

# Testing scalar-tensor theory from observations of gravitational wave bursts with a network of interferometric detectors

Project page :

<http://www.aei.mpg.de/~kahaya/RIDGE-SCALAR/index.html>

Kazuhiro Hayama (NAOJ, AEI Hannover)

Atsushi Nishizawa (NAOJ)

Atsushi Taruya (U. Tokyo)

**LIGO Document G0901084**



# Motivation

- Testing relativistic gravity theory is important for fundamental physics and cosmology e.g. dark matter, dark energy accelerating the Universe.
- One of plausible gravity theories is scalar-tensor theory. Significant difference from the general relativity is the existence of a scalar field which is connected with the gravity field with coupling parameters, and a resulting scalar gravitational wave.
- In general, Scalar-tensor theories has a term of matter in the action in which physical metric:

$$\tilde{g}_{\mu\nu} := A^2(\phi)g_{\mu\nu}^*$$

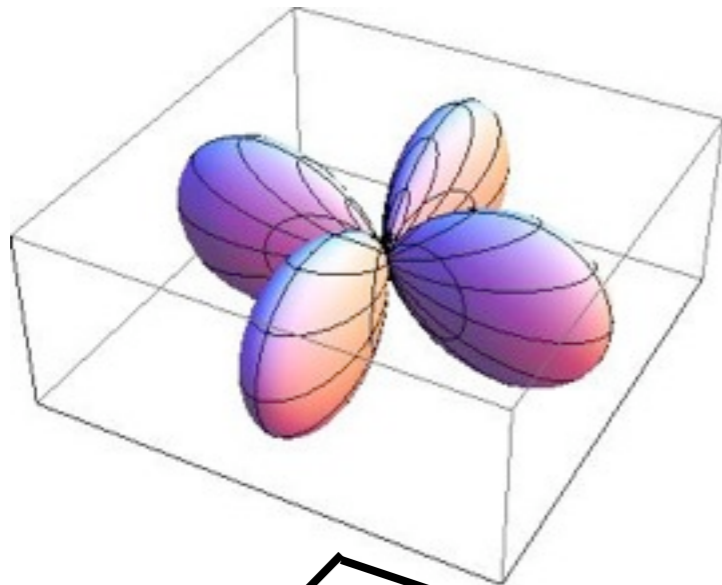
$$A(\phi) = e^{\alpha_0(\phi-\phi_0) + \frac{1}{2}\beta_0(\phi-\phi_0)^2}$$

$\alpha_0$  and  $\beta_0$  are parameters which characterize

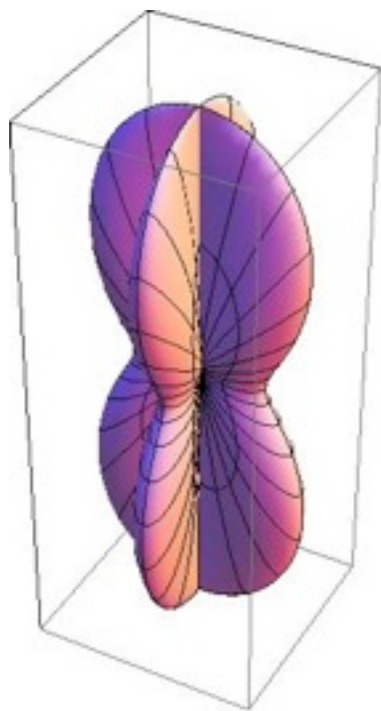
- $\alpha_0$  : weak gravity field (test in solar system,  $\alpha < 3.5 \times 10^{-3}$  by Cassini)
- $\beta_0$  : strong gravity field (can constrain by GW observation ( $\beta = 0$  in Brans–Dicke))
- Tensor GW search might miss some type of source, e.g. highly spherical core collapse.. In this sense, search for SGW is complementary to current GWsearch.
- we present how to extract a scalar gravitational wave signal using a network of world wide interferometric detectors.

# Antenna pattern to polarization

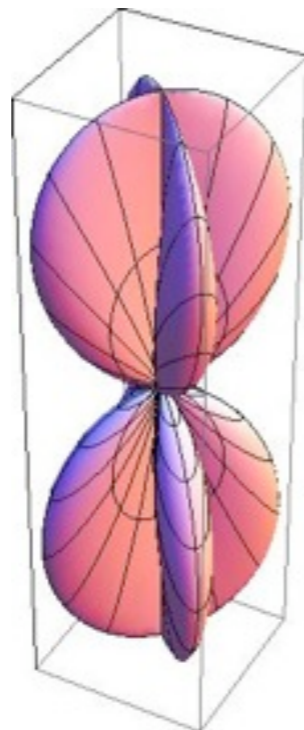
$$F_o(\vartheta, \varphi)$$



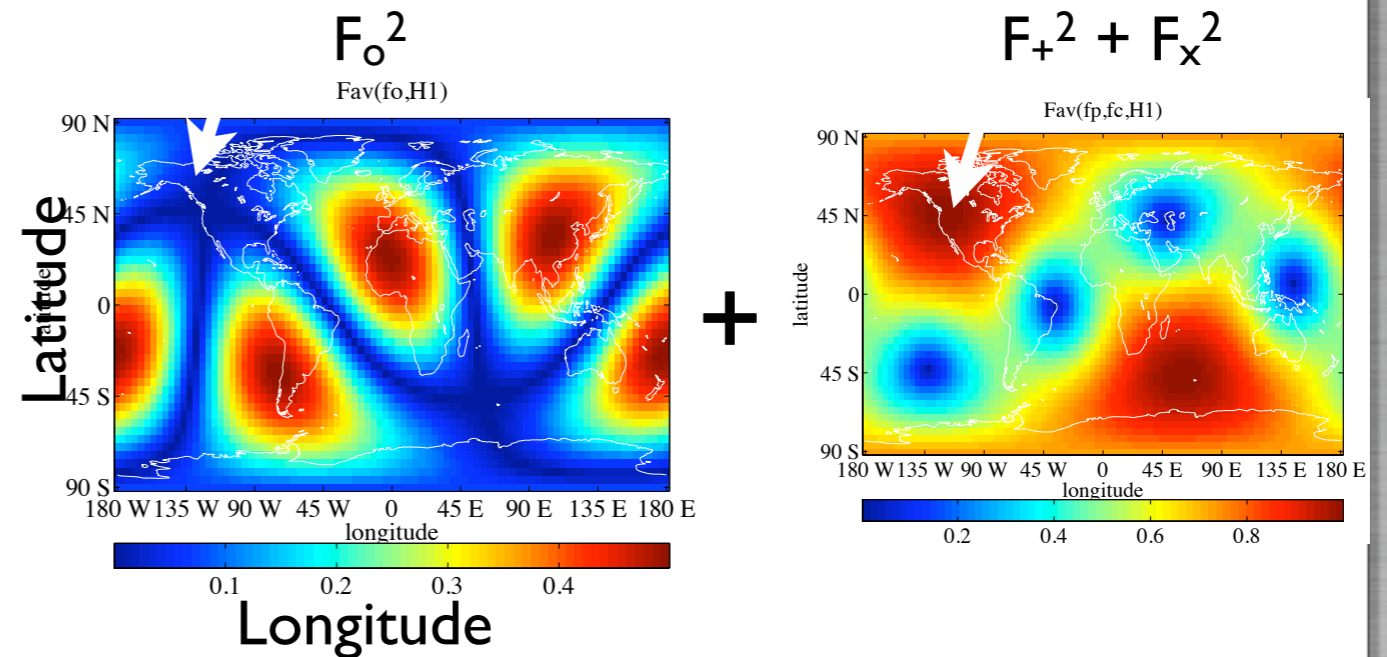
$$F_+(\vartheta, \varphi)$$



$$F_x(\vartheta, \varphi)$$



## Antenna pattern sky-map (LIGO Hanford)



- SGWs have Maximum sensitivity to the direction in which TGWs don't have sensitivity.  
----> complementary to TGW search

# Search for scalar gravitational waves

- Coherent network analysis can extract scalar gravitational wave with more than 3 world-wide detectors.
- This approach combines data taking account of the sky position  $(\vartheta, \varphi)$ , arrival time difference  $\tau(\vartheta, \varphi)$  coherently, and calculates all polarization components on a sky position which has largest likelihood statistic.

