

# Direct Search for the Neutrino Mass and the KATRIN experiment

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**Introduction**

**Direct Neutrino Mass determination**

**The Karlsruhe Tritium Neutrino experiment KATRIN**

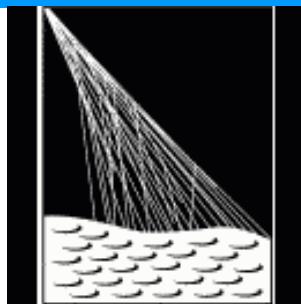
**Options for sterile neutrinos with KATRIN**

**Outlook on other approaches: ECHo, HOLMes and Project 8**

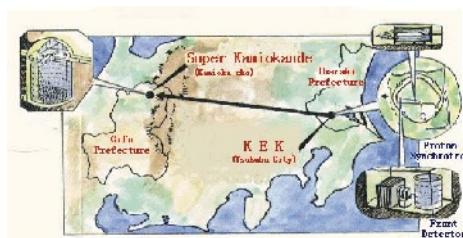
**Conclusions**

# Positive results from $\nu$ oscillation experiments

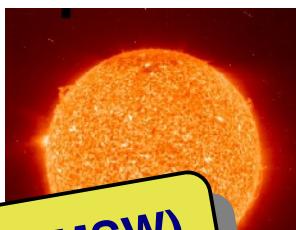
atmospheric neutrinos  
(Kamiokande,  
Super-Kamiokande, ...)



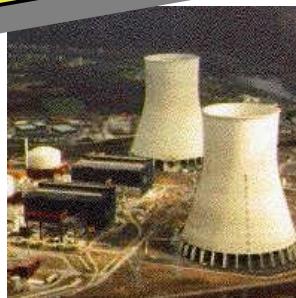
accelerator neutrinos  
(K2K, T2K, MINOS,  
OPERA, MiniBoone)



solar neutrinos  
(Homestake, Gallex,  
Sage, Super-Kamiokande,  
SNO, Borexino)



reactor neutrinos  
(KamLAND, CHOOZ, Daya Bay,  
DoubleCHOOZ, RENO, ...)



⇒ non-trivial  $\nu$ -mixing

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \cdot \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

with:

0.37 <  $\sin^2(\theta_{23})$  < 0.63 maximal!

0.26 <  $\sin^2(\theta_{12})$  < 0.36 large !

0.018 <  $\sin^2(\theta_{13})$  < 0.030  $8.9^\circ$

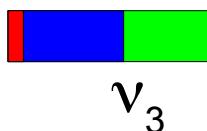
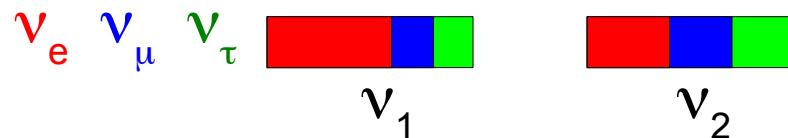
$7.0 \cdot 10^{-5} \text{ eV}^2 < \Delta m_{12}^2 < 8.2 \cdot 10^{-5} \text{ eV}^2$

$2.2 \cdot 10^{-3} \text{ eV}^2 < |\Delta m_{13}^2| < 2.6 \cdot 10^{-3} \text{ eV}^2$

⇒  $m(\nu_j) \neq 0$ , but unknown !

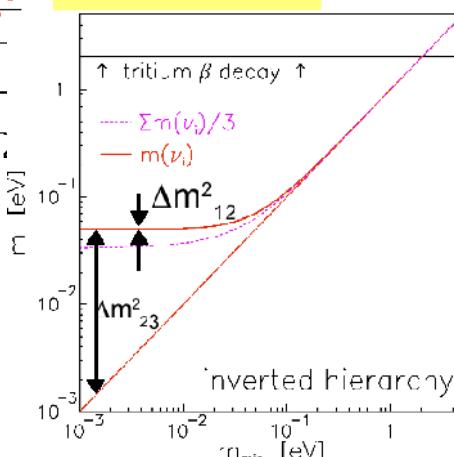
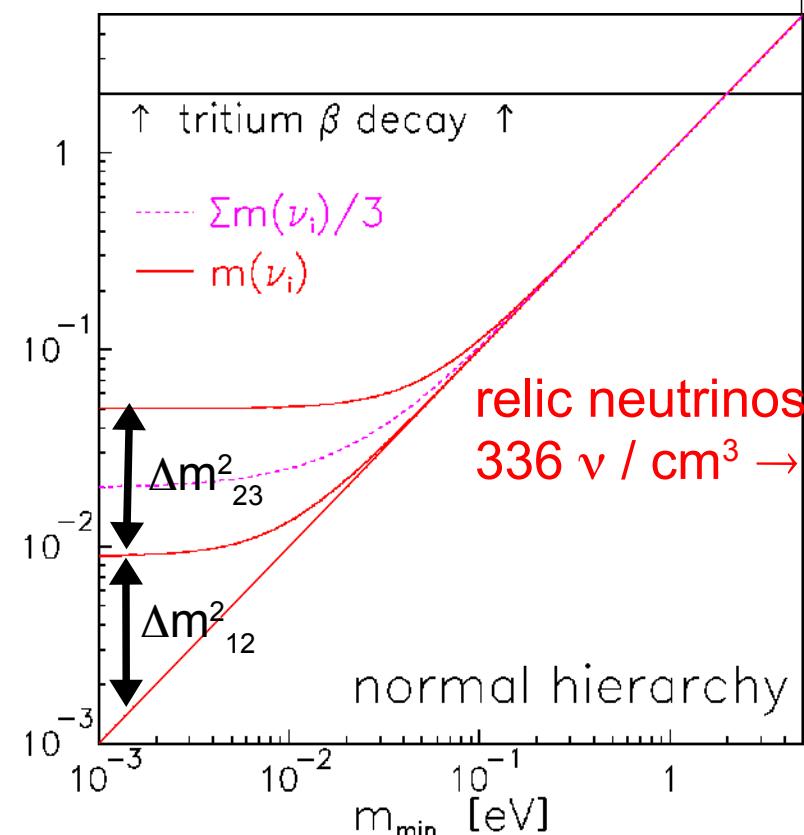
# Need for the absolute $\nu$ mass determination

Results of recent oscillation experiments:  $\Theta_{23}$ ,  $\Theta_{12}$ ,  $\Theta_{13}$ ,  $\Delta m^2_{23}$ ,  $\Delta m^2_{12}$



degenerated masses  
cosmological relevant  
e.g. seesaw mechanism type 2

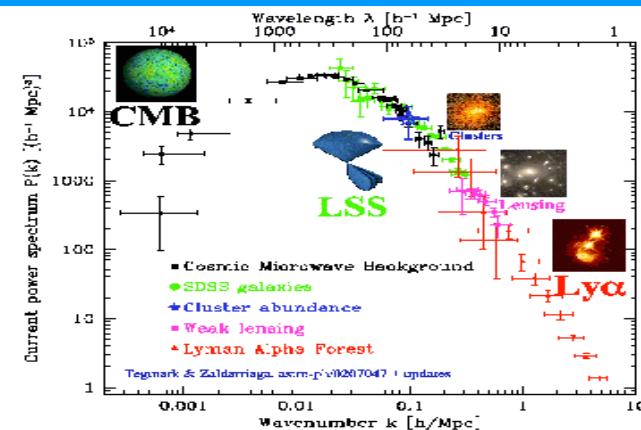
hierarchical masses  
e.g. seesaw mechanism type 1  
explains smallness of masses,  
but not large (maximal) mixing



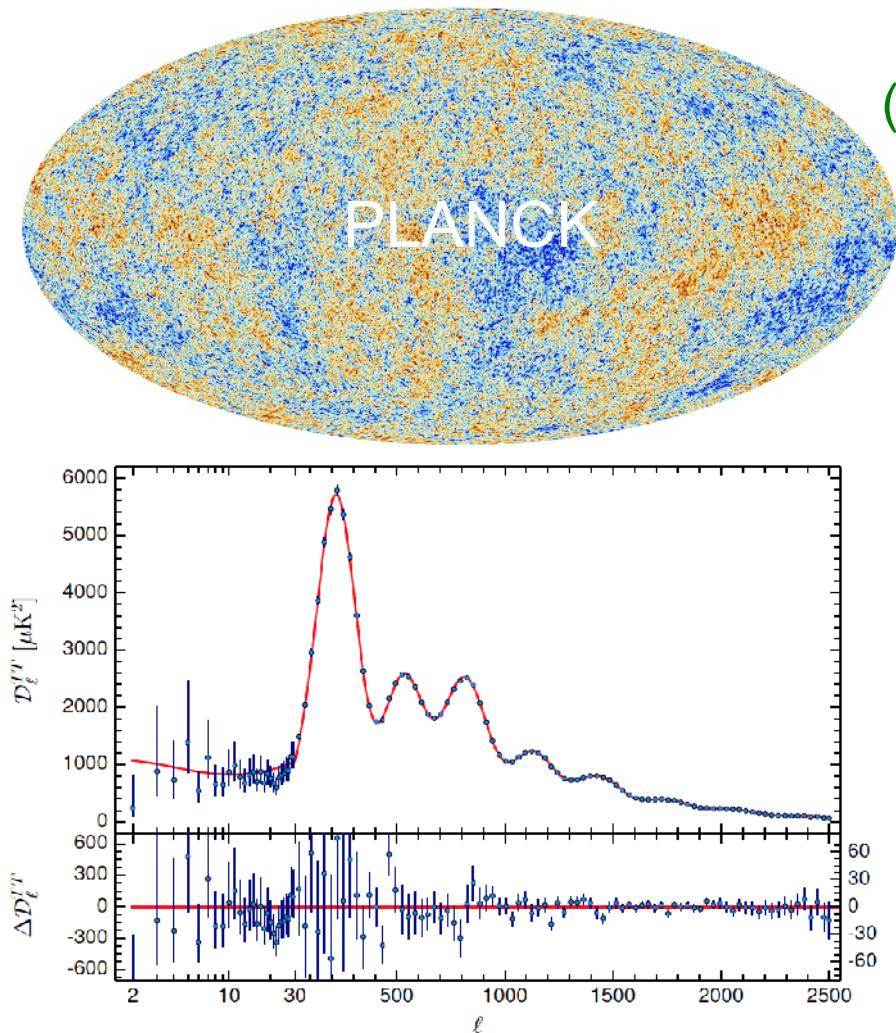
# Three complementary ways to the absolute neutrino mass scale

## 1) Cosmology

very sensitive, but model dependent  
compares power at different scales  
current sensitivity:  $\sum m(\nu_i) \approx 0.23 \text{ eV}$



# Neutrino mass from cosmology



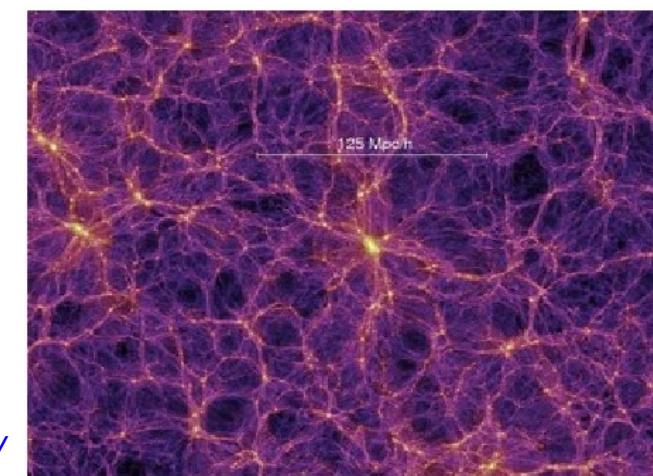
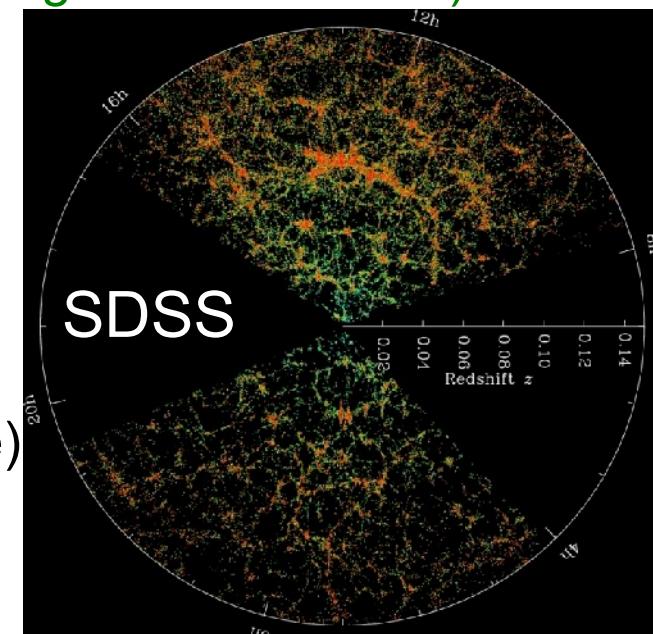
Planck Collaboration:  
P. A. R. Ade et al., arXiv:1502.01589

measurement of CMBR  
(Cosmic Microwave Background Radiation)

measurement of  
matter density  
distribution LSS  
(Large Scale Structure)  
by 2dF, SDSS, ...

compare to  
numeric. models  
including relic  
neutrino density  
of  $336 \text{ cm}^{-3}$

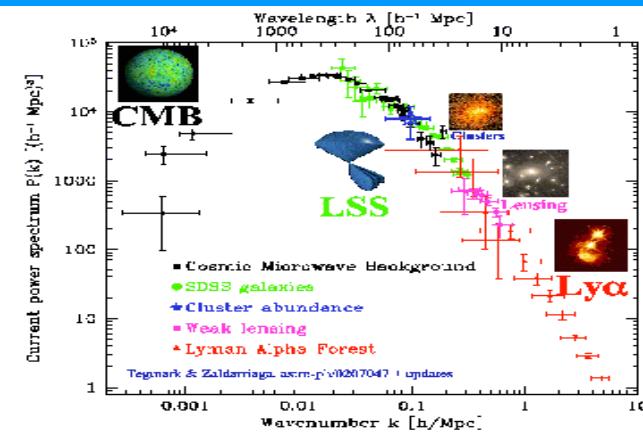
Millenium simulation →  
<http://www.mpa-garching.mpg.de/galform/presse/>



# Three complementary ways to the absolute neutrino mass scale

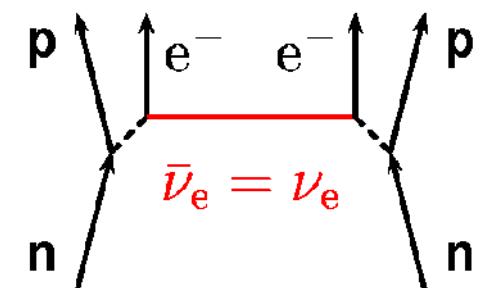
## 1) Cosmology

very sensitive, but model dependent  
compares power at different scales  
current sensitivity:  $\sum m(\nu_i) \approx 0.5$  eV



## 2) Search for $0\nu\beta\beta$

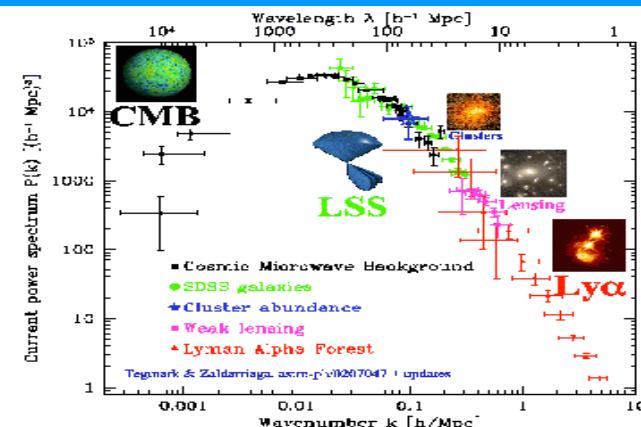
Sensitive to Majorana neutrinos  
First upper limits by EXO-200, KamLAND-Zen, GERDA



# Three complementary ways to the absolute neutrino mass scale

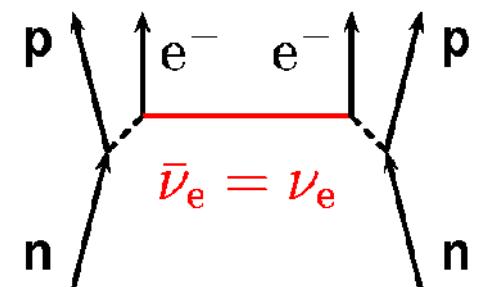
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Sensitive to Majorana neutrinos  
 First upper limits by EXO-200, KamLAND-Zen, GERDA



## 3) Direct neutrino mass determination:

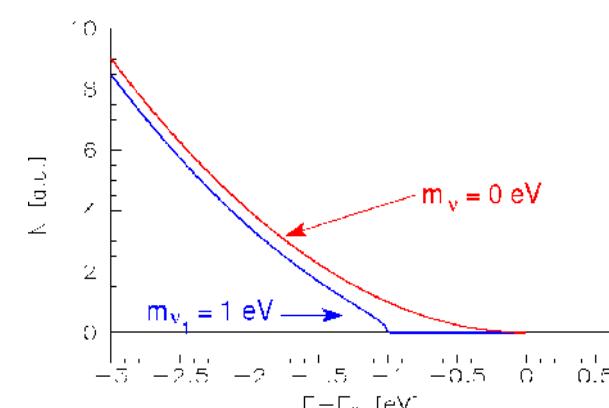
No further assumptions needed, use  $E^2 = p^2c^2 + m^2c^4 \Rightarrow m^2(\nu)$  is observable mostly  
**Time-of-flight measurements** ( $\nu$  from supernova)

SN1987a (large Magellan cloud)  $\Rightarrow m(\nu_e) < 5.7$  eV

### Kinematics of weak decays / beta decays

measure charged decay prod., E-, p-conservation

- $\beta$ -decay searches for  $m(\nu_e)$ 
  - tritium  $\beta^-$  spectrometers
  - $^{187}\text{Re}$ ,  $^{163}\text{Ho}$  bolometers



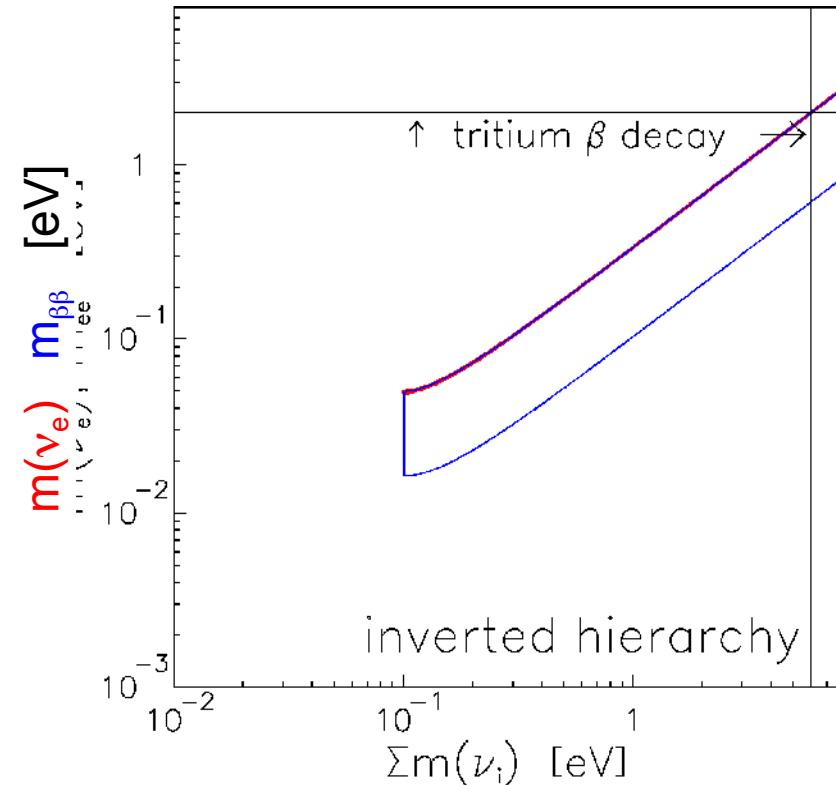
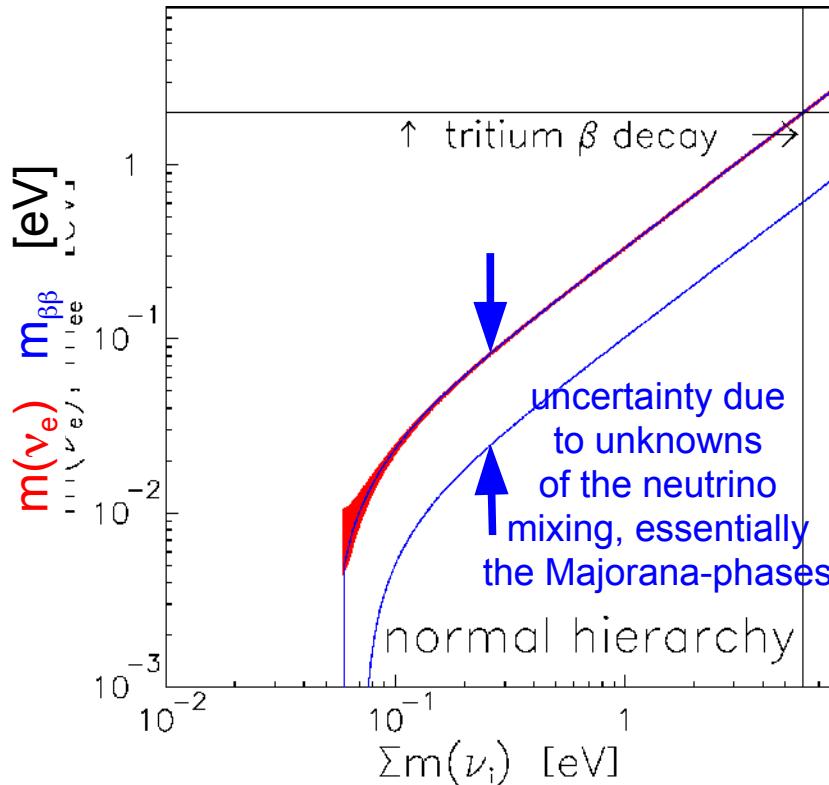
# Comparison of the different approaches to the neutrino mass

Direct kinematic measurement:  $m^2(\nu_e) = \sum |U_{ei}|^2 m^2(\nu_i)$  (incoherent)

Neutrinoless double  $\beta$  decay:  $m_{\beta\beta}(\nu) = |\sum |U_{ei}|^2 e^{i\alpha(i)} m(\nu_i)|$  (coherent)

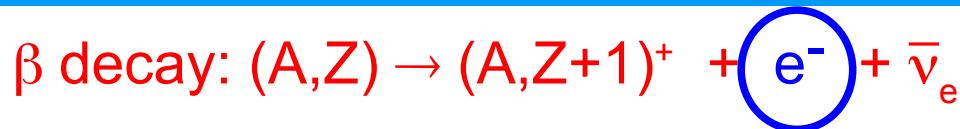
if no other particle is exchanged (e.g. R-violating SUSY)

problems with uncertainty of nuclear matrix elements



⇒ absolute scale/cosmological relevant neutrino mass in the lab by single  $\beta$  decay

# Direct determination of $m(\nu_e)$ from $\beta$ decay

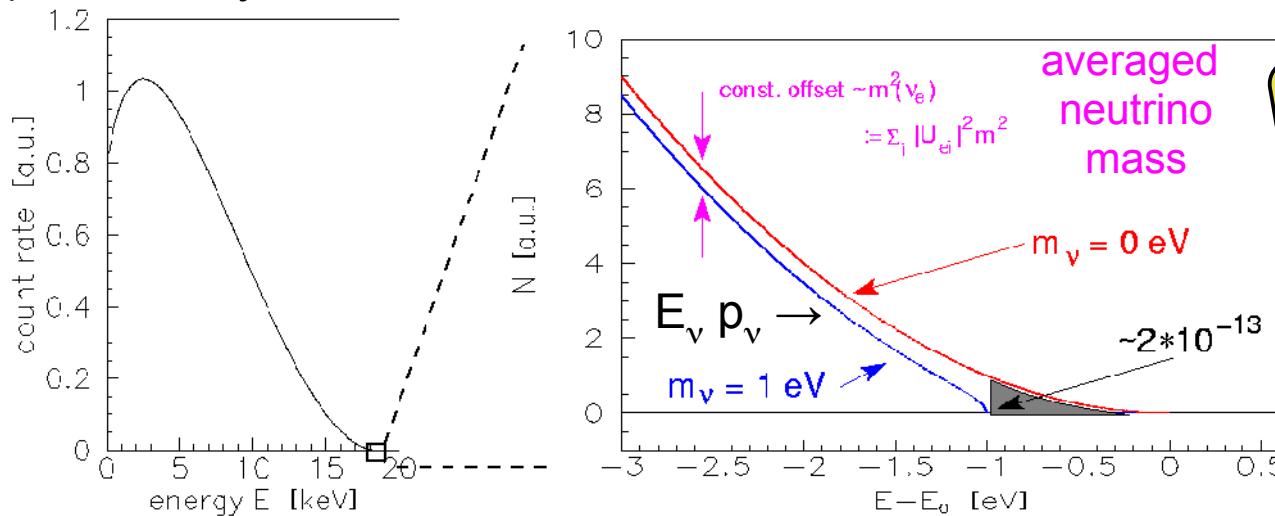


$$\beta: dN/dE = K F(E,Z) p \underbrace{E_{\text{tot}}}_{p} \underbrace{(E_0 - E_e)}_{E_\nu} \underbrace{\sqrt{(E_0 - E_e)^2 - "m(\nu_e)"^2}}_{p_\nu}$$

phase space:  $p_e$   $E_e$   $E_\nu$   $p_\nu$

Complementary to  $0\nu\beta\beta$   
and cosmology

(modified by electronic final states, recoil corrections, radiative corrections)

