

Scalar Dark Matter: A revision of the inert doublet model

Laura Lopez Honorez

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based on:

A new viable region of the inert doublet model: JCAP 1101:002
The inert doublet model of dark matter revisited: JHEP 1009:046

in collaboration with C. Yaguna



Les Rencontres de Physique de la Vallée d'Aoste
La Thuile 2011 - Italy

Scalar WIMP as dark matter

Minimal DM spirit : SM + extra $SU(2)_L$ doublet see Cirelli *et al* '05-'09

- DM = **neutral** member of the extra doublet
Higgs : $H_1 \rightsquigarrow h$ and DM : $H_2 \rightsquigarrow H_0, A_0, H^\pm$
- stability \rightsquigarrow extra Z_2 symmetry
SM \rightarrow SM, $H_1 \rightarrow H_1$ and $H_2 \rightarrow -H_2$

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Possible Interactions

- (co)annihilation through *known* $SU(2)_L \times U(1)$ gauge bosons exchange
- quartic coupling λ_i to Higgs H_1

$$\lambda_3 |H_1|^2 |H_2|^2 + \lambda_4 |H_1^\dagger H_2|^2 + \frac{\lambda_5}{2} \left[(H_1^\dagger H_2)^2 + h.c. \right]$$

Scalar WIMP as dark matter

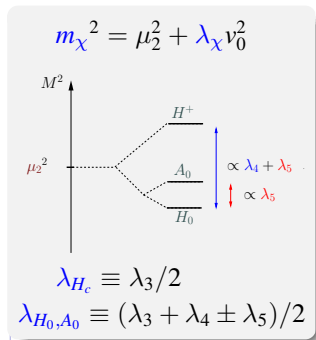
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$$\frac{1}{2} (\lambda_{H_0} H_0^2 + \lambda_{A_0} A_0^2 + 2\lambda_{H_c} H^+ H^-) (2v_0 h + h^2)$$



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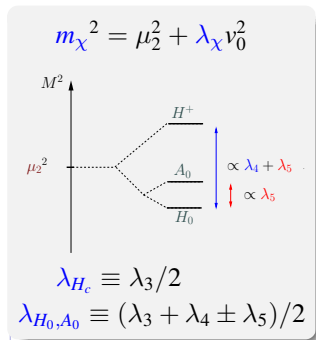
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Here we take H_0 as the LZ₂P and $\lambda_L \equiv \lambda_{H_0}$

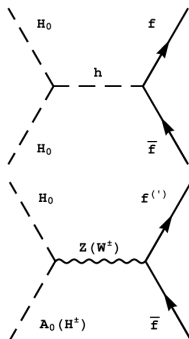


Inert doublet model parameter space

- $m_{H_0} \lesssim m_W$: GeV range ($\sim 10 - 80$ GeV) :

$$H_0 H_0 \rightarrow h^* \rightarrow \bar{f} f \text{ and } H_0 A_0 \rightarrow Z^* \rightarrow \bar{f} f$$

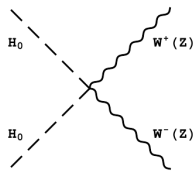
Barbieri PRD06, LLH JCAP06, Gustafsson PRL07, Cao PRD07, Andreas JCAP08,...



LARGE MASS GAP DUE TO EFFICIENT
 WW AND ZZ ANNIHILATION

- $m_{H_0} \gg m_W$: TeV range ($\gtrsim 530$ GeV) :
 $H_0 H_0 \rightarrow ZZ, WW, hh$ and coannihil into bosons

Cirelli NPB06, Hambye JHEP09



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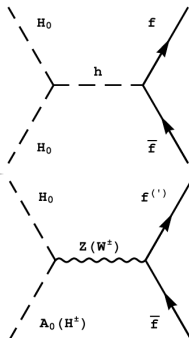
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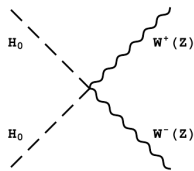
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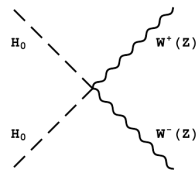
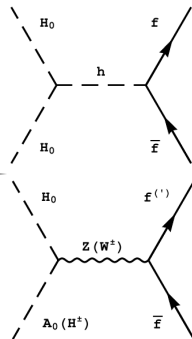
- Above W -threshold : cancellations

$$H_0 H_0 \rightarrow WW \text{ vs } H_0 H_0 \rightarrow h \rightarrow WW$$

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Extra contributions from 3 body annihilations

