

Status and prospects of inclusive $b \rightarrow uv$ theory

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Why precision CKM studies?

- The SM accomodates flavour & CP violation, but **we have no theory of flavour**
- We expect New Physics at the EW scale, and most models predict additional flavour and CP violation.
- The CKM mechanism is very successful \Rightarrow **flavour and CP problem** (NP must preserve agreement with data)
- To uncover small signals of physics beyond CKM, we need precision tests, in many ways a challenge for our QCD understanding

Outline

- Inclusive decays, inputs, pert corrections
- Shape Function(s)
- High q^2 tail and WA
- Existing approaches
- A trial comparison
- Conclusions

Inclusive semileptonic B decays: basic features

- **Simple idea:** inclusive decay do not depend on final state, factorize long distance dynamics of the meson. OPE allows to express it in terms of matrix elements of local operators
- The Wilson coefficients are perturbative, matrix elements of local ops parameterize non-pert physics: **double series in $\alpha_s, \Lambda/m_b$**
- Lowest order: decay of a free b , linear Λ/m_b absent. Depends on $m_{b,c}$, 2 parameters at $O(1/m_b^2)$, 2 more at $O(1/m_b^3)$...

$$\mu_\pi^2(\mu) = \frac{1}{2M_B} \left\langle B \left| \bar{b} (i\vec{D})^2 b \right| B \right\rangle_\mu \quad \mu_G^2(\mu) = \frac{1}{2M_B} \left\langle B \left| \bar{b} \frac{i}{2} \sigma_{\mu\nu} G^{\mu\nu} b \right| B \right\rangle_\mu$$

The total s.l. width in the OPE

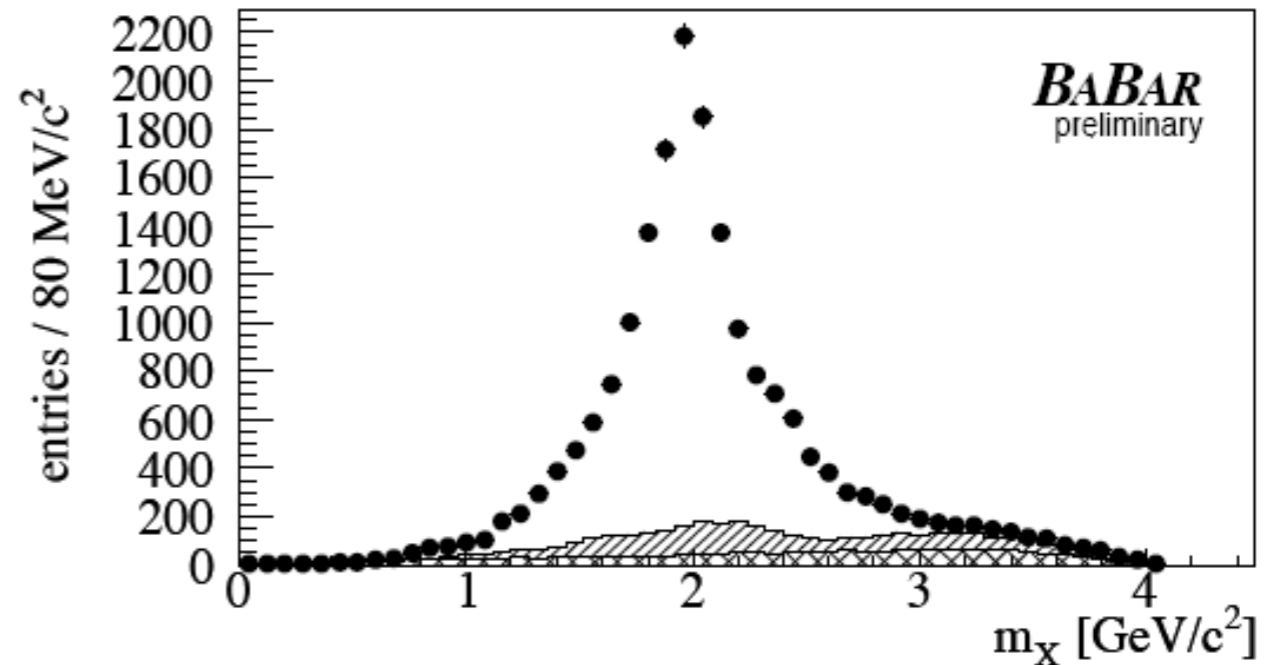
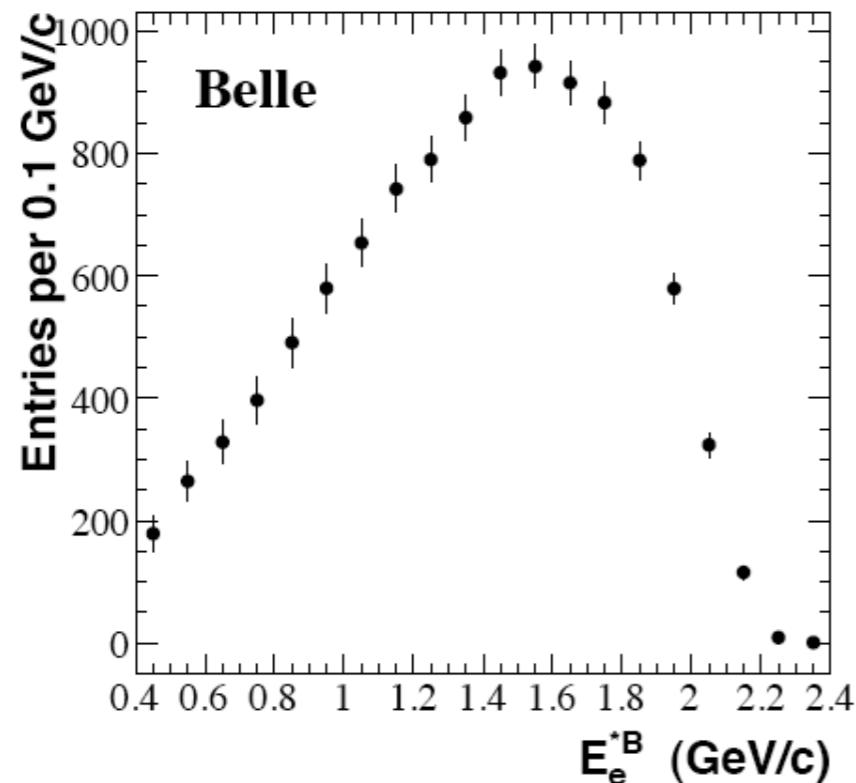
$$\Gamma[\bar{B} \rightarrow X_u e \bar{\nu}] = \frac{G_F^2 m_b^5}{192\pi^3} |V_{ub}|^2 \left[1 + \frac{\alpha_s}{\pi} p_u^{(1)}(\mu) + \frac{\alpha_s^2}{\pi^2} p_u^{(2)}(r, \mu) - \frac{\mu_\pi^2}{2m_b^2} - \frac{3\mu_G^2}{2m_b^2} \right. \\ \left. + \left(\frac{77}{6} + 8 \ln \frac{\mu_{\text{WA}}^2}{m_b^2} \right) \frac{\rho_D^3}{m_b^3} + \frac{3\rho_{LS}^3}{2m_b^3} + \frac{32\pi^2}{m_b^3} B_{\text{WA}}(\mu_{\text{WA}}) \right] \\ + O(\alpha_s/m_b^2 \dots)$$

