# Search for neutral Higgs bosons decaying into four $\tau$ at LEP2





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#### Outline

Introductory comments

- Motivations for a "new" Higgs search at LEP
- ALEPH data and Montecarlo
- Event Selections
- Interpretations and Results
- Conclusions

#### **NEW RESULTS: FIRST TIME PUBLIC** HERE @ La Thuile !!!!

- LEP stopped in 2000: after 10 years there are still some uncovered analyses
- Theoretical hints for a missed low mass Higgs in a region NOT accessible to the LHC program
- Hard to look at the archived data. It was an archaeological exercise
- Results are the subject of a recently submitted preprint (2 days ago!!!) by the ALEPH Collaboration: arXiv:1003.0705v1 (to be submitted to JHEP)

#### LEP

#### LEP operated from 1989-2000

- LEP1 running at the Z resonance (<1996)
- LEP2 running from  $\sqrt{s} = 183 207 \,\mathrm{GeV}$









At 95% CL,  $m_{h_{
m SM}} < 157~{
m GeV}$  and the  $\Delta\chi^2$  minimum is near 85  ${
m GeV}$  when all data are included.

# Higgs Searches

- Electroweak fits prefer a Higgs lighter than 150 GeV
   No indication of new Physics below the TeV scale
- Introduce fine tuning problems for SM and MSSM

#### LEP paradox:

 Hard maintain naturalness if m<sub>H</sub><150 GeV and new Physics (SUSY) at the TeV scale

This has motivated theories with extended Higgs sectors or next-to-minimal supersymmetric extension to the SM

#### How could we have missed the Higgs at LEP?

- if the production cross-section were smaller than expected
  - this has direct implications on how the Higgs couples to the Z and its role in EWSB
- or maybe it decayed into something exotic that the standard analysis missed
  - Is that difficult to achieve?
  - No, the Hbb coupling is quite small. It doesn't take much for a new decay mode to dominate the bb mode.
  - would the existing analyses have seen it?
    - it depends: in some existing searches may still be sensitive
    - but there is no easy and accurate way to determine the efficiency of existing analyses to alternative models

#### LEP limits for $H_2Z \rightarrow H_1H_1$ in the $H_1H_2$ plane (for issance H1 = pseudoscalar *a*)



(factor x SM cross section that corresponds to 95% exclusion)

Here we see that Higgs bosons produced via Higgsstrahlung decaying to 4b are highly constrained

+  $4\tau$  are less constrained with a notable hole for m<sub>h</sub>>85 &  $2m_{\tau} < m_a < 10 \, {\rm GeV}$ 

# Higgs reconstruction in 4 taus



4-vector constraint for neutrinos

# $\begin{array}{c} \mathsf{ALEPH} \\ \textbf{Simulated event} \end{array} e^+e^- \to ZH \to 2e4\tau \end{array}$

2 back-to-back electrons clearly distinguished from 2 back-to-back jets. not much else in the event (about 50 GeV of missing energy)



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## Monte Carlo simulation

- More than 10 years of running in a very clean environment and tuning MC to data
- Very good description of SM processes in the ALEPH MC
- 2 important background processes for this search:
  - 4 fermions





#### Monte Carlo simulation

Table 2: Details on Standard Model background processes and their categorisation. Fragmentation, hadronization and final state radiation were simulated with PYTHIA 6.1 [15]. See Ref. [16] for more details.