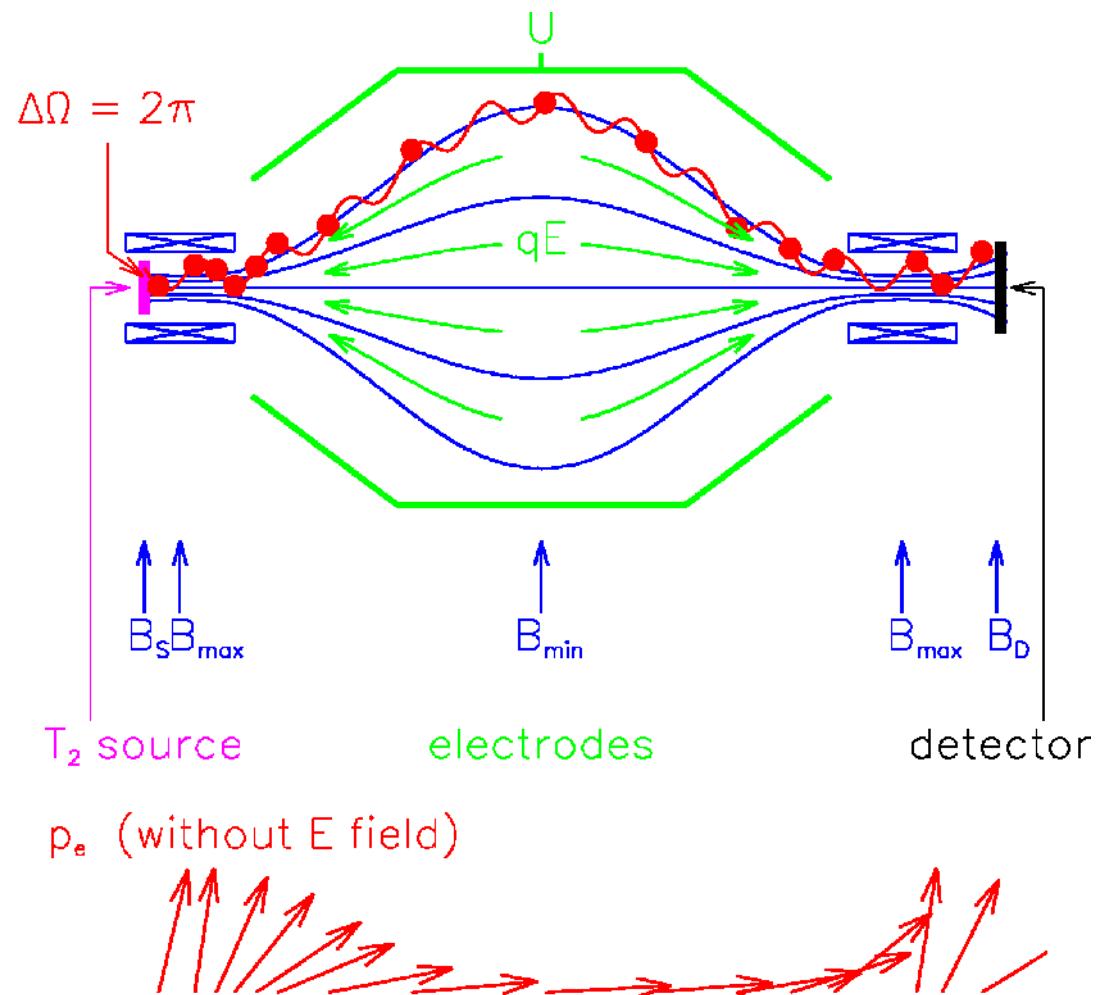


E.W. Otten & C. Weinheimer  
Rep. Prog. Phys.  
71 (2008) 086201

G. Drexlin, V. Hannen, S. Mertens,  
C. Weinheimer, Adv. High Energy  
Phys., 2013 (2013) 293986

Need: **low endpoint energy**  $\Rightarrow$  **Tritium  ${}^3\text{H}$  ( ${}^{187}\text{Re}$ ,  ${}^{163}\text{Ho}$ )**  
**very high energy resolution &**  
**very high luminosity &**  
**very low background**  $\Rightarrow$  **MAC-E-Filter**  
**(or bolometer for  ${}^{187}\text{Re}$ ,  ${}^{163}\text{Ho}$ )**

# The classical way: Tritium $\beta$ -spectroscopy with a MAC-E-Filter



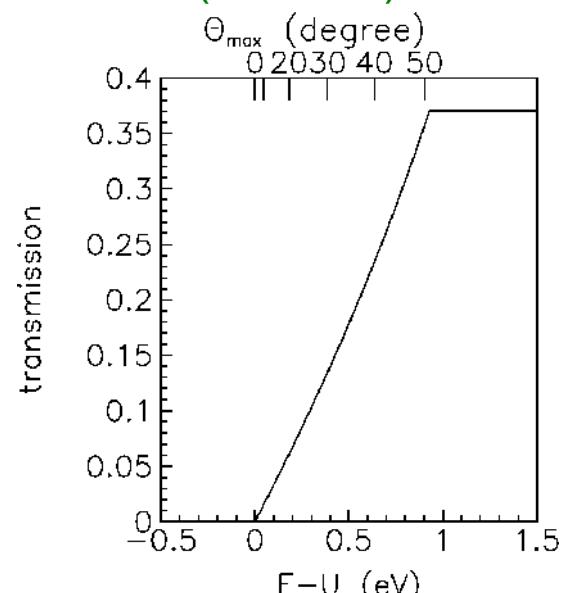
⇒ sharp integrating transmission function without tails →

Magnetic Adiabatic Collimation + Electrostatic Filter  
(A. Picard et al., Nucl. Instr. Meth. 63 (1992) 345)

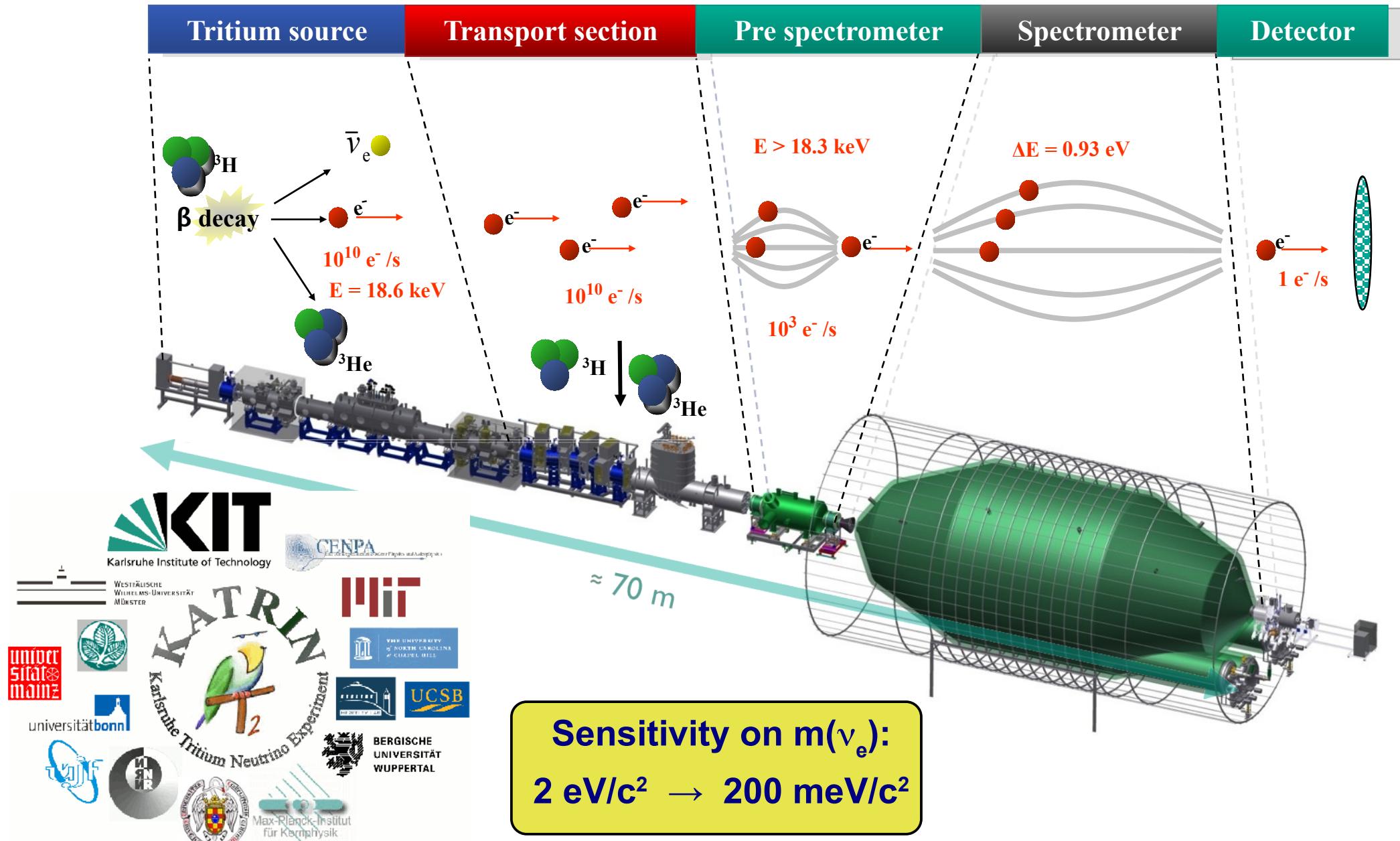
- Two supercond. solenoids compose magnetic guiding field
- adiabatic transformation:  
 $\mu = E_{\perp} / B = \text{const.}$   
⇒ parallel  $e^-$  beam
- Energy analysis by electrostat. retarding field  

$$\Delta E = E \cdot B_{\min} / B_{\max}$$

$$= 0.93 \text{ eV (KATRIN)}$$



# The Karlsruhe Tritium Neutrino Experiment KATRIN - overview

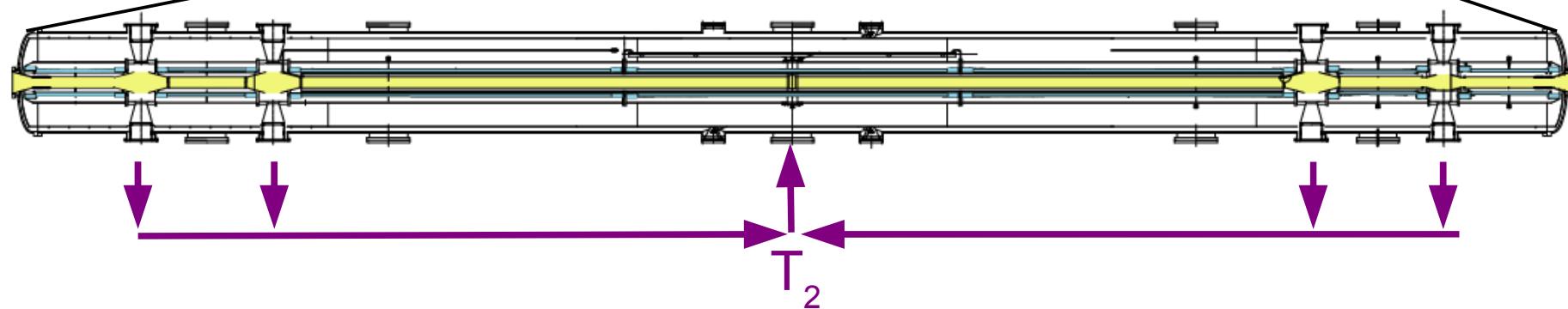
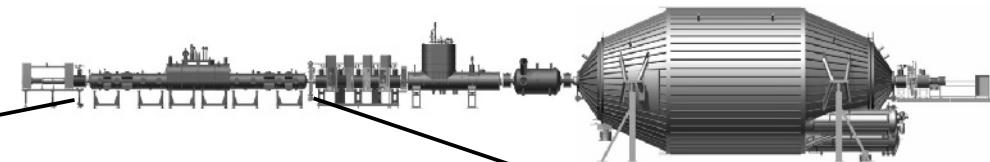


# Molecular Windowless Gaseous Tritium Source WGTS

per mill stability source strength request:

$$\frac{dN}{dt} \sim f_T \cdot N / \tau \sim n = f_T \cdot p V / R T$$

tritium fraction  $f_T$  & ideal gas law



WGTS: tub in long superconducting solenoids  
 $\varnothing$  9cm, length: 10m,  $T = 30$  K

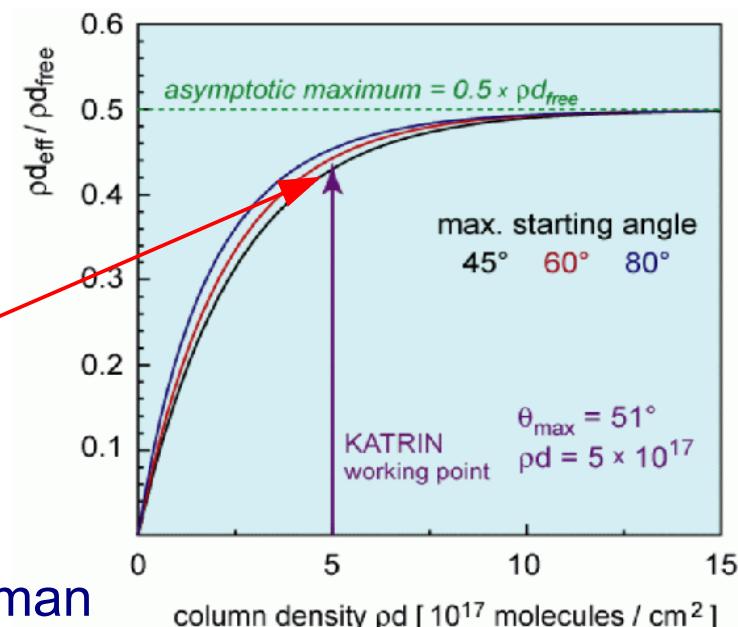
Tritium recirculation (and purification)

$$p_{\text{inj}} = 0.003 \text{ mbar}, q_{\text{inj}} = 4.7 \text{ Ci/s}$$

allows to measure with near to maximum count rate using

$$pd = 5 \cdot 10^{17} / \text{cm}^2$$

with small systematics



check column density by e-gun,  $T_2$  purity by laser Raman

# Status of Windowless Gaseous molecular Tritium Source WGTS

## Assembly of beam tube, magnets and cryostat:

Temperature stability tests of  
„demonstrator“ very successful

Management now in the hand of KIT  
progress according schedule

Arrived at KIT on Sept 10, 2015 !



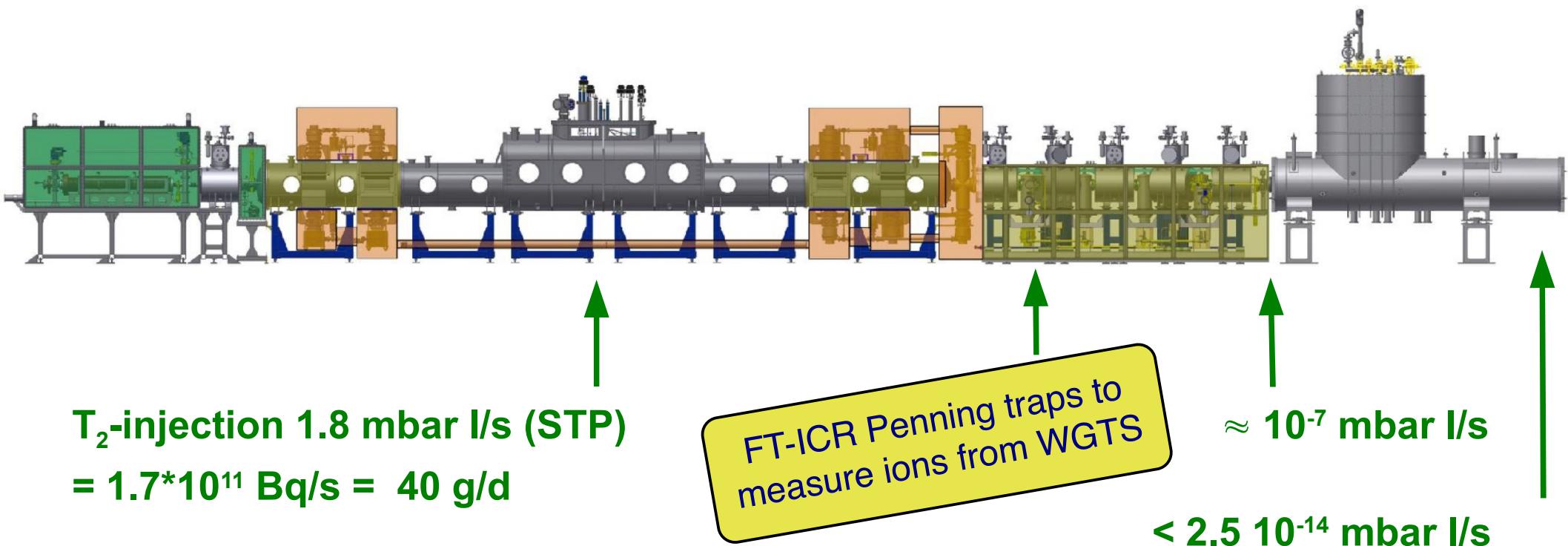
# Transport and differential & cryo pumping sections

Monitoring  
& calibration  
system

Molecular windowless  
gaseous tritium source

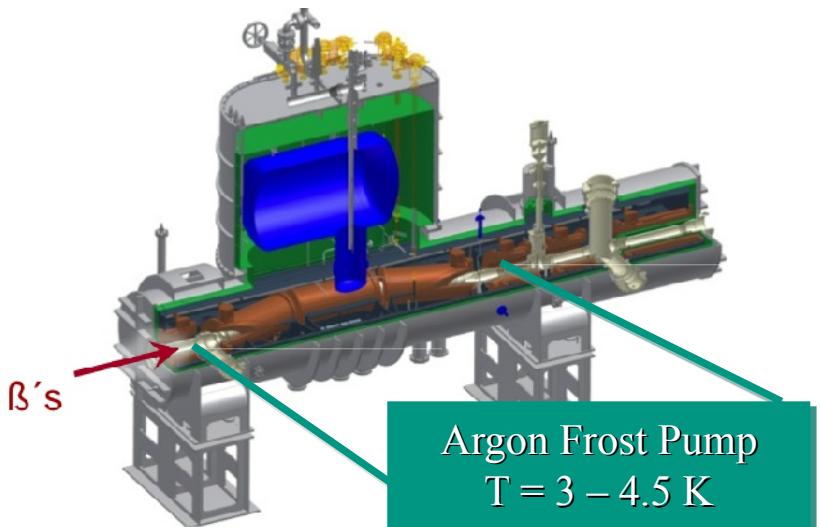
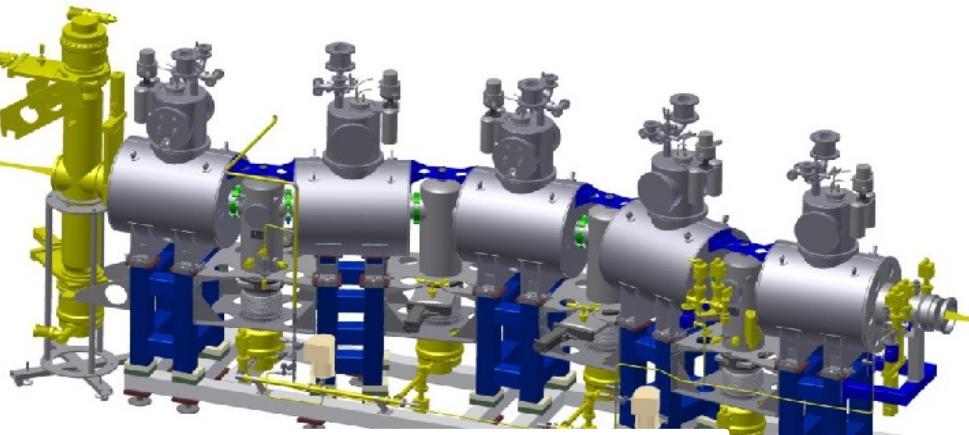
Differential  
pumping

Cryogenic  
pumping  
with Argon snow  
at LHe temperatures



⇒ adiabatic electron guiding & T<sub>2</sub> reduction factor of  $\sim 10^{14}$

# Differential and cryo pumping sections

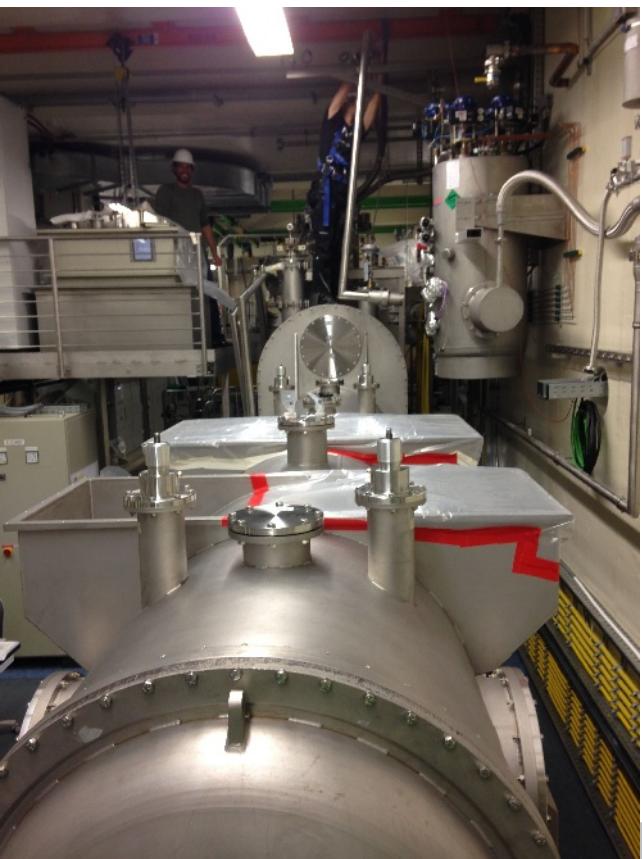
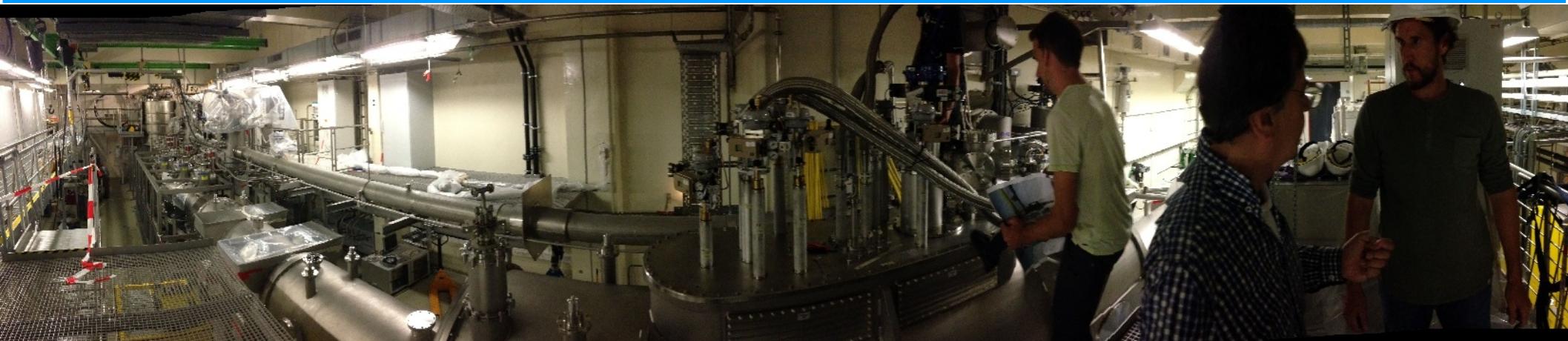


- active pumping: 4 TMPs
- Tritium retention:  $>10^7$
- magnetic field: 5.6 T
- **delivered to KIT: July 30, 2015**

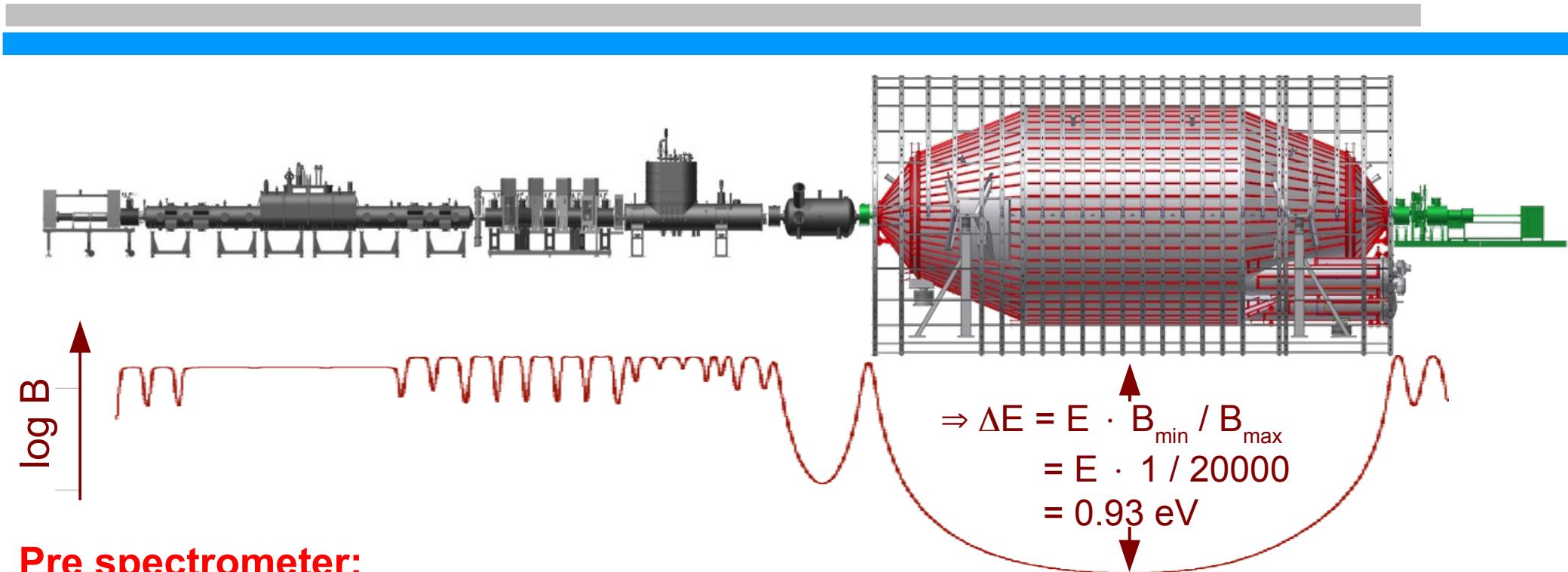
- under commissioning



# The source and transport section is really there !



# KATRIN spectrometers



## Pre spectrometer:

- successful tests & developments of new concepts

## Main spectrometer:

- huge size: 10m diameter, 24m length  
1240 m<sup>3</sup> volume, 690 m<sup>2</sup> inner surface
- ultra-high vacuum:  $p = O(10^{-11} \text{ mbar})$
- ultra-high energy resolution:  $\Delta E = 0.93 \text{ eV}$
- vacuum vessel on precise high voltage (ppm precision)



