

THE NON SM HIGGS

Michele Redi



La Thuile, 4 March

WHAT WE KNOW

SM is a gauge theory based on $SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$

$$\mathcal{L}_{Kinetic} = -\frac{1}{4}B_{\mu\nu}B^{\mu\nu} - \frac{1}{4}W_{\mu\nu}^a W^{a\mu\nu} - \frac{1}{4}W_{\mu\nu}^b W^{b\mu\nu} + i \sum_{j=1}^3 \left(\bar{\Psi}_L^j \not{D}\Psi_L^j + \bar{\Psi}_R^j \not{D}\Psi_R^j \right)$$

$$\Psi_{L,R} = (3, 2)_{\frac{1}{6}} \oplus (3, 1)_{\frac{2}{3}} \oplus (3, 1)_{-\frac{1}{3}} \oplus (1, 2)_{-\frac{1}{2}} \oplus (1, 1)_{-1} \quad (3 \text{ couplings})$$

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Unbroken gauge symmetry forbids mass terms:
vacuum must respect a smaller symmetry

$$SU(3)_c \otimes U(1)_{EM}$$

allowing to write mass terms,

$$\mathcal{L}_{mass} = \sum_{i,j=1}^3 \left[\bar{u}_L^i M_{i,j}^u u_R + \bar{d}_L^i M_{i,j}^d d_R + \bar{e}_L^i M_{i,j}^e e_R \right] + h.c.$$

$$+ m_W^2 W^2 + \frac{1}{2} m_Z^2 Z^2$$

$O(20)$ parameters

Mass for gauge bosons implies new degrees of freedom

$$m_1 = 0$$



$$m_1 \neq 0$$



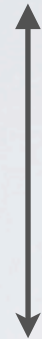
The extra degrees of freedom are Goldstone Bosons

$$SU(2)_L \otimes U(1)_Y \rightarrow U(1)_Q$$

which become longitudinal polarizations of W & Z

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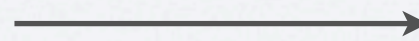
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Important hint:

$$\rho = \frac{m_W^2}{m_Z^2 \cos^2 \theta_W} \approx 1$$



Custodial Symmetry
 $SU(2)_c$

In the SM electro-weak symmetry is broken through a scalar doublet with $Y = 1/2$

$$V(H) = \lambda (|H|^2 - v^2)^2$$

$$H(x) = U(x) \begin{pmatrix} 0 \\ v + h(x) \end{pmatrix}, \quad v = 174 \text{ GeV}$$

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VEV breaks symmetry. The Goldstone Bosons in $U(x)$ are eaten giving mass to W & Z. Higgs sector respects custodial symmetry

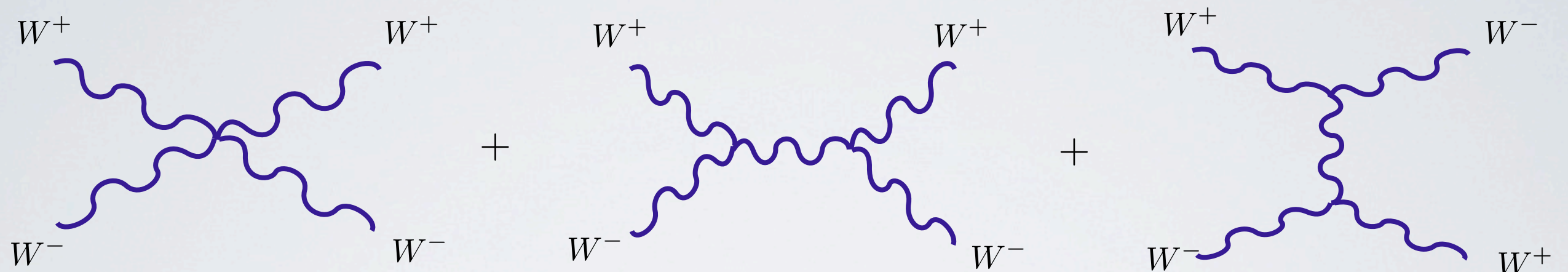
$$\frac{SU(2)_L \otimes SU(2)_R}{SU(2)_{L+R}} \longrightarrow \rho \approx 1$$

If SM is correct only unknown is the quartic/mass

$$m_h = \sqrt{\lambda} v$$

STRONG DYNAMICS

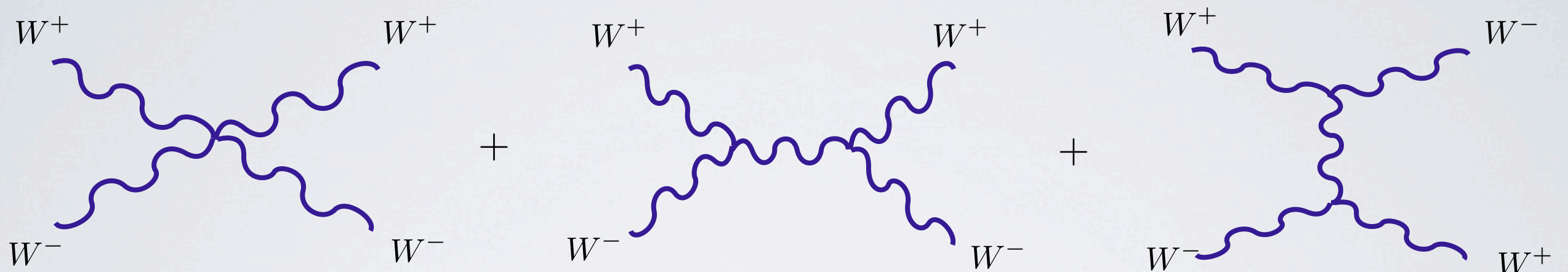
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Interactions become strongly coupled around TeV.
Perturbative unitarity is violated at

$$\Lambda \approx 1.7 \text{ TeV}$$

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$$\langle \bar{\Psi}_L^i \Psi_R^j \rangle = \Lambda_{QCD}^3 \delta_{ij} \longrightarrow \frac{SU(2)_L \otimes SU(2)_R}{SU(2)_{L+R}} \longrightarrow \mathcal{L} = f_\pi^2 \text{Tr} [\partial_\mu U \partial^\mu U^\dagger]$$

The 3 Goldstone bosons are the pions.

scattering is unitary because they are composite objects.

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The 3 pions would become longitudinal polarization of
 W & Z without other effects

$$m_W = g f_\pi \sim 50 \text{ Mev}$$

$$\rho \approx 1$$

Techni-color is a rescaled version of QCD,

$$f = v$$

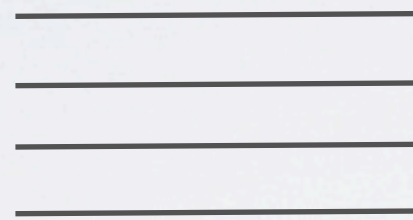
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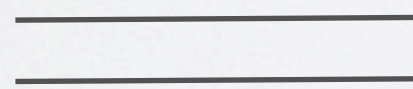
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Longitudinal polarizations of W & Z are composite states.