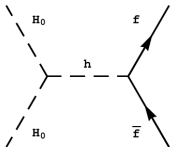
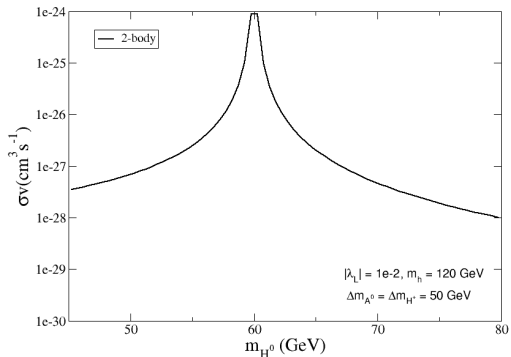
Analysis for fixed parameters

based on:

A new viable region of the inert doublet model: JHEP 1009:046

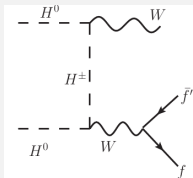
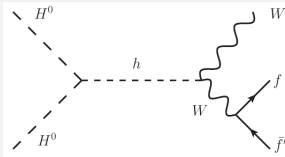
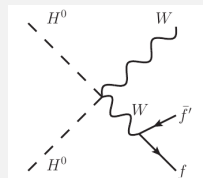
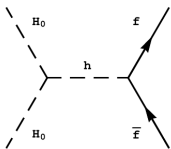
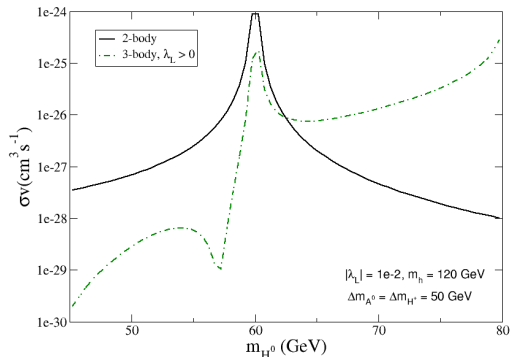
2-3 body annihilation cross section near m_W threshold

- σV_{2bdy} : higgs mediated,
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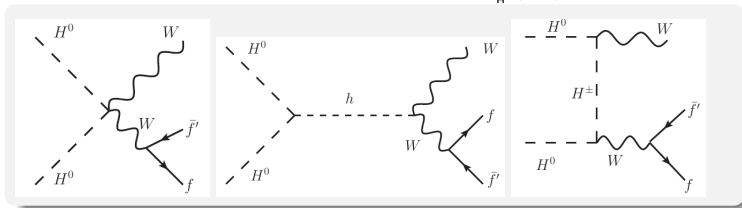
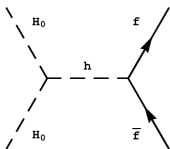
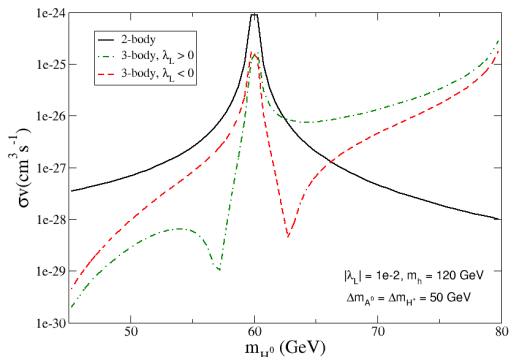
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- $\sigma v_{3bdy} = \sigma v(WW^*)$
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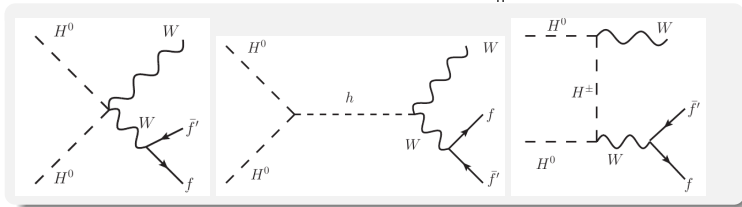
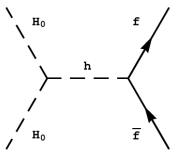
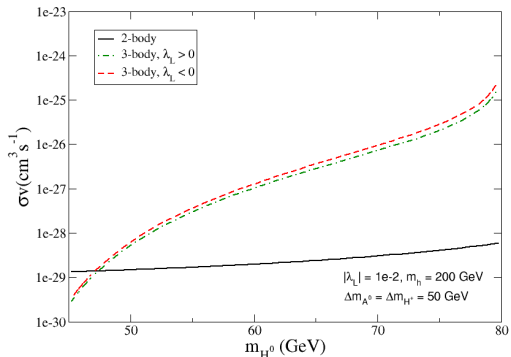
$\rightsquigarrow \sigma V_{2bdy}$ VS σV_{3bdy} depends on m_{H^0} , m_h , λ_L sign and amplitude



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Comparing 2-3 body relic density

- roughly $\Omega_{dm} \propto 1/\langle\sigma v\rangle$ with
 $\langle\sigma v\rangle = \langle\sigma v(2\text{-body})\rangle + \langle\sigma v(WW^*)\rangle$
- We expect
 $\Omega_{dm}(3\text{-body}) \lesssim \Omega_{dm}(2\text{-body})$

