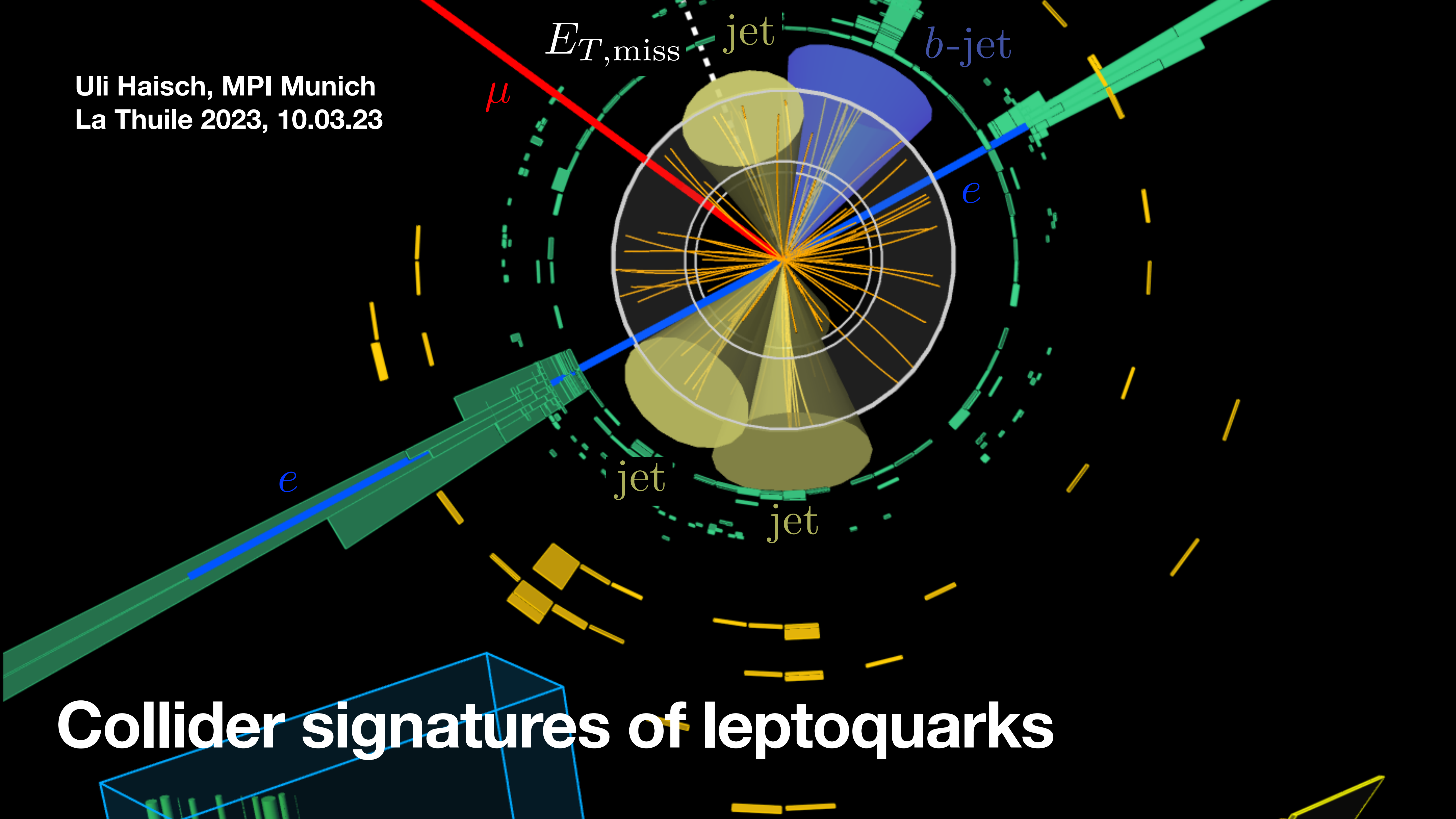
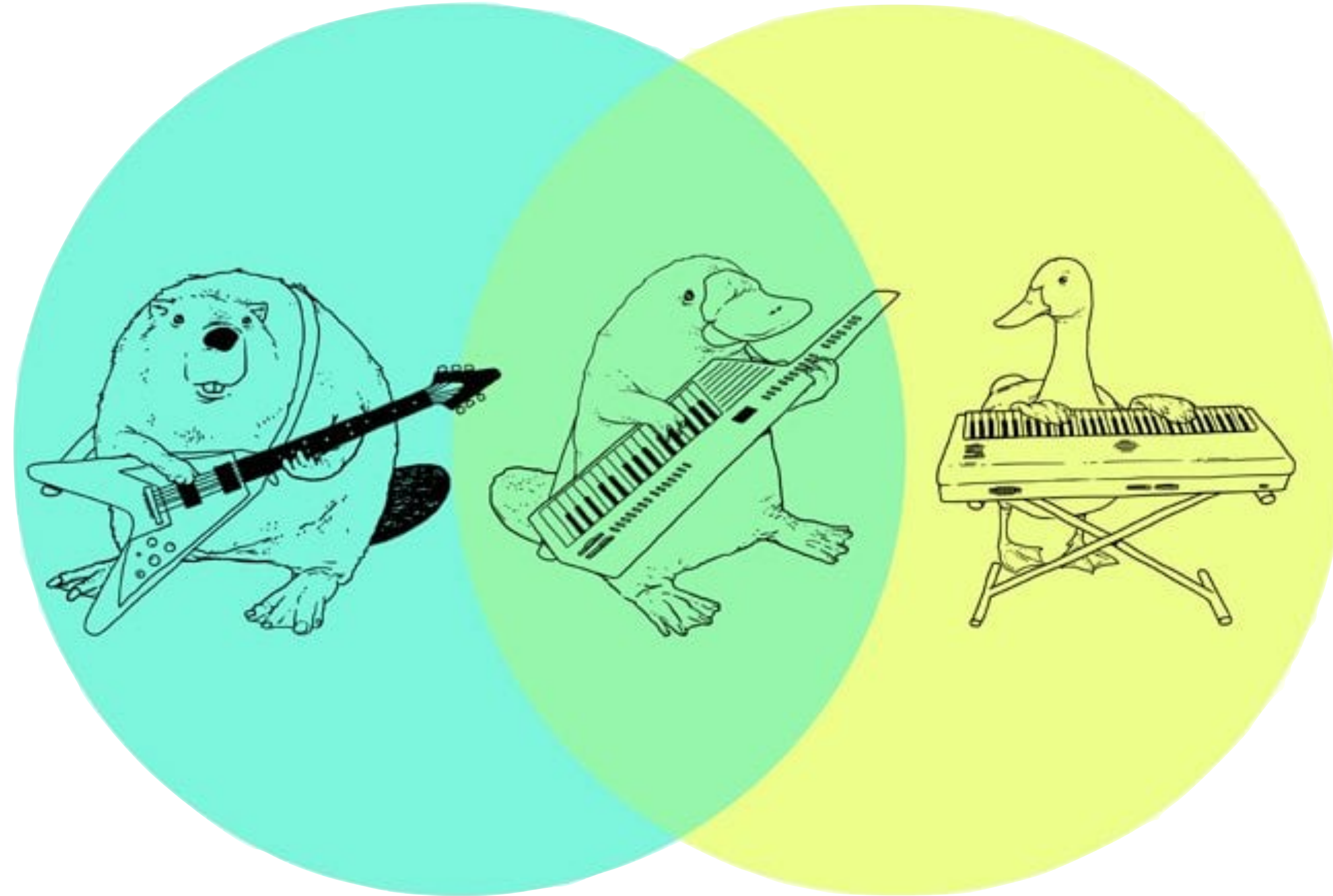


Uli Haisch, MPI Munich
La Thuile 2023, 10.03.23



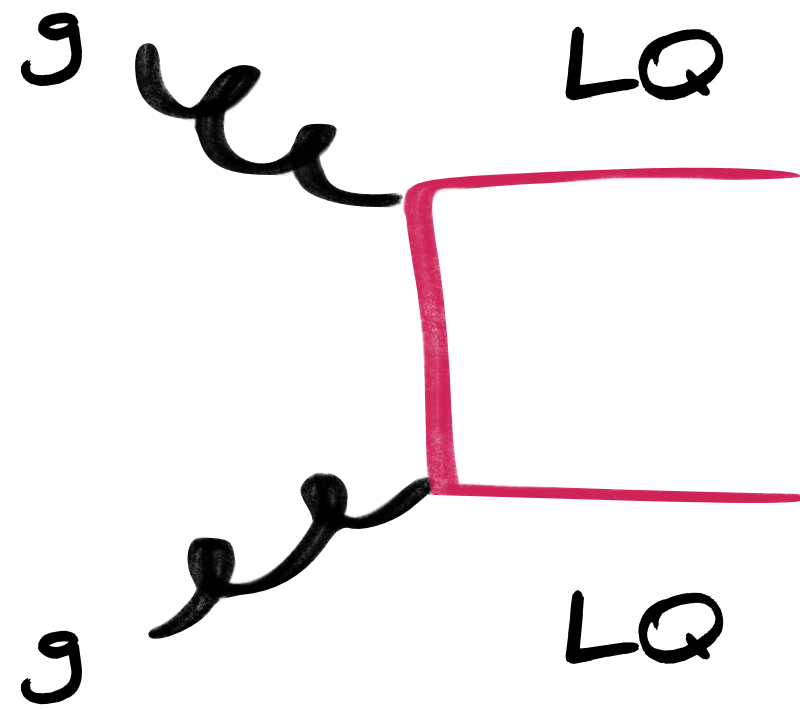
Collider signatures of leptoquarks

What is a leptoquark?

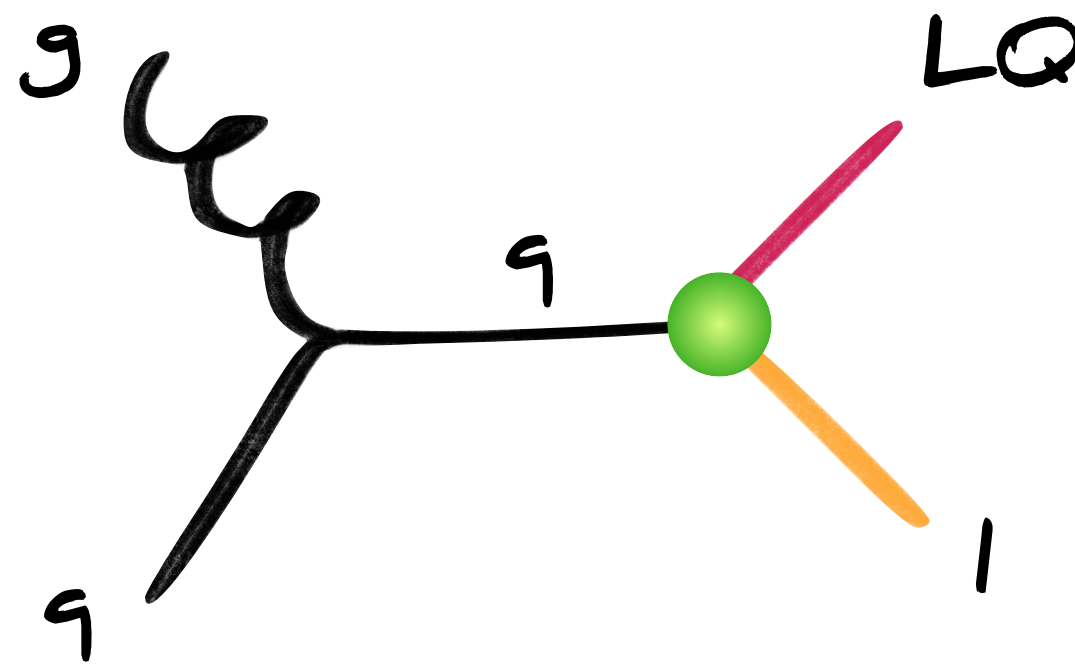


Like a cross between a beaver & a duck is a platypus or a cross between a keyboard & a guitar is a keytar, a cross between a lepton & a quark is a leptoquark (LQ)

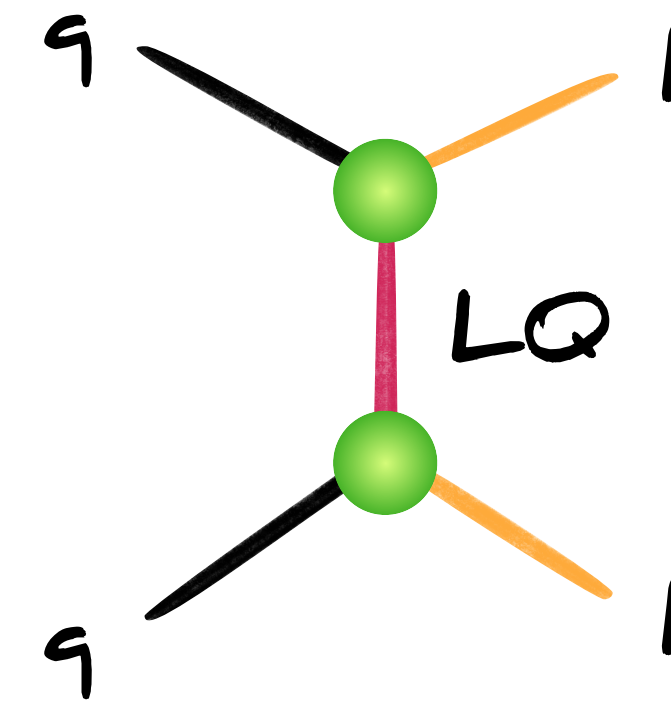
Existing LQ search strategies @ the LHC



pair production (PP)



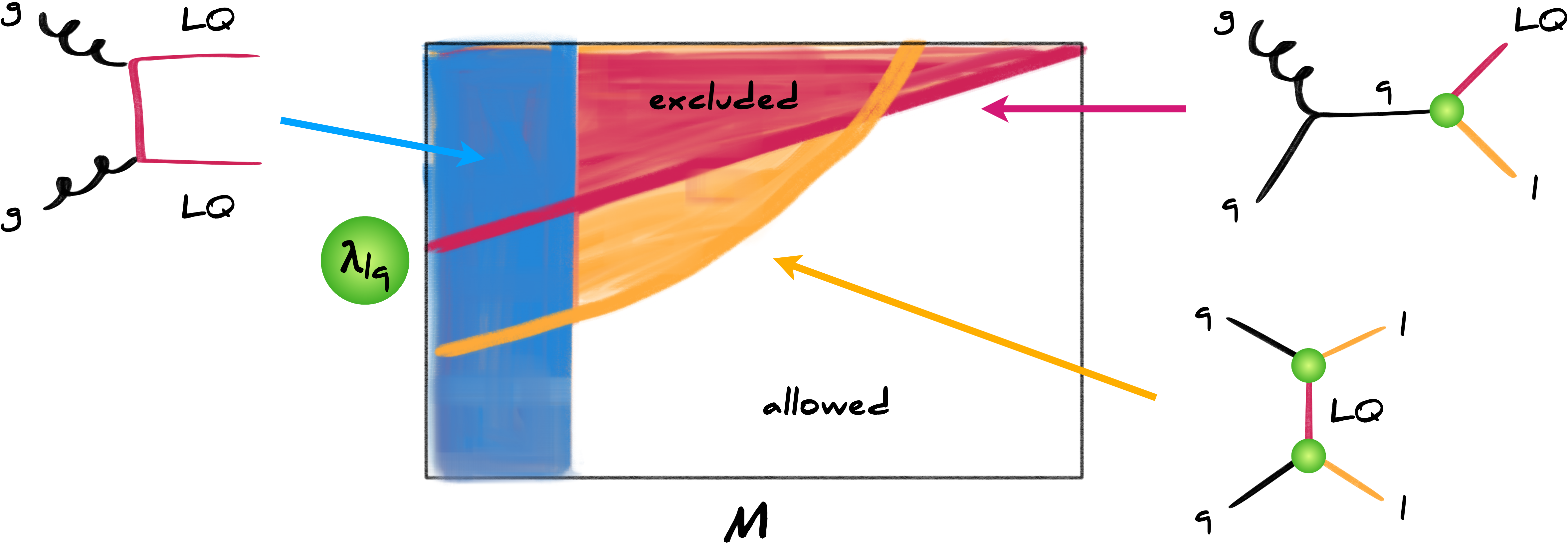
single production (SP)



t-channel Drell-Yan (DY)

Since they are coloured, TeV-scale LQs can be produced in pairs @ LHC with fb rates. Also single & t-channel DY production possible but only recently considered by ATLAS & CMS

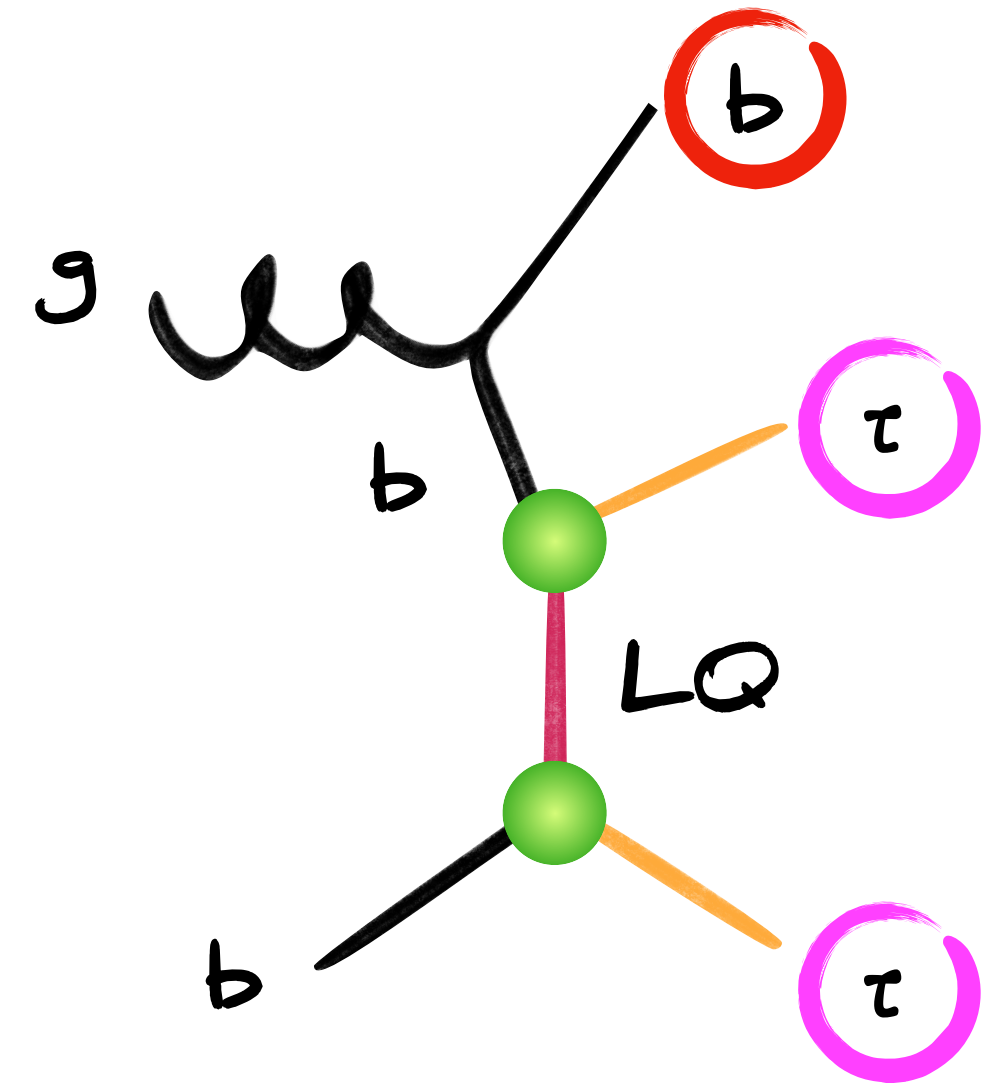
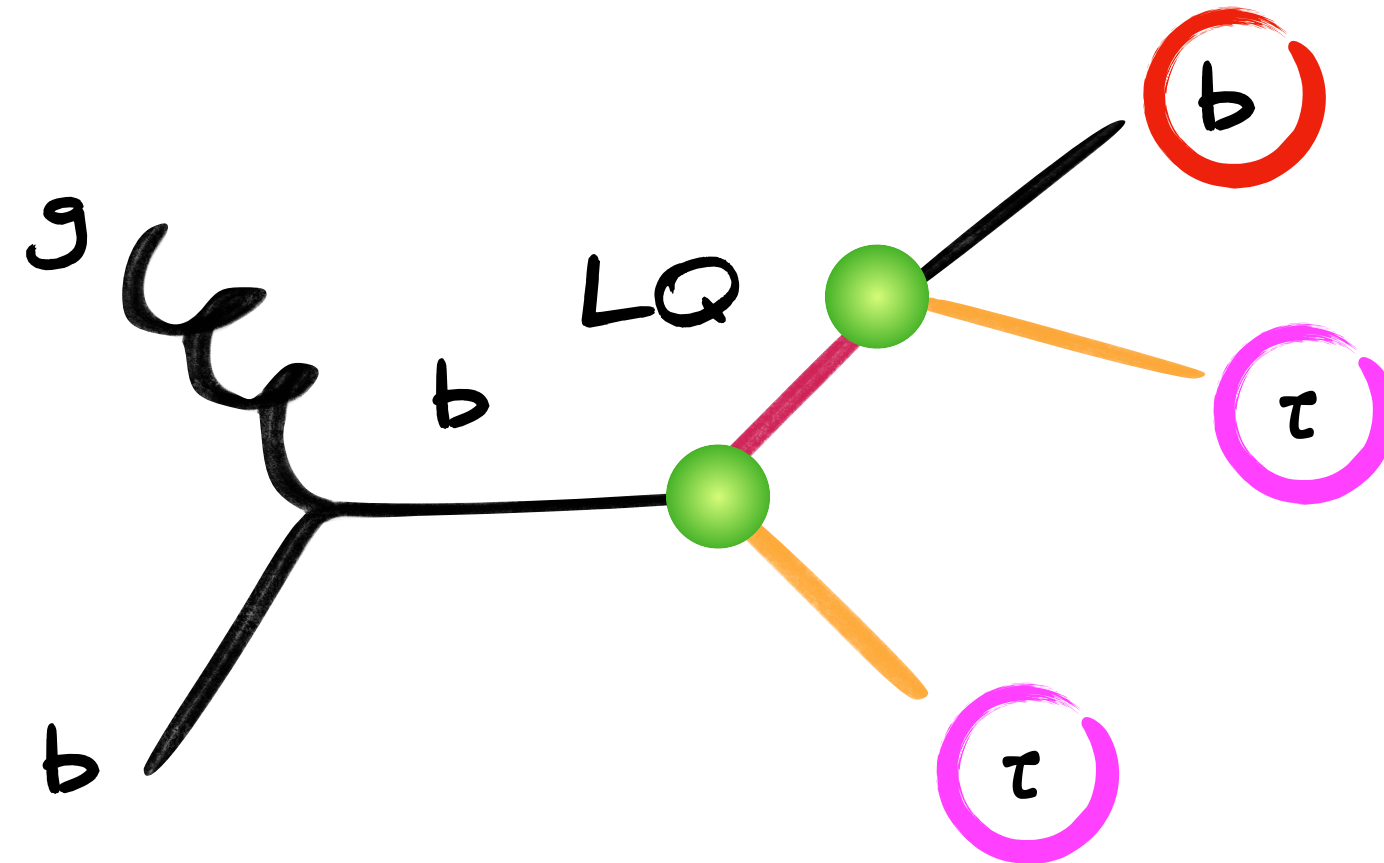
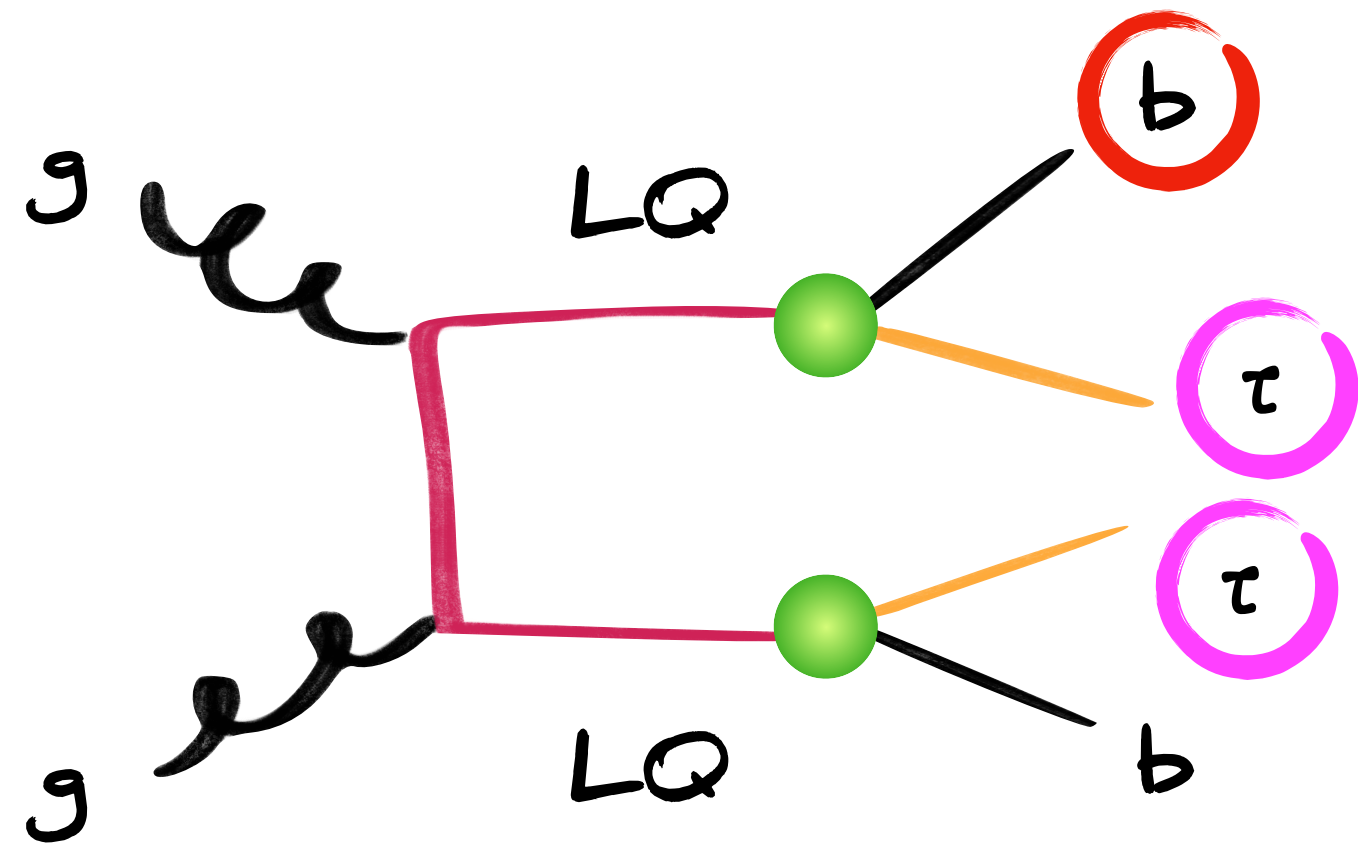
Existing LQ search strategies @ the LHC



Channels complementary because they provide different sensitivities on LQ parameter space

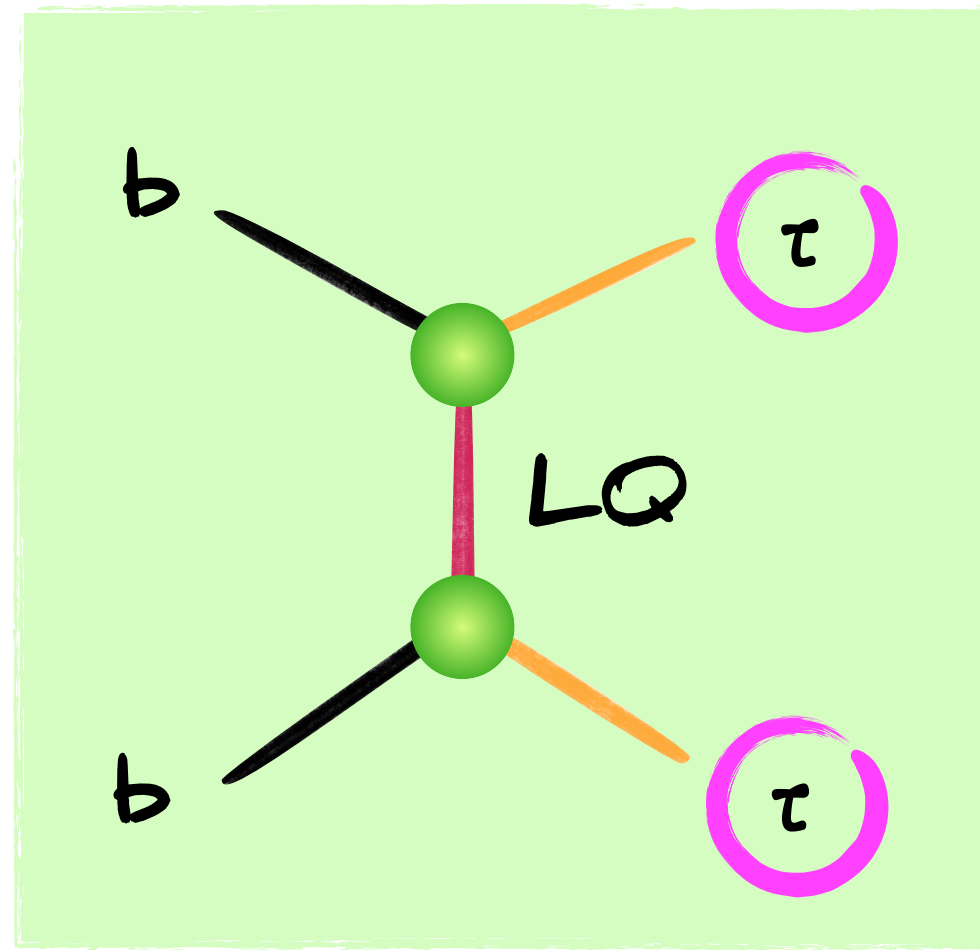
[sketch adopted from Dorsner & Greljo, 1801.07641]

Existing LQ search strategies @ the LHC

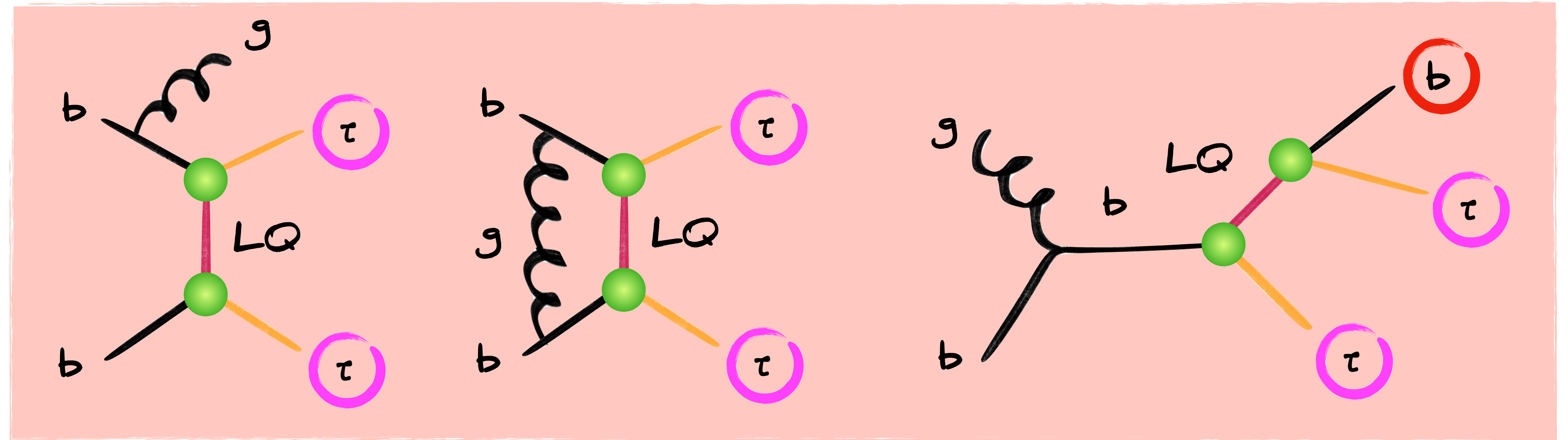


LQ production mechanisms studied by ATLAS & CMS lead to complicated multi-particle final states. Depending on experimental selections individual channels lead to overlapping contributions, which calls for a precise modelling of LQ signal at level of hadronic events

Existing LQ search strategies @ the LHC



tree level



real NLO

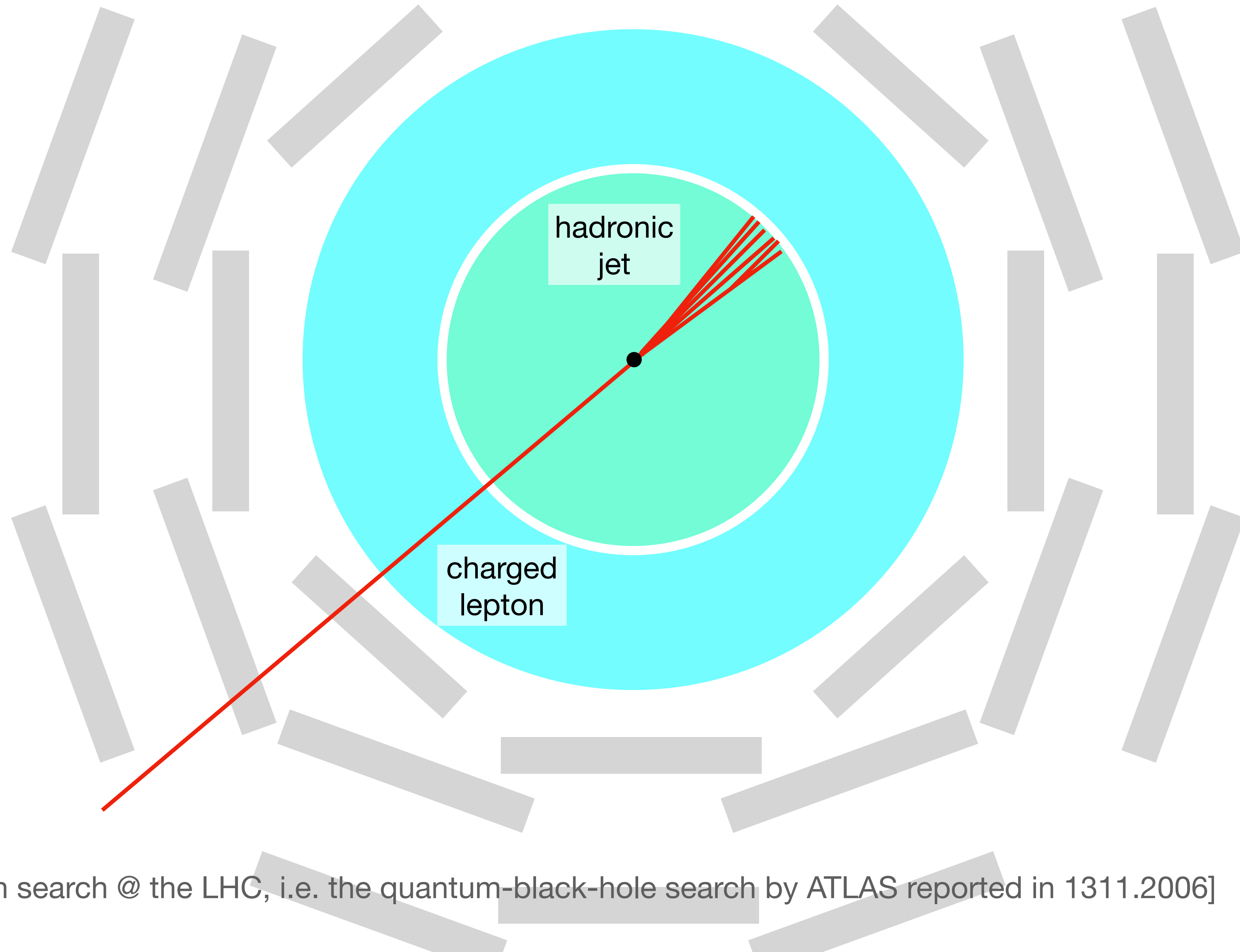
virtual NLO

finite real NLO

QCD effects particularly relevant for final states with b-jets due to small b-quark parton distribution function (PDF). LHC Run II results based on merging & matching but full next-to-leading order plus parton shower (NLO+PS) Monte Carlos now available in some cases

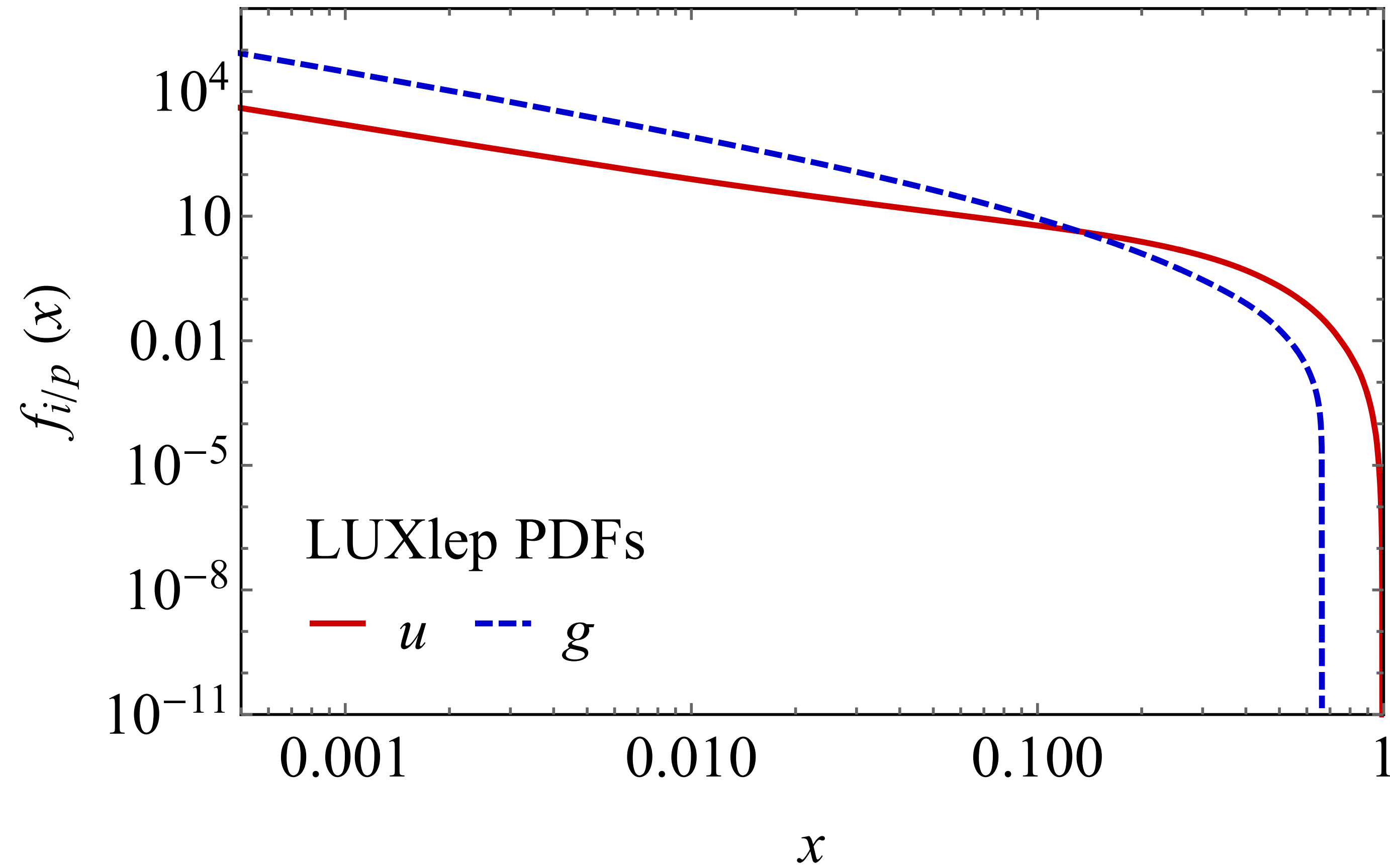
[see for instance NLO+PS POWHEG implementation of $pp \rightarrow l+l-$ based on UH, Schnell & Schulte, 2207.00356; 2209.12780]

Lepton-jet searches are much simpler ...

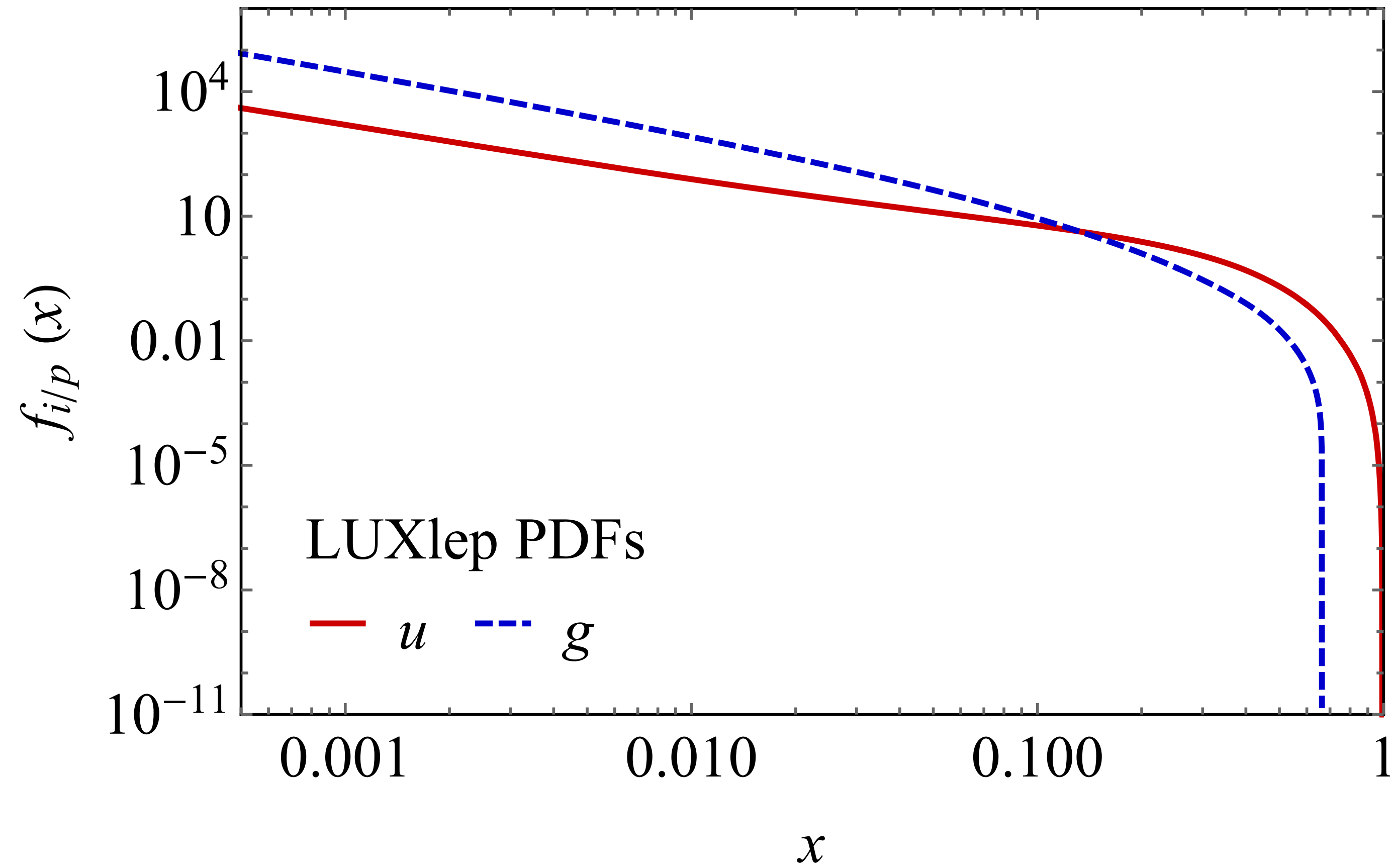


[there is only one such search @ the LHC, i.e. the quantum-black-hole search by ATLAS reported in 1311.2006]

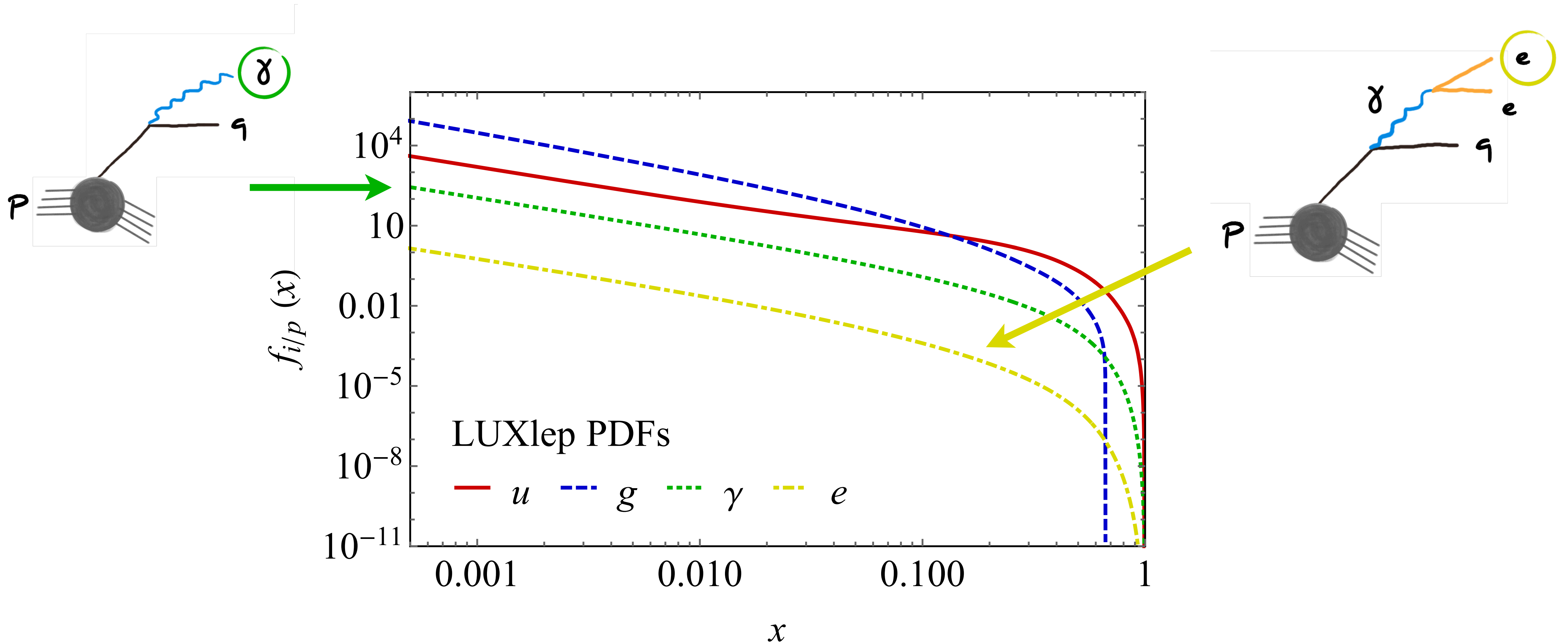
but LHC collides protons ...