	Category	Process	Software		
	2f	$e^+e^- \rightarrow \mathbf{Z}/\gamma^* \rightarrow q\bar{q}(\gamma)$	KK 4.14 [17]		
		Bhabha and $e^+e^- \to \mathbf{Z}/\gamma^* \to e^+e^-(\gamma)$	BHWIDE 1.01 [18]		
		$e^+e^- \to \mathbf{Z}/\gamma^* \to \mu^+\mu^-(\gamma)$	KK 4.14 [17]		
		$e^+e^-  ightarrow { m Z}/\gamma^*  ightarrow  au^+  au^-(\gamma)$	KK 4.14 [17]		
	4f	$e^+e^- \rightarrow \mathbf{Z}/\gamma^* \rightarrow W^+W^-$	KoralW 1.51 [19]		
		$e^+e^- \rightarrow ZZ$	Pythia 6.1 [15]		
		$e^+e^- \rightarrow {\rm Z}  e^+e^-$	Pythia 6.1 [15]		
		$e^+e^- \rightarrow Z  \nu \bar{\nu}$	Рутніа 6.1 [15]		
		$e^+e^- \rightarrow W^{\pm}e^{\mp}\nu$	Pythia 6.1 [15]		
		$\gamma\gamma  ightarrow \ell^+\ell^-$	PHOT02 [20]		
	ſΎ́Υ	$\gamma\gamma \to q\bar{q}$	PHOT02 [20]		
	$\mathrm{n}\gamma$	$e^+e^- \to \mathbf{Z} \to \nu \bar{\nu}(\gamma)$	Pythia 6.1 [15]		
		$e^+e^-  ightarrow n\gamma$	Pythia 6.1 [15]		



## $H \rightarrow 4\tau$ and $Z \rightarrow \mu^+\mu^-$ , $e^+e^-$

#### LOOSE

- Two muons /electrons identified with opposite charge
- $|\cos\theta_{j1}| < 0.9$ ,  $|\cos\theta_{j2}| < 0.9$  for a proper containment of the jet in the tracking volume
- $\cos\theta_{jl} < 0.95$  for a good lepton isolation
- Missing Energy > 20 GeV
- Additional photons from f.s.r. added

#### **FINAL**

80 < M<sub>Z</sub> < 102 GeV/c<sup>2</sup>
 cos θ<sub>j1j2</sub> < 0 jets back to back</li>
 N<sub>track</sub> = 2 or 4 for each jet



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# Our signal efficiency is pretty good, but clearly we have very few events in lepton channels

but we also have almost no background in lepton channels



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# Expectation for $m_h = 100 \text{ GeV}$



#### Background contributions for ee channel

Total bg	0.52
$1 \mathrm{ph}$	0.00
2ph-Gss	0.00
4f	0.05
Bhabha	0.23
2ph-Gtt	0.00
PZZ	0.04
KZZ	0.07
2ph-Gbb	0.00
2ph-Gud	0.00
PZe	0.01
2ph-Gcc	0.00
mu+mu-	0.00
Znn	0.00
2ph-Gee	0.00
qqbar	0.07
2ph-Gmm	0.00
tau+tau-	0.05
$\operatorname{multiph}$	0.00

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# Discovery potential: $5\sigma$ sensitivity in an un-excluded region



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## Loose selection: $H \to 4\tau$ , $~Z \to \mu^+\mu^-$



### Final selection: $H \to 4 \tau \ , \ Z \to \mu^+ \mu^-$



#### Loose selection $H{\rightarrow}~4\tau$ , $~~Z{\rightarrow}\nu\nu$



#### Final selection $H{\rightarrow}~4\tau$ , $~Z{\rightarrow}\nu\nu$



#### Selected events: No Excess observed

Table : Number of events passing loose and final selections in each channel, in data, simulated background, and simulated signal ( $m_h = 100$ ,  $m_a = 4 \text{ GeV}/c^2$ ). The numbers of events passing the final selection are categorised by track-multiplicity.