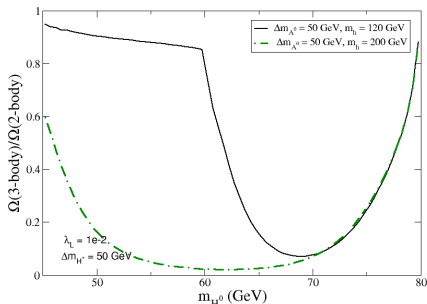
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$$\Omega_{dm}(3\text{-body}) \lesssim \Omega_{dm}(2\text{-body})$$

\rightsquigarrow confirmed numerically using
modified micrOMEGAs



\rightsquigarrow **3-body final states significantly affect predictions for Ω_{dm}**

Extra contributions from 3 body annihilations

Parameters for $\Omega_{H_0} = \Omega_{dm}^{WMAP}$

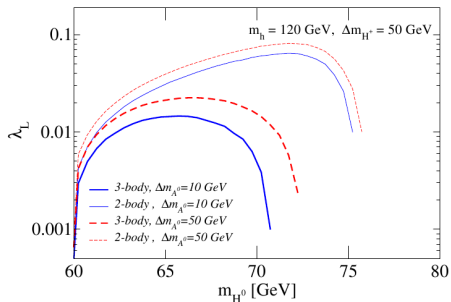
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Viable parameter space

Derive the $\lambda_L - m_{H^0}$ compatible with $\Omega_{dm}^{WMAP} h^2 = 0.11$

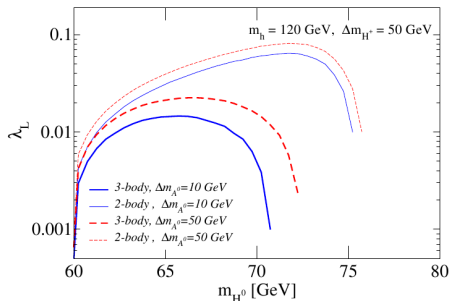
Going from 2bdy only to 2+3bdy with or without coannihilations :



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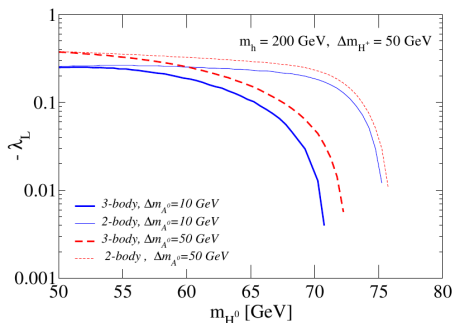


- correct $|\lambda_L|$ is reduced up to $\sim \mathcal{O}(10)$.
- W -threshold at lower m_{H^0}
 - 2bdy settled by the onset of W^+W^- annihilations
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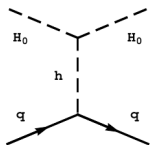
\rightsquigarrow rather generic feature of the Inert doublet model independently of m_h

\rightsquigarrow modify prospects for DM detection

Implications for Direct Detection

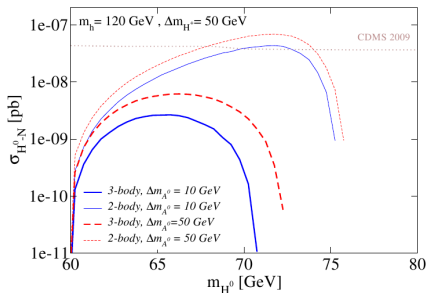
Direct detection through Elastic Scattering

Prospects along the viable parameter space :



$$\sigma_{H_0-N} \propto \left(\frac{\lambda_{H_0}}{M_{H_0} M_h^2} \right)^2$$

\rightsquigarrow predictions for σ_{H_0-N} reduced
up to $\sim \mathcal{O}(100)$

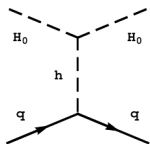


\rightsquigarrow better compatibility with present bounds

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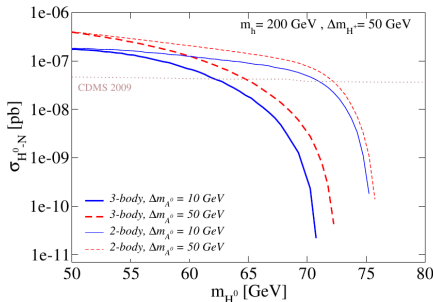
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Cancellations above W threshold

based on:

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Illustration of cancellation in 2 body processes

Above W threshold, for **fixed** parameters, one obtains $\Omega_{H_0} = \Omega_{dm}^{WMAP}$

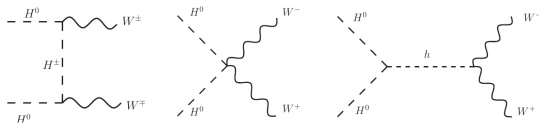


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 $\sigma v \propto g^4/m_{H^0}^2 \rightsquigarrow m_{H^0} > 534 \text{ GeV}$