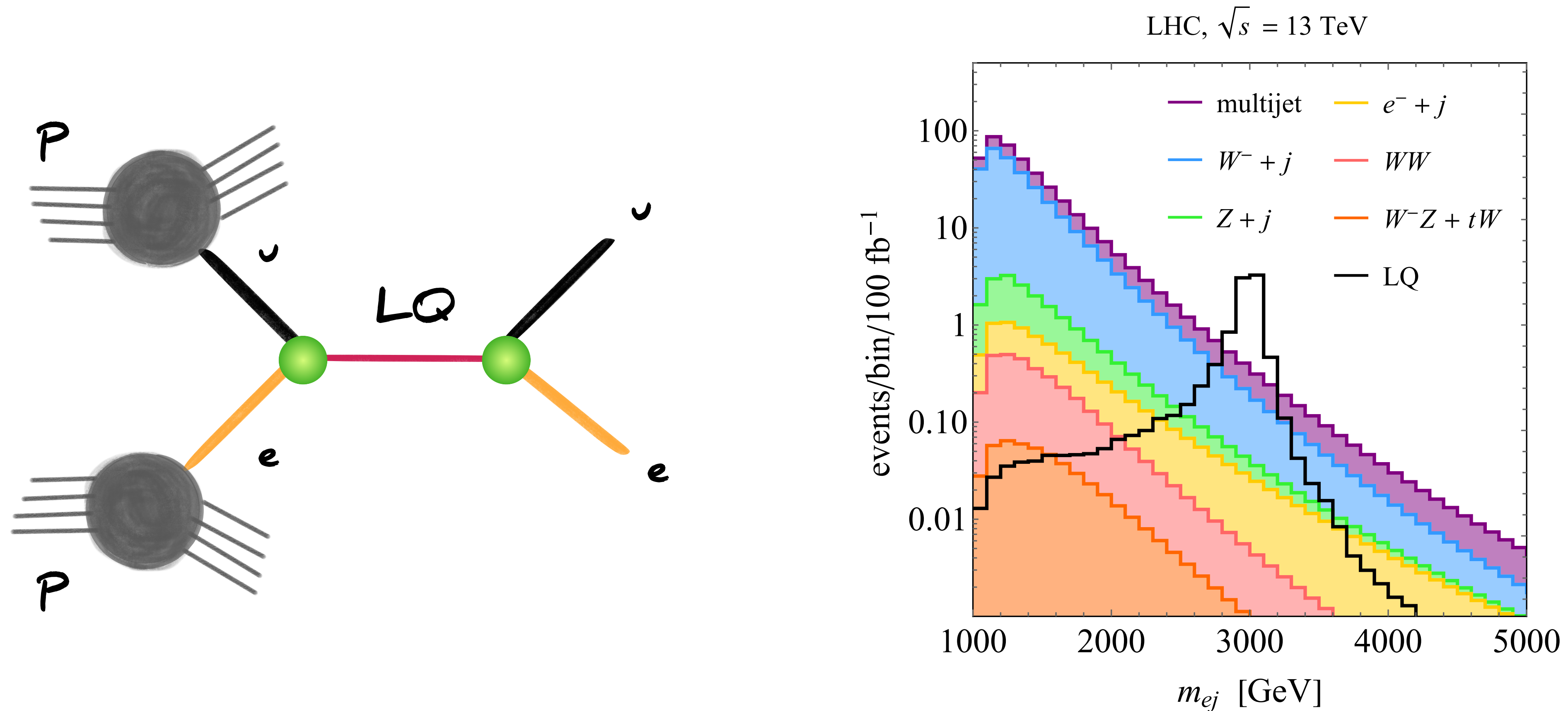
& protons consist of quarks & gluons



Quantum field theory to the rescue!



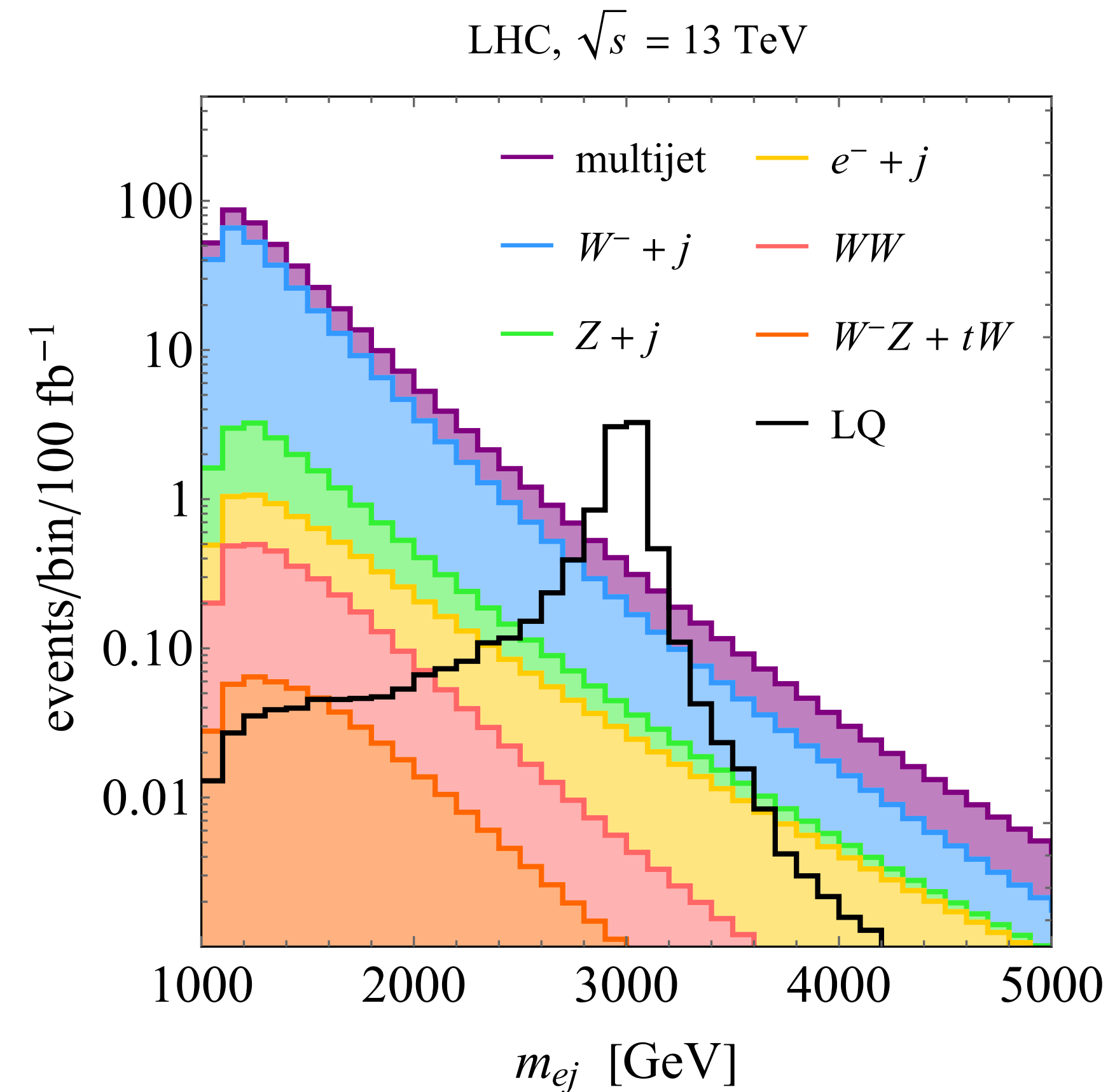
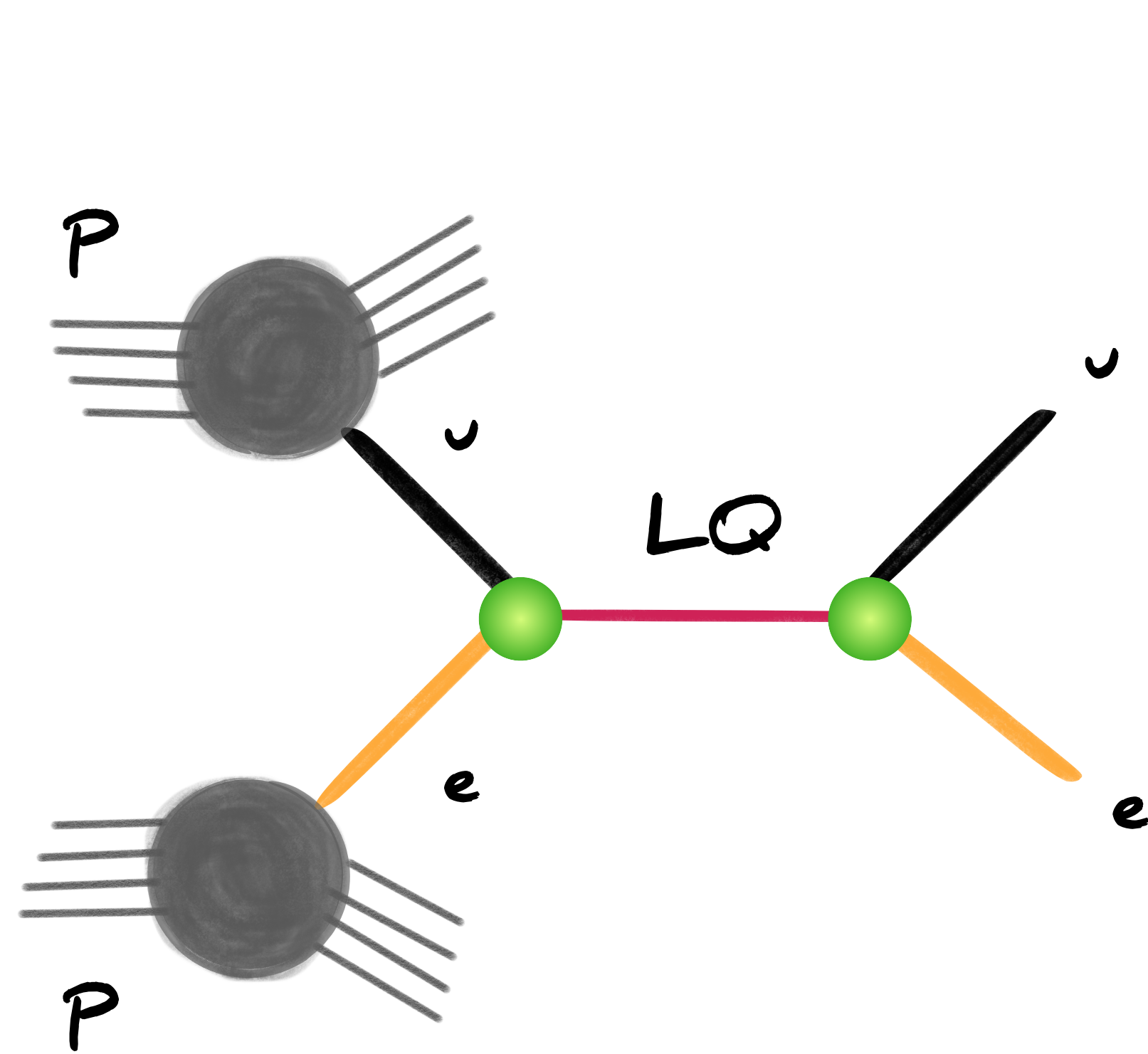
Resonant LQ production @ LHC



Non-zero lepton PDFs allow for resonant LQ production @ pp machines such as LHC

[search strategy proposed in Buonocore et al., 2005.06475 & refined in Greljo & Selimovic, 2012.02092; Buonocore et al., 2209.02599]

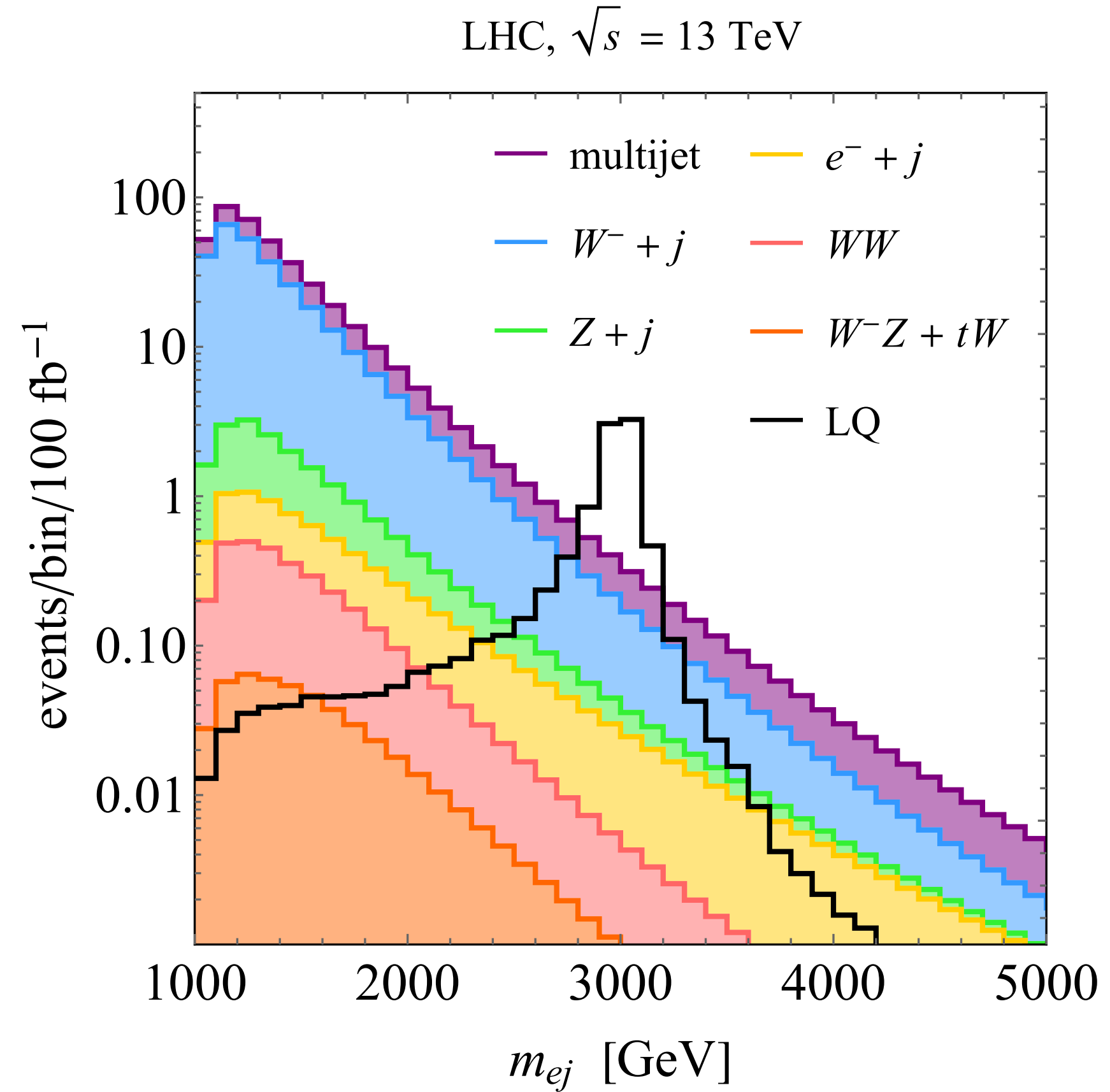
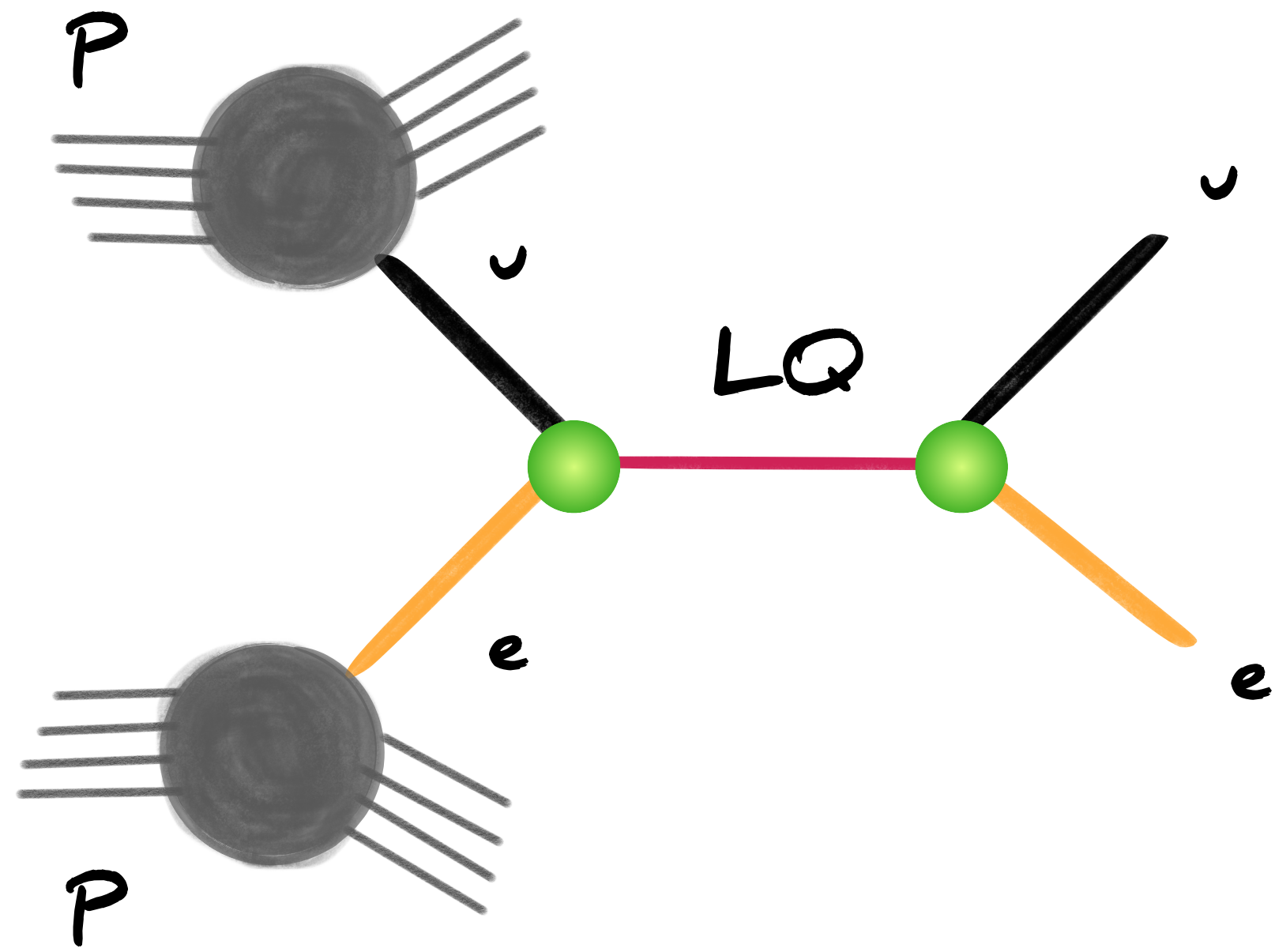
Resonant LQ production @ LHC



Sum over backgrounds is a steeply falling distribution, while signal exhibits a narrow peak

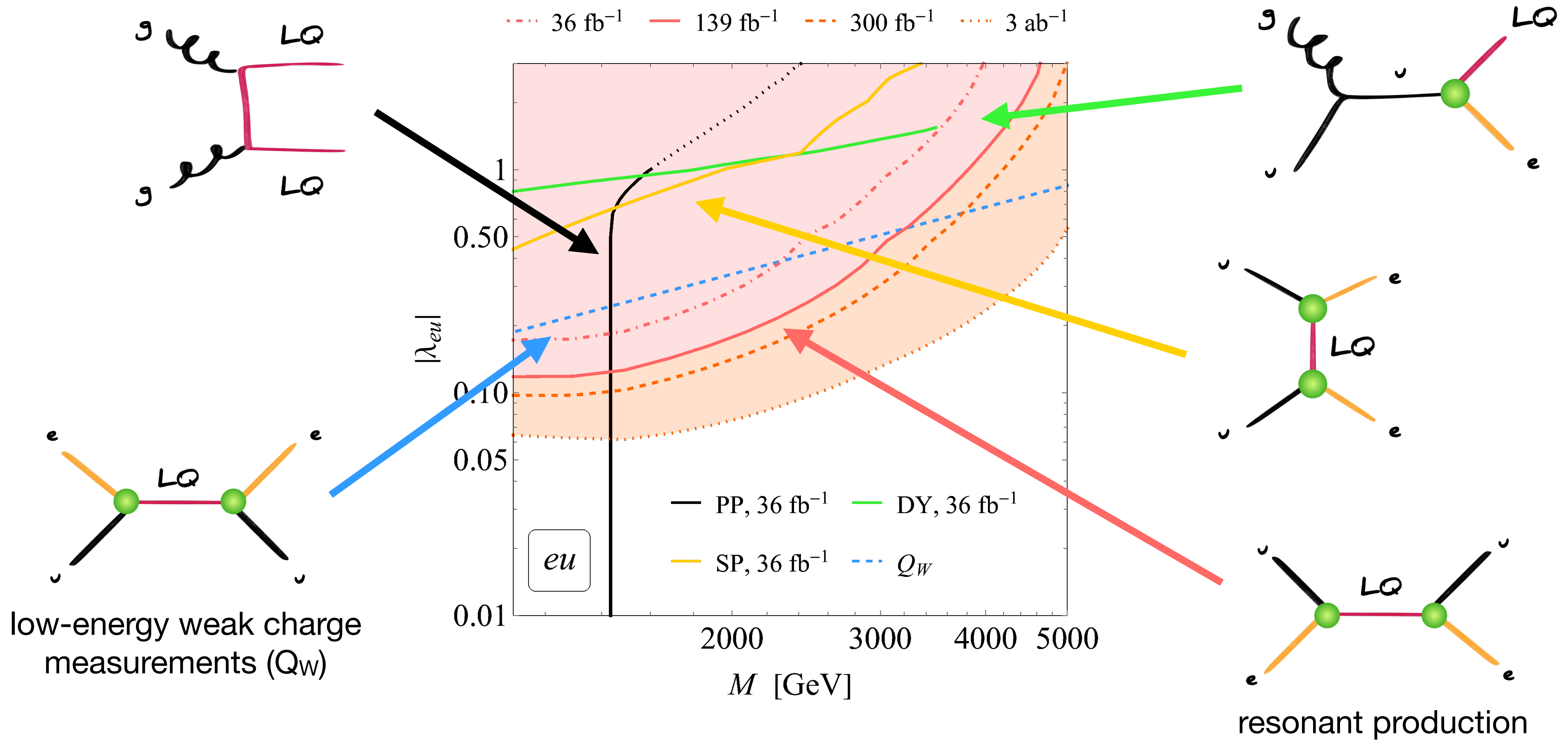
[search strategy proposed in Buonocore et al., 2005.06475 & refined in Greljo & Selimovic, 2012.02092; Buonocore et al., 2209.02599]

Resonant LQ production @ LHC



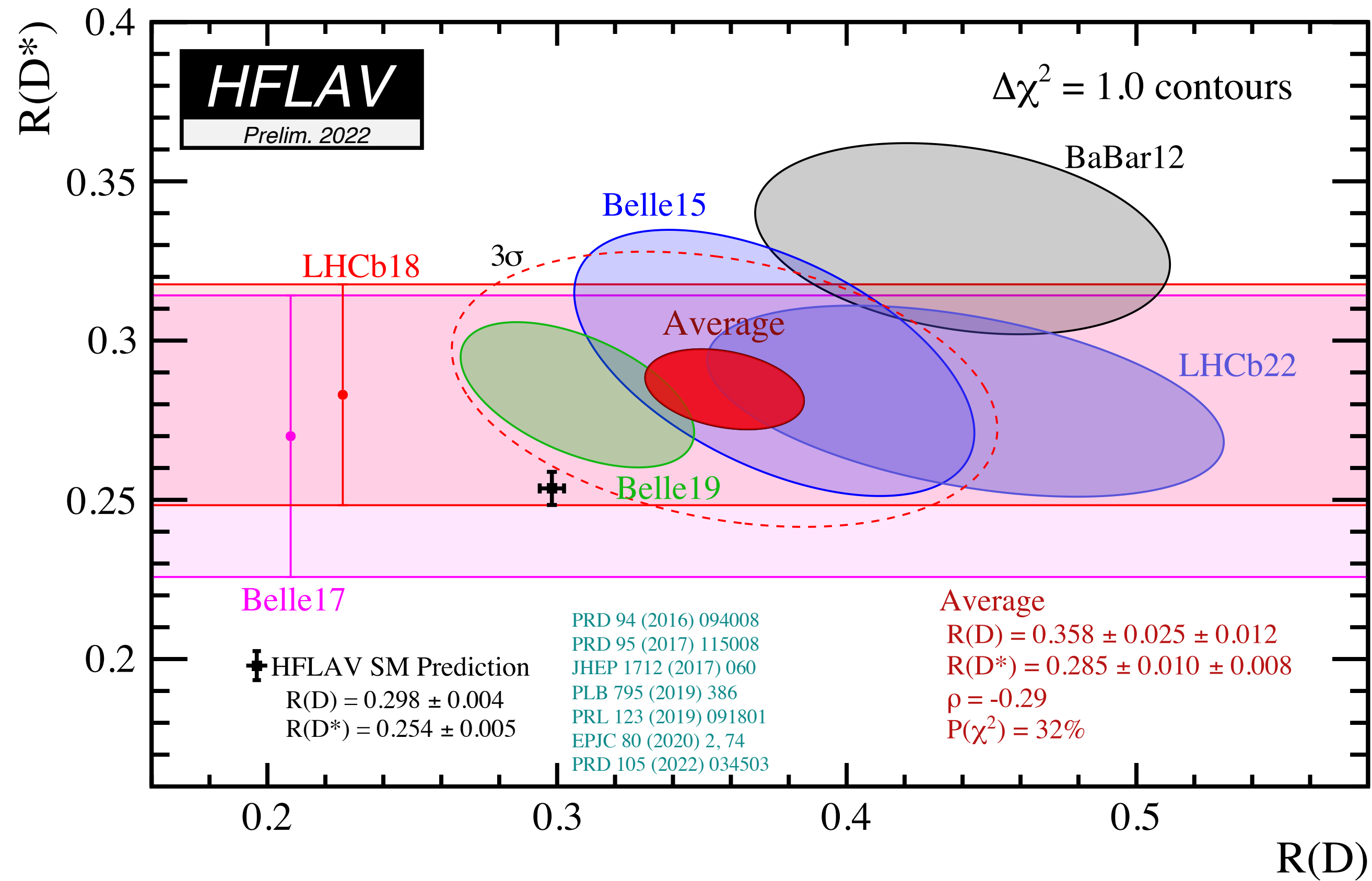
Compared to 1 background event, 9 events per 100 fb⁻¹ for LQ of $M = 3$ TeV & $\lambda_{eu} = 1$ @ 13 TeV

LHC limits on 1st & 2nd generation LQs



[Buonocore, UH, Nason, Tramontano & Zanderighi, 2005.06475]

Lepton flavour non-universality in $b \rightarrow c\ell\nu$



BaBar, Belle & LHCb measurements of lepton flavour universality ratios $R_{D^{(*)}}$ show a 3σ tension

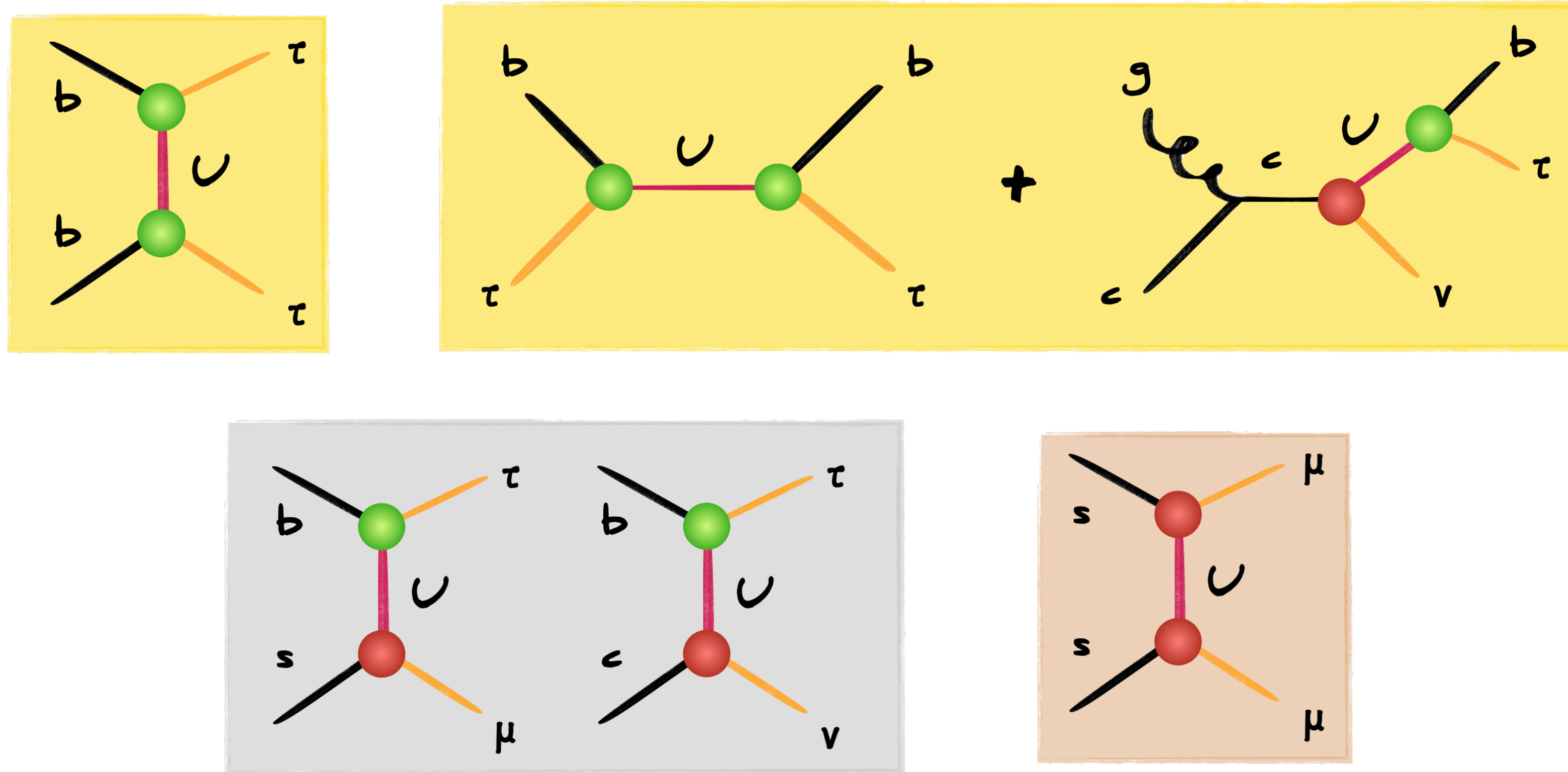
Singlet vector LQ models for B anomalies

$$\mathcal{L} \supset \frac{g_U}{\sqrt{2}} \left[\beta_L^{ij} \bar{Q}_L^{i,a} \gamma_\mu L_L^j + \beta_R^{ij} \bar{d}_R^{i,a} \gamma_\mu \ell_R^j \right] U^{\mu,a} + \text{h.c.}$$

$$|\beta_L^{22}| \lesssim |\beta_L^{32}| \ll |\beta_L^{23}| \lesssim |\beta_L^{33}| = \mathcal{O}(1)$$

Parameters		Branching ratios			
β_L^{33}	β_L^{23}	BR ($U \rightarrow b\tau^+$)	BR ($U \rightarrow t\bar{\nu}_\tau$)	BR ($U \rightarrow s\tau^+$)	BR ($U \rightarrow c\bar{\nu}_\tau$)
1	0	51%	49%	0%	0%
1	1	25%	22%	25%	27%

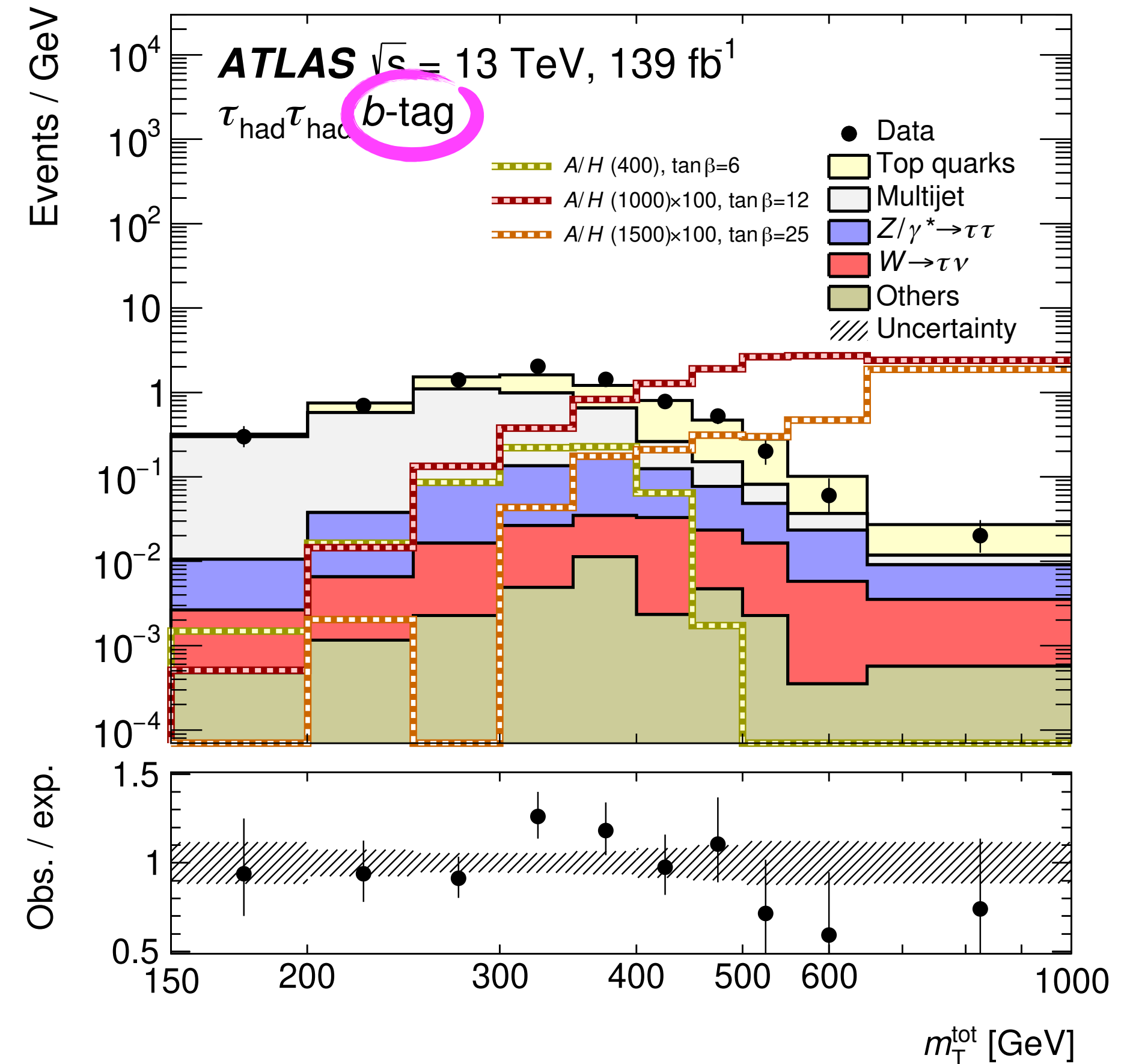
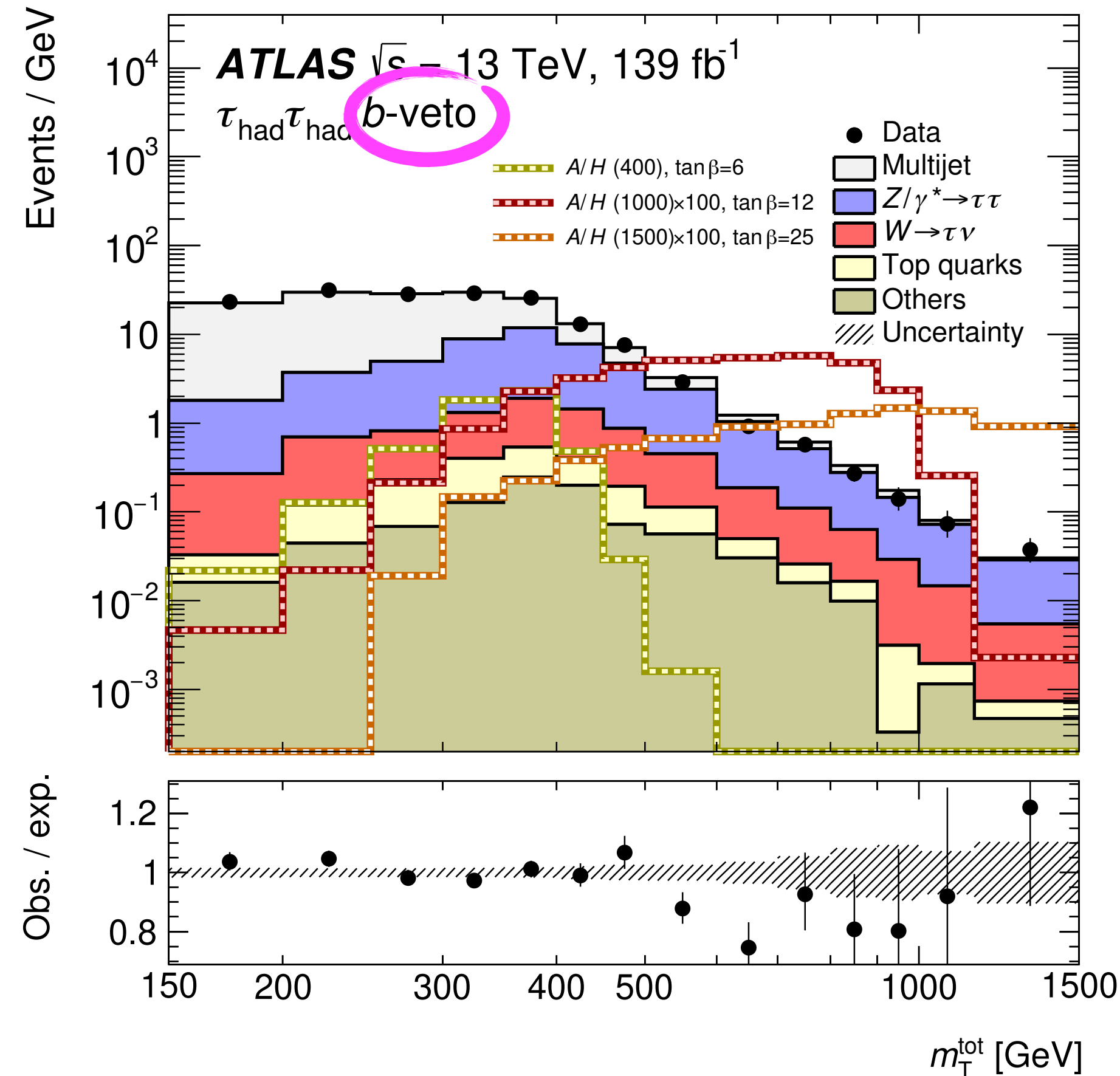
Possible singlet vector LQ signatures



$b \rightarrow cl\nu$ anomalies single out $pp \rightarrow \tau^+\tau^-$ as most interesting channel. After latest LHCb measurements of $R_{K^{(*)}}$, $pp \rightarrow b\tau$, $\tau\mu$, $\tau\nu$ & $\mu^+\mu^-$ production seem less relevant

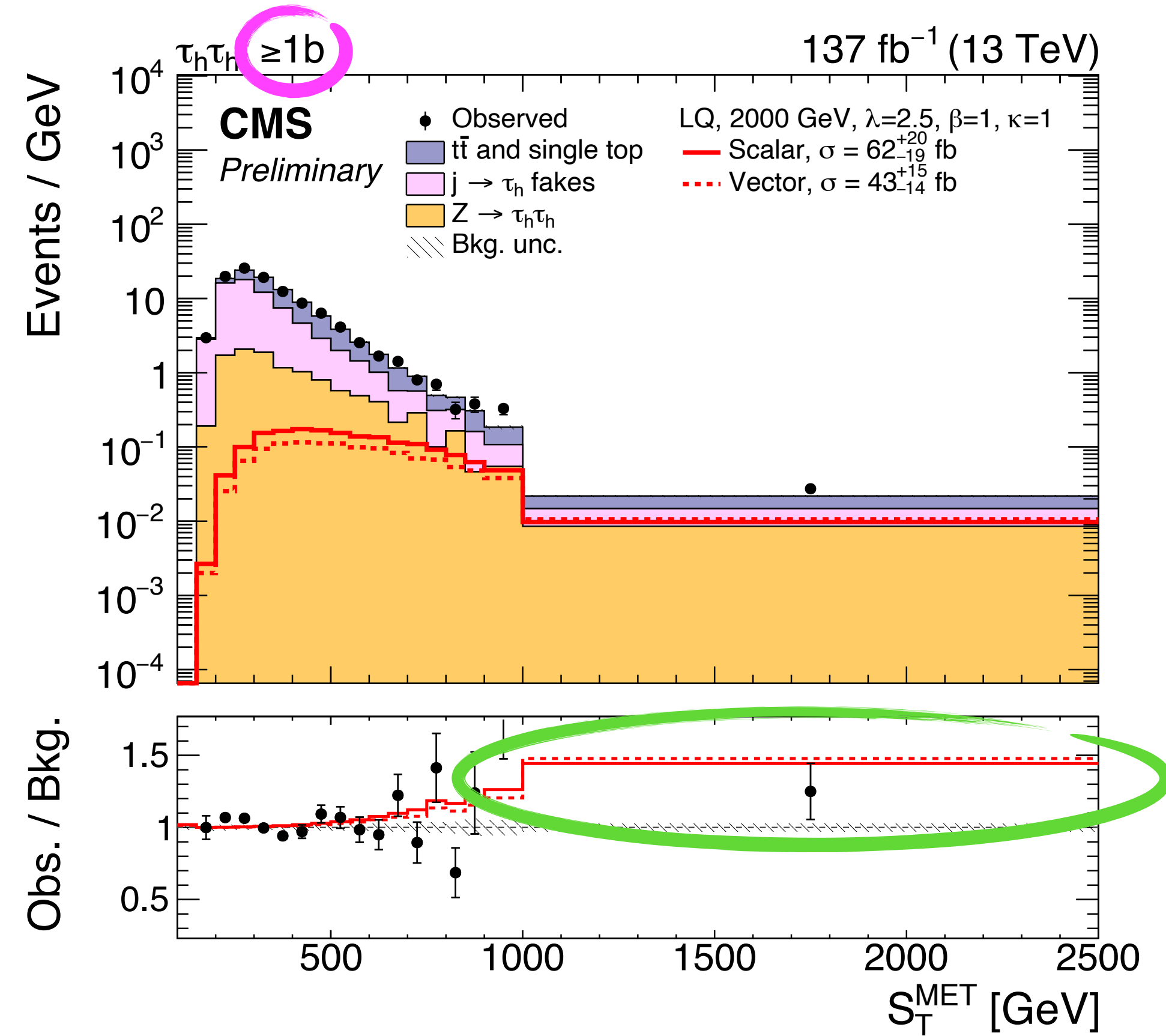
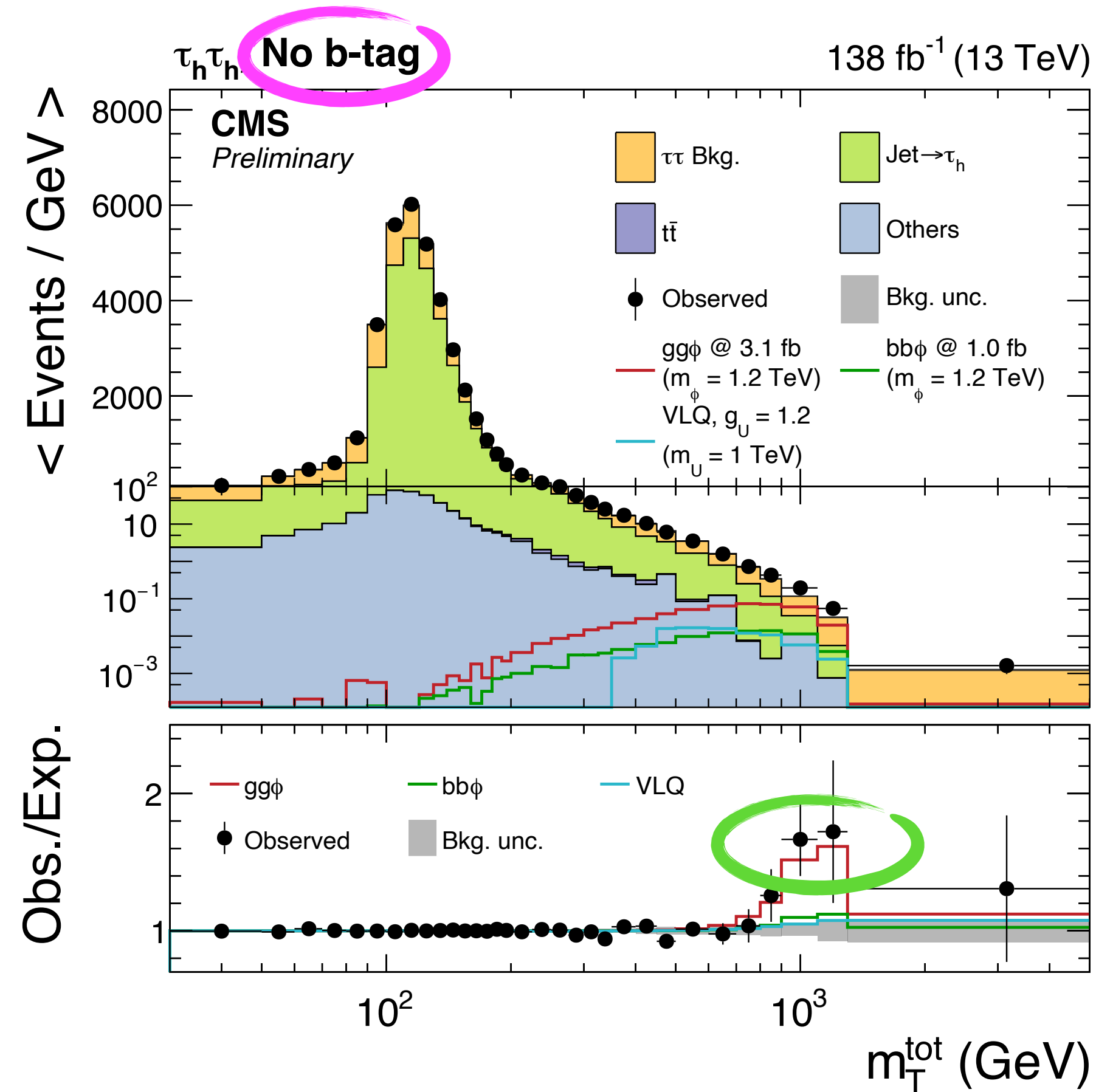
[singlet vector LQ effects in $pp \rightarrow b\tau$, mono-top & mono-jet production have been studied in UH & Polesello, 2012.11474]