OPE valid for inclusive enough measurements, away from perturbative singularities.

m_b dependence is up to twice stronger in the cut rate

Most OPE parameters (quark masses etc) from sl decays into charm

Fitting OPE parameters to the moments



Total **rate** gives $|V_{cb}|$, global **shape** parameters (moments of the distributions) tell us about B structure

OPE parameters describe universal properties of the B meson and of the quarks

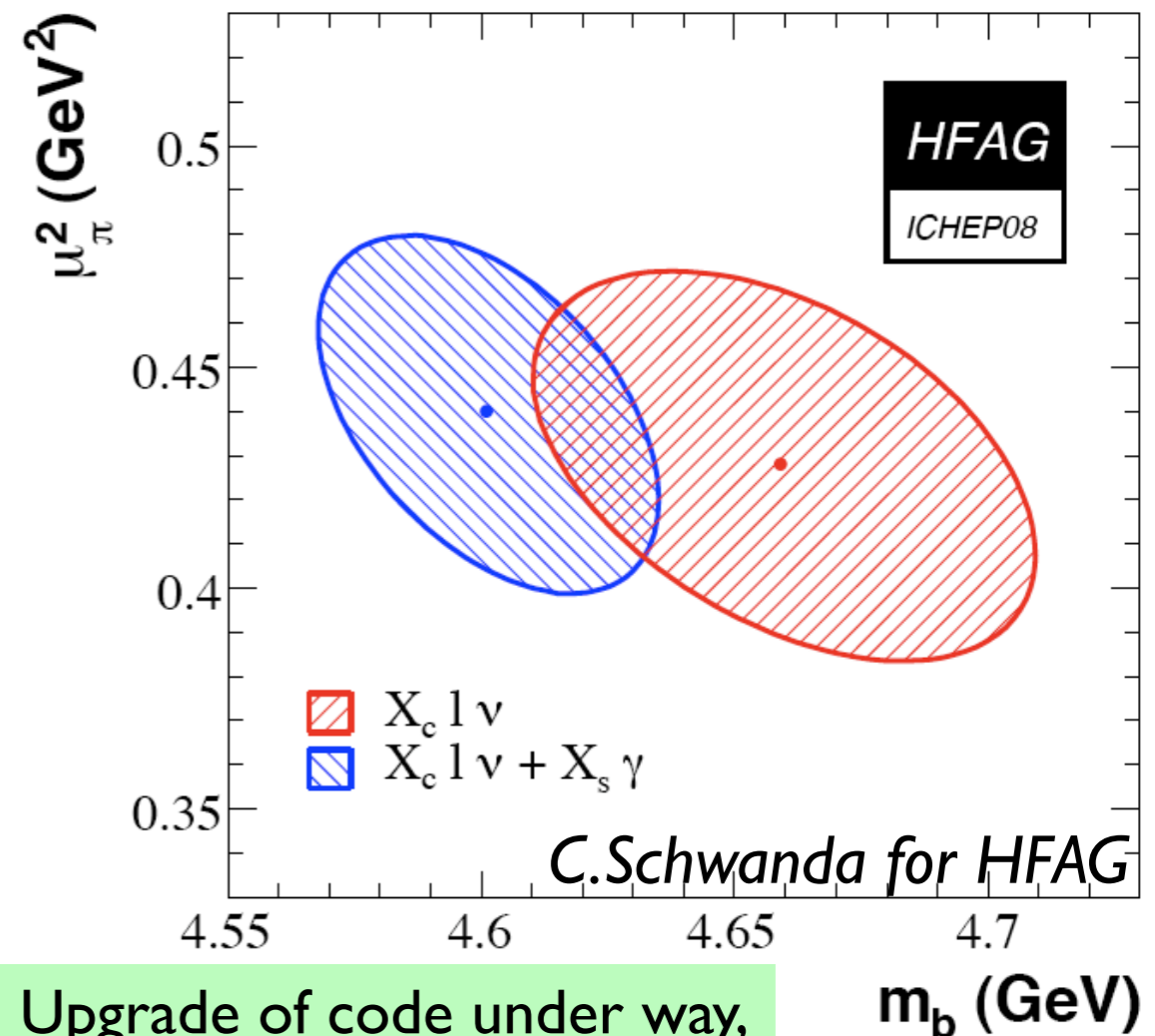
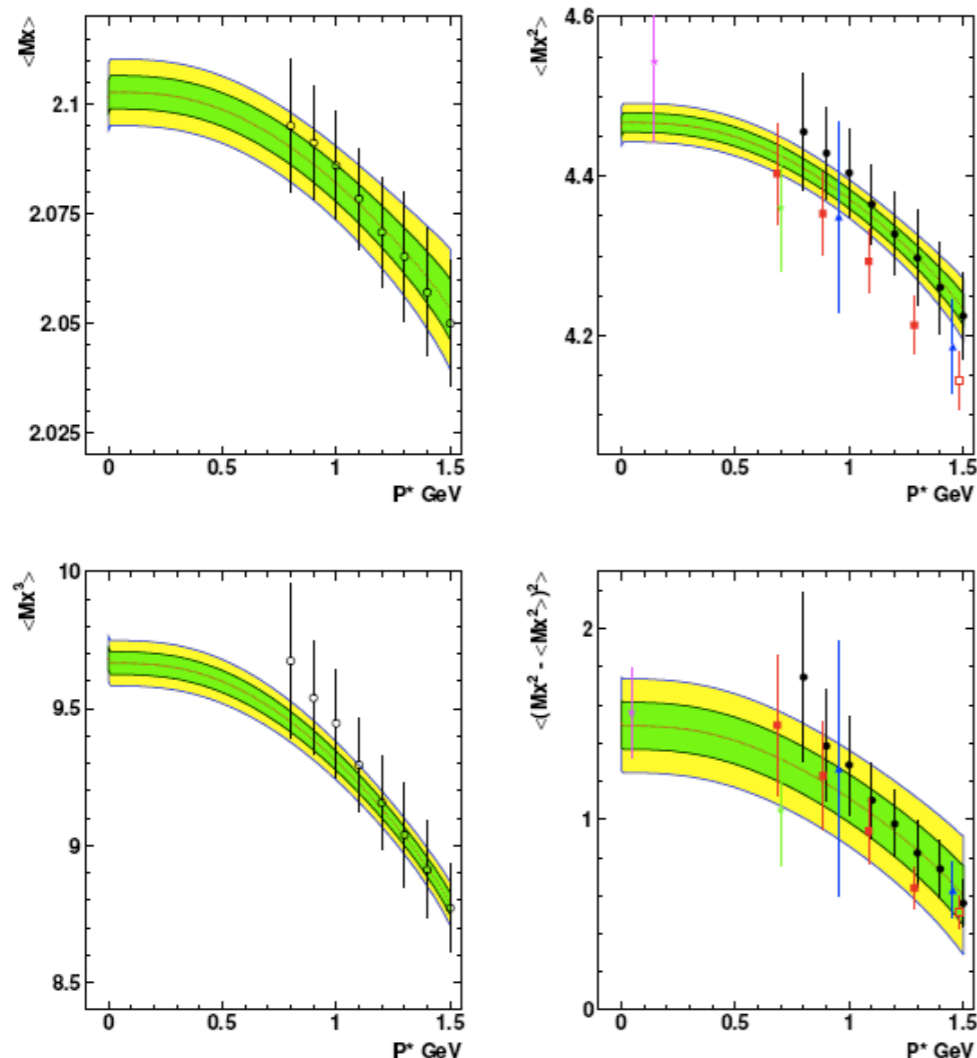
Global fit (kinetic scheme)

