



# Classical and exotic spectroscopy at LHCb

#### **Elisabeth Niel on behalf of the LHCb Collaboration**

École Polytechnique Fédérale de Lausanne - EPFL Les Rencontres de Physique de la Vallée d'Aoste La Thuile - Italy 5 – 11 March 2023

Mont Blanc from my home

### Introduction

Classical hadrons  $\rightarrow$  mesons [quark + antiquark] or baryons [3 quarks]

Exotic hadrons  $\rightarrow$  in principle, anything else : Glueballs, hybrids, tetraquarks, pentaquarks, hexaquarks

Studied by different experiments:

LHCb, BESIII, ATLAS, CMS, Belle, Belle II, BaBar, CDF, D0, ALICE ...







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#### Today: a selection of recent LHCb results

### Last year LHCb talk – La Thuile 2022

### **Overview**

Selection of recent LHCb results on classical and exotic spectroscopy

#### **Conventional hadrons**

- $\Lambda_c^+ \rightarrow p K^- \pi^+$  amplitude analysis &  $\Lambda_c^+$  polarisation measurement (**NEW!**)
- Observation of new excited  $\Xi_b^0$  states in  $\Lambda_b^0 K^- \pi^+$
- Observation of excited  $\Omega_c^0$  baryons in  $\Omega_b^- \to \Xi_c^+ K^- \pi^-$  decays
- Study of charmonium contributions in  $B^+ \rightarrow J/\psi \, \eta K^+$

#### **Exotics**

- $\chi_{c1}$  (3872) production in *pp* collisions at  $\sqrt{s} = 8, 13 \text{ TeV}$
- Observation of exotic tetraquark  $T_{cc}^+$  in  $D^0 D^0 \pi^+$
- Evidence of new pentaquark structure in  $B_s^0 \rightarrow p\bar{p}J/\psi$  decays

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Selection of recent LHCb results on classical and exotic spectroscopy

#### **Conventional hadrons**

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**Conventional hadrons for today** 

 $\succ \Lambda_c^+ \rightarrow p K^- \pi^+$  polarimetry

$$\succ D_s^+ \to \pi^+ \pi^- \pi^+$$

$$\succ D^+ \rightarrow \pi^+ \pi^- \pi^+$$

- Observation of new  $\Omega_c^0$  states decaying into  $\Xi_c^+ K^-$
- ➢ Observation of  $\mathcal{Z}_{cc}^{++}$  →  $\mathcal{Z}_{c}^{\prime+}\pi^{+}$  [talk by Gabriele Martelli]

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### **Overview**

### Exotics for today

- > Observation of the  $B_s^0 \rightarrow (\chi_{c1}(3872) \rightarrow J/\psi \pi^+ \pi^-)\pi^+\pi^- \text{ decay}$
- >  $B^0 \rightarrow J/\psi \phi K_S^0$  new state :  $T^{\theta}_{\psi s1}(4000)^0$  and  $T_{\psi s1}(4220)^0$
- New doubly charged and neutral open charmed tetraquarks observed in
   B<sup>0</sup> → D
   <sup>0</sup>D<sup>+</sup><sub>s</sub>π<sup>-</sup> and B<sup>+</sup> → D<sup>-</sup>D<sup>+</sup><sub>s</sub>π<sup>+</sup>
   New state from the amplitude analysis of
- ➢ New state from the amplitude analysis of  $B^+ → D_s^+ D_s^- K^+$

> Strange pentaquark candidate in  $B^- \rightarrow J/\psi \Lambda \, \overline{p}$ 

al and exotic spectroscopy

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### The LHCb experiment

#### [IJMPA 30 (2015) 1530022] [JINST 3 (2008) S08005]



Single arm forward spectrometer with excellent vertexing, tracking, PID

# Conventional spectroscopy at LHCb

#### heavy baryons:

- $\Omega_c^0 \to \Xi_c^+ K^-$
- doubly heavy baryons
  - Observation of  $\Xi_{cc}^{++} \to \Xi_{c}^{\prime+}\pi^+$
- > amplitude structures of charm hadron decays
  - $D_s^+ \to \pi^+ \pi^- \pi^+$
  - $D^+ \rightarrow \pi^+ \pi^- \pi^+$
  - $\Lambda_c^+ \to p K^- \pi^+$  polarimetry
  - Observation of the  $B_s^0 \rightarrow (\chi_{c1}(3872) \rightarrow J/\psi \pi^+ \pi^-)\pi^+\pi^-$  decay

- > Single charmed baryons: one charm quark + 2 lighter quarks  $\rightarrow$  large mass difference
- > Described by heavy quark effective theory HQET but masses and quantum numbers diverges among different theories, some examples:
  - o lattice quantum chromodynamics predicts invariant-mass spectrum with D or F-wave excited states PRL 119, 042001
  - baryon-meson molecular (quasi-bound) states interpretation for  $\Omega_c^0(3050)$  and  $\Omega_c^0(3090)$  PRD 97 (2018) 094035, EPJ. A54 (2018) 64, Few Body Syst. 61 (2020) 34
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- > NEW update of PRL 118 (2017) 182001  $\rightarrow$  5 times larger data sample



