index ?	15	15

1. Leptonic Decay of Heavy Flavors: A Theoretical Update.

⁽⁷⁹⁹⁾ [Guido Altarelli](#), [N. Cabibbo](#), [G. Corbo](#), [L. Maiani](#) (Rome U. & INFN, Rome), [G. Martinelli](#) (Frascati). ROME-302-1982. Jun 1982. 30 pp.
Published in **Nucl.Phys. B208 (1982) 365-380**

1. Chiral Symmetry on the Lattice with Wilson Fermions.

⁽³⁷³⁾ [Marco Bochicchio](#), [Luciano Maiani](#) (Rome U. & INFN, Rome), [Guido Martinelli](#) (Frascati), [Gian Carlo Rossi](#) (Rome U., Tor Vergata & INFN, F), [Massimo Testa](#) (INFN, Rome). ROME-452-1985. Mar 1985. 45 pp.
Published in **Nucl.Phys. B262 (1985) 331**

1. Current Algebra and Quark Masses from a Monte Carlo Simulation with Wilson Fermions.

⁽¹⁵⁶⁾ [L. Maiani](#), [G. Martinelli](#) (CERN). CERN-TH-4467/86. Jun 1986. 17 pp.
Published in **Phys.Lett. B178 (1986) 265**

2. A lattice computation of the decay constant of the B meson.

⁽¹⁴⁷⁾ [C.R. Allton](#), [Christopher T. Sachrajda](#) (Southampton U.), [V. Lubicz](#), [L. Maiani](#), [G. Martinelli](#) (Rome U. & INFN, Rome). SHEP-89/90-11. Jun 1990. 15 pp.

3. Heavy Flavor Weak Transitions on the Lattice.

⁽¹²⁶⁾ [M.B. Gavela](#) (Madrid, Autonoma U.), [L. Maiani](#), [S. Petrarca](#) (Rome U. & INFN, Rome), [G. Martinelli](#) (CERN), [O. Pene](#) (Orsay, LPT). CERN-TH-4905/87. Dec 1987. 15 pp.
Published in **Phys.Lett. B206 (1988) 113**

4. The Kaon B Parameter and K- pi and K- pi pi Transition Amplitudes on the Lattice.

⁽¹²³⁾ [M.B. Gavela](#) (Madrid, Autonoma U.), [L. Maiani](#), [S. Petrarca](#), [F. Rapuano](#) (Rome U. & INFN, Rome), [G. Martinelli](#) (CERN), [O. Pene](#) (Orsay, LPT), [Christopher T. Sachrajda](#) (Southampton U.). CERN-TH-4905/87. Nov 1987. 26 pp.
Published in **Nucl.Phys. B306 (1988) 677**

Blois '87 (1987)



*Luciano writing a paper
with Massimo Giancarlo Marco and Guido M.*

&

The sign of ε (or the missed dinner !!)

SUMMARY

WORK IN PROGRESS