Inputs	$ V_{cb} 10^3$	m_b^{kin}	χ^2/ndf
$b \rightarrow c$ & $b \rightarrow s\gamma$	41.67(44)(58)	4.601(34)	29.7/57
$b \rightarrow c$ only	41.48(48)(58)	4.659(49)	24.1/46

Based on PG, Uraltsev & Benson et al

In the kinetic scheme the contributions of gluons with energy below $\mu \approx 1 \text{ GeV}$ are absorbed in the OPE parameters

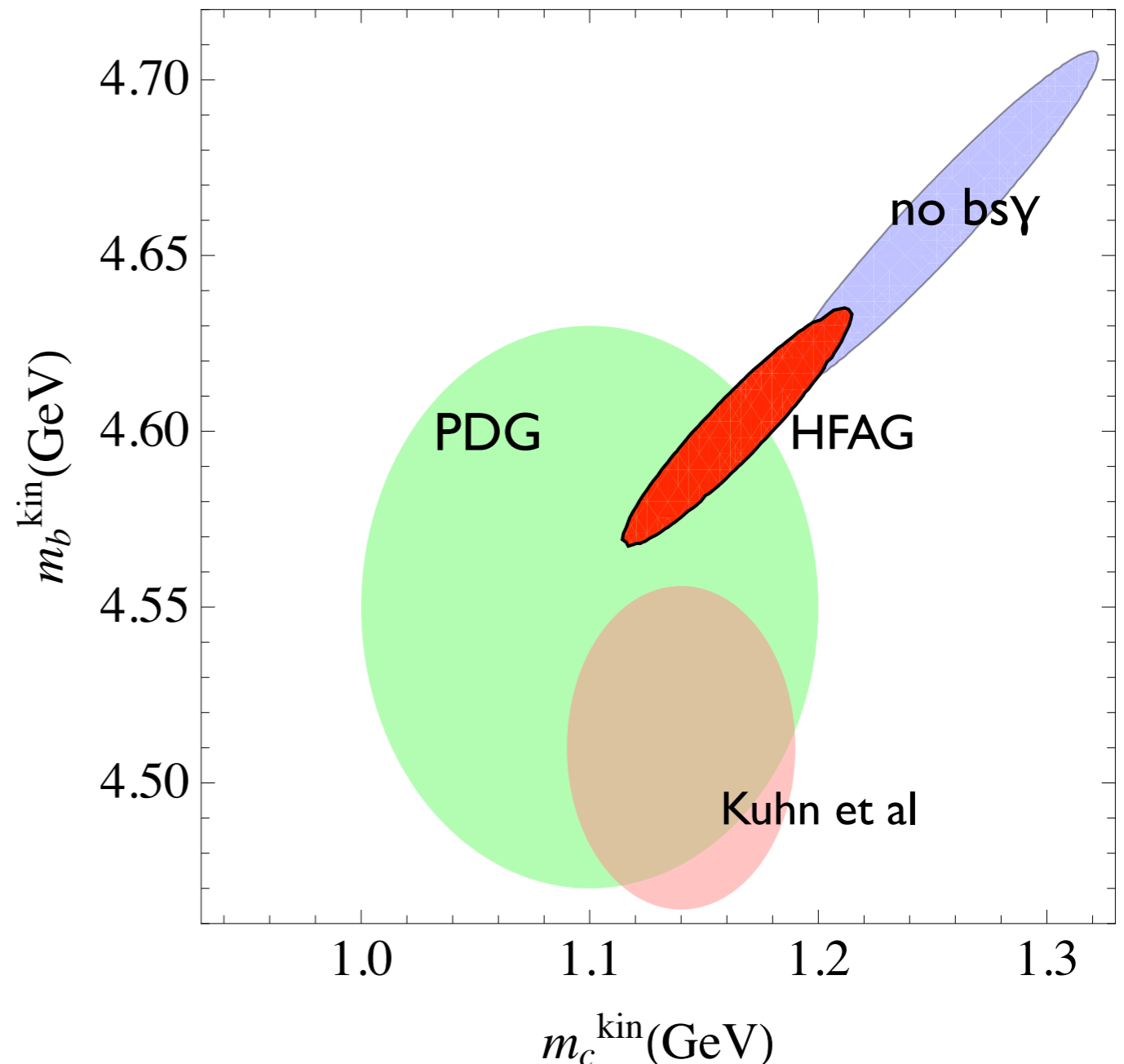
Here scheme means also a number of different assumptions and a recipe for theory errors



Upgrade of code under way,
discrepancy reduced

Fits & Quark Masses

- ▶ Assumes duality but it self-consistently checks it
- ▶ Very close result for $|V_{cb}|$ in $1S$ scheme (Bauer et al)
- ▶ **Higher order** power corr. under control Mannel et al
- ▶ new pert $O(\alpha_s^2) \Rightarrow -0.5\%$ in $|V_{cb}|$ Melnikov, Czarnecki, Pak
- ▶ part of $O(\alpha_s/m_b^2)$ Becher et al
- ▶ Semileptonic moments identify a strip in (m_b, m_c) plane along which the minimum is **shallow**.
- ▶ In the global HFAG fit the $B \rightarrow X_s \gamma$ moments **change significantly** $m_{b,c}$ determination. Without radiative moments the masses are too high!



