

Status of V_{us} determination and perspectives

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LNF-INFN

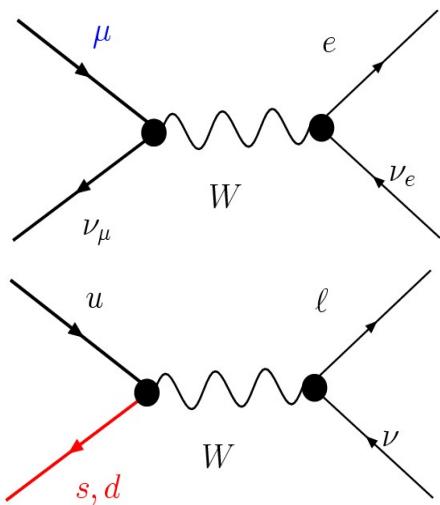
XVII SuperB Workshop and Kick Off Meeting

May 28-June 2, 2011, La Biodola(Isola d'Elba) Italy

1st raw unitarity: G_F universality

$$|V_{ud}|^2 + |V_{us}|^2 + \cancel{|V_{ub}|^2} \equiv 1$$

Universality of Weak coupling- $G_F = (g_w/M_w)^2$
 $G_F^2 \equiv G_{CKM}^2 = (|V_{ud}|^2 + |V_{us}|^2) G_F^2$



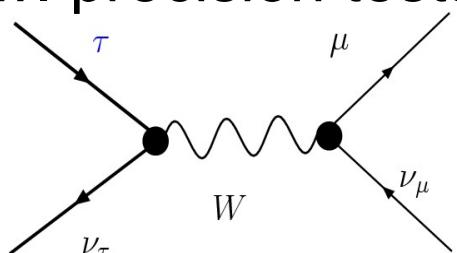
$$G_F = 1.166371(6) \times 10^{-5} \text{ GeV}^{-2}$$

$$G_{CKM} = 1.16633(35) \times 10^{-5} \text{ GeV}^{-2}$$

$$G_{ew} = 1.1655(12) \times 10^{-5} \text{ GeV}^{-2}$$

$$G_\tau = 1.1678(26) \times 10^{-5} \text{ GeV}^{-2}$$

$\alpha + M_w + S_w$
[e. w. precision tests]



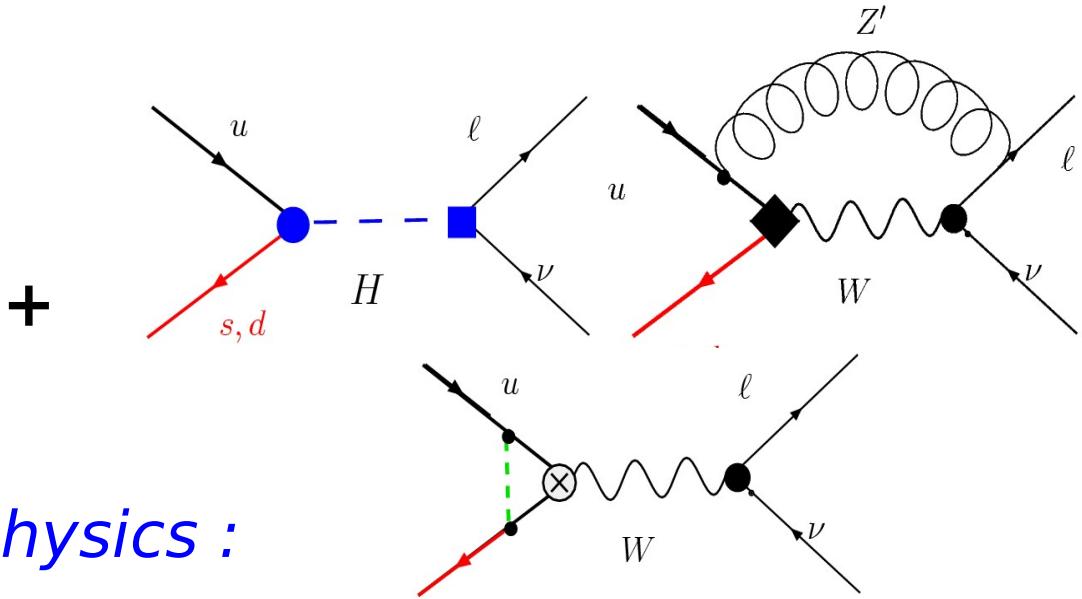
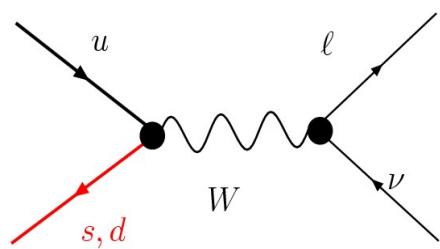
[Marciano]

G_F universality violation

$$|V_{ud}|^2 + |V_{us}|^2 + |\cancel{V_{ub}}|^2 \equiv 1$$

Universality of Weak coupling- $G_F = (g_w/M_w)^2$

$$G_F^2 \equiv G_{CKM}^2 = (|V_{ud}|^2 + |V_{us}|^2) G_F^2$$



Sensitivity to new physics :
naively

$$G_{CKM} = G_F [1 + a(M_w/M_M)^2]$$

Tree level $a \sim 1$

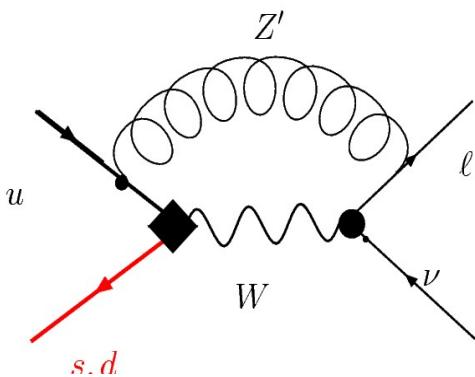
$M_M \sim 10 \text{ TeV}$

loops $a \sim g_w^2 / (16\pi^2)$

$M_M \sim 1 \text{ TeV}$

sensitivity to NP: Z'oology

1)



$$G_F = G_{CKM} [1 - 0.007 Q_e (Q_{\mu L} - Q_d) \frac{2 \ln(m_Z/m_W)}{(m_Z^2/m_W^2 - 1)}]$$