Ditau searches @ LHC Run II



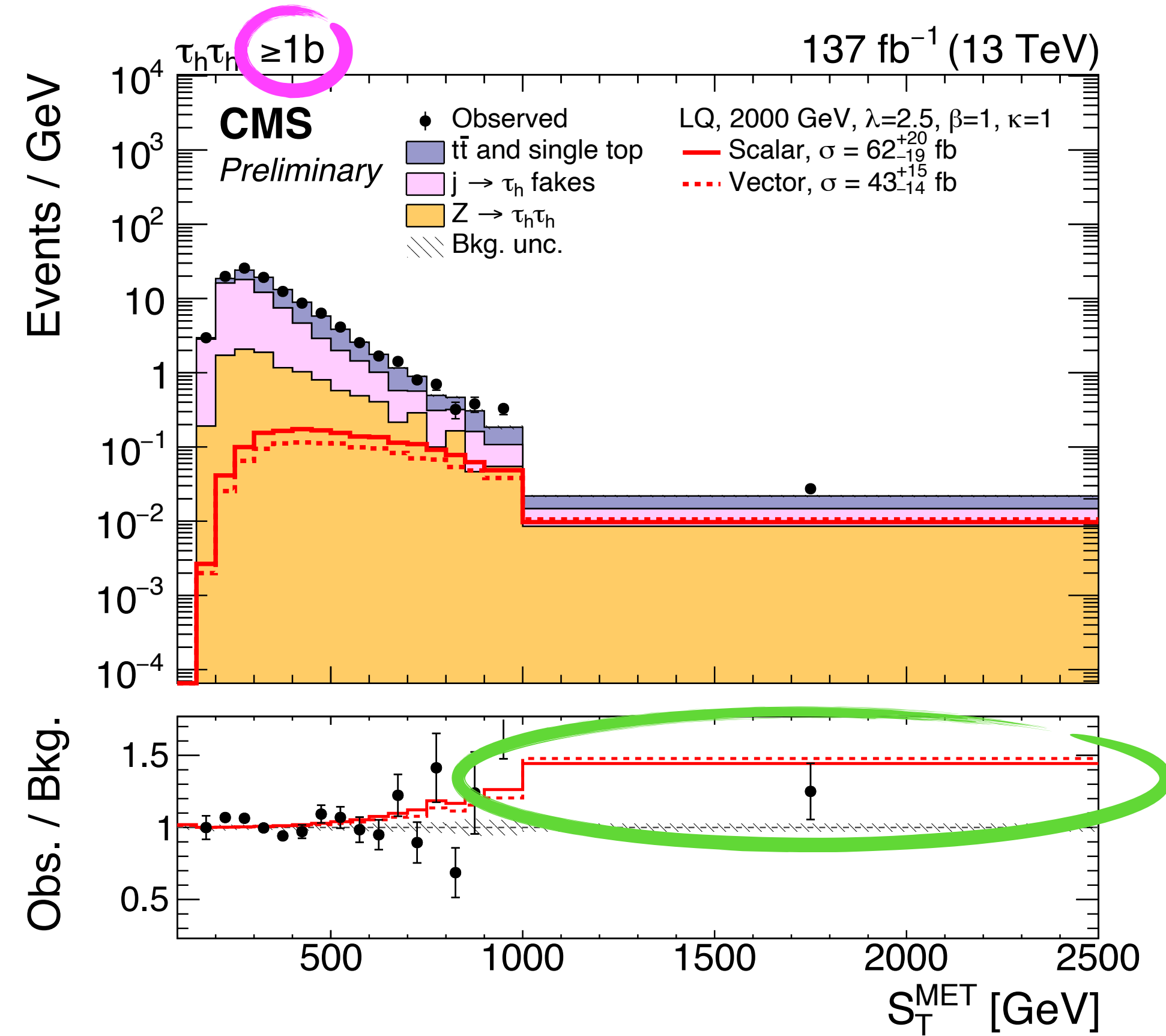
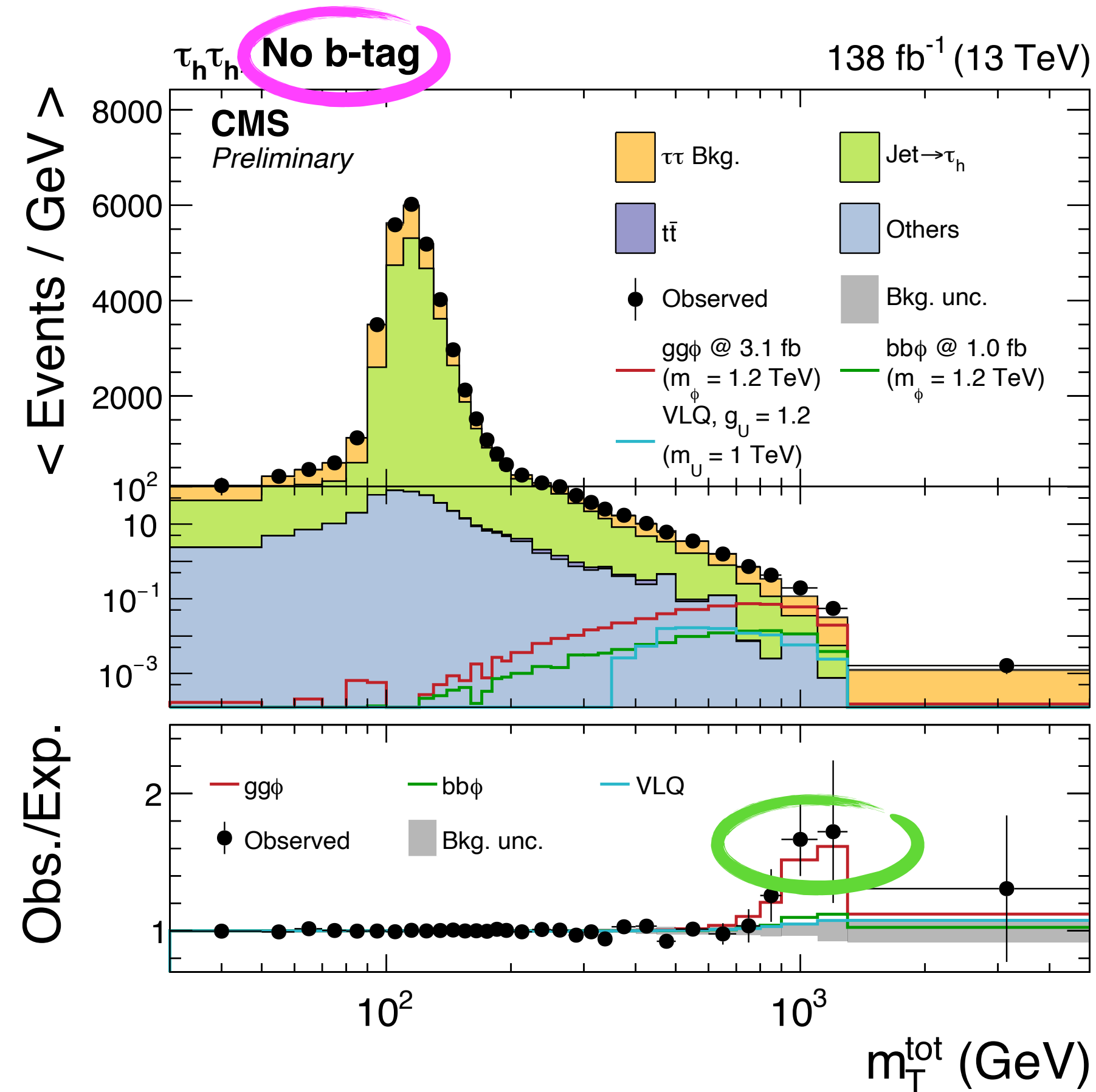
Three different ditau LHC Run II analyses, all considering events without & with an extra b-jet

Ditau searches @ LHC Run II



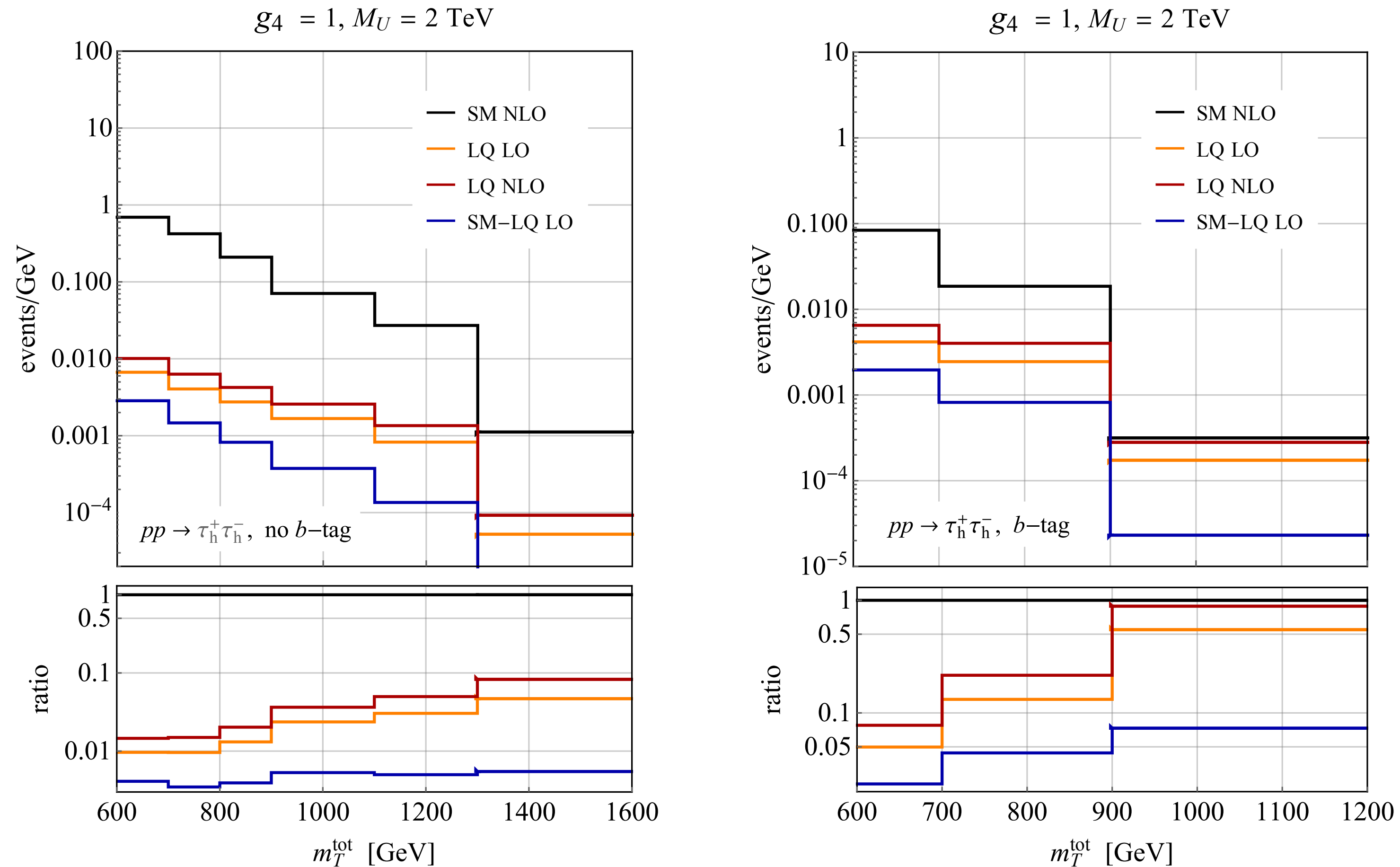
ATLAS data agrees with background predictions but both CMS analyses see a 3σ excess

Ditau searches @ LHC Run II



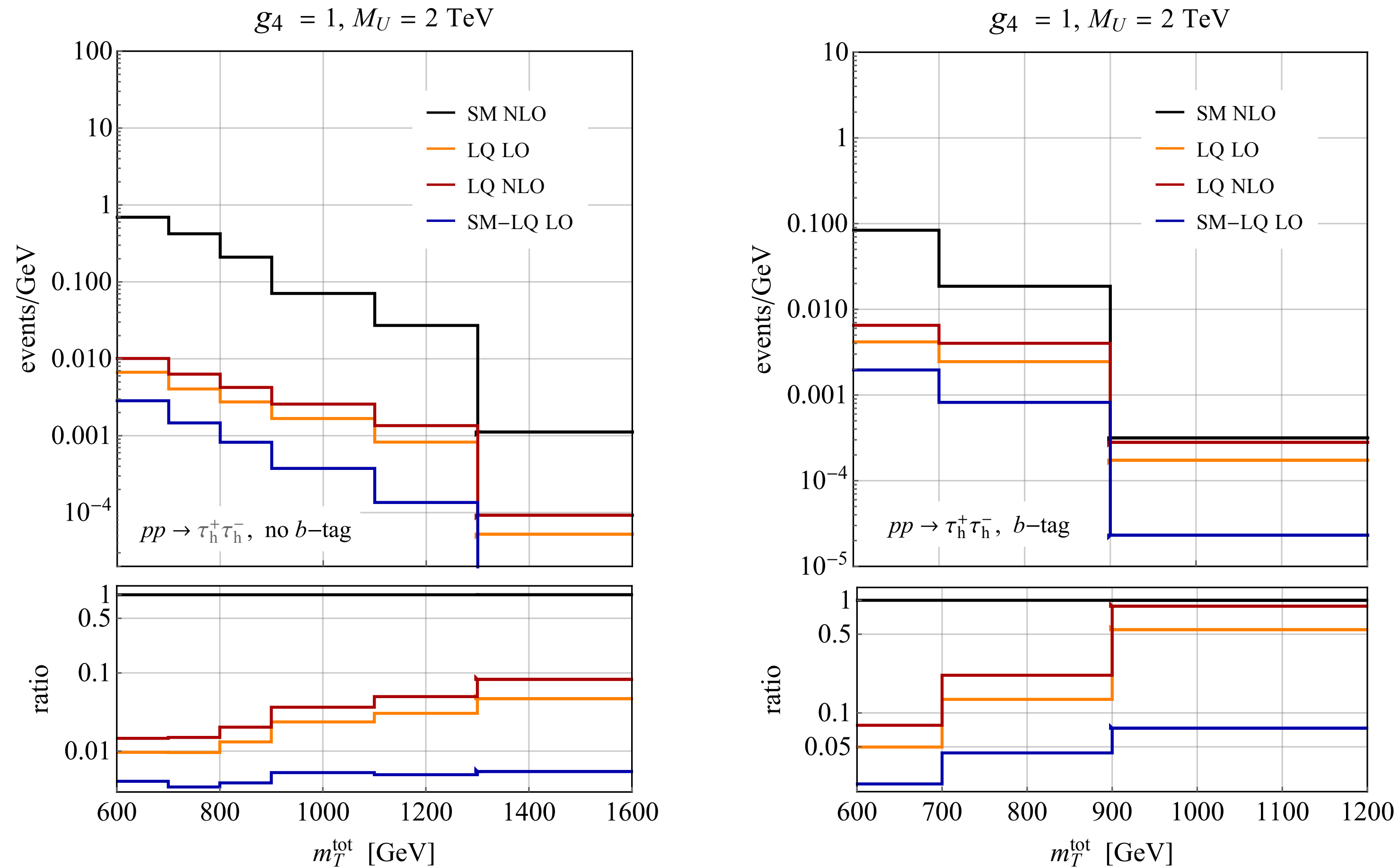
Non-resonant (resonant) excess in b-tag (b-veto) sample fits (does not fit) LQ explanation

Size of NLO & interference effects



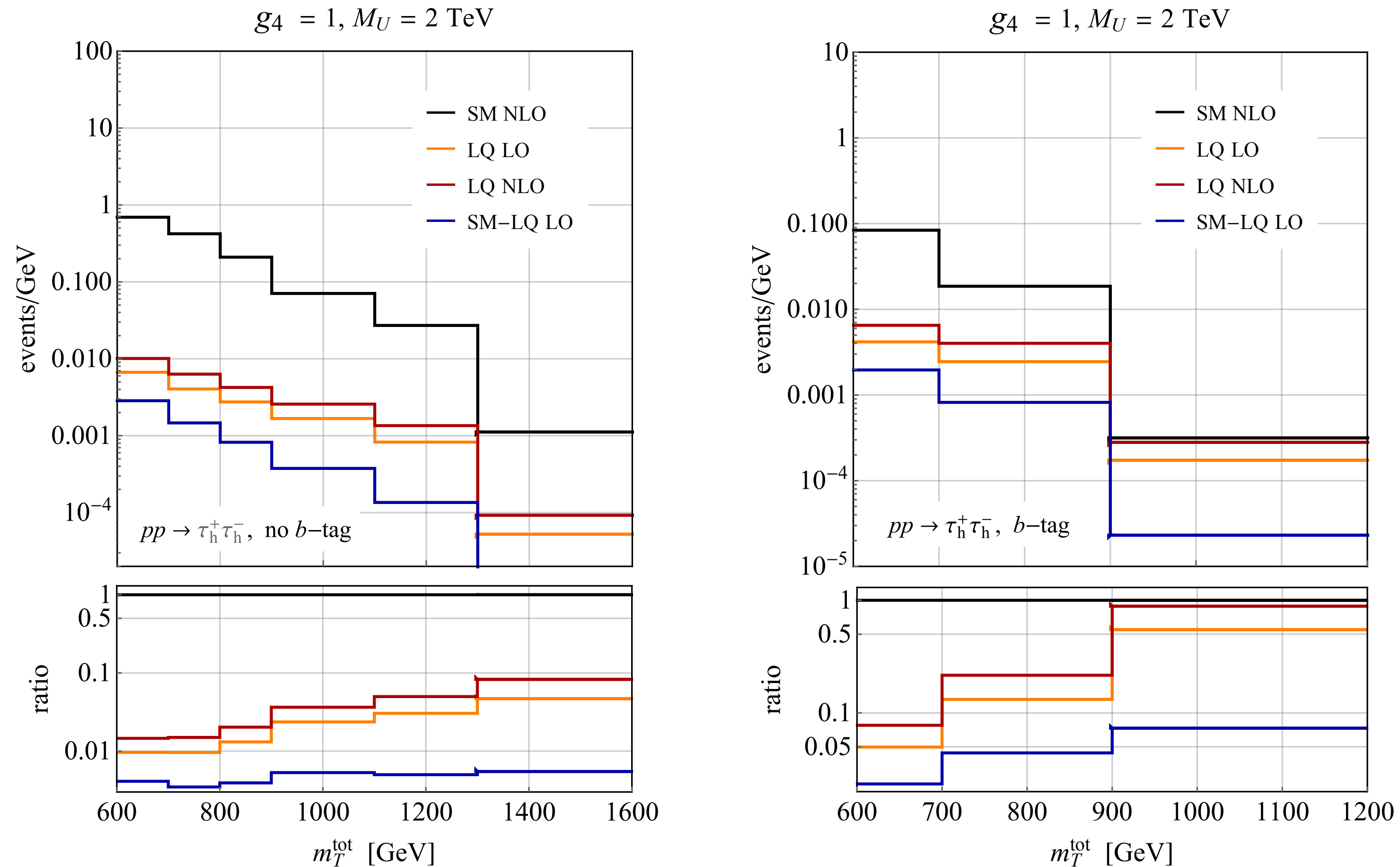
In b-veto (b-tag) sample, LQ corrections amount to 10% (85%) compared to DY background

Size of NLO & interference effects



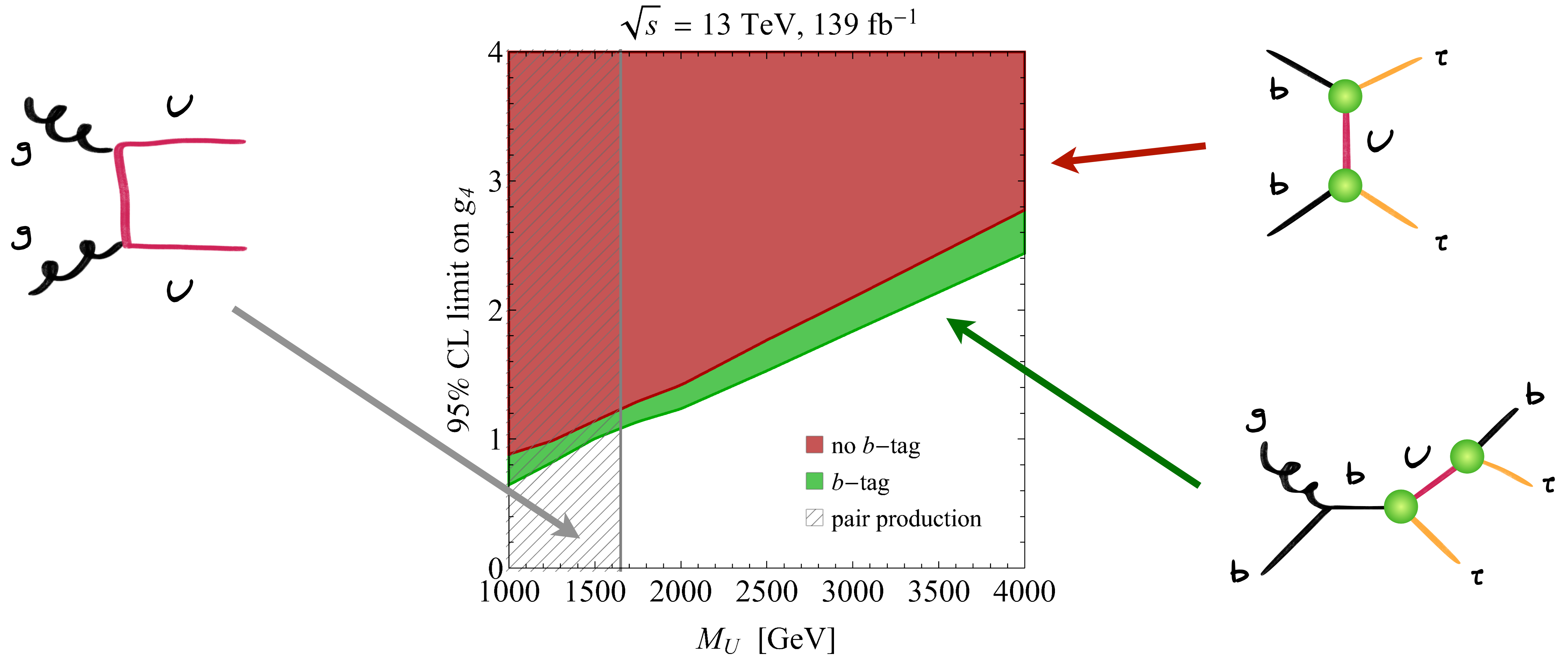
NLO QCD corrections exceed 50% in signal regions & grow in size with transverse mass

Size of NLO & interference effects



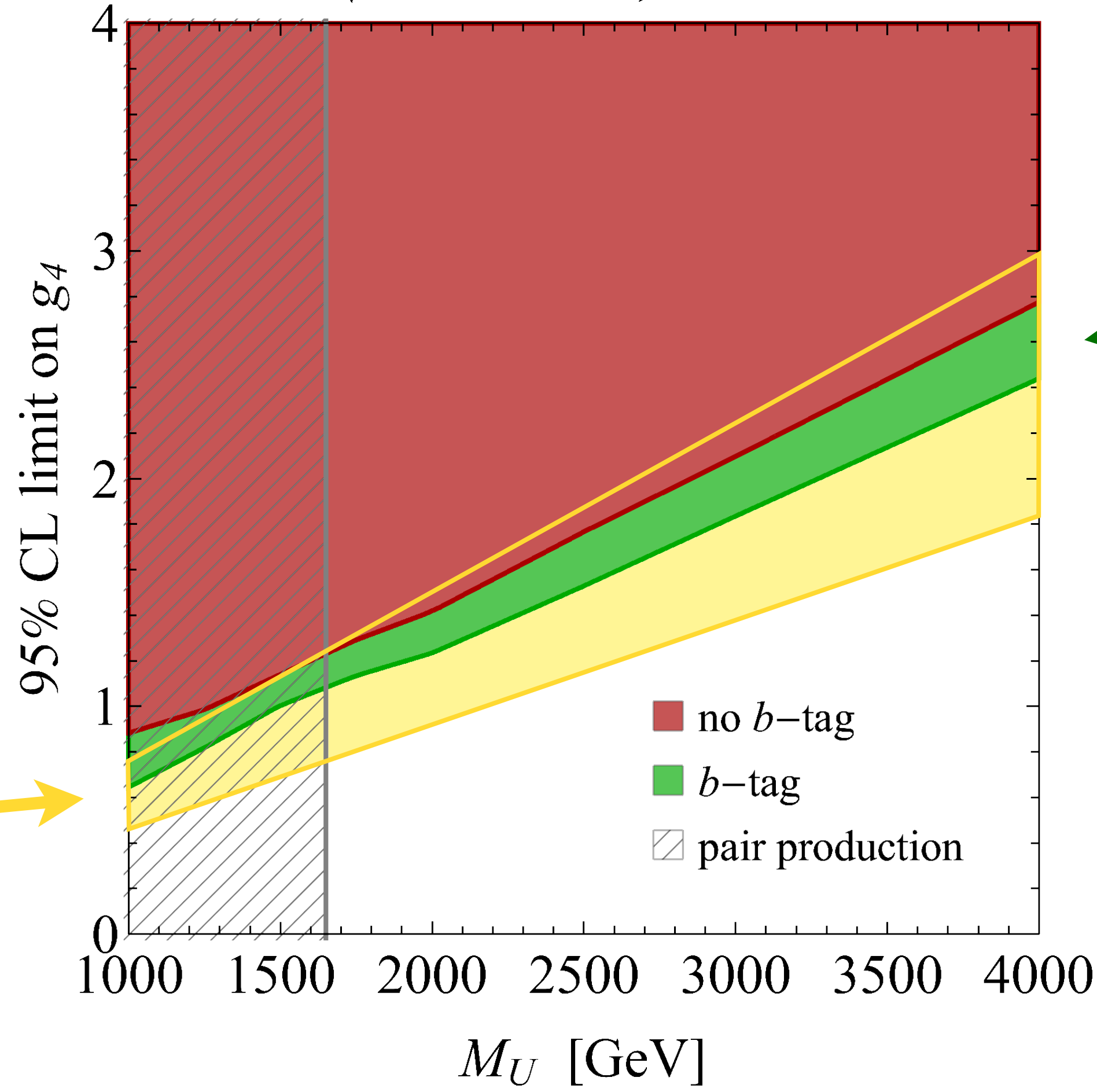
Relative to LQ tree-level contributions interference effects do not exceed level of 10%

ATLAS ditau limits on singlet vector LQs



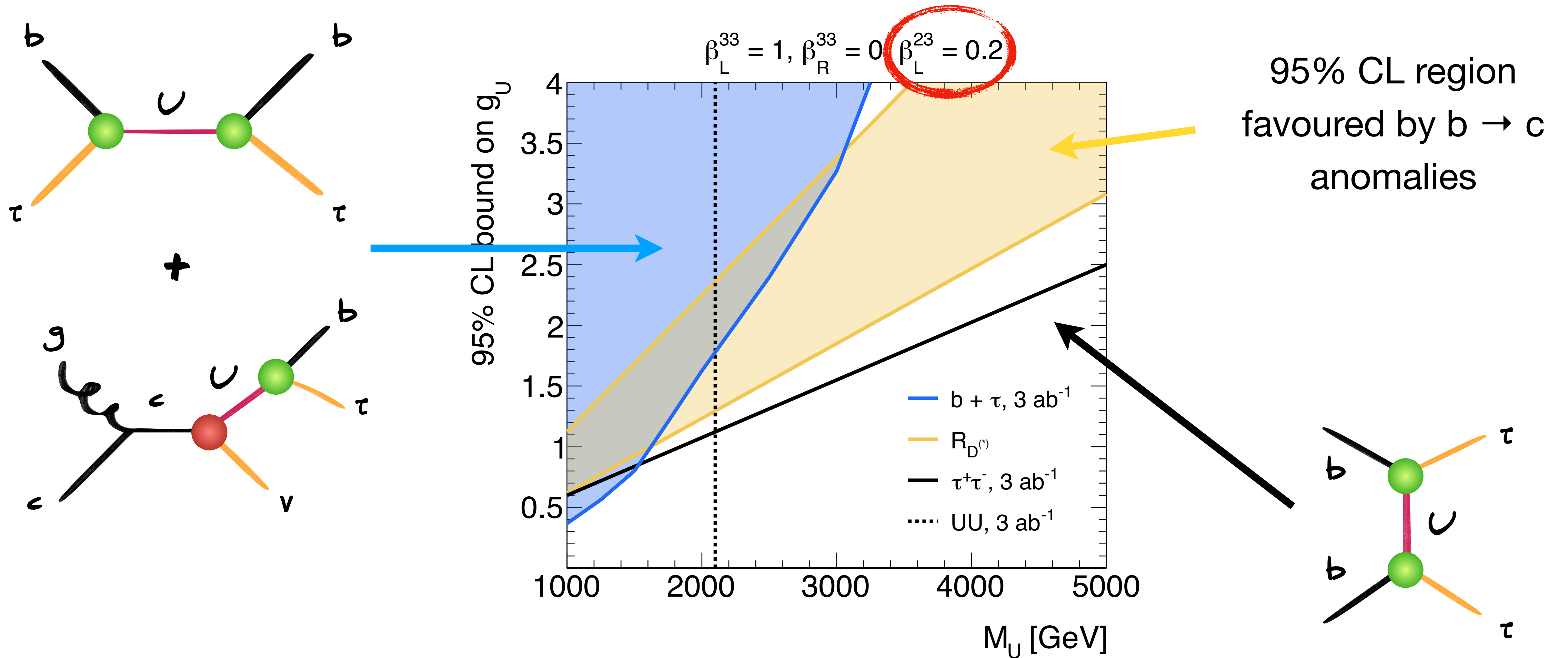
ATLAS ditau limits on singlet vector LQs

$\sqrt{s} = 13 \text{ TeV}, 139 \text{ fb}^{-1}$



ditau constraints start to test LQ explanations of $b \rightarrow c$ anomalies — leading limits arise from events with b-jets

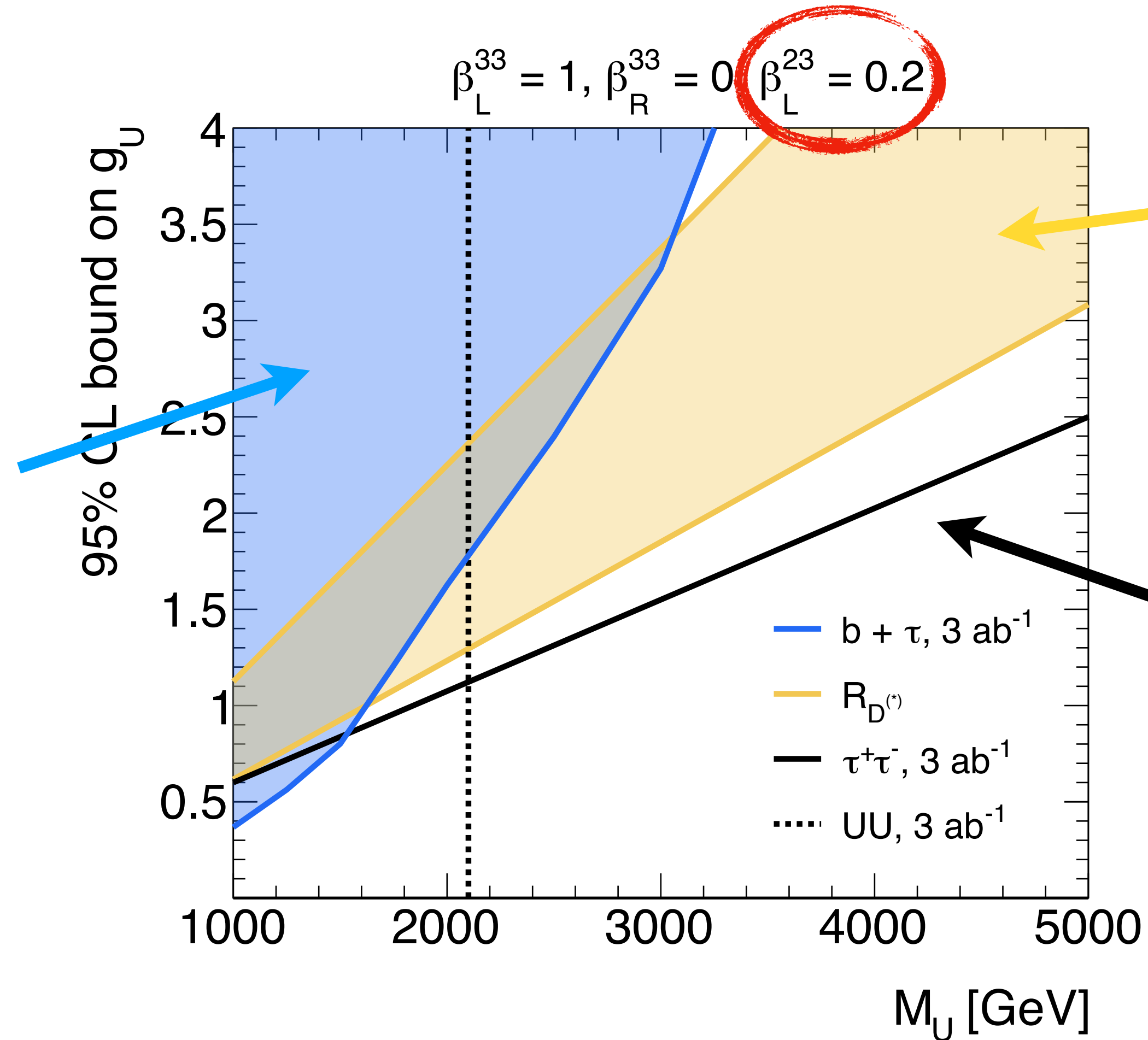
LHC 3 ab⁻¹ projections for singlet vector LQs



[UH & Polesello, 2012.11474; Cornella, Faroughy, Fuentes-Martin, Isidori & Neubert, 2103.16558]

LHC 3 ab⁻¹ projections for singlet vector LQs

weaker but complementary information provided by searches for resonant 3rd-generation LQ signatures



95% CL region favoured by $b \rightarrow c$ anomalies

probably all singlet vector LQ explanations of $b \rightarrow c$ anomalies can be tested with 3 ab⁻¹ of data via ditau searches

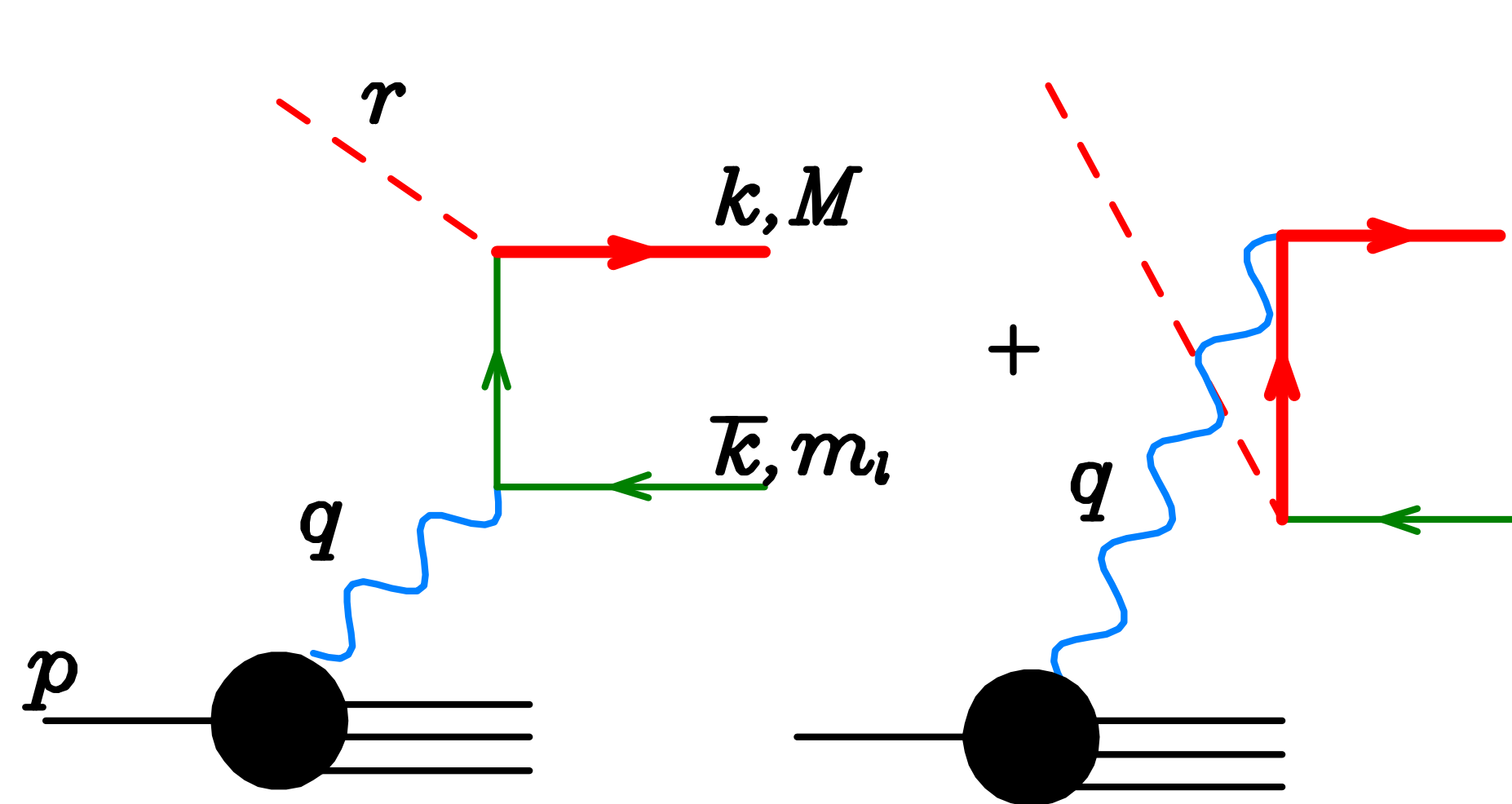
Conclusions & outlook

- Precision determination of lepton PDFs opens up new ways to test Standard Model (e.g. $l^\pm l^\pm$ production) & to search for new physics @ the LHC. For instance, resonant LQ production allows to probe so far unexplored parameter space & has discovery potential
- Models of new physics such as singlet vector LQs that explain $b \rightarrow clv$ anomalies generically lead to signatures (e.g. $pp \rightarrow \tau^+ \tau^-$, $b\tau$, $t\bar{t}$ & high-multiplicity final states with τ , b , t & $E_{T,miss}$) testable @ the LHC. High-luminosity LHC needed to cover full theory space

Backup



Computation of lepton PDFs



leptonic tensor

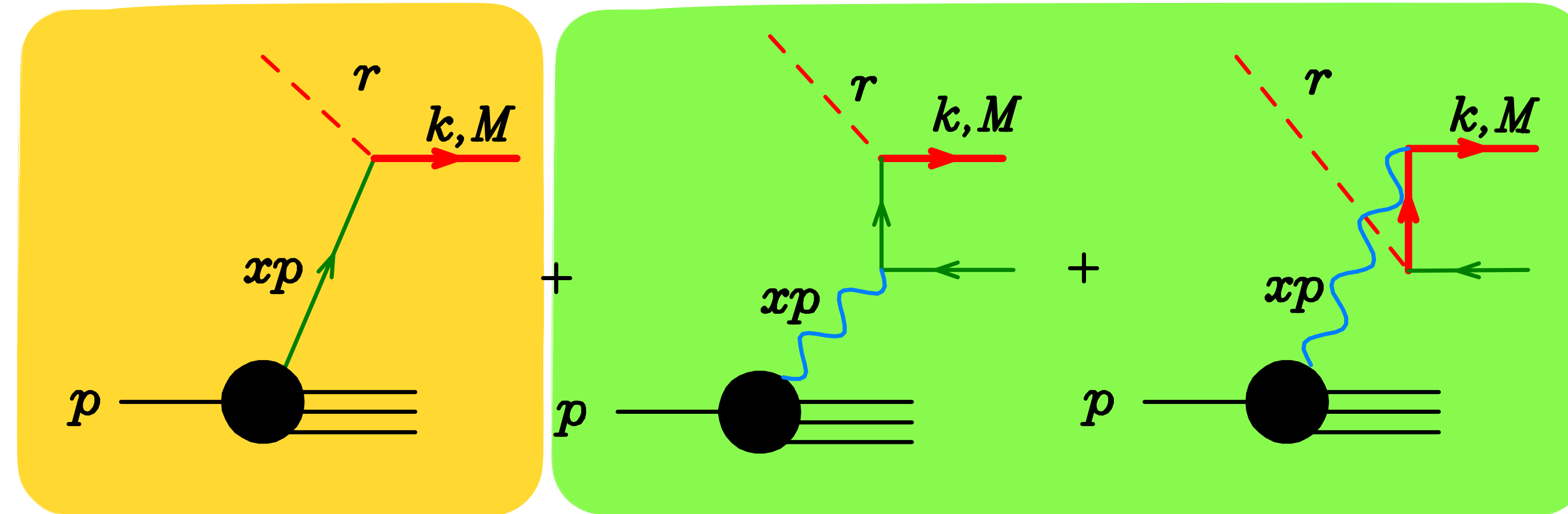
hadronic tensor

$$\sigma = \frac{1}{4p \cdot r} \int \frac{d^4 q}{(2\pi)^4} \frac{1}{Q^4} L^{\mu\nu}(r, q) (4\pi) W_{\mu\nu}(p, q)$$

proton structure functions that can be precisely extracted from low-energy ep scattering data

$$W_{\mu\nu}(p, q) = F_1 \left(-g_{\mu\nu} + \frac{q_\mu q_\nu}{q^2} \right) + \frac{F_2}{p \cdot q} \left(p_\mu - \frac{p \cdot q}{q^2} q_\mu \right) \left(p_\nu - \frac{p \cdot q}{q^2} q_\nu \right)$$

Computation of lepton PDFs

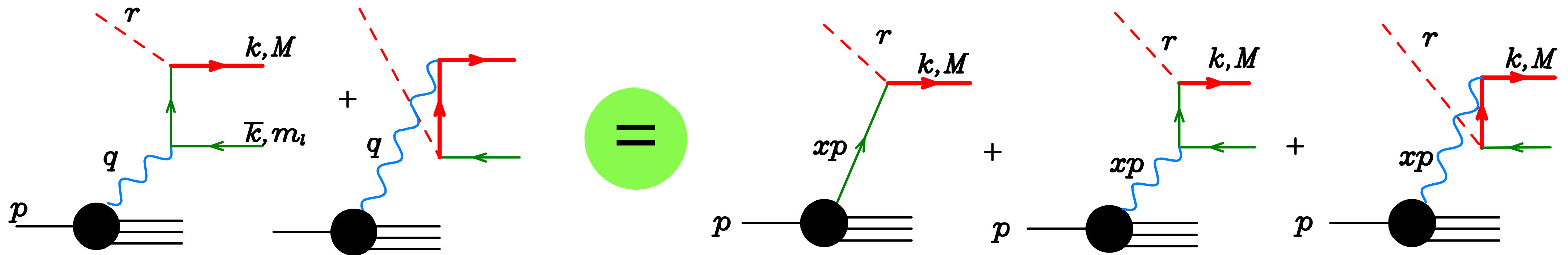


lepton PDF

photon PDF

$$\frac{\sigma}{\sigma_B} = \int dx f_\ell(x, \mu_F^2) \delta(Sx - M^2) + \frac{\alpha}{2\pi} \frac{1}{M^2} \int_{\frac{M^2}{S}}^1 dx f_\gamma(x, \mu_F^2) \times \left\{ z_\ell P_{\ell\gamma}(z_\ell) \left[\log \frac{M^2}{\mu_F^2} + \log \frac{(1 - z_\ell)^2}{z_\ell^2} \right] + 4z_\ell^2(1 - z_\ell) \right\},$$