Fits & Quark Masses

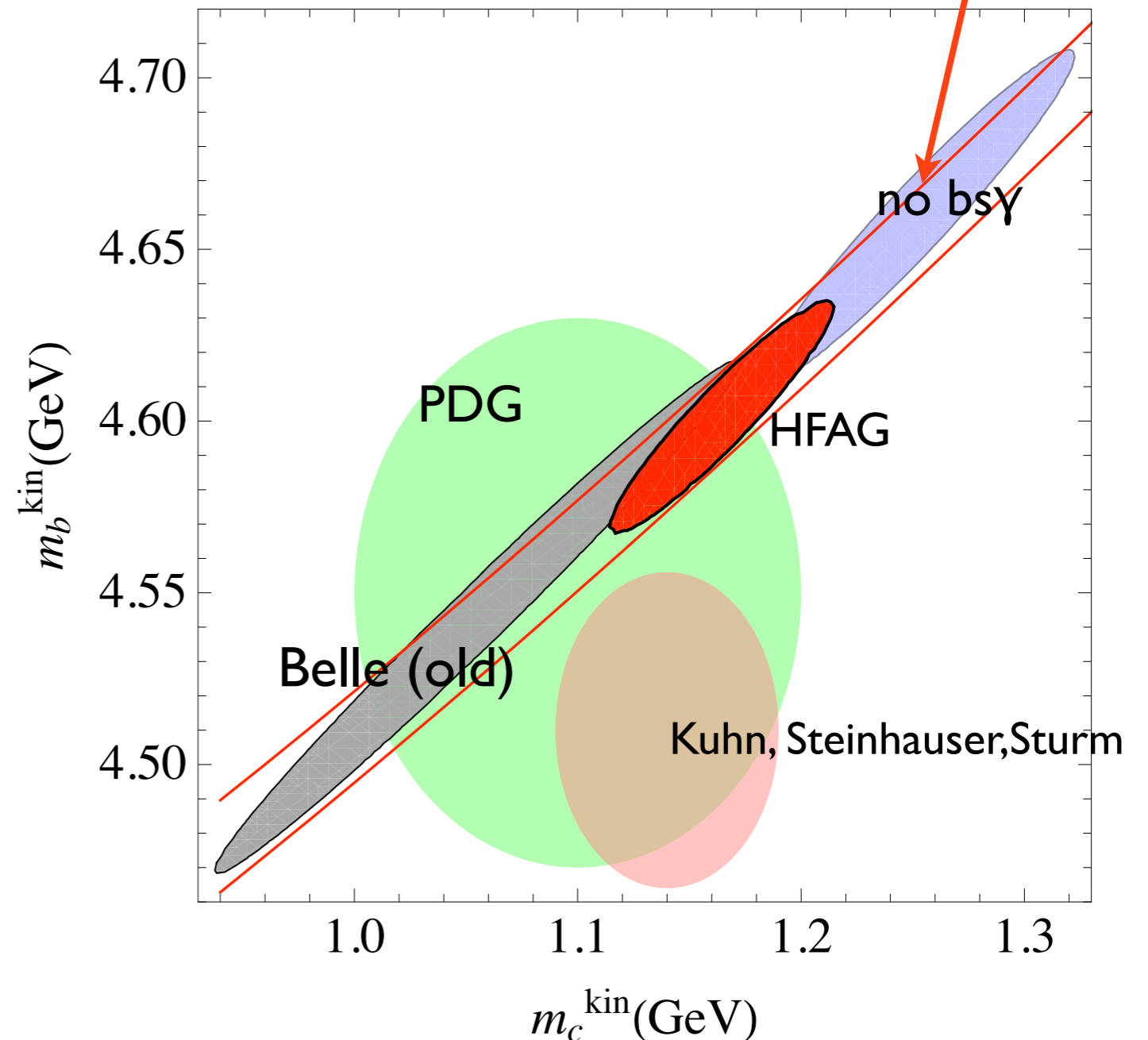
Fitted $|V_{cb}|$ stable, but 2.4σ from latest lattice exclusive determination.

Mass determinations not so stable

OPE fails for $bs\gamma$, but only at $O(\alpha_s)$ with operators $\neq O_7$. Unlikely to be relevant for normalized moments, but it must be studied

At the moment the role of radiative moments in the fits is similar to using PDG bound on m_b .