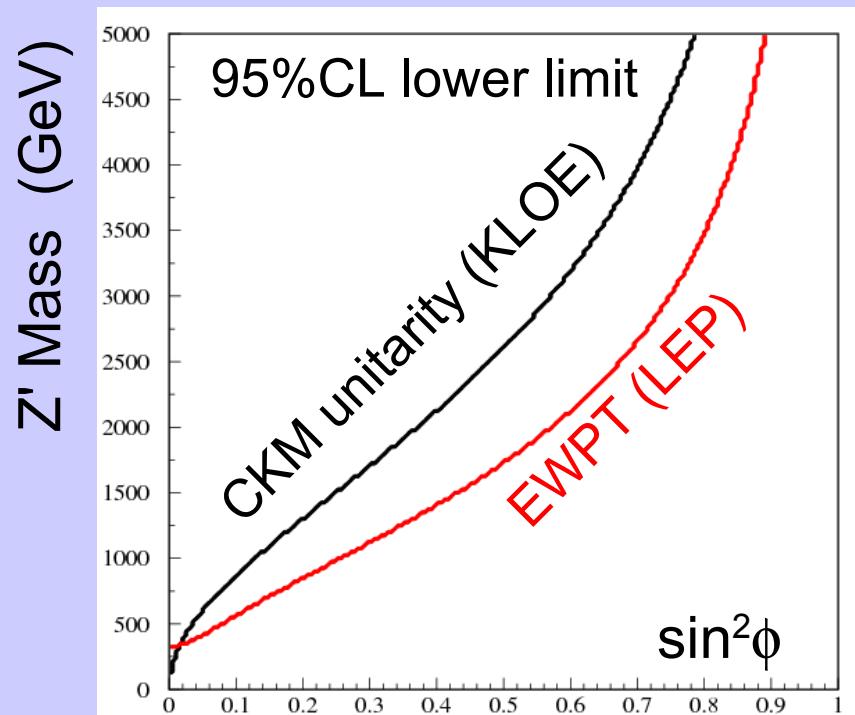
SO(10) Z_χ Boson: $Q_e = Q_{\mu L} = -3Q_d = 1$ [Marciano]

$m_{Z_\chi} > 750 \text{ GeV } 95\% \text{CL}$

2)

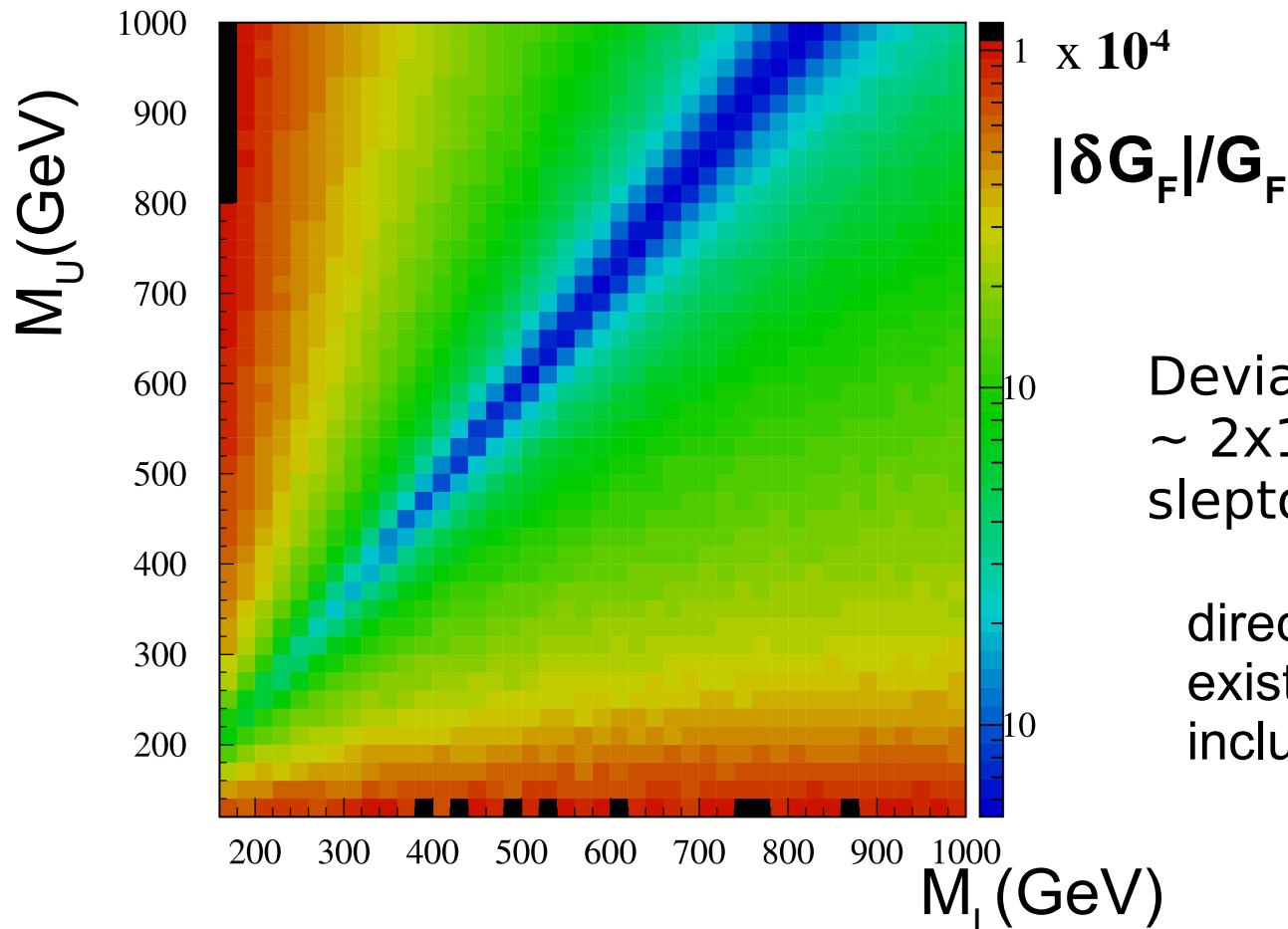
[K.Y. Lee]

Tree level breaking of unitarity in models with non-universal gauge interaction



sensitivity to NP: MSSM

sensitive to squark-slepton mass difference



[R. Barbieri '85,
K. Hagiwara et al
'95, A. Kurylov
et al '00]

Deviations up to
 $\sim 2 \times 10^{-4}$ for small
slepton mass

direct and indirect
existing limits
included

sensitivity to NP: charged Higgs

Pseudoscalar currents, e.g. due to H^\pm , affect the K width:

JHEP
0804:059

$$\frac{\Gamma(M \rightarrow \ell\nu)}{\Gamma_{SM}(M \rightarrow \ell\nu)} = \left[1 - \tan^2 \beta \left(\frac{m_{s,d}}{m_u + m_{s,d}} \right) \frac{m_M^2}{m_H^2} \right]^2 \quad \text{for } M = K, \pi$$

Hou, Isidori-Paradisi

The observable

$$R_{\ell 23} = \left| \frac{V_{us}(K_{\mu 2})}{V_{us}(K_{\ell 3})} \times \frac{V_{ud}(0^+ \rightarrow 0^+)}{V_{ud}(\pi_{\mu 2})} \right|$$

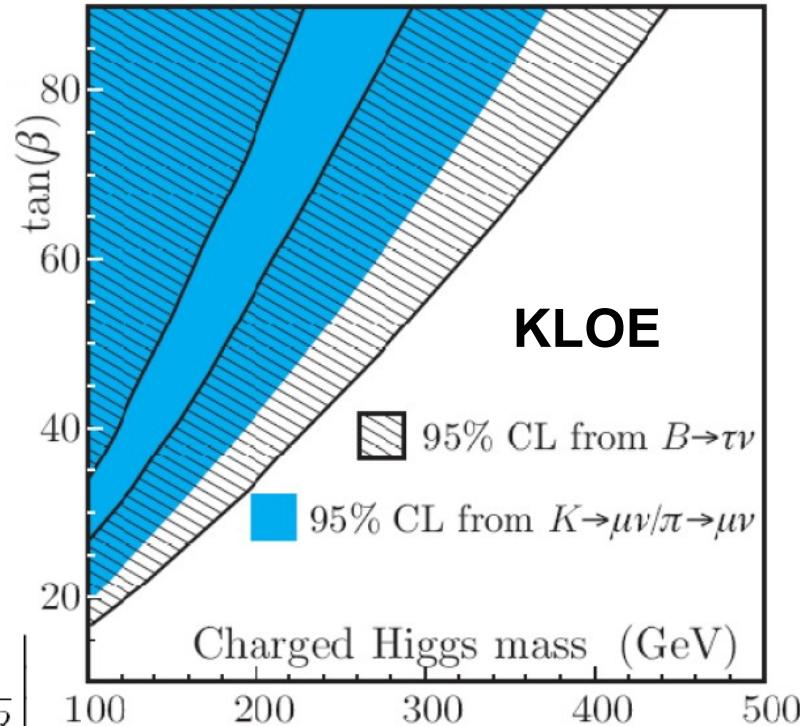
KLOE:

- $R_{\ell 23} = 1.008(8)$

(unitarity for K_B and β -decays is used)

$R_{\ell 23}$ sensitivity to H^\pm exchange

$$R_{\ell 23} = \left| 1 - \frac{m_{K^+}^2}{m_{H^+}^2} \left(1 - \frac{m_{\pi^+}^2}{m_{K^+}^2} \right) \frac{\tan^2 \beta}{1 + \epsilon_0 \tan \beta} \right|$$



Status of V_{ud} determination (superallowed β -decays only)

V_{ud} from Fermi transitions

$$V_{ud}^2 = \frac{K}{2G_F^2 \mathcal{F}t(1 + \Delta_R)}$$

$$\mathcal{F}t = ft(1 + \delta'_R)(1 - (\delta_C - \delta_{NS})) = \text{constant}$$

Measured on 13 Nuclei:

$t = t_{1/2}/\text{BR}$ = partial half life

f = statistical rate function $f(Z, Q_{ec})^*$

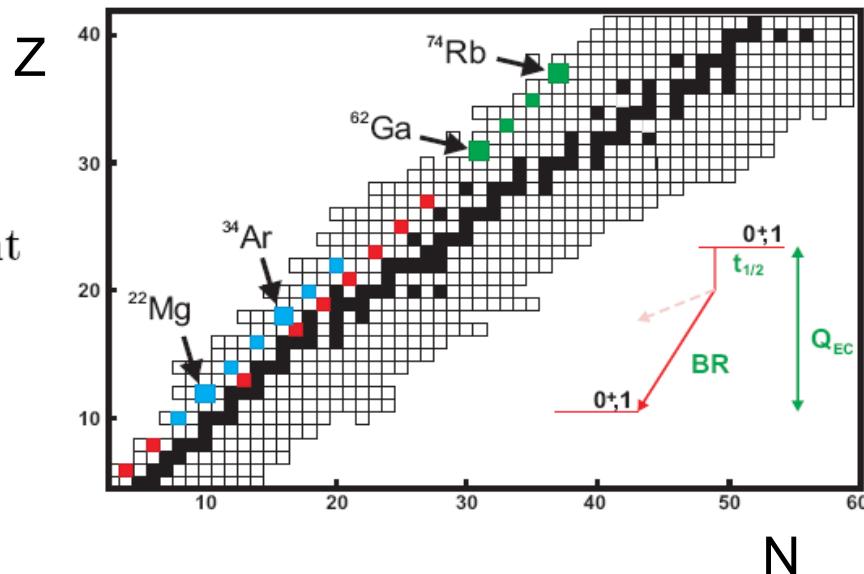
Radiative and isospin breaking corrections:

$\Delta_R = 2.361(38)\%$ Nucleus-independent

[Marciano Sirlin]