Computation of lepton PDFs



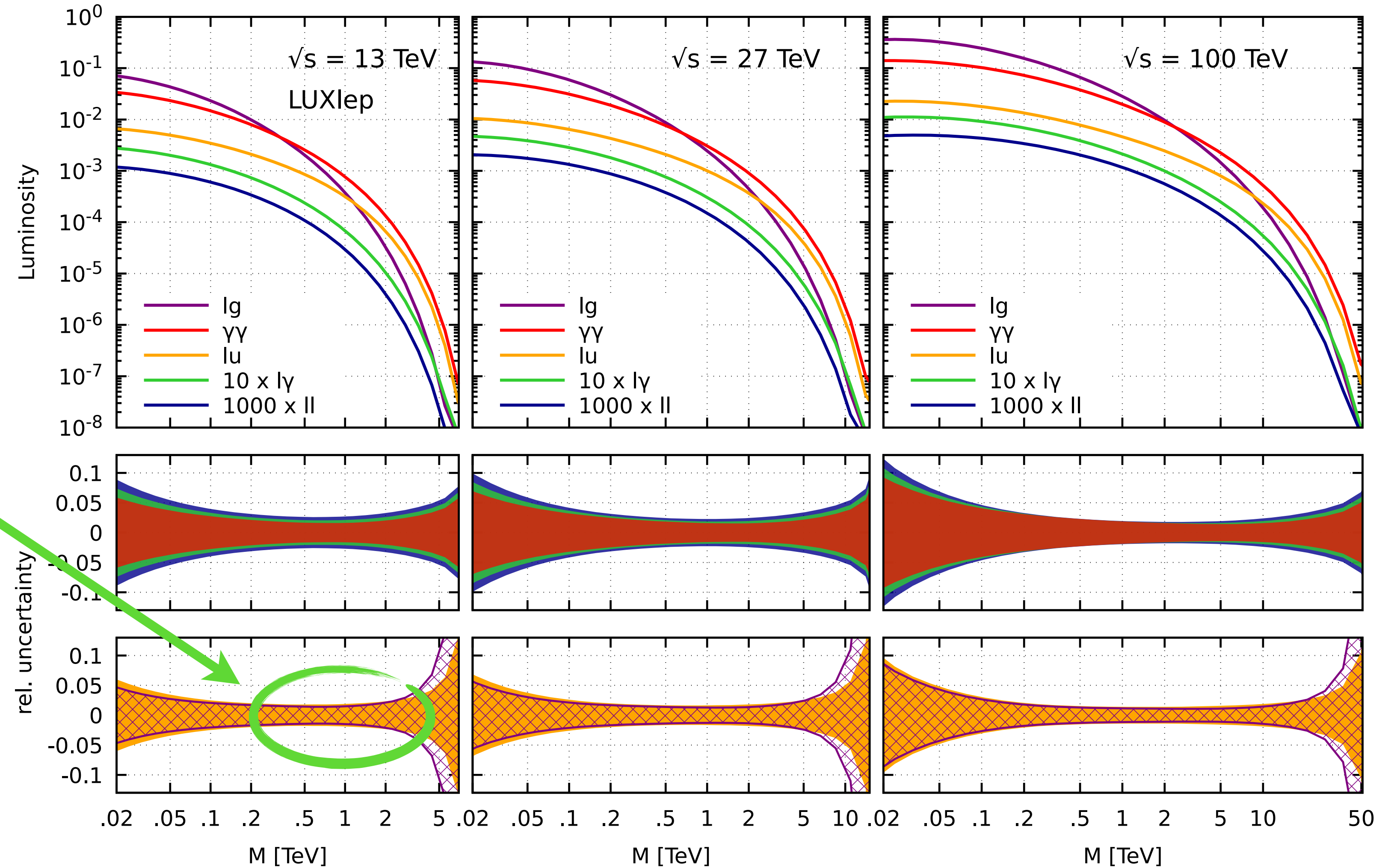
lepton PDF

$$x_l f_l(x_l, \mu_F^2) =$$

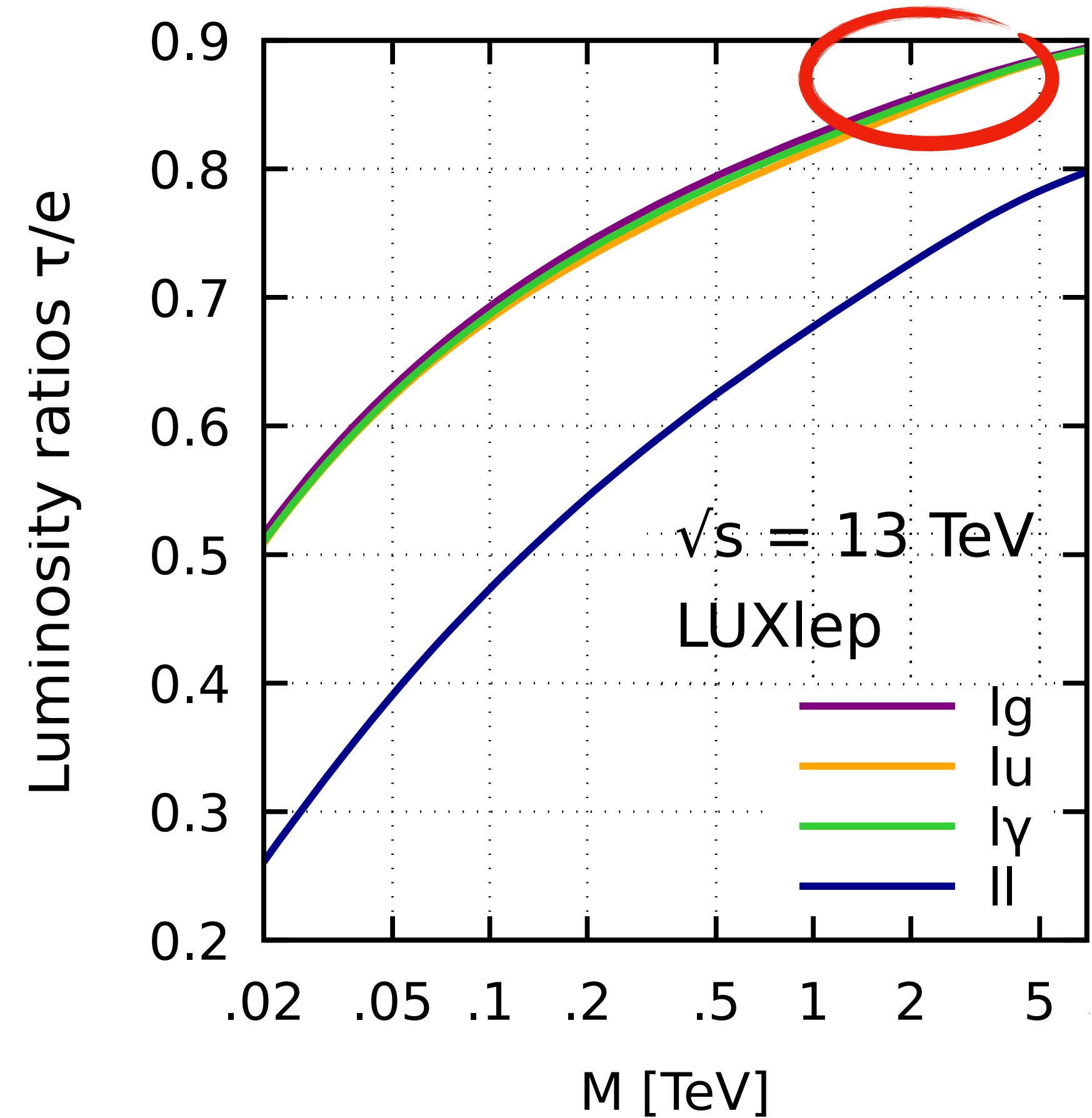
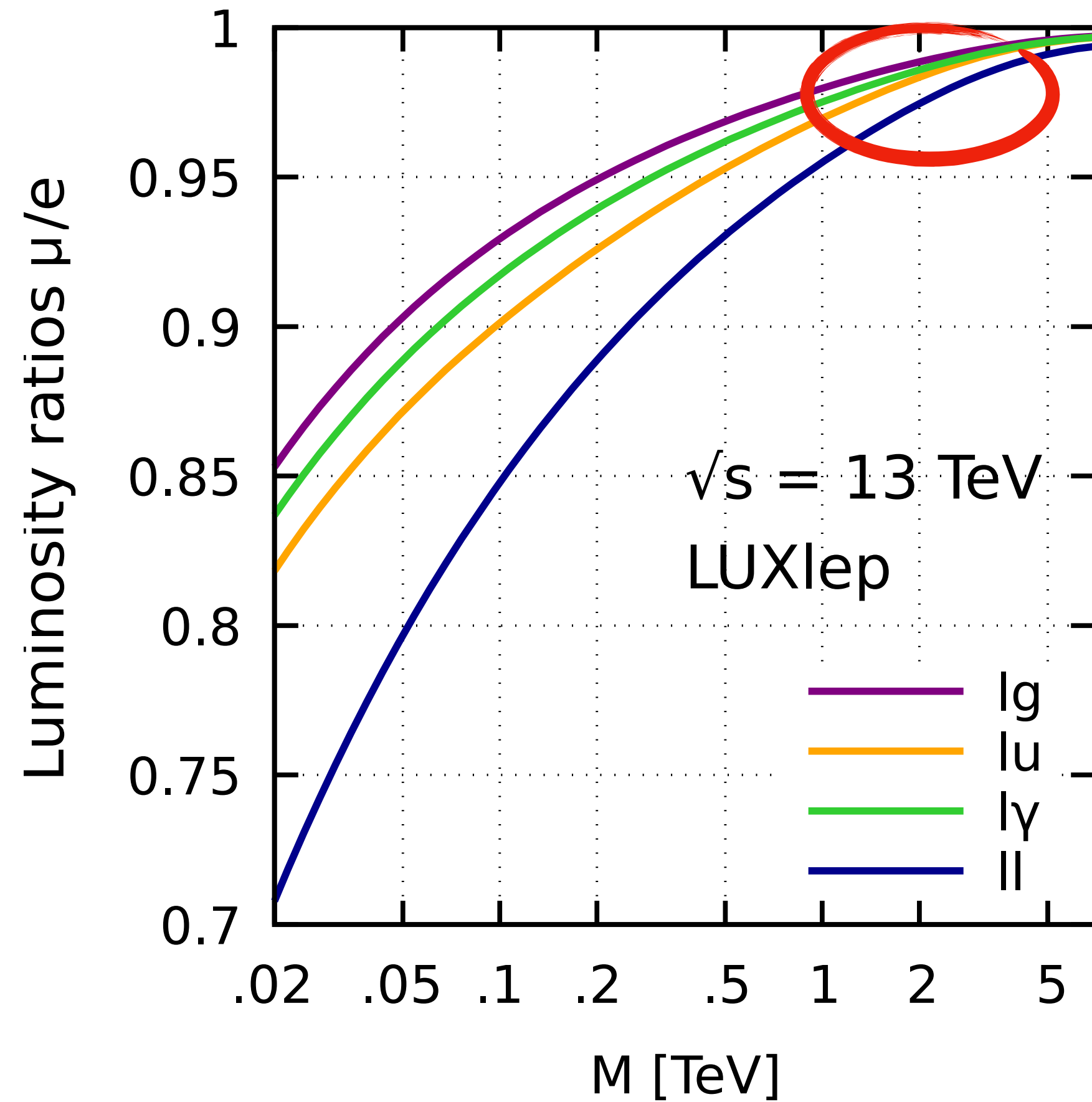
convolution involving photon PDF & $\gamma \rightarrow l+l^-$ splitting function +
convolution involving proton structure functions & splitting functions

Luminosities @ LHC & beyond

uncertainties
below 5% for a
wide range of
invariant masses

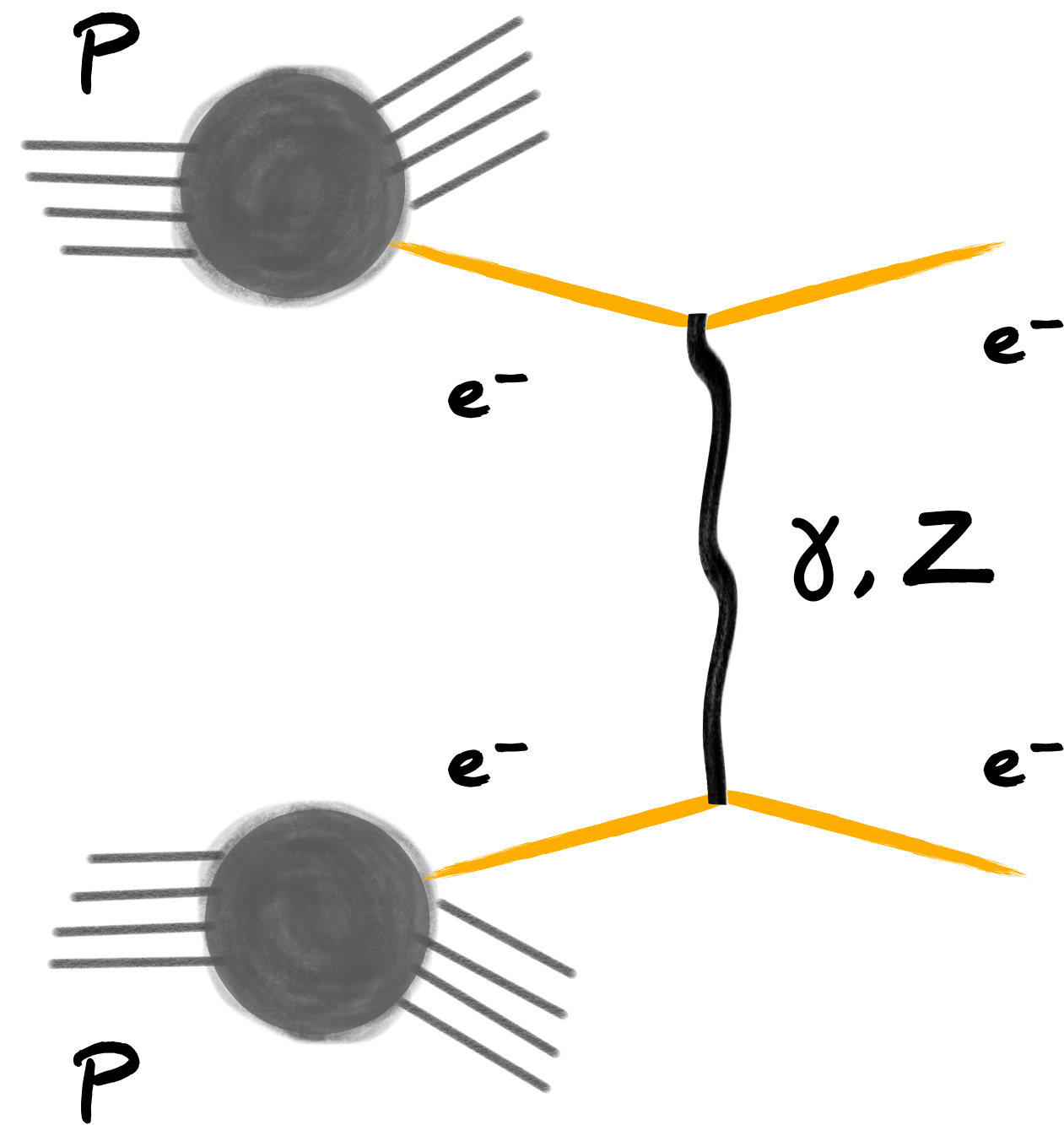


Luminosities ratios



For TeV-scale resonances, μ/e (τ/e) luminosity ratio above 95% (around 85%)

Same sign lepton-pair production @ LHC



Signal events after cuts:

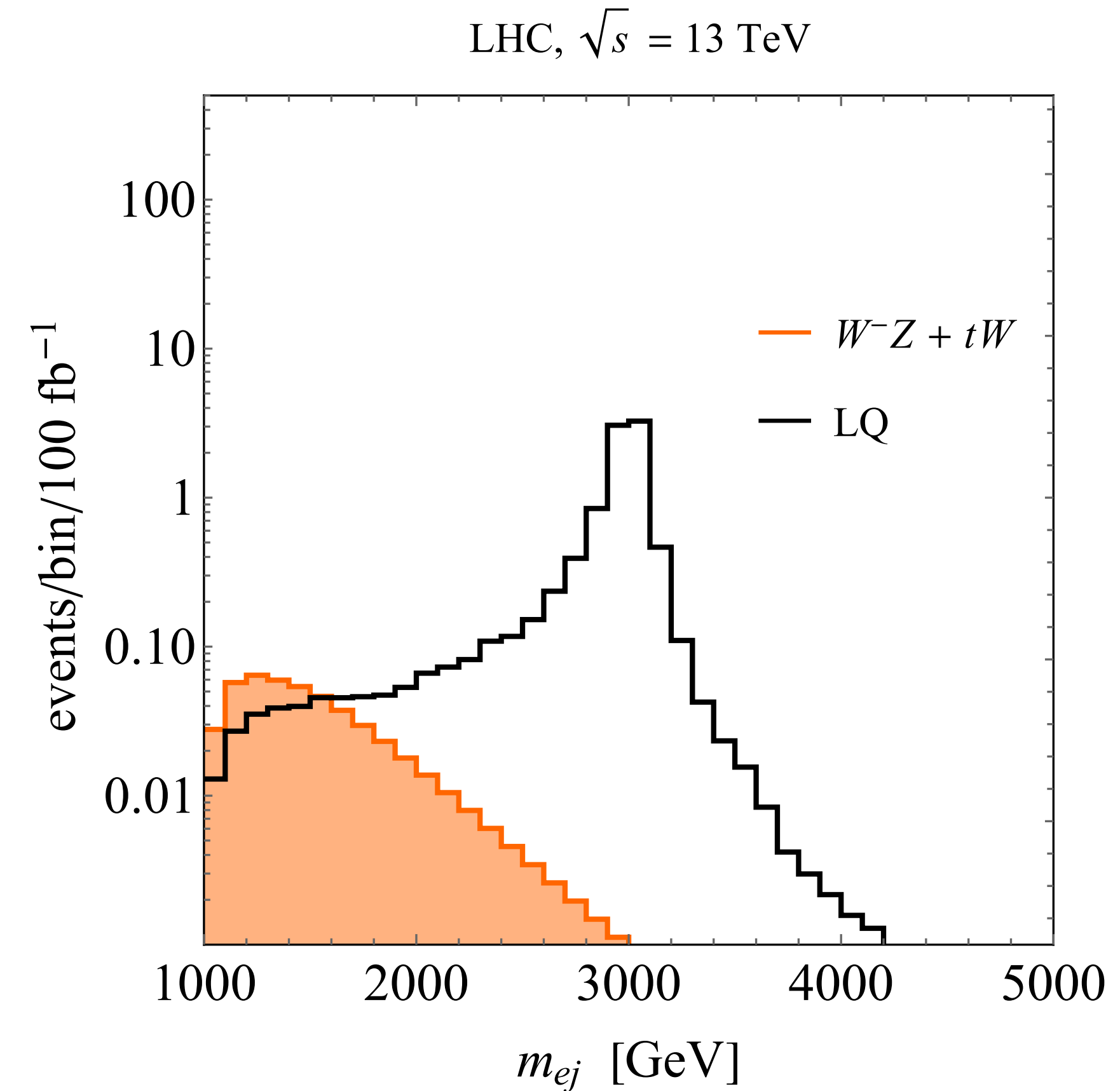
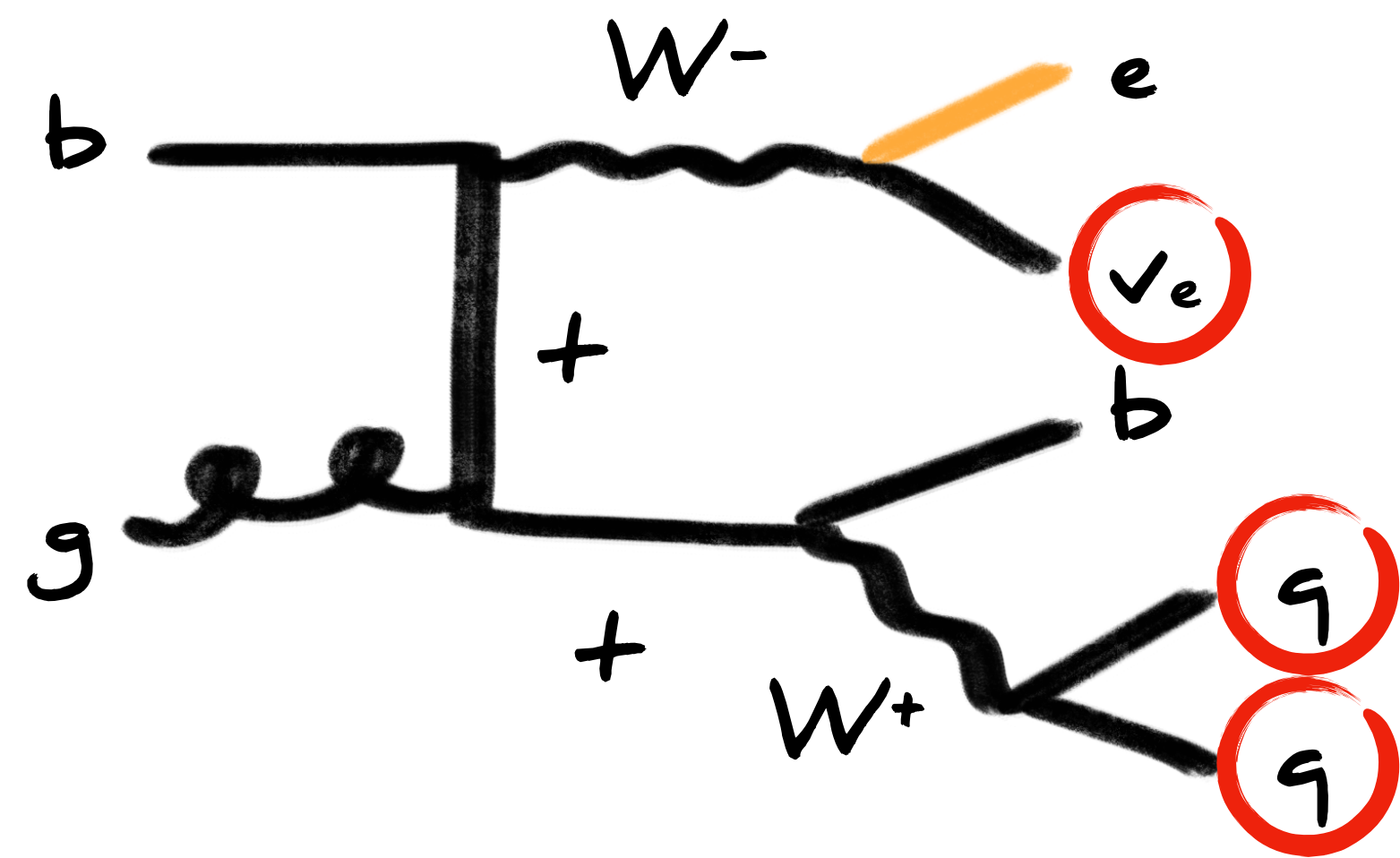
$$N_{\text{HL-LHC}}(e^-e^-) \simeq 700,$$

$$N_{\text{HL-LHC}}(\mu^-\mu^-) \simeq 550,$$

$$N_{\text{HL-LHC}}(\tau^-\tau^-) \simeq 250$$

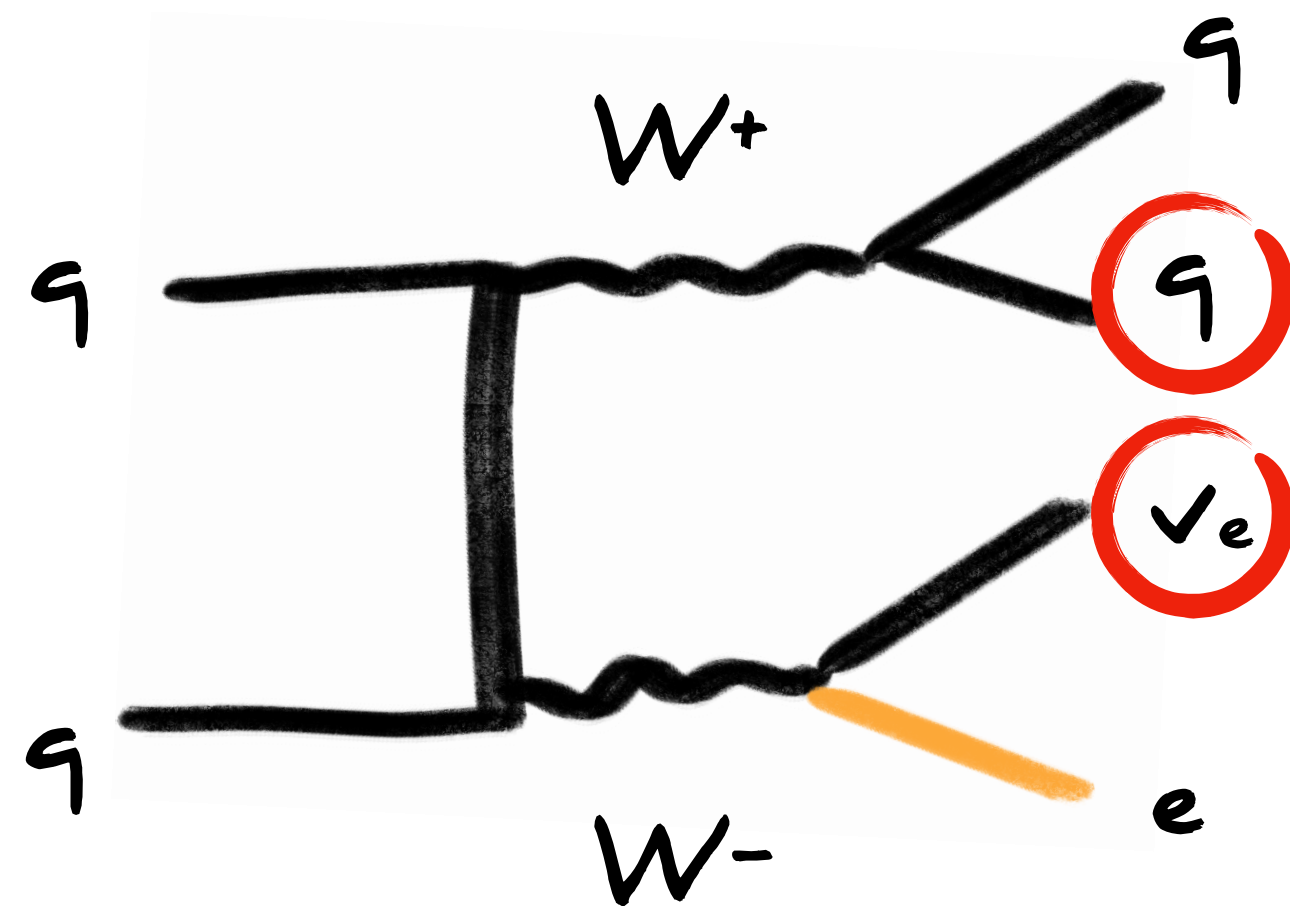
Dominant SM background from W-W-
production after same cuts close to 0

Resonant LQ production @ the LHC



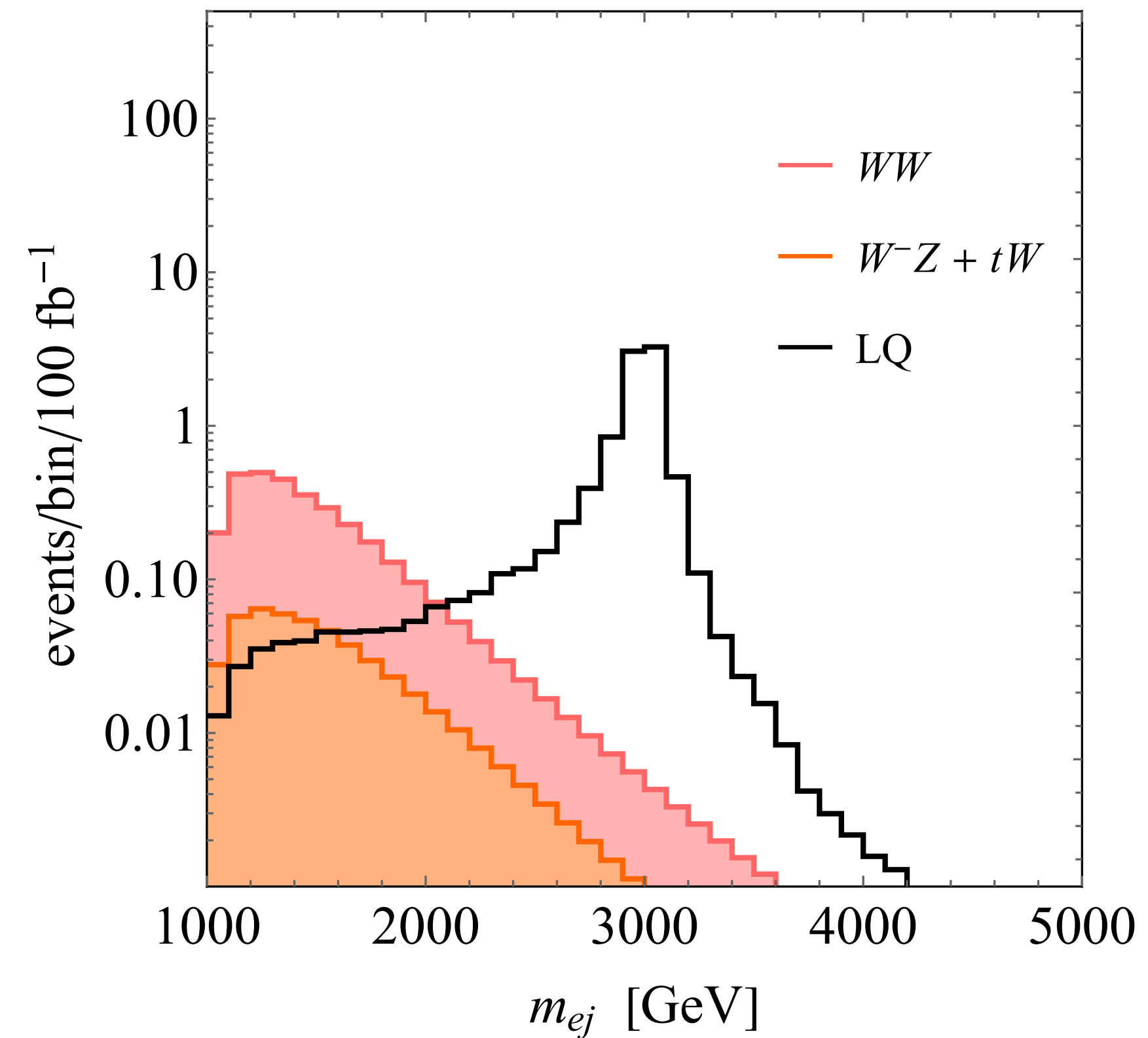
suppressed by $E_{T,miss}$ requirement & jet veto