## Mathematical expression of the coherent network analysis

Expression of d-detectors can be taken as

$$\begin{bmatrix} x_1 \\ \vdots \\ x_d \end{bmatrix} = \begin{bmatrix} F_{1+} & F_{1\times} & F_{1\circ} \\ \vdots & \vdots & \vdots \\ F_{d+} & F_{d\times} & F_{d\circ} \end{bmatrix} \begin{bmatrix} h_+ \\ h_\times \\ h_\circ \end{bmatrix} + \begin{bmatrix} n_1 \\ \vdots \\ n_d \end{bmatrix}$$

The reconstruction of a gravitational wave is an inverse problem.  
Maximum likelihood method to solve the inverse problem:

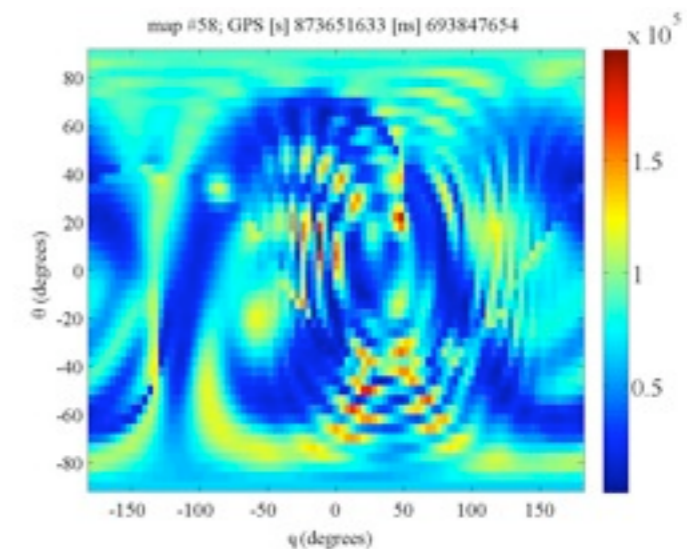
$$L[\mathbf{h}] := \|\mathbf{x} - \mathbf{F}\mathbf{h}\|^2$$

Changing sky position  $(\vartheta, \varphi)$ , time difference  $\tau(\vartheta, \varphi)$ .

The mathematical formula of the reconstructed scalar gravitational wave is

$$\begin{aligned} h_\circ &= \frac{1}{\det(\mathbf{M})} \left( ((F_+ \times F_\times) \cdot (F_\times \times F_\circ)) \cdot F_+ \right. \\ &\quad - ((F_+ \times F_\times) \cdot (F_+ \times F_\circ)) \cdot F_\times \\ &\quad \left. + ((F_+ \times F_\times) \cdot (F_+ \times F_\times)) \cdot F_\circ \right) \cdot \mathbf{x} \end{aligned}$$

Likelihood residual sky-map

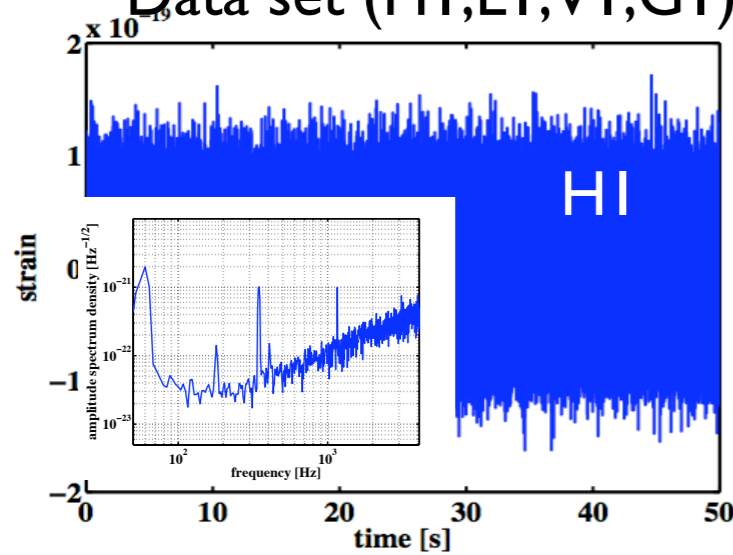




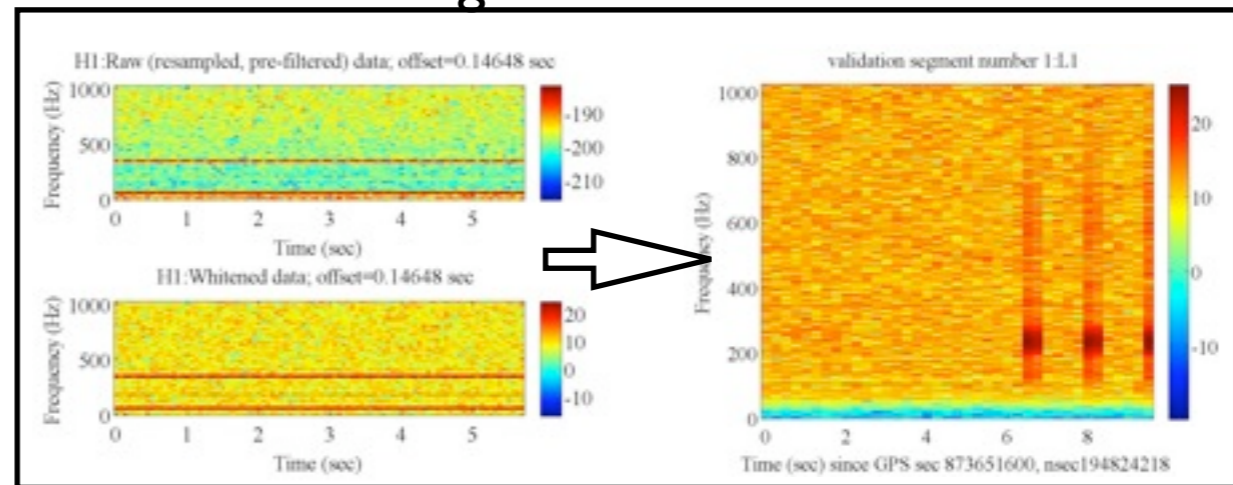
# Scalar pipeline

- Full featured coherent network analysis pipeline (Data conditioning, detection stat., Veto analysis)