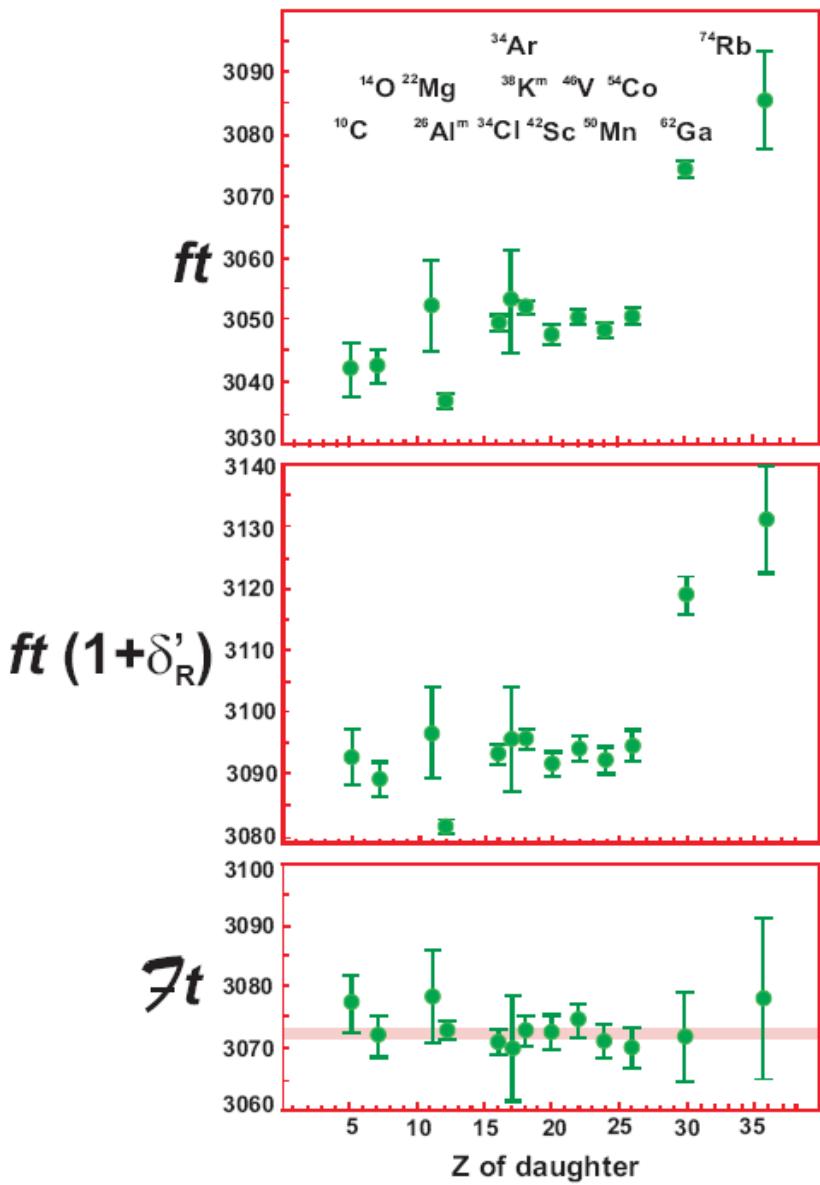
δ'_R, δ_{NS} Nucleus-dependent

δ_C Nucleus-dependent isospin breaking



* Z dependence account for e wave function

V_{ud} from Fermi transitions

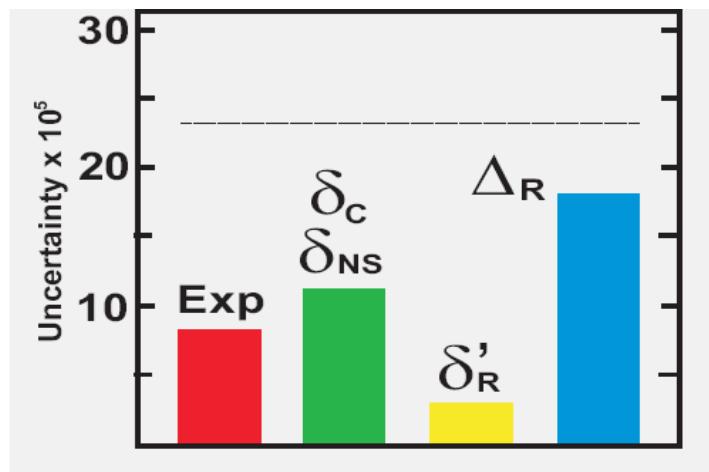


[Towner, Hardy
2008]

$$V_{ud}^2 = \frac{K}{2G_F^2 \mathcal{F}t(1 + \Delta_R)}$$

$$V_{ud} = 0.97425(23)$$

Error budget:



Status of V_{us} and V_{us}/V_{ud} determination

$\tau \rightarrow K\pi\nu$, $K_{\ell 3}$ decays

Vector transition protected against ~~SU(3)~~ corrections:

[Ademollo
Gatto]

$$\Gamma \propto M^{5(3)} S_{EW} G_F^2 |V_{us}|^2 |f_+^{K^0\pi^-}(0)|^2 I_K (1 + 2\Delta_K^{SU(2)} + 2\Delta_{K\ell}^{EM})$$

S_{EW} Universal SD EW correction (1.0232)

Inputs from theory:

$f_+^{K^0\pi^-}(0)$ Hadronic matrix element (form factor) at zero momentum transfer ($t = 0$)

$\Delta_K^{SU(2)}$ Form-factor correction for $SU(2)$ breaking

$\Delta_{K\ell}^{EM}$ Form-factor correction for long-distance EM effects

Inputs from experiment:

I_K Rates with well-determined treatment of radiative decays:

- Branching ratios
- lifetimes

I_K Integral of dalitz density (includes ff) over phase space:

- $\tau \rightarrow K\pi\nu, K_{\ell 3}$

$\tau \rightarrow P \nu$, $P_{\ell 2}$ decays ($P=K,\pi$)

$$\Gamma \propto G_F^2 |V_{uq}|^2 f_P^2 (1 + C_{P\ell})$$

Inputs from theory:

f_P decay constants

$C_{P\ell}$ Radiative inclusive
electroweak corrections

Inputs from experiment:

Γ Rates with well-determined
treatment of radiative decays:

- Branching ratios
- lifetimes

Used to determine pseudoscalar decay constants

Small uncertainties for ratios:

$\Gamma(K_{\mu 2(\gamma)})/\Gamma(\pi_{\mu 2(\gamma)})$ f_K/f_π from lattice \rightarrow determine V_{us}/V_{ud}
[Marciano]

$R_P = \Gamma(P_{e2(\gamma)})/\Gamma(P_{\mu 2(\gamma)})$ no f_P \rightarrow test lepton universality

[Cirigliano, Rosell]

$|V_{us}| f_+(0)$ from K_{l3} data

| | $ V_{us} f_+(0)$ | % err | BR | τ | Δ | Int |
|---------------|-------------------|-------|------|--------|----------|------|
| $K_L e3$ | 0.2163(6) | 0.26 | 0.09 | 0.20 | 0.11 | 0.06 |
| $K_L \mu 3$ | 0.2166(6) | 0.29 | 0.15 | 0.18 | 0.11 | 0.08 |
| $K_S e3$ | 0.2155(13) | 0.61 | 0.60 | 0.03 | 0.11 | 0.06 |
| $K^\pm e3$ | 0.2160(11) | 0.52 | 0.31 | 0.09 | 0.40 | 0.06 |
| $K^\pm \mu 3$ | 0.2158(14) | 0.63 | 0.47 | 0.08 | 0.39 | 0.08 |
| | 0.2163(5) | | | | | |