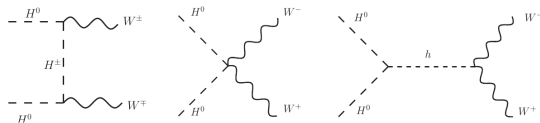


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- i.e.
- $\lambda_L < 0$ for $m_{H^0} > M_h/2$
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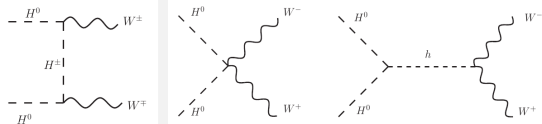


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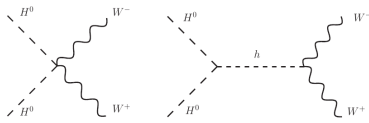
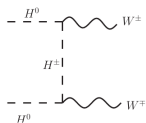
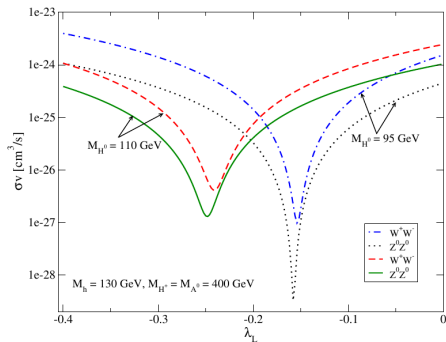


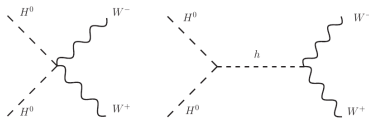
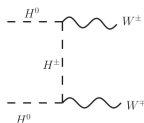
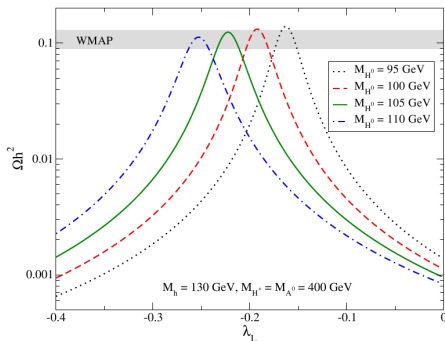
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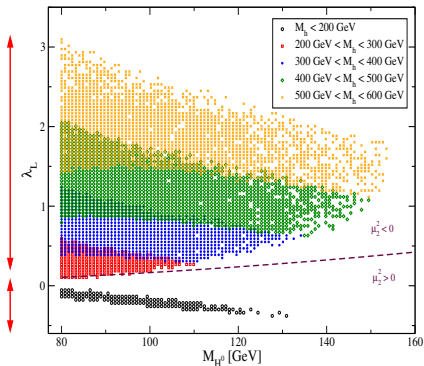


New viable parameter space thanks to cancellations

Result of a scan for $m_{H^0} > m_W$ and $\Omega_{H^0} = \Omega_{WMAP}$

$\lambda_L > 0$ $M_{H^0} < M_h/2$
 200 GeV $< M_h < 600$ GeV
 up to $M_{H^0} \sim 160$ GeV

$\lambda_L < 0$ $M_{H^0} > M_h/2$
 for $M_h < 200$ GeV
 up to $M_{H^0} = 130$ GeV

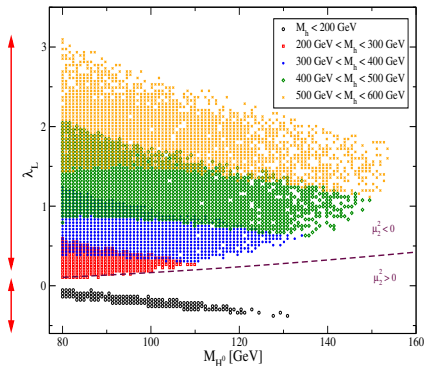


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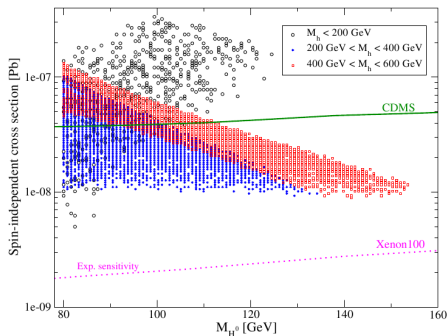


NB : IDM can comply with EWPT measurements for large M_h as Barbieri '06 :

$\Delta T_{H^0, H^+, A^0}$ can compensate negative T_h for $m_{H^+} > m_{A^0}, m_{H^0}$

Direct Detection searches

... A very efficient probe of the $m_{H^0} > m_W$ parameter space :