No Higgs scalar but techni-resonances (spin 0, 1/2, 1 etc.).



$$m_\rho \leq 1 \text{ TeV}$$



$$m_W = 80 \text{ GeV}$$

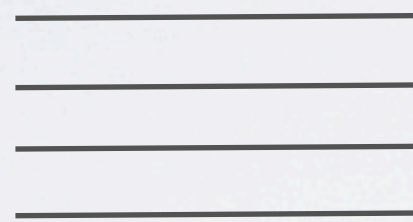
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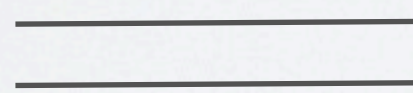
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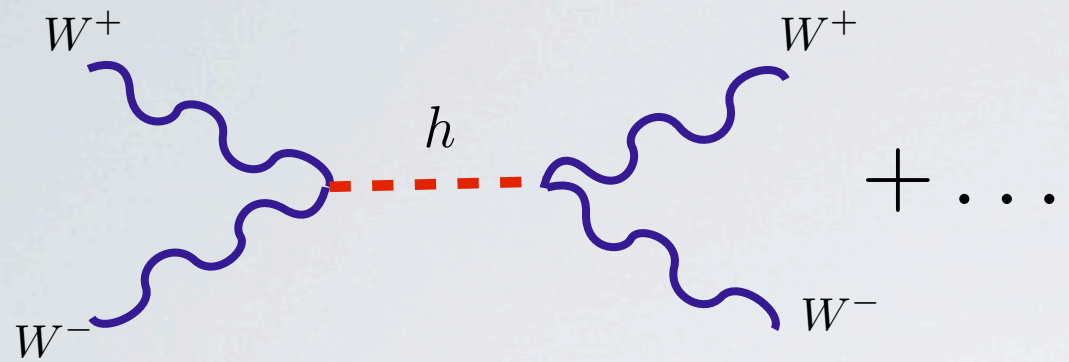


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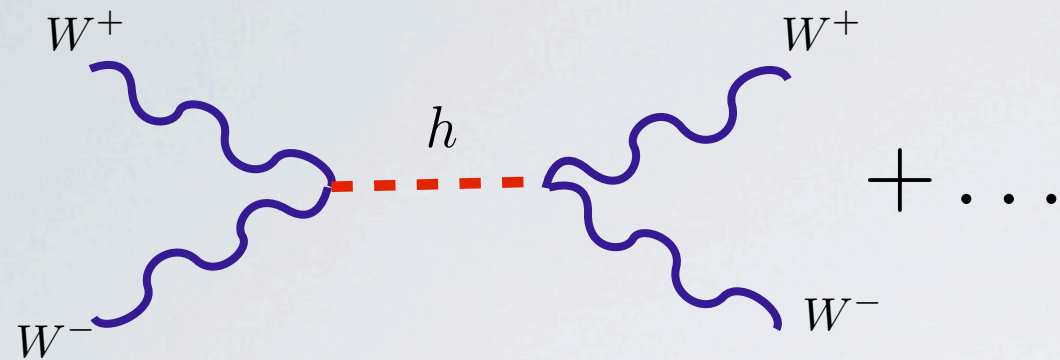
Sadly phenomenologically ruled out:
precision electro-weak measurements and flavor

In the SM:



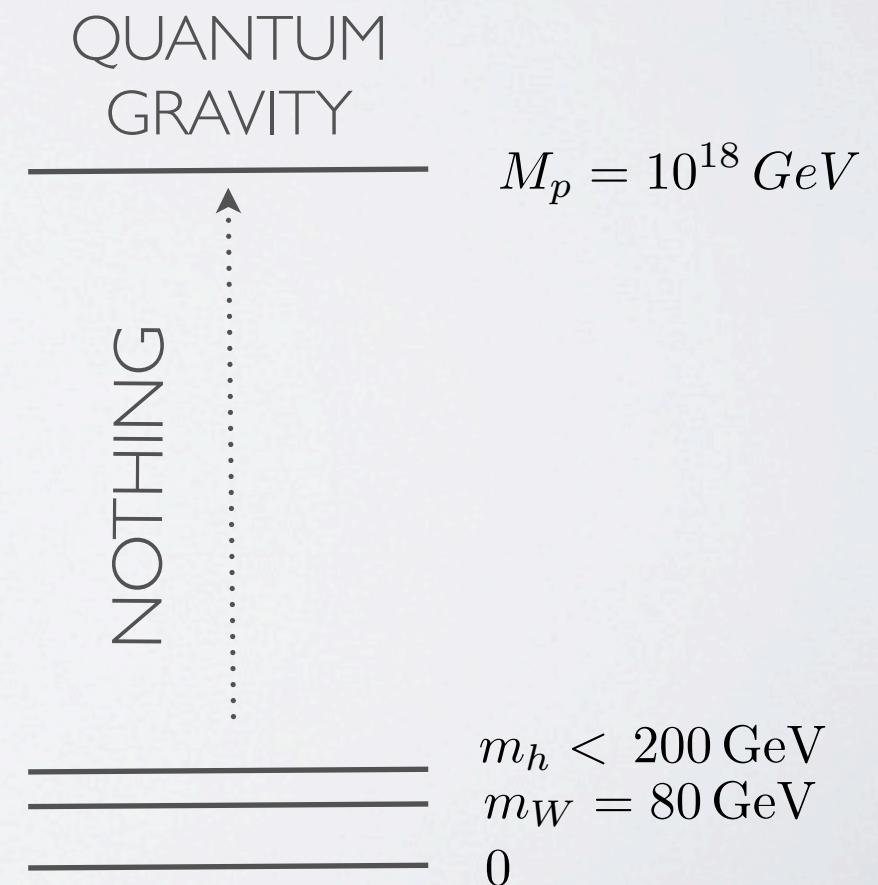
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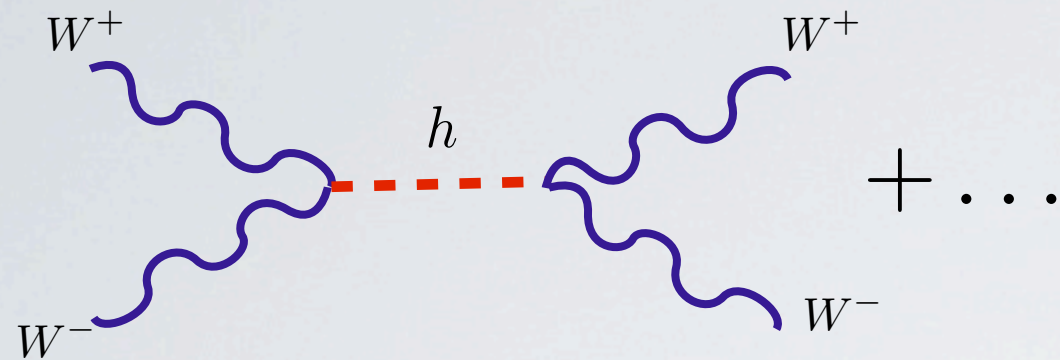


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Amplitude does not grow so SM can be valid up to the Planck scale.