Channel	Selection	data	total	background category		signal		
	$(n_1^{\mathrm{track}}, n_2^{\mathrm{track}})$		background	2f	4f	$\gamma\gamma$	$\mathrm{n}\gamma$	
	Loose	299	332	183	137	12.31	0.65	2.27
Z sate	(2,2)	0	0.034	0.034	0.000	0.000	0.000	0.689
$\Sigma \rightarrow e^+ e^-$	(2,4)+(4,2)	0	0.055	0.014	0.005	0.037	0.000	0.610
	(4,4)	0	0.031	0.019	0.013	0.000	0.000	0.126
	Loose	83	74.50	12.79	60.64	1.07	0.00	2.37
$7 \rightarrow u^+ u^-$	(2,2)	0	0.058	0.005	0.053	0.000	0.000	0.800
$\Sigma \rightarrow \mu^+ \mu$	(2,4)+(4,2)	0	0.005	0.000	0.005	0.000	0.000	0.676
	(2,2)	0	0.006	0.000	0.006	0.000	0.000	0.127
	Loose	206	193	135	43.93	10.41	3.74	12.63
7	(2,2)	0	1.312	0.663	0.408	0.240	0.000	5.097
$\Delta \rightarrow \nu \nu$	(2,4)+(4,2)	0	1.948	0.528	0.575	0.845	0.000	4.741
	(4,4)	2	2.569	0.461	0.820	1.288	0.000	1.089

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$$\xi^2 = \frac{\sigma \operatorname{BR}(h \to aa) \operatorname{BR}(a \to \tau \tau)^2}{\sigma_{SM}}$$



Contour of observed 95% confidence level limit on  $\xi^2$  in the  $m_a - m_h$  plane.

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- A search for the Higgs boson produced via higgsstrahlung at LEP2 energies has been performed, where  $h \rightarrow aa \rightarrow \tau^+\tau^- \tau^+\tau^-$  and  $Z \rightarrow \ell \ell = \mu^+\mu^-$ ,  $e^+e^-$ ,  $\nu\nu$
- No evidence for an excess of events above background was observed limit on the combined production cross section times branching ratio,  $\sigma(e^+e^- \rightarrow Zh) \times B(h \rightarrow aa) \times B(a \rightarrow \tau^+\tau^-)^2$  is presented.
- For  $m_h < 107 \text{ GeV/c}^2$  and  $4 < m_a < 12 \text{ GeV/c}^2$  the quantity  $\xi^2 < 1$  is excluded at 95% CL
  - This analysis covers a region of parameter space previously left unexplored: namely  $86 < m_h < 107 \text{ GeV/c}^2$  and  $4 < m_a < 12 \text{ GeV/c}^2$ , which further constrains models with non-standard Higgs decays, such as the NMSSM

# Acknowledgments

- Thanks to Neal Weiner and Riccardo Barbieri who triggered this work
- Thanks to Kyle, Itay and James from NYU who did most of the job
- Thanks to the old ALEPH Collaboration, especially the editorial board, for help, suggestions and review of the paper

# Back-up slides.....

# Limits vs Theory



# Setting Confidence Limits: CL<sub>b</sub>

$$CL_b = (\sum_{i=0}^{n} [e^{-(b)} \frac{(b)^i}{i!}])$$

CL<sub>b</sub> is close to one in presence of big fluctuation of bkg (signal?)

It turns out that  $CL_b$  is also a useful tool to analyse results where an excess of events is observed. Indeed in this case  $1-CL_b$  gives the probability that the observed excess is due to a positive fluctuation of the background and a very low value of this quantity gives evidence of a new phenomenon.

## Setting Confidence Limits

$$CL_{s+b} = \left(\sum_{i=0}^{n} \left[e^{-(s+b)} \frac{(s+b)^{i}}{i!}\right]\right)$$

$$CL_s = \frac{CL_{s+b}}{CL_b}$$

is used to set the exclusion limits. A signal hypothesis is therefore excluded at a certain confidence level CL when

 $1 - CL_s \leq CL.$ 

CL<sub>s</sub> is close to one when we are not able anymore to distinguish signal from background = signal NOT excluded

### Setting Confidence Limits

• Method: 3 final states  $f = \mu^+\mu^-$ ,  $e^+e^-$ ,  $\nu\nu$ 

three permutations  $(n_1^{track}, n_2^{track}) \in \{(2, 2), (2, 4) \text{ or } (4, 2), (4, 4)\} \equiv \mathcal{M}$ 

event count  $N_{m,f}$  in each of these nine categories is modeled with a Poisson distribution about the sum of the uncertain background  $b_{m,f}$  and the expected signal  $s_{m,f}$  scaled by  $\xi^2$ . A normal distribution is used to model the relationship between the uncertain background, the Monte Carlo-based background estimate  $b_{m,f}^{MC}$ , and its systematic uncertainty  $\Delta_f$ . This procedure leads to the following joint probability density for the event counts:

 $P(N_{m,f}|\xi^2, b_{m,f}) = \prod_{m \in \mathcal{M}} \prod_{f \in \{ee,\mu\mu,\nu\nu\}} \operatorname{Pois}(N_{m,f}|\xi^2 s_{m,f} + b_{m,f}) \cdot N(b_{m,f}^{MC}|b_{m,f}, \Delta_f).$ (1)

# Choice of Jet Algorithms

At LEP, the dominant jet algorithms were DURHAM and JADE.

- both are iterative recombination type algorithms: merge if  $m_{ij}^2/E_{tot}^2 < y_{cut}$ 
  - $y_{\text{cut}}$  is an adjustable parameter and  $E_{\text{tot}}$  was often chosen to be the visible energy in the event
- DURHAM defines  $m_{ij}^2$  in a way that is more robust to soft radiation, which is good if you are interested in bona fide hadronic showers.
  - But we are looking for a purely electroweak decay, so the straight invariant mass combination of JADE is more natural.
  - Furthermore, we know that we are interested in  $m_a < 10 \,{\rm GeV}$  which leads to an obvious choice for y<sub>cut</sub> if we use a fixed E<sub>tot</sub>.

By choosing this approach our s/b was significantly higher than forcing to two jets with DURHAM and cutting on the jet mass

• Additionally we have track multiplicity in jets as a handle

# **ALEPH Higgs Searches**

	$\sqrt{s}$ (GeV)	$\mathcal{L}\left(\mathbf{pb}^{i+1}\right)$	Mass range $(GeV/c^2)$	Ref.
$\mathcal{H}_1\mathbb{Z} \to (\dots) \ (\dots)$			m <sub>H</sub> .	i i
(bb)(qq), (bb, ct, $\tau\tau$ , gg)( $\nu\nu$ )	189	176.2	75 110	[32]
$(any)(e^-e^-, \mu^-\mu^-)$	189	176.2	75 110	32
(35)(= = ) (= = )(ag)	189	176.2	65 110	
(bb)(9 <b>9</b> , <i>PP</i> )	192 202	236.7	60 120	<b>委</b> 日
$(bb, \tau' \tau , c\bar{c}, gg)(c^{+}c^{-}, \mu^{-}\mu^{-})$	192 202	236.7	60 120	j351j
$(b\bar{b}, \tau^+\tau^-), v\bar{v}, gg\}(\tau^+\tau^-), (\tau^+\tau^+)(q\bar{q})$	192 202	236.7	60 120	[33]
(bb)(qq)	190 - 200	217.2	75 120	11.31
$(b\bar{b}, \tau^{+}\tau^{-}, c\bar{c}, gg, WW)(\tau^{-}\tau^{-}, \nu\rho)$	199 - 209	217.2	75 120	11.34
$(bb, \tau^* \tau^*, cb, gg)(c^* c^*, \mu^* \mu^*)$	160 - 200	217.2	70 120	11.34
(bb, cř, si, gg)(qq)	189	176.2	40 100	35
(bb, cf, si, gg)(v9)	189	176.2	60 100	[35]
$(b\bar{b},c\bar{c},s\bar{s},gg)(c^+c^+,\mu^+\mu^+)$	189	176.2	60 115	32.35
(T'T)(qq)	189	176.2	65 110	[32]
(bb, sī, sī, sī, gg)(gē)	192 202	236.7	40 110	[35]
(bb, cf, st, gg)(v9)	192 202	236.7	60 118	[35]
(bb, ct, st, gg)(c * c * , μ* μ* )	192 202	236.7	60 115	33. 35
(⊤` <b>⊤</b> )(qq̃)	192 202	236.7	60 120	[33i]
(bb, ct, st, gg)(gg)	199 200	217.2	40 115	[325]
(bb, cč, sč, gg)(v9)	199 200	217.2	75 120	[35]
$(bb, c\bar{c}, s\bar{s}, g\bar{\mu})(c - c - \mu^{\dagger}\mu^{-})$	199 200	217.2	70 120	[11,34.35]
( <i>τ</i> = ')(qq)	196 - 200	217.2	60 120	11.34