Previous LHCb measurement

Resonance	Mass (MeV)	$\Gamma$ (MeV)
$\Omega_{c}(3000)^{0}$	$3000.4 \pm 0.2 \pm 0.1^{+0.3}_{-0.5}$	$4.5\pm0.6\pm0.3$
$\Omega_{c}(3050)^{0}$	$3050.2 \pm 0.1 \pm 0.1 \substack{+0.3 \\ -0.5}$	$0.8\pm0.2\pm0.1$
	0.5	<1.2 MeV, 95% C.L.
$\Omega_{c}(3066)^{0}$	$3065.6 \pm 0.1 \pm 0.3^{+0.3}_{-0.5}$	$3.5\pm0.4\pm0.2$
$\Omega_{c}(3090)^{0}$	$3090.2 \pm 0.3 \pm 0.5^{+0.3}_{-0.5}$	$8.7\pm1.0\pm0.8$
$\Omega_{c}(3119)^{0}$	$3119.1 \pm 0.3 \pm 0.9^{+0.3}_{-0.5}$	$1.1\pm0.8\pm0.4$
	-0.5	<2.6 MeV, 95% C.L.
$\Omega_{c}(3188)^{0}$	$3188 \pm 5 \pm 13$	$60\pm15\pm11$
$\Omega_{c}(3066)^{0}_{fd}$		
$\Omega_{c}(3090)_{fd}^{0}$		
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- Measurement of BR would help to discriminate internal structure (molecular, pentaquarks,...)
- > NEW update of PRL 118 (2017) 182001  $\rightarrow$  5 times larger data sample
- ≻ Production from  $\Omega_b^- \to \Xi_c^+ K^- \pi^-$  also studied by LHCb <u>Phys. Rev. D 104, L091102</u>



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> Two new excited states  $\Omega_c(3185)^0$  and  $\Omega_c(3327)^0$ 

≻ Five states from previous analysis confirmed → masses and widths measured with highest precision to date



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Partially reconstructed decays with photons determined from simulation (yields vary in the fit)

- > Fit at threshold enanchement with BW with or w/o feed down  $\Omega_c(3065)^0$ , not possible to separate two states but the existence of another hidden state cannot be excluded.
- $\geq \Omega_c(3185)^0$  could be also described by two states

$\Omega_c(3000)$	$())^0 \rightarrow \Xi_c^+ K^-$	$\Omega_c(3065)^0 \to \Xi_c^{+}(\to \Xi_c^+ \gamma) K^-$
$\ldots \Omega_c(3050)$	$(0)^0 \rightarrow \Xi_c^+ K^-$	$\Omega_c(3090)^0 \to \Xi_c^{+}(\to \Xi_c^+ \gamma) K^-$
$\Omega_c(3065)$	$(5)^0 \rightarrow \Xi_c^+ K^-$	$\Omega_c(3119)^0 \to \Xi_c^{+}(\to \Xi_c^+ \gamma) K^-$
$\Omega_{c}(3090)$	$\tilde{D}^{0} \rightarrow \Xi^{+} K^{-}$	$\Omega_c(3185)^0 \to \Xi_c^+ K^-$
$\Omega_{c}(3119)$	$\tilde{D}^{0} \rightarrow \Xi^{+} K^{-}$	$\Omega_c(3327)^0 \rightarrow \Xi_c^+ K^-$
	,	
Resonance	$m \; ({ m MeV})$	$\Gamma ~({ m MeV})$
$\Omega_c(3000)^0$	$3000.44 \pm 0.07 \ ^{+0.07}_{-0.13} \pm 0.23$	$3.83 \pm 0.23 \stackrel{+1.59}{_{-0.29}}$
$\Omega_c(3050)^0$	$3050.18 \pm 0.04  {}^{+0.06}_{-0.07} \pm 0.23$	$0.67 \pm 0.17 \ \substack{+0.64 \\ -0.72}$
		$< 1.8\mathrm{MeV}, 95\%$ C.L.
$\Omega_c(3065)^0$	$3065.63 \pm 0.06  {}^{+0.06}_{-0.06} \pm 0.23$	$3.79 \pm 0.20  {}^{+0.38}_{-0.47}$
$\Omega_c(3090)^0$	$3090.16 \pm 0.11  {}^{+0.06}_{-0.10} \pm 0.23$	$8.48 \pm 0.44  {}^{+0.61}_{-1.62}$
$\Omega_c(3119)^0$	$3118.98 \pm 0.12  {}^{+0.09}_{-0.23} \pm 0.23$	$0.60 \pm 0.63 \ ^{+0.90}_{-1.05}$
		$< 2.5 \mathrm{MeV}, 95\%$ C.L.
$\Omega_c(3185)^0$	$3185.1 \pm 1.7 \stackrel{+1.4}{-0.9} \pm 0.2$	$50 \pm 7 \begin{array}{c} +10 \\ -20 \end{array}$
$\Omega_{c}(3327)^{0}$	$3327.1 \pm 1.2  {}^{+0.1}_{-1.3} \pm 0.2$	$20 \pm 5 \begin{array}{c} +13 \\ -1 \end{array}$