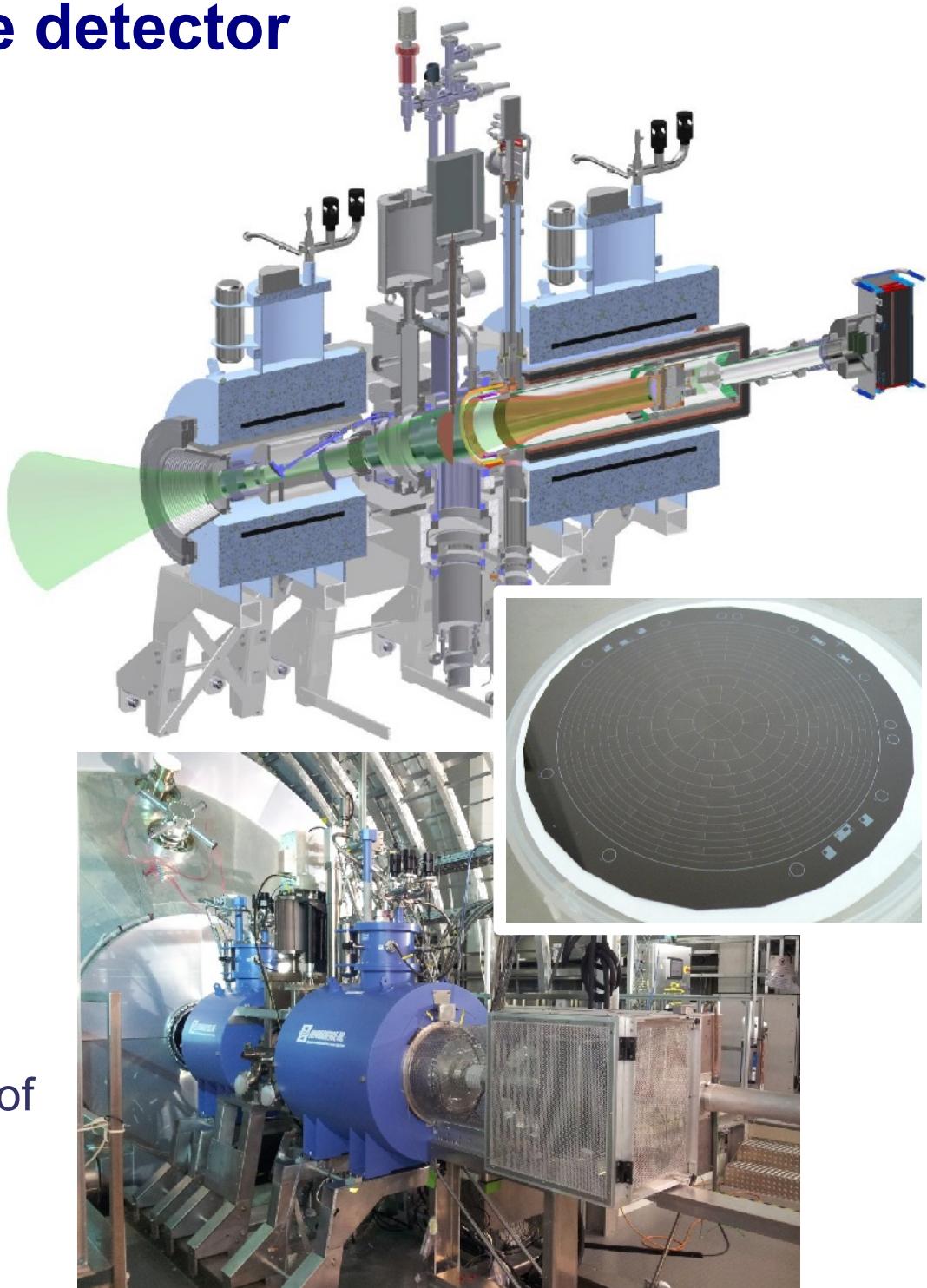
# The detector

## Requirements

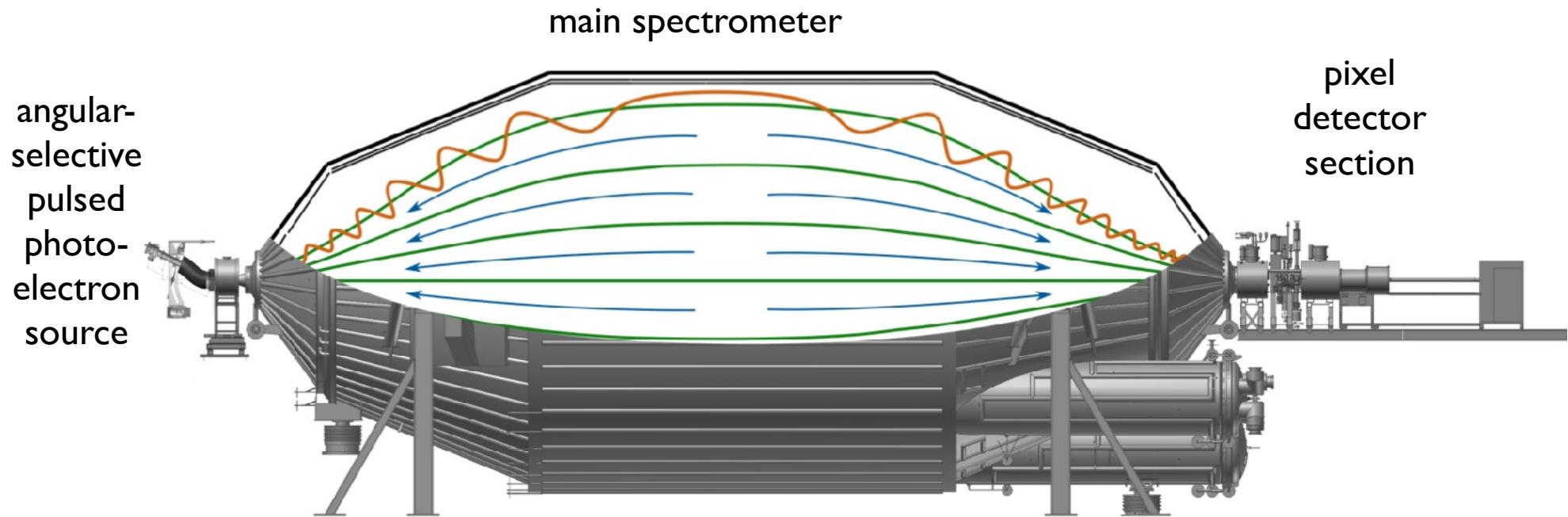
- detection of  $\beta$ -electrons (mHz to kHz)
- high efficiency (> 90%)
- low background (< 1 mHz)  
(passive and active shielding)
- good energy resolution (< 1 keV)

## Properties

- 90 mm Ø Si PIN diode
- thin entry window (50nm)
- detector magnet 3 - 6 T
- post acceleration (30kV)  
(to lower background in signal region)
- segmented wafer (148 pixels)
  - record azimuthal and radial profile of the flux tube
  - investigate systematic effects
  - compensate field inhomogeneities



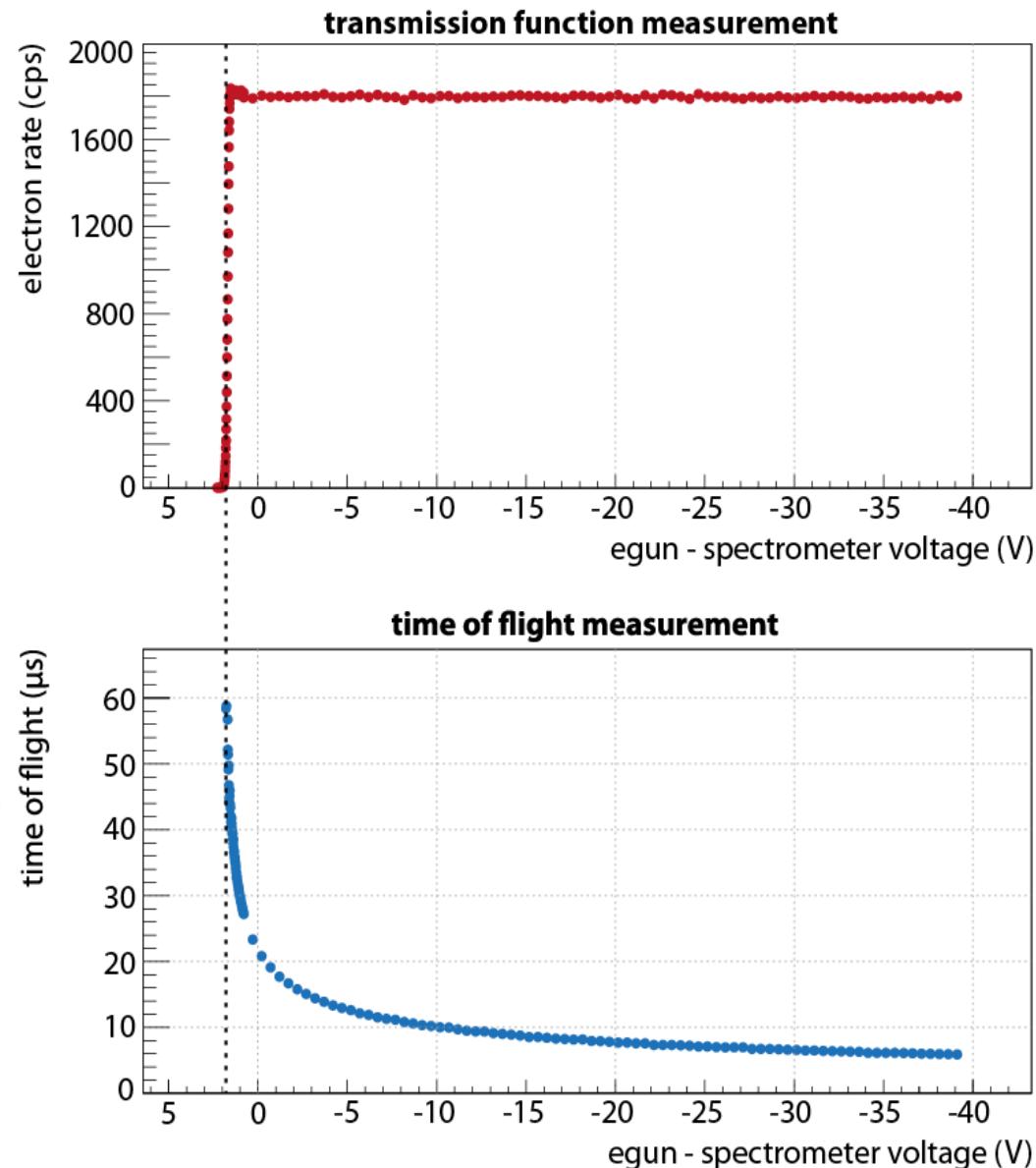
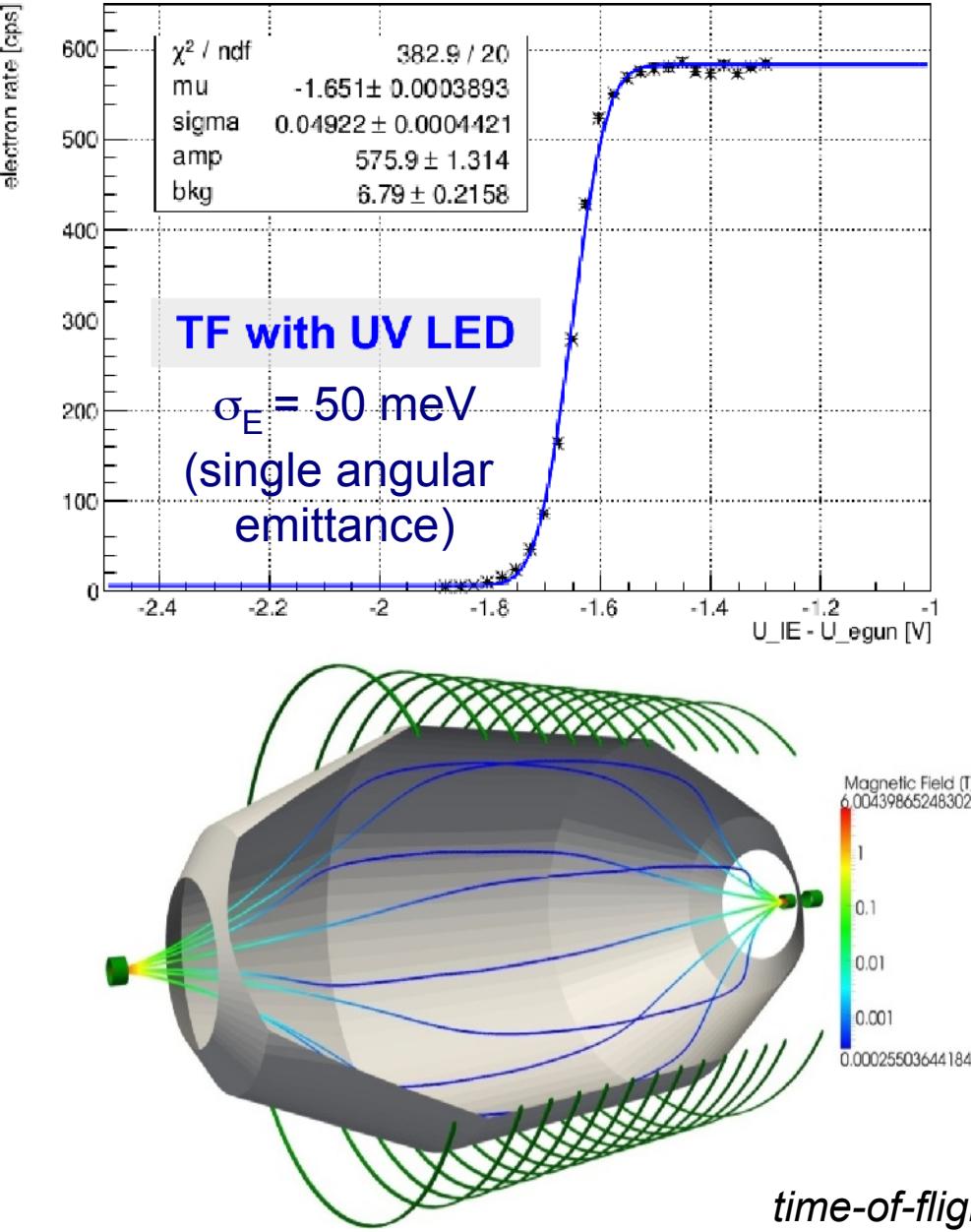
# Main spectrometer and detector commissioning – objectives



## Primary objectives:

- test of individual hardware, software and slow control components
- provide ultra high vacuum conditions at the  $p \approx 10^{-11}$  mbar level
- detailed understanding of the transmission properties of this MAC-E-Filter ( $E = 18.6$  keV with  $\Delta E = 0.93$  eV resolution) and compare to simulation with Kasseiopia
- detailed understanding and passive & active control of background processes

# Commissioning of main spectrometer and detector



# Suppress secondary electron background from walls on high potential

Secondary electrons from wall/electrode

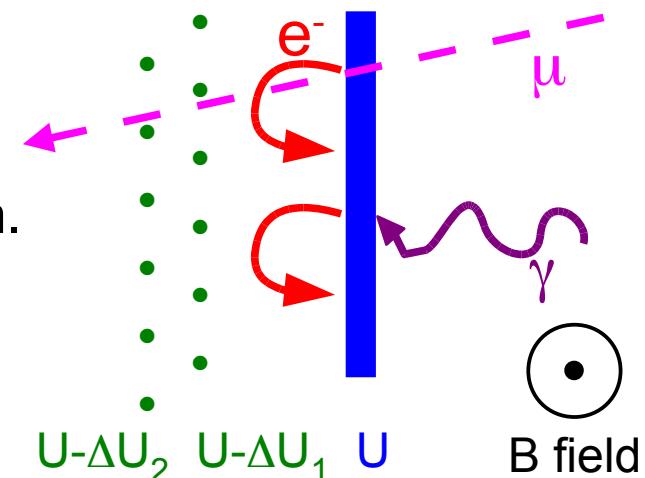
by cosmic rays, environmental radioactivity, ...

Excellent magnetic shielding by nearly perfect axial sym.

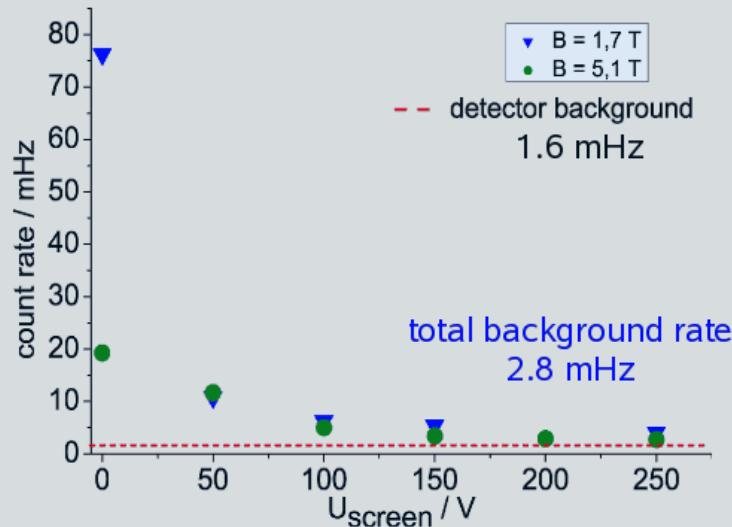
Additionally double layer wire electrode

on slightly more negative potential

(ca. 23,000 wires, 200  $\mu\text{m}$  precision, UHV compatible)



Background suppression successfully tested at the Mainz MAC-E filter:



Dipl. thesis B. Ostrick (U Mainz, 2002),  
PhD thesis B. Flatt (U Mainz, 2004)

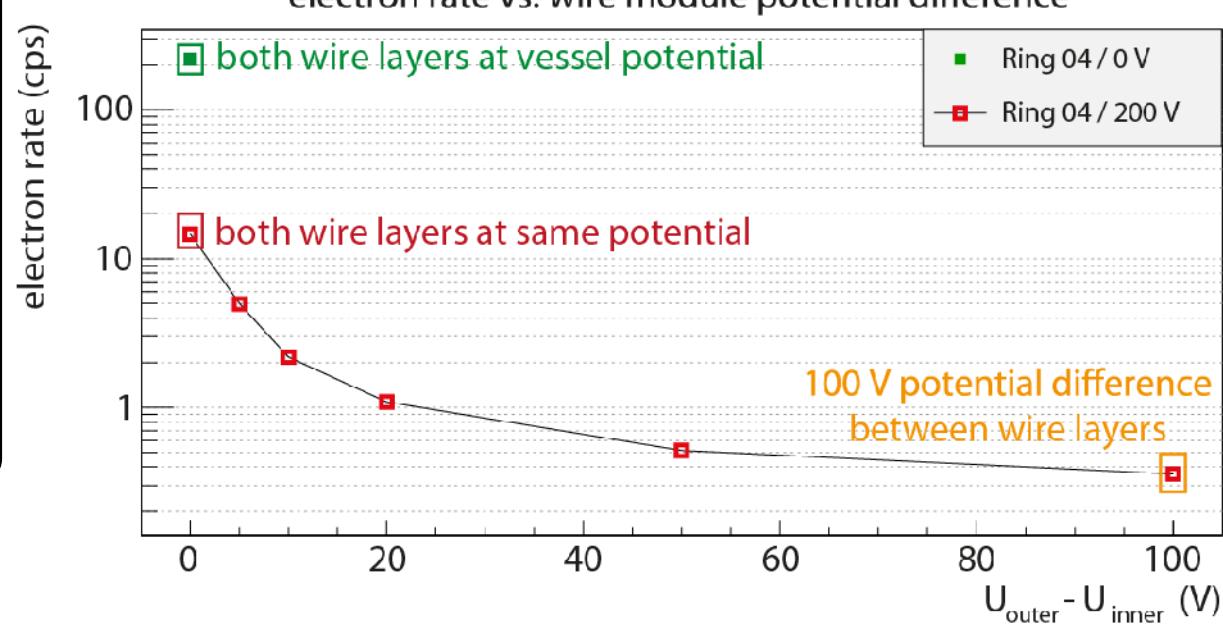
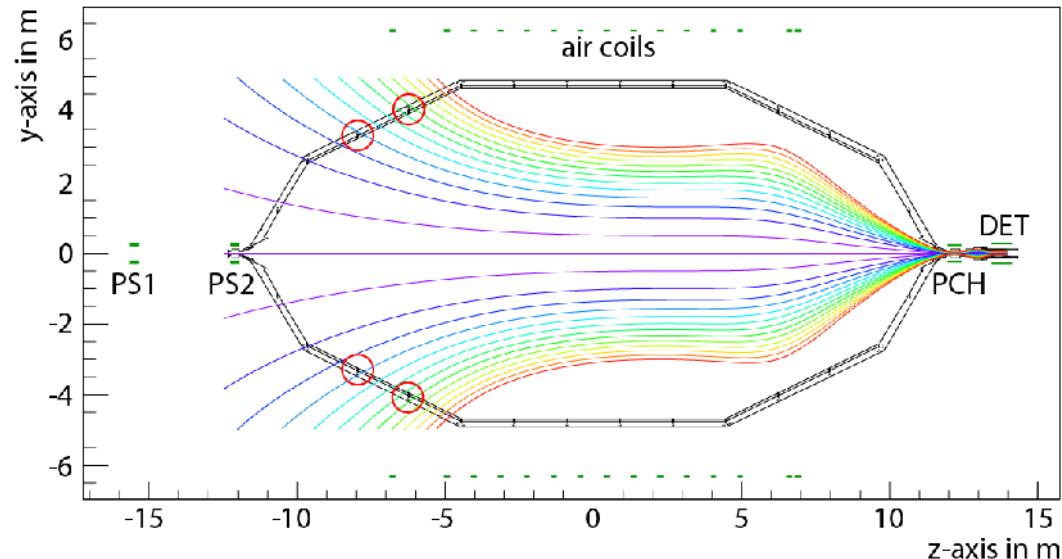


# Background suppression by wire electrode

Investigating secondary electrons from the walls by directly looking to the wall and applying an asymmetric magnetic field

→ there are many secondary electrons from the walls, which are shielded by the wire electrode

→ but there is a very efficient magnetic shielding as well



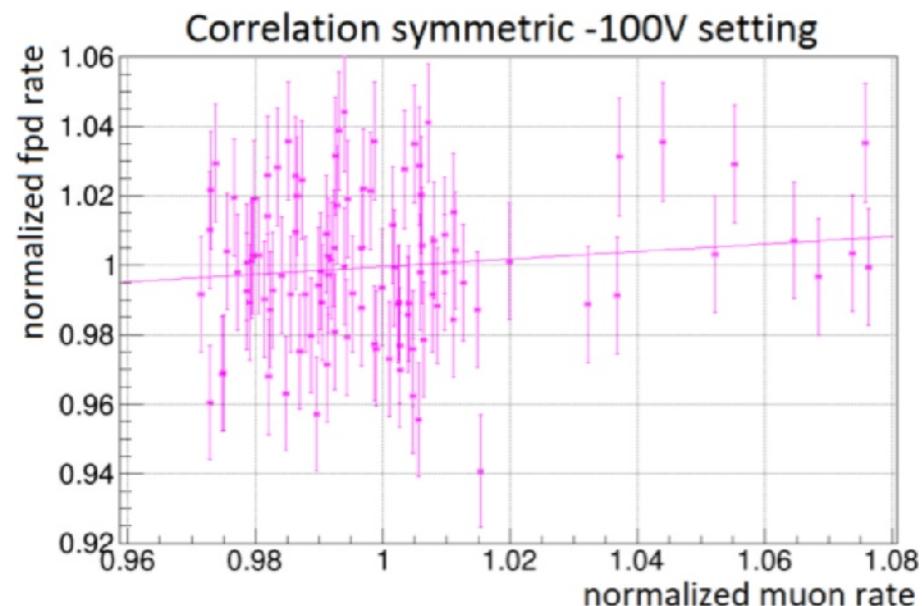
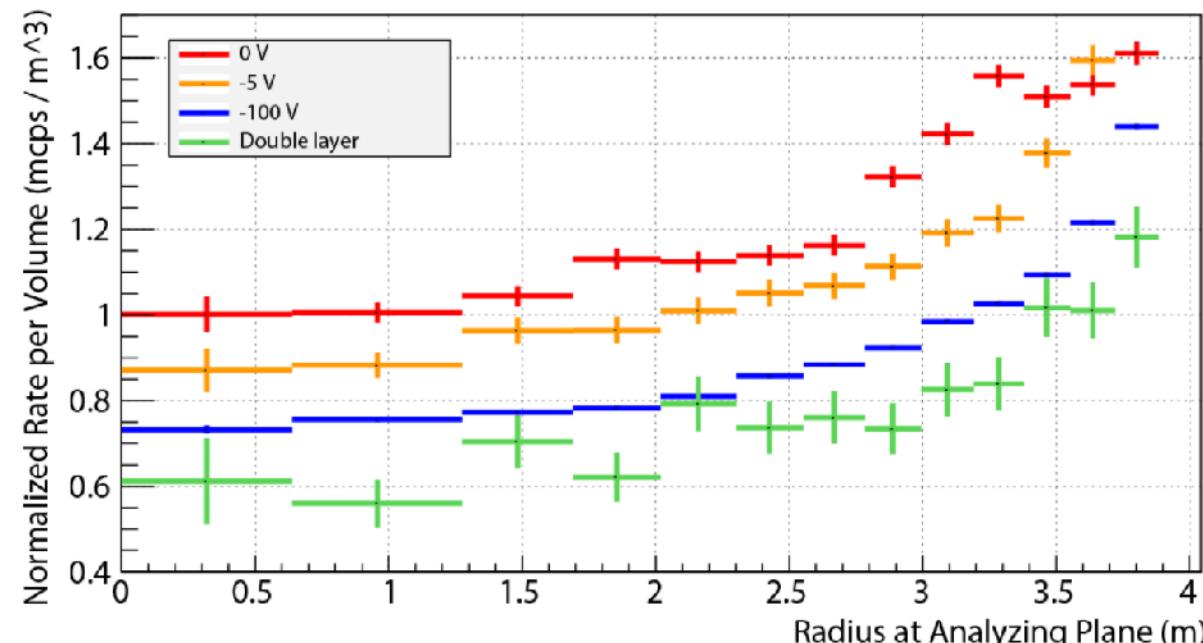
# Background suppression by wire electrode

Significant background reduction by single & dual layer wire electrode  
(some electric shorts prohibit full dual layer operation)

but the remaining background does not seem to originate from cosmic muons or other charged particles, but from neutrals ...

slope a:  $0.11 \pm 0.07$

correlation factor:  $0.11 \pm 0.09$



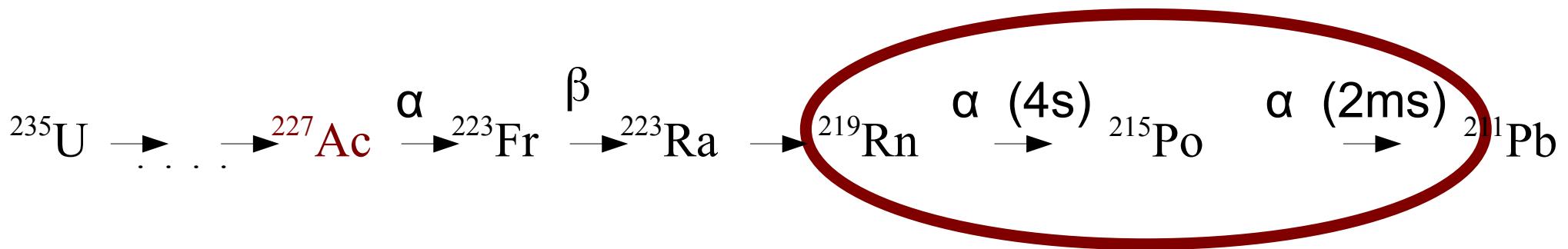
# Radon background in the pre spectrometer from the non-evaporable getter (NEG) pumps

Getter strips (SAES ST707) adsorbing residual gas ( $H_2$ )  
are essential to reach  $10^{-11}$  mbar, composition:

70 % Zirconium: contains  $^{227}Ac$

25 % Vanadium

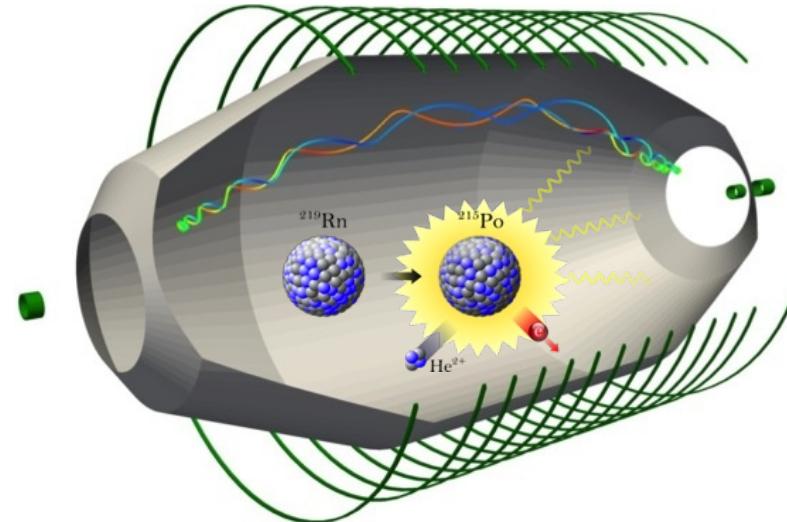
5 % Iron



$^{219}Ra$  is gaseous  $\rightarrow$  spectrometer  $\rightarrow$  ionizations  $\rightarrow$  **background**