Remember :

$$\sigma_{H_0-N} \propto \left(\frac{\lambda_{H_0}}{M_{H_0} M_h^2} \right)^2$$

and for cancellations,

λ_L is necessarily **non zero**

\rightsquigarrow a large fraction of the parameter space is already ruled out by CDMS Ahmed '10
the remaining viable param. space is within the reach of Xenon 100 Aprile '10

Conclusion

The **Inert Doublet** is a **WIMP** with a rich **Scalar** DM phenomenology

- We have shown that for $m_{H_0} \lesssim m_W$, annihilation into **3 bdy final states (WW^*) MUST be taken** into account
- **Cancellations** between diagrams contributing to $H_0 H_0 \rightarrow ZZ, WW \rightsquigarrow$ **new viable region** from W -threshold up to $m_{H^0} \sim 130$ (160) GeV for $M_h < 200$ GeV (600 GeV)
- **The genuine DM viable parameter space** next to W threshold was derived
 - below W -threshold : λ_L is **reduced** up to $\mathcal{O}(10)$
 - above W -threshold : $\lambda_L > 0$ (< 0) for $m_{H^0} < M_h/2$ ($m_{H^0} > M_h/2$)
- **Prospects for DM detection are modified**
 - **Direct detection** :
 - below W -threshold σ_{H_0-N} **decreases** by a factor $\mathcal{O}(100)$
 - above W -threshold within the reach of Xenon 100

This is the End
Thank you for your attention !!

Backup

Important 3-body processes : Is that so surprising ?

3-body processes can take over 2-body processes

3-body \equiv real + virtual massive particle

e.g. $WW^* \rightarrow W\bar{f}f'$

Important 3-body processes : Is that so surprising ?

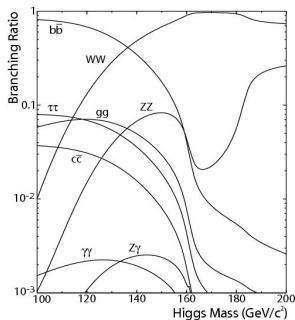
3-body processes can take over 2-body processes

3-body \equiv **real + virtual** massive particle

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Well known example : **higgs decay**

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3-body processes can enhance DM annihilation/decay :

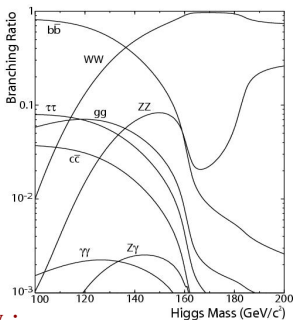
\rightsquigarrow Affect **relic abundance, viable parameter space, detection**

\rightsquigarrow **Significant effect** on : neutralino LSP [Chen & Kamionkowski JHEP '98, Yaguna PRD'10],

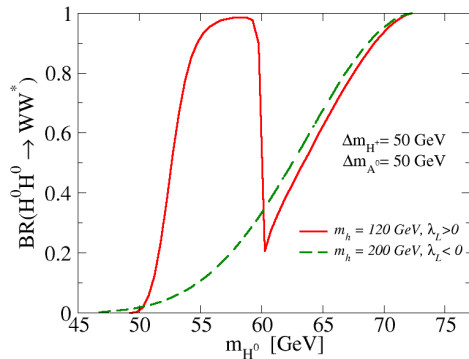
gravitino LSP [Choi & Yaguna '1003, & all '1007],

Higgs DM [Hosotani, Ko & Tanaka PLB'09], singlet scalar DM [Yaguna PRD'10],

Inert Doublet Model [LLH & Yaguna JHEP'10]



For Indirect Detection



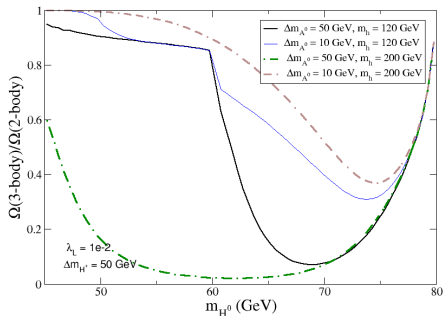
annihilations **no more $\bar{b}b$ dominated**

$$\rightsquigarrow BR(H_0H_0 \rightarrow WW^*) \sim 1$$

for m_{H^0} near W threshold

work in progress...

Coannihilation-fixed parameter



Above W threshold - fine tuning ?