Resonant LQ production @ the LHC

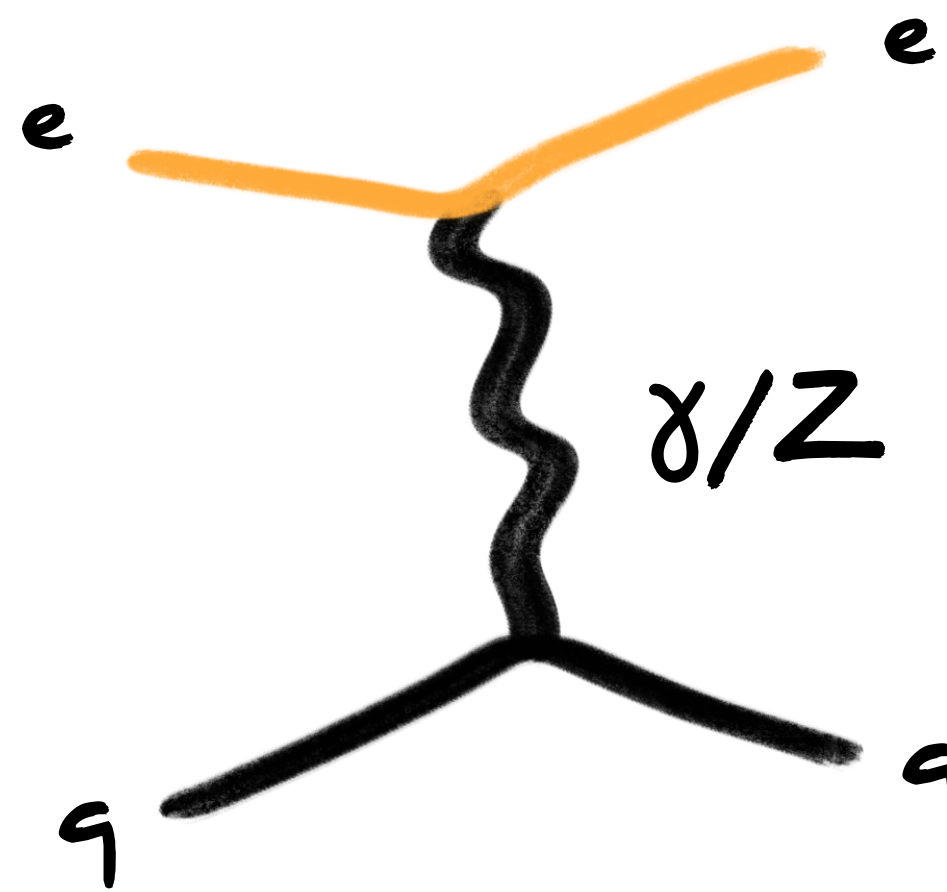


suppressed by $E_{T,miss}$ requirement & jet veto

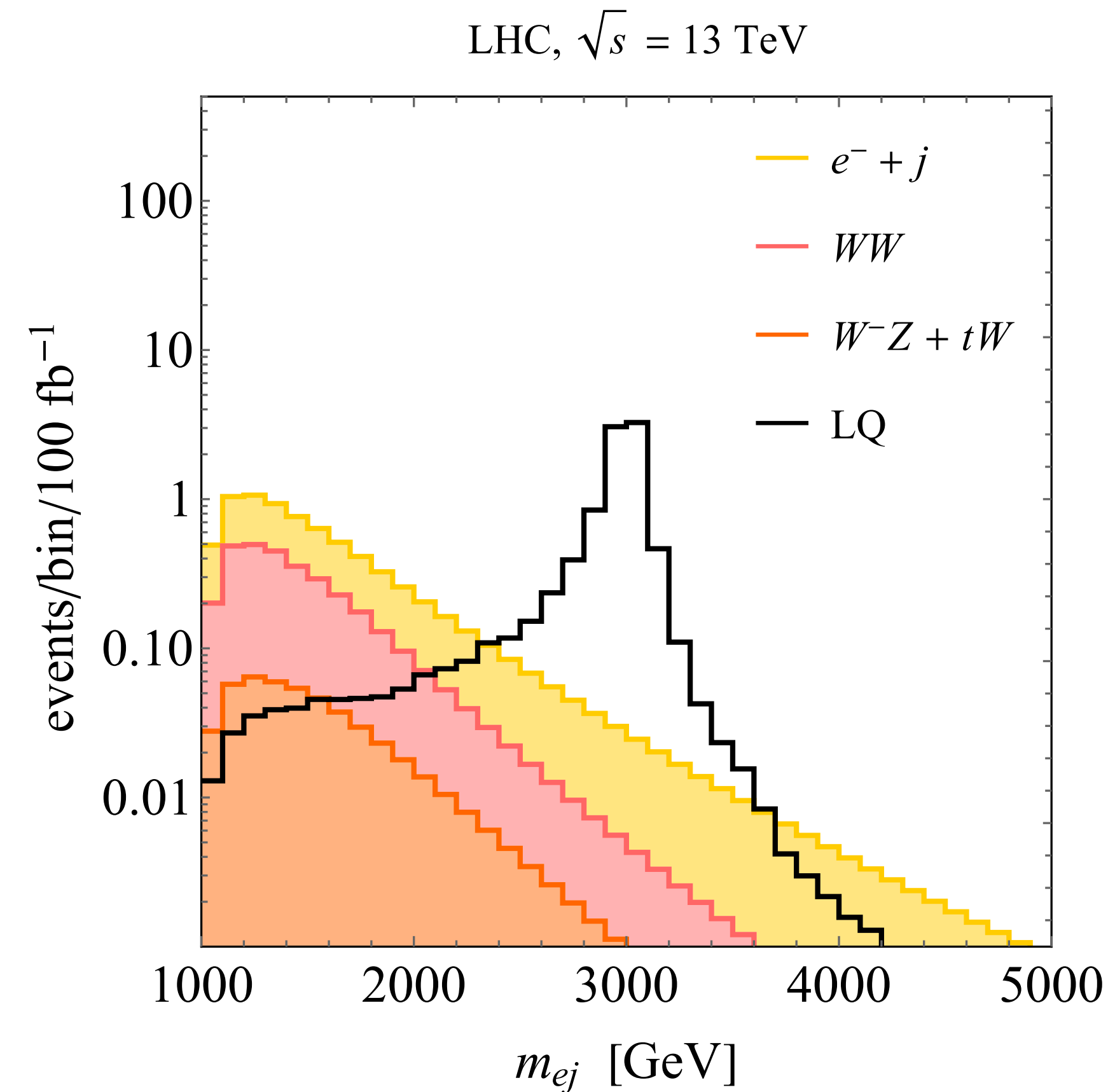
LHC, $\sqrt{s} = 13$ TeV



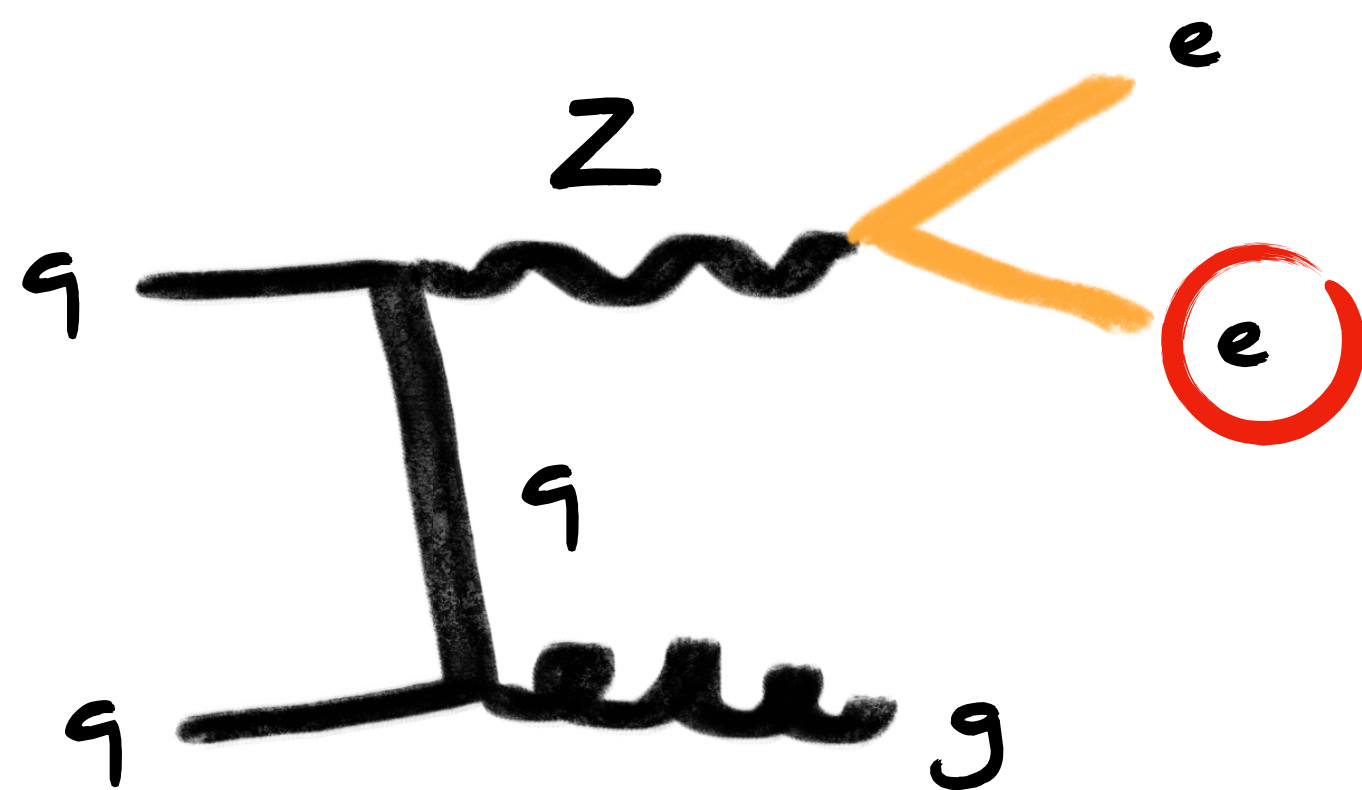
Resonant LQ production @ the LHC



irreducible background particularly relevant
@ high invariant lepton-jet mass

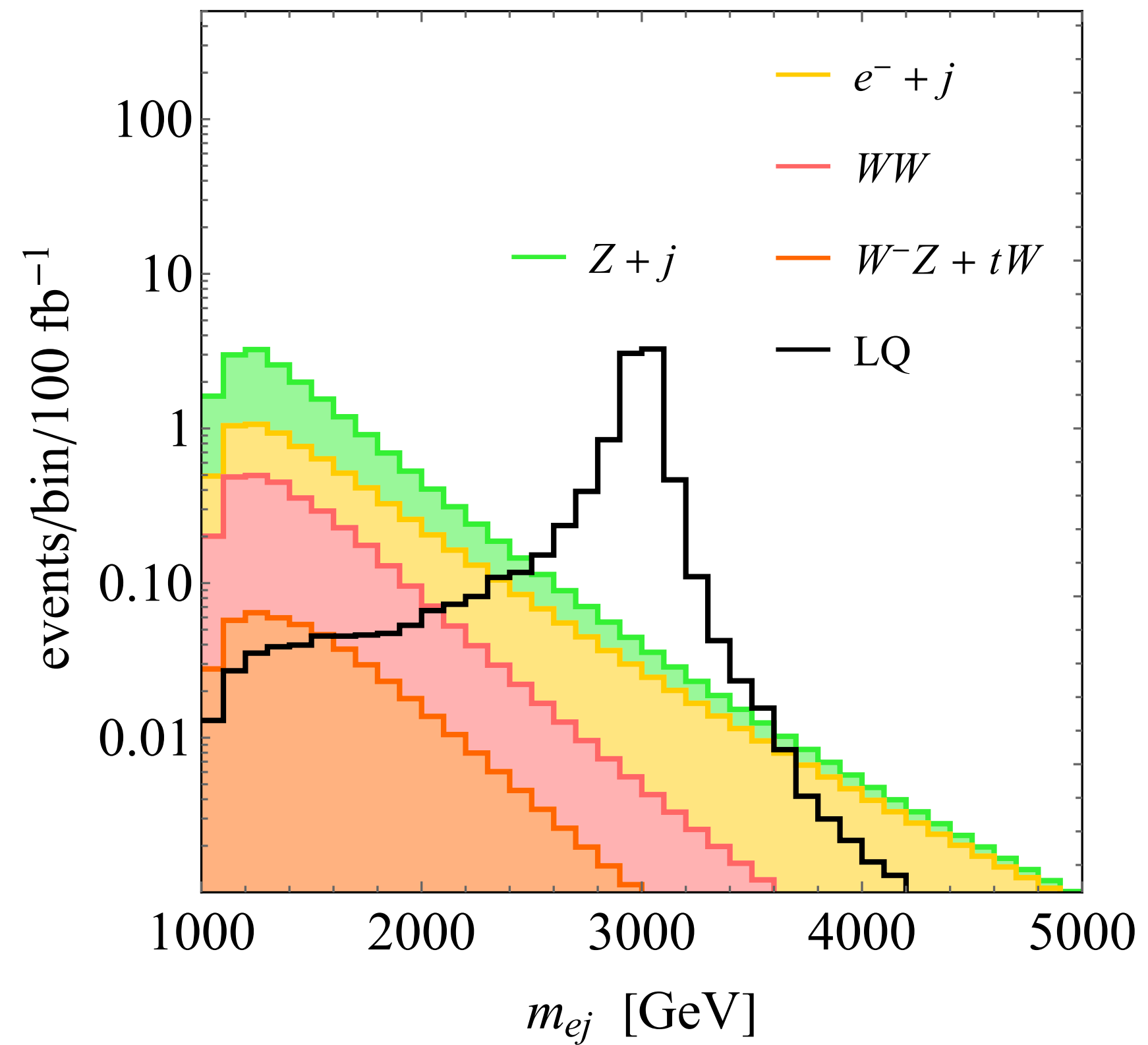


Resonant LQ production @ the LHC

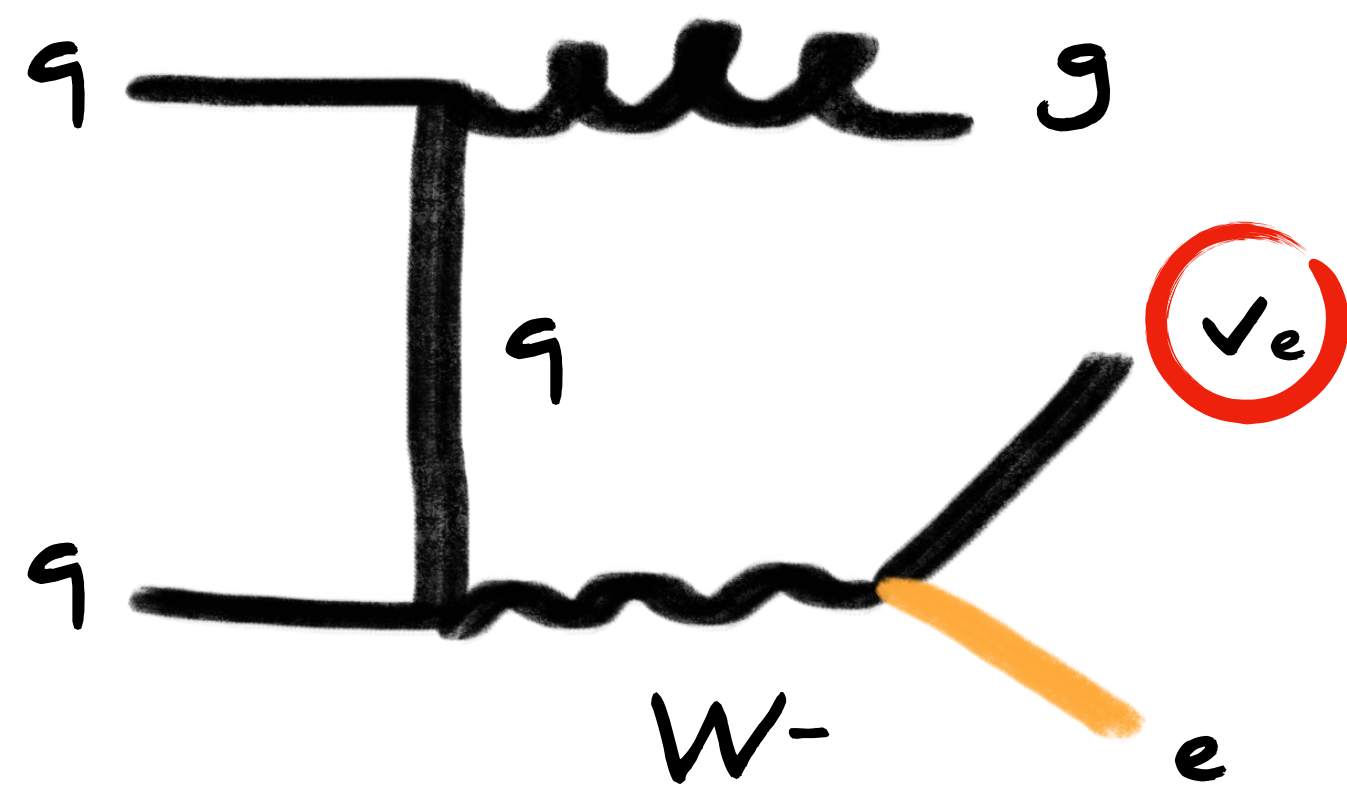


suppressed by lepton veto

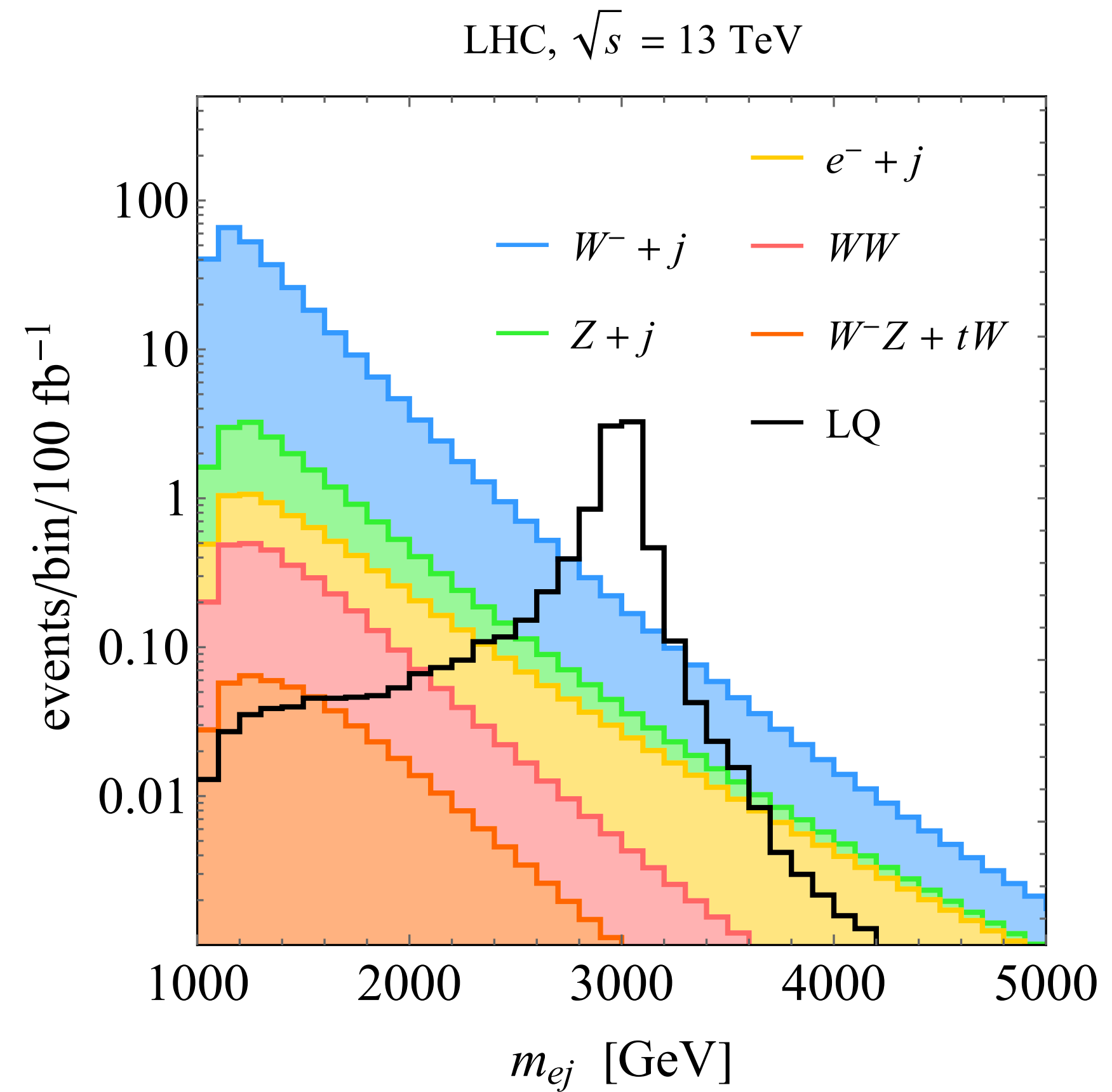
LHC, $\sqrt{s} = 13$ TeV



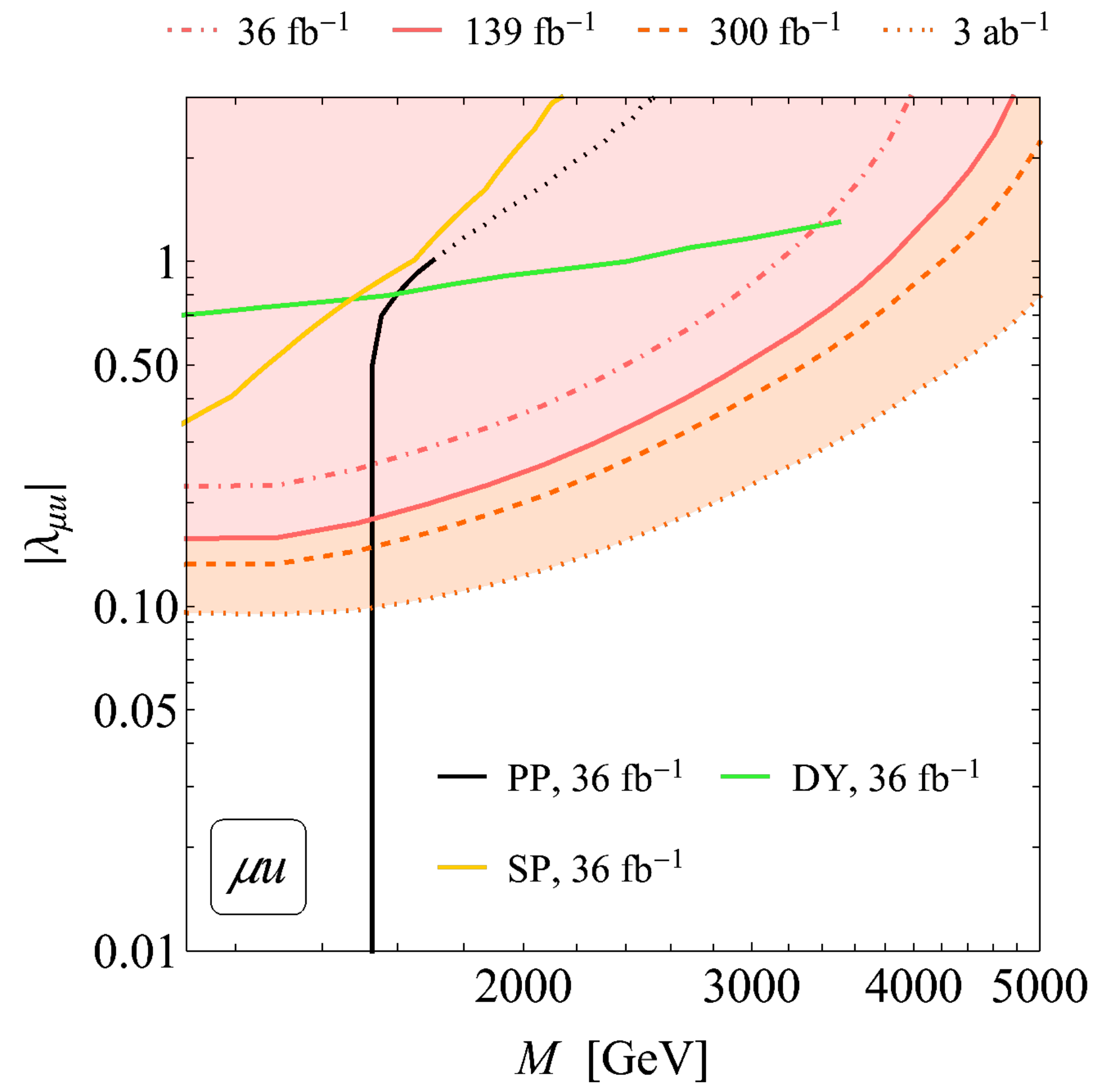
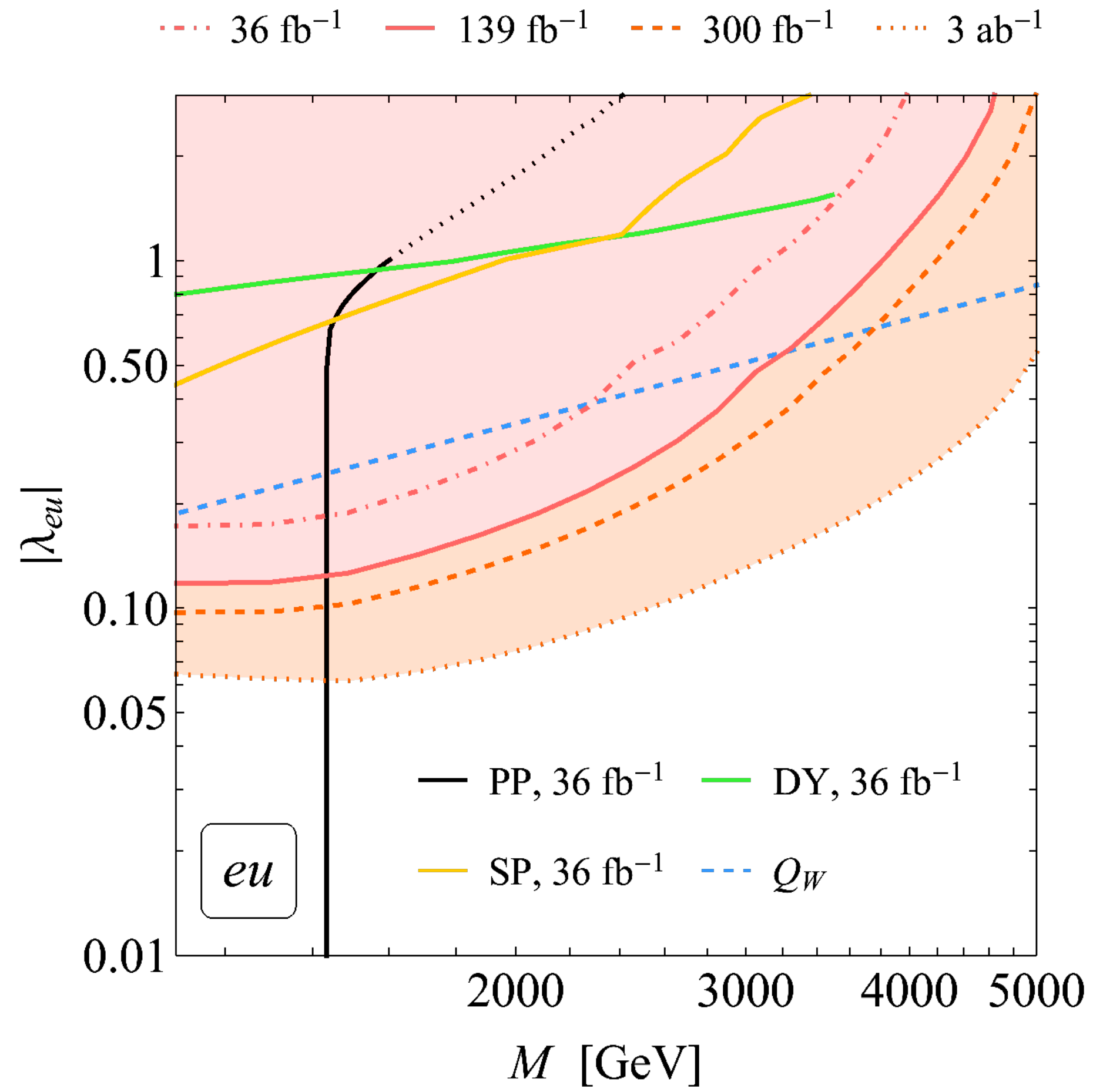
Resonant LQ production @ the LHC



suppressed by $E_{T,miss}$ requirement

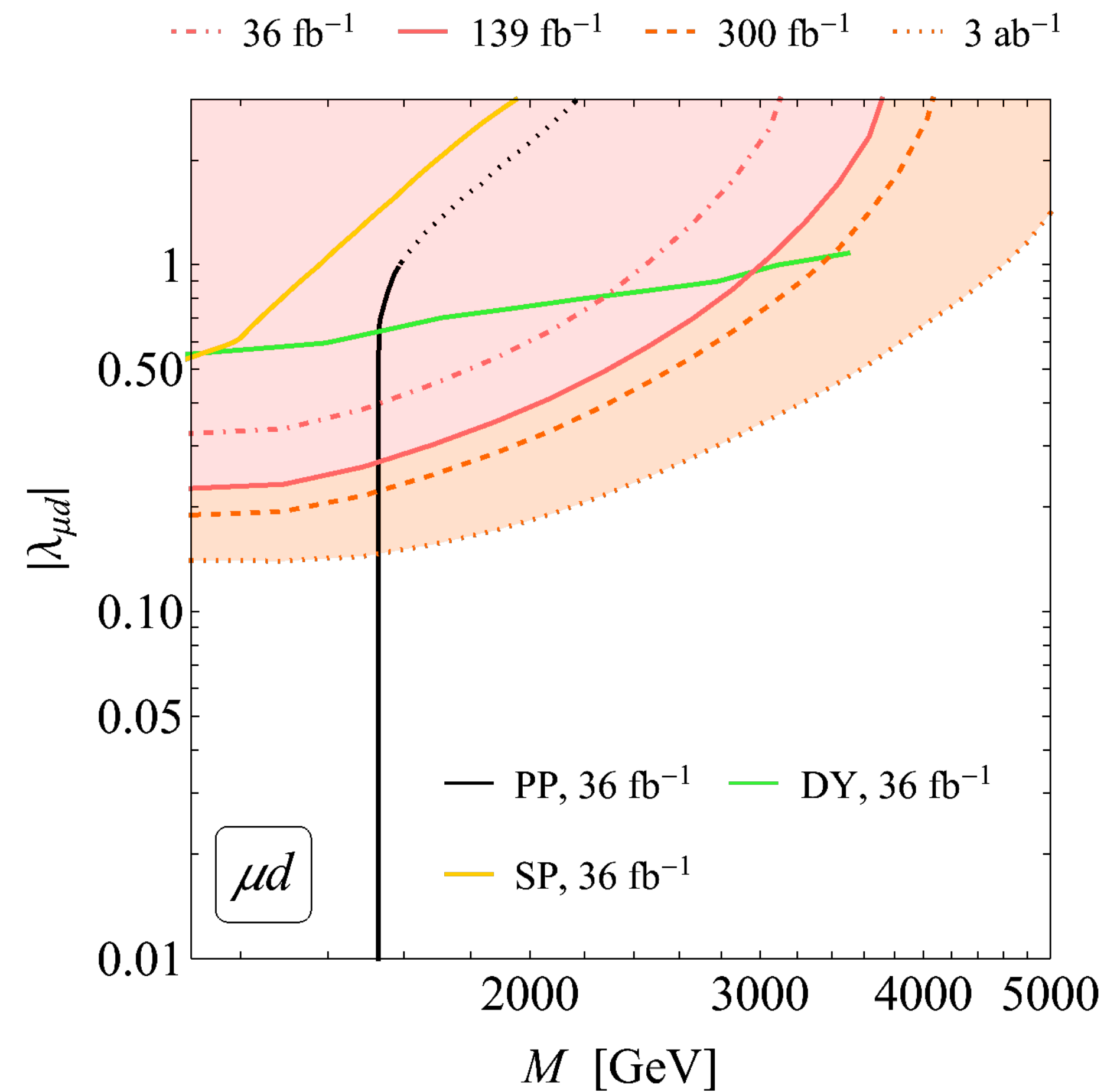
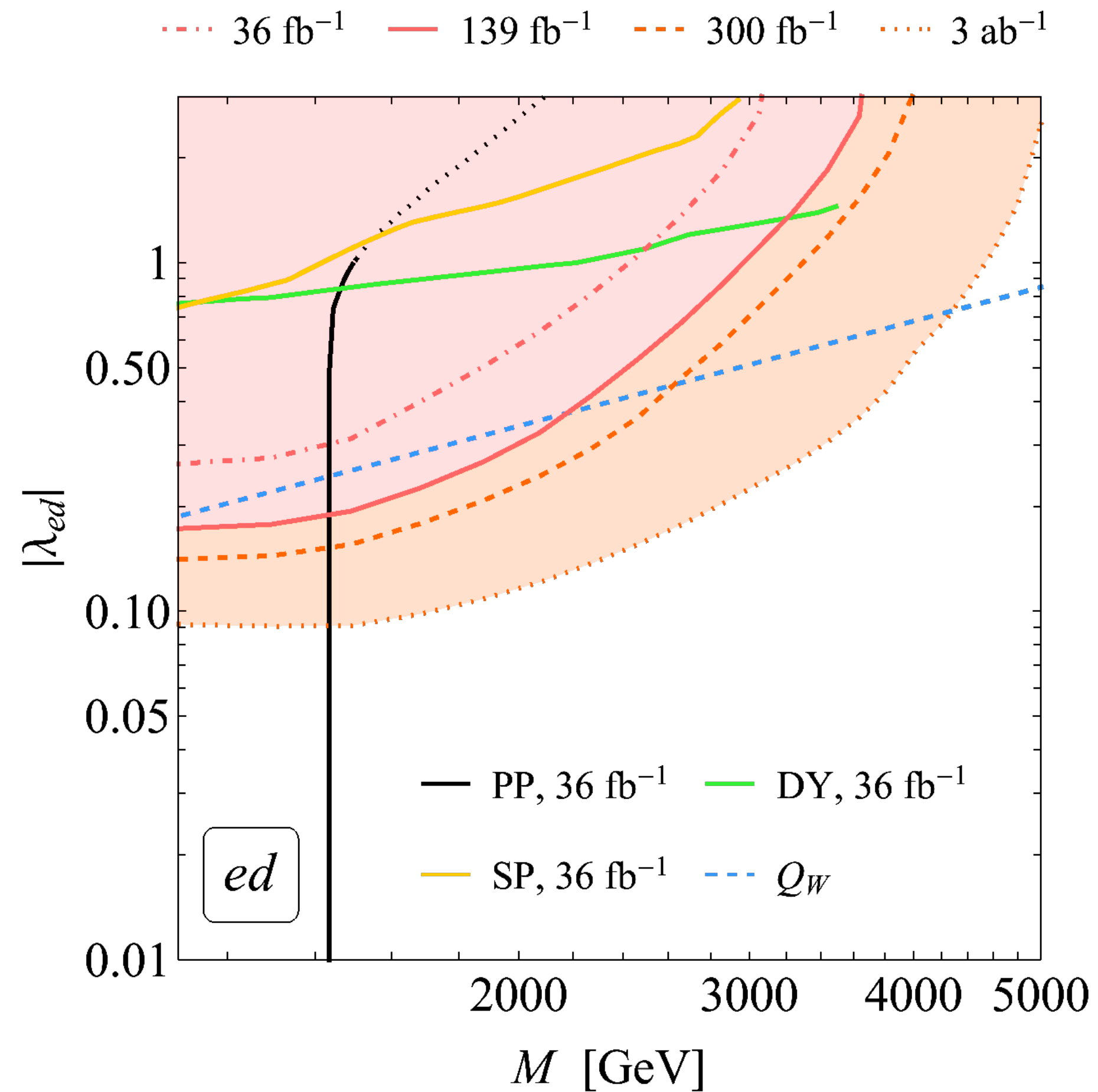


LHC limits on 1st & 2nd generation LQs

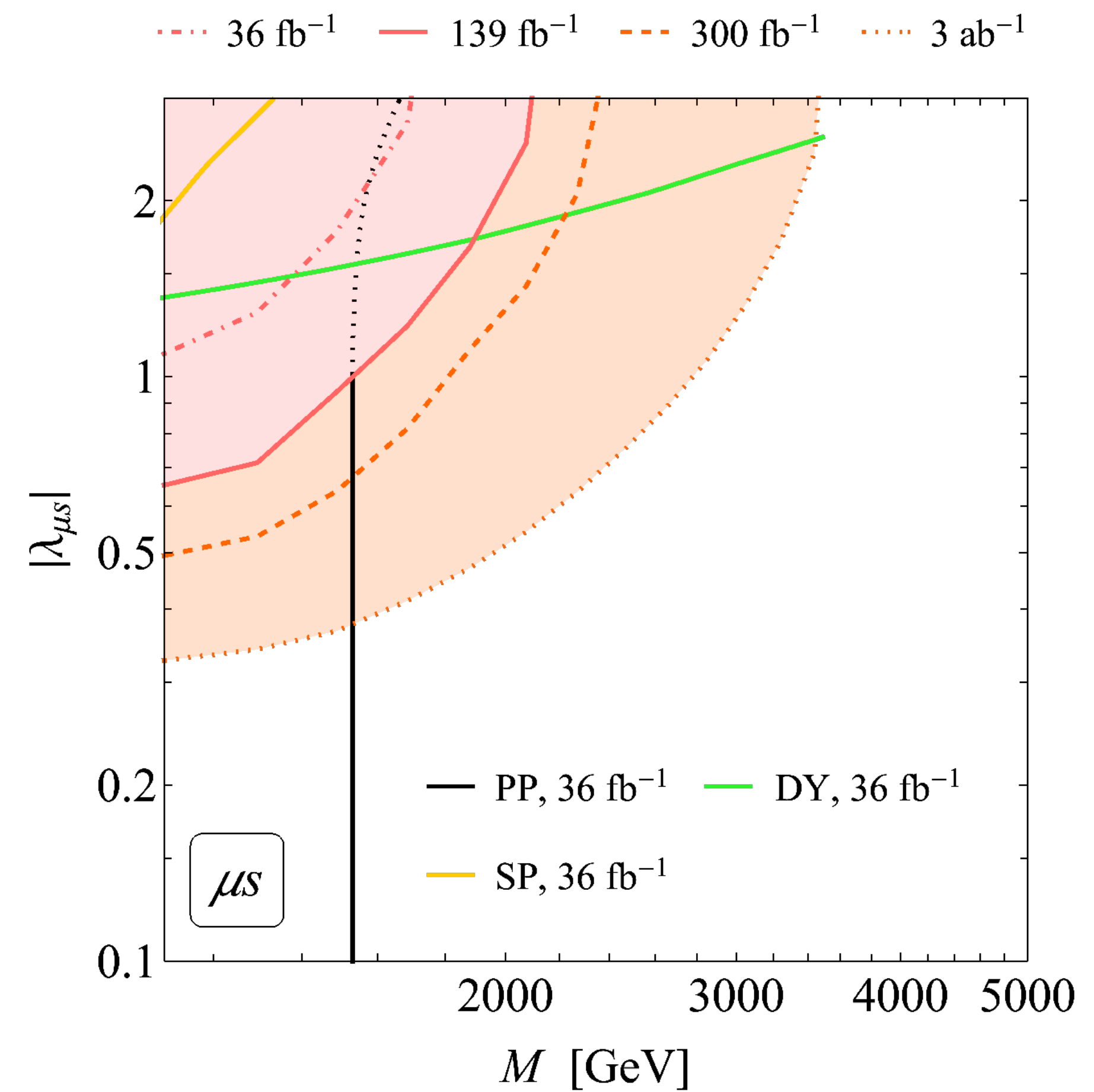
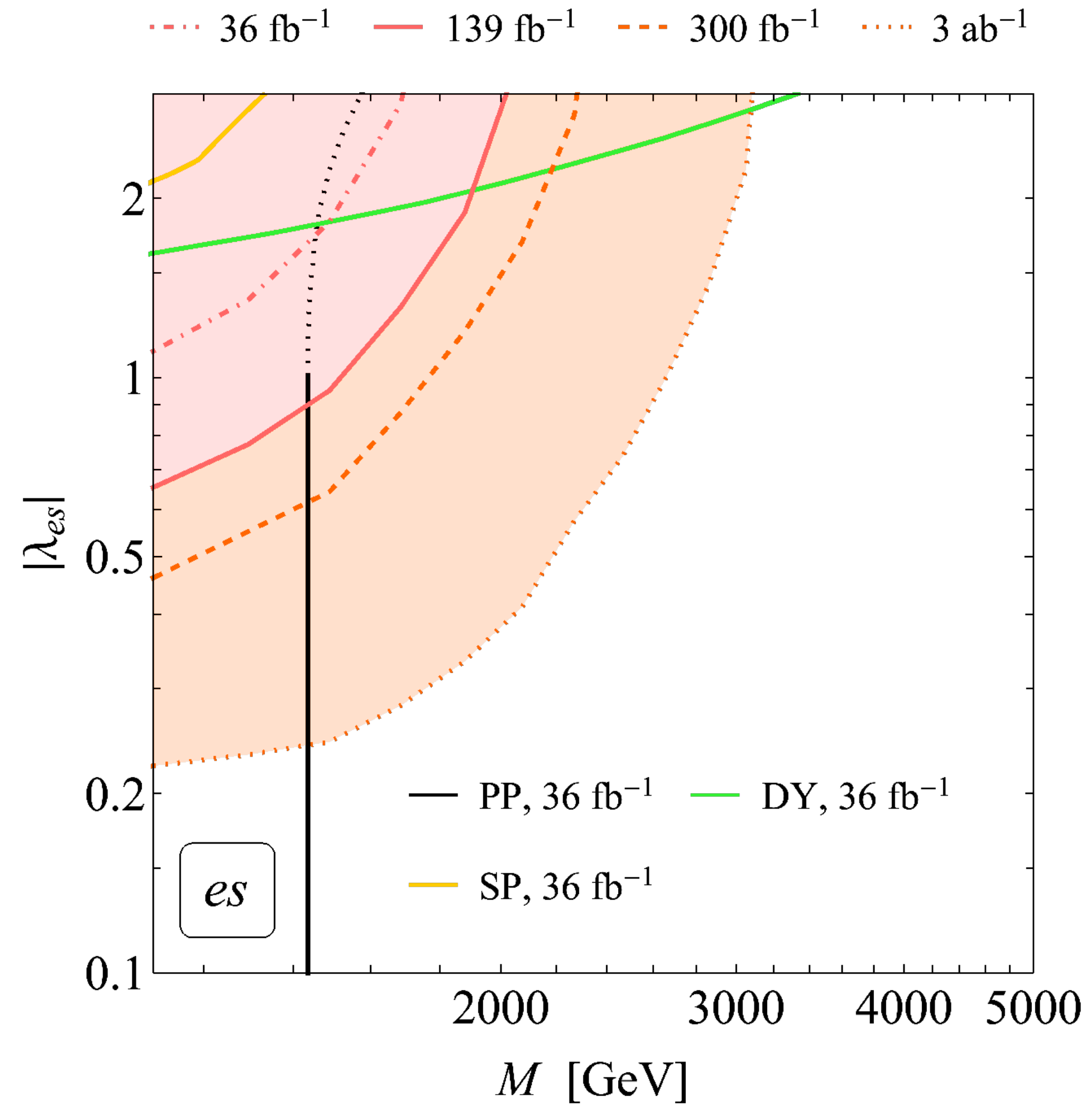


[Buonocore, UH, Nason, Tramontano & Zanderighi, 2005.06475]

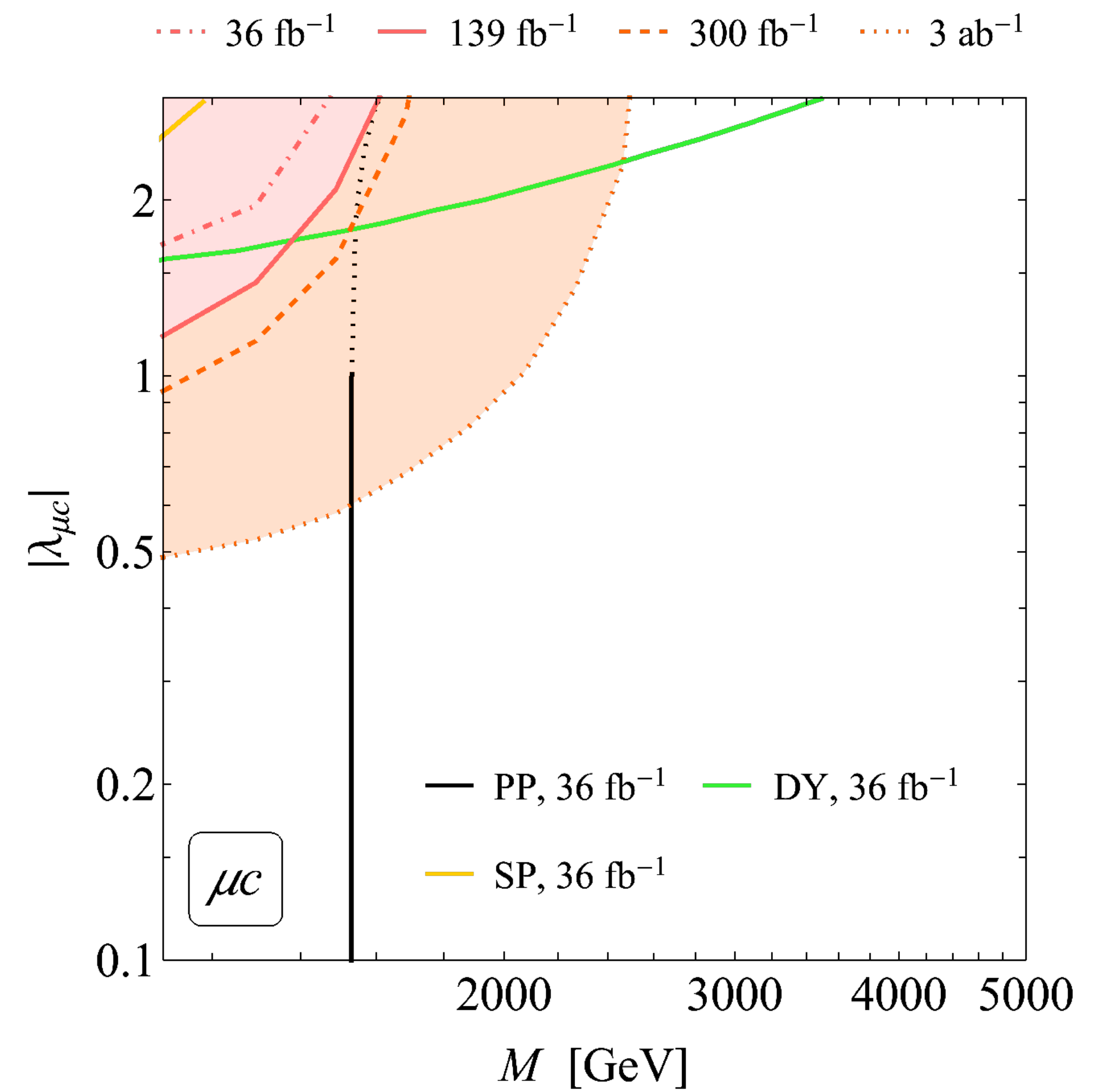
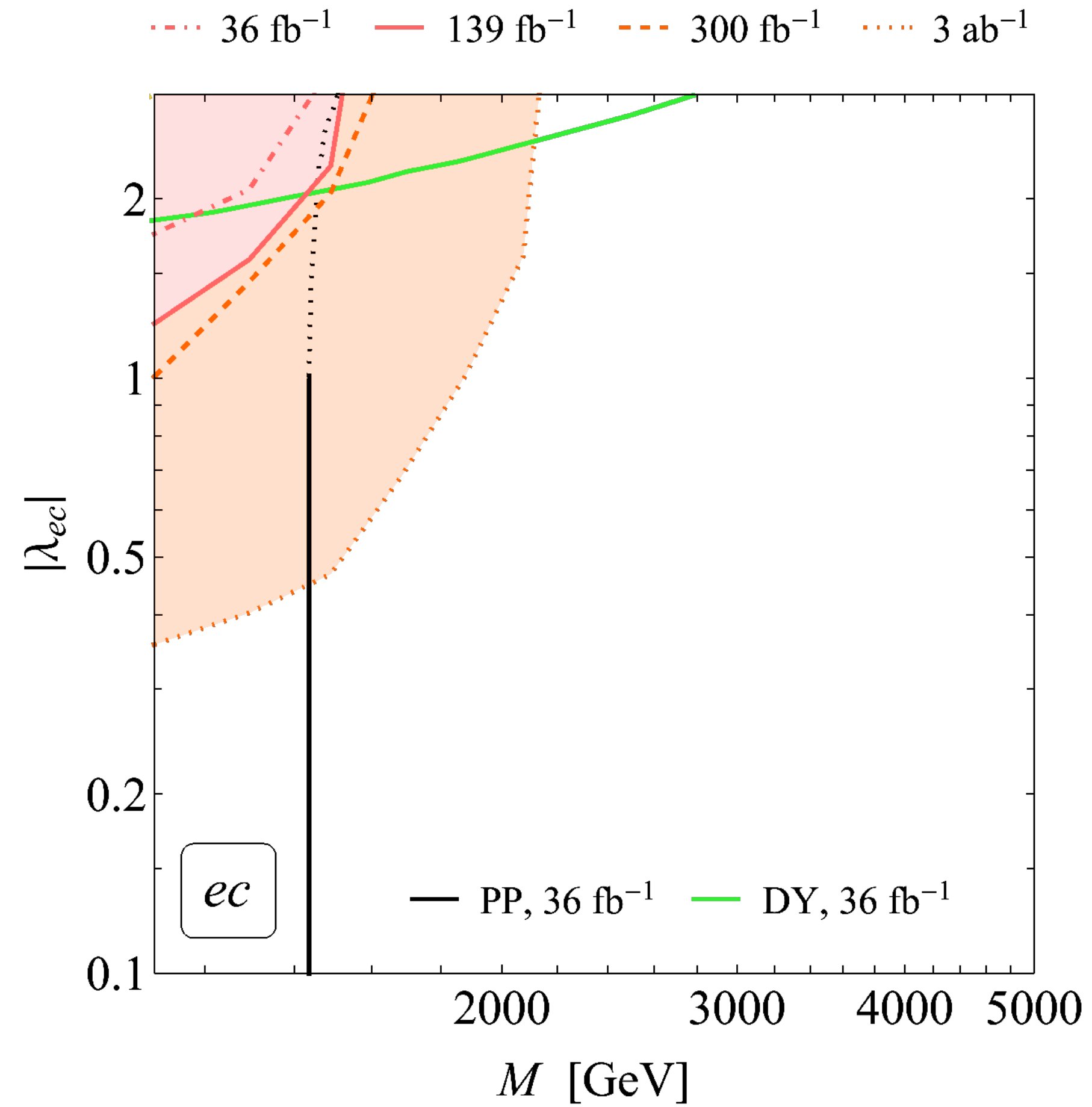
LHC limits on 1st & 2nd generation LQs



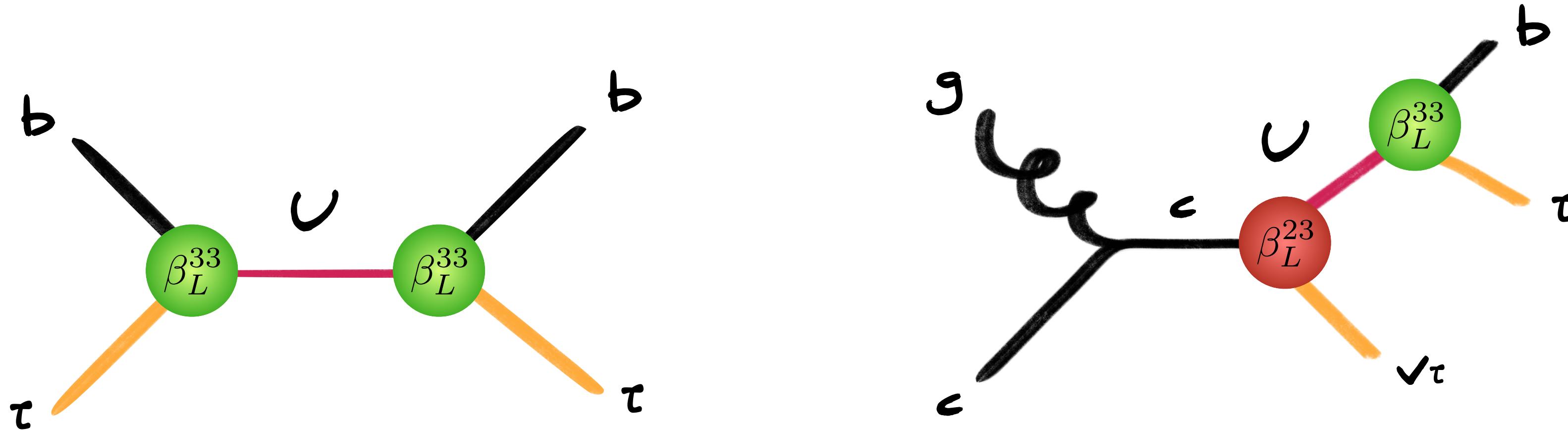
LHC limits on 1st & 2nd generation LQs



LHC limits on 1st & 2nd generation LQs



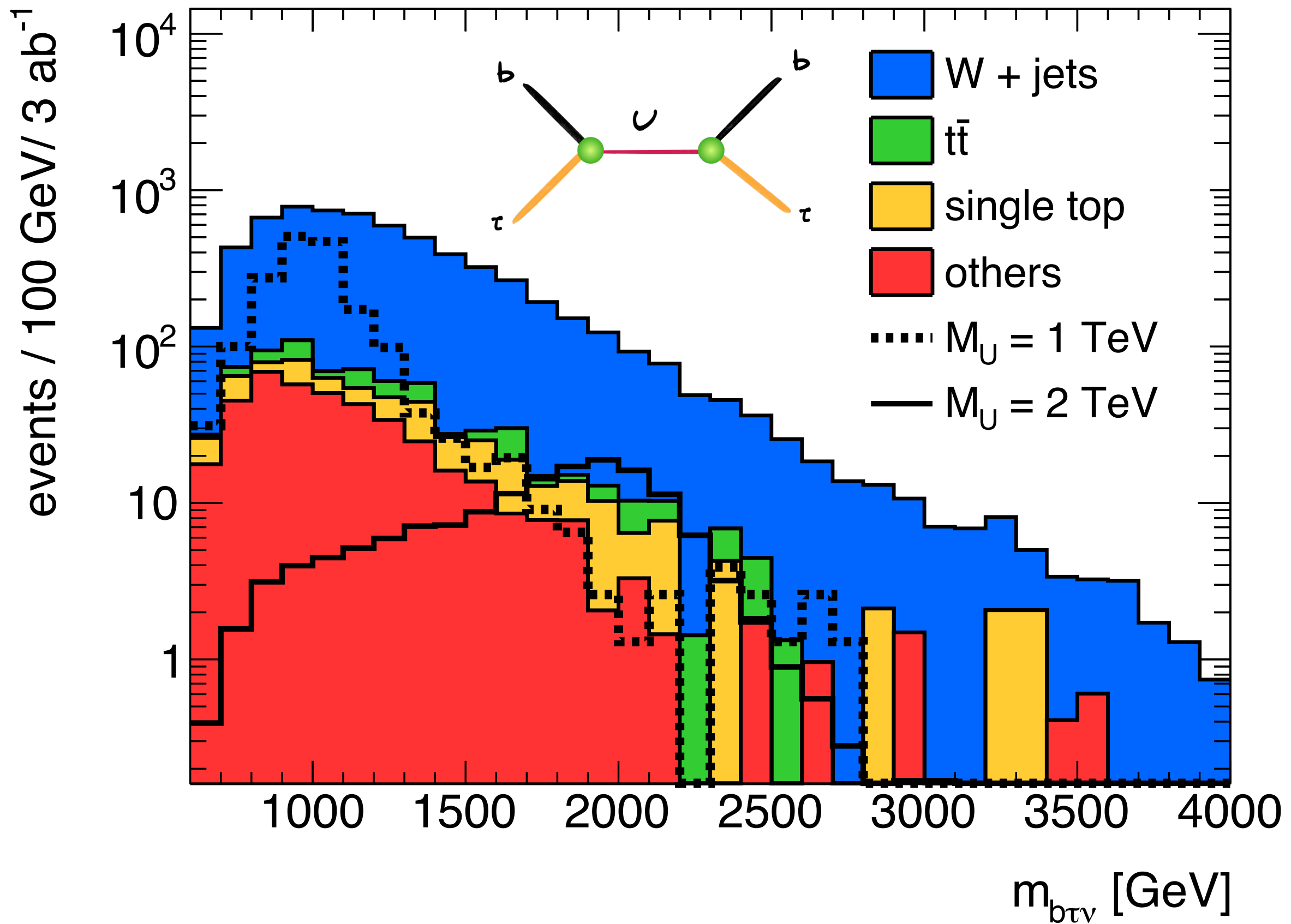
LQ contributions to $pp \rightarrow b\tau$ signature



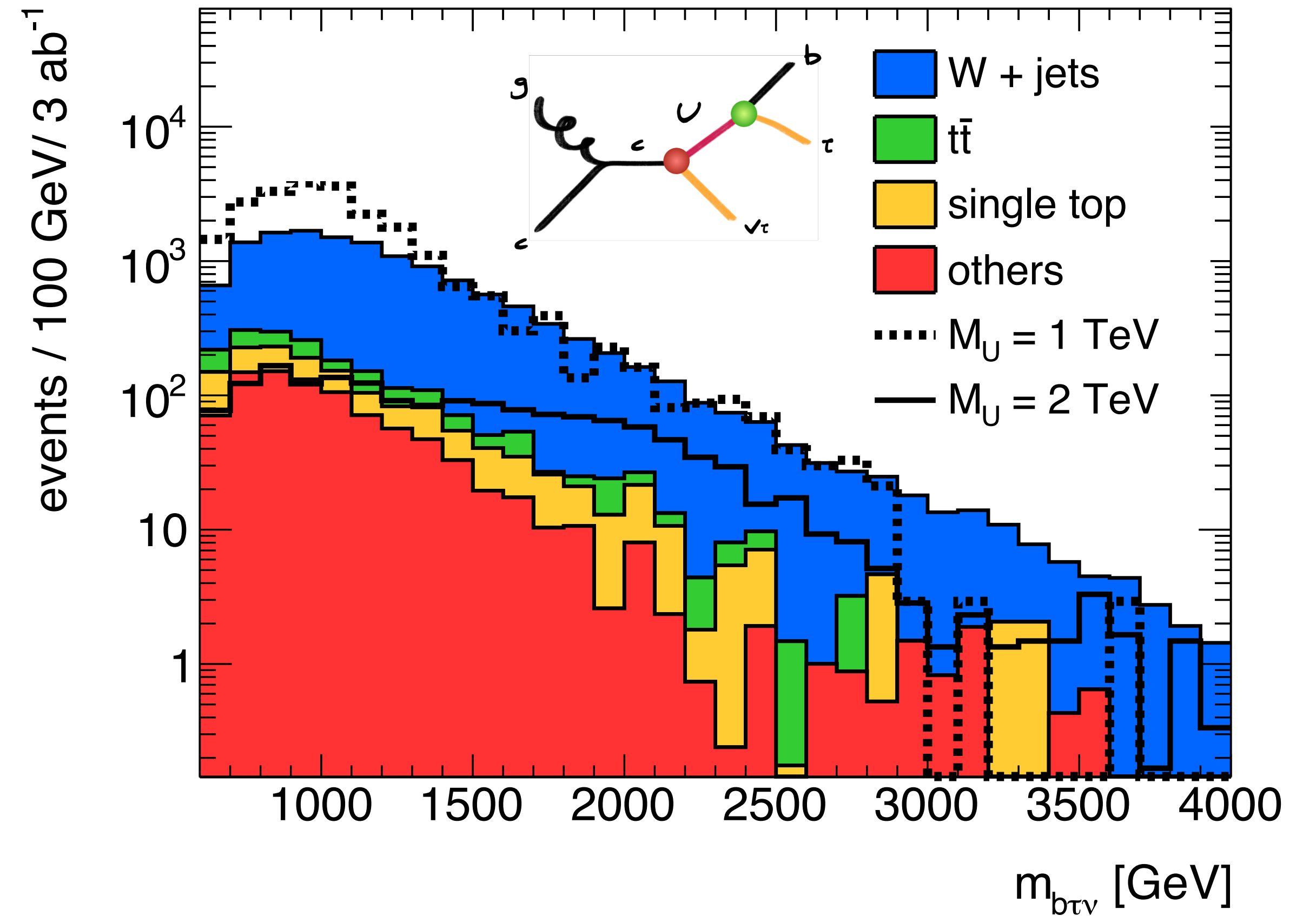
For $\beta_L^{23} = 0$, $b\tau$ signal arises only from $2 \rightarrow 2$ process, while for $\beta_L^{23} \neq 0$ also $2 \rightarrow 3$ scattering is relevant. Since two topologies lead to final states with very different kinematic features, it is essential to develop two separate search strategies for them