# Amplitude analysis $D^+(D_s^+) \rightarrow \pi^+\pi^-\pi^+$

- Spectroscopy of scalar states, test glueballs interpretation J. Phys. G: Nucl. Part. Phys. 40 043001
- > 2012 data,  $\mathcal{L} = 0.75$  fb<sup>-1</sup>. Promptly produced D mesons:  $N(D^+) \sim 600k$  and  $N(D_s^+) \sim 700k$
- > Amplitude description: S-wave with Quasi-Model Independent approach (QMIPWA) and isobar model for spin-1, spin-2 components

$$\mathcal{A}_{S}(s_{12}, s_{13}) = \mathcal{A}_{S}(s_{12}) + \mathcal{A}_{S}(s_{13}) \qquad \qquad \mathcal{A}_{S}^{k}(s_{\pi^{+}\pi^{-}}) = c_{k}e^{i\phi_{k}}$$

 $c_k, \phi_k$ : Generic functions determined by fit to data

Component	Magnitude	Phase [°]		Fit fra	action [%	6]
$\rho(770)^0\pi^+$	1 [fixed]	0 [fixed]	26.0	$\pm 0.3$	$\pm 1.6$	$\pm 0.3$
$\omega(782)\pi^+$	$(1.68 \pm 0.06 \pm 0.15 \pm 0.02) \times 10^{-2}$	$-103.3 \pm 2.1 \pm 2.6 \pm 0.4$	0.103	$3 \pm 0.003$	$8 \pm 0.014$	$4 \pm 0.002$
$ ho(1450)^0\pi^+$	$2.66 \pm 0.07 \pm 0.24 \pm 0.22$	$47.0 \pm 1.5 \pm 5.5 \pm 4.1$	5.4	$\pm 0.4$	$\pm 1.3$	$\pm 0.8$
$ ho(1700)^0\pi^+$	$7.41 \pm 0.18 \pm 0.47 \pm 0.71$	$-65.7 \pm 1.5 \pm 3.8 \pm 4.6$	5.7	$\pm 0.5$	$\pm 1.0$	$\pm 1.0$
$f_2(1270)\pi^+$	$2.16 \pm 0.02 \pm 0.10 \pm 0.02$	$-100.9 \pm 0.7 \pm 2.0 \pm 0.4$	13.8	+0.2	+0.4	+ 0.2
S-wave			61.8	$\pm 0.5$	$\pm 0.6$	$\pm 0.5$
$\sum_{i} FF_{i}$				1	12.8	
$\chi^2/\text{ndof}$ (range)	[1.47 - 1.78]			$-2\log L$	$\mathcal{C} = 8056$	522

Dominated by S-wave, followed by  $\rho(770)^0\pi^+$  and  $f_2(1270)^0\pi^+$ Contribution from  $(\omega(782) \rightarrow \pi^+\pi^-)\pi^+$  observed for the first time



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 $c_k, \phi_k$ : Generic functions determined by fit to data

Resonance	Magnitude	Phase $[^{\circ}]$	Fit fraction (FF) $[\%]$
S-wave			$84.97\pm0.14$
$ ho(770)^{0}$	$0.1201 \pm 0.0030$	$79.4 \pm 1.8$	$1.038 \pm 0.054$
$\omega(782)$	$0.04001 \pm 0.00090$	$-109.9 \pm 1.7$	$0.360\pm0.016$
$ ho(1450)^{0}$	$1.277\pm0.026$	$-115.2\pm2.6$	$3.86 \pm 0.15$
$ ho(1700)^{0}$	$0.873 \pm 0.061$	$-60.9\pm6.1$	$0.365\pm0.050$
combined	-	_	$6.14\pm0.27$
$f_2(1270)$	1  (fixed)	0 (fixed)	$13.69\pm0.14$
$f'_2(1525)$	$0.1098 \pm 0.0069$	$178.1\pm4.2$	$0.0455 \pm 0.0070$
sum of fit fractions			104.3
$\chi^2/\text{ndof}$ (range)	[1.45 - 1.57]		

Dominated by S-wave, followed by spin-2 resonances Contribution from  $(\omega(782) \rightarrow \pi^+\pi^-)\pi^+$ observed for the first time

 $D_s^+ \rightarrow \pi^- \pi^+ \pi^+$ 

# $\Lambda_c^+ \rightarrow p K^- \pi^+$ polarimetry

- > Polarization measurement: insights on quark hadronisation mechanism, hadron spectroscopy, BSM searches in  $\Lambda_b \rightarrow \Lambda_c l^- \nu$  and with the measurement of EDM and MDM of charm baryons.
- $\geq \text{Recent LHCb } \Lambda_c^+ \text{ polarization measurement in } \Lambda_c^+ \rightarrow p \ K^- \pi^+ \text{ semileptonic decays and amplitude analysis:} \\ 6 \ \Lambda^*, 3 \ \Delta^{*++} \text{ states, and } 3 \ K^{*0} \text{ states} \qquad \underbrace{\text{arXiv:} 2208.03262}_{\text{arXiv:} 2208.03262}$
- $\triangleright$  Decay paremetrization is model dependent  $\rightarrow$  new method to express the polarized decay rate in model agnostic way

Reduction of uncertainties on polarization measurement!