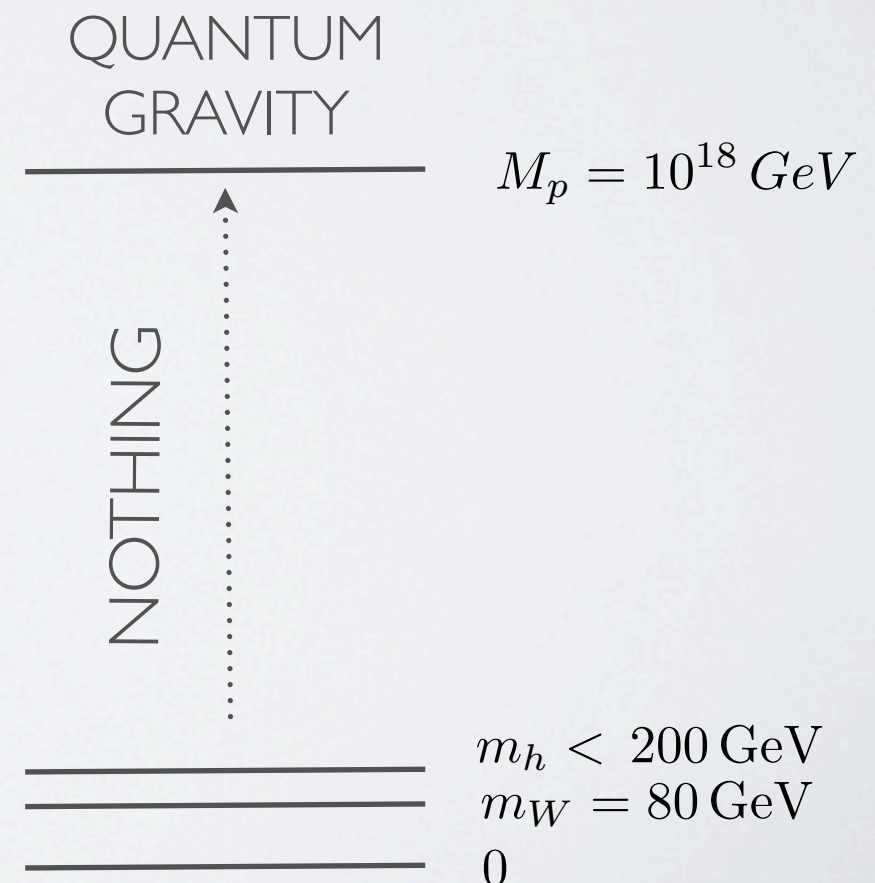
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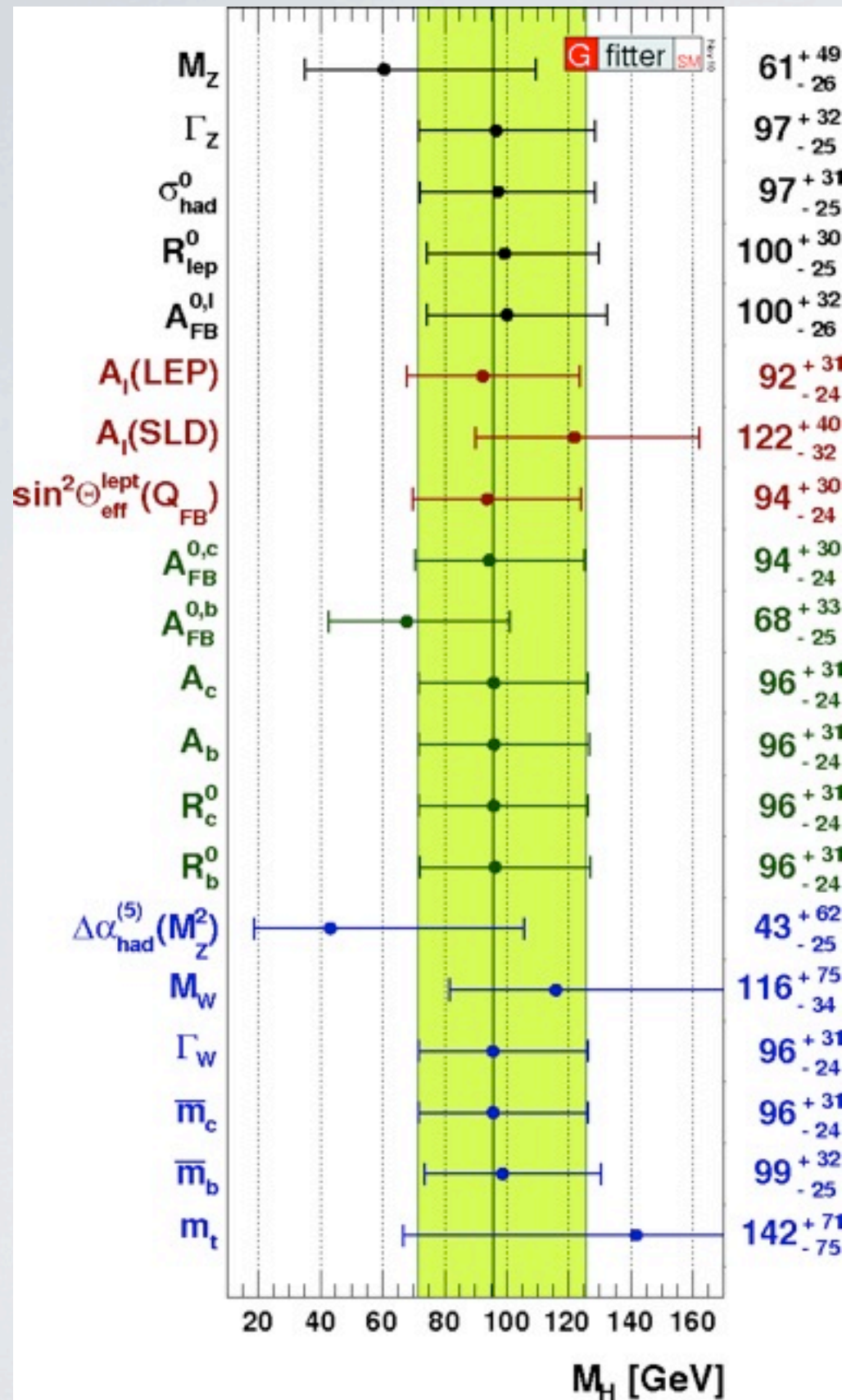


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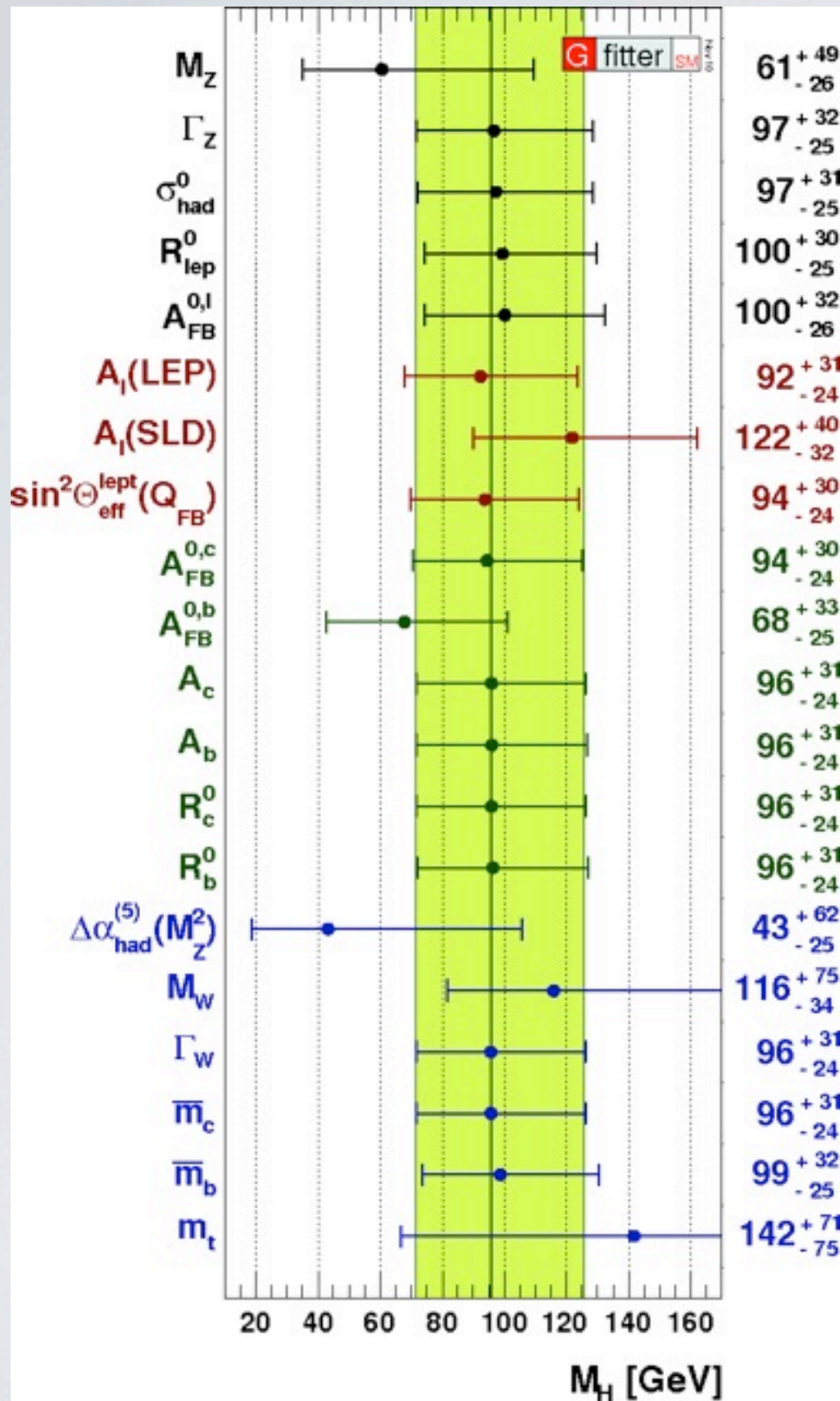
- Hierarchy problem
- Dark Matter
- Origin of Yukawas, CP
- Explains nothing



