## CPV and rare $B_{s}$ decays at the Tevatron



LaThuile<br>March 2011<br>Marj Corcoran<br>for the CDF and D0 collaborations

## Outline

## The $\mathrm{B}_{\mathrm{s}}$ system and CPV phases

Dimuon charge asymmetry (D0)

$$
\mathrm{B}_{\mathrm{s}} \rightarrow \mathrm{~J} / \psi \phi(\mathrm{D} 0 \text { and CDF })
$$

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B
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## CP violation the matter-antimatter asymmetry of the universe



In 1968 Andrei Sakharov showed that CP violation is necessary to generate the matter-antimatter asymmetry of the universe.

The CKM matrix gives rise to CP violation through the complex phase.

But the amount of CP violation we see in the CKM matrix is too small by 10 orders of magnitude.

## $B_{s}$ meson

In the $\mathrm{B}_{\mathrm{s}}$ system (as in the $\mathrm{B}_{\mathrm{d}}^{0}$ and $\mathrm{K}^{0}$ systems) CPV can arise through mixing via the box diagram:

or through interference of decay and mixing amplitudes through phases $\left(\phi_{s}\right)$,
or directly through the decay amplitudes (direct CPV).

## $\mathrm{B}_{\mathrm{s}}$ mixing

$$
i \frac{d}{d t}\binom{B_{\mathrm{s}}^{0}}{\bar{B}_{\mathrm{s}}^{0}}=\left(\begin{array}{ll}
M-\frac{i \Gamma}{2} & M_{12}-\frac{i \Gamma_{12}}{2} \\
M_{12}^{*}-\frac{i \Gamma_{12}^{*}}{2} & M-\frac{i \Gamma}{2}
\end{array}\right)\binom{B_{\mathrm{s}}^{0}}{\bar{B}_{\mathrm{s}}^{0}}
$$

The solution of the Schroedinger equation yields two mass eigenstates:

$$
\mathrm{B}_{\mathrm{s}}^{\mathrm{H}}=\mathrm{p}\left|\mathrm{~B}_{\mathrm{s}}>-\mathrm{q}\right| \overline{\mathrm{B}}_{\mathrm{s}}>\quad \text { and } \quad \mathrm{B}_{\mathrm{s}}^{\mathrm{L}}=\mathrm{p}\left|\mathrm{~B}_{\mathrm{s}}>+\mathrm{q}\right| \overline{\mathrm{B}}_{\mathrm{s}}>
$$

If $\mathrm{p} \neq \mathrm{q}$, the mass eigenstates are not CP eigenstates, and $\overline{\mathrm{B}} \rightarrow \mathrm{B} \neq \mathrm{B} \rightarrow \overline{\mathrm{B}}$

## $\mathrm{B}_{\mathrm{s}}$ mixing parameters

$\Delta \mathrm{m}_{\mathrm{s}}=\mathrm{M}_{\mathrm{H}}-\mathrm{M}_{\mathrm{L}}=2\left|\mathrm{M}_{12}\right|=17.77 \pm 0.12 \mathrm{ps}^{-1}$ (well measured by CDF)
$\phi_{\mathrm{s}}^{\mathrm{SM}}=\arg \left[-\mathrm{M}_{12} / \Gamma_{12}\right]=0.0042 \pm 0.0014$ (SM expectation)
$\Delta \Gamma=\Gamma_{\mathrm{L}}-\Gamma_{\mathrm{H}}=2\left|\Gamma_{12}\right| \cos \phi_{\mathrm{s}}=0.096 \pm 0.039 \mathrm{ps}^{-1}$ (SM expectation)
$\phi_{\mathrm{s}}^{\mathrm{SM}}$ is small, so measurements of $\phi_{\mathrm{s}}$ are sensitive to new physics

$$
\phi_{\mathrm{s}}=\phi_{\mathrm{s}}^{\mathrm{SM}}+\phi_{\mathrm{s}}^{\mathrm{NP}}
$$

## $\mathrm{B}_{\mathrm{s}} \rightarrow \mathrm{J} / \psi \phi$

$B_{s}^{0} \longrightarrow J / \psi \phi$
For $\mathrm{B}_{\mathrm{s}} \rightarrow \mathrm{J} / \psi \phi$ ，interference of direct decay and decay after mixing leads to a phase $\beta_{\mathrm{s}}$ or $\phi^{J / w \phi}$ which is also small in the SM．

$$
\phi_{\mathrm{s}}^{J / v \phi}=-2 \beta_{\mathrm{s}}=-2 \arg \left[\mathrm{~V}_{\mathrm{ts}} \mathrm{~V}_{\mathrm{tb}}^{*} / \mathrm{V}_{\mathrm{cs}} \mathrm{~V}_{\mathrm{cb}}^{*}\right]=0.038 \pm 0.002
$$