# Secondary electron background from radon decays in the volume

- $^{219,220}\text{Rn}$  emanation mainly from SAES getter pumps (zirconium vanadium iron alloy) conversion, Auger, shake-off electrons can get stored by magnetic mirror effect

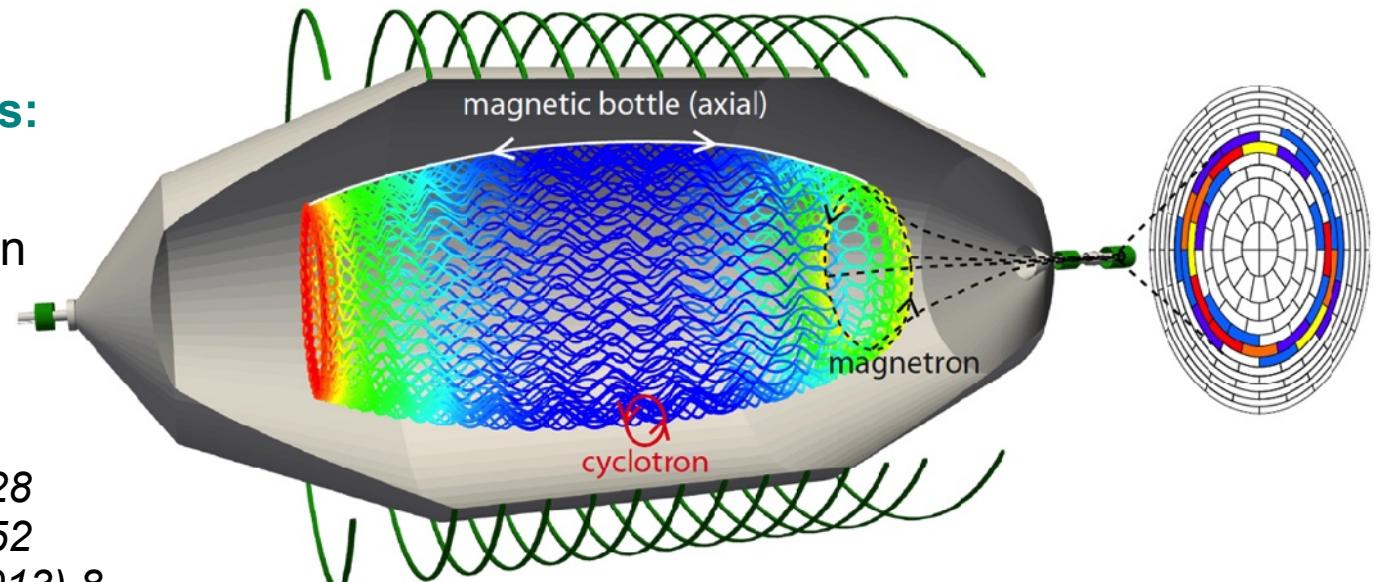


## ■ background process continues:

- ionization of residual gas → secondary electrons
- primary electron energies:  $100 \text{ eV} < E < 500 \text{ keV}$
- up to 5000 secondary electrons per stored primary
- significant background increase for hours

## ■ stored multi-keV electrons:

rapid cyclotron motion  
intermediate axial oscillation  
slow magnetron drift

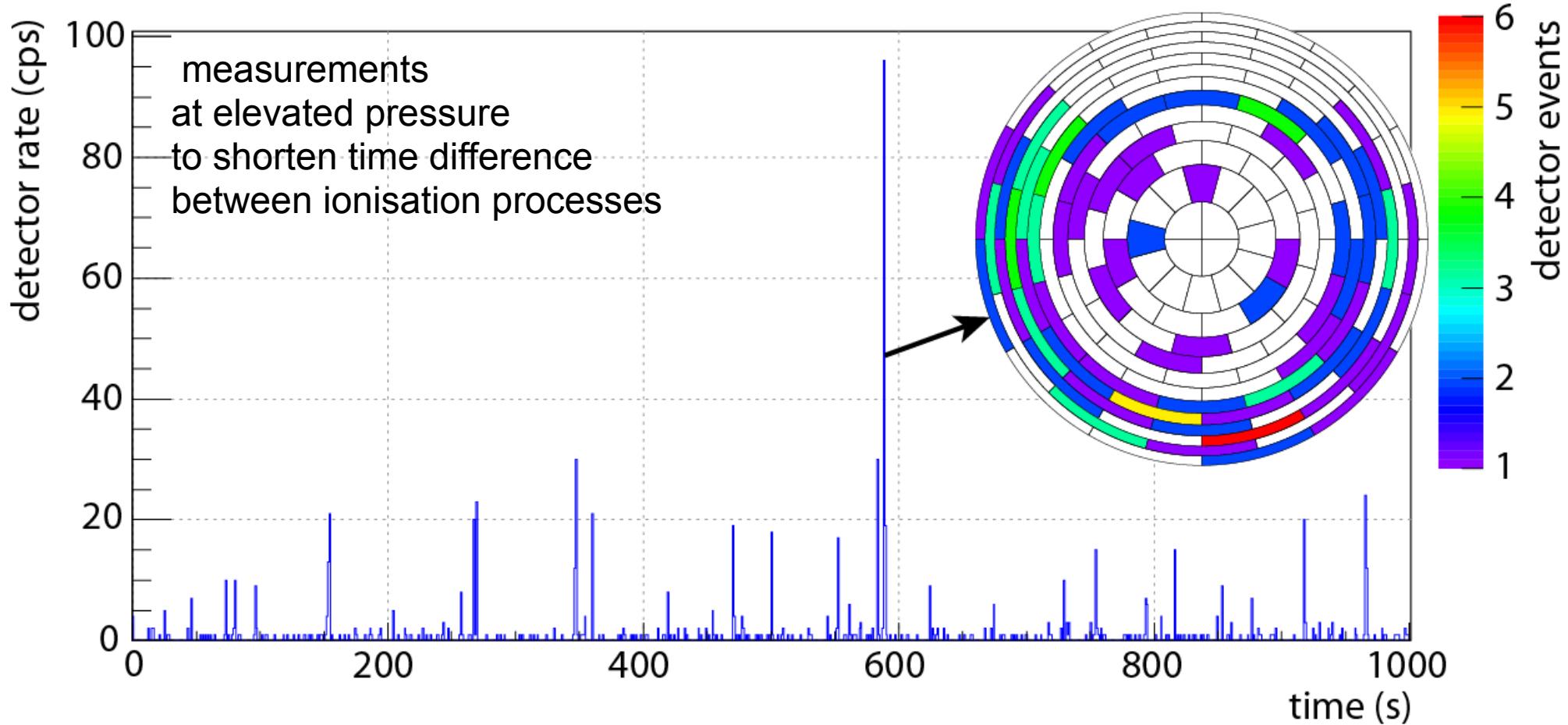


F. Fränkle et al., APP 35 (2011) 128

S. Mertens et al., APP 41 (2012) 52

N. Wandkowsky et al., NJP 15 (2013) 8

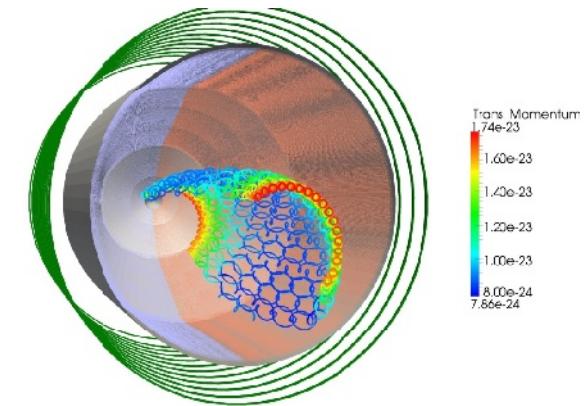
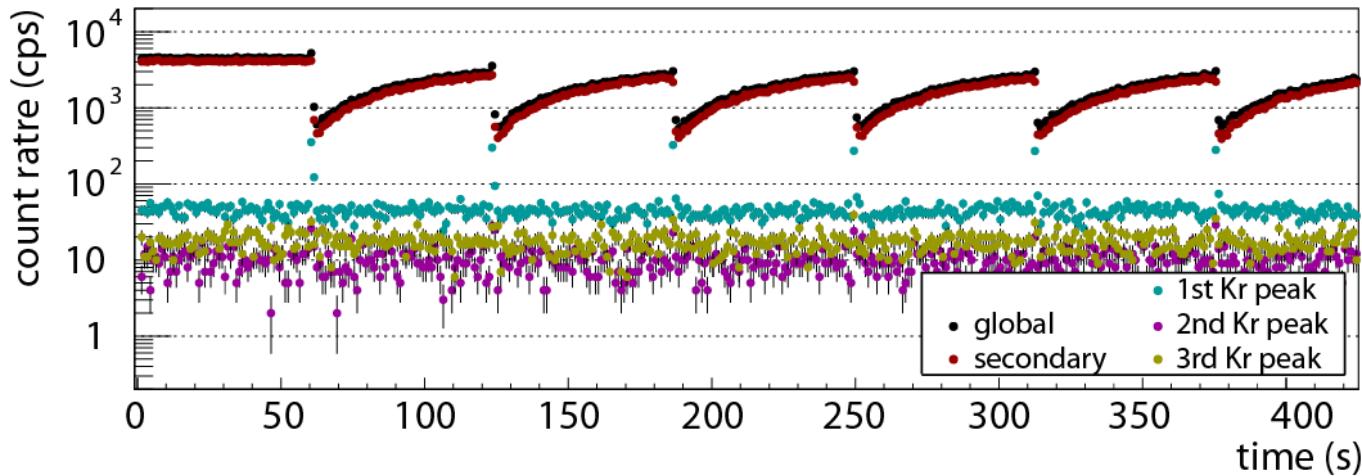
# Understanding the background: Radon and other background sources



**High energy electrons from radon decays:**  
→ multiple, spike-like events  
with a ring pattern at the detector (magnetron motion)

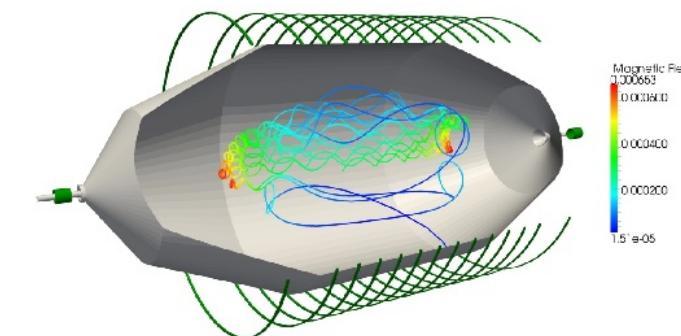
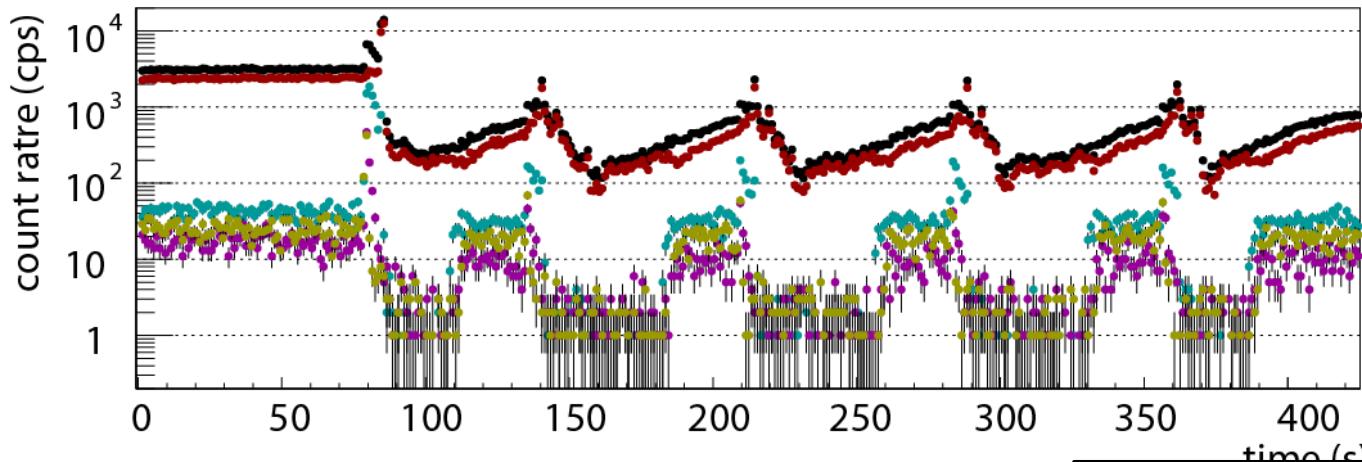
# Active stored particle removal by electric dipole and magnetic zeroing

## electric dipole



$$v_D = E \times B / B^2$$

## magnetic pulse

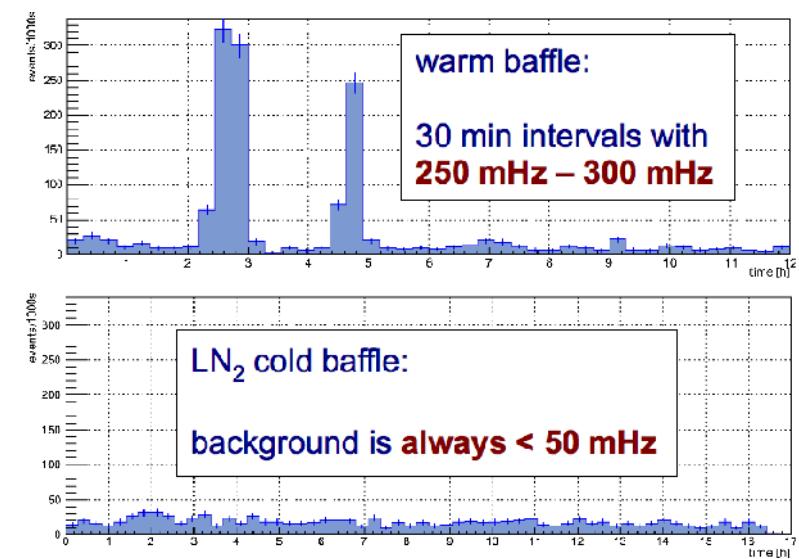
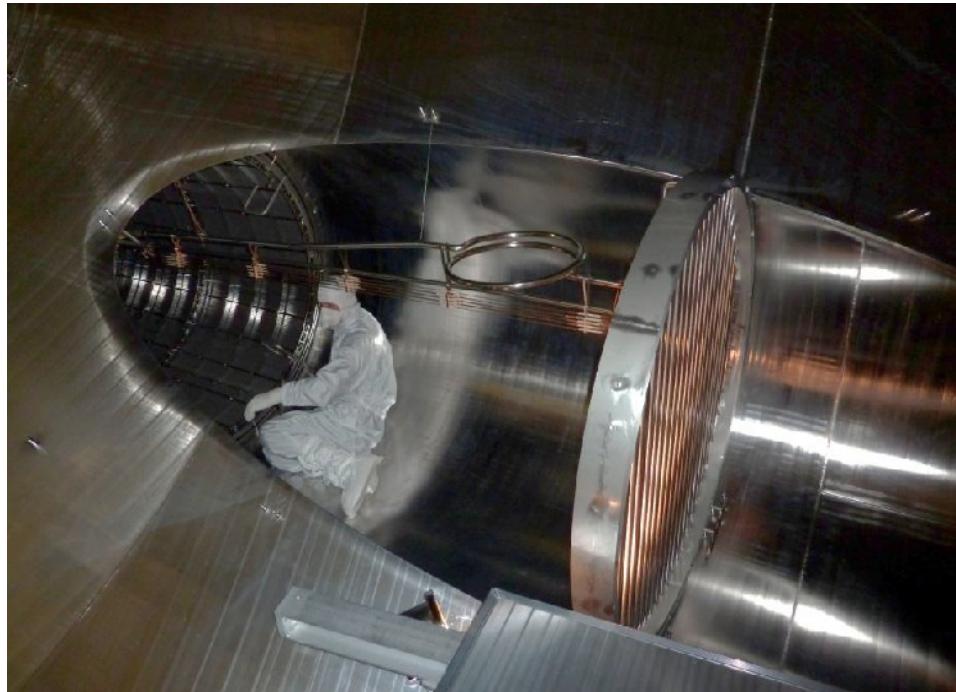
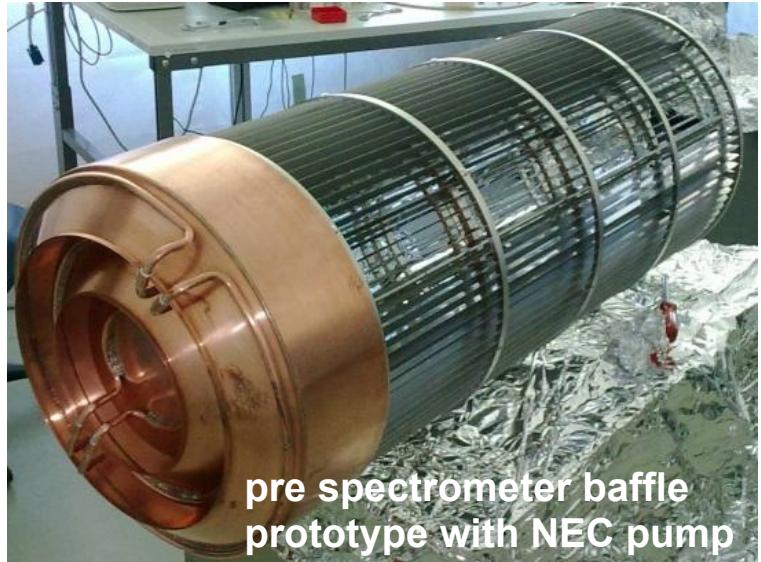


loss of magnetic guidance

$^{83m}\text{Kr}$  injected  
to enhance number of stored particles

2014: all 14 axial air coils are instrumented  
with fast magnet puls switches in SDS II  
(compared to 1 switch in SDS I in 2013)

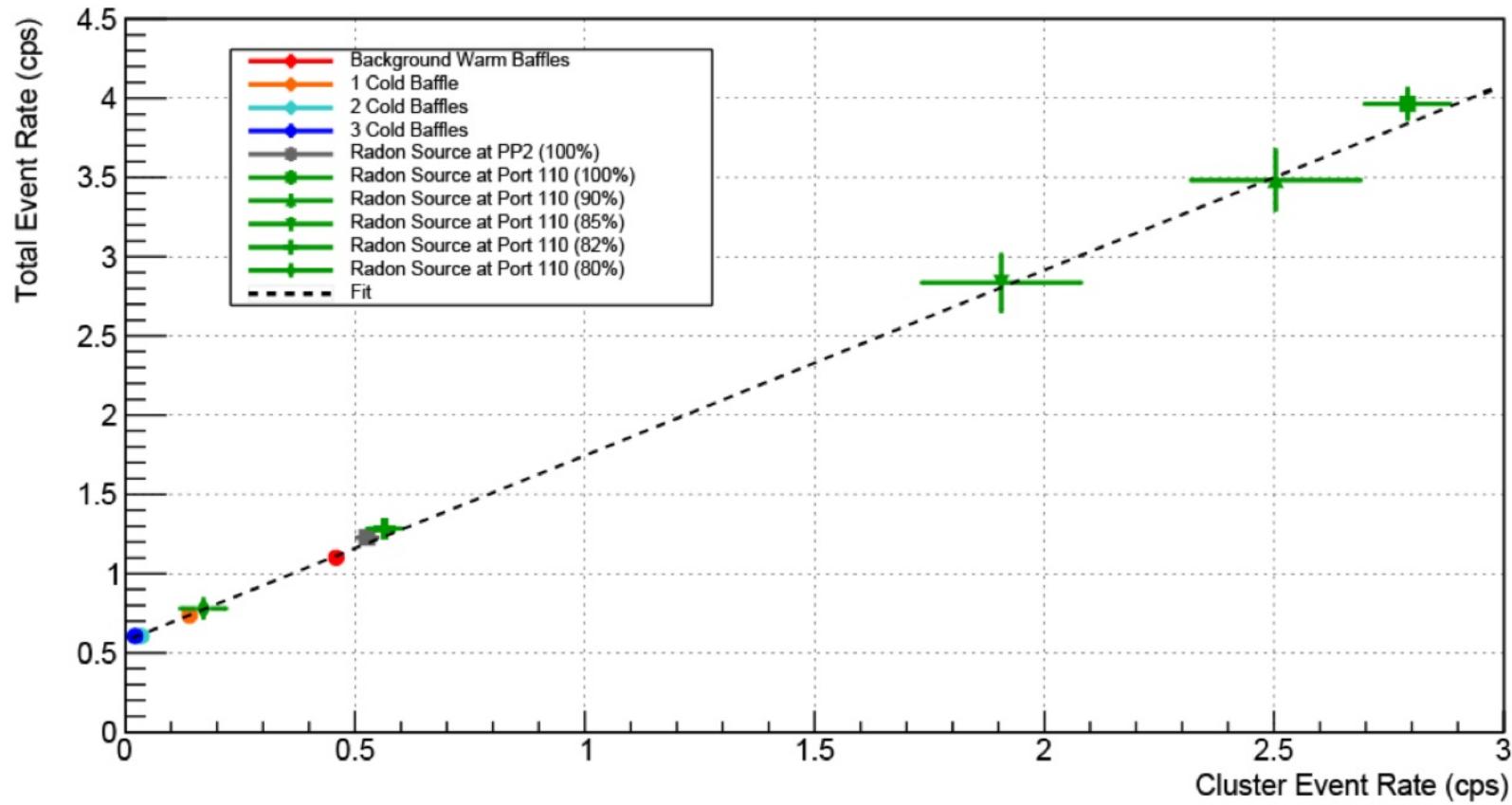
# Radon elimination by LN<sub>2</sub>-cooled baffles in the pre & main spectrometer



successful application at pre spectrometer



# Radon induced background $^{219}\text{Rn}$ from getter and artefical $^{220}\text{Rn}$ source

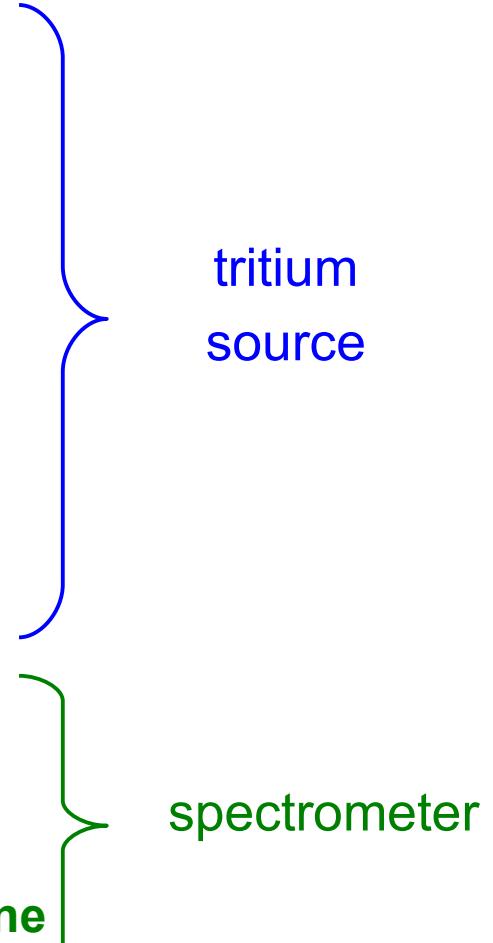


- radon-induced background is very efficiently eliminated by  $\text{LN}_2$  baffles
  - residual non-radon background of about 600 mcps
  - optimal magnetic field settings: 477 mcps = reference background rate (SDS2)
- SDS2b after baking to reach better understanding of residual background:  
bg < 300 mcps reached (July 2015, preliminary)

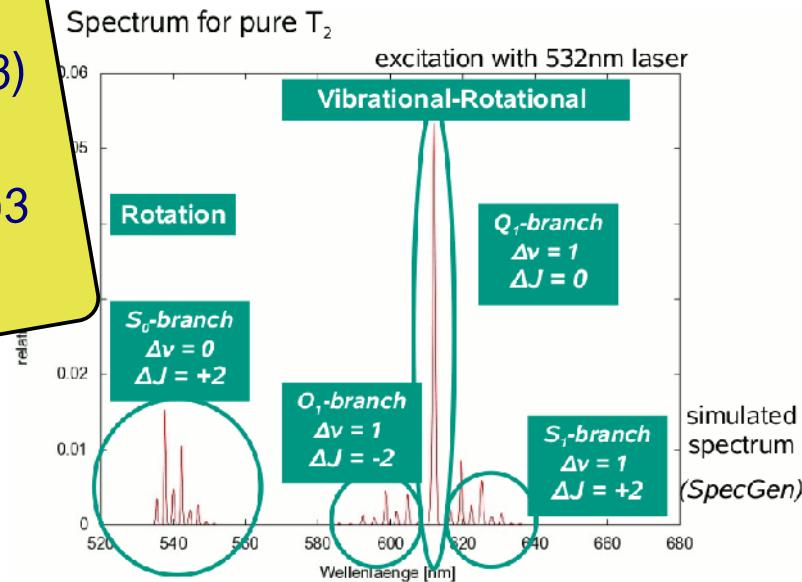
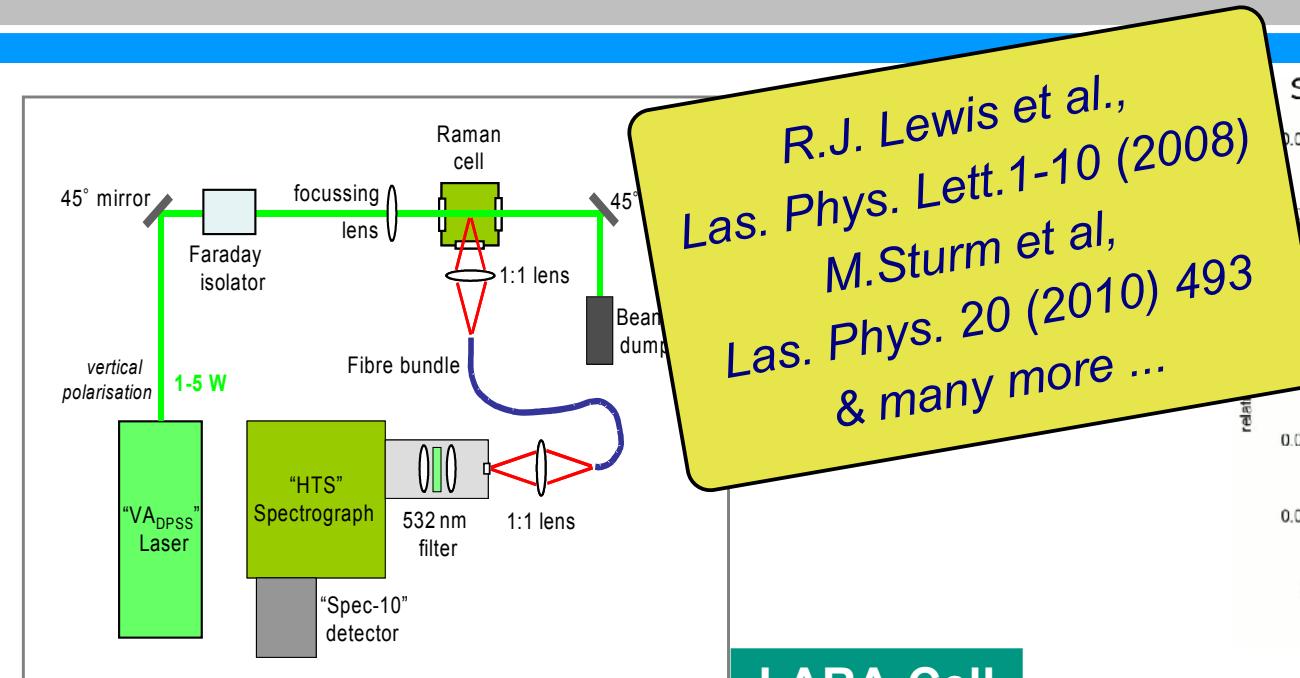
# Systematic uncertainties

As smaller  $m(v)$  as smaller the region of interest below endpoint  $E_0$   
→ quantum mechanical thresholds help a lot !