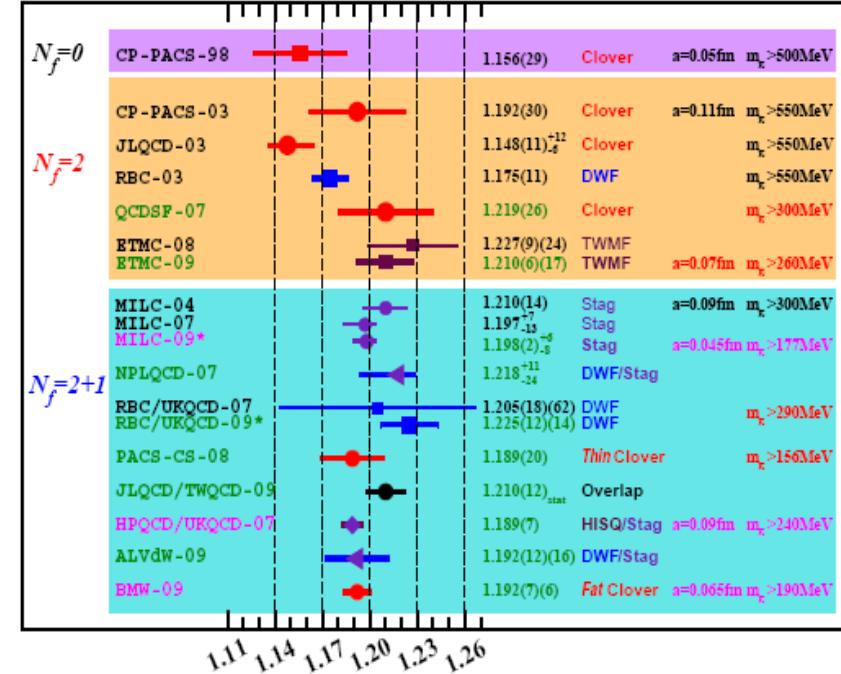
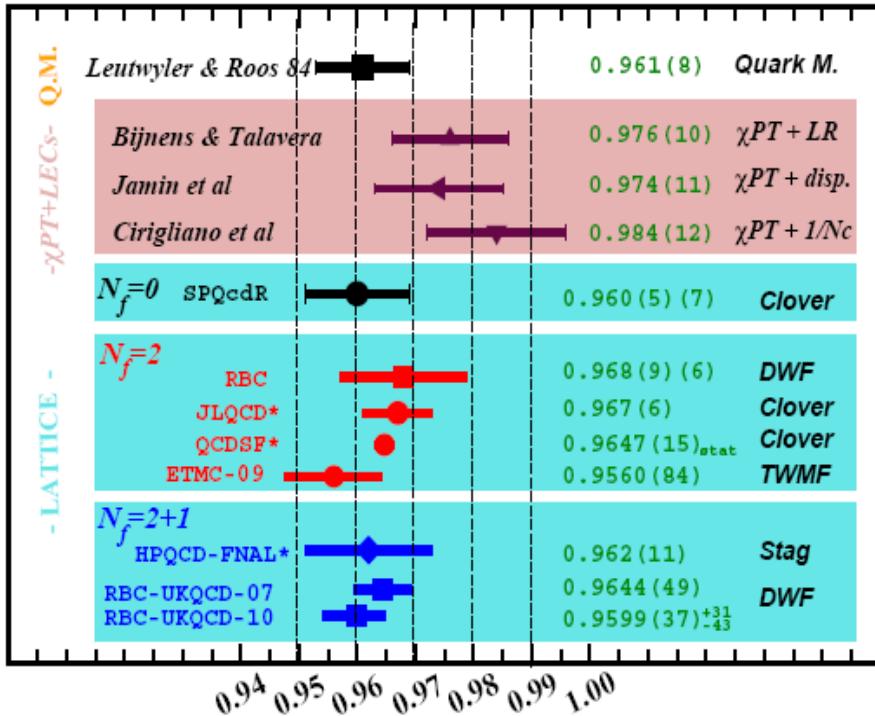
Average: $|V_{us}| f_+(0) = 0.2163(5)$ $\chi^2/\text{ndf} = 0.77/4$ (94%)

Evaluations of $f_+(0)$ and f_K/f_π

Lattice continuously improving

$$f_+^{K^0\pi^+}(0)$$

$$f_K/f_\pi$$



Use:

$f_+(0) = 0.964(5)$ RBC-UKQCD [FLAG 0.956(8)] [F. Mescia]

$f_K/f_\pi = 1.193(6)$ Kaon WG and FLAG average

$\tau \rightarrow P\nu, K_{\ell 2}$ decays

[Marciano]

Small uncertainties in f_K/f_π from lattice → determine V_{us}/V_{ud}
 Reduced uncertainty from e.m. Structure Dependence corrections

$$\frac{\Gamma(K \rightarrow \mu\nu)}{\Gamma(\pi \rightarrow \mu\nu)} = \frac{|V_{us}|^2}{|V_{ud}|^2} \times \frac{f_K^2}{f_\pi^2} \times \frac{m_K(1-m_\mu^2/m_K^2)^2}{m_\pi(1-m_\mu^2/m_\pi^2)^2} \times 0.9930(35)$$

$$|V_{us}|/|V_{ud}| f_K/f_\pi = 0.2760(6)$$

$$\frac{\Gamma(\tau \rightarrow K\nu)}{\Gamma(\tau \rightarrow \pi\nu)} = \frac{|V_{us}|^2}{|V_{ud}|^2} \times \frac{f_K^2}{f_\pi^2} \times \frac{(1-m_K^2/m_\tau^2)^2}{(1-m_\pi^2/m_\tau^2)^2} \times 1.0003(44)$$

$$|V_{us}|/|V_{ud}| f_K/f_\pi = 0.273(2)$$

$f_K/f_\pi = 1.193(6)$ [Kaon WG
FLAG WG]

$$V_{us}/V_{ud} = 0.2312(13) \quad [\text{Kaon WG}]$$

Inclusive V_{us} determination (more on D. Boito talk)

V_{us} from inclusive $\tau \rightarrow \nu X_{us}$ involves PQCD

$$|V_{us}|^2 = \frac{R_{\tau,S}^{00}}{R_{\tau,V+A}^{00} - \delta R_{\tau,\text{th}}^{00}}$$