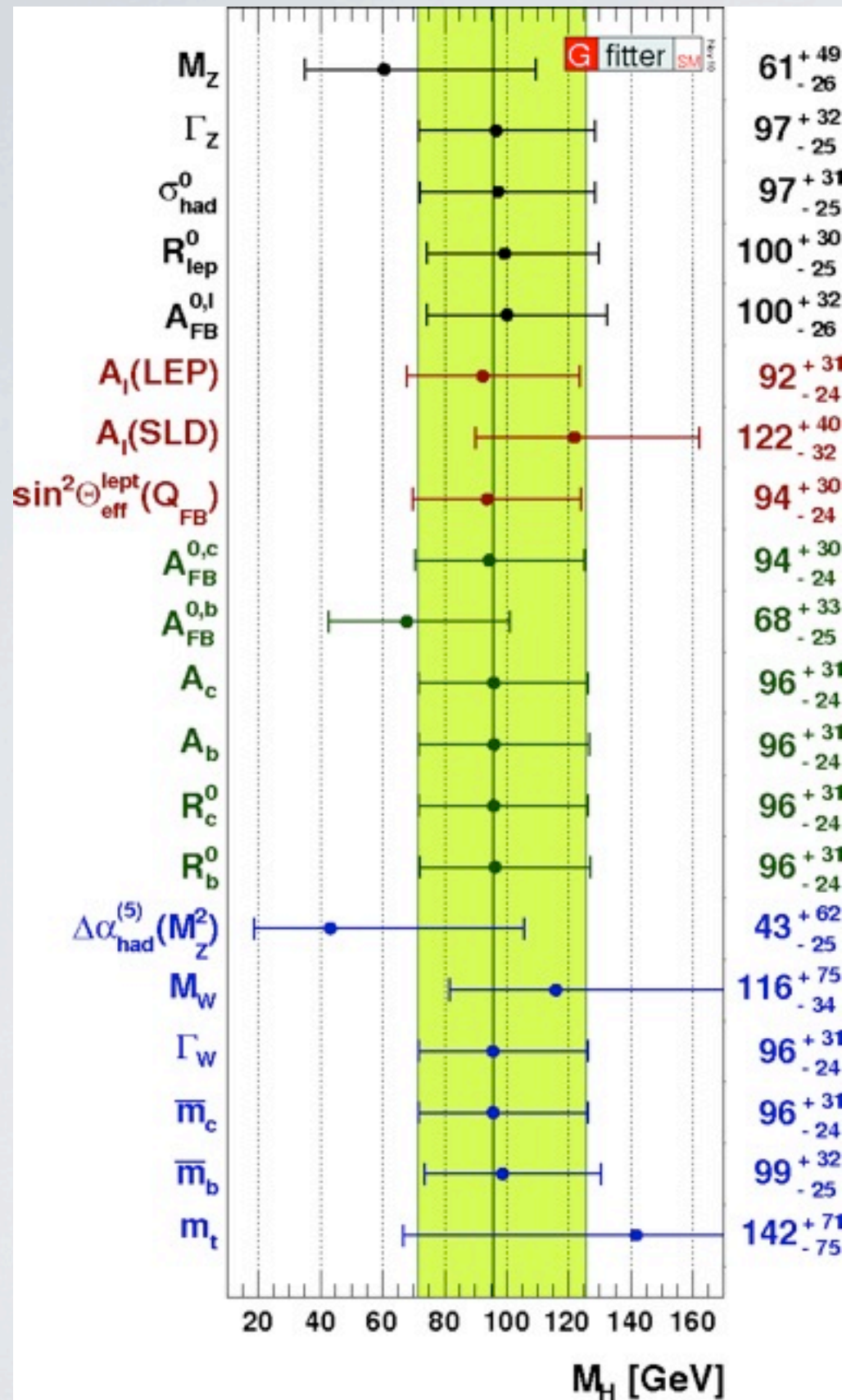


SM nice but not perfect

SM probability 15 %



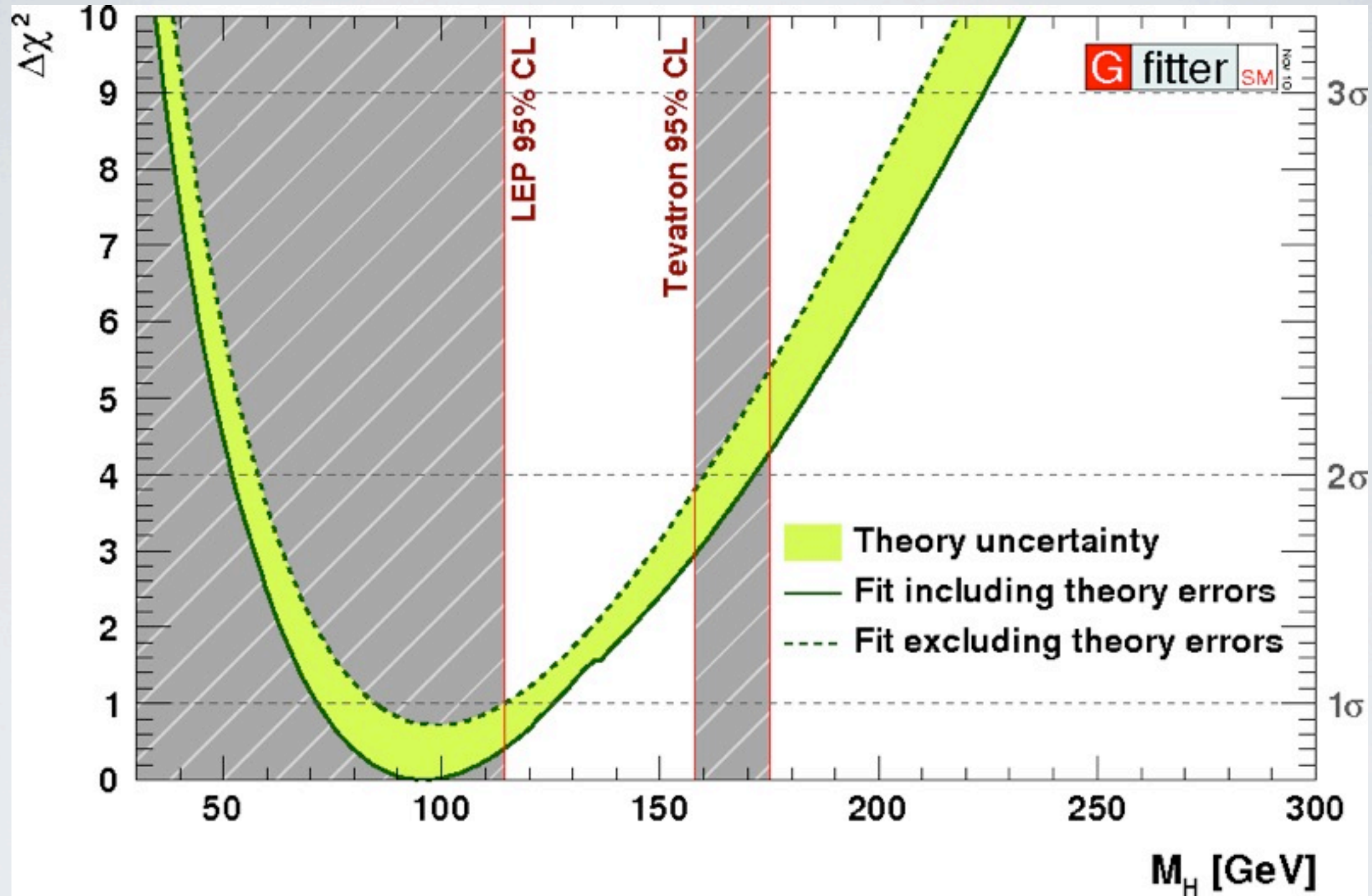
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SM probability 15 %