Kinematic distributions of $pp \rightarrow b\tau$ signal

LHC 14 TeV, $b + \tau$

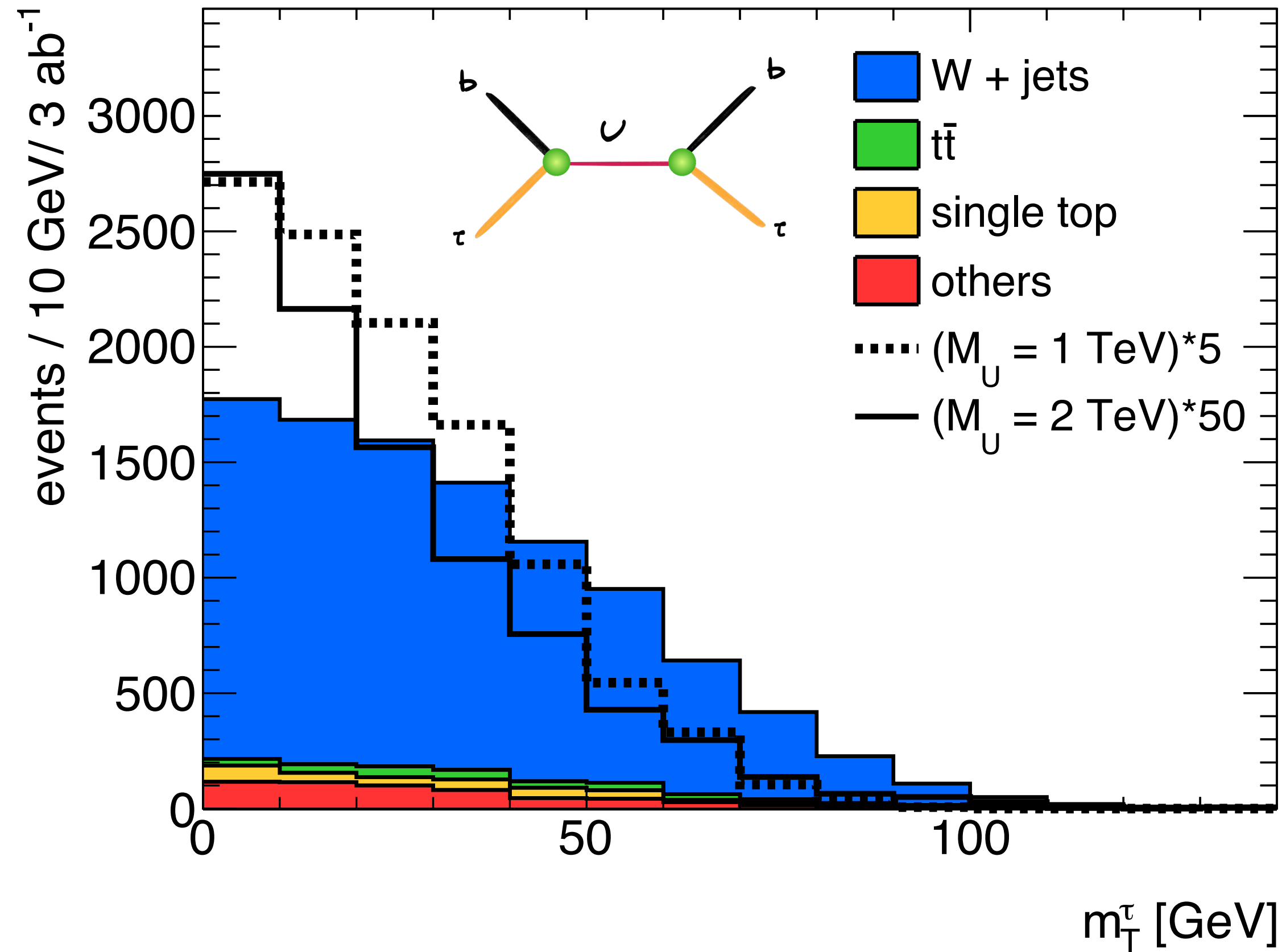


LHC 14 TeV, $b + \tau$

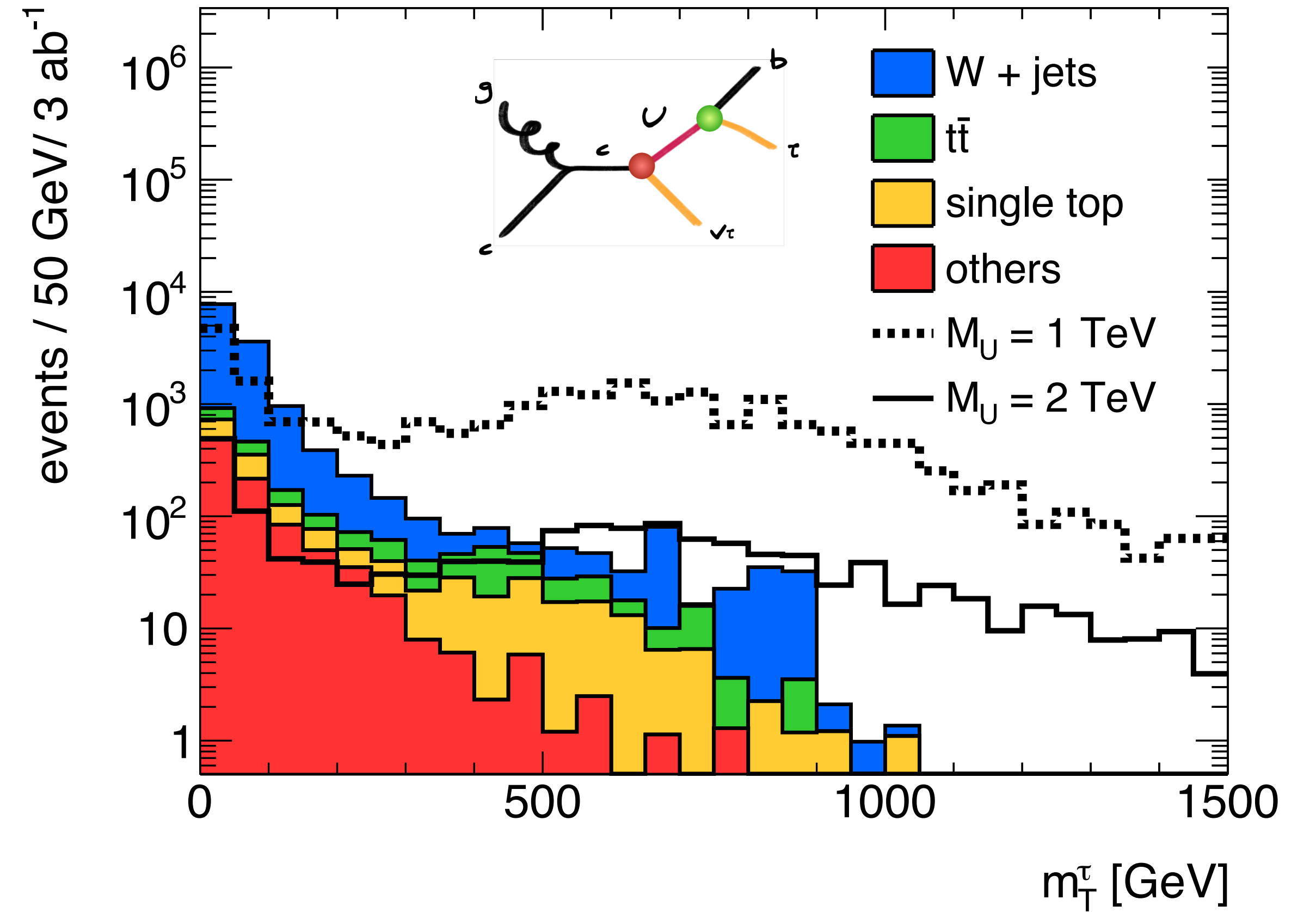


Kinematic distributions of $pp \rightarrow b\tau$ signal

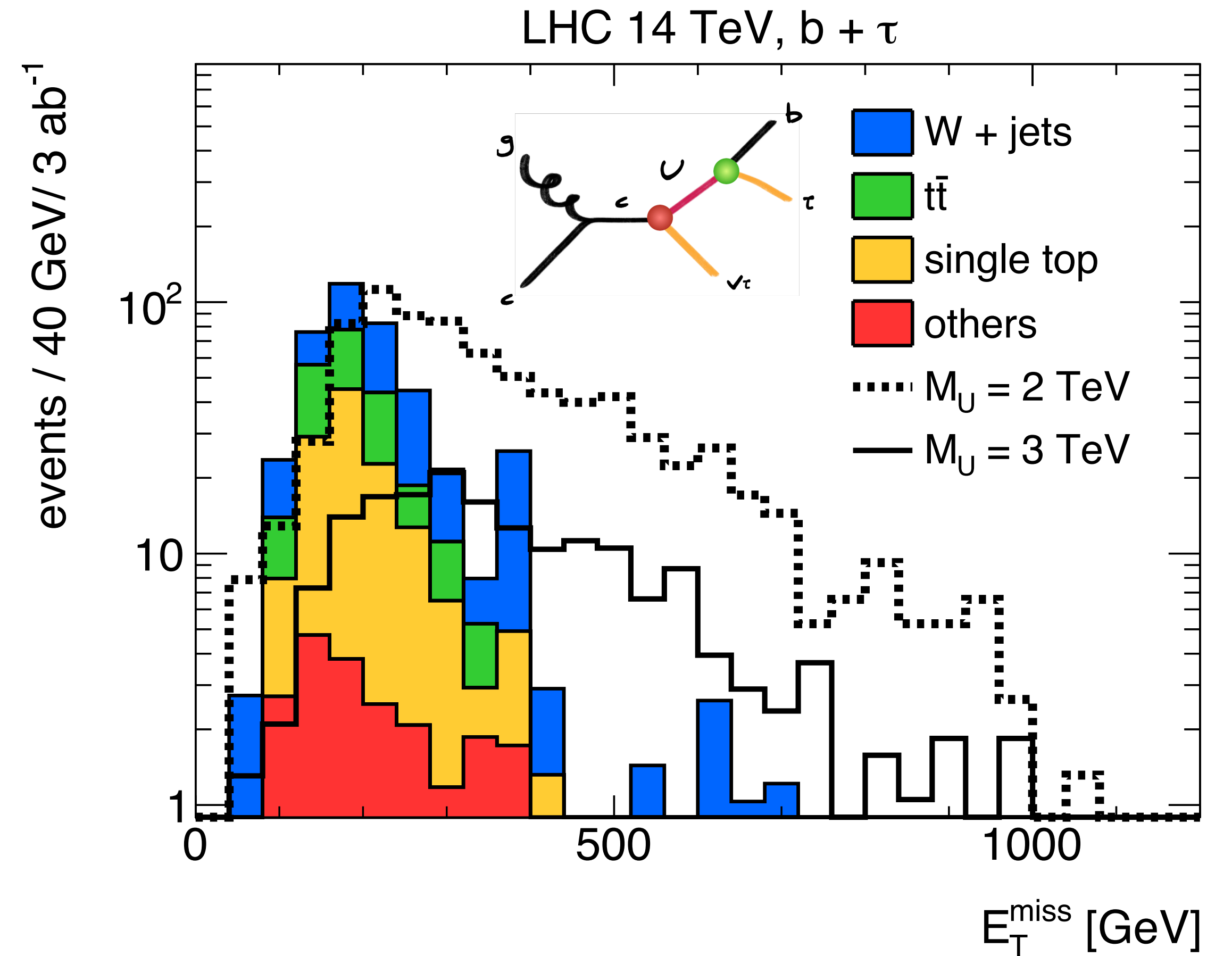
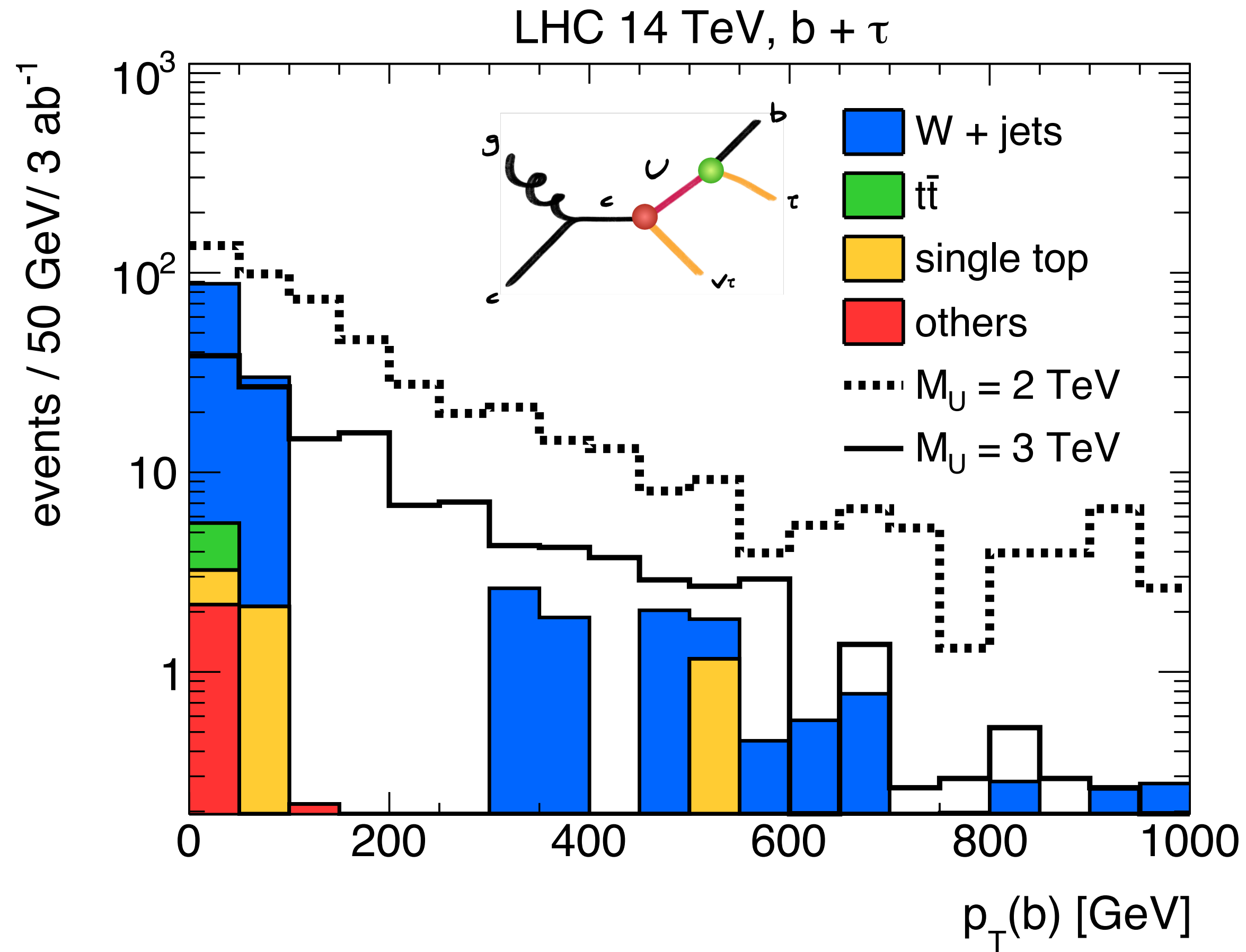
LHC 14 TeV, $b + \tau$



LHC 14 TeV, $b + \tau$

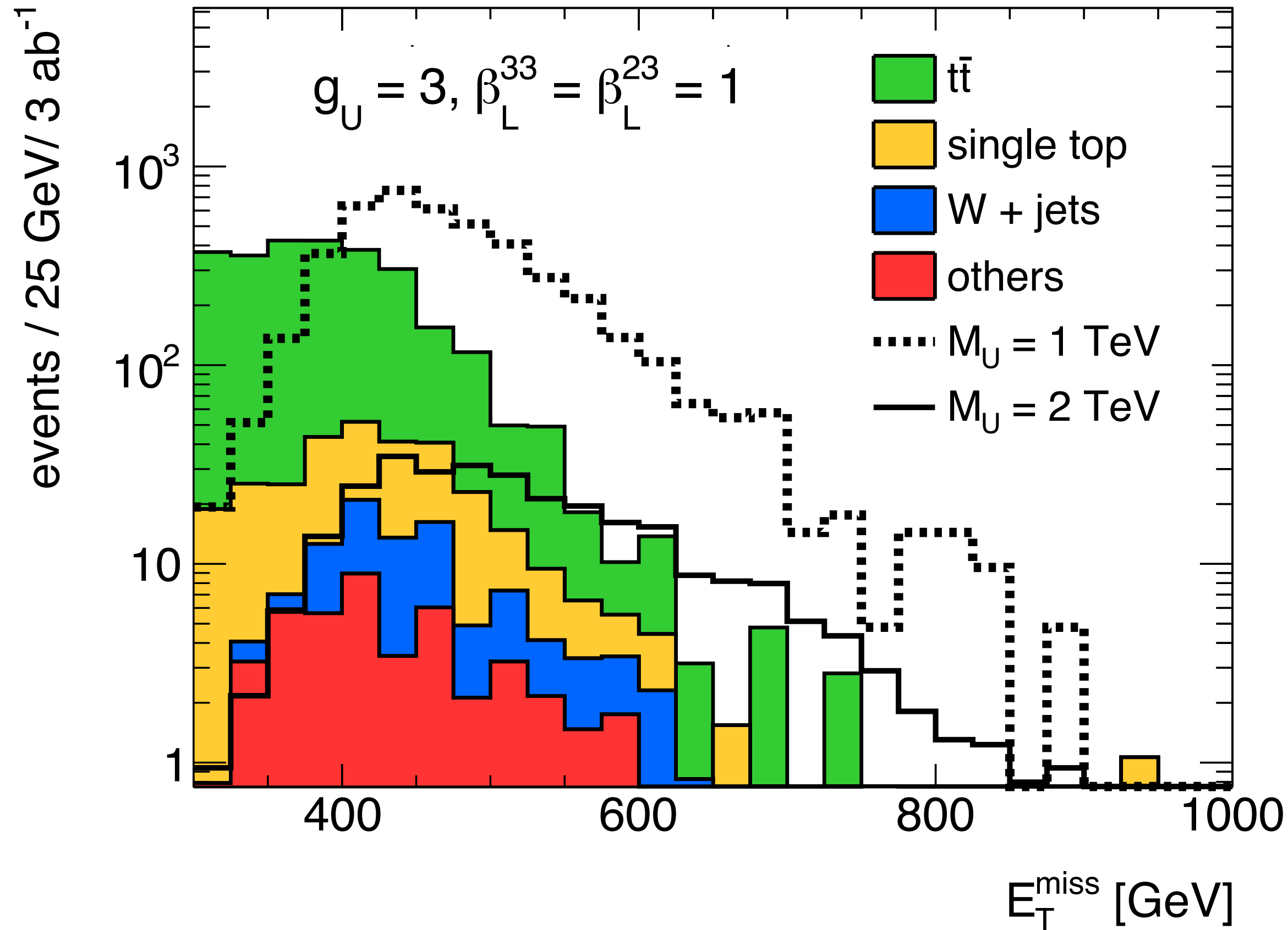


Kinematic distributions of $pp \rightarrow b\tau$ signal

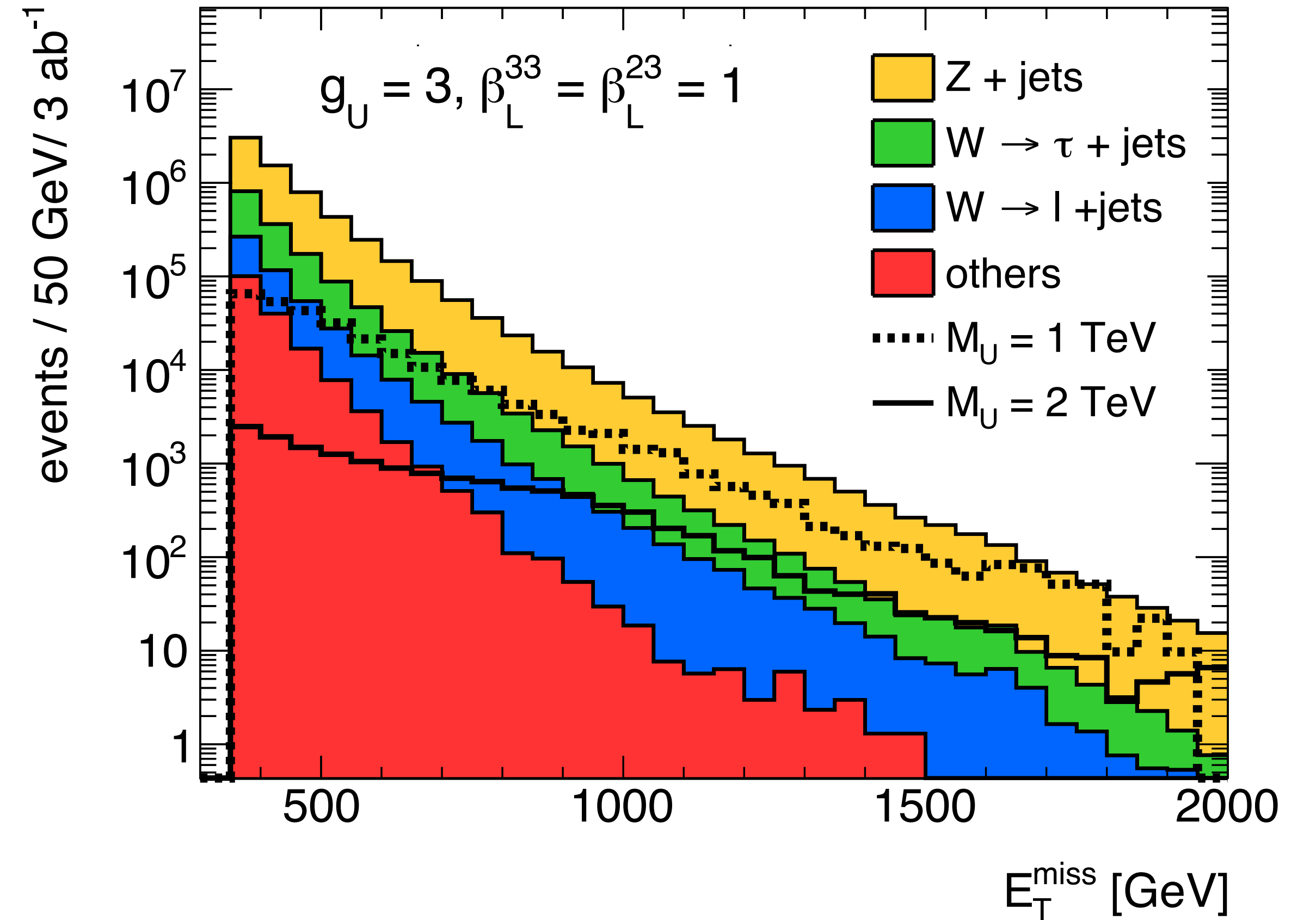


Mono-top & mono-jet distributions

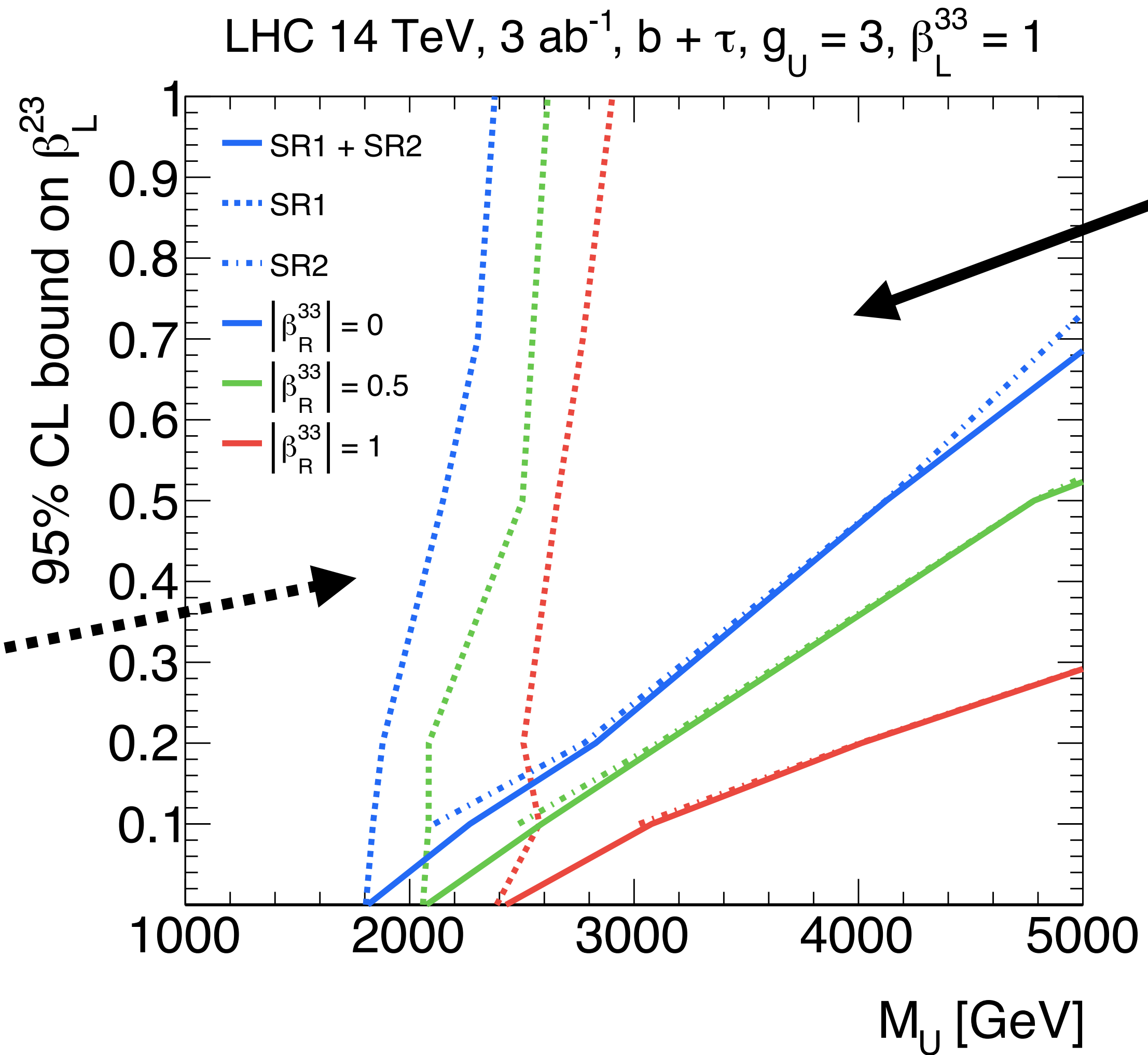
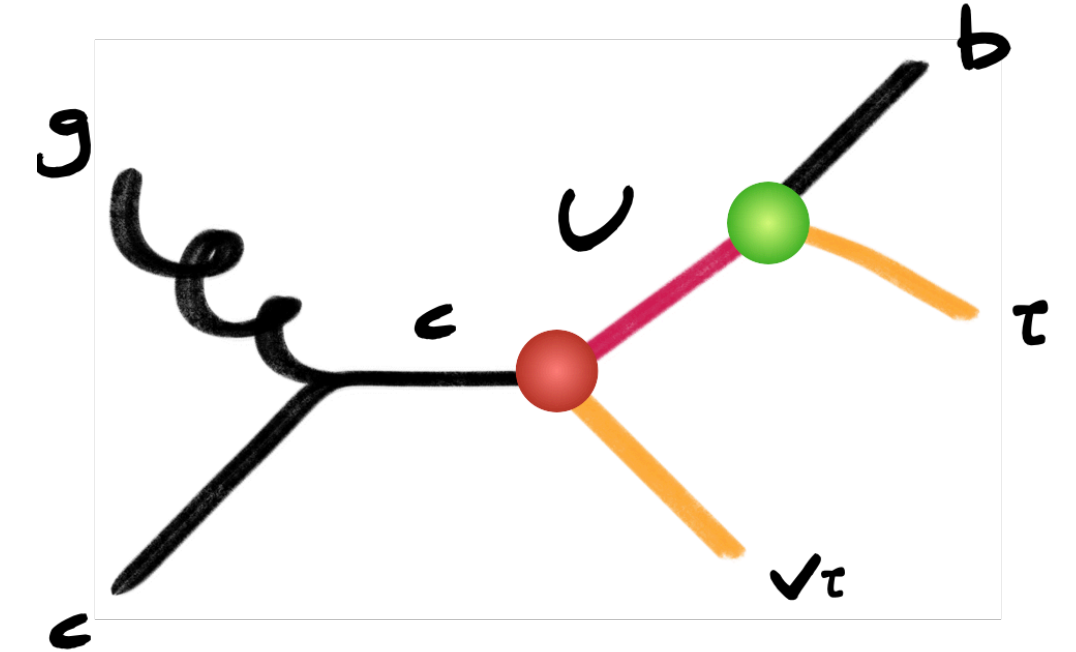
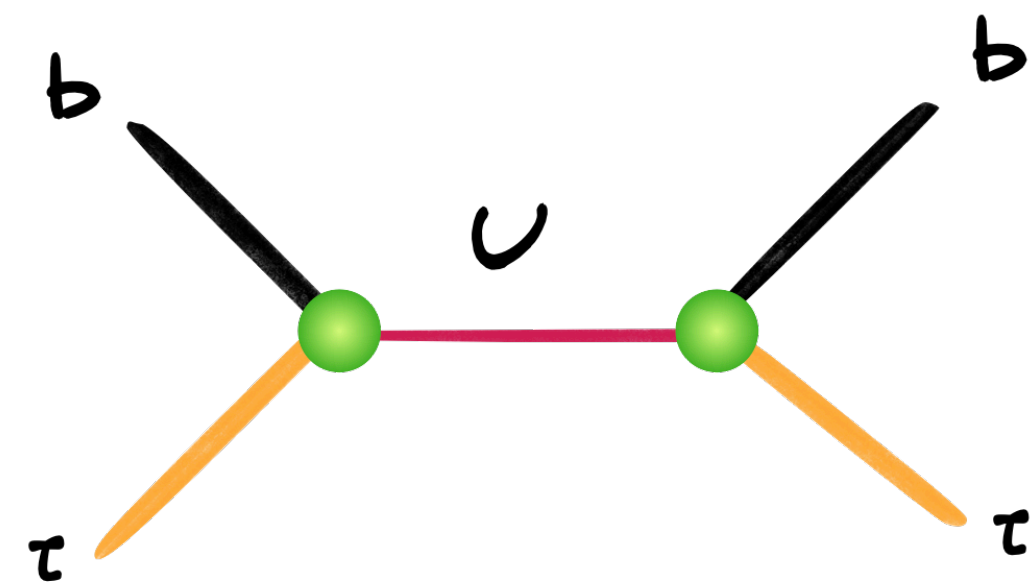
LHC 14 TeV, mono-top



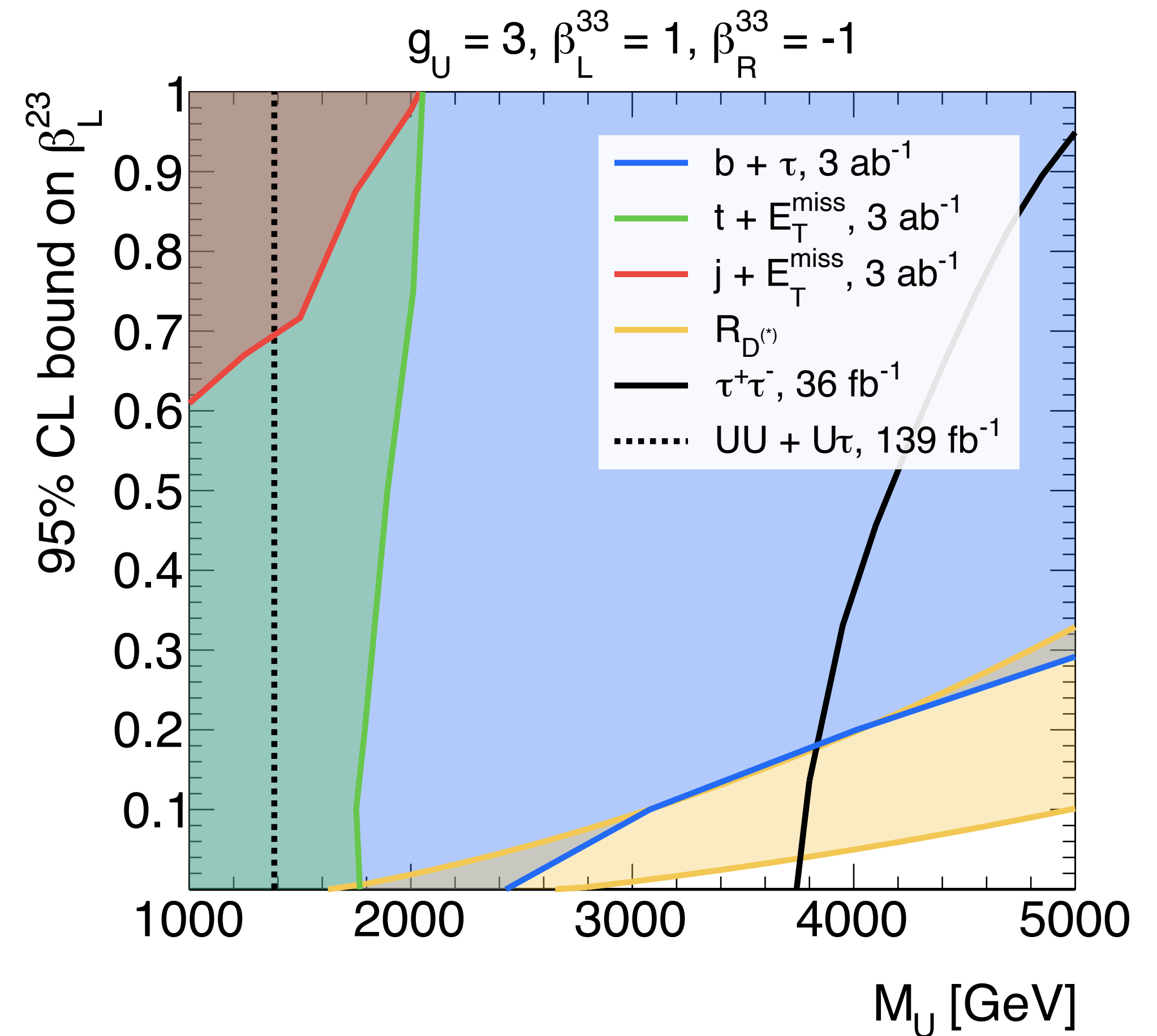
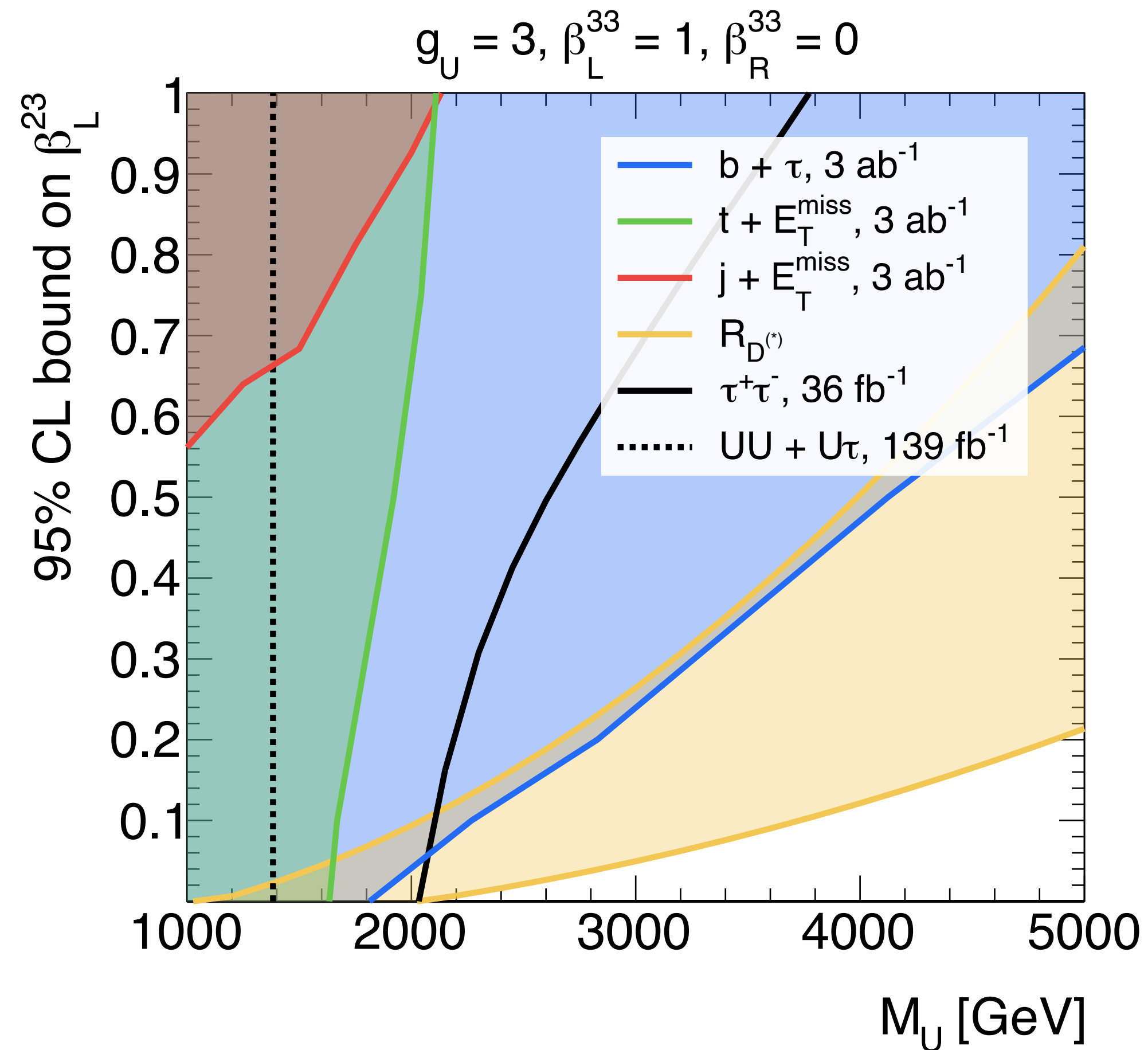
LHC 14 TeV, mono-jet



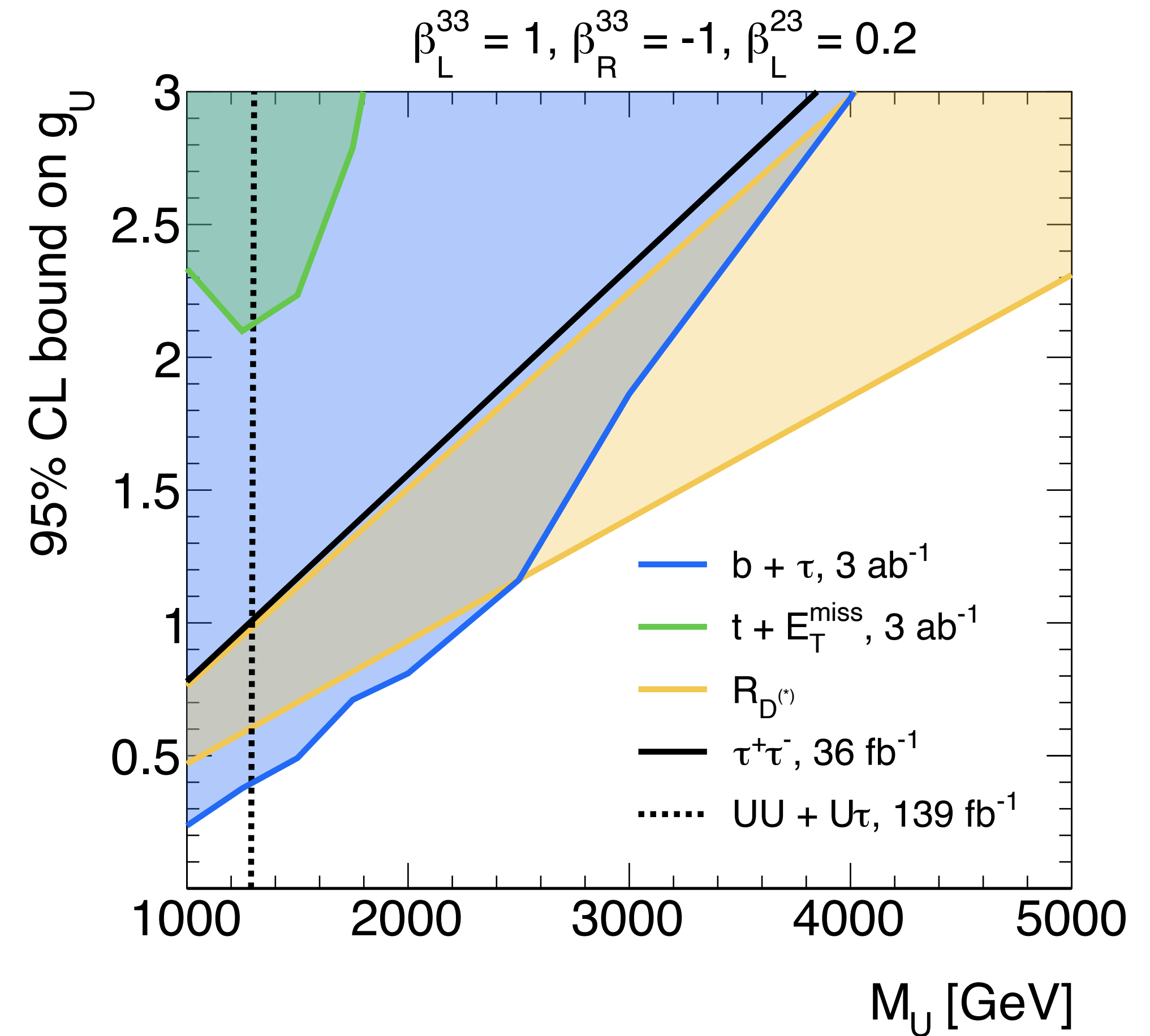
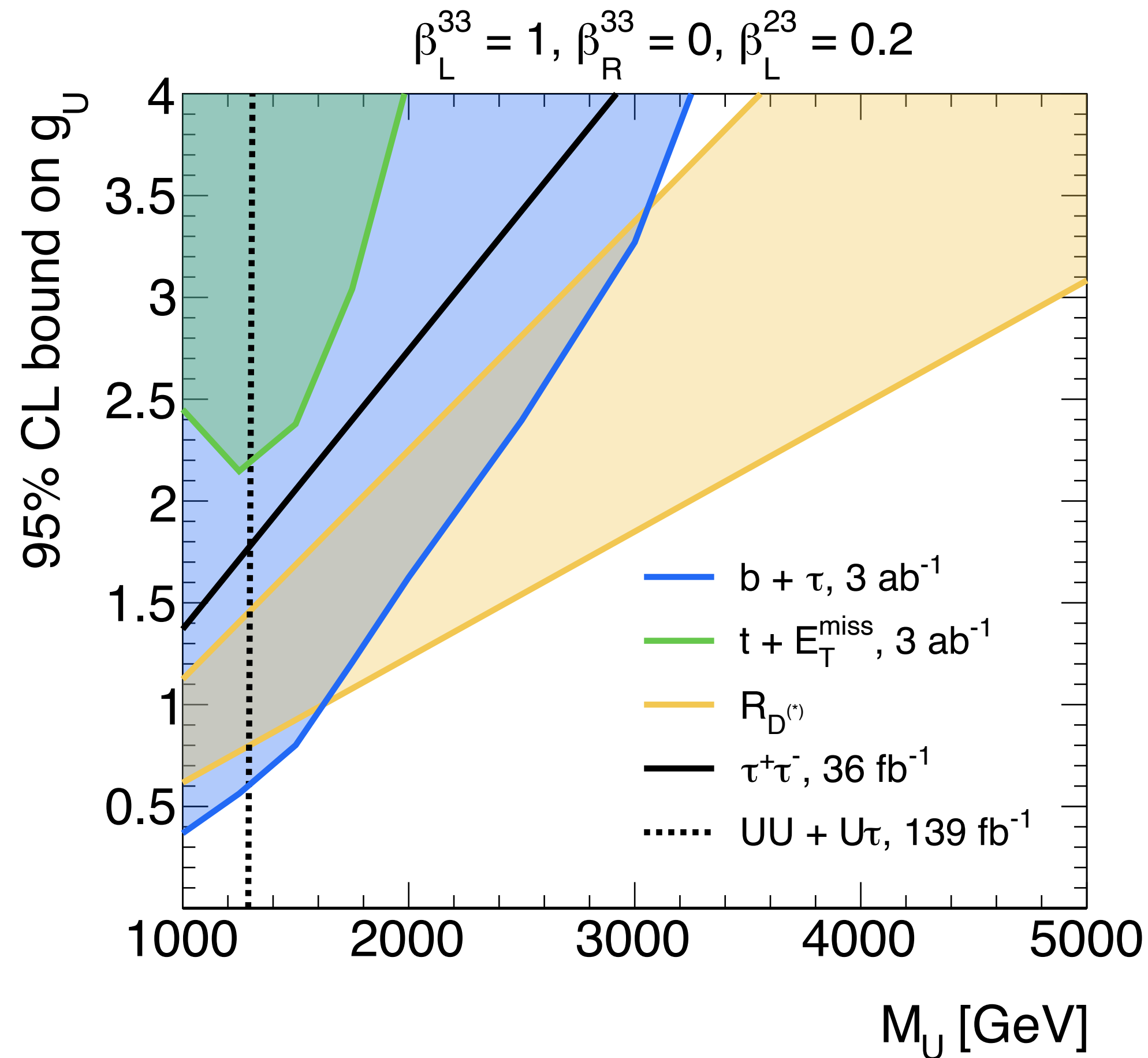
pp → bτ constraints from 2 → 2 & 2 → 3 signal



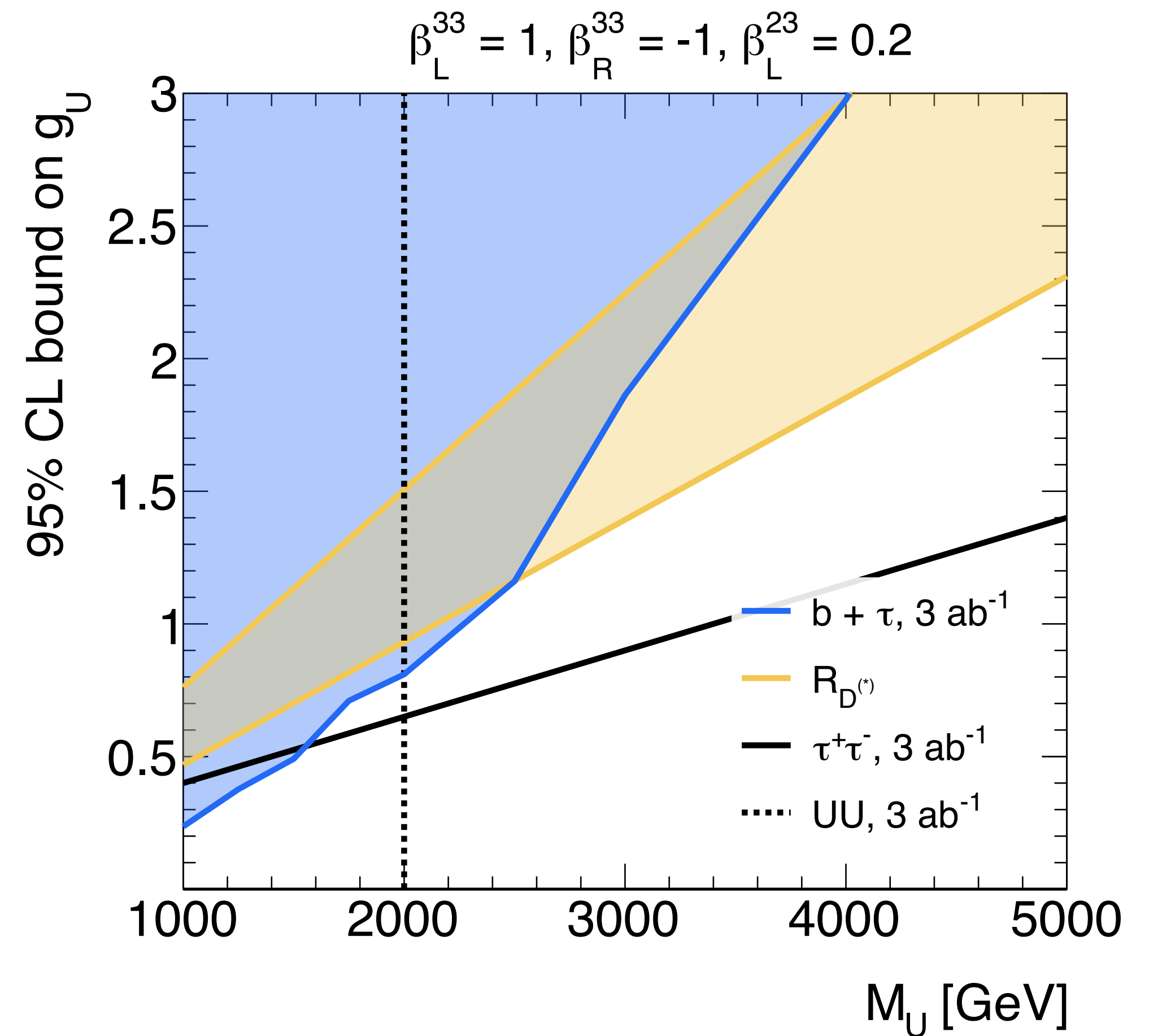
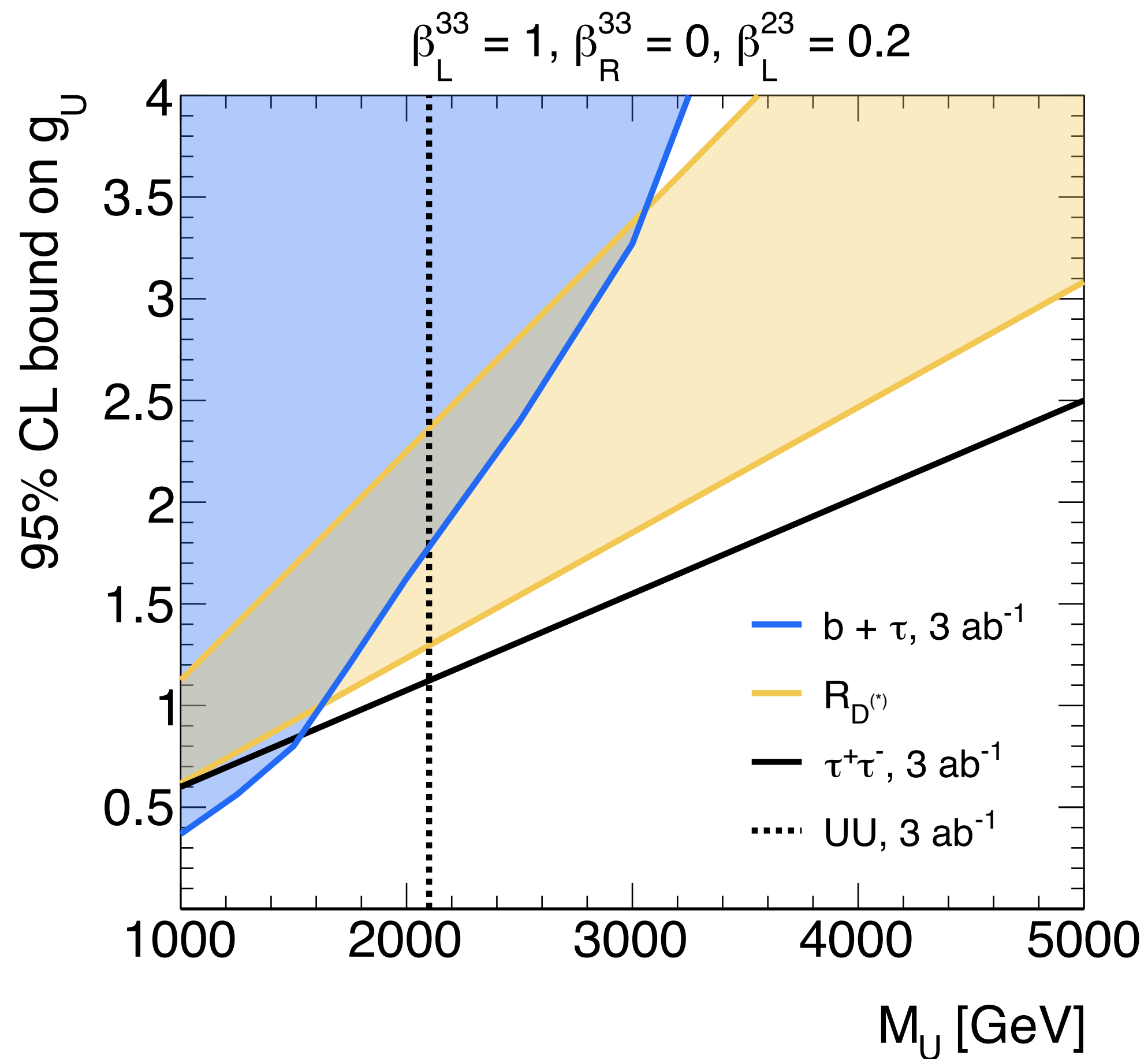
Constraints from $b\tau$, mono-top & mono-jet