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unpolarized intensity aligned polarimeter vector  

$$|\mathcal{M}(\phi,\theta,\chi,\kappa)|^{2} = I_{0}(\kappa) \left(1 + \sum_{i,j} P_{i}R_{ij}(\phi,\theta,\chi)\alpha_{j}(\kappa)\right)$$

>  $\alpha_j(\tau)$  can be computed for any model and represented as an arrow over the x-z plane, where  $\kappa$  denotes the kinematic variables

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- >  $\alpha_j(\tau)$  can be computed for any model and represented as an arrow over the x-z plane, where  $\kappa$  denotes the kinematic variables
  - Length of polarization vector = color axis, > 0.5 over the Dalitz → significant contribution of both parity conserving and violating currents!
  - > Averaged on angles  $\rightarrow$  worse uncertainty on  $P_i$
  - Code available at <u>https://lc2pkpi-polarimetry.docs.cern.ch</u>

Observation of the 
$$B_s^0 \rightarrow (\chi_{c1}(3872) \rightarrow J/\psi \pi^+ \pi^-)\pi^+\pi^- \text{decay}$$
 19

arXiv:2302.02127

- Puzzling difference between the branching fractions  $B^+ \to \chi_{c1}(3872)K^+$  and  $B^0 \to \chi_{c1}(3872)K^0$
- Compact-tetraquark interpretation explaining also similar BR  $B_s^0 \rightarrow \chi_{c1}(3872)\phi$  and  $B^0 \rightarrow \chi_{c1}(3872)K^0$  Phys. Rev. D102 (2020) 034017
- The ratio of branching fractions measured:

$$\frac{\mathcal{B}\left(\mathrm{B}^{0}_{\mathrm{s}} \rightarrow \chi_{\mathrm{c1}}(3872)\pi^{+}\pi^{-}\right) \times \mathcal{B}\left(\chi_{\mathrm{c1}}(3872) \rightarrow \mathrm{J/\psi}\pi^{+}\pi^{-}\right)}{\mathcal{B}\left(\mathrm{B}^{0}_{\mathrm{s}} \rightarrow \psi(2\mathrm{S})\pi^{+}\pi^{-}\right) \times \mathcal{B}\left(\psi(2\mathrm{S}) \rightarrow \mathrm{J/\psi}\pi^{+}\pi^{-}\right)} = (6.8 \pm 1.1 \pm 0.2) \times 10^{-2}$$

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• The signal yields two-dimensional extended unbinned maximum-likelihood fit to the  $J/\psi \pi^+ \pi^- \pi^-$  and  $J/\psi \pi^+ \pi^+$  mass distributions

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#### Fit components:

- 1. Signal template  $B_s^0 \rightarrow (\chi_{c1}(3872) \rightarrow J/\psi \pi^+ \pi^-) \pi^+ \pi^-)$
- 2. J/ $\psi \pi \pi$  combination does not originate from a  $\chi_{c1}(3872)$  meson (no interference with first accounted for)
- 3. Random combinations of the  $\chi_{c1}(3872)$  state with a  $\pi + \pi \text{pair}$ ,
- 4. Polynomial background

```
\mbox{Parameter} \qquad B^0_s \rightarrow \chi_{c1}(3872) \pi^+\pi^- \quad B^0_s \rightarrow \psi(2S) \pi^+\pi
```

Ν	$155 \pm 23$	$1301 \pm 47$
$m_{\chi_{-1}(3872)} [\text{MeV}/c^2]$	$3871.57 \pm 0.09$	
$m_{\psi(2S)}$ [MeV/ $c^2$ ]	_	$3686.08 \pm 0.0'$
$m_{B_0^0}$ [MeV/ $c^2$ ]	5366.9	$97 \pm 0.23$
$s_{B_s^0}$	1.(	$06 \pm 0.03$
$s_{J/\psi\pi^+\pi^-}$	1.1	$2\pm0.03$

 $\mathcal{B}\left(\mathrm{B_{s}^{0}}
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ight)$ 

 $imes \mathcal{B} \left( \chi_{c1}(3872) 
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### arXiv:2302.02127

Phys. Rev. D102 (2020) 034017  $3.85 < m_{{\rm J}/\psi\pi^+\pi^-} < 3.90\,{\rm GeV}/c^2$ 



arXiv:2302.10629

Using PDG input for  $\mathcal{B}(B^0_s \to \psi(2S)\pi^+\pi^-) \times \mathcal{B}(\psi(2S) \to J/\psi\pi^+\pi^-)$ 

### Observation of the $B_s^0 \rightarrow \chi_{c1}(3872)\pi^+\pi^-$ decay

- > Mass spectra ( $\pi^+\pi^-$ ) system recoiling against the  $\chi_{c1}(3872)$  and  $\psi(2S)$  states obtained using the *sPlot technique* from mass fits
- From  $B_s^0 \rightarrow J/\psi \pi^+ \pi^-$  analysis  $f_0(980)$  and  $f_0(1500)$  contributions dominant confirmed Phys. Rev. D89 (2014) 092006
- > Describe the spectra with coherent sum of scalar  $f_0(980)$  and  $f_0(1500)$  + non resonant component

$$F(m) \propto n q p^{3} \left[ f \mathcal{A}_{f_{0}(980)}(m) + e^{i\varphi} \mathcal{A}_{f_{0}(1500)}(m) \right]^{2}$$
Phase space



- ➢ Fit parameter shared in the simultaneous fit
- → Di-pion mass spectrum from  $B_s^0 \rightarrow \psi(2S)\pi^+\pi^-$  decays decribed → see dominant contribution of 2 S-wave resonances
- ► Large component  $B_s^0 \rightarrow \chi_{c1}(3872) f_0(980)$  with significance exceeding 7 standard deviations.