Just observables related to the Higgs < 2%

SM nice but not perfect



Higgs must be light:

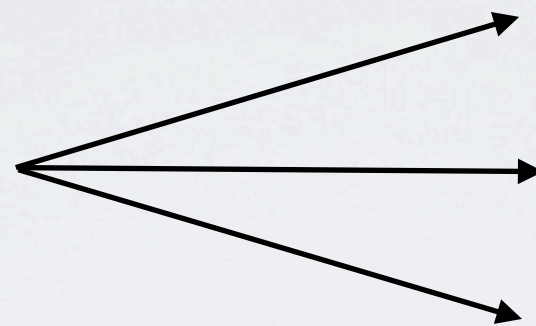
SM ruled out @ 99.7 C.L. if $M > 225$ GeV (worse in SUSY)

COMPOSITE HIGGS

Georgi, Kaplan '80s

Higgs doublet could be a light remnant of strong dynamics.

Strong sector:
resonances +
Higgs bound state



spin 1

spin 1/2

spin 0.... $2\frac{1}{2}$

Two parameters:

m_ρ

g_ρ

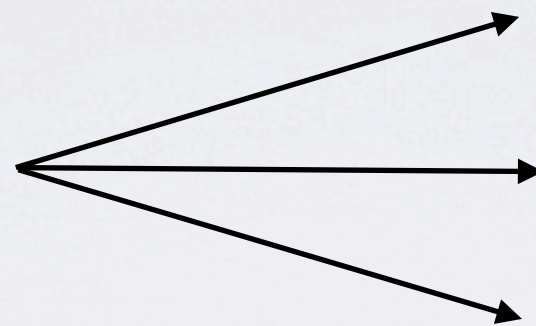
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- Relieves hierarchy problem

$$\delta m_h^2 \sim \frac{3 \lambda_t^2}{4\pi^2} m_\rho^2$$

Particularly natural if Higgs is a Goldstone Boson:
Massless at leading order.

Ex: $\frac{SO(5)}{SU(2)_L \otimes SU(2)_R} \longrightarrow GB = (2, 2)$

Agashe , Contino,
Pomarol, '04

Or: $\frac{SO(6)}{SO(5)} \quad \frac{SO(6)}{SO(4) \otimes U(1)} \quad + \dots$

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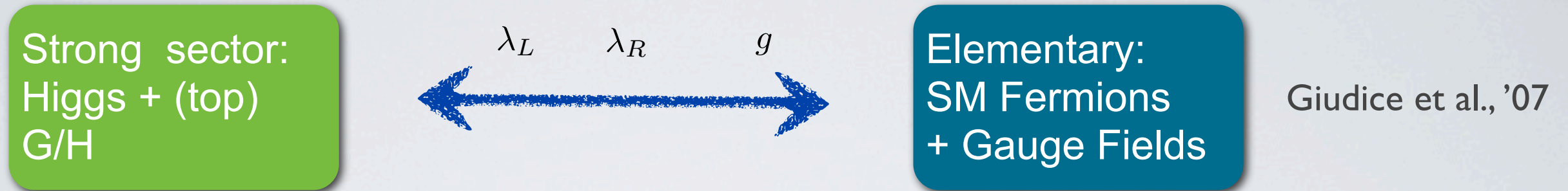
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Low energy lagrangian

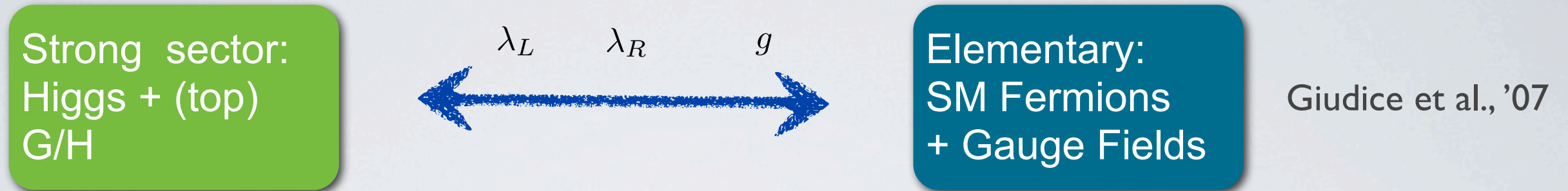
$$\mathcal{L} = f^2 D_\mu \Sigma^i D^\mu \Sigma^i + \dots \xrightarrow{SU(2)_L \otimes SU(2)_R} \rho \approx 1$$

$$m_\rho = g_\rho f$$

Goldstone symmetry is broken explicitly by SM couplings



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Two sectors talk through gauging of $SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$

$$\mathcal{L}_{gauge} = g A_\mu J^\mu$$

And through mixing to fermionic states

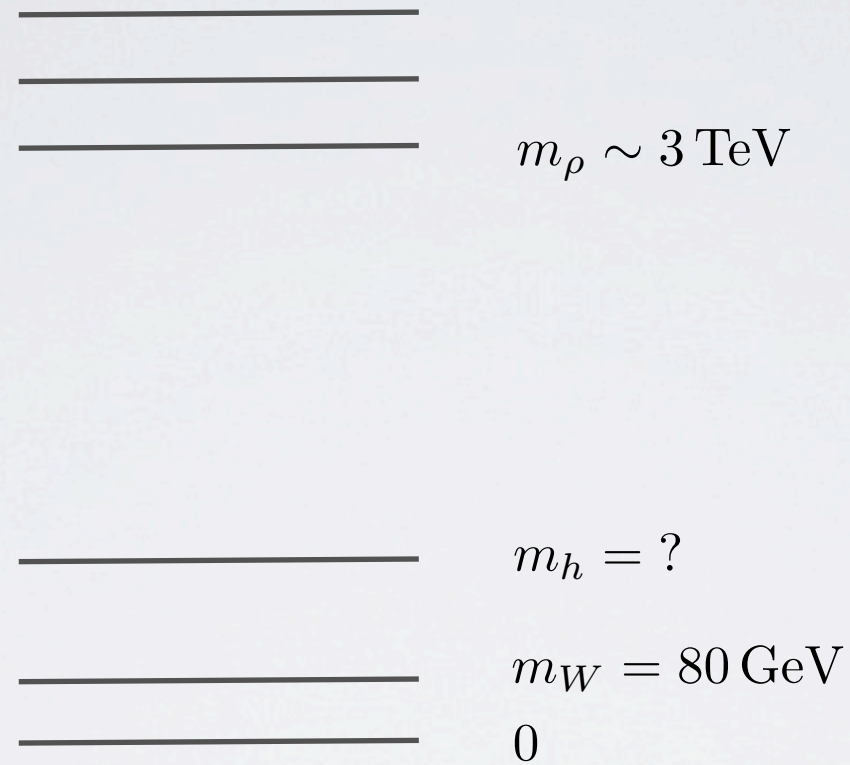
$$\mathcal{L}_{mixing} = \lambda_L \bar{f}_L O_R + \lambda_R \bar{f}_R O_L \quad \longrightarrow \quad y_f \sim \frac{\lambda_L \lambda_R}{g_\rho}$$