**A few contributions with  $\Delta m_v^2 \leq 0.007 \text{ eV}^2$  each:**

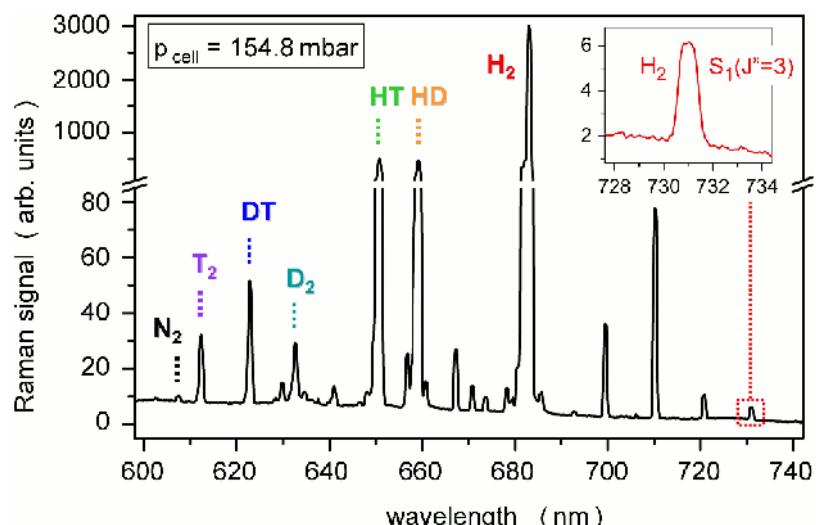
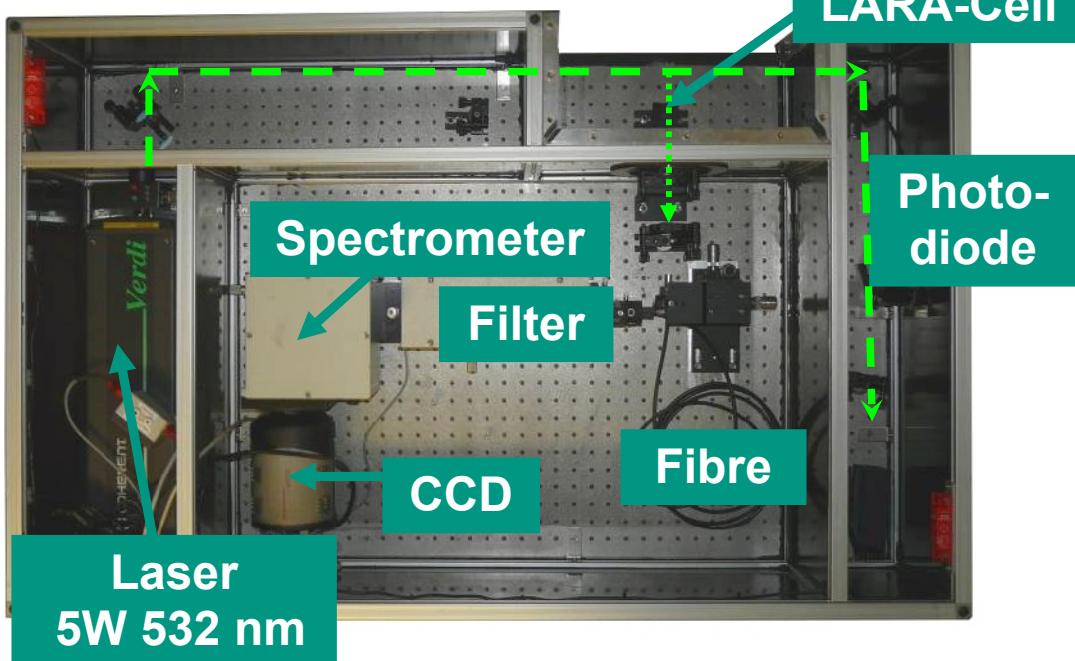
1. inelastic scatterings of  $\beta$ 's inside WGTS
    - **dedicated e-gun measurements**, unfolding of response fct.
  2. fluctuations of WGTS column density (required < 0.1%)
    - rear detector, Laser-Raman spectroscopy, T=30K stabilisation,  
**e-gun measurements**
  3. WGTS charging due to remaining ions (MC:  $\varphi < 20\text{mV}$ )
    - **monocrystalline rear plate short-cuts potential differences**
  4. final state distribution
    - **reliable quantum chem. calculations**
  5. transmission function
    - detailed simulations, **angular-selective e-gun measurements**
  6. HV stability of retarding potential on ~3ppm level required
    - **precision HV divider (with PTB), monitor spectrometer beamline**
- 

# Measurement of tritium concentration by laser Raman spectroscopy

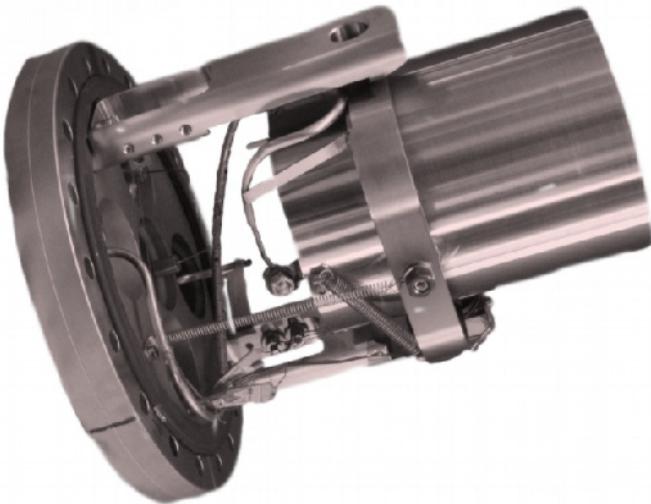


LARA-Cell

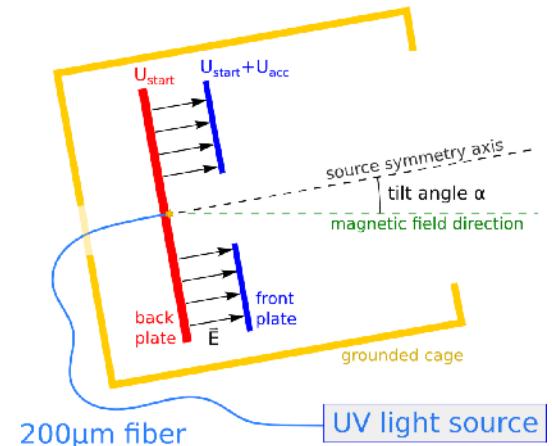
$$\begin{aligned} \text{H}_2 / \text{HD} / \text{T}_2 / \text{DT} / \text{HT} \\ = 0.820 / 0.083 / 0.003 / 0.005 / 0.085 \end{aligned}$$



# An angular-selective, pulsed UV photoelectron e-gun

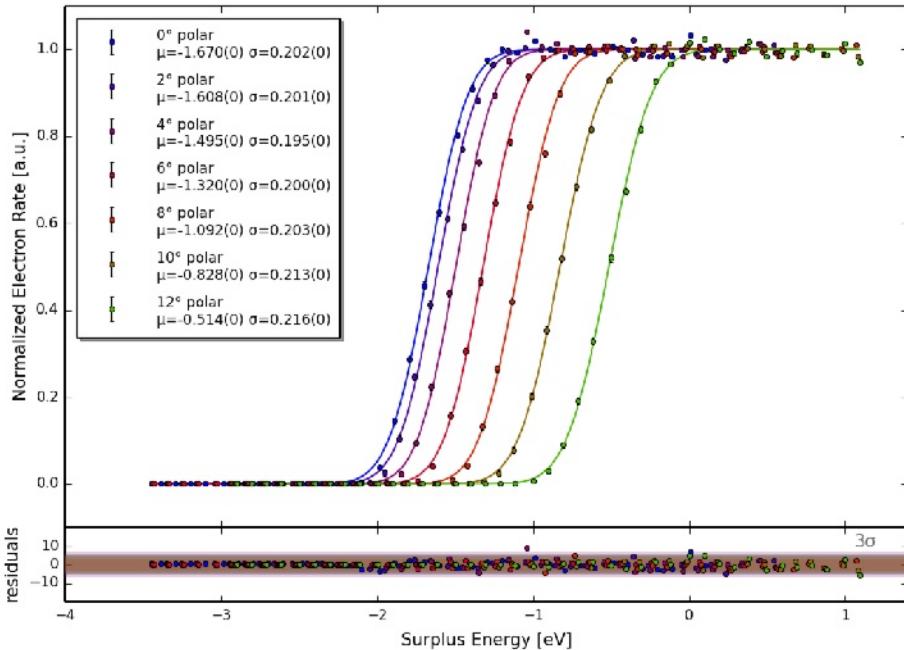


see poster by Jan Behrens



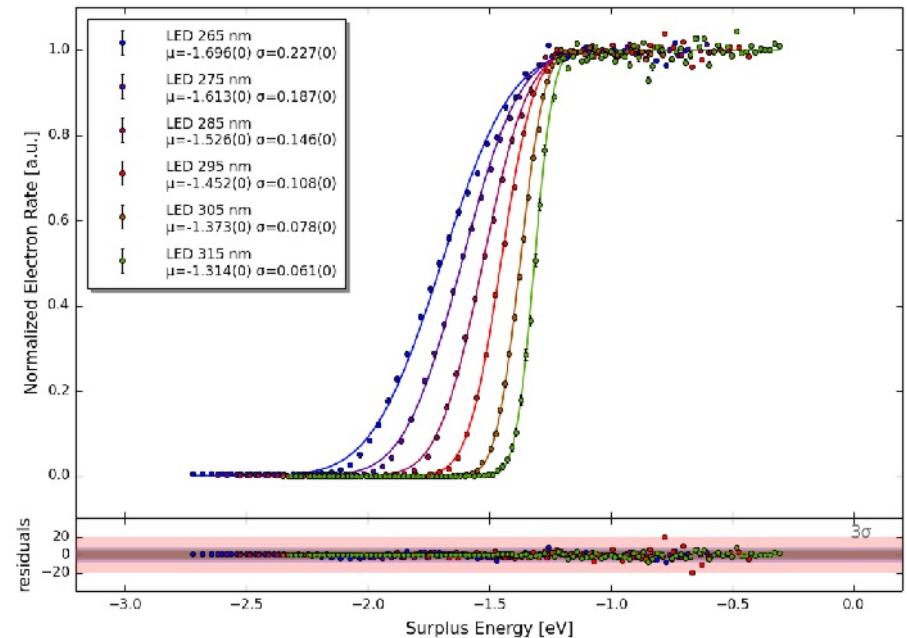
## Different tilding angles of egun

Transmission Function at different plate angles  
(Au photocathode, Mainspec at 18.6 kV / 3.8 Gs)



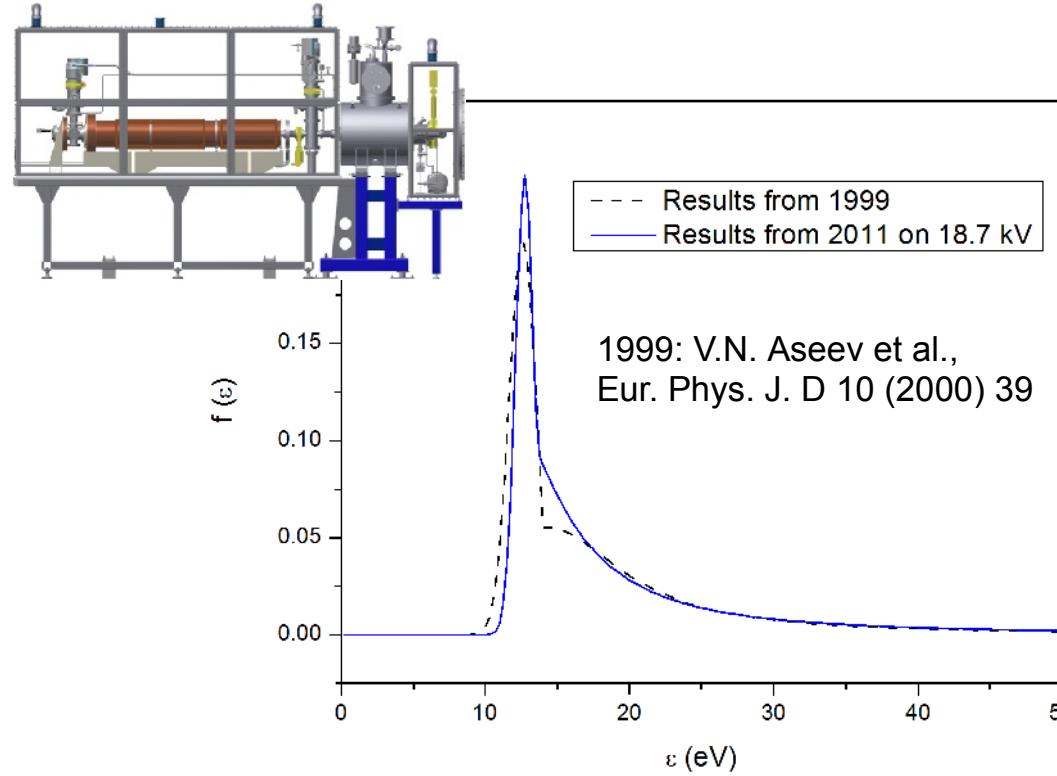
## Different UV wavelengths

Transmission Function at different LED wavelengths  
(Au photocathode, Mainspec at 200 V / 3.8 Gs)

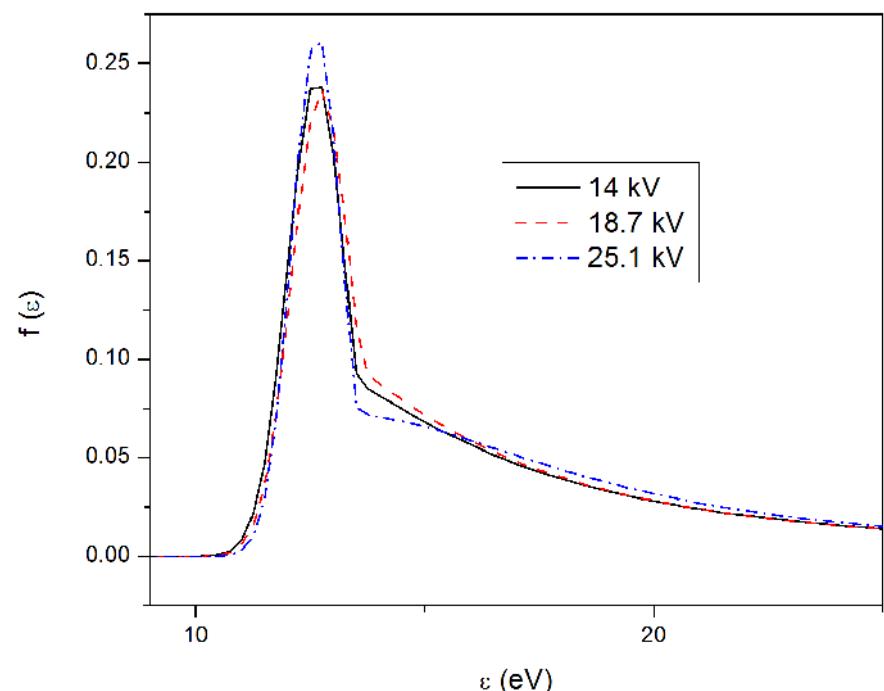


# Recently achieved results at Troitsk

Measurement of electron scattering on  $H_2$  at 14, 18, and 25 keV. New data on excitation and ionization spectra obtained with spectrometer resolution of about 1 eV.

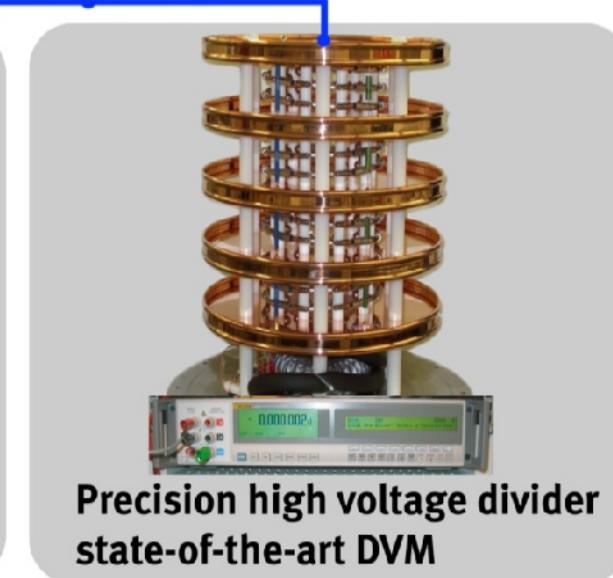
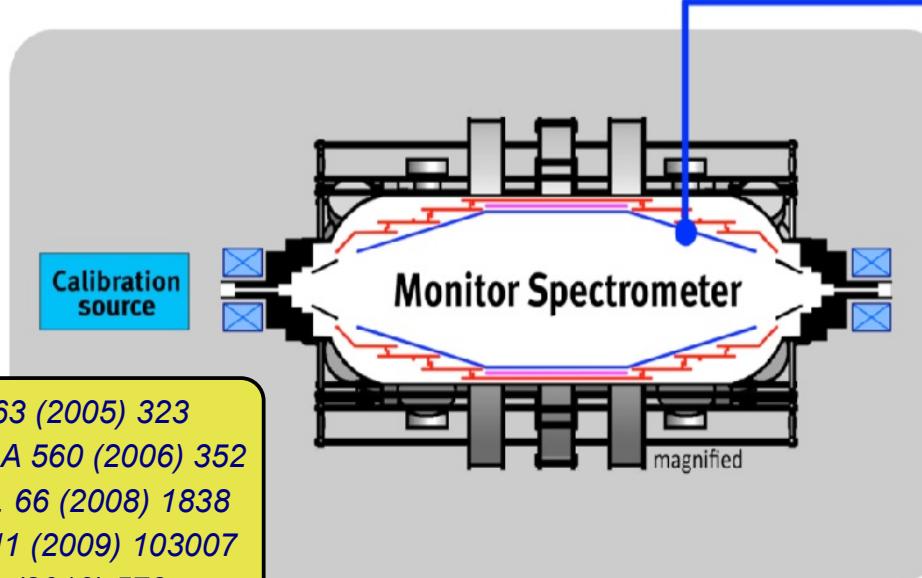
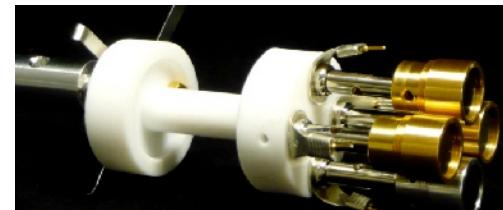
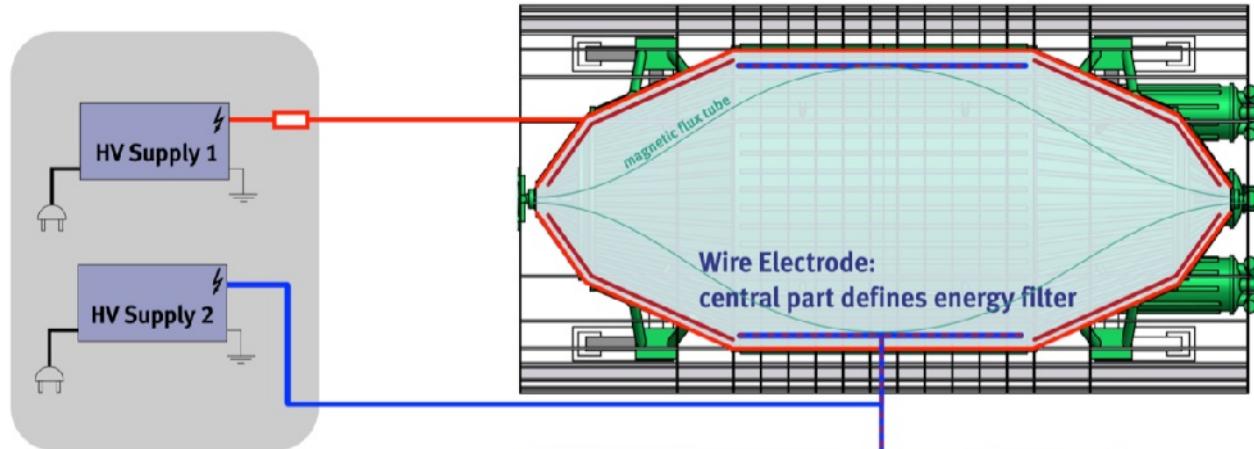


Electron energy losses by scattering in  $H_2$



The same, at different energies

# KATRIN's high voltage concept: stability and monitoring: ppm at -18.6 kV



D. Venos et al., *Appl. Rad. Iso.* 63 (2005) 323

D. Venos et al., *Nucl. Instrum. Meth. A* 560 (2006) 352

M. Rasulbaev et al., *Appl. Rad. Iso.* 66 (2008) 1838

Th. Thümmler et al., *New J. Phys.* 11 (2009) 103007

D. Venos et al., *Meas. Tech.* 53 (2010) 573

M. Slezák et al., *EPJ A48* (2012) 12

M. Zboril et al., *JINST* 8 (2013) P03009

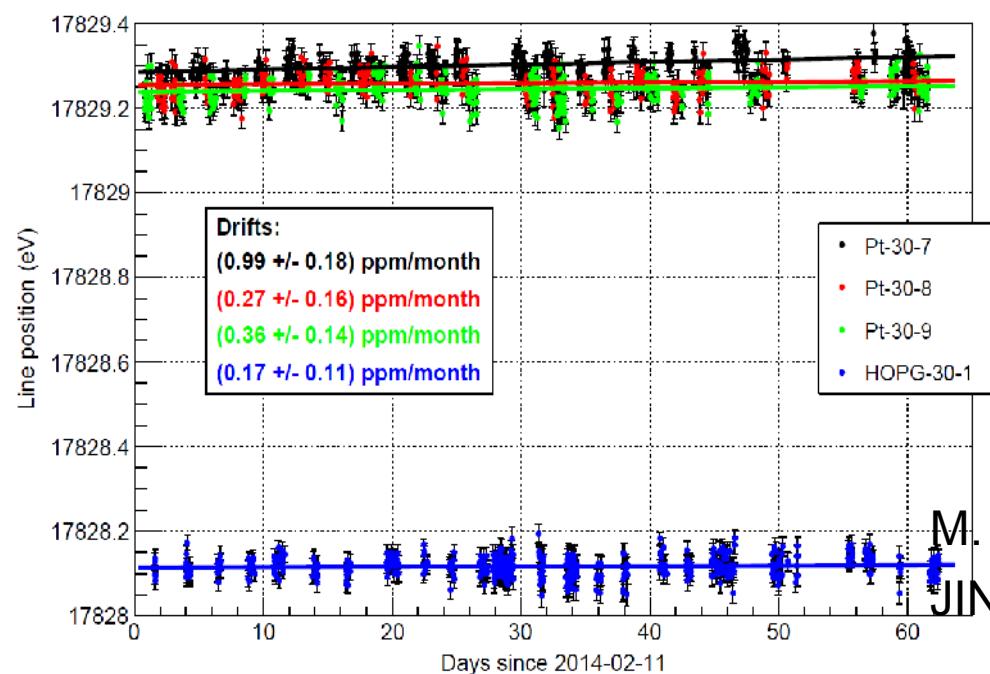
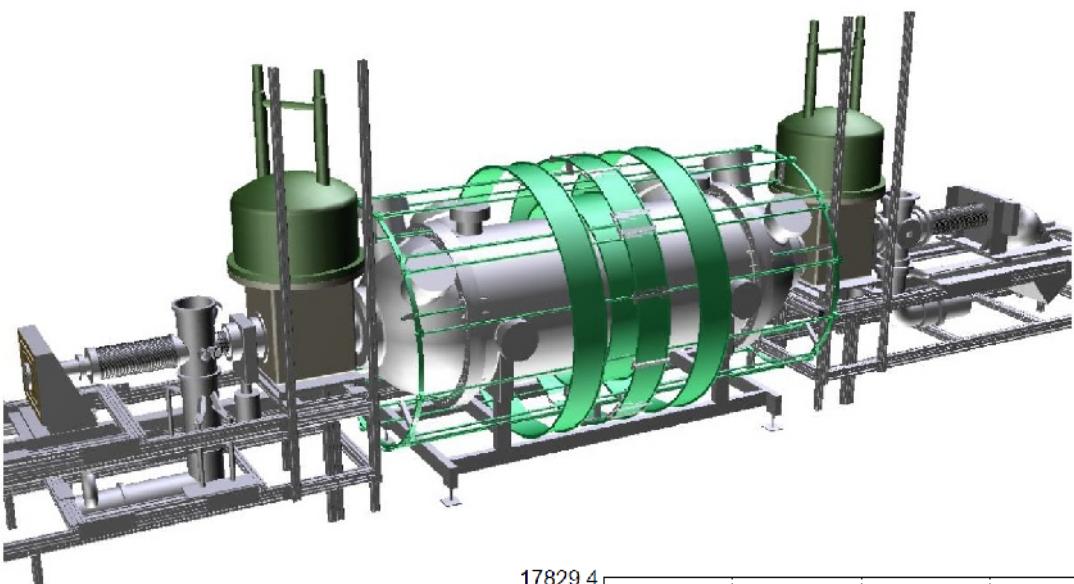
J. Bauer et al., *JINST* 8 (2013) P10026

M. Slezak et al., *JINST* 8 (2013) T12002

M. Erhard et al., *JINST* 9 (2014) P06022

courtesy: S. Bauer

# The former Mainz Spectrometer still contributes as monitor spectrometer for ppm HV precision



# Systematic uncertainties

As smaller  $m(\nu)$  as smaller the region of interest below endpoint  $E_0$   
→ quantum mechanical thresholds help a lot !