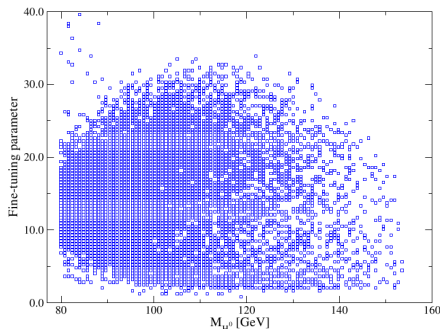
a_i ($i = 1, \dots, n$) are free parameters of a given dark matter model, the fine-tuning parameter

$$\Delta_{\Omega h^2, a_i} \equiv \frac{\partial \log \Omega h^2}{\partial \log a_i}.$$

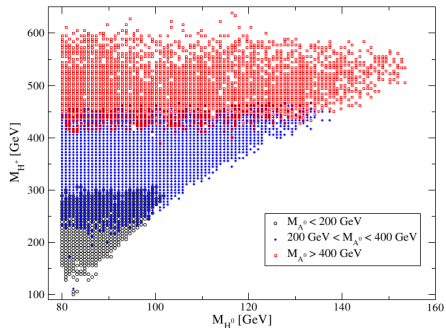
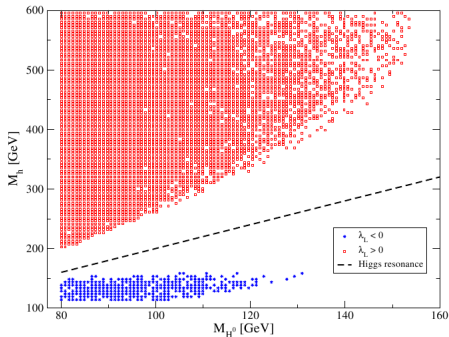
And the total fine-tuning, $\Delta_{\Omega h^2, \text{total}}$, is obtained by summing in quadrature the contributions of the different parameters of the model

$$\Delta_{\Omega h^2, \text{total}} = \sqrt{\Delta_{\Omega h^2, m_{H^0}}^2 + \Delta_{\Omega h^2, m_{H^\pm}}^2 + \Delta_{\Omega h^2, m_{A^0}}^2 + \Delta_{\Omega h^2, \lambda_L}^2 + \Delta_{\Omega h^2, M_h}^2}.$$

The fine-tuning parameter is large if a small variation in the parameters of the models leads to a large modification of the relic density. If, for instance, $\Delta_{\Omega h^2, \text{total}} \lesssim 10$, a measurement of the parameters of the model at the 10% level will enable to compute Ωh^2 to within a factor $\mathcal{O}(2)$. As a rule of thumb, one can say that the fine-tuning is small if $\Delta_{\Omega h^2, \text{total}} \lesssim 10$.

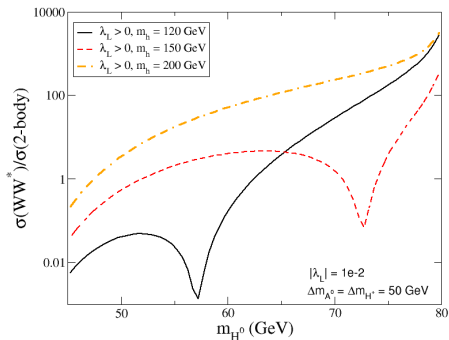


Cancellations require increasing M_h, m_{H^+}, m_{A^0}



Comparing 2-3 body annihilation cross section

3bdy annihilation dominates over 2 bdy on

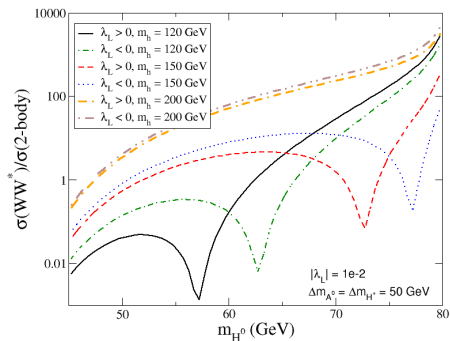


- $\lambda_L = 10^{-2}$

... a **significant** range of the parameter space, depend on m_h

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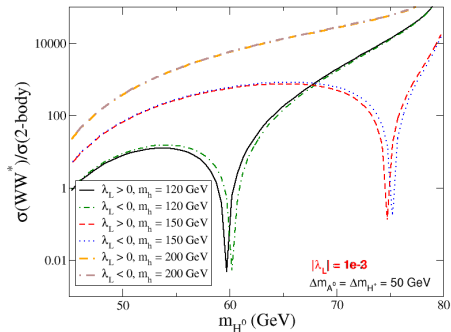


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Comparing 2-3 body annihilation cross section

3bdy annihilation **dominates** over 2 bdy on



- $\lambda_L = 10^{-2}$
... a **significant** range of the parameter space, depend on m_h
- $\lambda_L = 10^{-3}$
... the entire mass range independently of m_h
but not representative for $H_0 \equiv \text{DM}$

IDM : Mass Ranges

Mass Ranges	main contributions to σ_{eff}	mass splittings	main Refs
$m_{H_0} \ll m_W (\mathcal{O}(GeV))$	$H_0 H_0 \rightarrow h^* \rightarrow \bar{f} f$	$\Delta m_{ij} \gtrsim m_Z - m_{H_0} \sim 90 \text{ GeV}$	Andreas <i>et al</i> '08
$m_{H_0} \lesssim m_W$	$H_0 H_0 \rightarrow h^* \rightarrow \bar{f} f$ $H_0 A_0(H^+) \rightarrow Z^*(W^*) \rightarrow \bar{f} f^{(\prime)}$	$\Delta m_{ij} \gtrsim m_Z - m_{H_0} \gtrsim 7 \text{ GeV}$	Barbieri <i>et al</i> '06 LLH <i>et al</i> '06
$m_{H_0} \gg m_W (\mathcal{O}(TeV))$	$H_0 H_0 \rightarrow ZZ, WW, hh$ coannihil into bosons	$\Delta m_{ij} \lesssim 17.6 \text{ GeV}$	Hambye <i>et al</i> '09

How to conciliate Heavy Higgs and EWPT measurements ?