Potential also generated

$$V(H) \sim N_c \frac{m_\rho^4}{g_\rho^2} \frac{\lambda_{t_L, t_R}^2}{16\pi^2} \hat{V} \left(\frac{H}{f} \right)$$

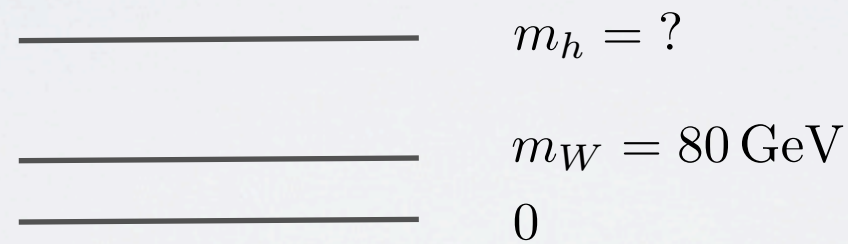
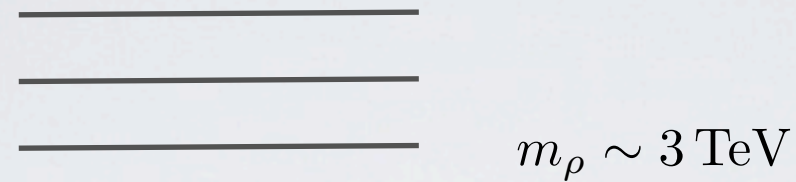
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Spectrum:

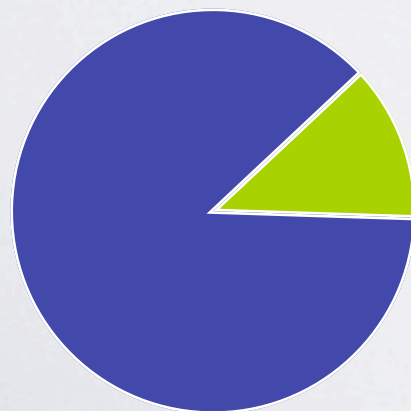


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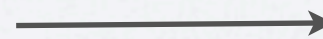
Spectrum:



Higgs is angle,



$$0 < h < 2\pi f$$



Tuning

$$\xi = \frac{v^2}{f^2}$$

Scenario approaches SM for $f \gg v$.

Precision measurements can be reproduced for $m_\rho > 2.5\text{TeV}$

Flavor problems greatly reduced,

$$C_4^K \bar{d}_R^\alpha s_L^\alpha \bar{d}_L^\beta s_R^\beta$$

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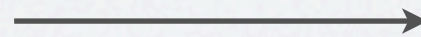
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Acceptable pheno at the price of mild tuning,

$$f = 500 \text{ GeV} \quad m_\rho = 3 \text{ TeV}$$

$$(g_\rho \sim 6 \sim \text{QCD})$$



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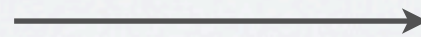
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Overall models are far from perfect but general picture is compelling and worth taking seriously.

Composite Higgs can be distinguished from SM Higgs:

- Heavy resonances

Mostly coupled to third generation quarks

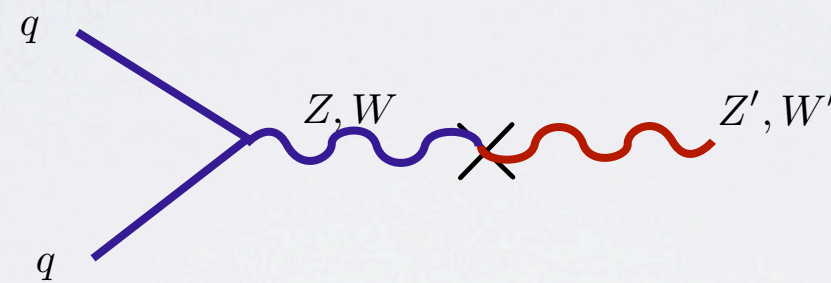
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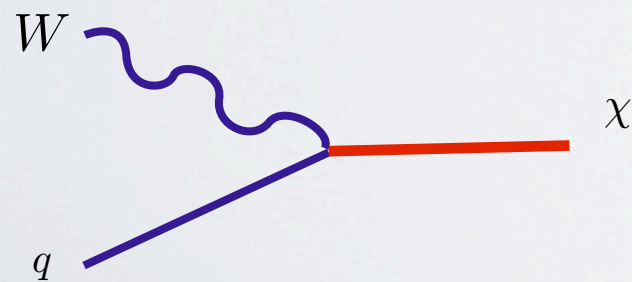
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Spin-1: electro-weak and gluon resonances

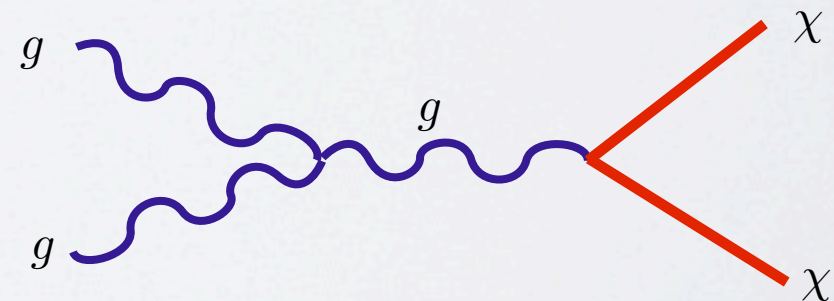
Ex:



Spin-1/2: Top partners could be lighter + exotic charges

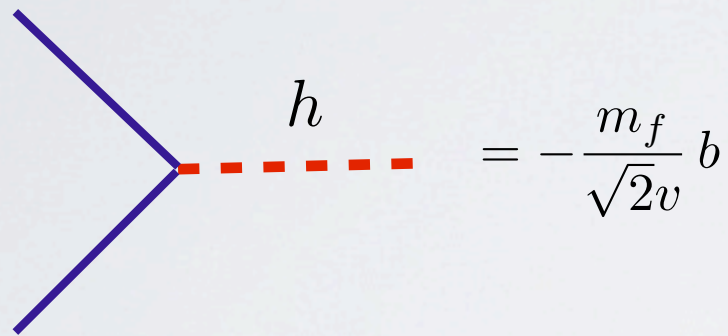
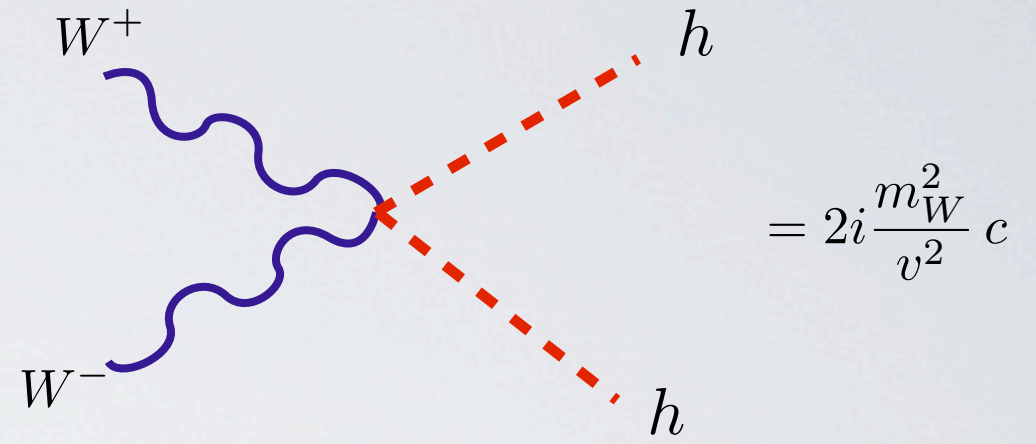
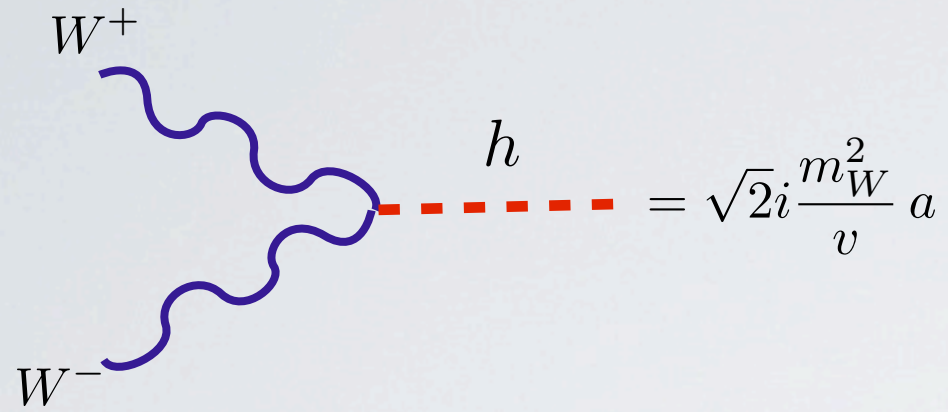


Single



Double

• Modified Couplings



SM : $a = b = c = 1$

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$$= \sqrt{2}i \frac{m_W^2}{v} a$$

$$= 2i \frac{m_W^2}{v^2} c$$

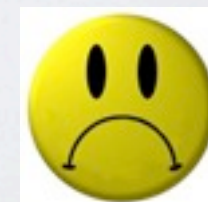
$$= -\frac{m_f}{\sqrt{2}v} b$$

SM : $a = b = c = 1$

WW scattering does not unitarize completely.
 Unfortunately very hard to see at LHC unless $v/f \sim 1$.

Contino et al. '10

Deviations from $a, b, c = 1$. Wait ILC...



Higgs quartic increases with strength of the strong dynamics. For example for $\lambda_{tL} \sim \lambda_{tR}$

$$\lambda \sim N_c \frac{g_\rho^3 \lambda_t}{16\pi^2} \longrightarrow m_h^2 \sim N_c \frac{g_\rho^3 \lambda_t}{16\pi^2} v^2$$

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Higgs can be heavy for a strong coupling similar to QCD.

$$100 \text{ GeV} < m_h < 400 \text{ GeV}$$

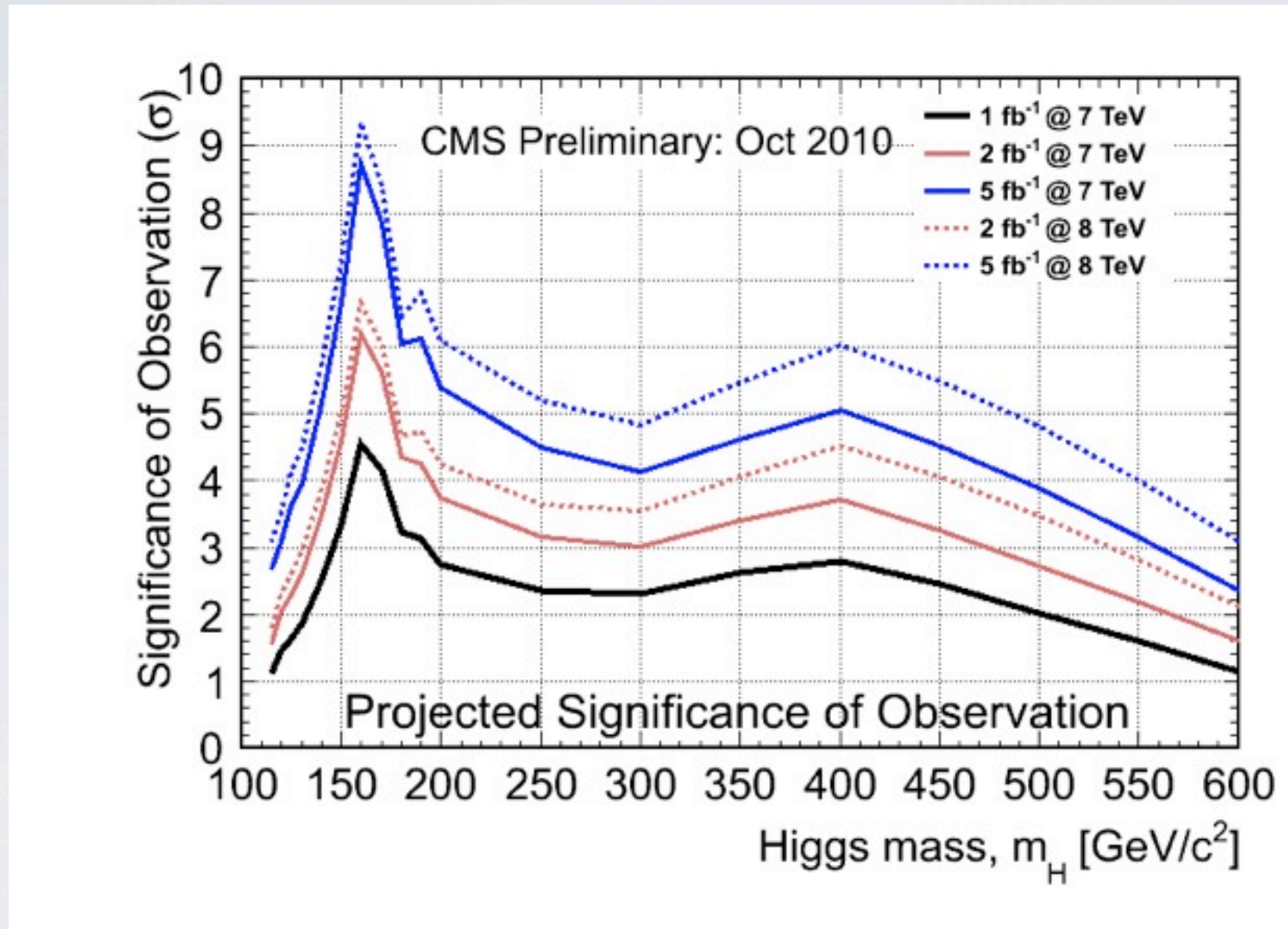
A large coupling is also hinted to minimize the tuning,

$$\xi \sim \frac{v^2}{f^2} \sim g_\rho^2 \frac{v^2}{m_\rho^2}$$

Barbieri et al., '07

Precision measurements could be ok with the contributions from extra-states (positive \mathbb{T}).

COULD WE GET LUCKY?



A heavy Higgs would rule out SM and hint to compositeness!

CONCLUSIONS

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- Finding soon something like the Higgs is likely.