### Exotic Spectroscopy at LHCb

#### Tetraquarks

- $B^+ \rightarrow J/\psi \phi K^+$  two states :  $T^{\theta}_{\psi s1}(4000)^+$  and  $T_{\psi s1}(4220)^0$
- New doubly charged and neutral open charmed tetraquarks observed in  $B^0 \to \overline{D}{}^0 D_s^+ \pi^-$  and  $B^+ \to D^- D_s^+ \pi^+$
- New state from the amplitude analysis of  $B^+ \rightarrow D_s^+ D_s^- K^+$

#### > Pentaquarks

• Strange pentaquark candidate in  $B^- \rightarrow J/\psi \Lambda \bar{p}$ 

## LHCb naming scheme

#### We proposed a new exotic hadron naming convention arXiv:2206.15233v1

- *T* for tetraquark
- *P* for pentaquark
- **superscript**: based on existing symbols, to indicate isospin, parity and G-parity
- **subscript**: heavy quark content

2	T states			T s	tates	
zero net $S, C, B$		n	non-zero net $S, C, B$			
(P,G)	I = 0	I = 1	(P)	I = 0	$I = \frac{1}{2}$	I = 1
(-,-)	$\omega$	$\pi$	(-)	$\eta$	au	$\pi$
(-,+)	$\eta$	ho	(+)	f	$\theta$	a
(+, +)	f	b				
(+, -)	h	a				

P states

I = 0	$I = \frac{1}{2}$	I = 1	$I = \frac{3}{2}$
Λ	$N^{-}$	$\boldsymbol{\Sigma}$	$\Delta^{-}$

Minimal quark content	Current name	$I^{(G)}, J^{P(C)}$	Proposed name	Reference
	$\chi_{c1}(3872)$	$I^G = 0^+, J^{PC} = 1^{++}$	$\chi_{c1}(3872)$	[24, 25]
$car{c}uar{d}$	$Z_c(3900)^+$	$I^G = 1^+, \ J^P = 1^+$	$T^b_{\psi 1}(3900)^+$	[26-28]
$car{c}uar{d}$	$X(4100)^+$	$I^{G} = 1^{-}$	$T_\psi(4100)^+$	[29]
$car{c}uar{d}$	$Z_c(4430)^+$	$I^G = 1^+, \ J^P = 1^+$	$T^{b}_{\psi 1}(4430)^{+}$	[30, 31]
$car{c}(sar{s})$	$\chi_{c1}(4140)$	$I^G = 0^+, J^{PC} = 1^{++}$	$\chi_{c1}^{-1}(4140)$	[32 - 35]
$car{c}uar{s}$	$Z_{cs}(4000)^+$	$I=rac{1}{2},\ J^P=1^+$	$T^{ heta}_{\psi s1}(4000)^+$	[7]
$car{c}uar{s}$	$Z_{cs}(4220)^+$	$I = \frac{1}{2}, \ J^P = 1^?$	$T_{\psi s1}(4220)^+$	[7]
$c\bar{c}c\bar{c}$	X(6900)	$I^G = \bar{0^+}, \ J^{PC} = ?^{?+}$	$T_{\psi\psi}(6900)$	[4]
$csar{u}ar{d}$	$X_0(2900)$	$J^P = 0^+$	$T_{cs0}(2900)^0$	[5,6]
$csar{u}ar{d}$	$X_1(2900)$	$J^{P} = 1^{-}$	$T_{cs1}(2900)^0$	[5,6]
$ccar{u}ar{d}$	$T_{cc}(3875)^+$		$T_{cc}(3875)^+$	[8,9]
$bar{b}uar{d}$	$Z_b(10610)^+$	$I^G = 1^+, \ J^P = 1^+$	$T^b_{\Upsilon 1}(10610)^+$	[36]
$car{c}uud$	$P_c(4312)^+$	$I = \frac{1}{2}$	$P_{\psi}^{N}(4312)^{+}$	[3]
$c\bar{c}uds$	$P_{cs}(4459)^0$	$I = \tilde{0}$	$P_{\psi s}^{^{ au}}(4459)^0$	[20]

# Evidence of $T_{\psi s1}^{\theta}$ (4000)<sup>0</sup> state

Previously  $B^+ \to J/\psi \phi K^+$  we observed two states :  $T^{\theta}_{\psi s1}(4000)^+$  and  $T_{\psi s1}(4220)^+$  Phys. Rev. Lett. 127, 082001 Here we looked for the isospin partner decay, is there a  $T^{\theta}_{\psi s1}(4000)^0$  isospin partner in  $B^0 \to J/\psi \phi K^0_S$ ?