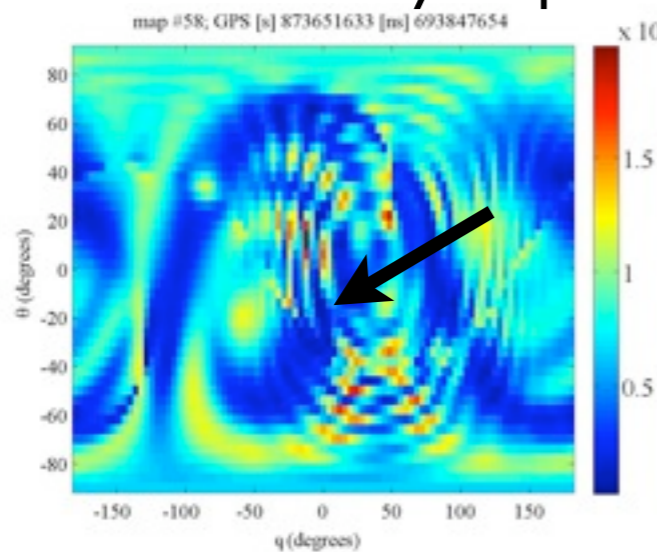
Data set (HI,LI,VI,GI)



Data conditioning

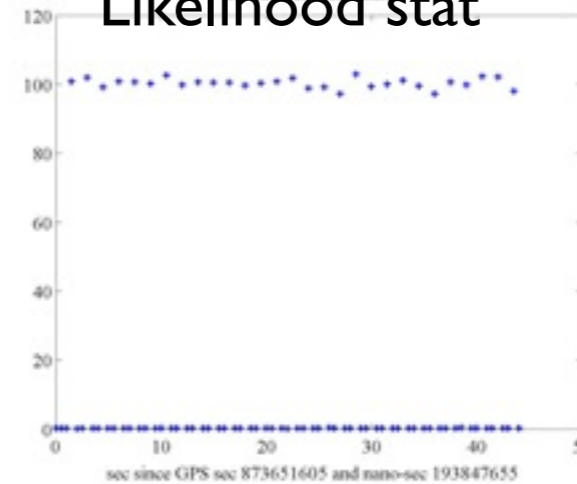


Likelihood sky-map



Coherent network analysis

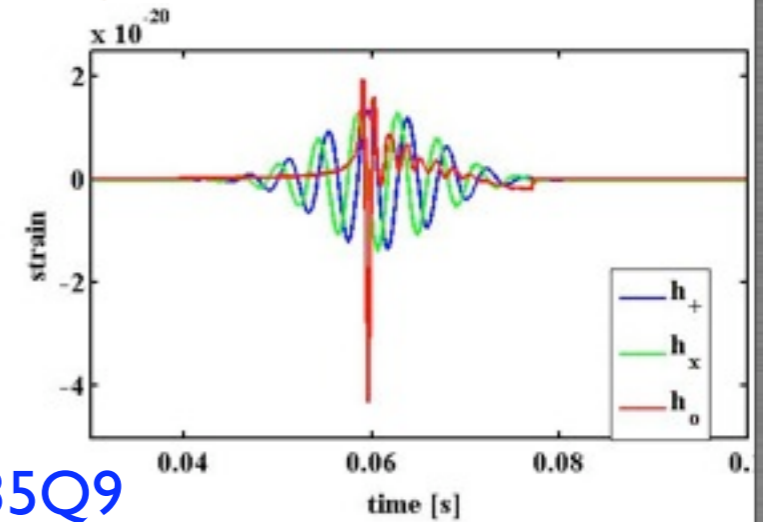
Likelihood stat



# Reconstruction of pol. modes

## Reconstruction of $h_+$ , $h_x$ , $h_o$

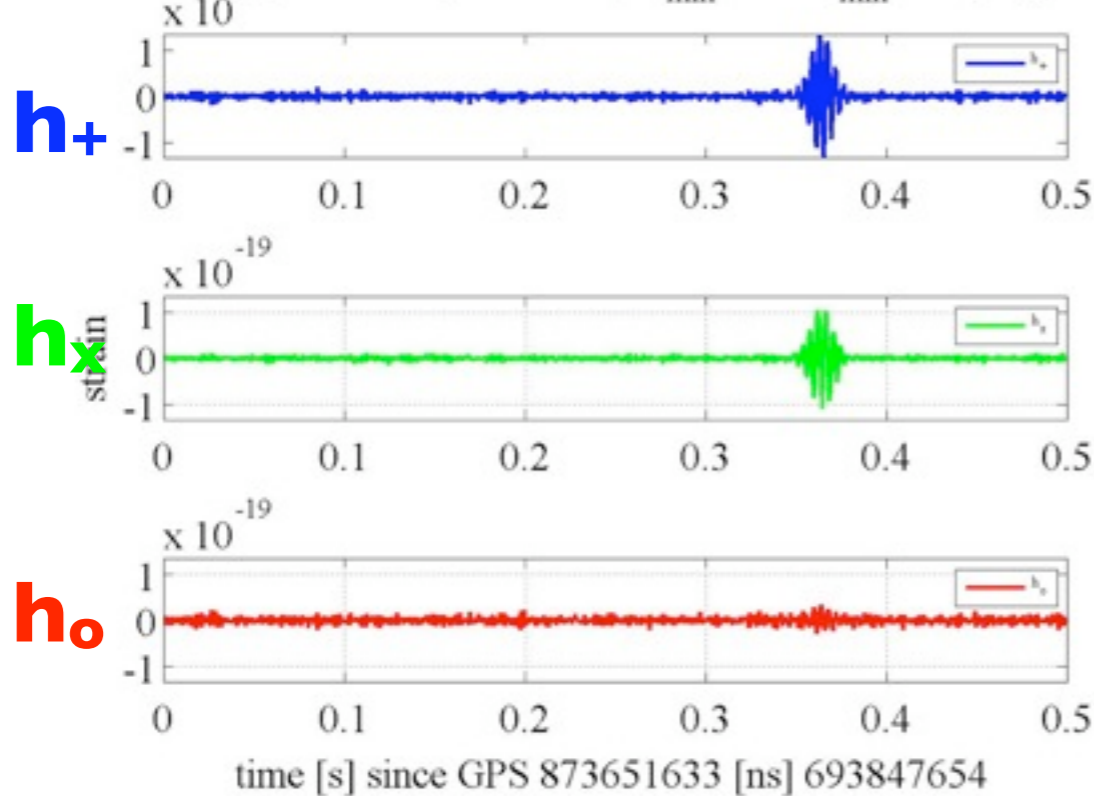
- Although the grid of lat-lon map is coarse ( $4^\circ \times 4^\circ$ ) in the simulation,  $h_o$  is reconstructed clearly.