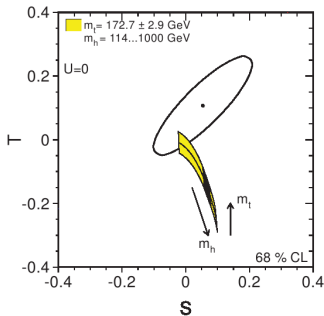
New physics affect EW observables

Contributions to EWPT measurement variable T from :

- Higgs : $T(M_h) = -\frac{3}{8\pi \cos^2 \theta_W} \ln \frac{M_h}{M_Z}$.

- H_2 scalars :

$$\Delta T \approx \frac{1}{24\pi^2 \alpha v^2} (M_{H^+} - M_{A_0})(M_{H^+} - M_{H_0})$$



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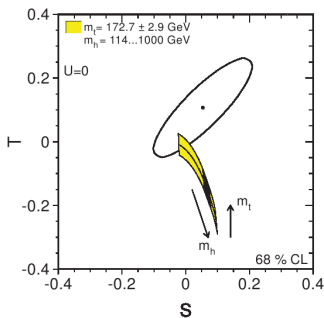
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↪ When $M_{H^+} > M_{A_0}, M_{H_0}$ positive contributions from ΔT can **compensate** the too large negative contributions from $T(M_h)$ due to heavy Higgs.

↪ With H_2 new physics one may push M_h up to **500-600 GeV** [Barbieri *et al* '06]



IDM : Potential - constraints

- Full Potential

$$V(H_1, H_2) = \mu_1^2 |H_1|^2 + \mu_2^2 |H_2|^2 + \lambda_1 |H_1|^4 + \lambda_2 |H_2|^4 \\ + \lambda_3 |H_1|^2 |H_2|^2 + \lambda_4 |H_1^\dagger H_2|^2 + \frac{\lambda_5}{2} \left[(H_1^\dagger H_2)^2 + h.c. \right]$$

- Dark scalars couplings to Higgs and masses :

$$\frac{1}{2} (\lambda_{H_0} H_0^2 + \lambda_{A_0} A_0^2 + 2\lambda_{H_c} H^+ H^-) (2v_0 h + h^2) \\ m_h^2 = 2\lambda_1 v_0^2, \quad m_i^2 = \mu_2^2 + \lambda_i v_0^2.$$

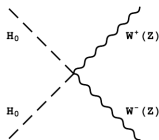
- Stability constraint

$$\lambda_{1,2} > 0, \\ \lambda_{H_0}, \quad \lambda_{A_0}, \quad \lambda_{H_c} > -\sqrt{\lambda_1 \lambda_2}.$$

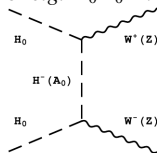
- EWPT measurements : $\Delta T \approx \frac{1}{12\pi^2 \alpha v^2} (m_{H^+} - m_{A_0})(m_{H^+} - m_{H_0})$

IDM in the High mass regime : Relic abundance

Quartic couplings ON : extra Higgs processes and $m_{H_0} \neq m_{A_0} \neq m_{H_c}$