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Similar analysis: reconstruct pairs of muons, kaons and pions ( + require long flight distance on the  $K_S^0$ )

A simultaneous fit on the two modes is performed.





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Similar analysis: reconstruct pairs of muons, kaons and pions ( + require long flight distance on the  $K_S^0$ )

A simultaneous fit on the two modes is performed.



Default model:

 $\succ K^* \text{ resonances in } (\phi K_S^0) : 9 \text{ in total}$  $\succ X (\chi_{c 0,1}, \eta_c \text{ or } T_{\psi \phi 1}^{\eta}) \text{ in } (J/\psi \phi) : 7 \text{ in total}$  $\succ T_{\psi s 1}^{\theta} \text{ in } (J/\psi K_S^0)$ 



Assumption of isospin symmetry  $\rightarrow$  mass, width and helicity couplings for all the components except  $T_{\psi s1}^{\theta}(4000)^{0}$  constraint to be identical.

#### $T_{\psi s1}(4220)^0$ constraints to be identical to $T_{\psi s1}(4220)^+$ due to limited size of $B^0$ sample

# Evidence of $T_{\psi s1}^{\theta}$ (4000)<sup>0</sup> state



> Difference in mass with  $T_{\psi s1}^{\theta}(4000)^+$  is small  $\rightarrow$  confirm the isospin partnership  $\Delta M = -12 \frac{+11}{-10} \frac{+6}{-4} \text{ MeV}$ 

Significance computed with likelihood ratio method with :

 $t \equiv -2\ln[\mathcal{L}(H_0)/\mathcal{L}(H_1)]$   $H_{0,1}$ : default model without or with the new state

> Significance is : 4  $\sigma$ , goes up to 5.4  $\sigma$  assuming isospin symmetry for the  $T_{\psi s1}^{\theta}(4000)$  states

### Doubly charged tetraquark

 $\succ$  New doubly charged and neutral open charmed tetraquarks observed in  $B^0 \rightarrow \overline{D}{}^0 D_s^+ \pi^-$  and  $B^+ \rightarrow D^- D_s^+ \pi^+$  arXiv:2212.02717v1

▷ Predicted by diquark-antiquark model from X(2900) in  $B \rightarrow DD^+K^-$ 

Name	Decay	Content
$T^{a}_{c\bar{s}0}(2900)^{0}$	$B^0 \to \overline{D}{}^0 D_s^+ \pi^-$	csūd
$T^a_{c\bar{s}0}(2900)^{++}$	$B^+ \rightarrow D^- D_s^+ \pi^+$	csūud

#### Model:

- > Only expect  $\overline{D^*}$  resonances
- > Spin parity of  $\overline{D}(3000)^0$  determined : 4<sup>+</sup>
- $\triangleright$   $\overline{D}\pi$  S-wave with quasi Model Independent splines



### Expected resonant content

Resonance	$J^P$	Mass (GeV)	Width $(GeV)$	Comments
$\overline{D}^*(2007)^0$	$1^{-}$	$2.00685 \pm 0.00005$	$< 2.1  imes 10^{-3}$	Width set to be $0.1\mathrm{MeV}$
$D^{*}(2010)^{-}$	$1^{-}$	$2.01026 \pm 0.00005$	$(8.34 \pm 0.18)  imes 10^{-5}$	
$\overline{D}_0^*(2300)$	$0^+$	$2.343 \pm 0.010$	$0.229 \pm 0.016$	#
$\overline{D}_{2}^{*}(2460)$	$2^{+}$	$2.4611 \pm 0.0007$	$0.0473 \pm 0.0008$	#
$\overline{D}_{1}^{*}(2600)^{0}$	$1^{-}$	$2.627 \pm 0.010$	$0.141 \pm 0.023$	#
$\overline{D}_3^*(2750)$	$3^{-}$	$2.7631 \pm 0.0032$	$0.066 \pm 0.005$	#
$\overline{D}_{1}^{*}(2760)^{0}$	$1^{-}$	$2.781 \pm 0.022$	$0.177 \pm 0.040$	#
$\overline{D}_{J}^{*}(3000)^{0}$	$?^{?}$	$3.214 \pm 0.060$	$0.186 \pm 0.080$	# $J^P = 4^+$ is assumed



### Doubly charged tetraquark

 $B^+ \rightarrow D^- D_s^+ \pi^+$ 

Simultaneous fit of both channels with isospin relation between two states enforced  $\rightarrow$  Model with  $T^a_{c\bar{s}0}(2900)^0$  and  $T^a_{c\bar{s}0}(2900)^{++}$  preffered

 $\blacktriangleright$  Fit with only D\* resonances any spin combinations discarded, fit with different masses for the two states compatible results