Table 6: Summary of the ALEPH searches for the Higgs trabing process  $e^+e^- \rightarrow H_1Z$ . The top part of the table lists the searches originally developed for the Standard Model Higgs boson. The bottom part lists flavour-independent searches where the decays of the Higgs boson into a quark pair of any flavour, a gluon pair or a tau pair were considered; the signal officiencies were evaluated for all indicated hadronic decays of the Higgs boson. In the cases of the  $(7^{-7})(qq)$ and leatonic changels listed in the flavour-independent part, the ownt selections of the Standard Model Hittes boson searches were used. La Thuile

# Summary of similar LEP searches

$e^{i}e^{-i}\mathcal{H}_{2}Z^{-}(\mathcal{H}_{1}\mathcal{H}_{1})Z^{-}()()$			. m <sub>Ma</sub>	271 14 1	
(any)(a))	91	18.2	12 70	< 0.21	(4 <b>4</b> )
$(\nabla^{0}\nabla^{0})$ iany but $\gamma^{+}\gamma^{+}$ )	91	8.7	0.5 55	< 0.21	(ui -
(17)(am)	91	12.6	0.5 80	< 0.21	[ (ui -
(4 pringe)(any)	91	12.9	UB 80	0.21 13	િલ્હો -
(havirons)(1/2)	91	25.1	1 例	6.21 20	(44
$(\tau - \tau - \tau' + )(\nu P)$	91	25.1	2 78	2.5 12	( <b>4</b> 4)
(may)(a), e0)	181,179	20.0	- 43 - 79	211 - 365	(w  -
(pa)(dulle)	2.63	54.H		12 -61	[ (a)
(5646, 5668, sect)(96)	192-208	484 A	- 30 - 112	12 50	43.64
(coci)/ci)	1955-5004	484.4	- 10 - 110	4 12	[ [47] ] _
(9t, b5.se.ag)(qg)	160 230	625.9	ં સગ્ર ફ્રક્ત	10 42	[56]
(q0cd)(AP)	91	48.3	19 75	0 35	84.65
(0000)(	185	24.1	49 80	10.5 38	_ [`[00]`
(1651-5120-51)	189	172.3	40 109	10.5 48	[ <u>84</u> ]
(6666)((6)	162 208	421.2	80 129	i 12 m <sub>16</sub> /2	a   [10]
(b8b6)(ur)	183	<b>5</b> 610	10 95	:0.5 m <sub>7.9</sub> /	2 [80]
(qqqq)(AP)	189	171.4	50 109	:0.5 m <sub>2m</sub> /	2 [82]
(b6b5)(xP)	199-209	207.2	109 110	) 12 mm//	a   [10] -
(b6b5)(r ' r' )	183	56.7	30 109	:0.5 m <sub>2m</sub> /	/z  ja0j -
(b6b6;(r_r_)	1.69	168.7	30 109	:0.5 m <sub>Pe</sub> /	2 [82]
(b6b5, b5r "r", r"r" r*r")					1
(μ0, e'e , μ μ )	189 238	MN A	45 90	2 10.5	[M]

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## mh-ma limits



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#### **OPAL** low mass searches





#### $H \rightarrow aa \rightarrow 2\mu 2\tau$ at the Tevatron

#### FERMILAB-PUB-09-257-E

Search for NMSSM Higgs bosons in the  $h \rightarrow aa \rightarrow \mu\mu \ \mu\mu, \ \mu\mu \ \tau\tau$  channels using pp collisions at  $\sqrt{s} = 1.96$  TeV



FIG. 3: The expected and observed limits and  $\pm 1$  s.d. and  $\pm 2$  s.d. expected limit bands for  $\sigma(p\overline{p}\rightarrow h+X)\times \text{BR}(h\rightarrow aa)$ , for (a)  $M_h=100$  GeV and (b)  $M_a=4$  GeV. The signal for  $\text{BR}(h\rightarrow aa)=1$  is shown by the solid line. The region  $M_h<86$  GeV is excluded by LEP.