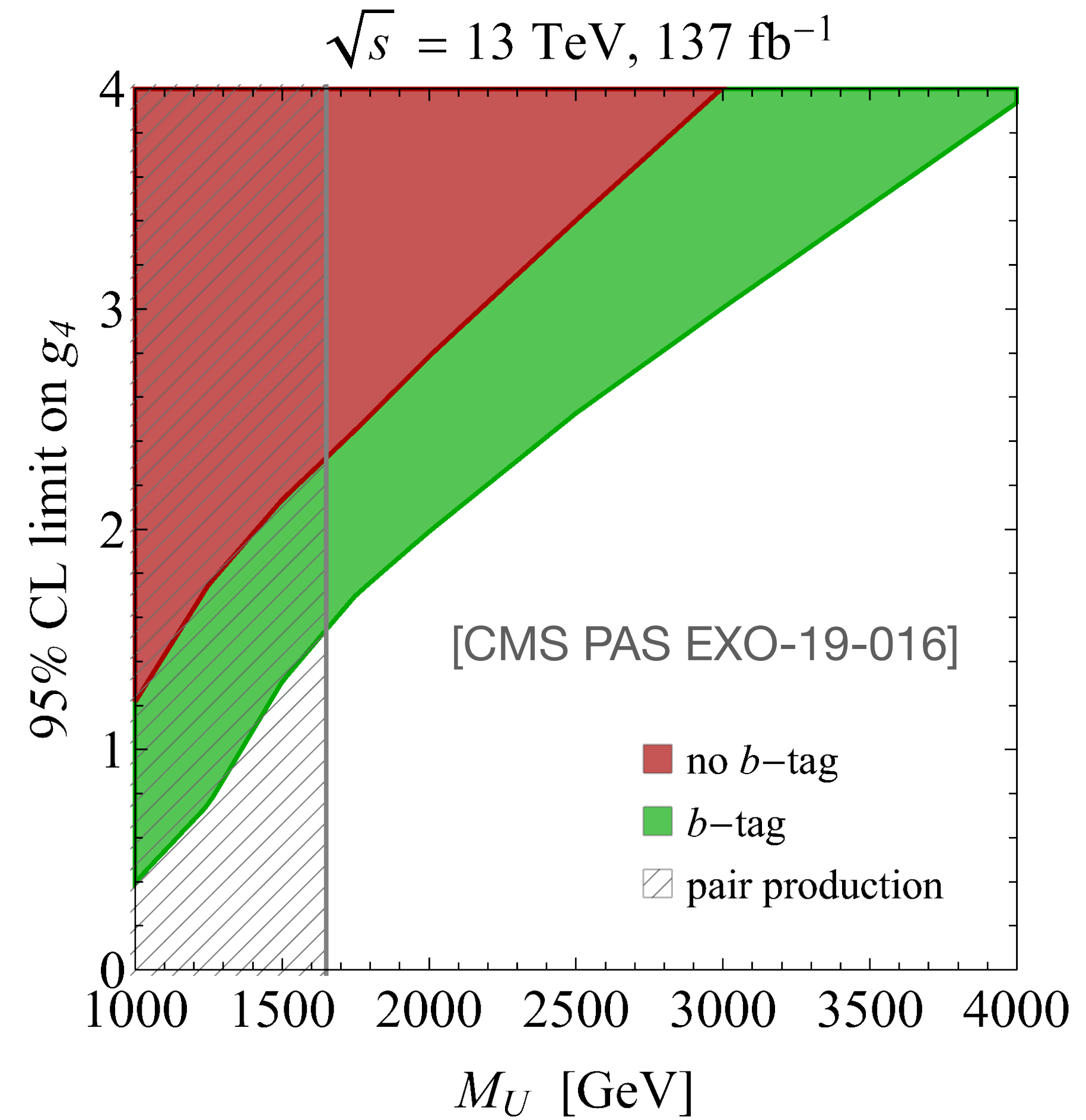
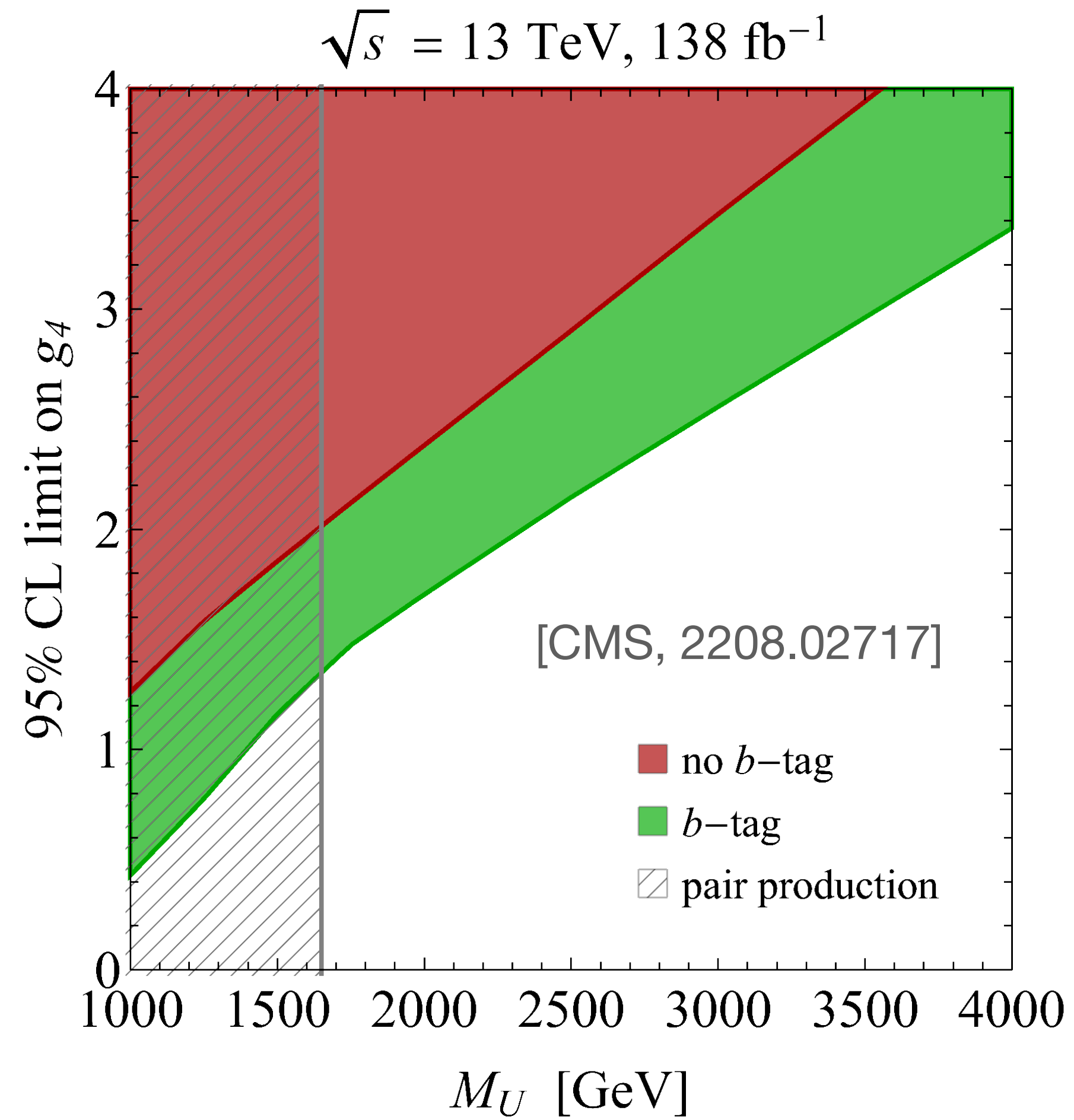
Constraints from $b\tau$, mono-top & mono-jet



Prospects of LQ search strategies

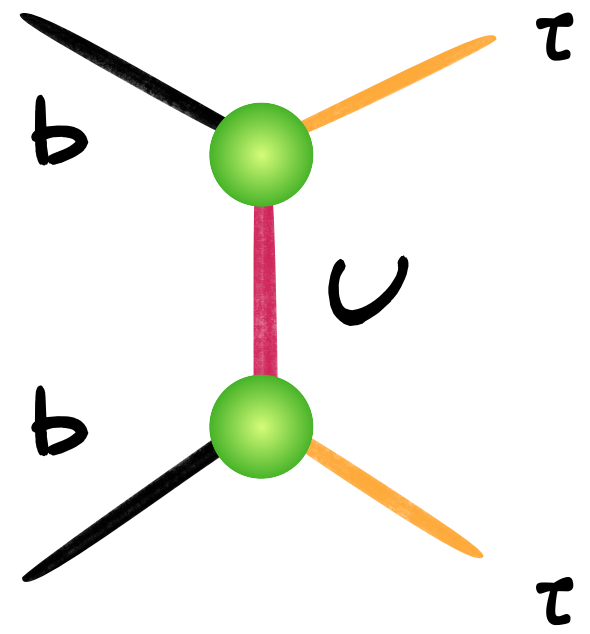
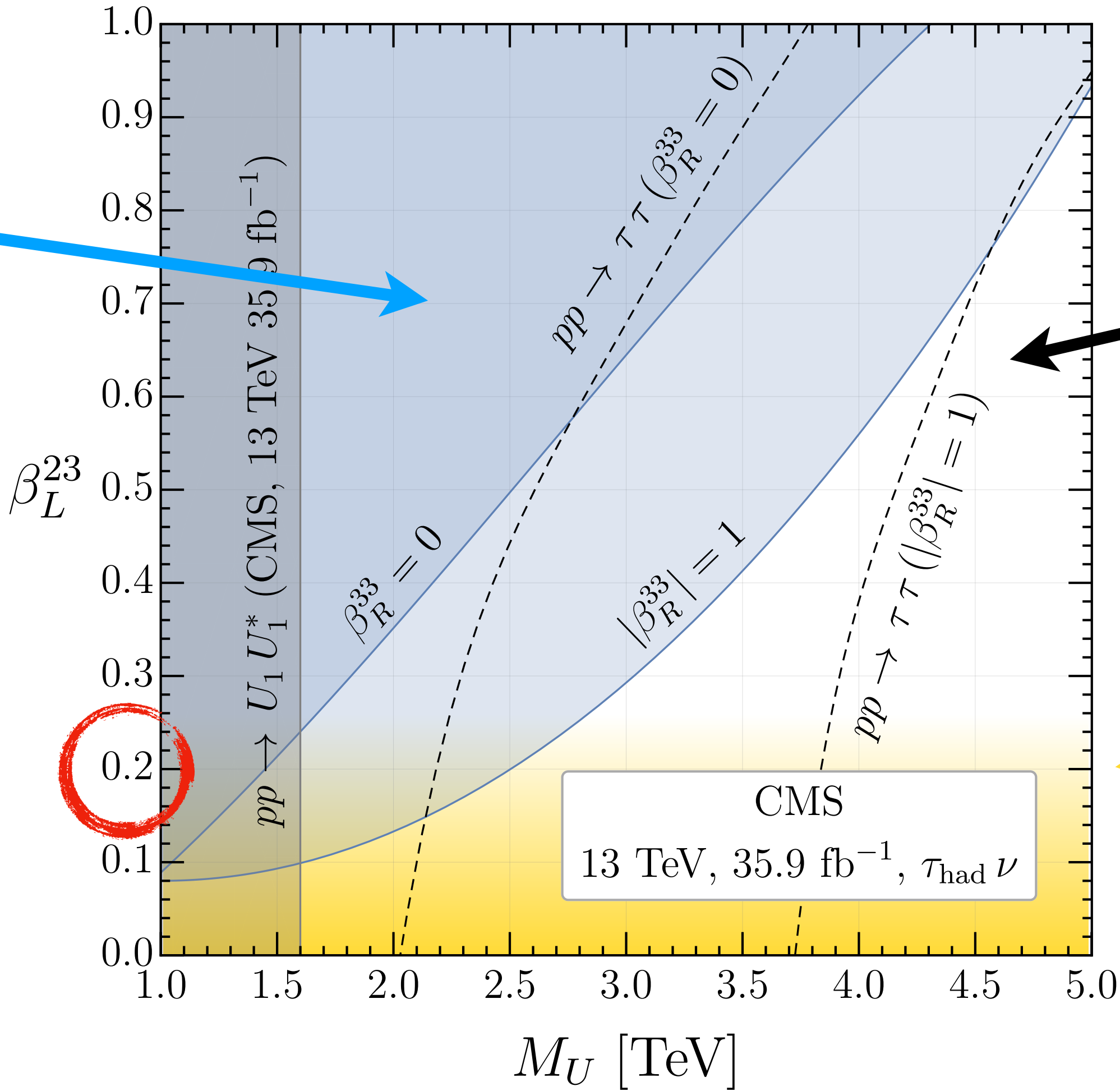
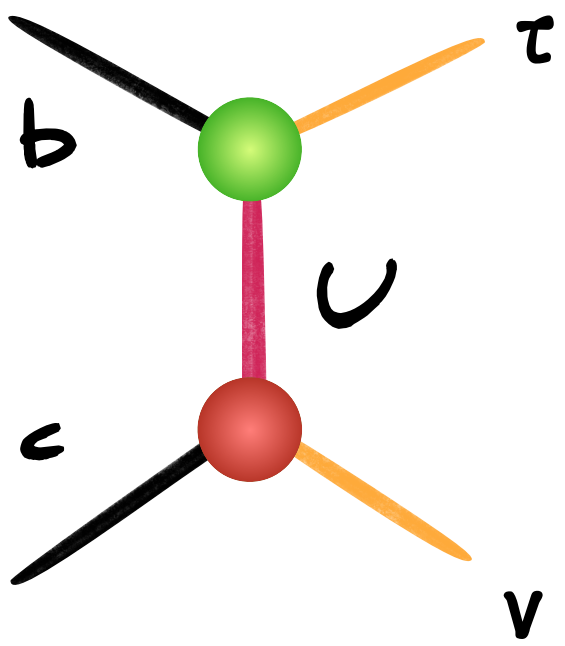


Ditau limits on singlet vector LQs from CMS



[NLO+PS accurate results for t-channel ditau production in LQ models obtained in UH, Schnell & Schulte, 2207.00356; 2209.12780]

LHC bounds: $pp \rightarrow \tau\tau$ vs. $pp \rightarrow \tau\nu$

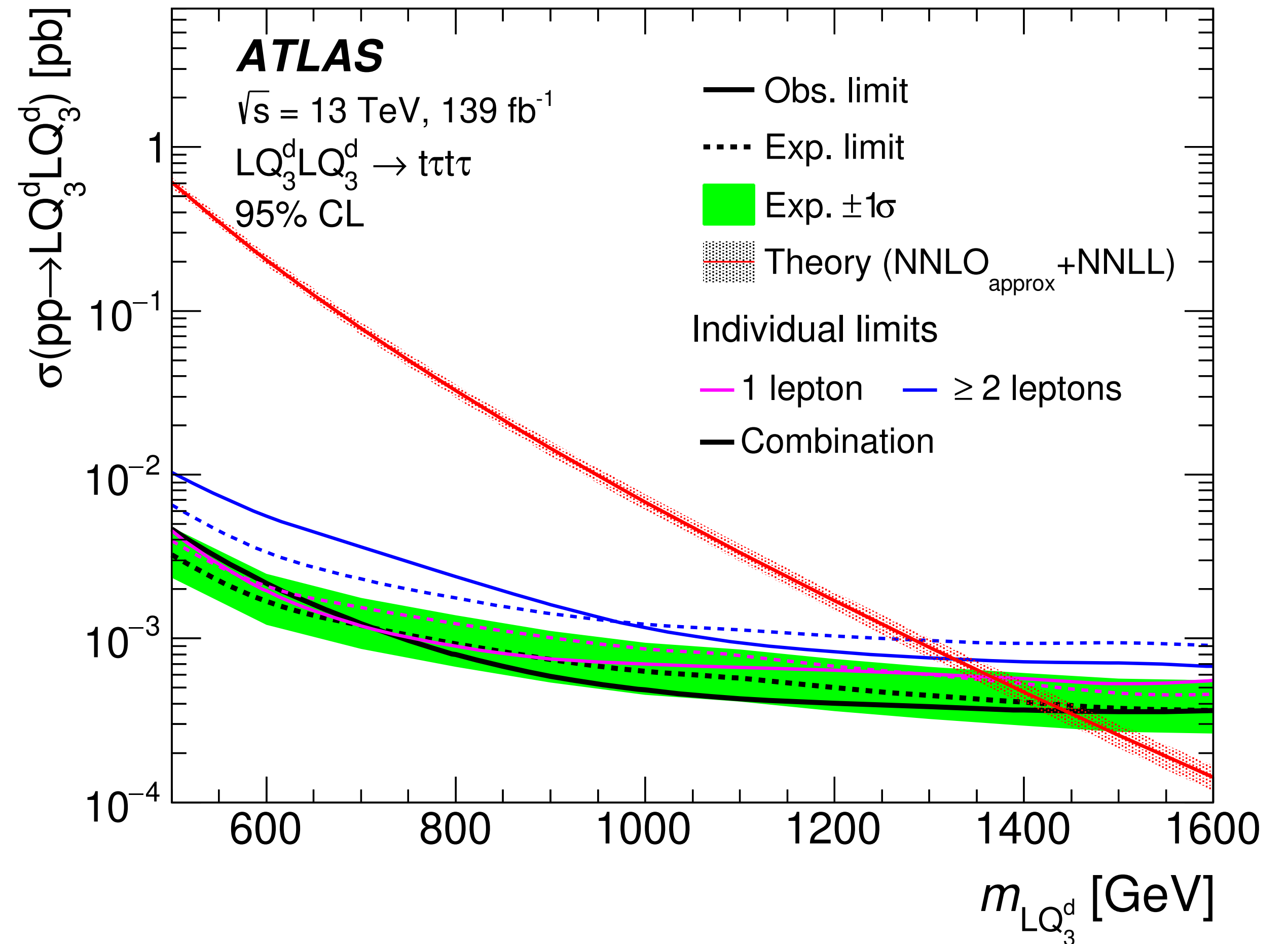
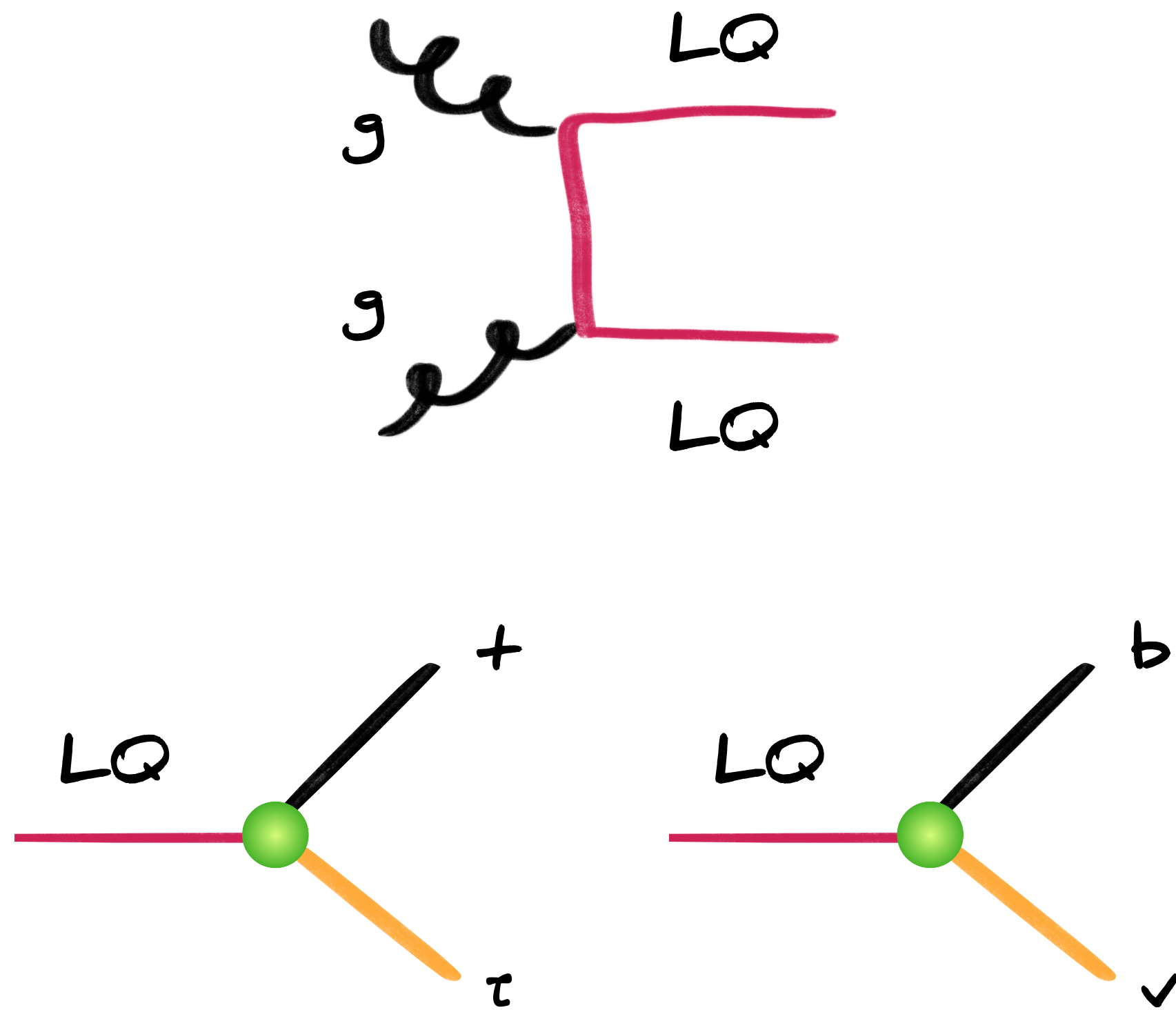


parameter space of singlet vector LQ with natural flavour structure

[Baker, Fuentes-Martin, Isidori & König, 1901.10480]

Another LQ search triggered by B anomalies

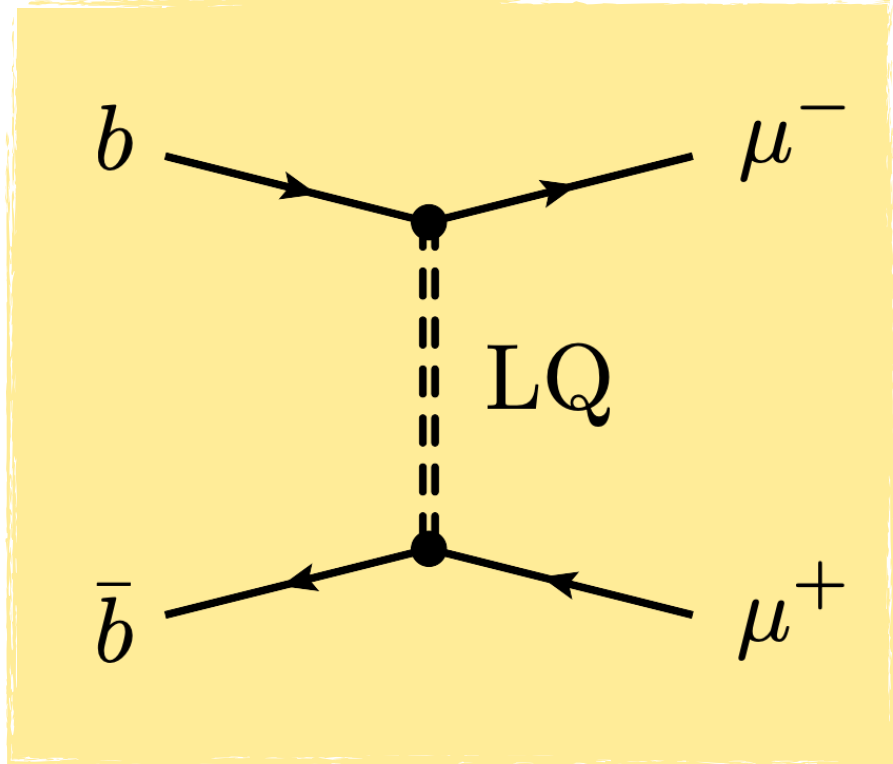
[ATLAS, 2101.11582]



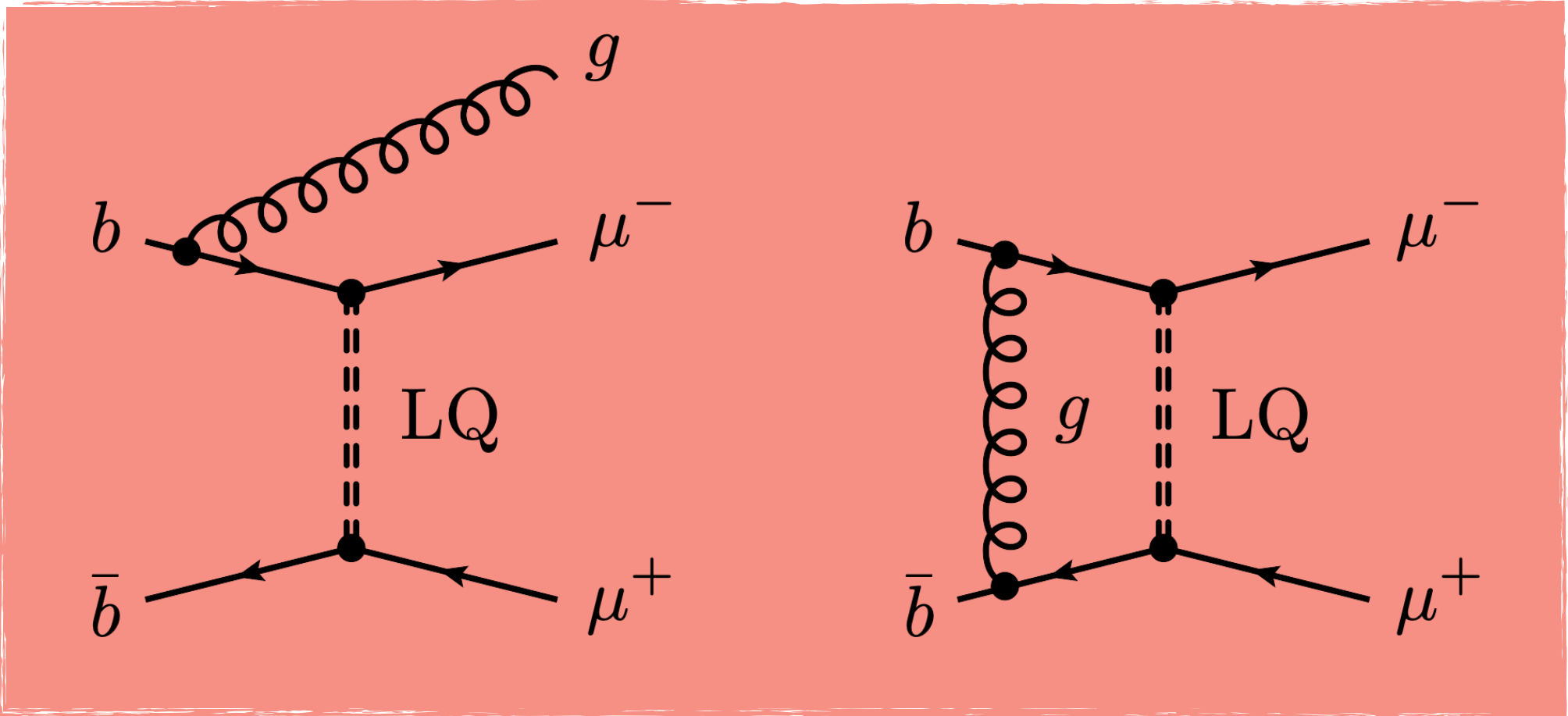
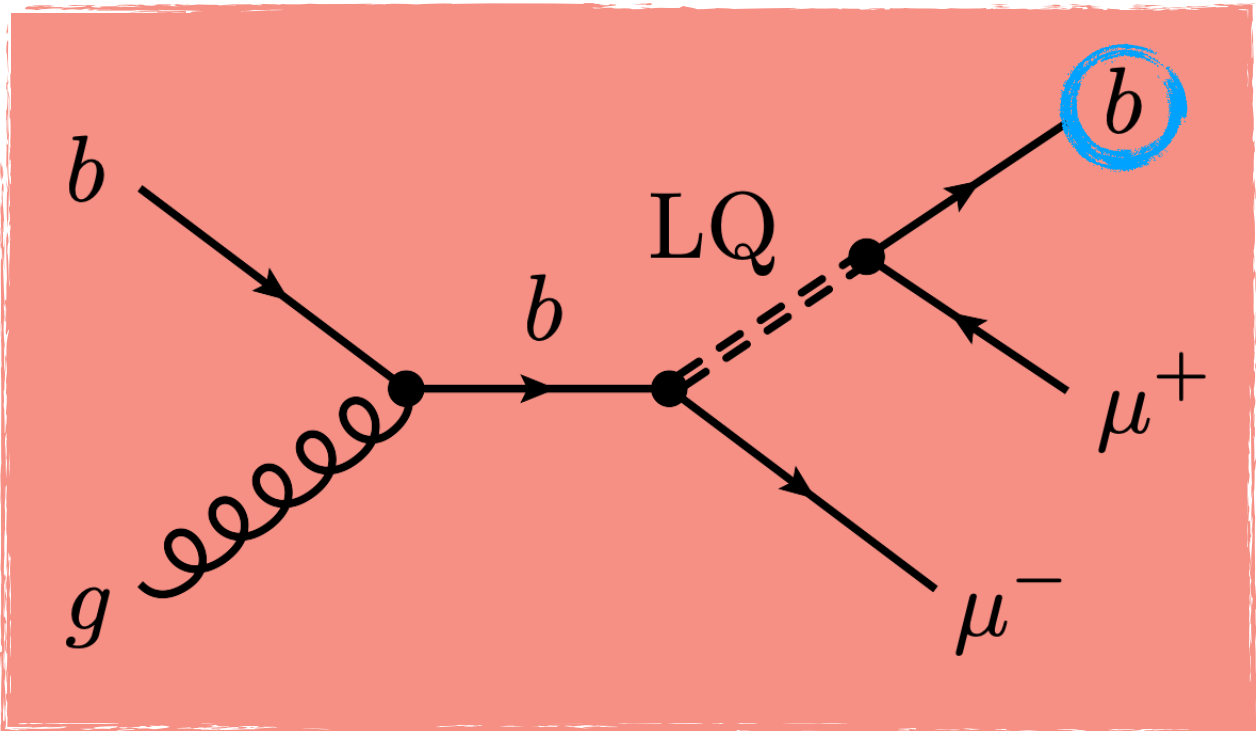
[Bauer & Neubert, 1511.01900]

NLO+PS calculation for scalar LQs

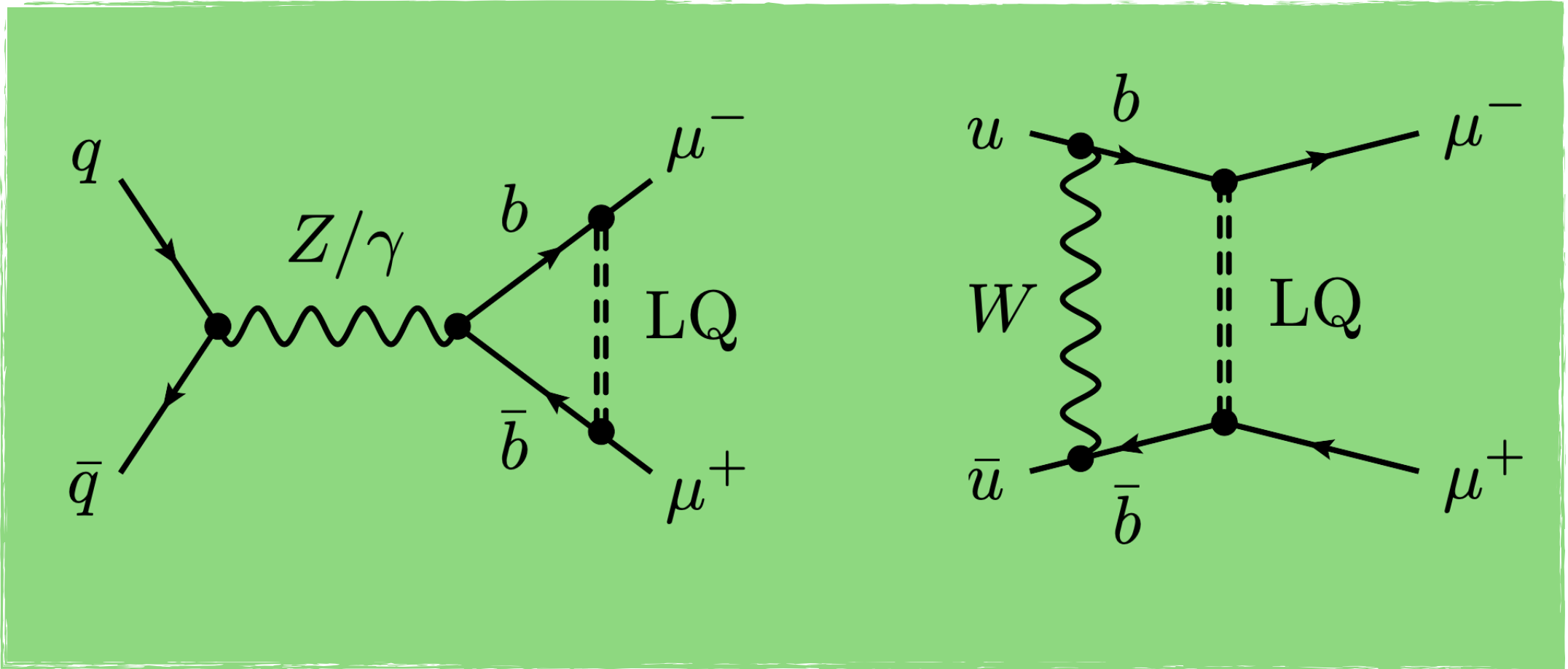
tree-level contribution



infrared finite QCD correction



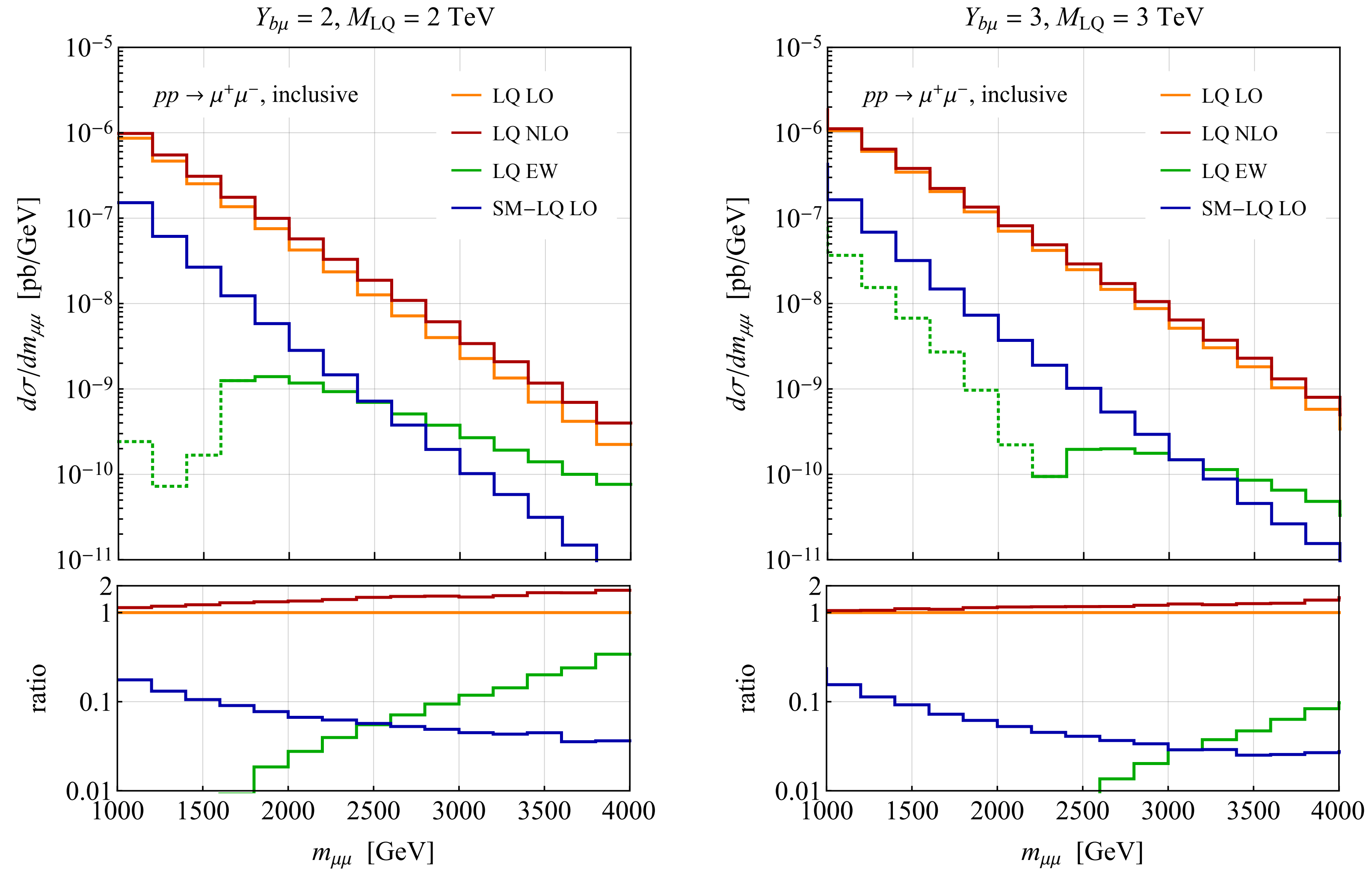
real & virtual correction to Born level



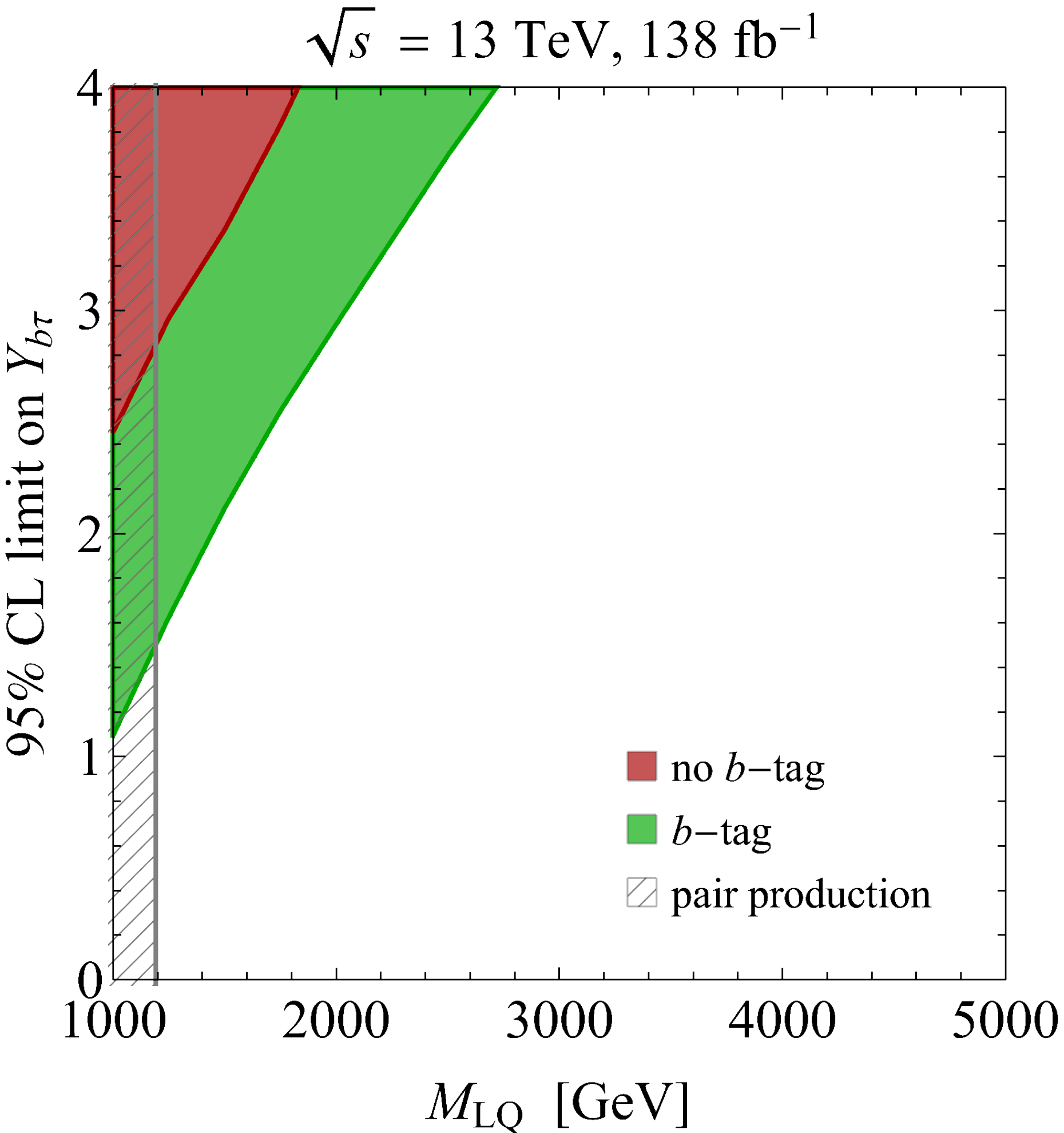
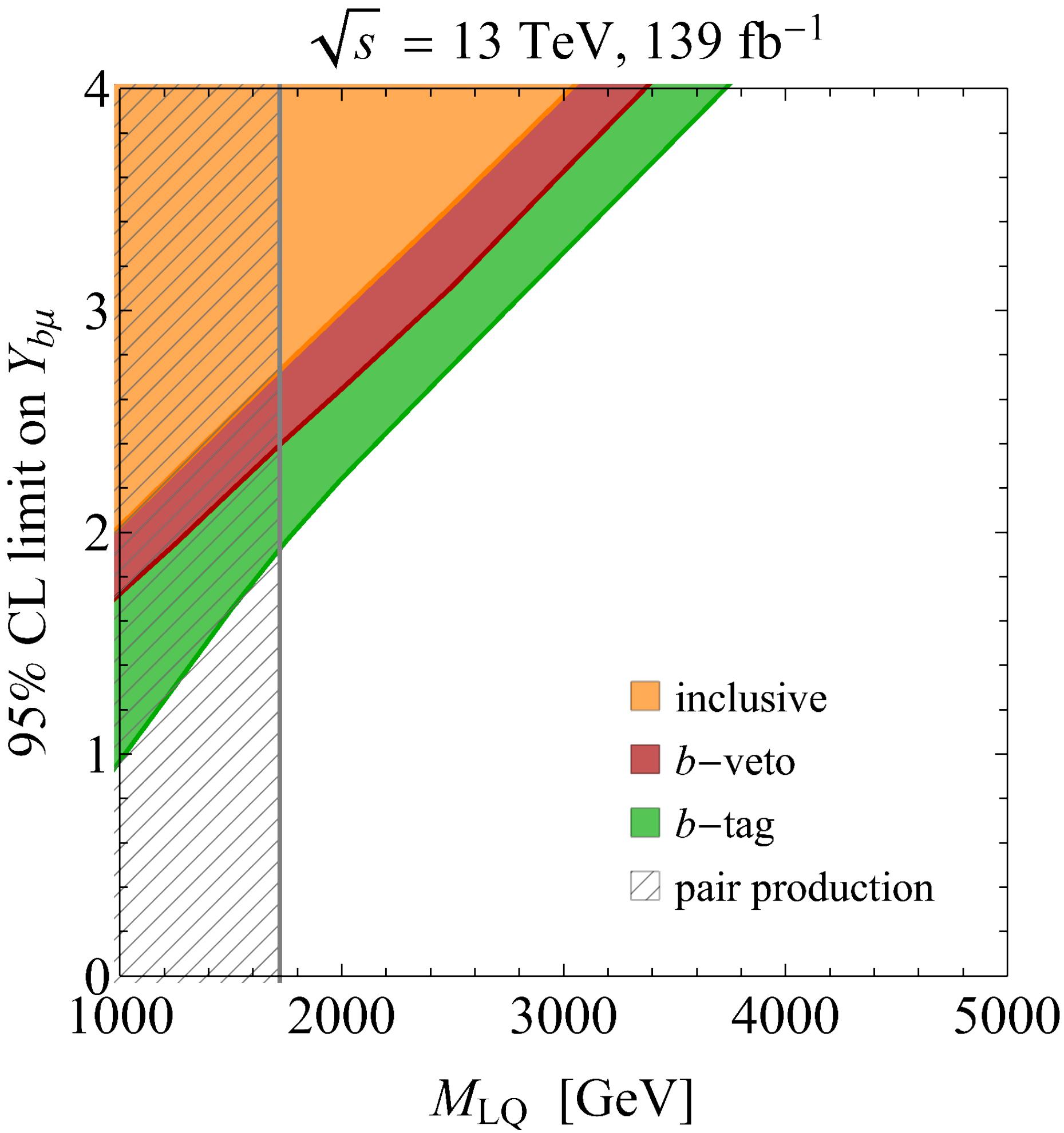
electroweak corrections

[UH, Schnell & Schulte, 2207.00356]

Size of NLO, EW & interference effects

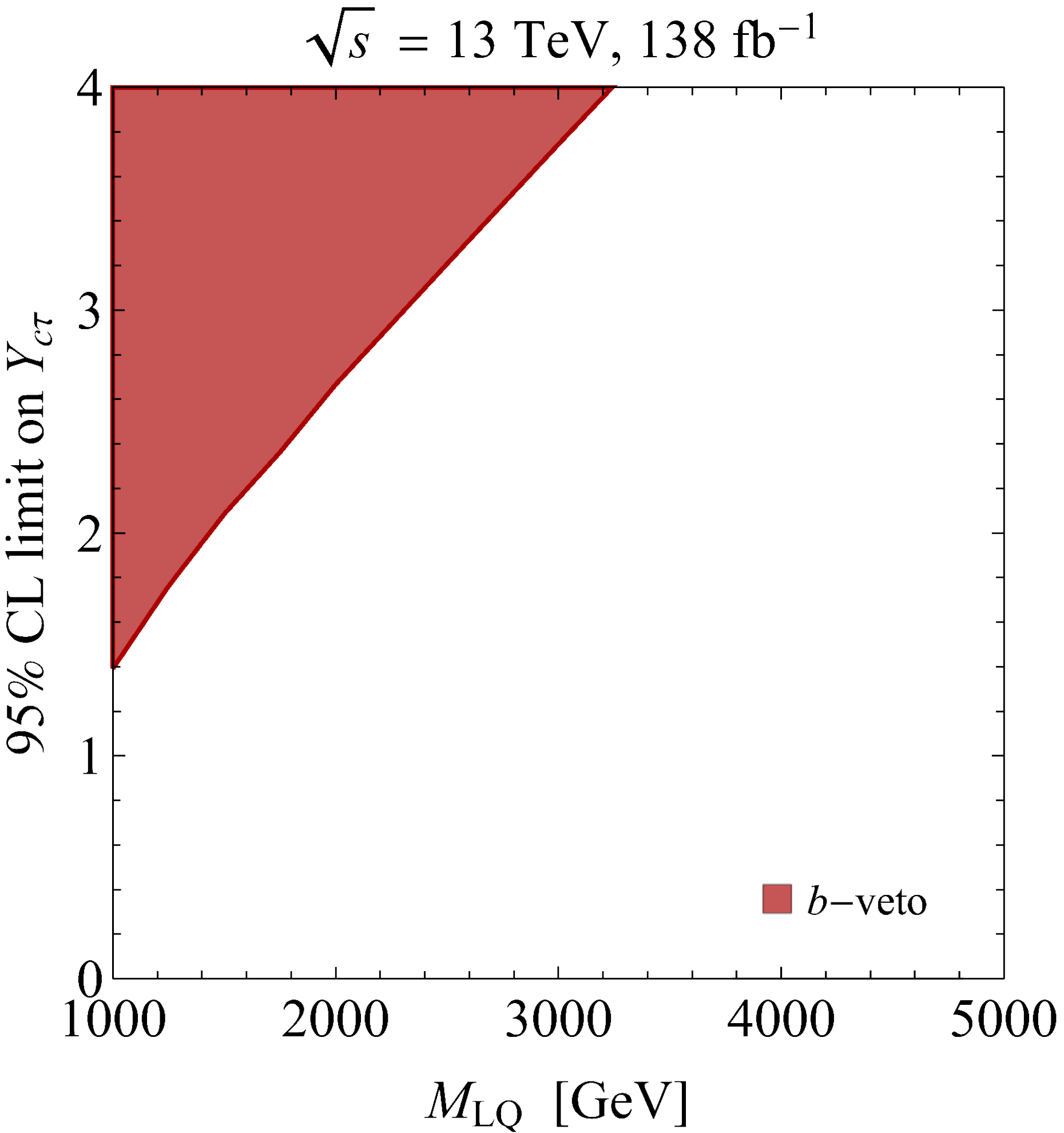
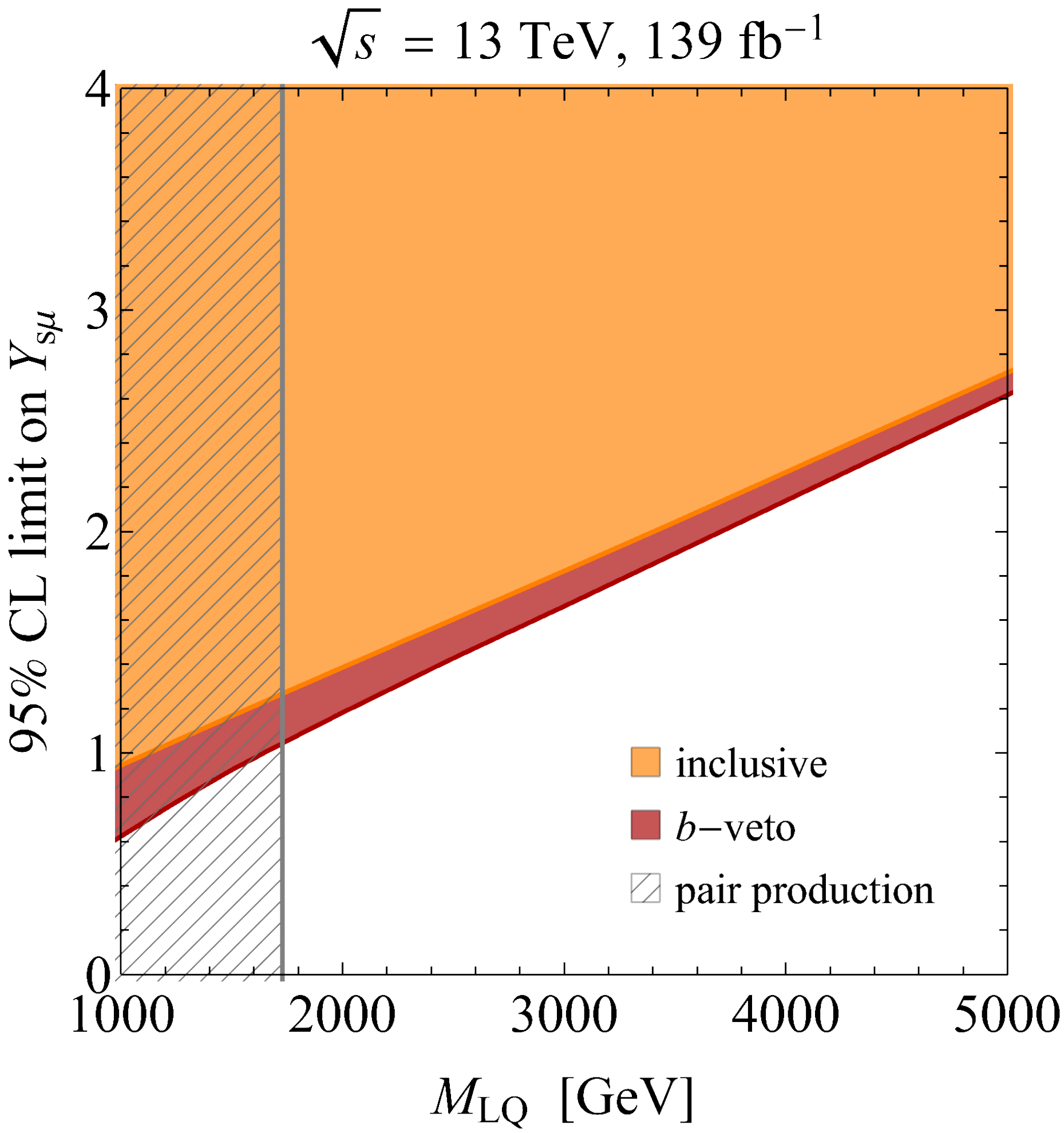