New physics could effect both $\phi_{s}$ and $\phi_{s}^{J / \varphi \phi}$

$$
\begin{aligned}
& \phi_{\mathrm{s}}^{\mathrm{J} / \varphi \phi}=-2 \beta_{\mathrm{s}}^{\mathrm{SM}}+\phi_{\mathrm{s}}^{\mathrm{NP}} \\
& \phi_{\mathrm{s}}=\phi_{\mathrm{s}}^{\mathrm{SM}}+\phi_{\mathrm{s}}^{\mathrm{NP}}
\end{aligned}
$$

Note that CDF uses $\beta_{\mathrm{s}}$ and D0 uses $\phi_{\mathrm{s}}^{\mathrm{J} / ⿲ 丨 丨 又}$

## The Tevatron performs!



Both CDF and D0 have received more than $10 \mathrm{fb}^{-1}$ of integrated luminosity, and both have been running at $\sim 90 \%$ efficiency.

## Dimuon Charge Asymmetry

## Dimuon charge asymmetry



Decay of B tags its flavor
The quantity being measured is simple, and can be nonzero from B mixing.

$$
A_{s l}^{b} \equiv \frac{N_{b}^{++}-N_{b}^{--}}{N_{b}^{++}+N_{b}^{--}}
$$

If the rate of $\mathrm{B} \rightarrow \overline{\mathrm{B}}$ is the same as $\overline{\mathrm{B}} \rightarrow \mathrm{B}$, this quantity will be zero. The SM expectation is nearly zero.

## Semileptonic muon charge asymmetry

We also measure the 'wrong sign' semileptonic charge asymmetry for inclusive muons:

$$
a_{\mathrm{sl}}^{b}=\frac{\Gamma\left(\bar{B} \rightarrow \mu^{+} X\right)-\Gamma\left(B \rightarrow \mu^{-} X\right)}{\Gamma\left(\bar{B} \rightarrow \mu^{+} X\right)+\Gamma\left(B \rightarrow \mu^{-} X\right)}=A_{\mathrm{sl}}^{b}
$$

The 'right sign' decays are $\mathrm{B} \rightarrow \mu^{+} \mathrm{X}$ and $\overline{\mathrm{B}} \rightarrow \mu^{-} \mathrm{X}$ The 'wrong sign' decays can only occur through mixing.

## Semileptonic and dimuon asymmetries

We measure a linear combination of $\mathrm{B}_{\mathrm{d}}$ and $\mathrm{B}_{\mathrm{s}}$ decays

$$
A_{\mathrm{sl}}^{b}=(0.506 \pm 0.043) a_{\mathrm{sl}}^{d}+(0.494 \pm 0.043) a_{\mathrm{sl}}^{s}
$$

dimuon semileptonic $\mathrm{B}_{\mathrm{d}}$ semileptonic $\mathrm{B}_{\mathrm{s}}$

The fractions of $\mathrm{B}_{\mathrm{d}}$ and $\mathrm{B}_{\mathrm{s}}$ are known from other measurements. Upper case A refers to dimuon quantities; lower case to single muon inclusive quantities.

## Analysis strategy

Measure the raw asymmetries

$$
\begin{array}{cc}
A=\frac{N\left(\mu^{+} \mu^{+}\right)-N\left(\mu^{-} \mu^{-}\right)}{N\left(\mu^{+} \mu^{+}\right)+N\left(\mu^{-} \mu^{-}\right)} & a=\frac{n\left(\mu^{+}\right)-n\left(\mu^{-}\right)}{n\left(\mu^{+}\right)+n\left(\mu^{-}\right)} \\
3.7 \times 10^{6} \text { events } & 1.5 \times 10^{9} \text { events }
\end{array}
$$

Correct for backgrounds and detector related asymmetries almost entirely from the data
Correct for non-B contributions and non-oscillating contributions from simulation
Use the two independent measurements to cancel some systematics

## Analysis strategy

$\mathrm{k}=0.041 \pm 0.003$


Determine dilution factors from simulation
Blind analysis! The central value of $\mathrm{A}_{\mathrm{sl}}{ }^{\mathrm{b}}$ from the full data set was not revealed until all the background subtraction and systematics had been finalized.