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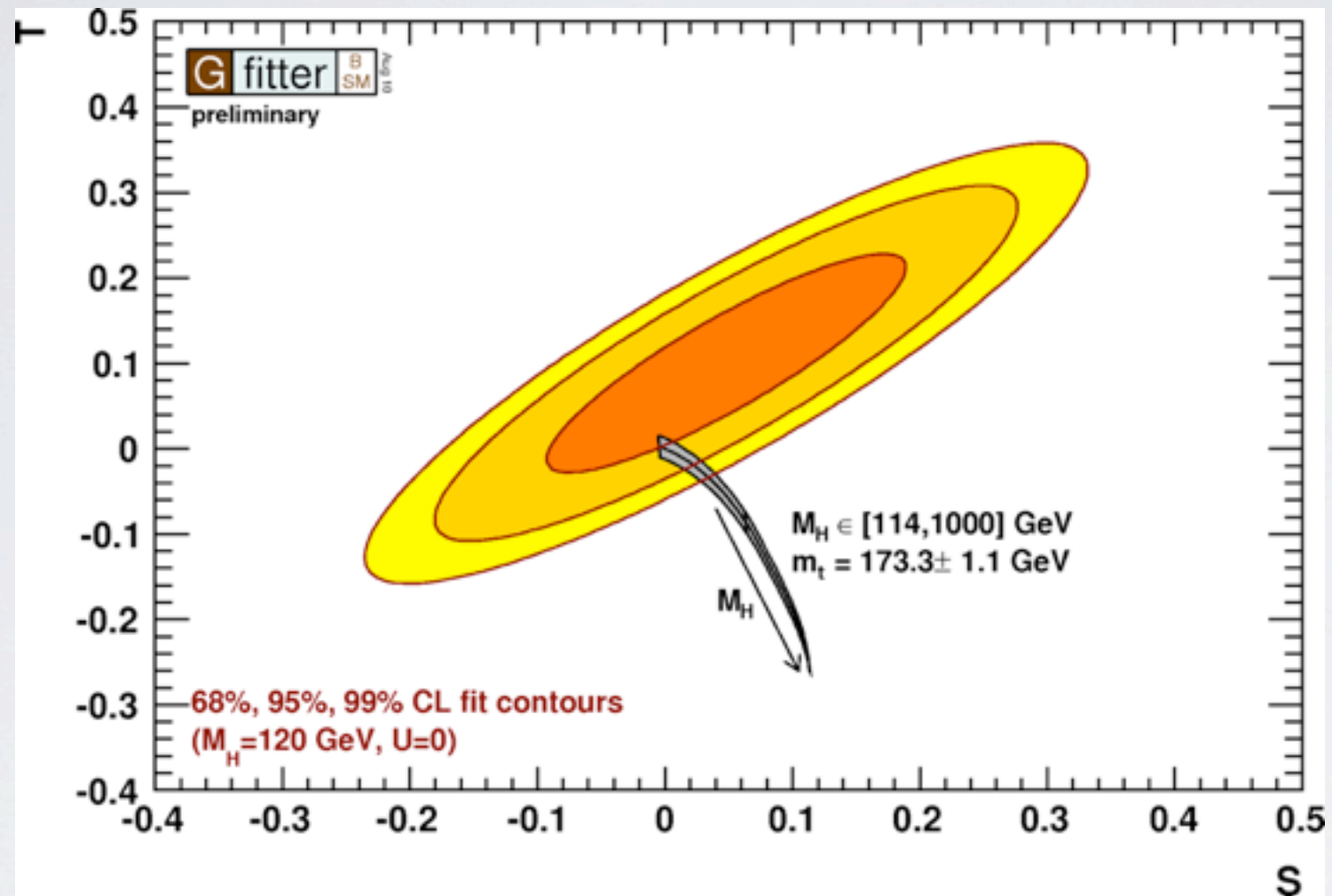
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- Distinguishing SM Higgs from composite Higgs scenarios will take time. Unless it is heavy.

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- Finding soon something like the Higgs is likely.
- Distinguishing SM Higgs from composite Higgs scenarios will take time. Unless it is heavy.
- 2011-12 will be VERY exciting!

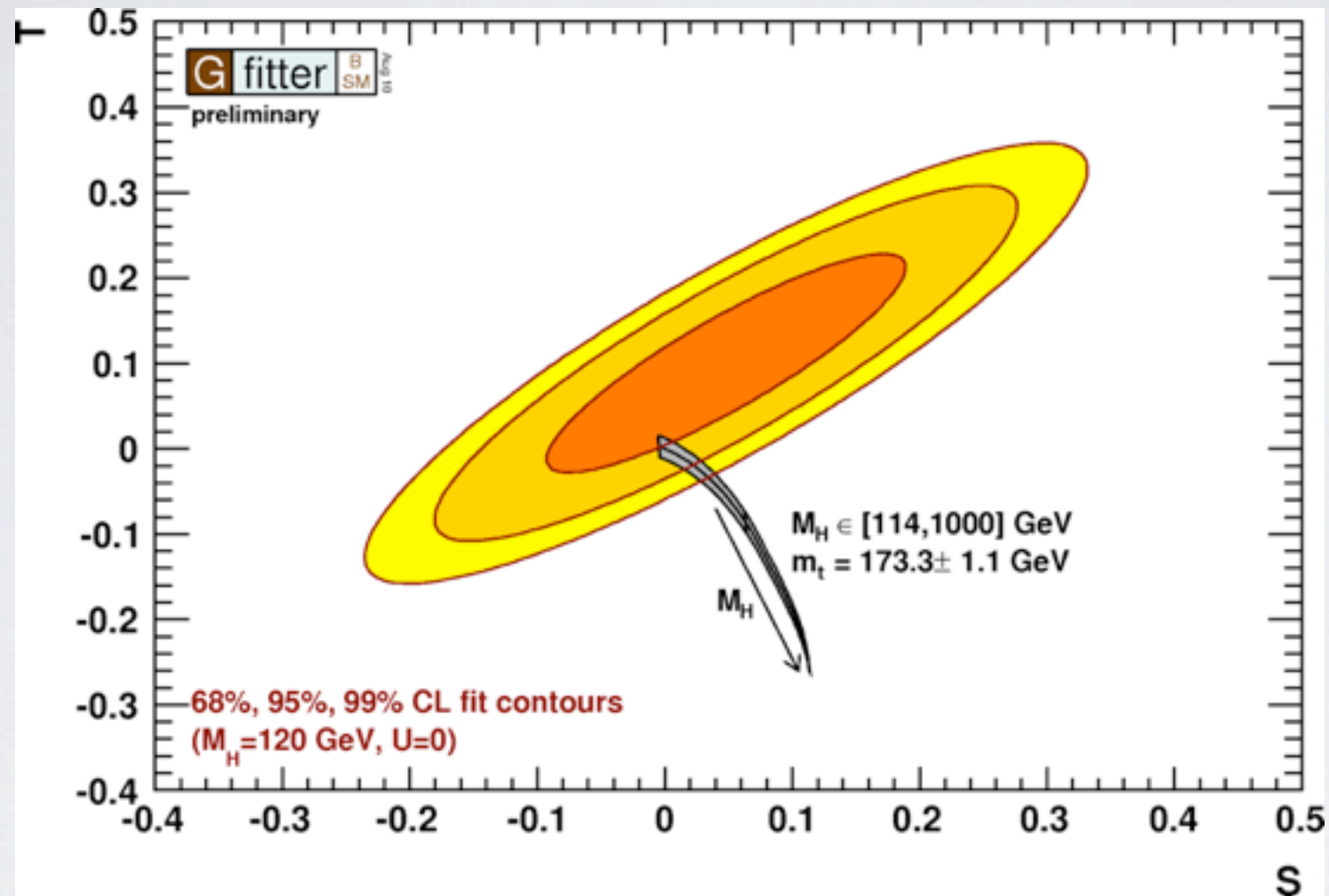
Technicolor issues:

- *Precision Measurements*



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- Flavor Bounds

Unacceptable Flavor Changing Neutral Currents generated unless miracles happen.