$h_+$ ,  $h_x$ : SG235Q9

$h_o$ : Not injected

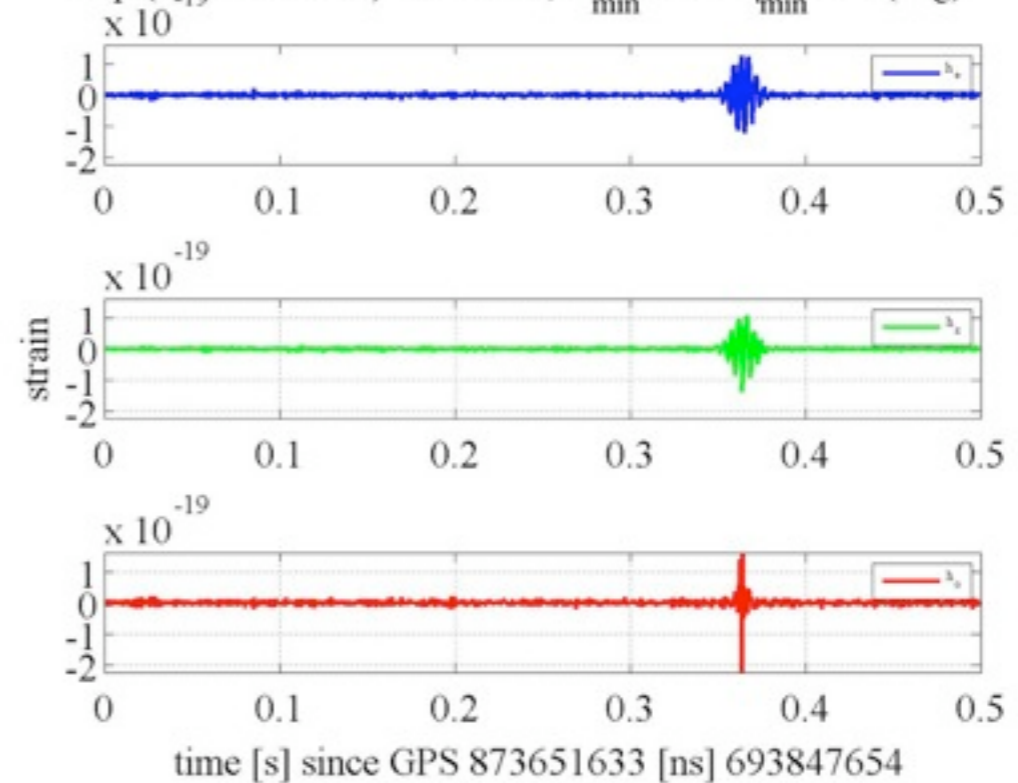
map (max radialdist) serial #58;  $\theta_{\min} = -10$   $\varphi_{\min} = -4$ (deg)



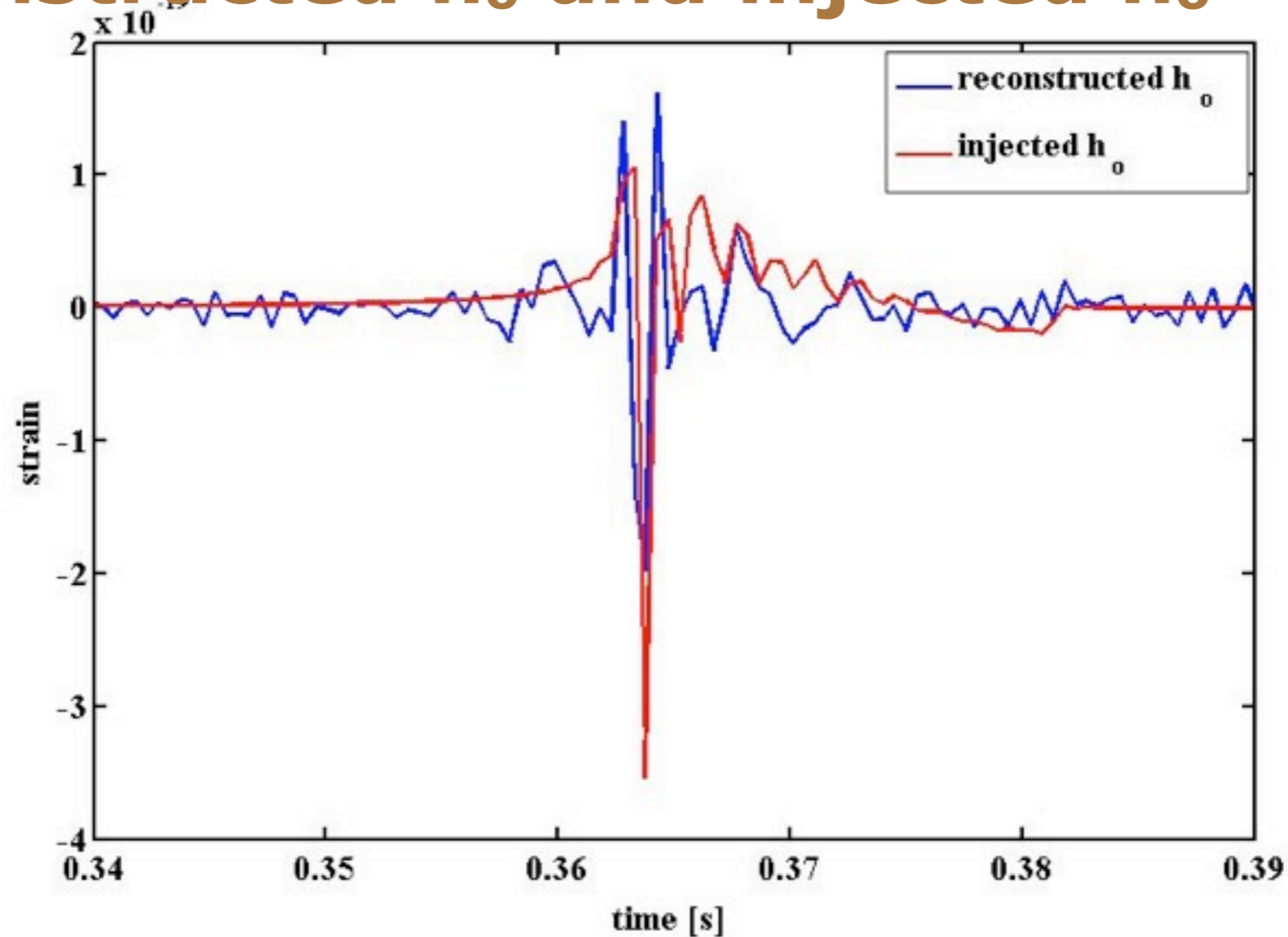
$h_+$ ,  $h_x$ : SG235Q9

$h_o$ : Spike-like burst

map (max radialdist) serial #58;  $\theta_{\min} = -10$   $\varphi_{\min} = -4$ (deg)