Gámiz-Jamin-Pich-Prades-Schwab

$$V_{us} = 0.2159 (30_{\text{exp}})(5_{\text{th}})$$

Different theoretical analysis give somewhat larger errors : $\sim 1\%$

[Maltman et al]

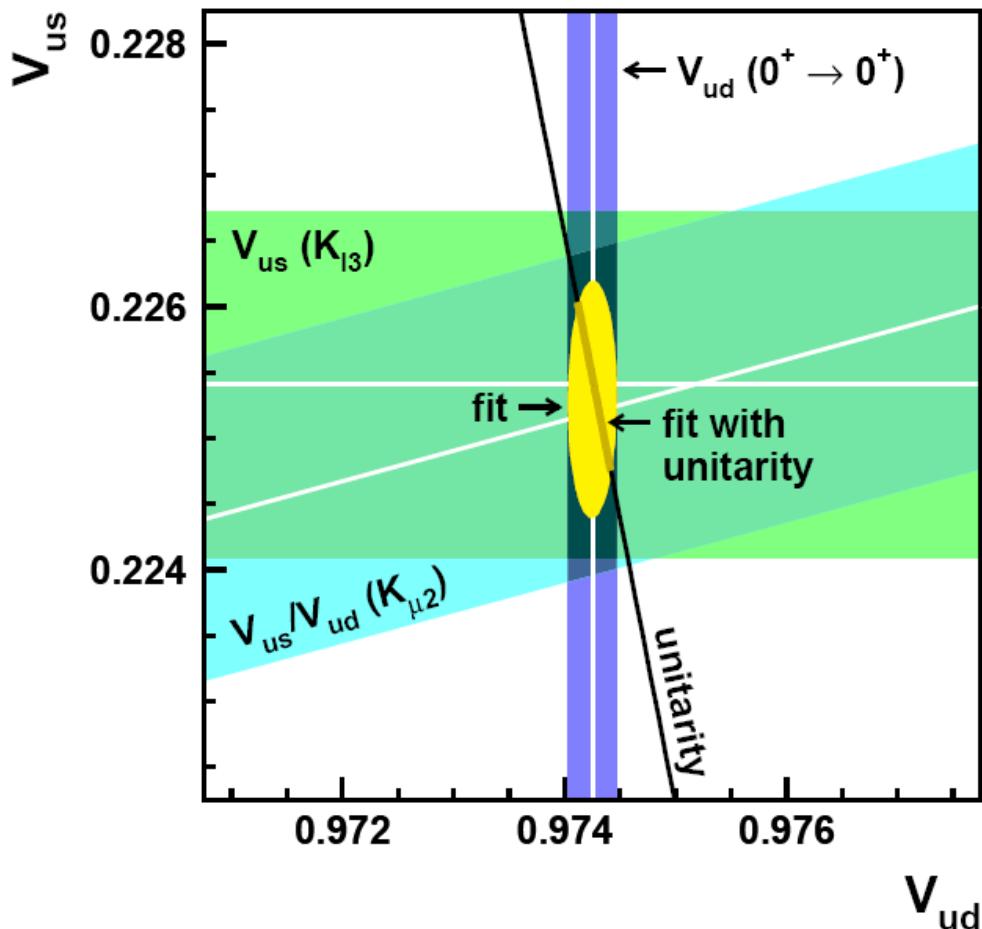
[HFAG]

| Base modes from τ^- decay | No B-Factory Data | With B-Factory Data |
|---|---------------------|---------------------|
| $K^-\nu_\tau$ | 0.686 ± 0.022 | 0.697 ± 0.010 |
| $K^-\pi^0\nu_\tau$ | 0.453 ± 0.027 | 0.431 ± 0.015 |
| $K^-2\pi^0\nu_\tau$ (ex. K^0) | 0.057 ± 0.023 | 0.060 ± 0.022 |
| $K^-3\pi^0\nu_\tau$ (ex. K^0, η) | 0.036 ± 0.022 | 0.039 ± 0.022 |
| $\overline{K}^0\nu_\tau$ | 0.888 ± 0.037 | 0.831 ± 0.018 |
| $\overline{K}^0\pi^-\nu_\tau$ | 0.358 ± 0.035 | 0.350 ± 0.015 |
| $\overline{K}^0\pi^-2\pi^0\nu_\tau$ | 0.027 ± 0.023 | 0.035 ± 0.023 |
| $\overline{K}^0h^-h^+\nu_\tau$ | 0.023 ± 0.020 | 0.028 ± 0.020 |
| $K^-\pi^-\pi^+\nu_\tau$ (ex. K^0, ω) | 0.334 ± 0.023 | 0.293 ± 0.007 |
| $K^-\pi^-\pi^+\pi^0\nu_\tau$ (ex. K^0, ω, η) | 0.039 ± 0.014 | 0.041 ± 0.014 |
| $K^-\phi\nu_\tau (\phi \rightarrow KK)$ | | 0.004 ± 0.001 |
| $K^-\eta\nu_\tau$ | 0.027 ± 0.006 | 0.016 ± 0.001 |
| $K^-\pi^0\eta\nu_\tau$ | 0.018 ± 0.009 | 0.005 ± 0.001 |
| $\overline{K}^0\pi^-\eta\nu_\tau$ | 0.022 ± 0.007 | 0.009 ± 0.001 |
| $K^-\omega\nu_\tau$ | 0.041 ± 0.009 | 0.041 ± 0.009 |
| Sum of strange modes | 3.0091 ± 0.0722 | 2.8796 ± 0.0501 |
| Sum of all modes | 100.00 | 100.00 |