## Kaon background

The largest background comes from $\mathrm{K} \rightarrow \mu \nu$, which is asymmetric for $\mathrm{K}^{+} / \mathrm{K}$. The resulting fake asymmetry is measured from the data, using $\phi \rightarrow \mathrm{KK}$ and $\mathrm{K}^{0^{*}} \rightarrow \mathrm{~K} \pi$.

Processes such as $\mathrm{K}^{-} \mathrm{p} \rightarrow \Lambda \pi^{0}$ gives $\mathrm{K}^{-}$a much larger cross section than $\mathrm{K}^{+}$. So more $\mathrm{K}^{+}$ survive to decay.


## Kaon asymmetry




The K asymmetry is measured two ways, in complete agreement. The average is used.

Background asymmetries from protons and pions are 20x smaller and are determined from $\Lambda \rightarrow \mathrm{p} \pi$ and $\mathrm{K}_{\mathrm{s}} \rightarrow \pi \pi$.

## Background fractions

The background fractions are determined mainly from the data, but a comparison with MC truth gives us confidence that everything hangs together.

|  | $\left(1-f_{\text {bkg }}\right)$ | $f_{K}$ | $f_{\pi}$ | $f_{p}$ |
| :---: | :---: | :---: | :---: | :---: |
| MC | $(59.0 \pm 0.3) \%$ | $(14.5 \pm 0.2) \%$ | $(25.7 \pm 0.3) \%$ | $(0.8 \pm 0.1) \%$ |
| Data | $(58.1 \pm 1.4) \%$ | $(15.5 \pm 0.2) \%$ | $(25.9 \pm 1.4) \%$ | $(0.7 \pm 0.2) \%$ |

The fractions of single muons coming from background sources, comparing data and Monte Carlo 'truth'.

Many other cross-checks were done.

## Track reconstruction



## Swapping Magnet Polarity <br> $\longrightarrow$

The detector can give a fake asymmetry. Reversal of the magnet polarities is important for cancelling these instrumental asymmetries.

## Cancellation of systematics

There are two measurements of the charge asymmetry-one from dimuons and one from single inclusive muons Reduce the systematic uncertainties by a cancellation using the dimuon and the single muon charge asymmetries.


## Dimuon charge asymmetry

## Measurement

$$
A_{s l}^{\mathrm{b}}=[-0.957 \pm 0.251 \text { (stat.) } \pm 0.146 \text { (syst.) }] \%
$$

Expectation $\rightarrow \mathbf{A}_{\text {sl }}^{\mathbf{b}}(\mathbf{S M})=\left[-\mathbf{0 . 0 2 3}{ }_{-0.006}^{+0.005}\right] \%$

The final result is $3.2 \sigma$ from the SM expectation.
CDF is working on a similar analysis (we eagerly await...) D0 is working on an update with more data, better background treatment, and asymmetry vs. impact parameter.


The quantity we measure is a linear combination of the asymmetry from $B_{d}$ and $\mathrm{B}_{\mathrm{s}}$ mesons.

$$
A_{\mathrm{sl}}^{b}=(0.506 \pm 0.043) a_{\mathrm{sl}}^{d}+(0.494 \pm 0.043) a_{\mathrm{sl}}^{s} .
$$

"A New Clue to Explain Existence" New York Times
Joe Lykken, a theorist at Fermilab, said, "So I would not say that this announcement is the equivalent of seeing the face of God, but it might turn out to be the toe of God."



Guennadi Borissov Lancaster Univ. UK

Bruce Hoeneisen
Universided San Franciseo de Ouito Ecuador
V. A. Abazov et al, Phys. Rev. D 82, 032001 (2010)
V. A. Abazov et al, Phys. Rev. Lett. 105, 081801 (2010)


## $\mathrm{B}_{\mathrm{s}} \rightarrow \mathrm{J} / \psi \phi$

A previous CDF/D0 combination had a $2.1 \sigma$ discrepancy with the SM .

See CDF public note 10206
See D0 public note 6098

## $\mathrm{B}_{\mathrm{s}} \rightarrow \mathrm{J} / \psi \phi$



$$
\frac{d^{3} \Gamma\left[B_{s}^{0}(t) \rightarrow J / \psi\left(\rightarrow \mu^{+} \mu^{-}\right) \phi\left(\rightarrow K^{+} K^{-}\right)\right]}{d \cos \theta d \varphi d \cos \psi} \propto
$$

$$
2 \cos ^{2} \psi\left(1-\sin ^{2} \theta \cos ^{2} \varphi\right)
$$

$$
\left|A_{0}(t)\right|^{2}
$$

$$
+\sin ^{2} \psi\left(1-\sin ^{2} \theta \sin ^{2} \varphi\right)
$$

$$
\left|A_{\|}(t)\right|^{2}
$$

$$
+\sin ^{2} \psi \sin ^{2} \theta
$$

$$
+\frac{1}{\sqrt{2}} \sin 2 \psi \sin ^{2} \theta \sin 2 \varphi \operatorname{Re}\left(\left|A_{0}^{*}(t) \| A_{\|}(t)\right|\right)
$$