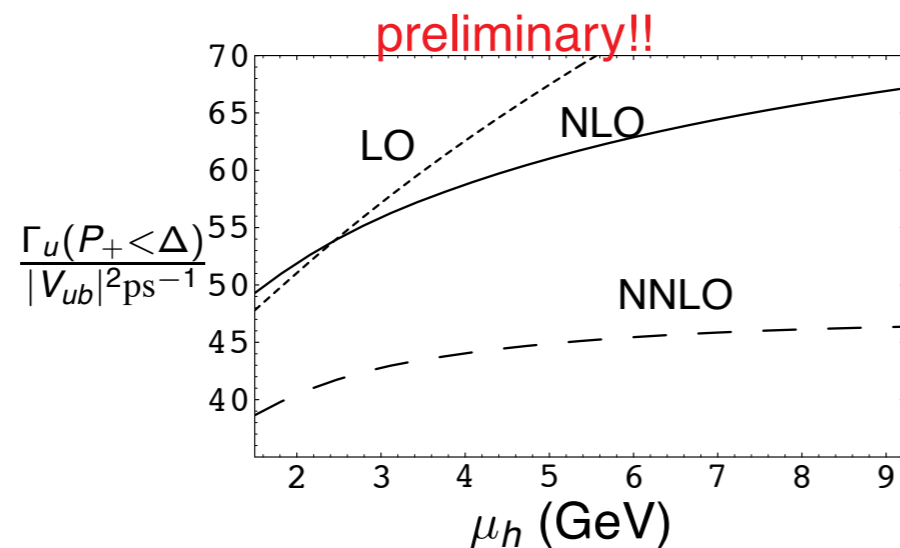
Inclusion of additional constraints?



Upgrade of moments analysis under way

Perturbative calculations

Partial rate for $P_+ < \Delta = M_D^2/M_B$



- ▶ NNLO result is smaller and less dependent on μ_h than NLO
- ▶ would lead to higher $|V_{ub}|$ compared to NLO (preliminary)

Some of the shift is due to different S at LO, NLO, NNLO

Ben Pecjak, ICHEP08

Complete $O(\alpha_s)$ implemented by all groups De Fazio-Neubert

Running coupling NNLO
 $O(\alpha_s^2 \beta_0)$ in GGOU & DGE lead to -5% & +2%, resp. in $|V_{ub}|$
 Gardi, Ridolfi, PG

2008

Asatrian, Greub, Pecjak
 Bonciani, Ferroglia,
 Beneke, Huber, Li
 G. Bell
 in SCET-HQET
 corresponds to fixed order
 $O(\alpha_s^2)$ in the SF region

The problems with cuts

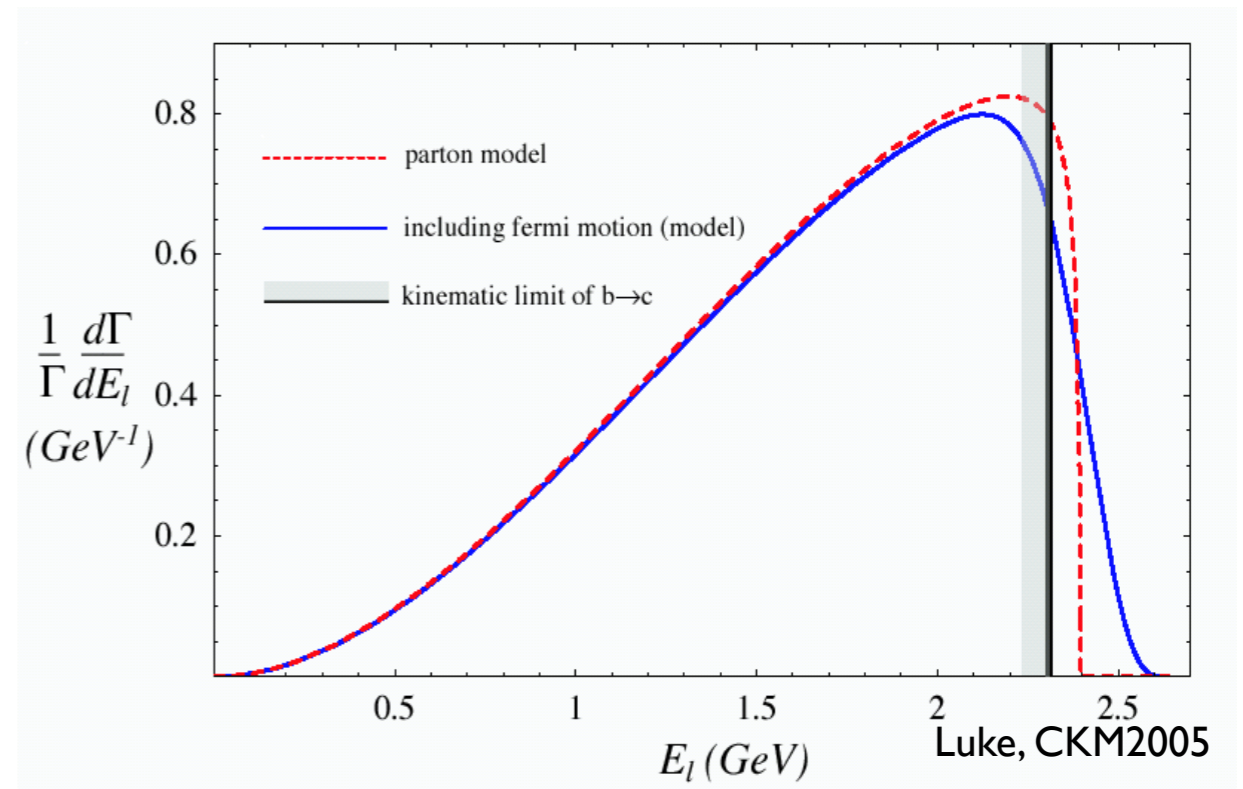
$|V_{ub}|$ from total BR($b \rightarrow ul\nu$) like incl $|V_{cb}|$ but we need kinematic cuts to avoid the $\sim 100x$ larger $b \rightarrow cl\nu$ background:

$$m_X < M_D \quad E_l > (M_B^2 - M_D^2)/2M_B \quad q^2 > (M_B - M_D)^2 \dots$$

or combined (m_X, q^2) cuts

*The cuts destroy convergence of the OPE that works so well in $b \rightarrow c$.
OPE expected to work only away from
pert singularities*

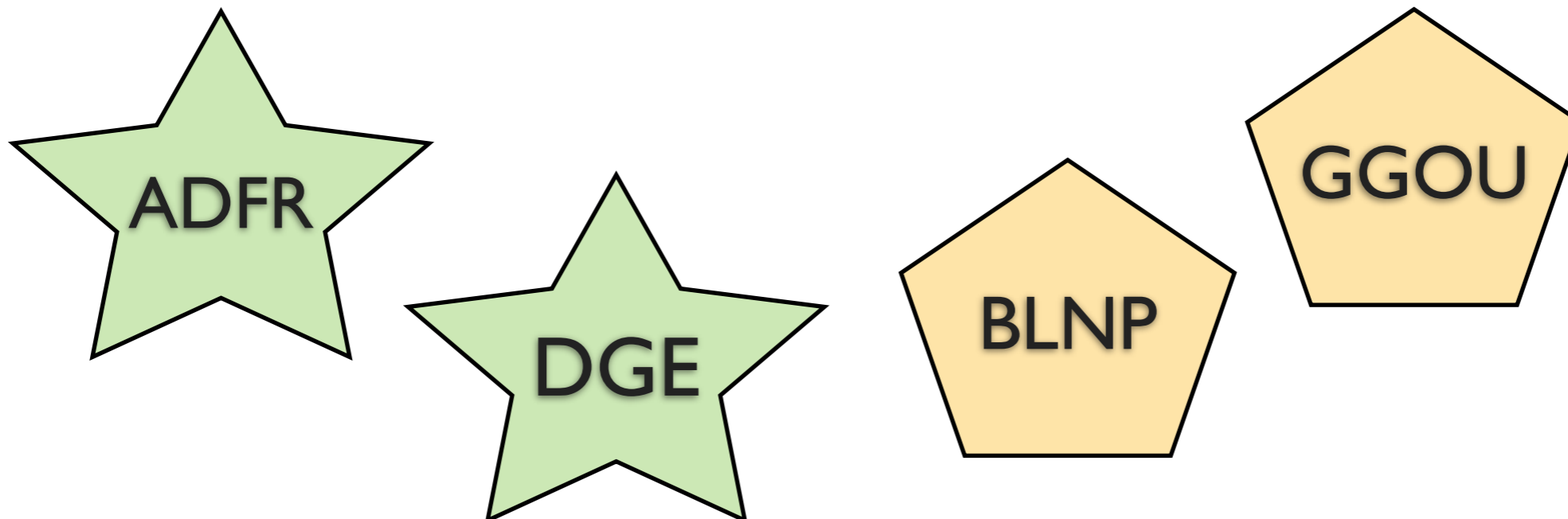
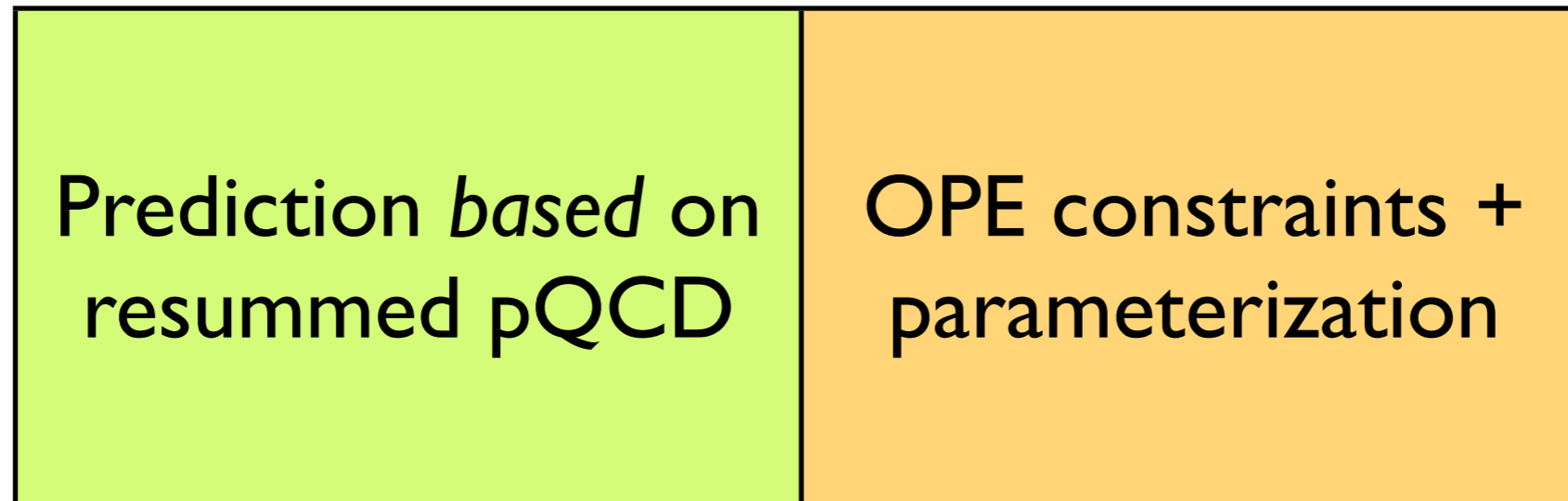
Rate becomes sensitive to “local”
b-quark wave function properties
like Fermi motion Dominant non-
pert contributions can be resummed
into a **SHAPE FUNCTION** $f(k_+)$