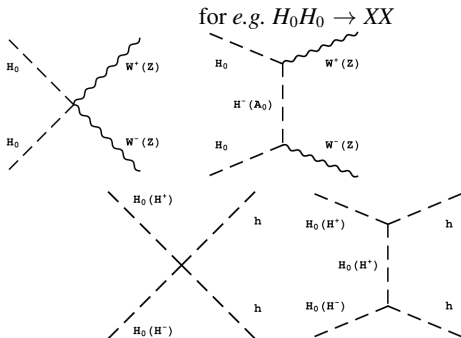


for e.g. $H_0H_0 \rightarrow XX$



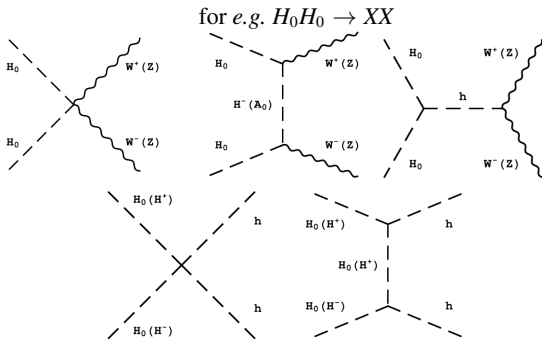
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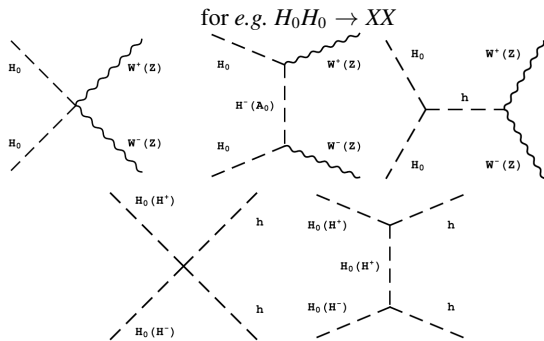


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$$\sigma_{\text{eff}} = \sum_{ij} \left(\sigma_g^{ij} + \sigma_\lambda^{ij} \right) \propto \frac{1}{m_{H_0}^2}$$

where $\sigma_\lambda^{ij} = \frac{\Lambda^{ij}}{m_{H_0}^2}$ with $\Lambda^{ij} \propto \lambda * \lambda$ and $\Lambda^{ij} > 0$



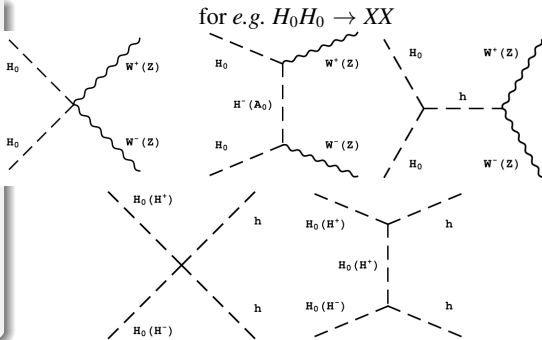
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- $m^* \sim 534 \text{ GeV}$ is minimal to satisfy WMAP



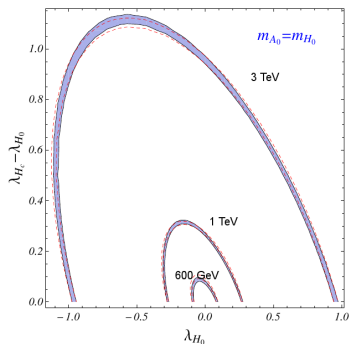
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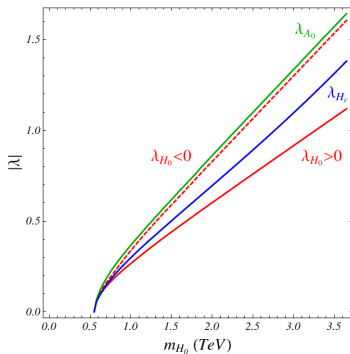
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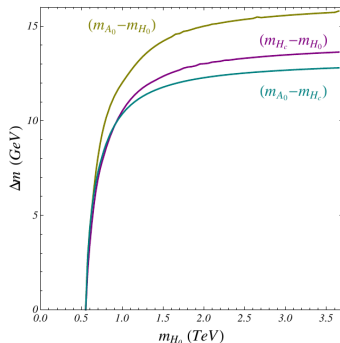
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- mass splittings are also bounded as
 $m_i - m_j \propto (\lambda_i - \lambda_j) / m_{H_0}$



3bdy effect on DM ?

3-body processes can enhance DM annihilation :

- supersymmetric dark matter :
 - neutralino LSP : Chen & Kamionkowski JHEP '98 study $\sigma v_{\nu \rightarrow 0}$ and impact on ν detection from annihilation in the Earth bellow WW and $\bar{t}t$ mass threshold
Yaguna PRD'10 demonstrate up to 10% effect on Ωh^2_{χ} for bino-like χ including $\bar{t}t^*$ (usually 2-bdy $\bar{b}b$ dom)
 - gravitino LSP : Choi & Yaguna '1003 W^*l and $Z^*\nu$ give significant (up to 90%) to \tilde{G} decay (usually 2-bdy $\gamma\nu$ dom)
Choi, Restrepo, Yaguna & Zapata '1007 gamma+antimatter signal [see Yaguna talk ! !]
- scalar DM
 - Higgs DM : Hosotani, Ko & Tanaka PLB'09 (gauge-Higgs unification)
 $\Omega_{DM} \rightsquigarrow m_{DM} = 75 \text{ GeV}$ (2bdy only) $\Rightarrow m_{DM} = 70 \text{ GeV}$ (including 3bdy)
 - singlet scalar DM : Yaguna PRD'10, $SS \rightarrow h \rightarrow WW^*$ enhance $\sigma v_{\nu \rightarrow 0}$ and reduce Ω_{DM} independently of S-higgs coupling