3.2 2.2 2.4 2.6 2.8 3.0 2.6 2.8 3.0 3.2 3.4 2.2 2.4 $M(D_s^+\pi^+)$  (GeV)  $B^0$  Fraction (%) Particle Amplitude Phase  $B^+$  Fraction (%)  $T^{a}_{c\bar{s}0}(2900)$  $0.149 \pm 0.031 \pm 0.031$  $-1.26 \pm 0.22 \pm 0.35$  $2.55 \pm 0.64 \pm 0.83$  $2.45 \pm 0.65 \pm 0.84$  $D^*(2007)^0$  $2.58 \pm 0.11 \pm 1.07$  $-3.01 \pm 0.06 \pm 0.31$  $14.0 \pm 1.1 \pm 2.7$  $D^{*}(2010)^{-}$  $17.0 \pm 1.0 \pm 2.4$  $3.05 \pm 0.11 \pm 0.48$  $-2.91 \pm 0.06 \pm 0.28$  $D_2^*(2460)$  $22.35 \pm 0.76 \pm 0.74$  $22.53 \pm 0.74 \pm 0.54$  $D_1^*(2600)$  $0.218 \pm 0.030 \pm 0.051$  $0.13 \pm 0.16 \pm 0.22$  $1.28 \pm 0.39 \pm 0.60$  $1.32 \pm 0.38 \pm 0.59$  $D_3^*(2750)$  $0.153 \pm 0.032 \pm 0.040$  $-2.80 \pm 0.19 \pm 0.60$  $0.32 \pm 0.15 \pm 0.21$  $0.33 \pm 0.14 \pm 0.20$  $D_1^*(2760)$  $0.119 \pm 0.044 \pm 0.153$  $-0.18 \pm 0.34 \pm 1.01$  $0.26 \pm 0.27 \pm 1.37$  $0.28 \pm 0.26 \pm 1.35$  $D_{J}^{*}(3000)$  $1.44 \pm 0.23 \pm 1.15$  $1.40 \pm 0.23 \pm 1.33$  $0.45 \pm 0.16 \pm 0.34$  $0.46 \pm 0.15 \pm 0.33$  $D\pi$  S-wave  $1.142 \pm 0.045 \pm 0.083$  $-0.972 \pm 0.045 \pm 0.084$  $44.9 \pm 1.9 \pm 3.6$  $48.3 \pm 1.8 \pm 3.5$ 

Common mass and widths is:

 $M = 2.908 \pm 0.011 \pm 0.020 \, \text{GeV}$  $\Gamma = 0.136 \pm 0.023 \pm 0.013 \,\text{GeV}$ 

30

arXiv:2212.02717v1

 $T^a_{c\bar{c}0}(2900)^0: M = (2.892 \pm 0.014 \pm 0.015) \,\text{GeV}$  $\Gamma = (0.119 \pm 0.026 \pm 0.013) \,\text{GeV}$  $T^a_{c\bar{s}0}(2900)^{++}: M = (2.921 \pm 0.017 \pm 0.020) \,\text{GeV}$  $\Gamma = (0.137 \pm 0.032 \pm 0.017) \,\text{GeV}$ 

Significance estimated from  $2\Delta LL$ 

New state has 9  $\sigma$ Spin parity: 7.5  $\sigma$  0<sup>+</sup> vs 1<sup>-</sup>



➢ Branching ratio measured:

 $\frac{\mathcal{B}(B^+ \to D_s^+ D_s^- K^+)}{\mathcal{B}(B^+ \to D^+ D^- K^+)} = 0.525 \pm 0.033 \pm 0.027 \pm 0.034, \qquad \text{arXiv:2211.05034 accepted by PRD}$ 

> Observed state: X(3960) with  $J^P = 0^{++}$  close to  $D_s^+ D_s^-$  threshold



> Branching ratio measured:

 $\frac{\mathcal{B}\left(B^{+} \to D_{s}^{+} D_{s}^{-} K^{+}\right)}{\mathcal{B}\left(B^{+} \to D^{+} D^{-} K^{+}\right)} = 0.525 \pm 0.033 \pm 0.027 \pm 0.034,$ 

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- > Observed state: X(3960) with  $J^{P}=0^{++}$  close to  $D_{s}^{+}D_{s}^{-}$  threshold
- ▷ Previously measured  $\chi_{c0}(3930)$  in  $D^+D^-$  does not fit in  $\chi_{c0}(2P)$  [4131–4292 MeV ] or  $\chi_{c0}(3P)$ [3842–3868 MeV ] spectrum.
- > If  $\chi_{c0}(3930)$  and X(3960) are the same particle:

$$\frac{\Gamma(X \to D^+ D^-)}{\Gamma(X \to D_s^+ D_s^-)} = \frac{\mathcal{B}^{(1)} \mathcal{F}_X^{(1)}}{\mathcal{B}^{(2)} \mathcal{F}_X^{(2)}} = 0.29 \pm 0.09 \pm 0.10 \pm 0.08$$

> If X(3960) had no  $s\bar{s}$  content  $\Gamma(D^+D^-) >> \Gamma(D_s^+D_s^-)$  contradicting the ratio above