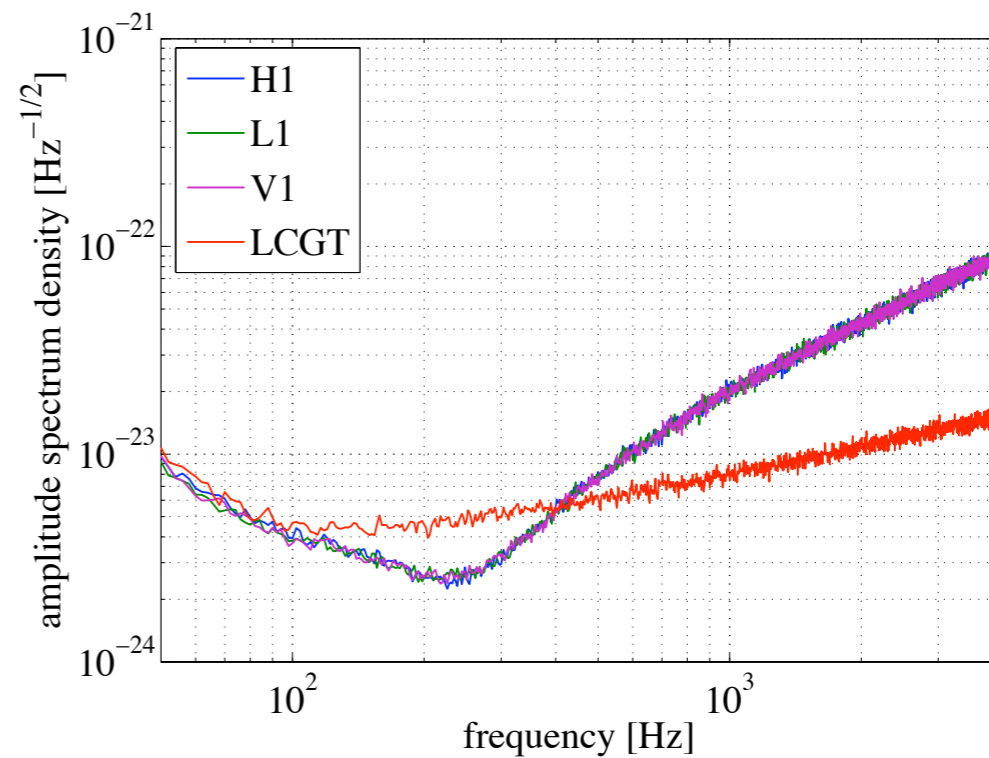
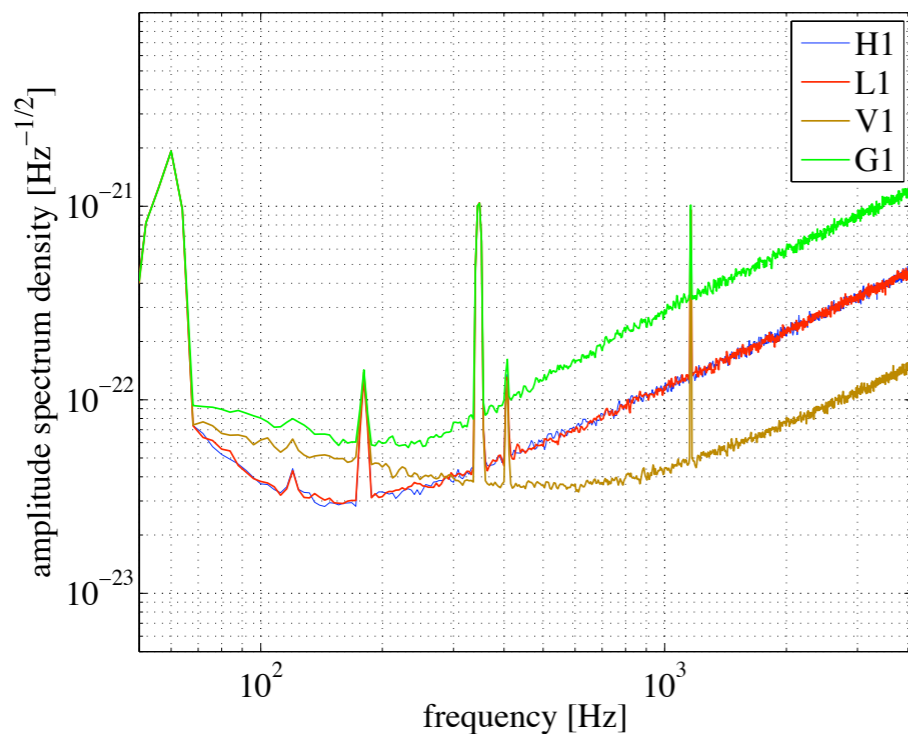
# Comparison of reconstructed $h_0$ and injected $h_0$



Red plot is injected  $h_0$  signal and blue plot is the reconstructed  $h_0$ .  
The difference at the low frequency region comes from the data conditioning step.  
Detector noise at low frequency is very high, such region is cut at the step.

# Reconstruction simulation

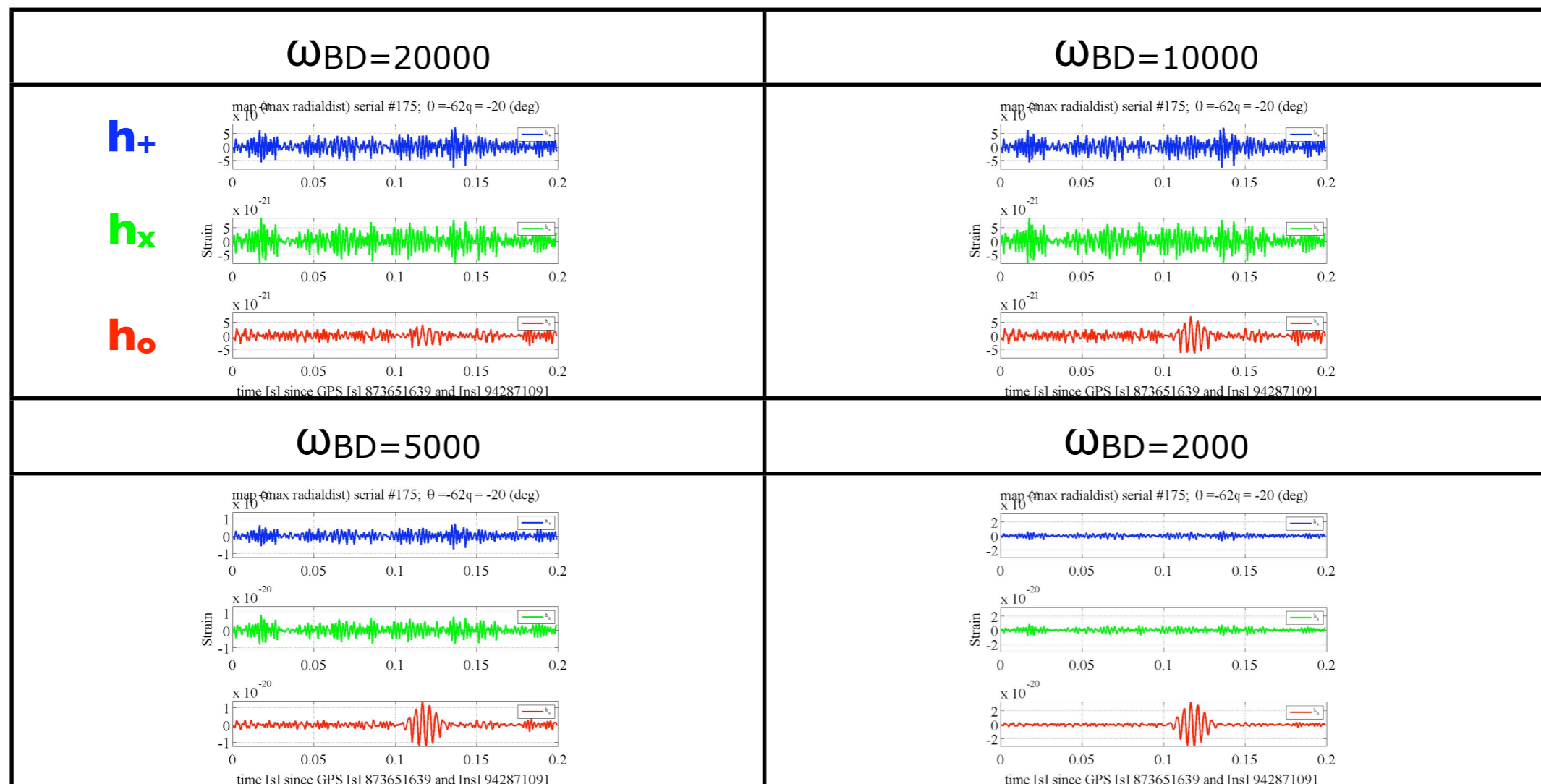
- Simulations were performed using simulated noise of current detector network (4km LIGO Hanford, Livingston, VIRGO, GEO600) and next generation (4km advLIGO Hanford, Livingston, advVIRGO, LCGT) For advVIRGO, the design sensitivity of advLIGO is used.
- We consider scalar GW in Brans-Dicke case ( $\alpha_0$  and  $\beta_0$ )  
(  $\alpha_0^2 = \frac{1}{3 + 2\omega_{BD}}$  ,  $\beta_0=0$  )