Scalar $\rightarrow$ two vectors, $L=0,1,2$ allowed
Complicated fit to time-dependent angular distributions.

$$
+\frac{1}{\sqrt{2}} \sin 2 \psi \sin 2 \theta \cos \varphi \quad \operatorname{Im}\left(\left|A_{0}^{*}(t) \| A_{\perp}(t)\right|\right)
$$

$-\sin ^{2} \psi \sin 2 \theta \sin \varphi \quad \operatorname{Im}\left(\left|A_{\|}(t) \| A_{\perp}(t)\right|\right)$

## $\mathrm{B}_{\mathrm{s}} \rightarrow \mathrm{J} / \psi \phi$




CDF has uses both same-side and opposite-side tagging. A neural net is used for signal selection, and a multidimensional max-likelihood fit extracts the parameters.

## $\mathrm{B}_{\mathrm{s}} \rightarrow \mathrm{J} / \psi \phi$




New CDF result agrees with the SM at about $1 \sigma$ They also have excellent measurements of $\mathrm{B}_{\mathrm{s}}$ lifetime, $\Delta \Gamma_{\mathrm{s}}$ and the polarization amplitudes. S-wave contribution in $\mathrm{KK}<6.7 \%$ at $95 \% \mathrm{CL}$.

## $\mathrm{B} \rightarrow \mathrm{J} / \psi \phi$

D0 analysis is similar: multidimensional unbinned likelihood fit, opposite- side tagging,strong phases fixed to $\mathrm{B}_{\mathrm{d}}$ values $3435 \pm 84$ signal events




## $\mathrm{B}_{\mathrm{s}} \rightarrow \mathrm{J} / \psi \phi$




D0s new result is a bit further from the SM than the older result, and in good agreement with the dimuon charge asymmetry.

## $\mathrm{B}_{\mathrm{s}} \rightarrow \mathrm{J} / \psi \phi$

Both experiments agree that results can vary a lot, even with data sets of the same size!


CDF results for different data periods-significant variation

## Pseudoexperiments

CDF pseudoexperiments equivalent to $1.4 \mathrm{fb}^{-1}$, all with the same initial parameters-a lot of variation.


From Gavril Giurgui's talk at ICHEP2010

## D0 pseudoexperiments



Distribution of errors from
2000 pseudoexperiments with
the same statistics as the data.
The two experiments are consistent.

D0 is working on finalizing the 2010 result with BDT selection and including $s$-wave in the KK system.

CDF result (translated to $\phi_{\mathrm{s}}^{\mathrm{J} / \mathrm{M}}$ )

$$
[-0.04,-1.04] \mathrm{U}[-2.16,-3.10] \text { at } 68 \% \mathrm{CL}
$$

## $\mathrm{B}_{\mathrm{s}} \rightarrow \mu \mu$



D0 has a result from 2010 on $6.1 \mathrm{fb}^{-1}$
CDF has a result from 2009 on $3.7 \mathrm{fb}^{-1}$

## $\mathrm{B}_{\mathrm{s}} \rightarrow \mu \mu$


(b)


SM expected branching ratio is well known, $\mathrm{BR}=3.6 \pm 0.3 \times 10^{-9}$ Many models of physics beyond the SM enhance this BR by large factors, lots of room for new physics.
Some models that lead to new physics in $\mathrm{B}_{\mathrm{s}}$ mixing also enhance this decay .

## Background suppression



## $\mathrm{B}_{\mathrm{s}} \rightarrow \mu \mu$

In highest sensitivity region, $51 \pm 4$ bkgnd expected, 55 events seen.
$\mathrm{BF}<51 \times 10^{-9}$ at $95 \% \mathrm{CL}$ Expected limit $38 \times 10^{-9}$ PLB 695, 539 (2010)

CDF preliminary 2009 $43 \times 10^{-9}$ observed $33 \times 10^{-9}$ expected


Both CDF and D0 are working on NN output updates.

## Summary

D0 has measured the dimuon charge asymmetry and finds a $3.2 \sigma$ discrepancy with the SM. D0 is working on an update, and CDF is working on a similar result.