 $\rightarrow$ either not the same resonance or they are the same non-conventional charmonium-like state [candidate containing the  $c\overline{c}s\overline{s}$  dominant constituents]



arXiv:2210.15153 accepted by PRL



- ➤ X(4140) produces the dip around 4140 MeV via destructive interference with the 0<sup>++</sup> NR and X(3960) components → with (-22.4 ± 6.4)% and (-5.2 ± 3.9)% interference fractions
- → If the dip produced by the opening of the nearby  $J/\psi \phi$  threshold →test with **K matrix** with coupled channels  $J/\psi \phi$  and  $D_s D_s$

arXiv:2210.15153 accepted by PRL



- ➤ X(4140) produces the dip around 4140 MeV via destructive interference with the 0<sup>++</sup> NR and X(3960) components → with (-22.4 ± 6.4)% and (-5.2 ± 3.9)% interference fractions
- ► If the dip produced by the opening of the nearby  $J/\psi\phi$  threshold → test with K matrix with coupled channels  $J/\psi\phi$  and  $D_sD_s$
- $\succ$  K-matrix parameterisation similar fit quality as baseline model  $\rightarrow$  no strong conclusion whether the dip is due to:
  - destructive interference with the  $X_0(4140)$  resonance
  - or caused by the  $J/\psi \phi \rightarrow D_s^+ D_s^-$  rescattering.

Need more data to confirm!





#### arXiv:2210.15153 accepted by PRL

# Exotic Spectroscopy at LHCb

#### > Tetraquarks

- $B^+ \rightarrow J/\psi \phi K^+$  two states :  $T^{\theta}_{\psi s1}(4000)^+$  and  $T_{\psi s1}(4220)^0$
- New doubly charged and neutral open charmed tetraquarks observed in  $B^0 \to \overline{D}{}^0 D_s^+ \pi^-$  and  $B^+ \to D^- D_s^+ \pi^+$
- New state from the amplitude analysis of  $B^+ \rightarrow D_s^+ D_s^- K^+$

#### > Pentaquarks

• Strange pentaquark candidate in  $B^- \rightarrow J/\psi \Lambda \bar{p}$ 

# Strange pentaquark candidate $P_{\psi s}^{\Lambda}$

- ≻ Amplitude analysis of  $B^- \rightarrow J/\psi \Lambda \bar{p}$
- ➢ A new narrow pentaquark state observed
- $\succ$  Close to the mass threshold  $\Xi_c^+ D^-$
- ➢ Spin ½ and preferred parity odd
- $\succ$  First observation of pentaquark with strange content
- > Previous analysis by CMS JHEP12(2019)100  $[c\bar{c}uds]$





# Strange pentaquark candidate $P_{\psi s}^{\Lambda}$

 $\succ$  Fit with only the known K\* resonances

This model does not describe the data well  $\chi^2/ndf = 123/46$ 





2.1

-0.5



### Conclusions

>LHCb leading contributions on conventional and exotic spectroscopy shown today:

- $\checkmark$  New tetraquark and pentaquark candidates shown
- $\checkmark$  New singly and doubly heavy baryons found with high significance
- $\checkmark$  New method for polarization measurement of spin  $\frac{1}{2}$  baryons
- Spectroscopy of heavy hadrons is crucial to understand QCD dynamics and binding rules
   these are valuable inputs for the theory community
- Collaboration theory/experiment is crucial here: need predictions for masses, widths, lifetimes and model to describe/discover exotics.

LHCb Run 3 data will help access decays not yet observed, and more analysis are on-going with LHCb Run1 and Run2 data, so stay tuned!



# Stange pentaquark candidate $P_{\psi s}^{\Lambda}$

Fit with  $\overline{P}^{N}\psi^{-}$  contribution

This model is discarded because of an increase in -2 log L around 80



Efficiency corrected distributions

40

arXiv:2210.10346v1

### Observation of $\Xi_{cc}^{++} \rightarrow \Xi_{c}^{\prime+}\pi^{+}$

Search for  $\Xi_{cc}^{++} \rightarrow \Xi_{c}^{\prime+}\pi^{+}$  and measure branching fraction

- $\geq E_c^{\prime +} \rightarrow E_c^+ \gamma$ , partial reconstructed signal in  $m(E_c^+ \pi^+)$
- > Similar efficiency & comparable production as  $\mathcal{Z}_{cc}^{++} \rightarrow \mathcal{Z}_{c}^{+} \pi^{+}$

$$\geq \frac{B(\Xi_{cc}^{++} \to \Xi_{c}^{+} \pi^{+})}{B(\Xi_{cc}^{++} \to \Xi_{c}^{+} \pi^{+})}$$
 prediction vary between **0**. **3** ~ **7**

1<sup>st</sup> observation of  $\Xi_{cc}^{++} \rightarrow \Xi_{c}^{\prime+}\pi^{+}$  with 2016-2018 data,  $\mathcal{L} \sim 5.4 \text{ fb}^{-1} \rightarrow \text{statistical significance} > 9 \sigma$ 



### Amplitude analysis $D^+(D_s^+) \rightarrow \pi^+\pi^-\pi^+$



 $\overline{s}$ 

ab  $\rightarrow \pi + \pi -$  scattering



Indicating  $f_0(500)$  as dynamical pole of  $\pi\pi \to \pi\pi$  rescattering ?



 $(s\bar{s})$ 

 $\sum s\bar{q}_i q_i \bar{s} = K^+ K^- + K^0 \overline{K}{}^0 + \frac{1}{3}\eta\eta$