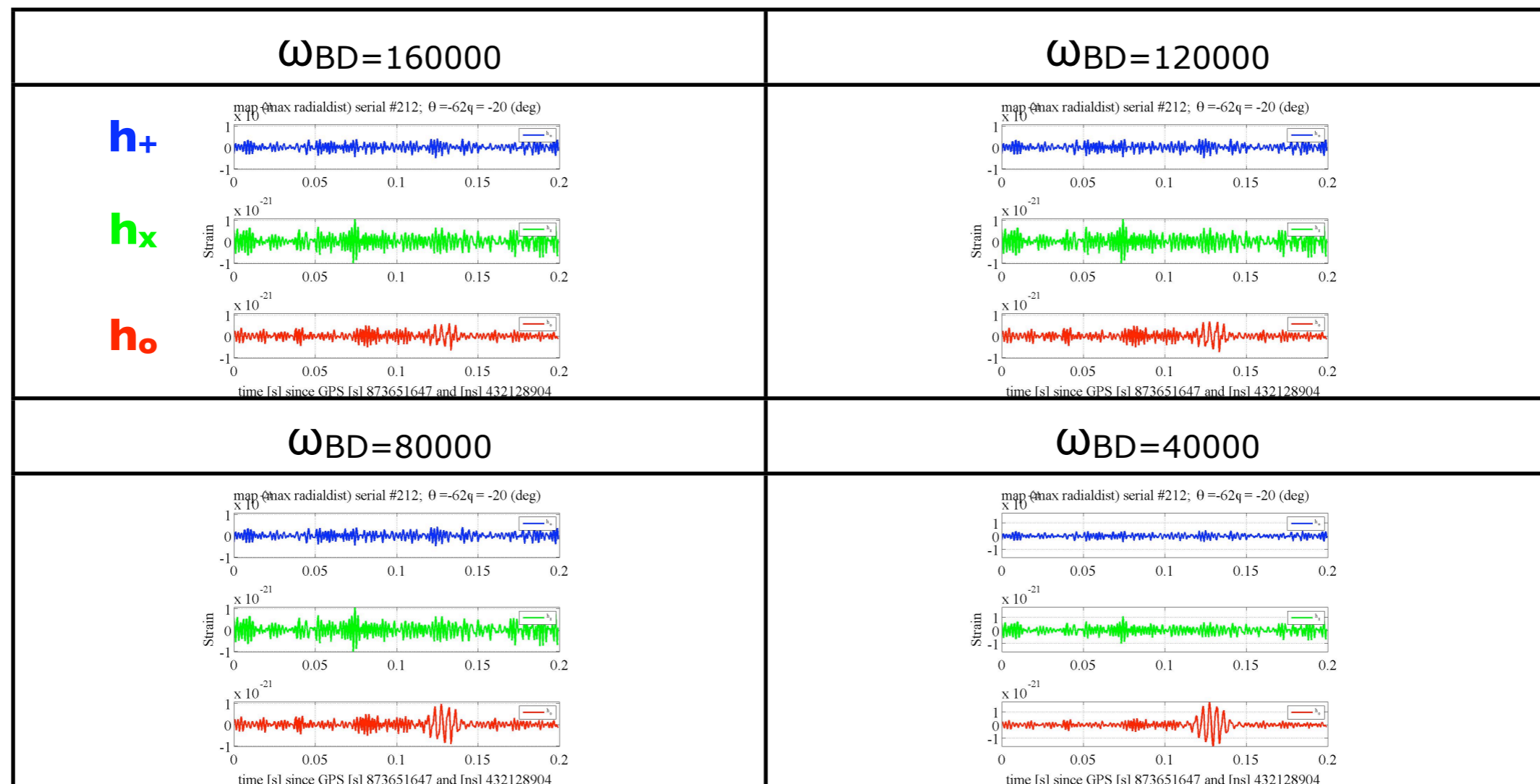
# Reconstruction of SGW (current detectors)

We performed simulations to reconstruct scalar gravitational waves with  $\omega_{BD} = 2000, 5000, 10000, 20000$ . Astrophysical model used is a spherically symmetric core collapse with  $10M_{\odot}$  at the distance of  $10\text{kpc}$  from the earth. Sine Gaussian with center frequency of  $235\text{Hz}$  and  $Q$  value of  $9$  is used as scalar gravitational wave. The maximum amplitude of the signal is set to  $3 \times 10^{-20} \times 500 / \omega_{BD}$  from Shibata et al(1994). From the result, we found the signal with  $\omega_{BD} \leq 10000$  can be detected and reconstructed clearly.



# Reconstruction of SGW(advanced detectors)

We performed simulations to reconstruct scalar gravitational waves with  $\omega_{BD} = 40000, 80000, 120000, 160000$ . This simulation uses the design sensitivity of advLIGO for LIGO, VIRGO, and the one of LCGT. Astrophysical model used is the same as the previous simulation.



# Detection statistics

We perform an excess power method to calculate detection statistic.

## Excess power method

- $h_0$  is divided into segments (each segment is overlapped 50%)
- In each segment, events are extracted in wavelet space by hard thresholding.

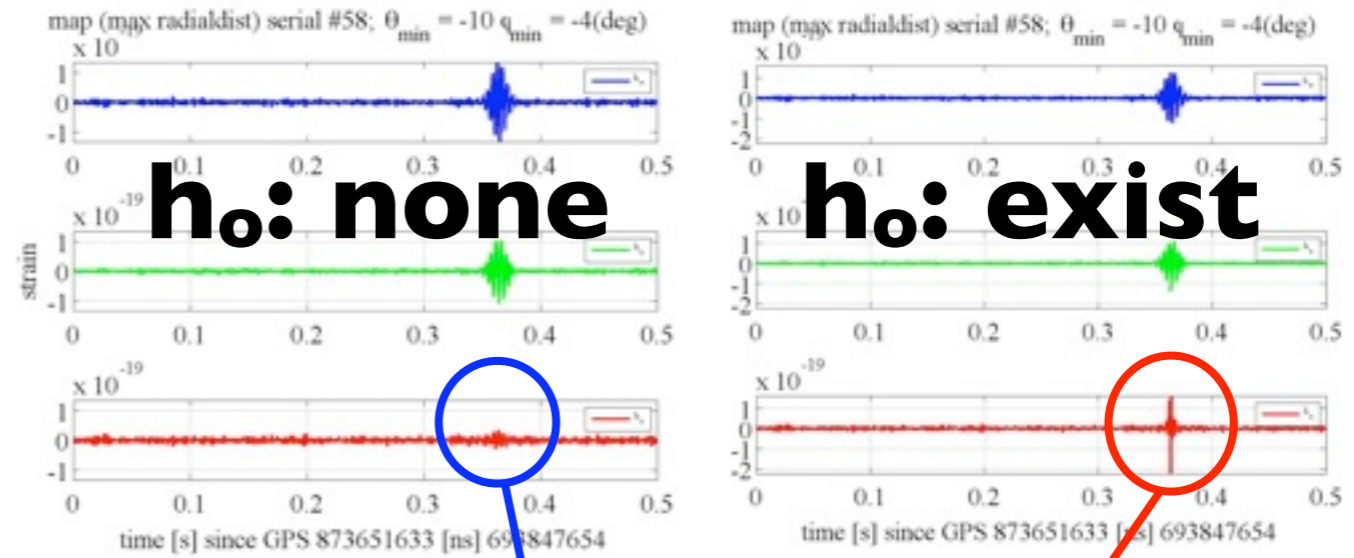
$$\hat{w}_{i,j} = w_{i,j} \text{ if } \text{abs}(w_{i,j}) > T$$