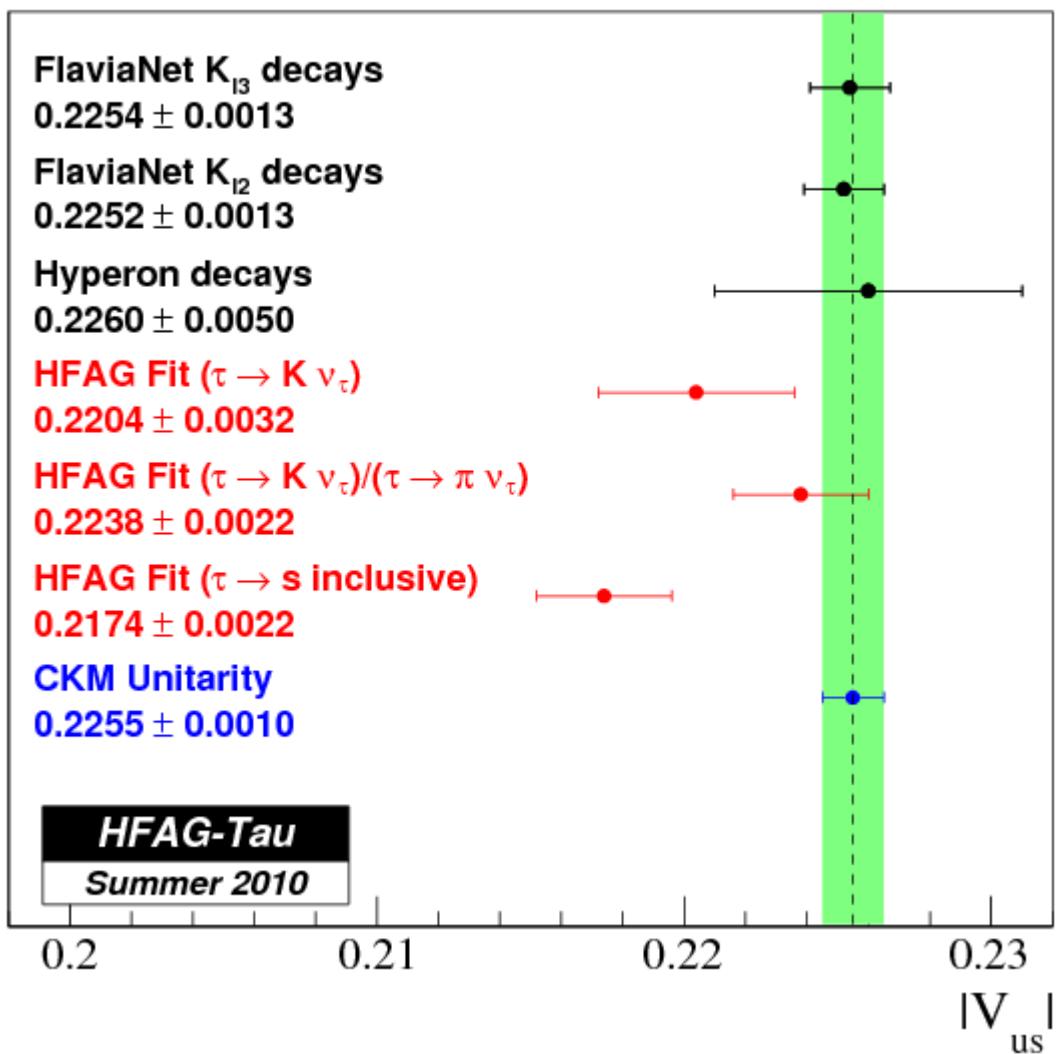
[HFAG]

$$V_{us} = 0.2169 (23)$$

CKM unitarity



τ and kaons summary



inclusive vs exclusive

inclusive

theory

Involve **PQCD** (~same method produced the best α_s determination)

Exp.

many sub-percent channels
New B-factory precision measurements not always
In agreement with old measurements from LEP

70% of r_s is made by

$\tau \rightarrow K\nu + \tau \rightarrow K\pi\nu$

Predictions from K might help

exclusive

theory

non **perturbative QCD** regime
Lattice + low energy theorems current dominating uncertainty

Exp.

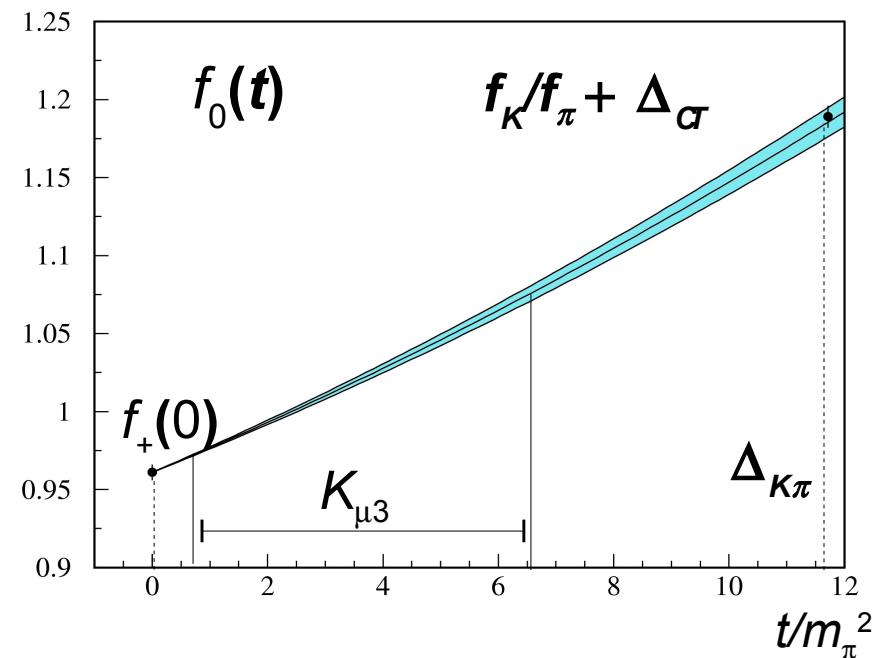
Very precise measurements form K, precise $B(\tau \rightarrow K\nu)$ from b-factory (2.6% below prediction from K) more complicated picture for $\tau \rightarrow K\pi\nu$