A few contributions with  $\Delta m_\nu^2 \leq 0.007 \text{ eV}^2$

1. inelastic scatterings of  $\beta$ 's inside WGTS
  - **dedicated e-gun measurements**, unfolding
2. fluctuations of WGTS column density (required  $\sim 10^{-3}$ )
  - rear detector, Laser-Raman spectroscopy,  $T=30\text{K}$  stabilisation,  
**e-gun measurements**
3. WGTS charging due to  $\beta$ 's
  - **monocrystalline readout**
4. final state distribution
  - **reliable quantum ch**
5. transmission function
  - detailed simulations
6. HV stability of retarding potential
  - **precision HV divider**

Measuring the last 25 or 30 eV only  
KATRIN becomes nearly  
a „single final state“ experiment  
as cryo-bolometers

tritium  
source

**sensitivity:**

$m_\nu < 0.2 \text{ eV (90\% CL)}$

**discovery potential:**

$m_\nu = 0.3 \text{ eV } (3\sigma)$

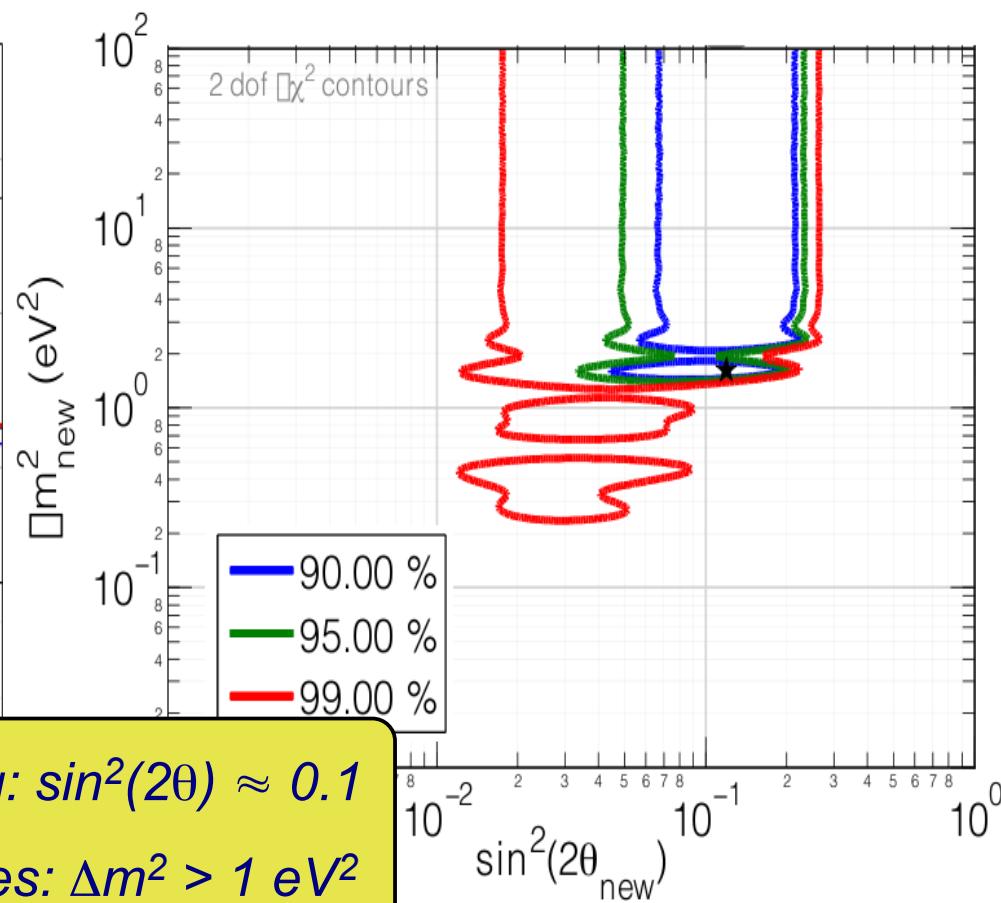
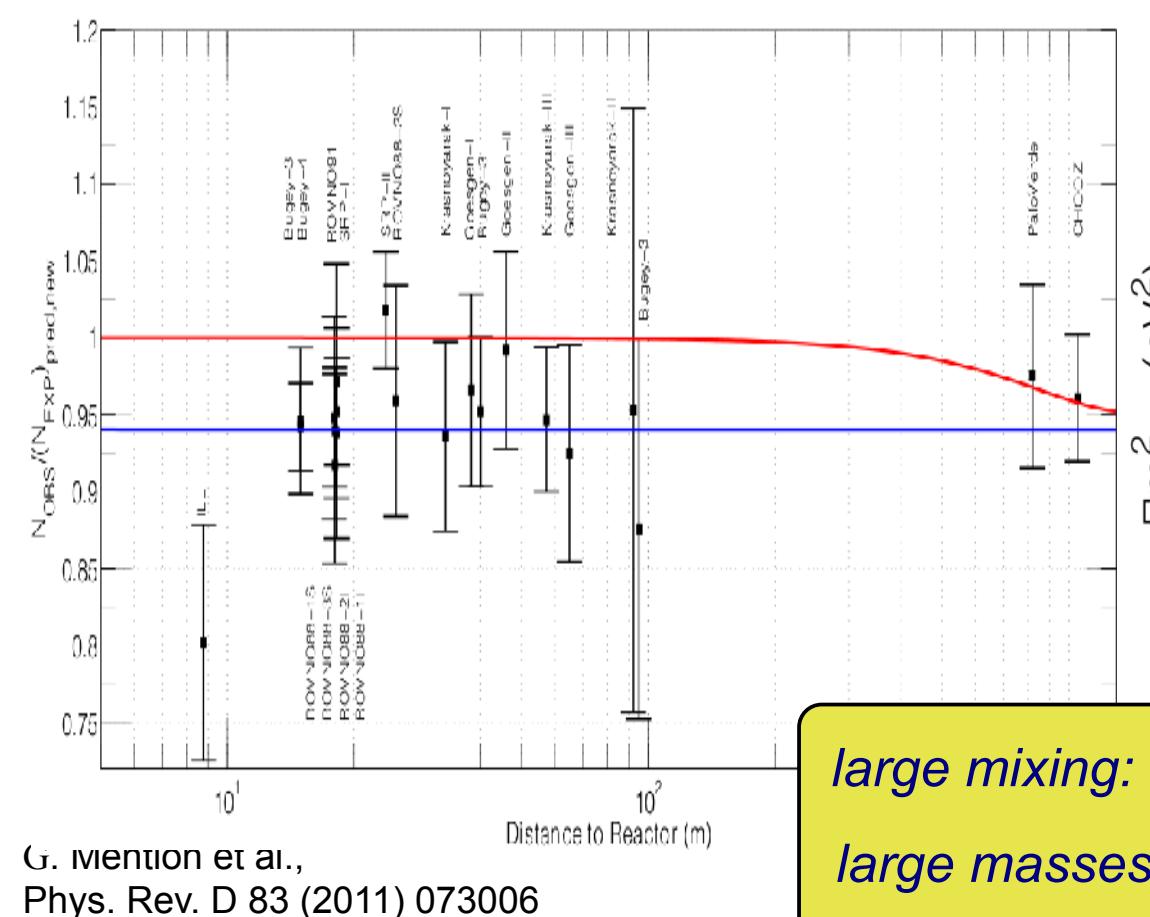
$m_\nu = 0.35 \text{ eV } (5\sigma)$

line

spectrometer

# Is there a fourth sterile neutrino state ?

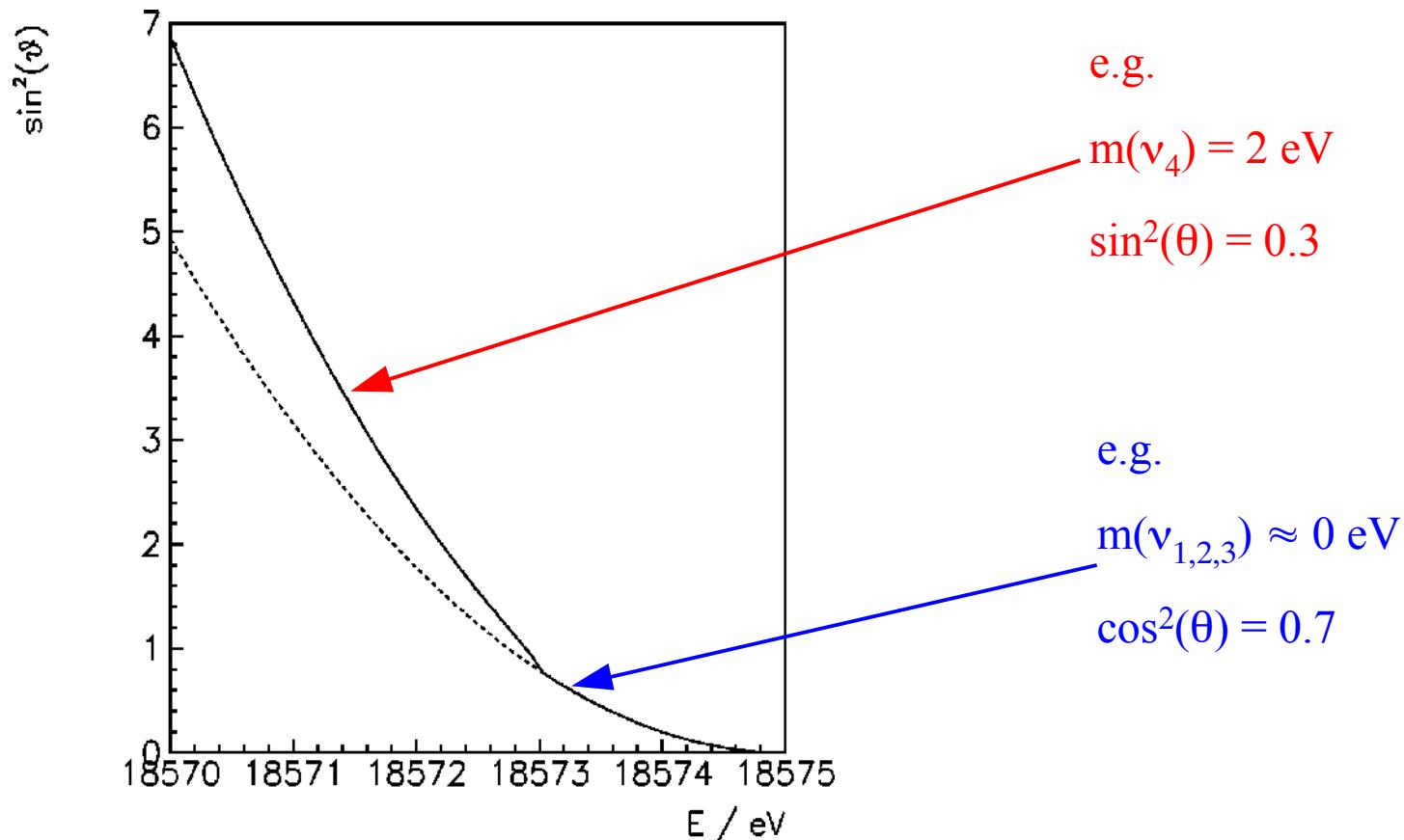
Re-evaluation of reactor neutrinos fluxes and use of  
GALLEX/SAGE calibration measurements:  
“reactor antineutrino anomaly”:  $P_{ee} = 0.943 \pm 0.023$



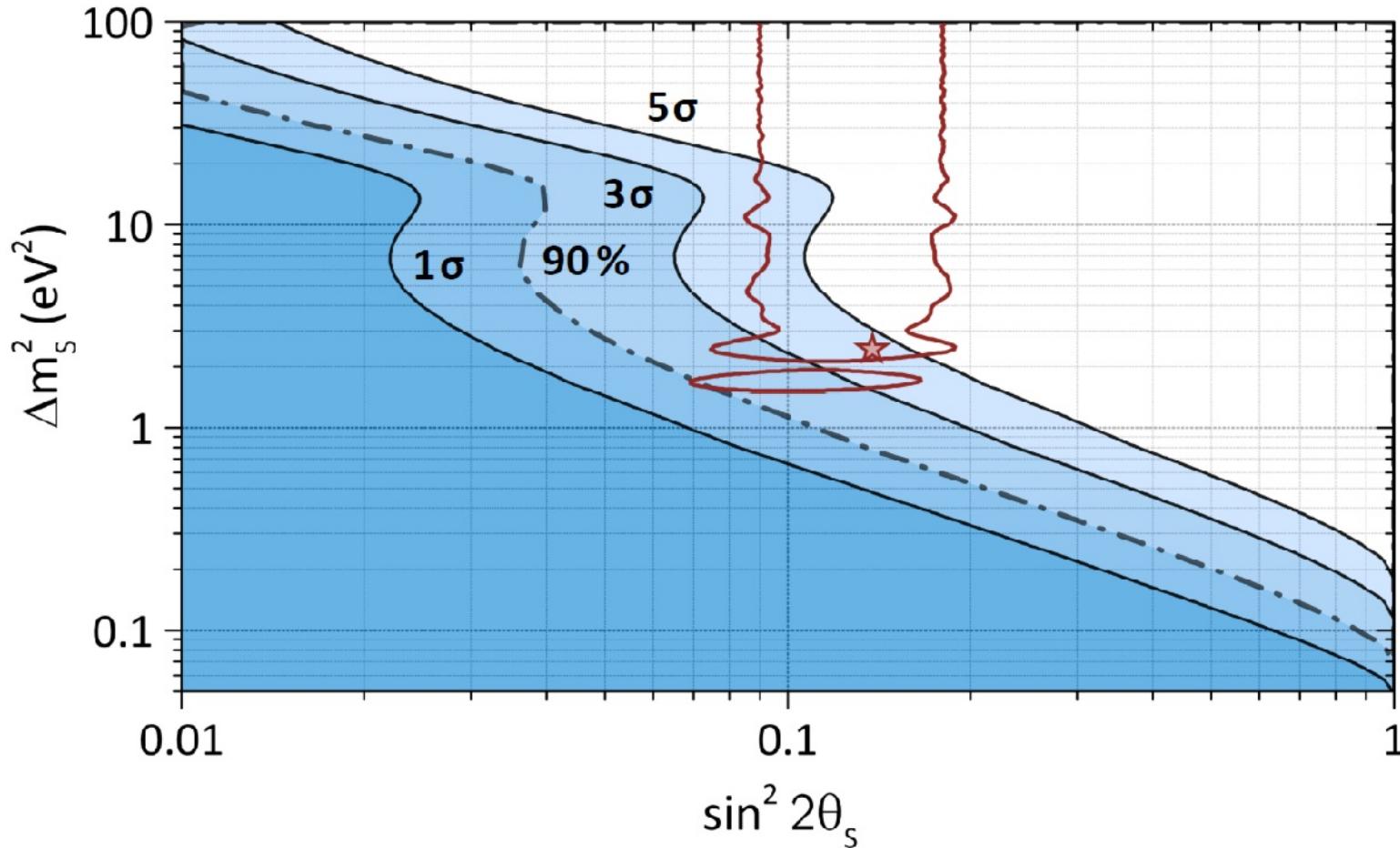
*large mixing:  $\sin^2(2\theta) \approx 0.1$*   
*large masses:  $\Delta m^2 > 1 \text{ eV}^2$*

# Influence of a 4<sup>th</sup> sterile neutrino near the endpoint $E_0$

$$dN/dE = K \cdot F(E, Z) \cdot p \cdot E_{\text{tot}} \cdot (E_0 - E_e) \left( \cos^2(\theta) \sqrt{(E_0 - E_e)^2 - m(v_{1,2,3})^2} + \sin^2(\theta) \sqrt{(E_0 - E_e)^2 - m(v_4)^2} \right)$$



# Sensitivity on sterile eV neutrinos



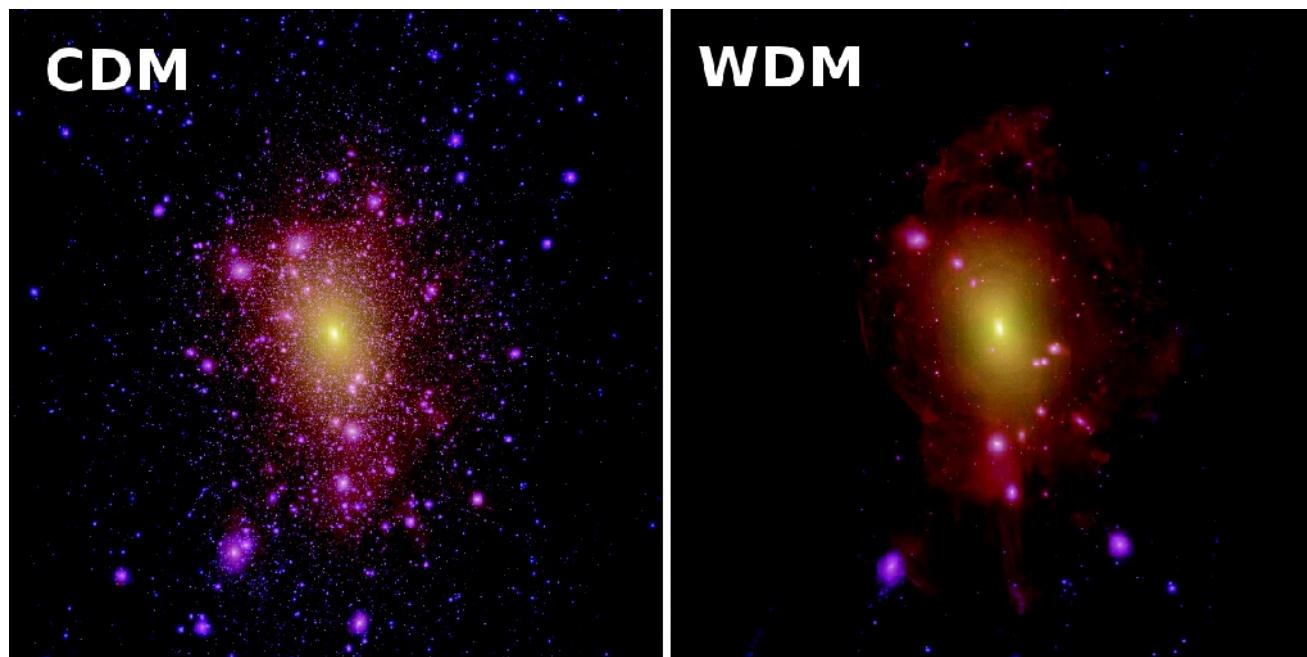
M.Kleesiek,  
PhD thesis,  
KIT (2014)

see also:

- J. A. Formaggio, J. Barret, PLB 706 (2011) 68
- A. Sejersen Riis, S. Hannestad, JCAP02 (2011) 011
- A. Esmaili, O.L.G. Peres, arXiv:1203.2632

# Hints for a 2<sup>nd</sup> sterile neutrino: Warm Dark Matter in the universe

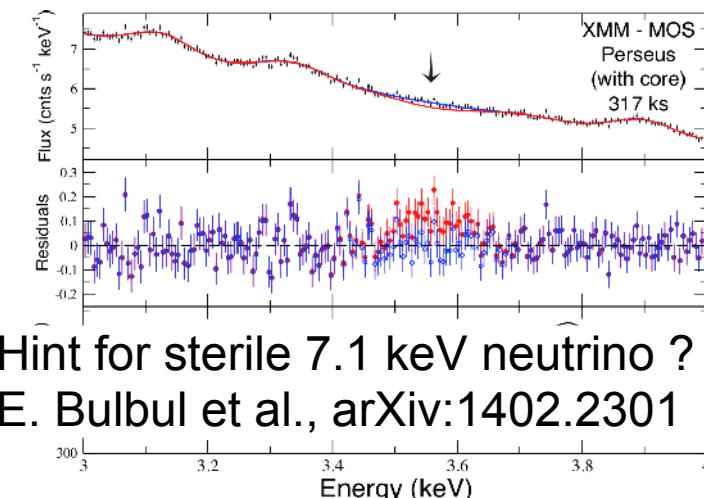
$\Lambda$ CDM (Cold Dark Matter with cosmological constant) models (masses of about 100 GeV) predict too much structure at galactic scales (too many satellite galaxies)



(e.g. Lovell et al. at Meudon Workshop 2012)

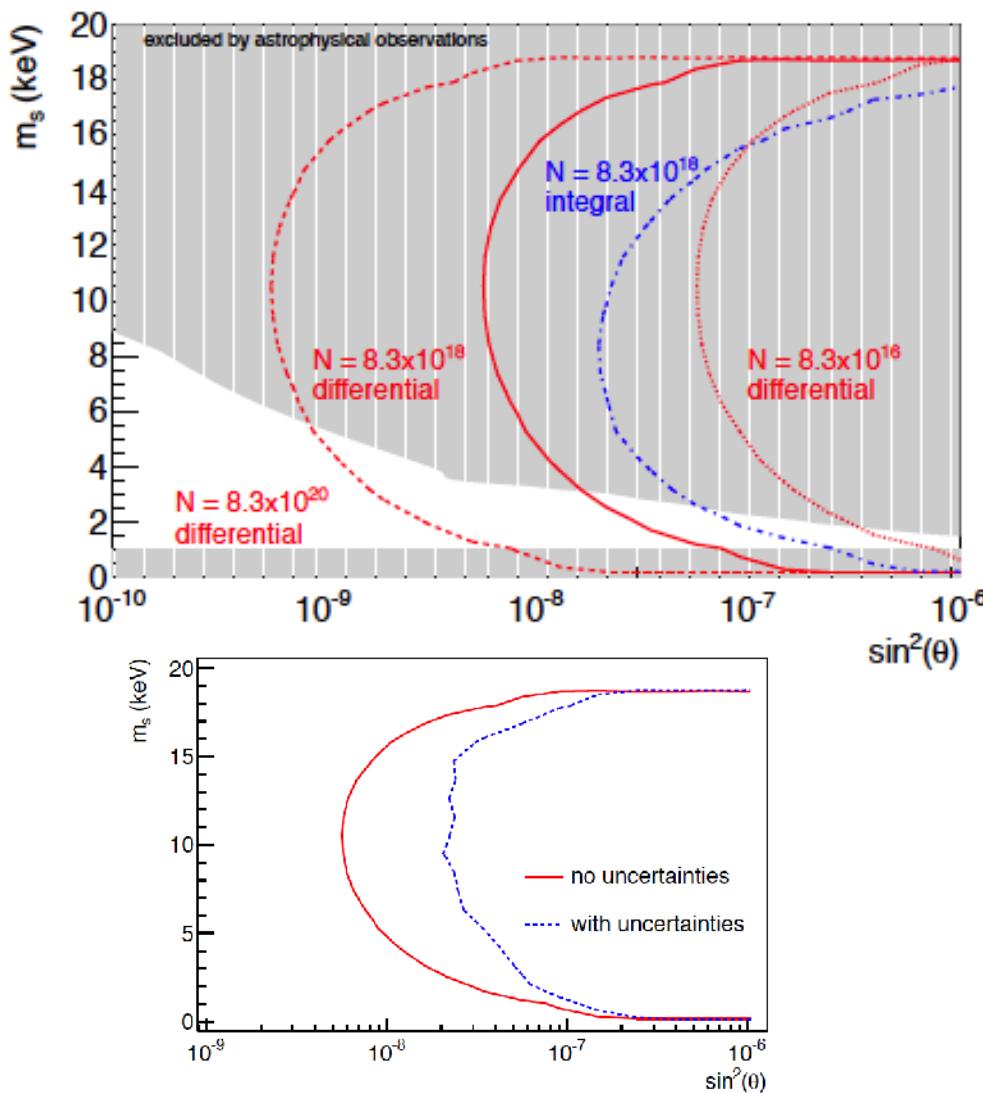
In contrast to observations ! (here only artist view on the right)

Warm Dark Matter (masses of a few keV, e.g. sterile neutrinos) would smear out these structures



[http://chandra.harvard.edu/graphics/resources/illustrations/milkyWay/milkyway\\_magellanic\\_clouds.jpg](http://chandra.harvard.edu/graphics/resources/illustrations/milkyWay/milkyway_magellanic_clouds.jpg)

# Statistical sensitivity for integral and differential measurement



S. Mertens et al., JCAP 02 (2015) 020  
 „Sensitivity of Next Generation Tritium  $\beta$ -Decay  
 Experiments for keV-Scale Sterile Neutrinos“

→ **statistical uncertainty  
 is not a problem for  $10^{-7}$   
 even including systematics  
 (1<sup>st</sup> investigation)!**

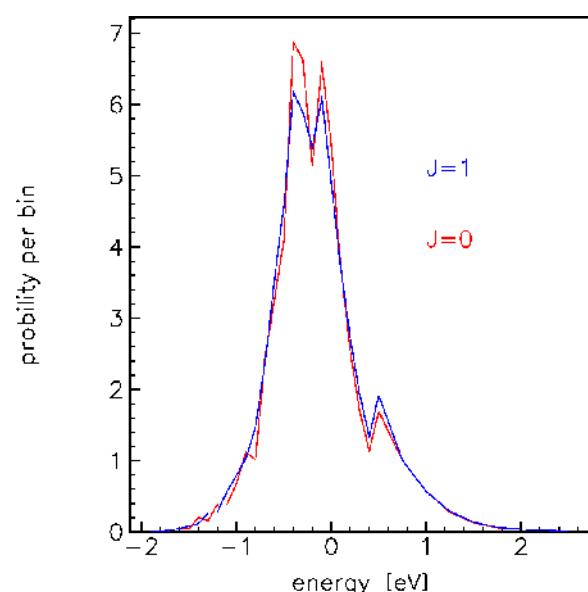
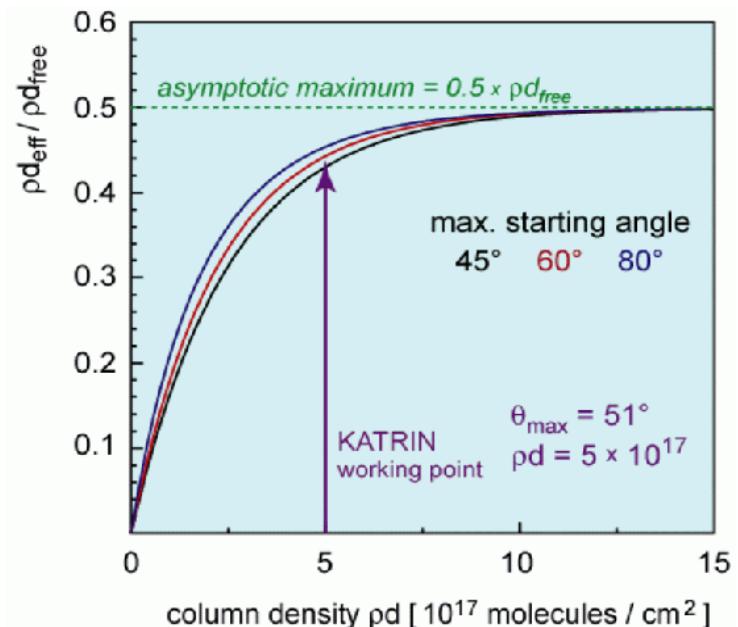
# Can KATRIN be largely improved ? Problems to be solved

- 1) The source is already opaque
  - need to increase size transversally
  - magnetic flux tube conservation
  - requests larger spectrometer too
  - but a Ø100m spectrometer is not feasible