$$0 \text{ if } \text{abs}(w_{i,j}) < T$$

- Calculate statistic defined as

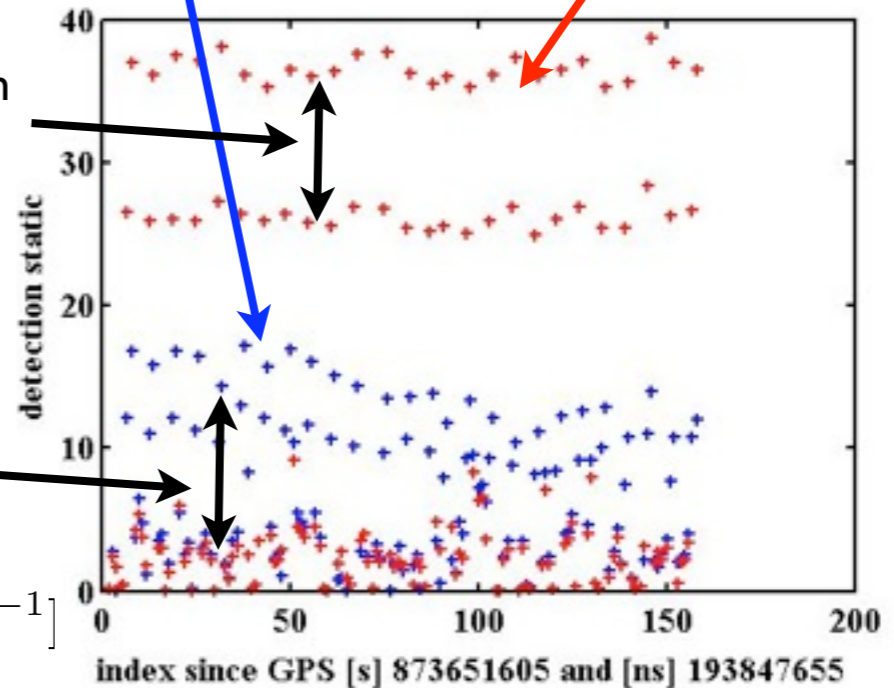
$$\text{statistic} = \sqrt{\frac{\sum \hat{w}_{i,j}^2}{\sigma^2}}$$

$$E\{\| h - E\{h\} \|^2\} = \sigma^2 \text{Tr}[(A^T A)^{-1}]$$



This gap comes from overlap of segments

Error of waveform estimation

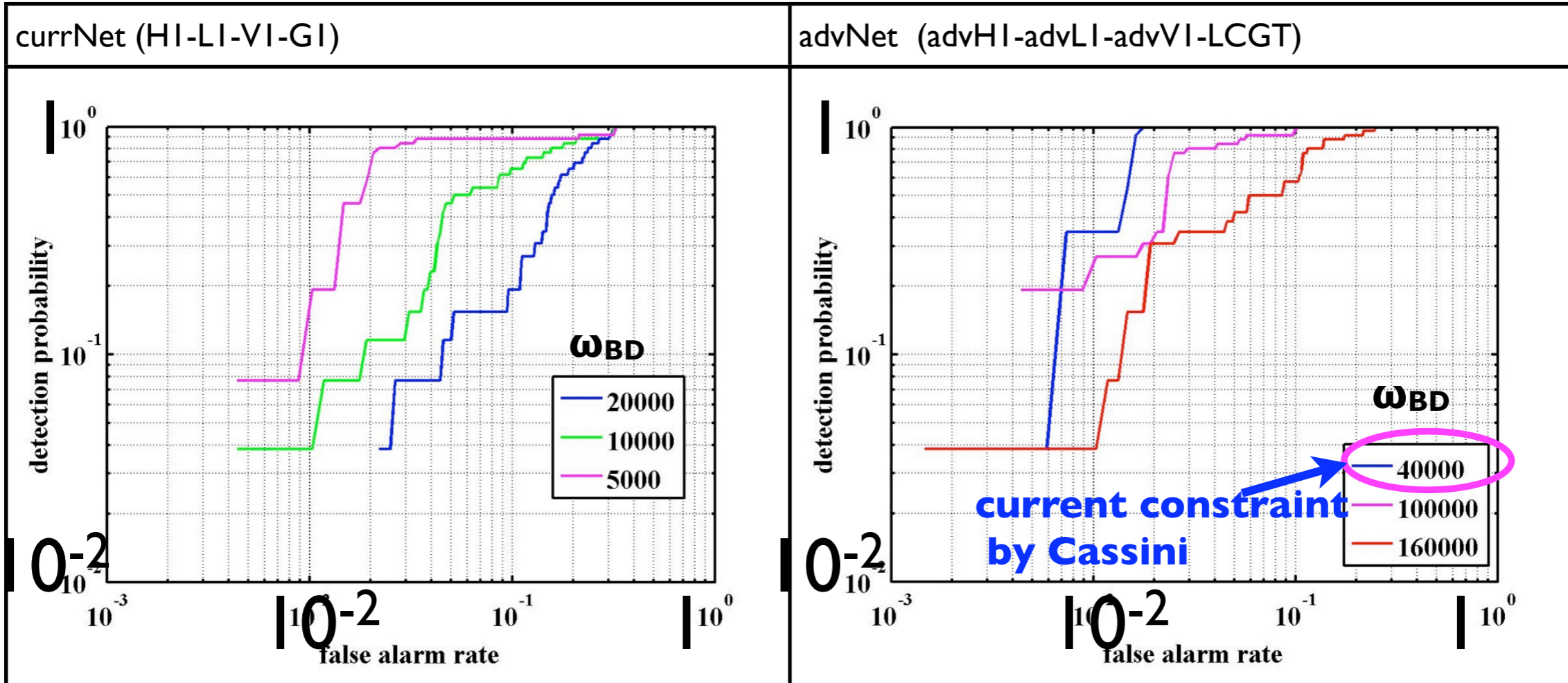




# Detection Efficiency

- For  $\omega_{BD}=40000$ (Cassini), 90% det. prob. at false alarm prob.=0.18 using advNet
- For  $\omega_{BD}=5000$ , 80% det. prob. at false alarm prob.=0.2 using currNet