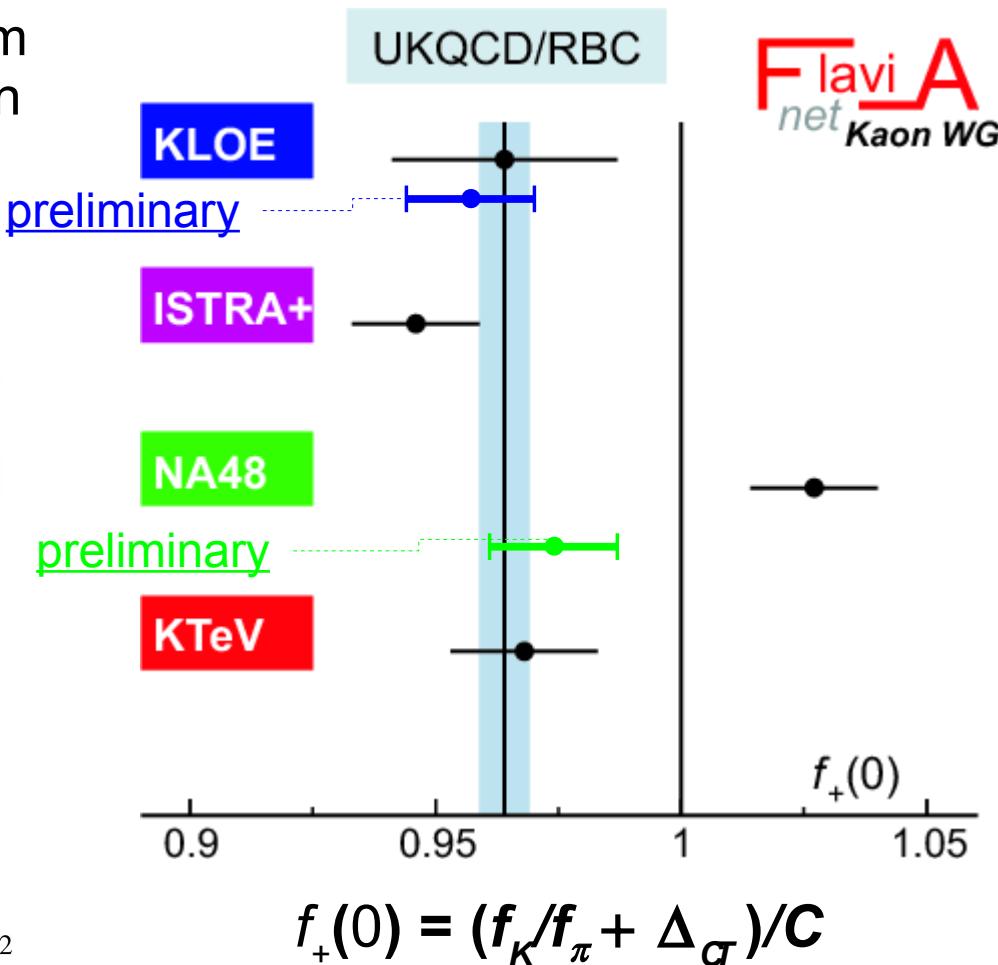
Must have a consistent picture between τ and K (no new physics in the game)

Callan-Treiman relation

Check/improve lattice results from measurement of scalar ff slopes in $K\mu 3$ and use of dispersive parametrization
 [Stern et al] [Pich et al]



$$f_K/f_\pi = 1.189(7) \quad \text{from HPQCD-UKQCD}$$



FlaviA
net Kaon WG

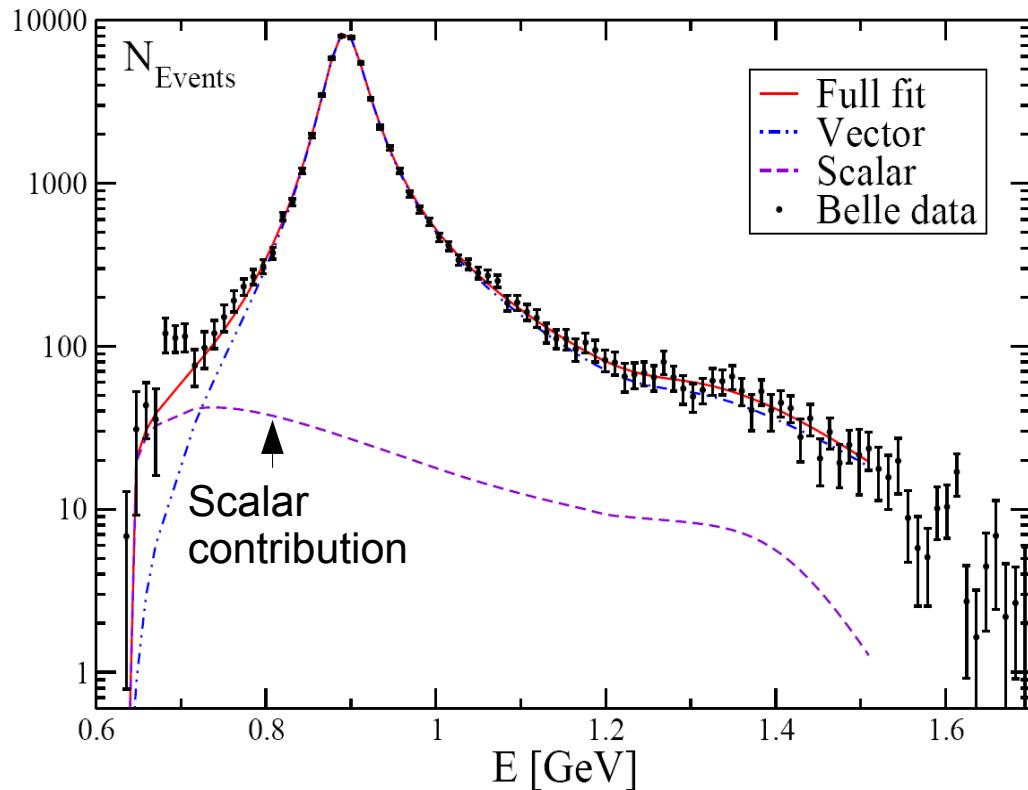
Relevance of $\tau \rightarrow K\pi\nu$ (more on D. Boito talk)

Present knowledge of the scalar ff parameter can be already improved by a factor ~ 2 with present B-factory data

Prediction for $B(\tau \rightarrow K\pi\nu)$ from KI3 and ff parameters important for R_s
Current accuracy still $\sim 2\%$
Need $100 \times$ data to match current exp. error on K BRs

[E. Passemar]

Jamin-Pich-Portolés 08, D. Boito, E. Passemar V. Bernard11 fit to **BELLE** data



Polarization and angular analysis helps a lot

Conclusion

1st raw unitarity status

$$V_{ud} = 0.97425(22)$$
$$V_{us} = 0.2253(9)$$

$$G_{\text{OM}} = 1.16633(35) \times 10^{-5} \text{ GeV}^2$$

0.03% accuracy on
 G_F from quarks

$$\sin\theta_c = \lambda = 0.2254(6)$$

Cabibbo angle at 0.3 % accuracy!

Present accuracy allows to test NP at ~ 1 TeV

V_{us} determination can be improved combining K and τ :

PQCD vs LATTICE QCD

ff knowledge can be improved with τ

Precise K BR measurements to predict 70% of Rs

Cabibbo angle must be measured at the Cabibbo Lab