$M_{a}$	$\sigma  imes BR$	Sample	$\sigma  imes 2  imes { m BR}$
$({ m GeV})$	[exp] obs $(fb)$	Data	
0.2143	[10.0] 10.0	$M_a = 3.6 \text{ GeV}$	[23.8] 19.1 fb
0.3	[9.5] 9.5	$M_a=4~{ m GeV}$	[23.9] 45.9 fb
0.5	[7.3] 7.3	$M_a{=}7~{ m GeV}$	[25.0] 24.6 fb
1	[6.1] 6.1	$M_a = 10 \text{ GeV}$	[24.7] 27.3 fb
3	[5.6] 5.6	$M_a {=} 19 \text{ GeV}$	[30.0] 33.7 fb
			_

#### Andy Haas and company collaborated with Wacker and Lisanti to look for these signatures at the Tevatron

Discovering the Higgs with Low Mass Muon Pairs

Mariangela Lisanti and Jay G. Wacker<sup>1</sup> <sup>1</sup> SLAC, Stanford University, Menlo Park, CA 94025 Physics Department, Stanford University, Stanford, CA 94305 (Dated: March 8, 2009)

These searches are probing ~ 10 times the expected production cross-section.

 there are not enough signal events at LEP to compete

However, the  $4\tau$  signature is significantly more difficult at hadron colliders than at LEP

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#### The ALEPH detector: a piece of history



Detector simulation based on Geant 3, analysis based on 10 year old fortran framework

# **Blind Analysis**

Because the LEP data is old and it is not possible to confirm anything with "next year's data", we had to be quite careful

- remember, we're shooting for a discovery!
- no one would believe a signal if we adjusted our cuts looking at data
  - Also, we don't want to spoil the other analyses that we might be interested in:  $a \rightarrow jets, \mu, ..$

But we do need to verify that our Monte Carlo is describing the data well.

- $\cdot$  So we did a blind blind analysis and defined 5 control samples
  - 1. exclude  $m_{ll}$  around  $M_Z$  , that kills our signal, but otherwise similar
  - 2. Select events if #tracks<2 for each jet (kills  $au au, \mu\mu, q\bar{q}, gg$  )
  - 3. in  $Z \rightarrow ll$  exclude events with  $M(j_1, j_2, \text{invisible}) > 60 \text{GeV}$
  - 4. in  $Z \rightarrow \nu \nu$  exclude events with missing mass > 80 GeV
  - 5. exclude events with #track>6 in both jets (to remove taus) AND if di-jet mass > 60 (to avoid seeing  $h \rightarrow aa \rightarrow q\bar{q}, gg$  if it exists)

## Blind Analysis: control samples



## **Systematics**

- Given the low numbers of selected events, the final measurements are statistically limited
- $\blacksquare Z \longrightarrow \mu^+ \mu^-, e^+ e^-$ 
  - total systematic uncertainties (from lepton Id) are 0.6%,
     2.6% and 7.5% for the signal, ZZ, and Zee backgrounds.
  - Overall 10% background uncertainty for the other channels
- $\blacksquare \ \mathsf{Z} \to \mathsf{v}\mathsf{v}$ 
  - Efficiency error was 5% for the signal and 10% for ZZ, and it is between 30% and 100% for the other background processes.
  - Overall 30% background uncertainty for the other channels
- With Loose cuts data and MC agree within these errors

# M(μμ): Data vs Montecarlo



#### Recoil Mass: Data vs Montecarlo



#### Loose selection: $H \rightarrow 4\tau$ , $Z \rightarrow e^+e^-$



#### Final selection: $H \rightarrow 4\tau$ , $Z \rightarrow e^+e^-$