Dimuon constraints on scalar LQs



[UH, Schnell & Schulte, 2207.00356]

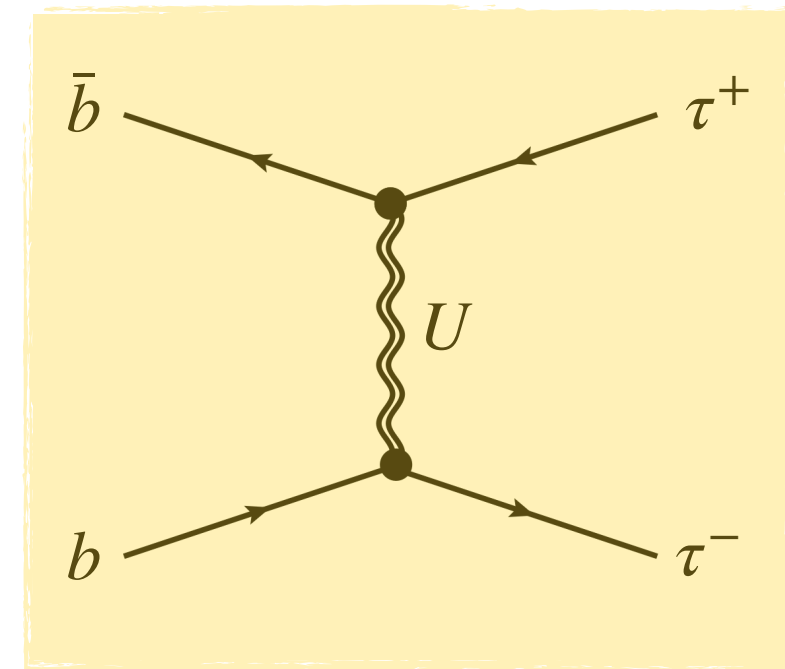
Dimuon constraints on scalar LQs



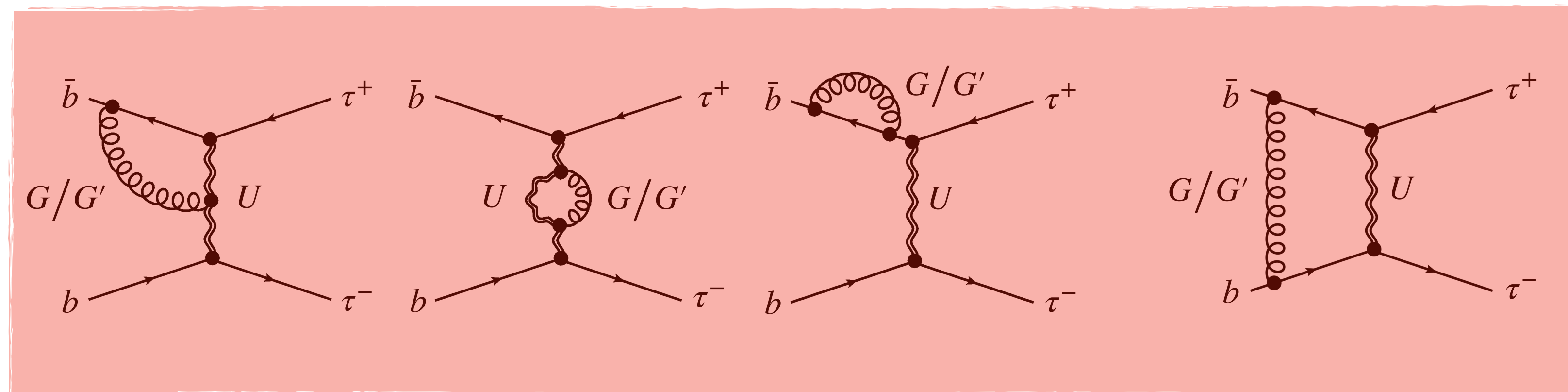
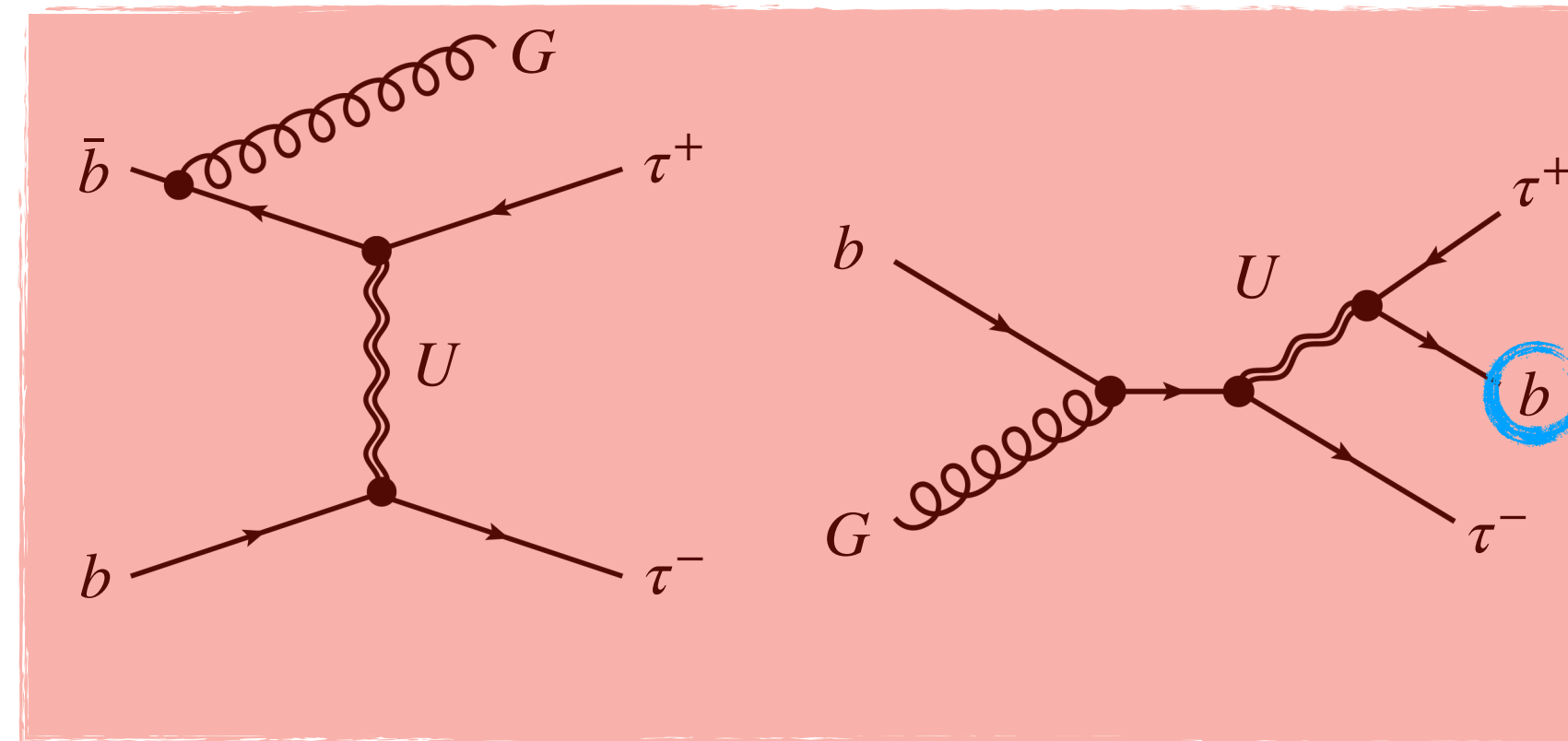
[UH, Schnell & Schulte, 2207.00356]

NLO+PS calculation for singlet vector LQs

tree-level
contribution

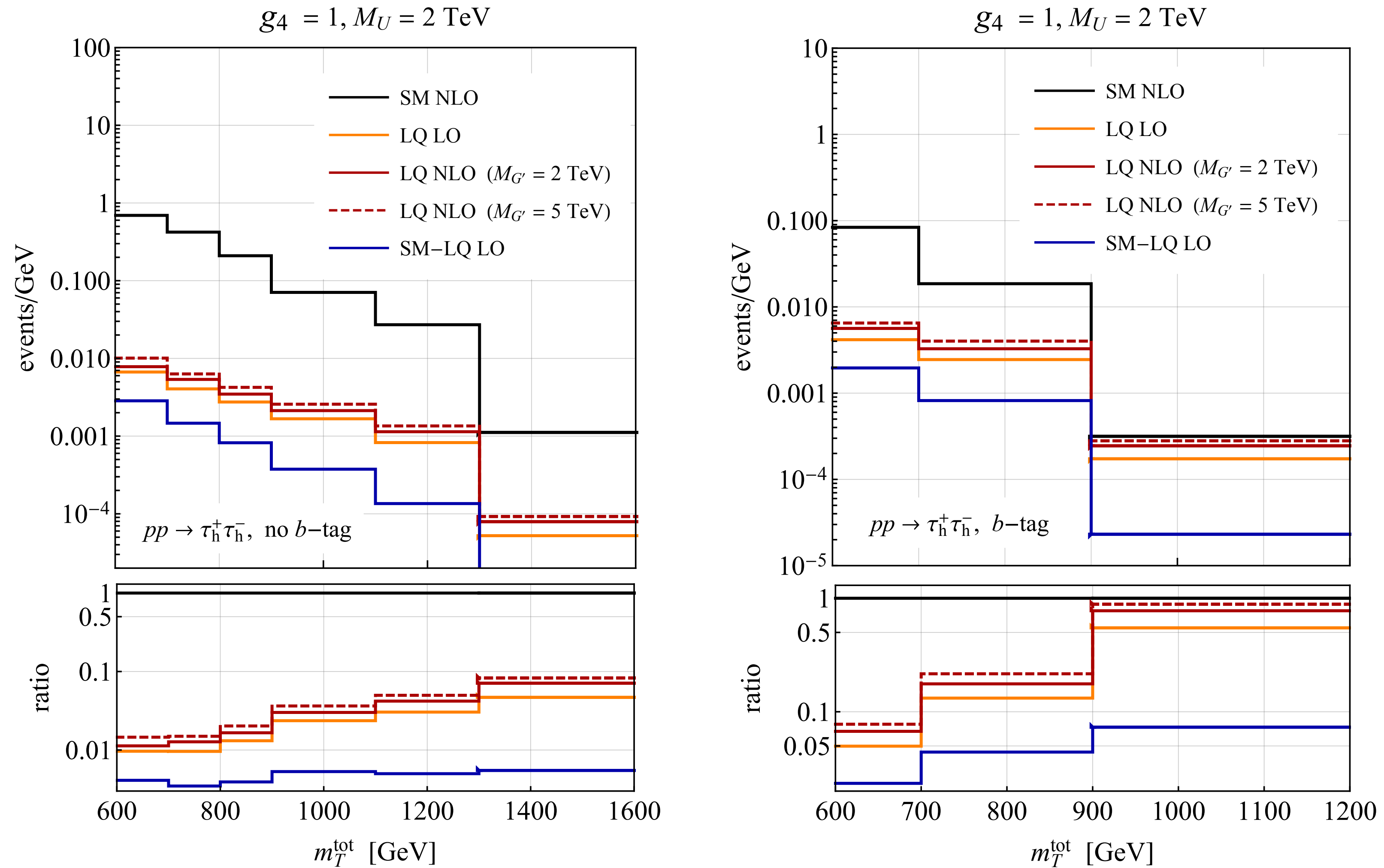


real QCD
corrections

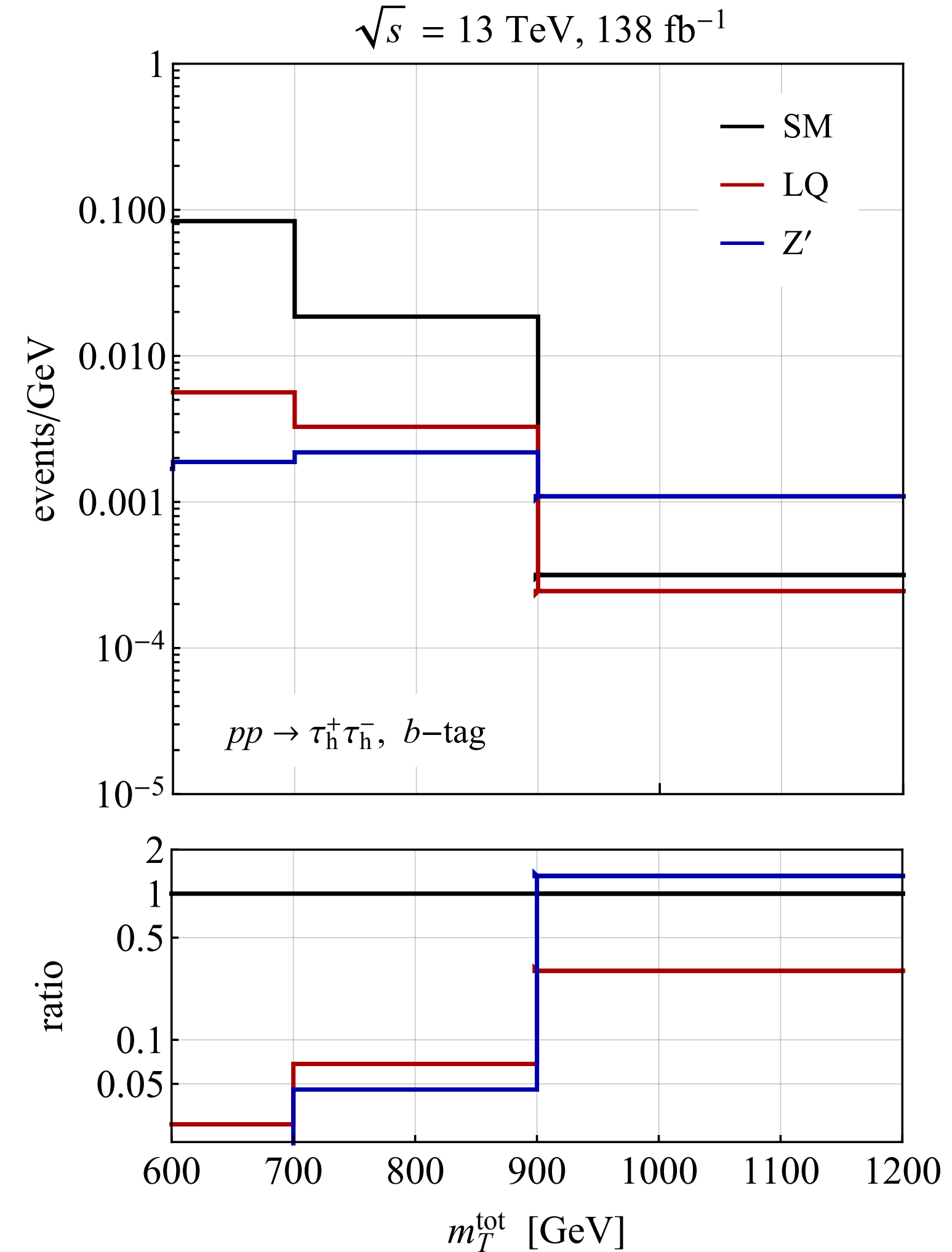
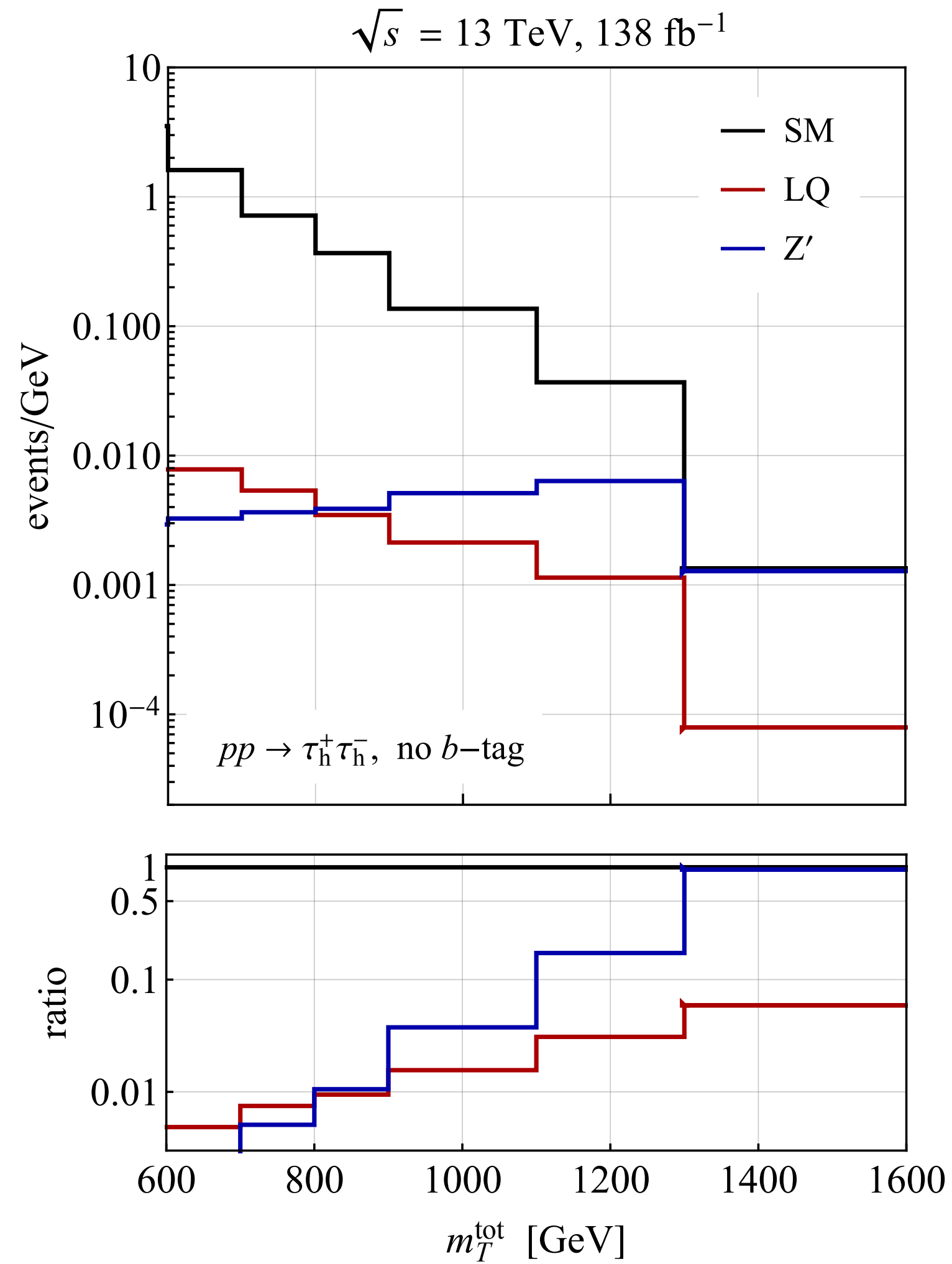


virtual QCD
effects involving
gluon (G) &
coloron (G')

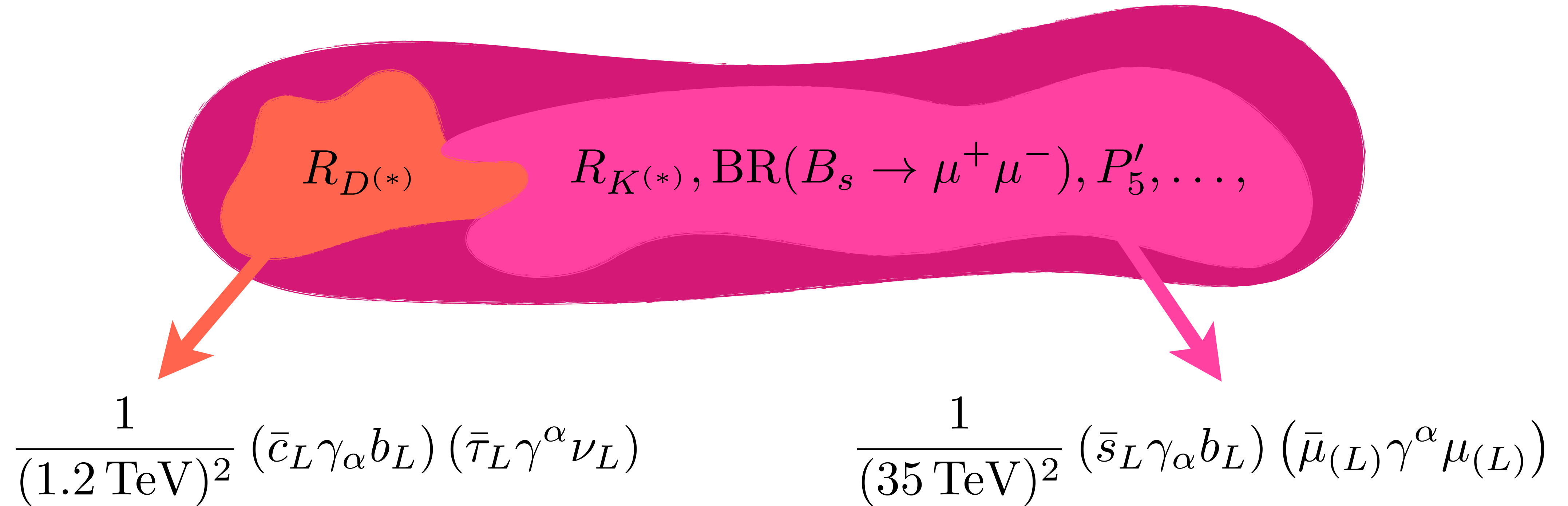
Size of NLO & interference effects



Z' contributions in 4-3-2-1 model

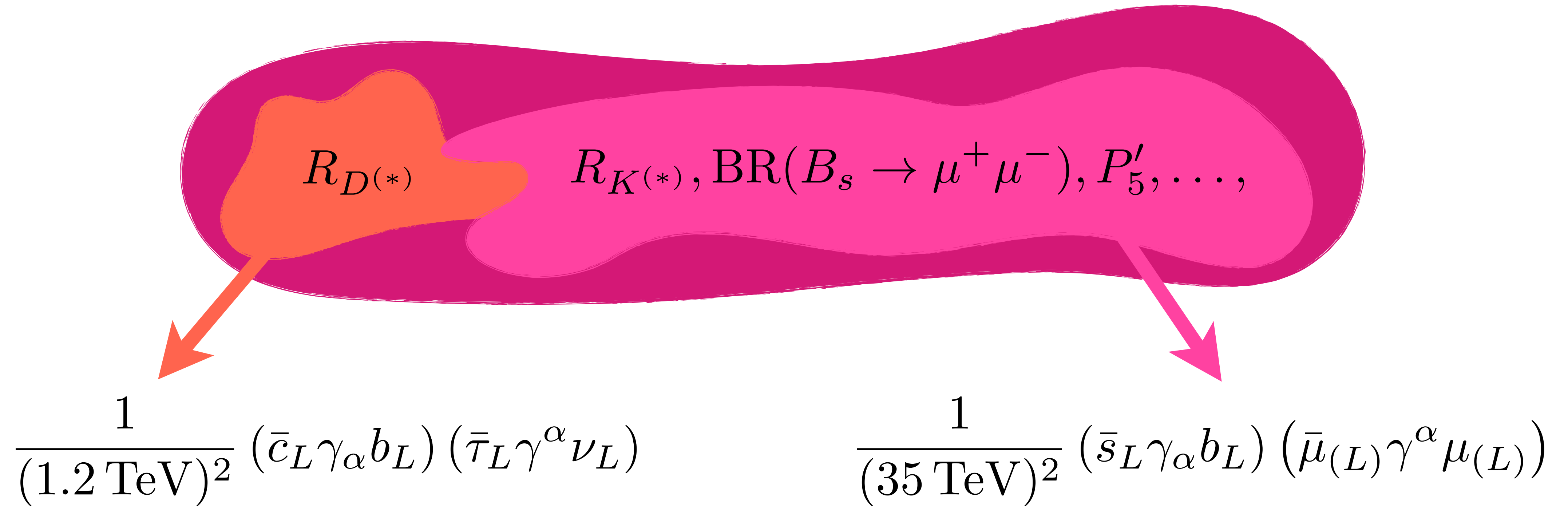


B anomalies in a nutshell



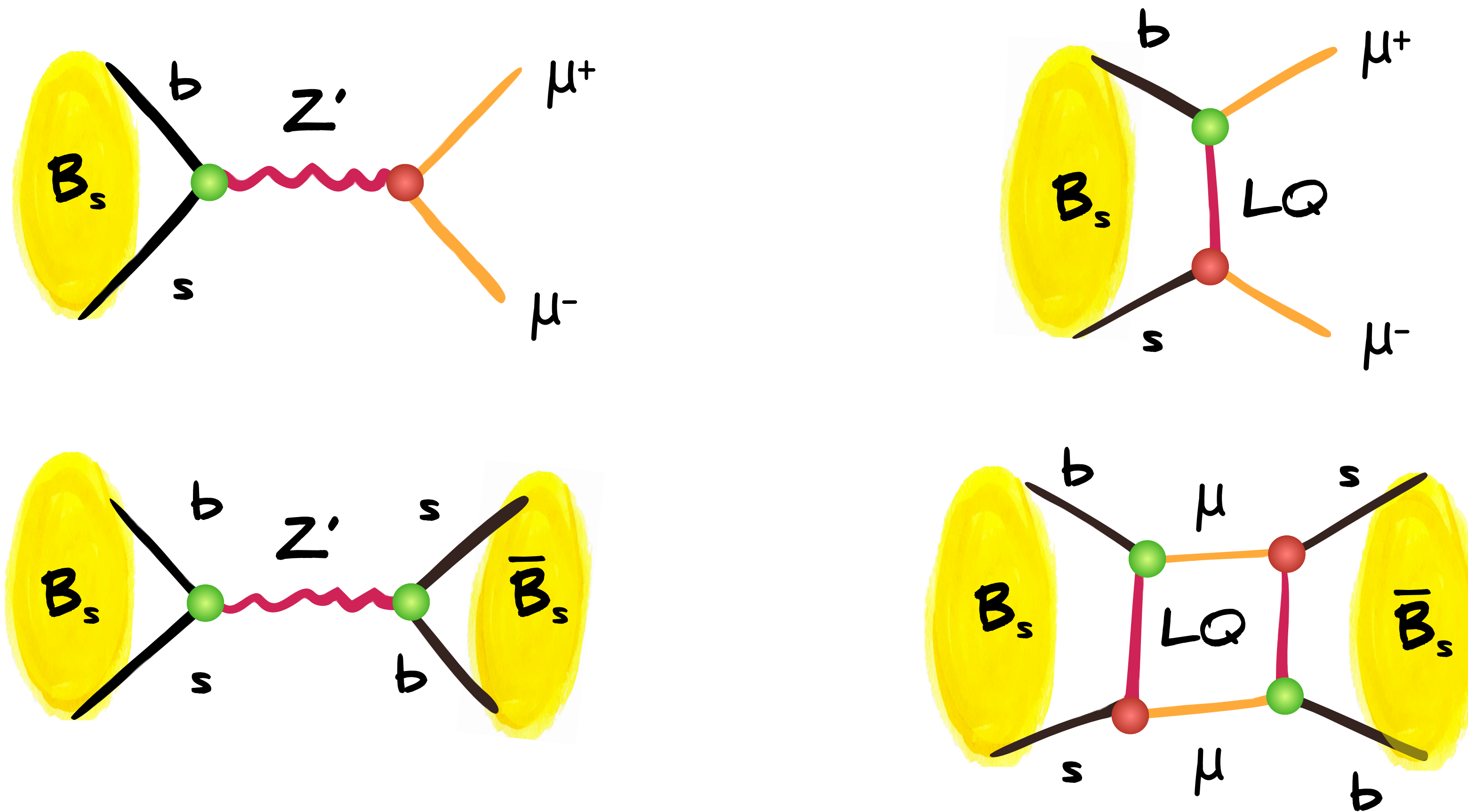
Both sets of B anomalies challenge(d) assumption of lepton flavour universality (LFU), which is usually taken for granted in high-energy physics

B anomalies in a nutshell



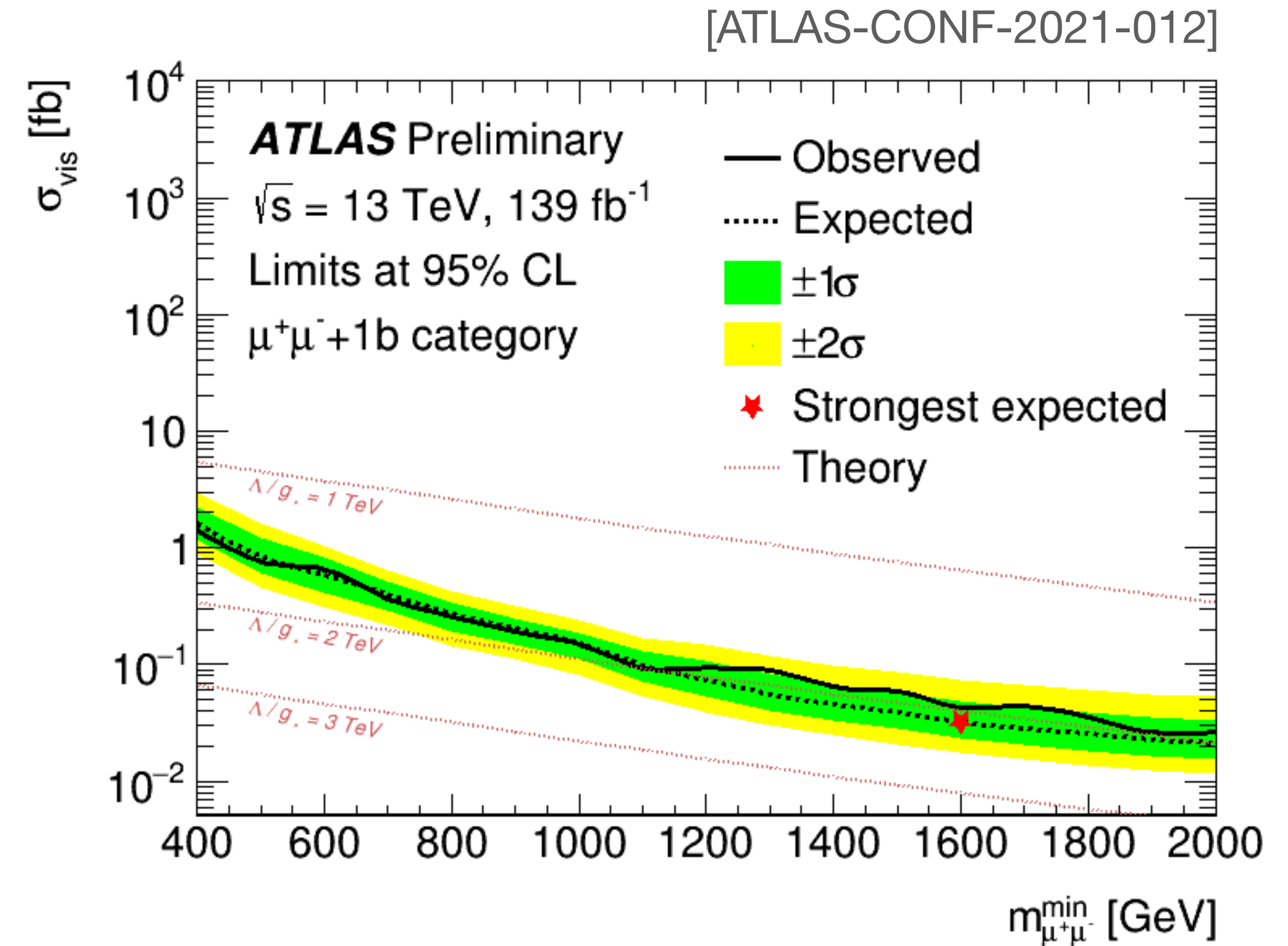
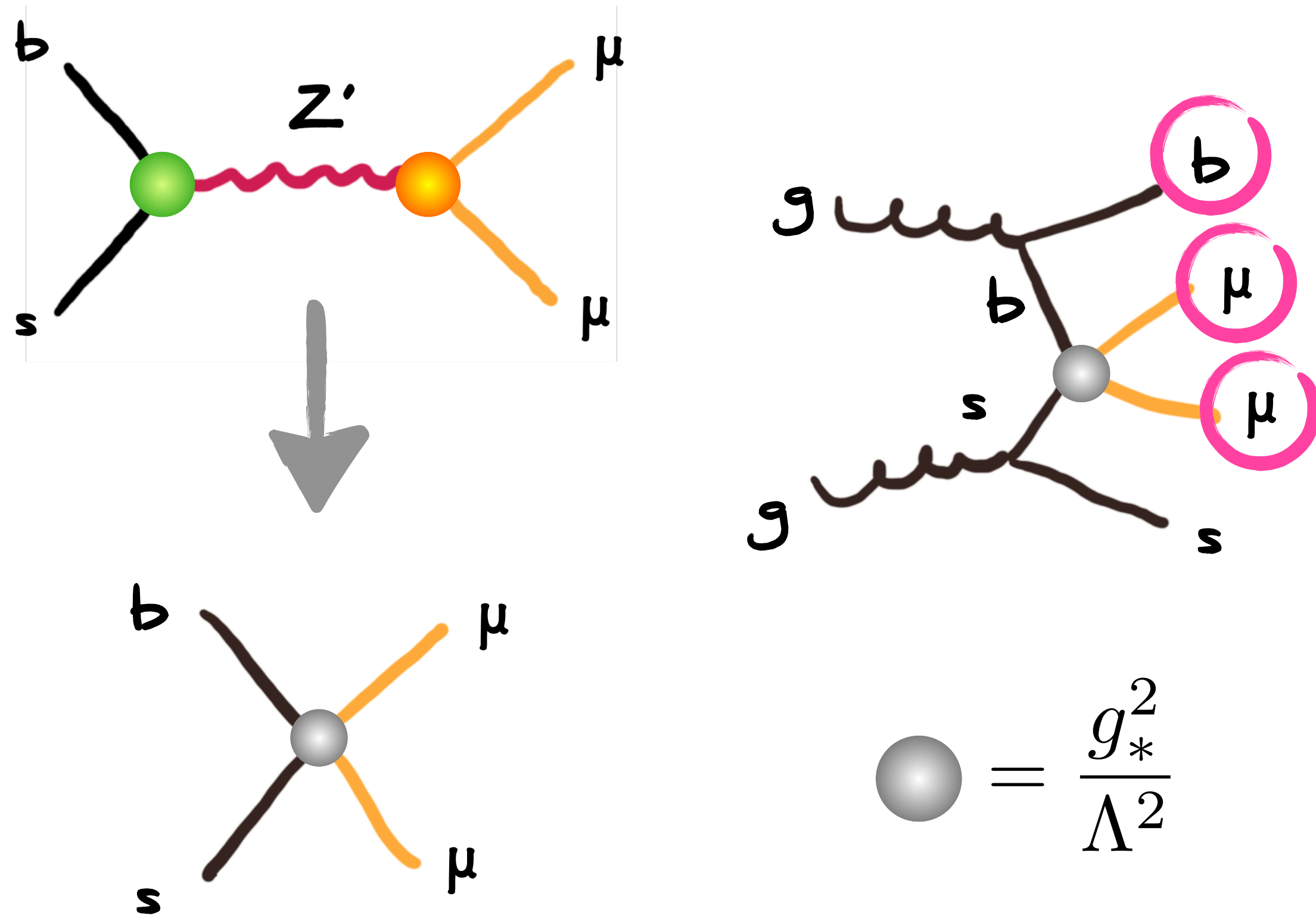
Suppression of operators suggests that explanations of $b \rightarrow c$ anomalies should lead to testable high- p_T signals, while $b \rightarrow s$ case looks grim

A digression on LQs



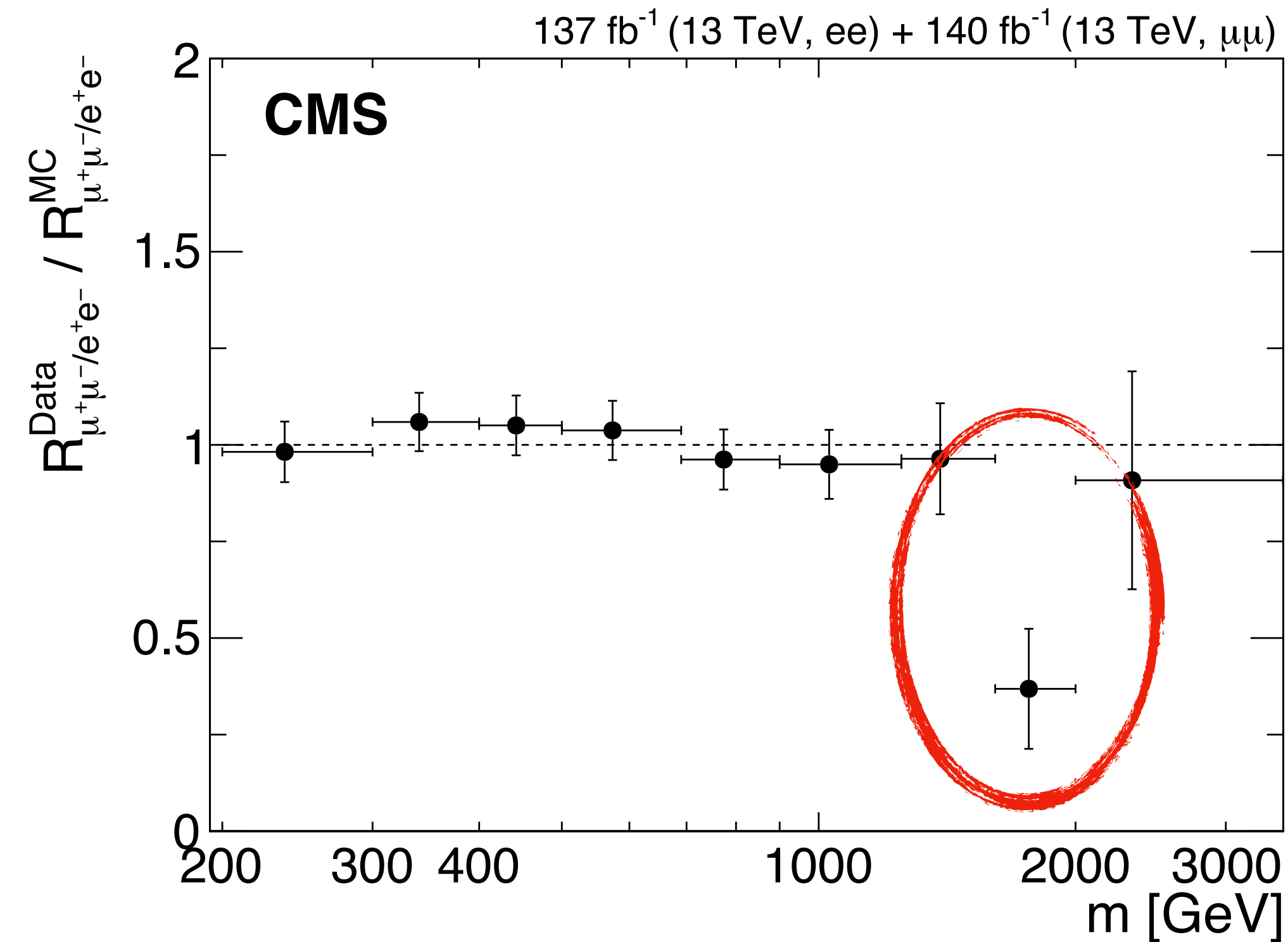
Both scalar & vector LQ have important advantage with respect to other tree-level mediators that they do not induce tree-level contributions to B mixing & $\tau \rightarrow \mu\nu\nu$

Searches for $bs\mu^+\mu^-$ contact interactions



First search for $bs\mu^+\mu^-$ four-Fermi operator by ATLAS, but bounds on suppression scale are a factor of $O(20)$ below sensitivity needed to test $b \rightarrow s$ anomalies model independently

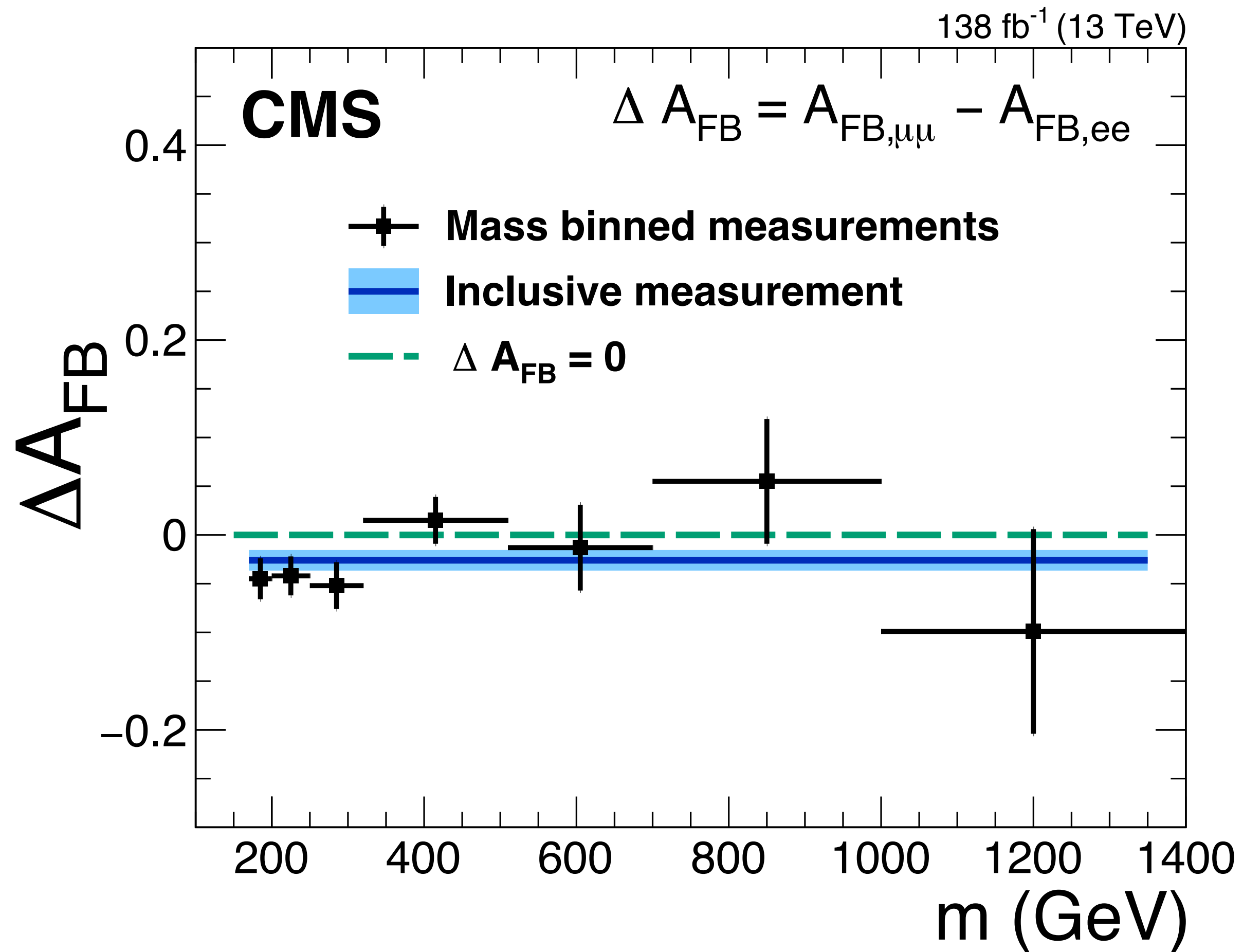
Testing LFU with dilepton events @ LHC



CMS observes good agreement with LFU up to masses of 1.5 TeV, but above 1.8 TeV there is slight excess in dielectron channel leading to a deviation of LFU ratio from 1

[CMS, 2103.02708 & for interpretations see for instance Crivellin et al., 2103.12003, 2104.06417]

Testing LFU with dilepton events @ LHC



CMS recently also measured difference between dimuon & dielectron forward-backward asymmetry (A_{FB}). Result is found to agree with zero within 2.4σ . Like rate measurement, also A_{FB} results show a slight dielectron excess