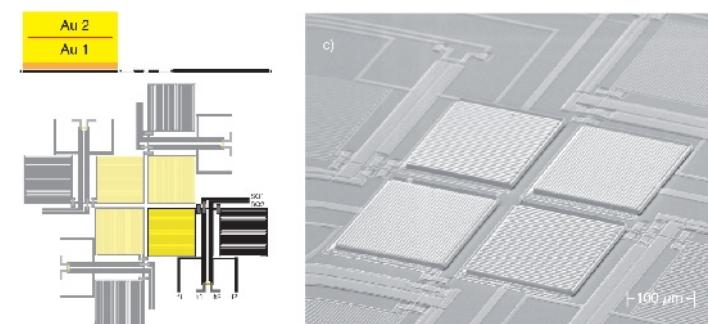
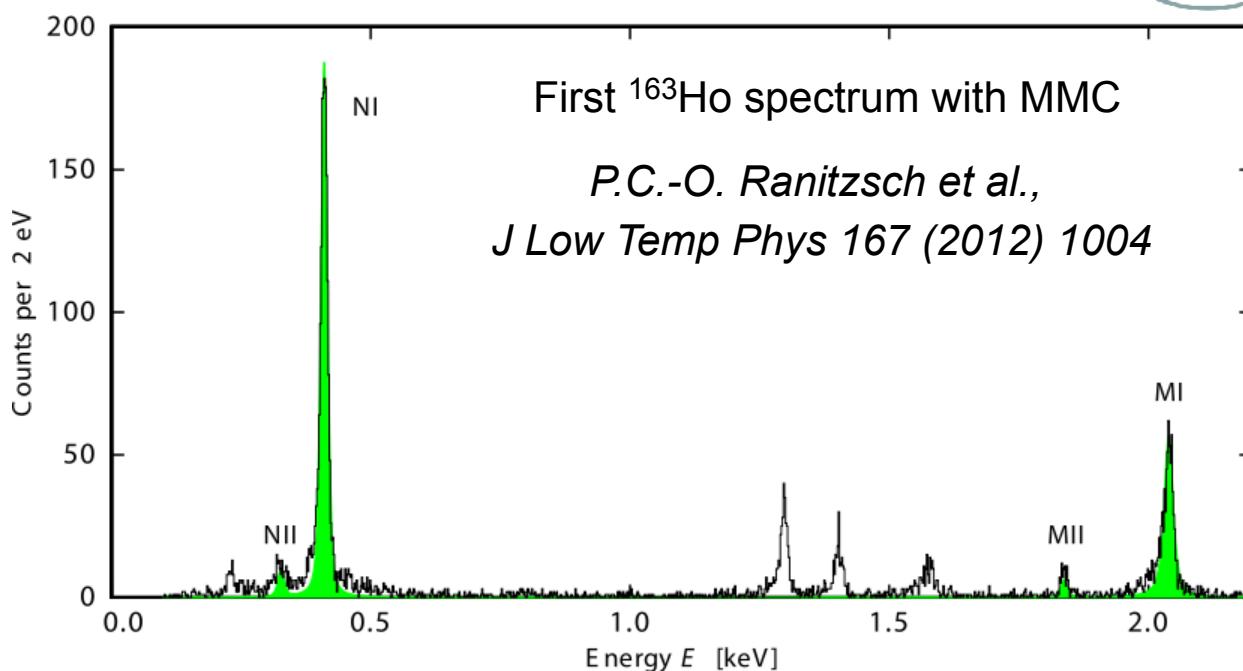
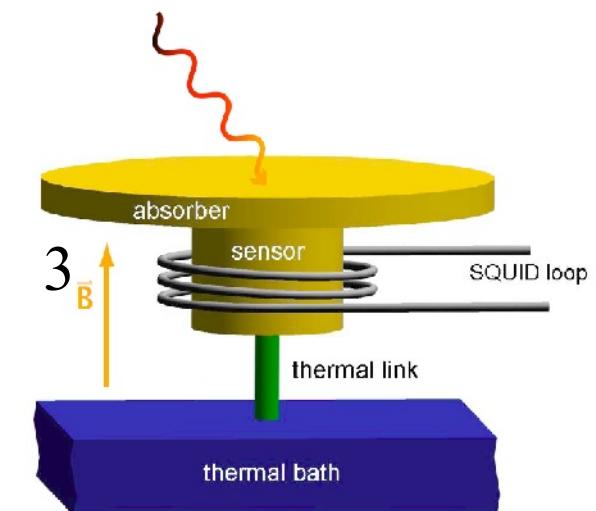
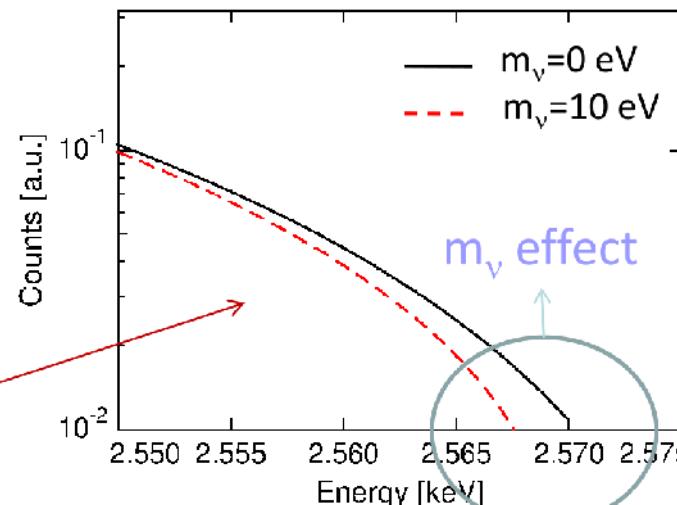
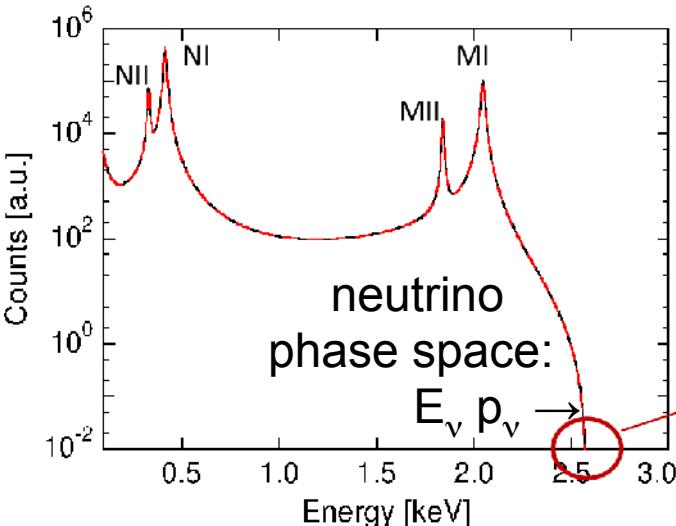
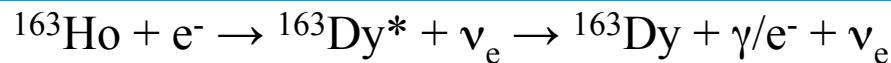
## Three possible ways out:

- a) source inside detector
  - using cryogenic bolometers (ECHO, HOLMES)
- b) hand-over energy information of  $\beta$  electron to other particle (radio photon), which can escape tritium source (Project 8)
- c) make better use of the electrons
  - time-of-flight spectroscopy

- 2) Resolution is limited to  $\sigma = 0.34$  eV when using molecular tritium by the excitation of ro-vibrational states in the final state



# ECHo neutrino mass project: $^{163}\text{Ho}$ electron capture with metallic magnetic calorimeters

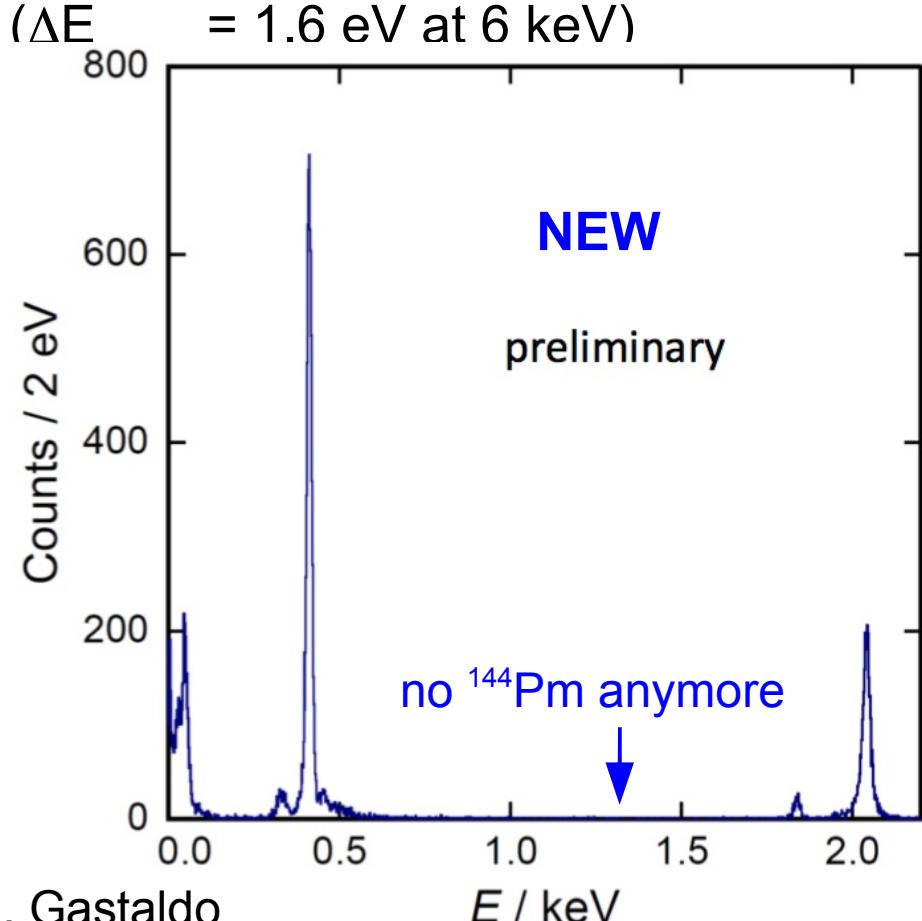
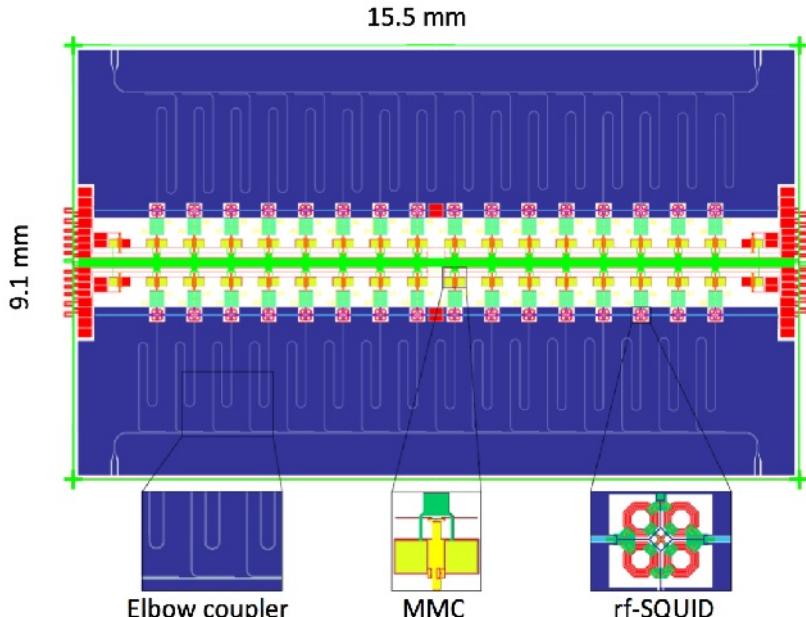


courtesy L. Gastaldo

# ECHO neutrino mass project: $^{163}\text{Ho}$ electron capture with metallic magnetic calorimeters

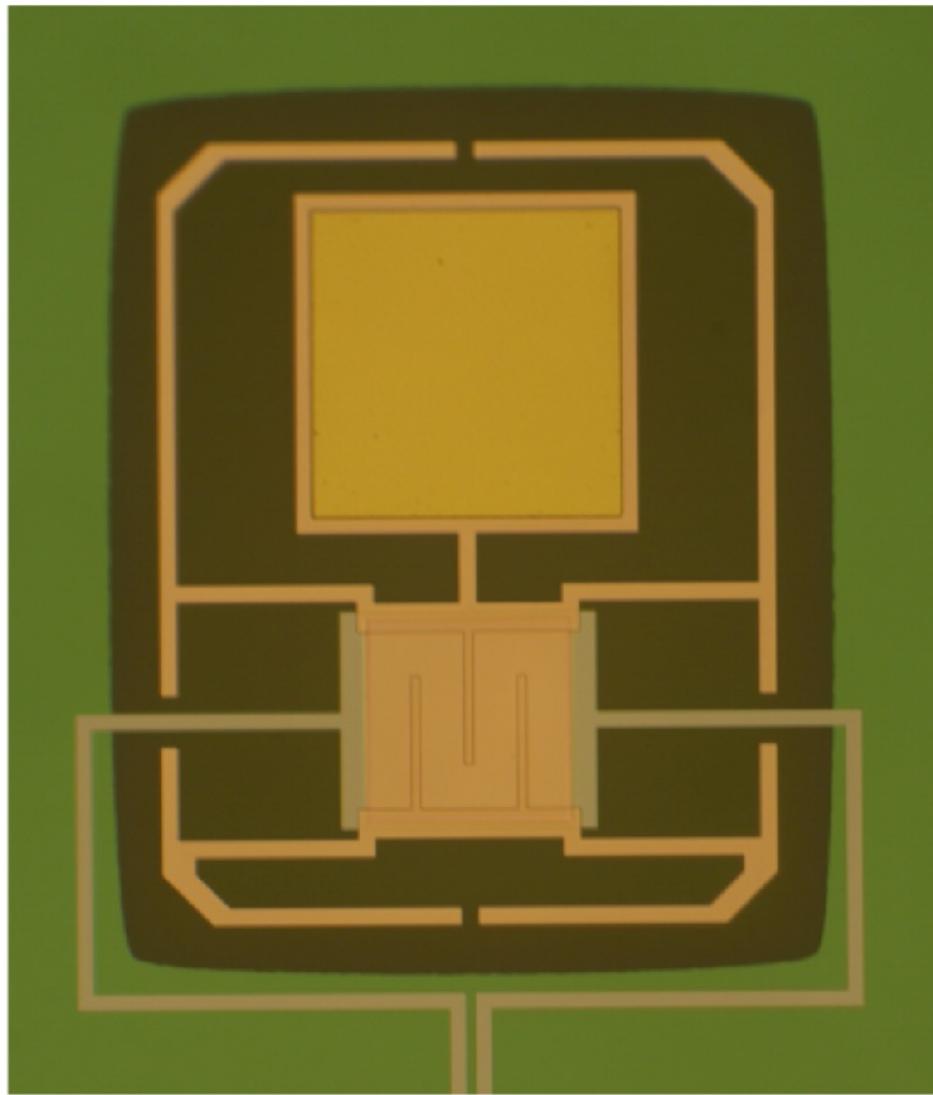
Recent achievements by ECHO:

- new Q-value: 2.8 keV (independently by MMC & Penning trap, was 2.5 keV before!)
- new source production: chemical purification + mass separation → no  $^{144}\text{Pm}$  or  $^{166\text{m}}\text{Ho}$
- very good energy resolution of this technology ( $\Delta E = 1.6 \text{ eV at } 6 \text{ keV}$ )
- ultra-short response (pile-up!): risetime 90 ns
- 128 pixels: microwave SQUID multiplexing
- fl



courtesy L. Gastaldo

# HOLMES: cryogenic calorimeter with transition edge sensor (TES) read-out

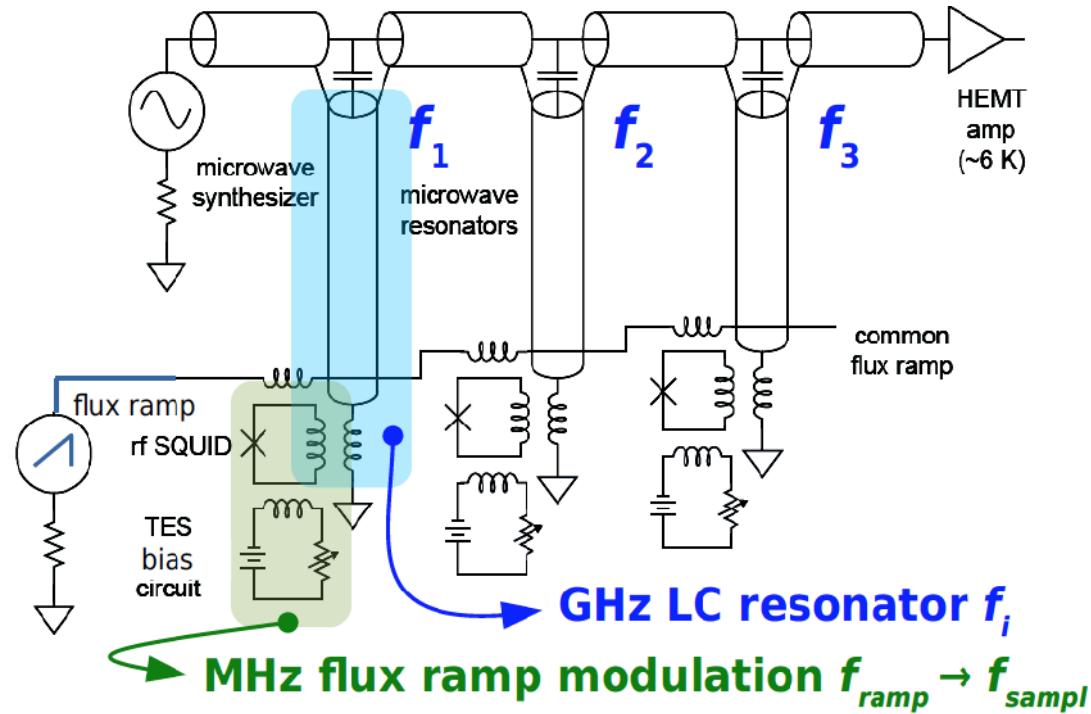


funding by ERC grant  
courtesy A. Nucciotti

Prediction from test measurements:

$$\Delta E_{FWHM} \approx 3 \text{ eV}, \tau_{\text{rise}} \approx 6 \mu\text{s}, \tau_{\text{decay}} \approx 130 \mu\text{s}$$

radiofrequency SQUID multiplexing



# HOLMES: cryogenic calorimeter with transition edge sensor (TES) read-out

Prediction from test measurements:

$$\Delta E_{FWHM} \approx 3 \text{ eV}, \tau_{\text{rise}} \approx 6 \mu\text{s}, \tau_{\text{decay}} \approx 130 \mu\text{s}$$

$^{163}\text{Ho}$  EC is investigated by 3 collaborations

Cryo-calorimetric multipixel detectors  
are a very interesting technology  
→ starts to become scalable

But many orders of magnitude to go for  
required statistics and background !

Systematics and show stoppers on the way?

→ We should stay tunned !

funding by ERC grant

courtesy A. Nucciotti



# Project 8's goal: Measure coherent cyclotron radiation of tritium $\beta$ electrons

General idea:

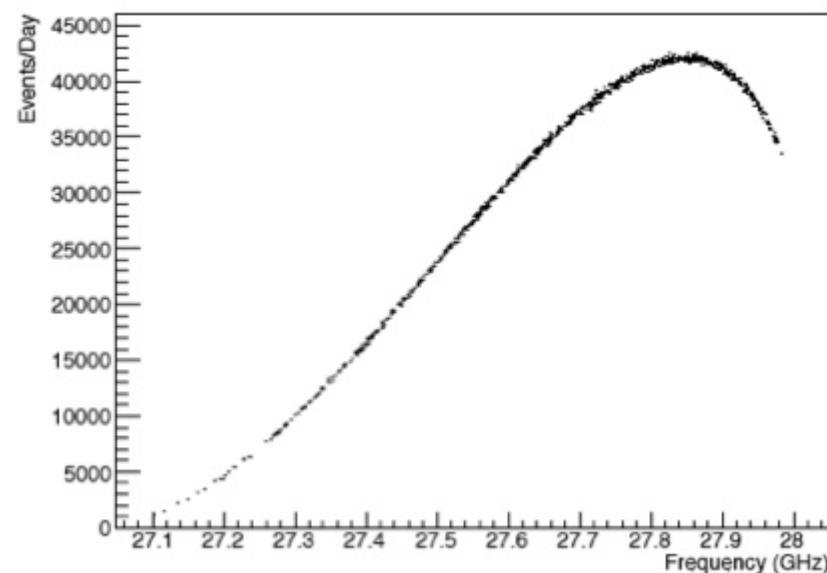
B. Montreal and J. Formaggio, PRD 80 (2009) 051301

- Source = KATRIN tritium source technology :

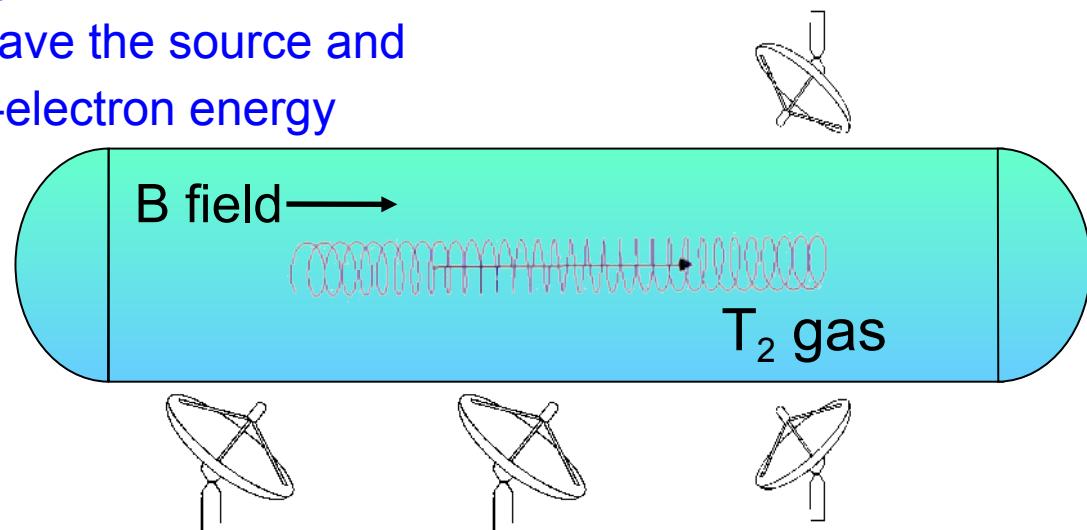
uniform B field + low pressure  $T_2$  gas

**$\beta$  electron radiates coherent cyclotron radiation**

$$\omega(\gamma) = \frac{\omega_0}{\gamma} = \frac{eB}{K + m_e}$$

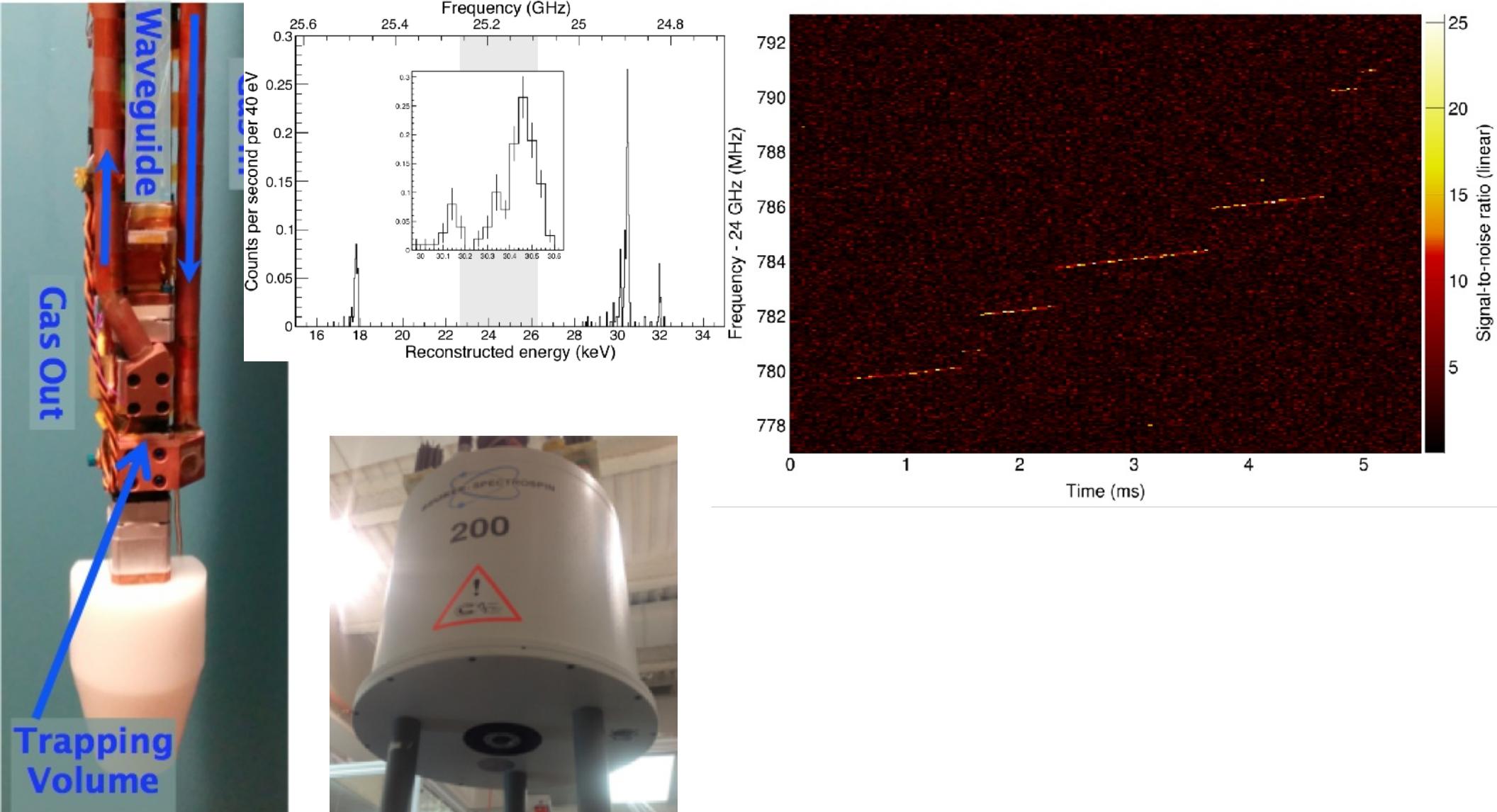


- Antenna array (interferometry) for cyclotron radiation detection since cyclotron radiation can leave the source and carries the information of the  $\beta$ -electron energy