Both CDF and D0 have updated results for the CP violating phase in $\mathrm{B}_{\mathrm{s}} \rightarrow \mathrm{J} / \psi \phi$. CDF is closer to the SM, D0 is a bit further from the SM, but within errors there is no discrepancy. Errors do not seem to scale with $1 / \sqrt{ } N$.

The current limit on $\mathrm{B}_{\mathrm{s}} \rightarrow \mu \mu$ is about 10 x SM expectation. Both CDF and D0 are working on updates for this decay.


## $\mathrm{B}_{\mathrm{s}} \rightarrow \mu \mu$ Signal and backgrounds



## Background K fractions


fraction along with MC is used to determine the $\pi$ and p fractions.

## Final background determination

(statistical uncertainties only)

| $\left(1-f_{\mathrm{bkg}}\right)$ | $f_{K}$ | $f_{\pi}$ | $f_{p}$ |
| :---: | :---: | :---: | :---: |
| $(58.1 \pm 1.4) \%$ | $(15.5 \pm 0.2) \%$ | $(25.9 \pm 1.4) \%$ | $(0.7 \pm 0.2) \%$ |
|  | $a_{K} f_{K}$ | $a_{\pi} f_{\pi}$ | $a_{p} f_{p}$ |
|  | $(+0.854 \pm 0.018) \%$ | $(+0.095 \pm 0.027) \%$ | $(+0.012 \pm 0.022) \%$ |
|  | $A_{K} F_{K}$ | $A_{\pi} F_{\pi}$ | $A_{p} F_{p}$ |
|  | $(+0.828 \pm 0.035) \%$ | $(+0.095 \pm 0.025) \%$ | $(+0.000 \pm 0.021) \%$ |

$$
\begin{aligned}
a-a_{\mathrm{bkg}} & =k A_{\mathrm{sl}}^{b} \\
A-A_{\mathrm{bkg}} & =K A_{\mathrm{sl}}^{b}
\end{aligned}
$$

- Multiple physics processes contribute to the inclusive and like-sign dimuon samples

| Process |  |  |
| :--- | :--- | :--- |
| $T_{1}$ | $b \rightarrow \mu^{-} X$ |  |
| $T_{1 a} \quad b \rightarrow \mu^{-} X$ (non-oscillating) |  |  |
| $T_{1 b} \quad \bar{b} \rightarrow b \rightarrow \mu^{-} X$ (oscillating) $\rightarrow A$ | $A_{\mathrm{sl}}^{b}$ |  |
| $T_{2} \quad b \rightarrow c \rightarrow \mu^{+} X$ |  |  |
| $T_{2 a} \quad b \rightarrow c \rightarrow \mu^{+} X$ (non-oscillating) |  |  |
| $T_{2 b} \quad \bar{b} \rightarrow b \rightarrow c \rightarrow \mu^{+} X$ (oscillating) | $A_{\mathrm{sl}}^{b}$ |  |
| $T_{3} \quad b \rightarrow c \bar{c} q$ with $c \rightarrow \mu^{+} X$ or $\bar{c} \rightarrow \mu^{-} X$ |  |  |
| $T_{4}$ | $\eta, \omega, \rho^{0}, \phi(1020), J / \psi, \psi^{\prime} \rightarrow \mu^{+} \mu^{-}$ |  |
| $T_{5}$ | $b \bar{b} c \bar{c}$ with $c \rightarrow \mu^{+} X$ or $\bar{c} \rightarrow \mu^{-} X \rightarrow A$ |  |
| $T_{6}$ | $c \bar{c}$ with $c \rightarrow \mu^{+} X$ or $\bar{c} \rightarrow \mu^{-} X$ |  |

$$
\begin{aligned}
k & =0.041 \pm 0.003 \\
K & =0.342 \pm 0.023
\end{aligned}
$$

## Sanity check



## Asymmetry vs mass



Measured and expected uncorrected asymmetry vs dimuon mass. The expected (red) values uses all the measured background contributions.

## Systematic checks



This plot shows the difference from the central value for 16 analysis variations.
The raw asymmetries changed by as much as $150 \%$ with these variations.
The variations include:
more stringent muon selection
variations in trigger requirements
different angular regions different $\mu \mu$ mass cuts


## D0 combination

$$
\begin{aligned}
\bar{\tau}_{s} & =1.45 \pm 0.04 \pm 0.01 \mathrm{ps} \\
\Delta \Gamma_{s} & =0.15 \pm 0.06 \pm 0.01 \mathrm{ps}^{-1} \\
\phi_{s}^{J / \psi \phi} & =-0.76_{-0.36}^{+0.38} \pm 0.02
\end{aligned}
$$