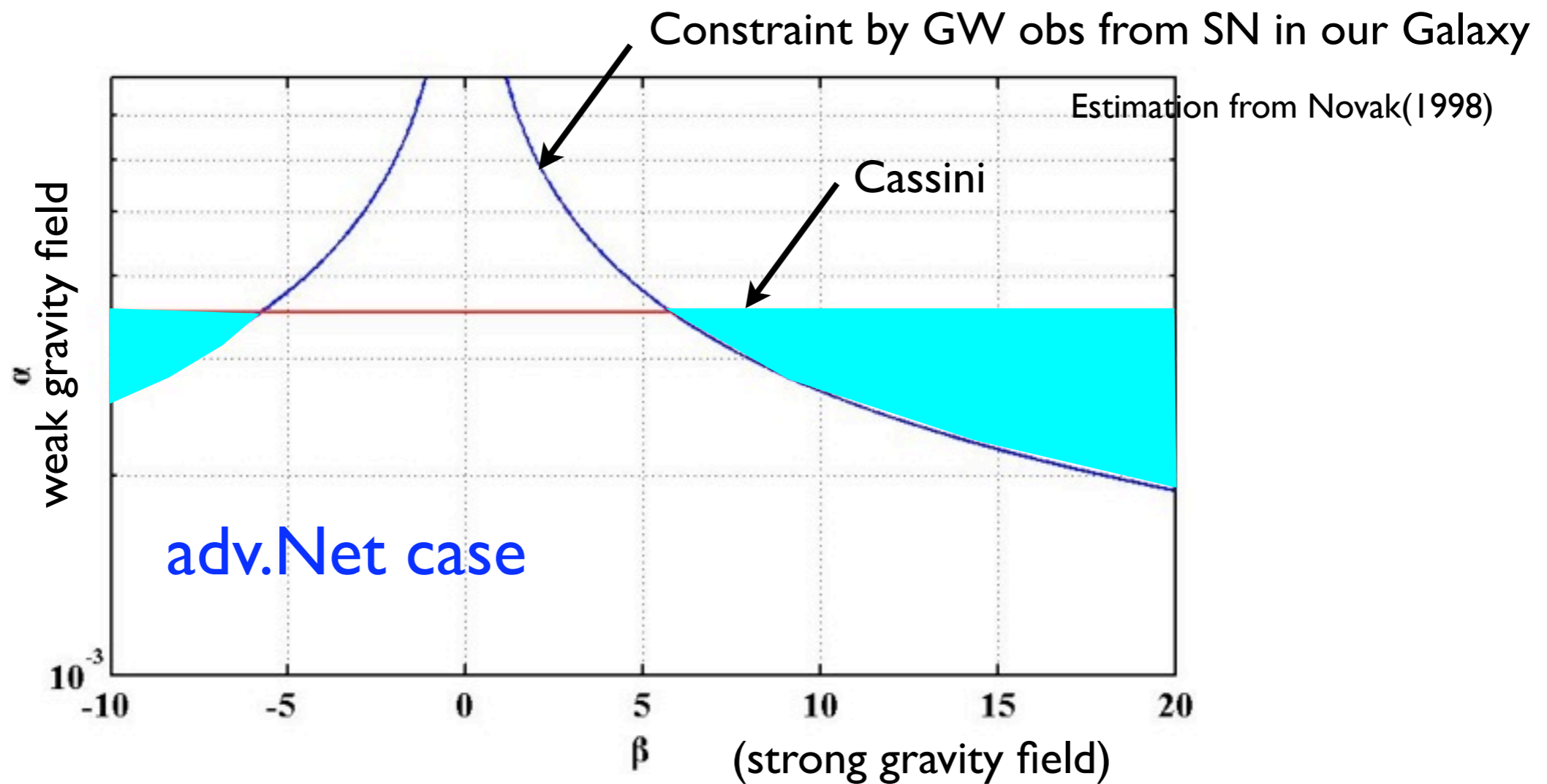
## ROC curve





# $\alpha$ - $\beta$ map

General Scalar-tensor theory which is characterized by  $\alpha$  and  $\beta$ ...



$\beta$  is constrained by the use of the constraint of  $\alpha$  by Cassini.

# Summary

- We discussed test of scalar-tensor theory from gravitational wave observations.
- We implemented a coherent network analysis pipeline for the detection of a scalar gravitational wave with simulated data from a network of interferometric detectors.
- $\alpha$   
Next generation of detectors (advLIGO,advVIRGO,LCGT) can put constraint much stronger than Cassini when a spherically symmetric core collapse occurs in our Galaxy.
- $\beta$   
Obs. in Solar system cannot constrain it. GW observation is a good tool.

# Search for scalar gravitational waves in Brans-Dicke Theory

- Main purpose of the search is the detection of a scalar gravitational wave. Particularly, in case of a spherically symmetric core collapse, a tensor gravitational wave cannot emit, only the scalar gravitational wave can emit.
- Even if the scalar gravitational wave is not detected, a constraint of  $\omega_{BD}$  is possible.
- Current constraint  $\omega_{BD}$  is  $\omega_{BD} > 4 \times 10^4$  from Cassini (Nature, 2003, astro-ph0709.0082)

How an interferometric detector can detect a scalar gravitational wave?

We assume the scalar gravitational wave  $h_0$  comes from z-axis. The geodesic equation of a mirror is

$$\frac{d}{d\tau} \left[ (1 + h_0) \frac{dx}{d\tau} \right] = \frac{d}{d\tau} \left[ (1 + h_0) \frac{dy}{d\tau} \right] = 0$$

$$\frac{d}{d\tau} \left[ (1 + h_0) \frac{dt}{d\tau} \right] = -\frac{1}{2} \frac{\partial_t(1 + h_0)}{1 + h_0}$$

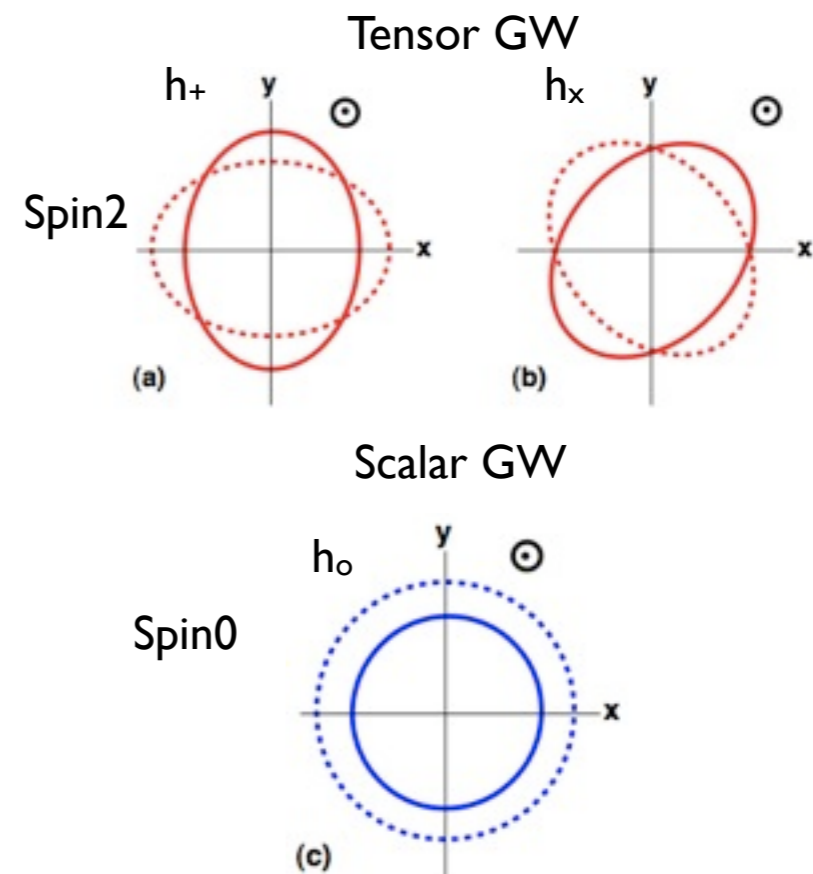
$$\frac{d}{d\tau} \left[ (1 + h_0) \frac{dz}{d\tau} \right] = -\frac{1}{2} \frac{\partial_z(1 + h_0)}{1 + h_0}$$

Introducing  $u = t - z$ , performing integration of  $z$ , we obtain

$$\Delta z \simeq \frac{1}{2} \int_{-\infty}^{t-z_0} h_0(u) du$$

This shows the scalar gravitational wave moves a mirror in z-axis.

Polarization of tensor, scalar gravitational wave



C. Will, Living Review (2006)