# Project 8's phase 1: detection single electrons from $^{83m}\text{Kr}$

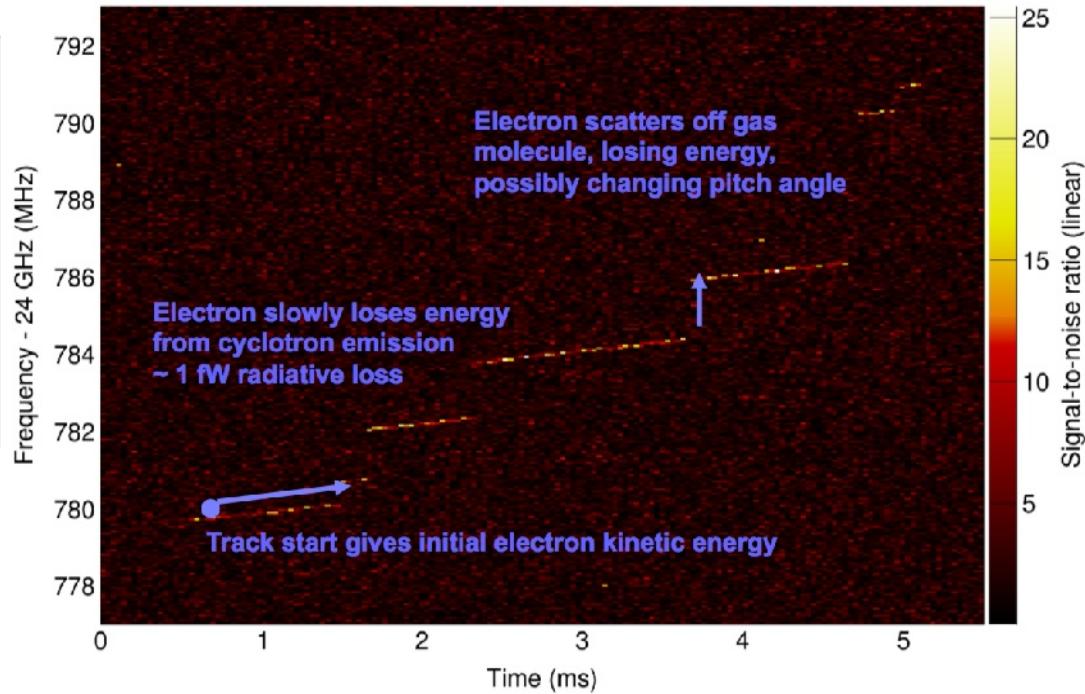
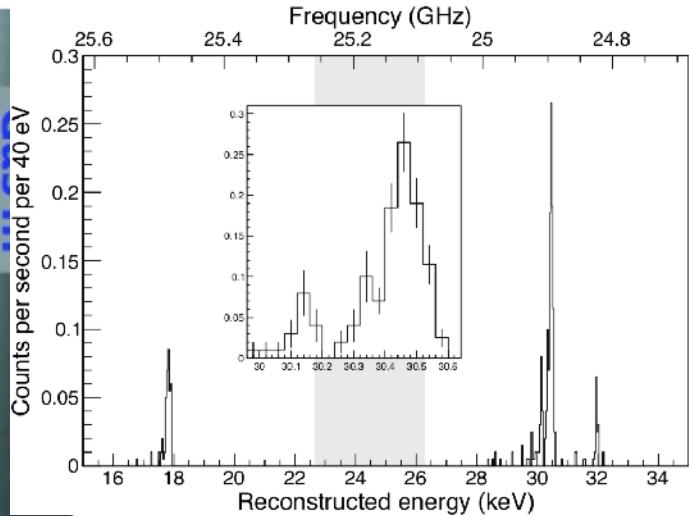
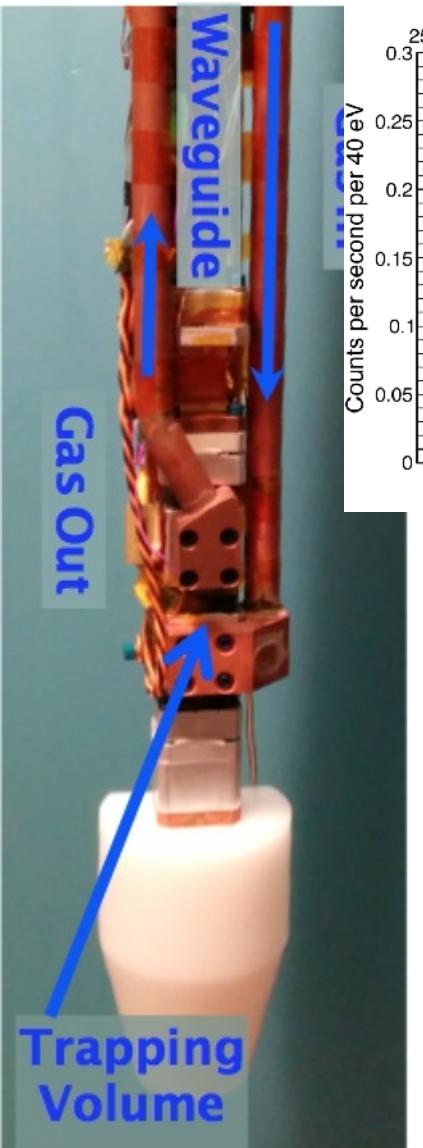
D. M. Asner et al., Phys. Rev. Lett. 114, 162501



courtesy J. Formaggio, RGH Robertson

# Project 8's phase 1: Detection single electrons from $^{83m}\text{Kr}$

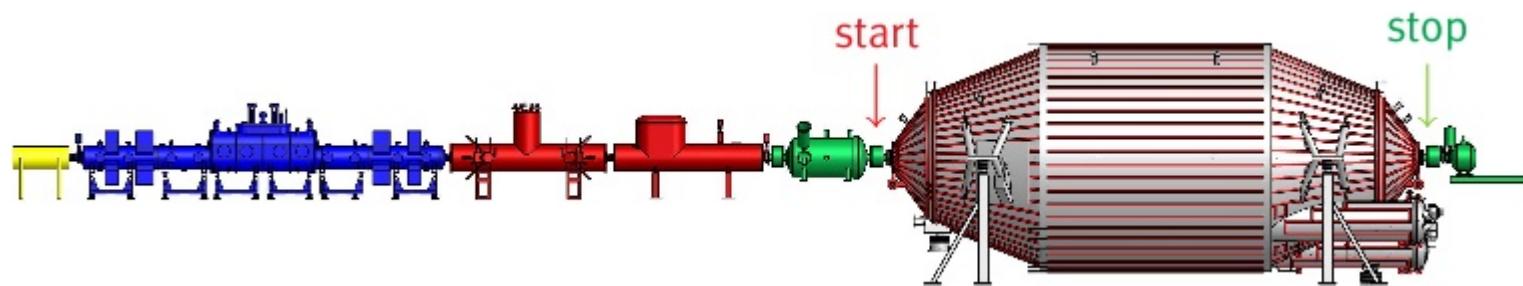
D. M. Asner et al., Phys. Rev. Lett. 114, 162501



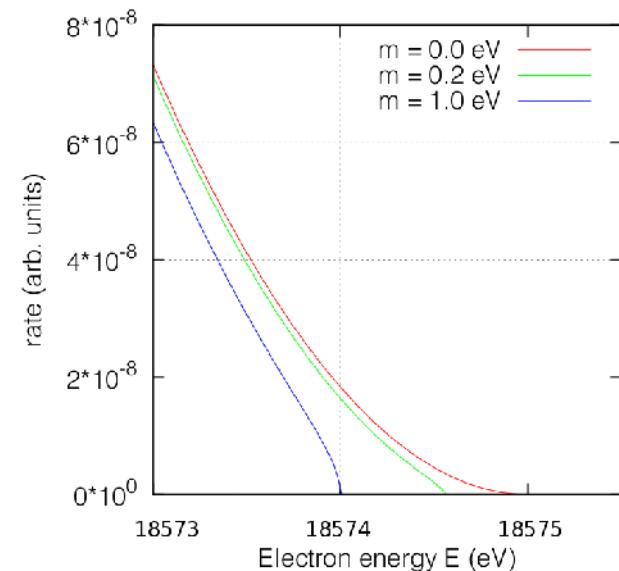
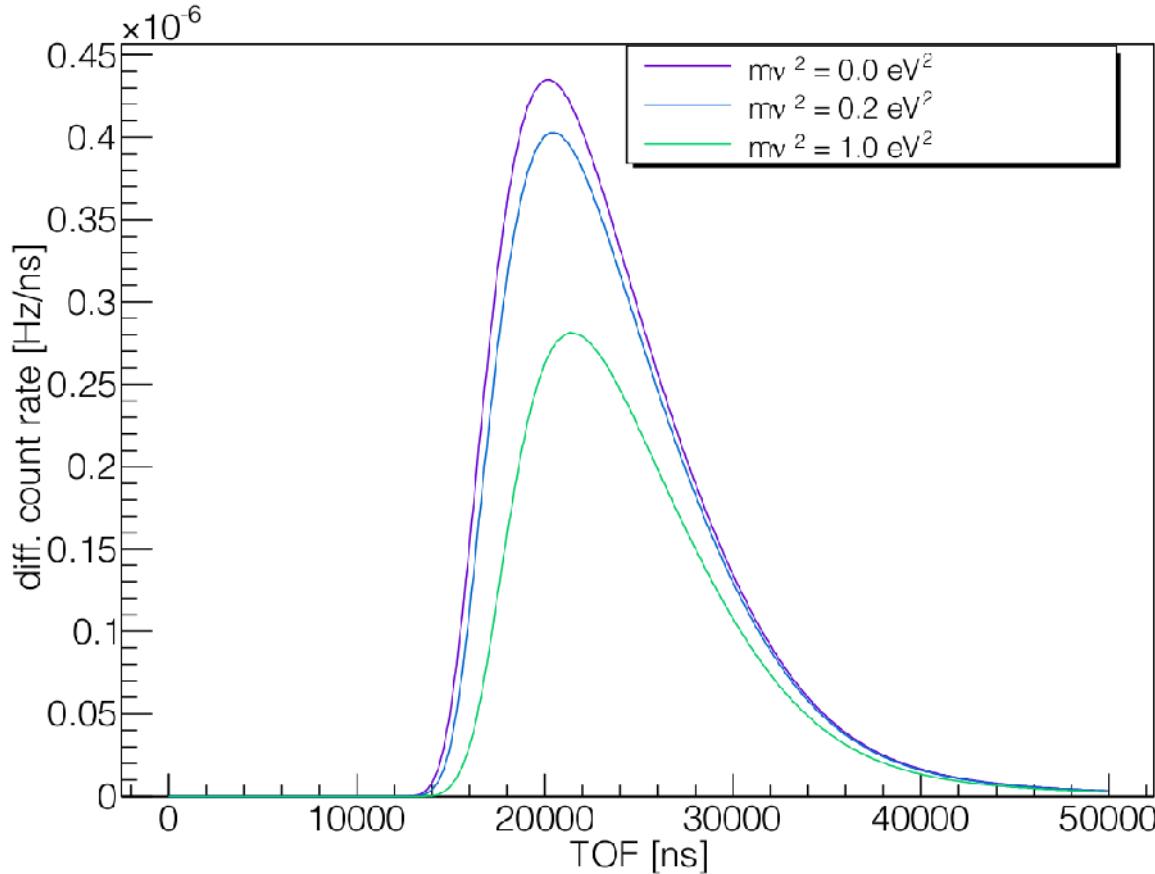
First detection of single electrons successfull  
but still a lot of R&D necessary  
- Is a large scale experiment possible ?  
- What are the systematic uncertainties  
& other limitations?

courtesy J. Formaggio, RGH Robertson

# Alternative spectroscopy: measure time-of-flight TOF through KATRIN spectrometer

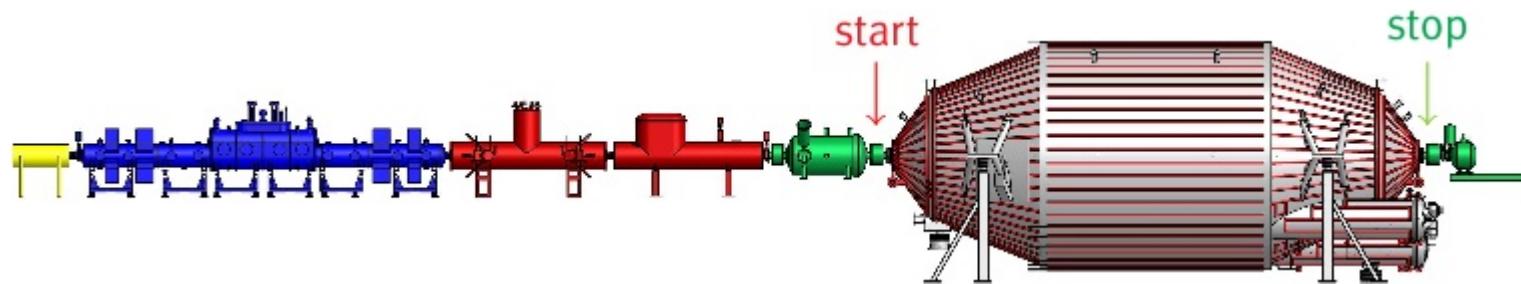


Comparison of TOF spectra for different neutrino masses for  $E_0 = 18571.0 \text{ eV}$ ,  $U_{\text{rel}} = -18570.0 \text{ eV}$

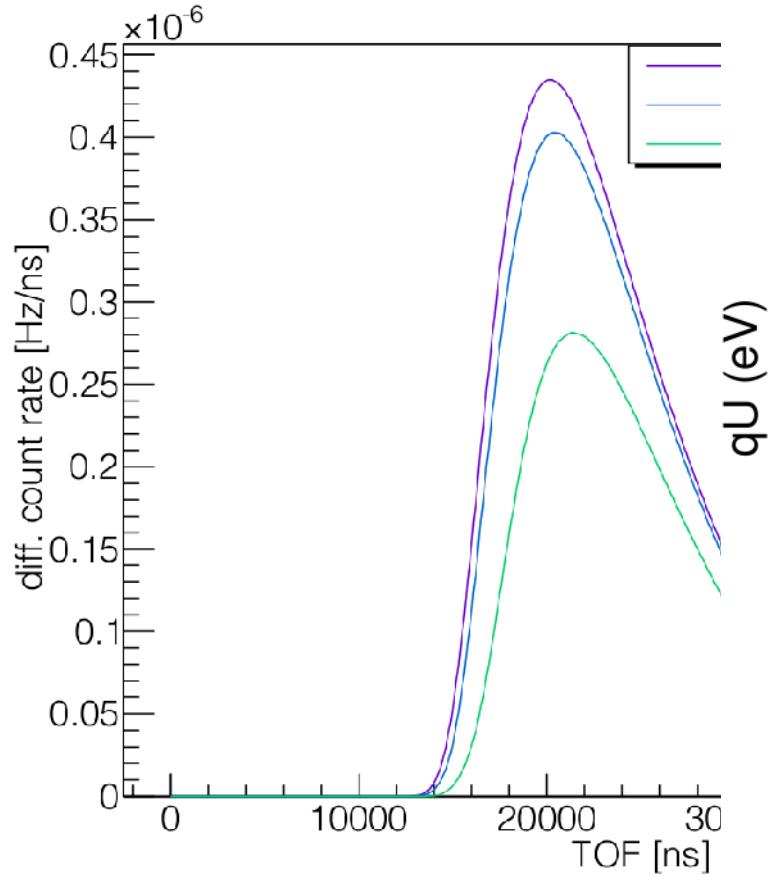


Time-of-flight spectrum is sensitive to the neutrino mass  
requires one retardation potential only  
not integral but differential  $\beta$ -spectrum

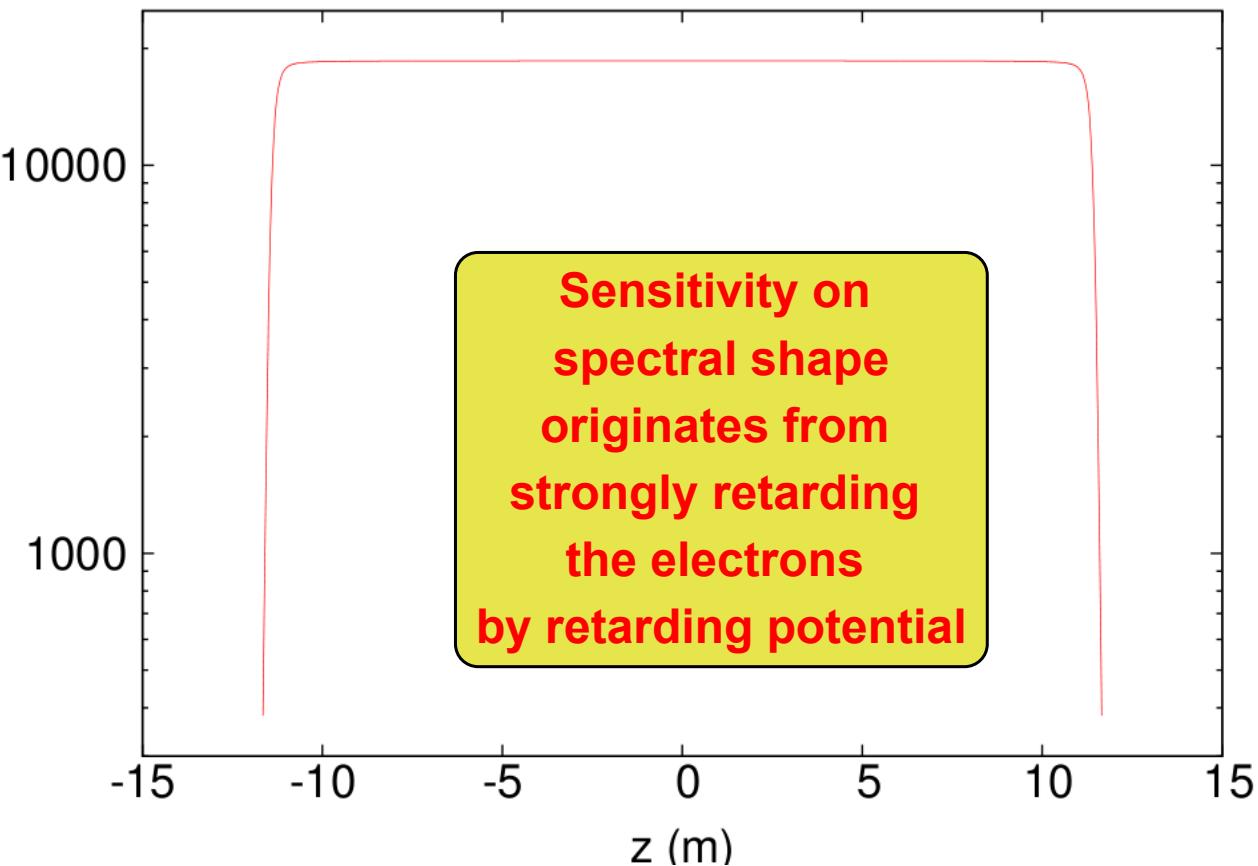
# Alternative spectroscopy: measure time-of-flight TOF through KATRIN spectrometer



Comparison of TOF spectra for different neutrino masses for



Electric potential on main spectrometer z axis

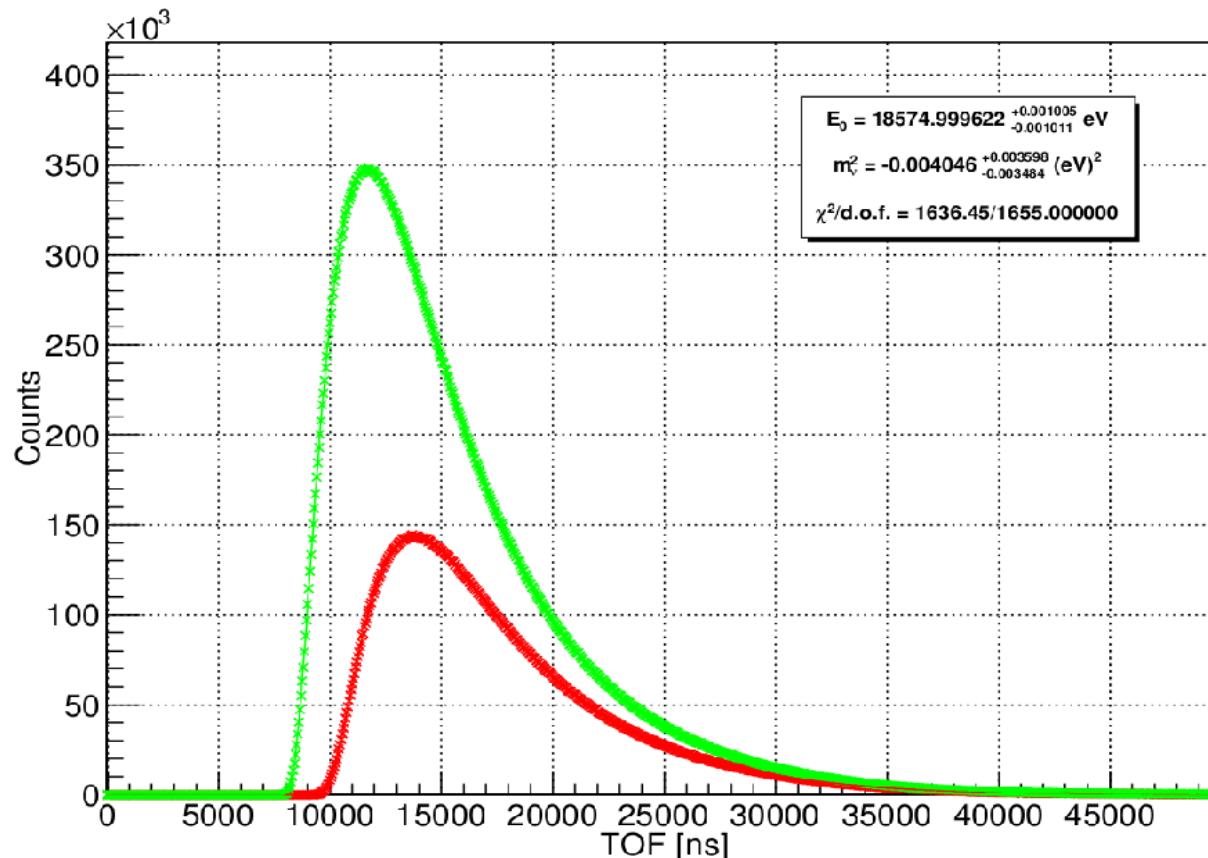


# Sensitivity improvement on $m^2(\nu_e)$ by ideal TOF determination

Measure at 2 (instead of  $\approx 30$ ) different retarding potentials  
since TOF spectra contain all the information

Coincidence request between start and stop signal → nice background suppression

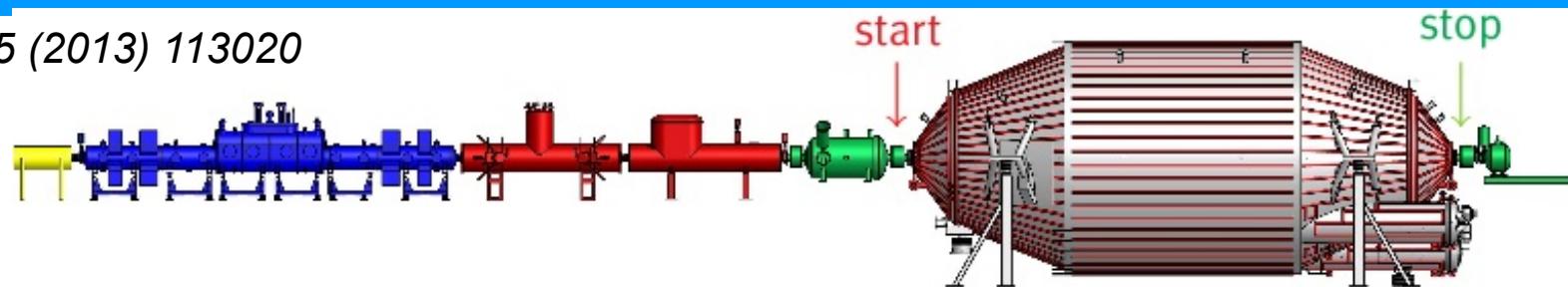
→ Factor 5 improvement in  $m_\nu^2$  w.r.t. standard KATRIN, but ideal case !



N. Steinbrink et al.  
NJP 15 (2013) 113020

# Alternative spectroscopy: measure time-of-flight through KATRIN spectrometer

N. Steinbrink et al., NJP 15 (2013) 113020



Advantage: measure full  $\beta$ -spectrum by time-of-flight at one (a few) retarding potential

Stop: Can measure time-of-arrival with KATRIN detector with  $\Delta t = 50$  ns  $\rightarrow$  ok

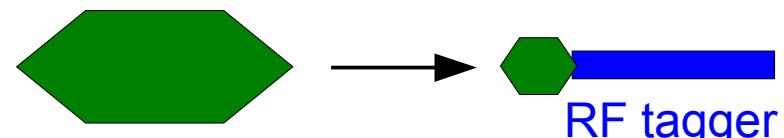
Start: **e<sup>-</sup>-tagger:** Need to determine time-of-passing-by of e<sup>-</sup> before main spectrometer without disturbing energy and momentum by more than 10 meV:

$\rightarrow$  Need „detector“ with 10 meV threshold

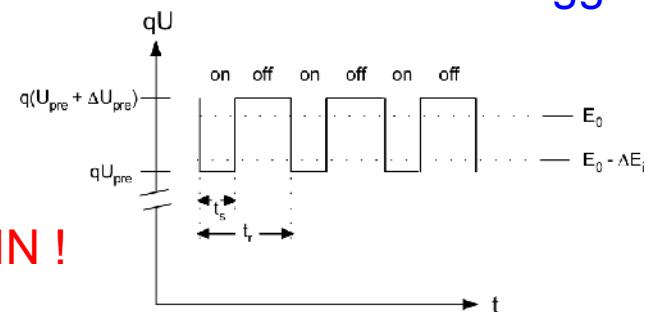
seems not to be forbidden but very difficult for the near future !

Added value: significant background reduction by coincidence !

An implementation: reduce pre spectrometer and add a Project 8-type tagger within a long solenoid



or: Use pre spectrometer as a „gated-filter“ by switching fast the retarding voltage  
 $\rightarrow$  As sensitive on the neutrino mass as standard KATRIN !



# Conclusions

KATRIN is the direct neutrino mass experiment complementary to cosmological analyses and  $0\nu\beta\beta$  searches

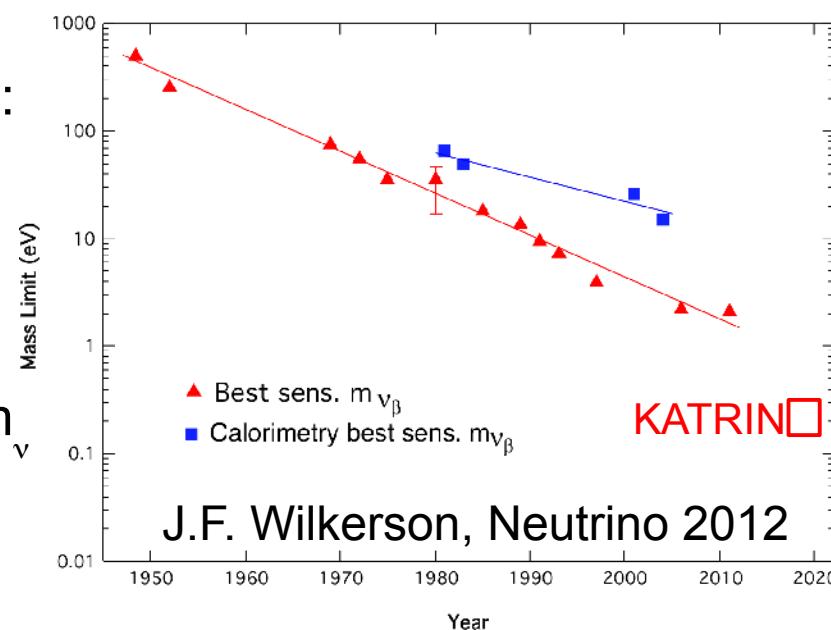
KATRIN can also look for sterile neutrinos (eV, keV)

KATRIN is close to start direct neutrino mass measurements

- Commissioning of spectrometer & detector SDS IIb finished in August 2015
- Commissioning of tritium source & transport section: up to summer 2016
- Tritium data taking: start in 2016 with small column density
- early 2017 operate with normal column density →  $m_\nu$

KATRIN's sensitivity: 200 meV

significant R&D on  $^{163}\text{Ho}$  micro calorimeters (ECHO, HOLMES, ...) new ideas like Project 8, ...



THANK YOU FOR YOUR ATTENTION !