How to access the SF?



How to access the SF?



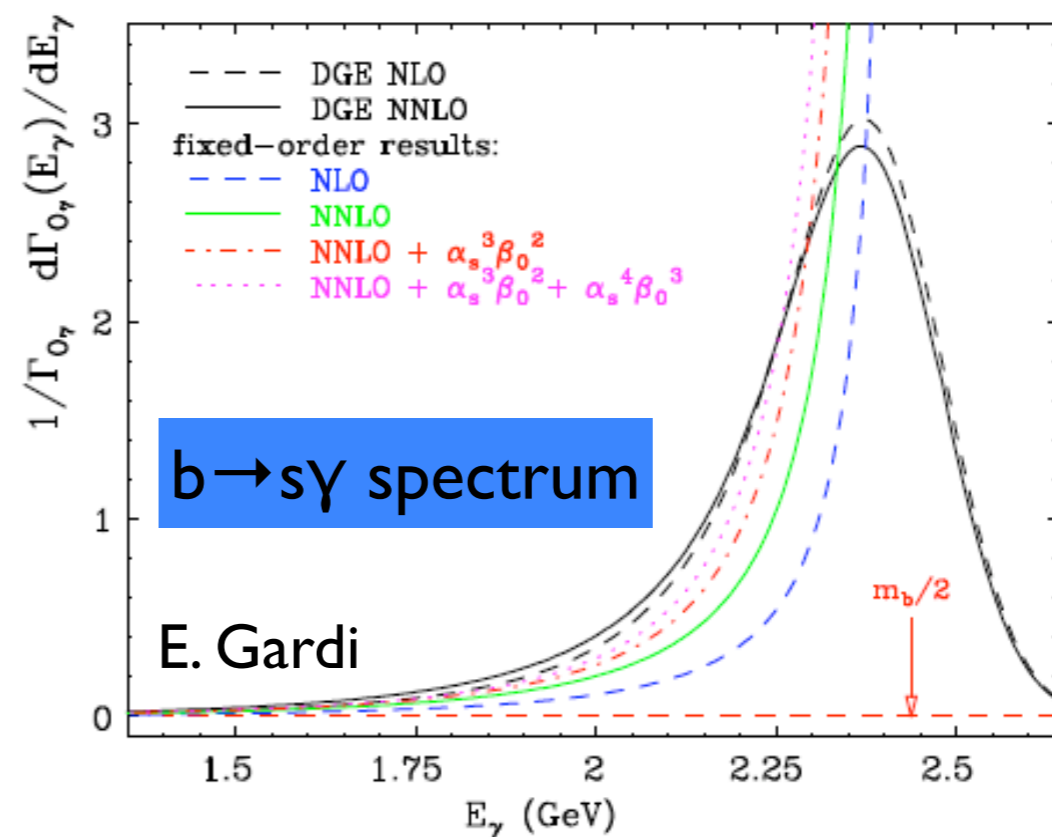
SF from perturbation theory

Resummed perturbation theory is qualitatively different: **Support properties; stability!** (E. Gardi)

b quark SF emerges from resummed pQCD but needs an IR prescription and power corrections for $b \rightarrow B$

Dress Gluon Exponentiation (DGE) by Gardi et al employs renormalon resummation to define Fermi motion. Power corrections can be partly accommodated.

Aglietti et al (ADFR) use Analytic Coupling in the IR



The SF in the OPE

Local OPE has also threshold singularities and SF can be equivalently introduced resumming dominant singularities Bigi et al, Neubert

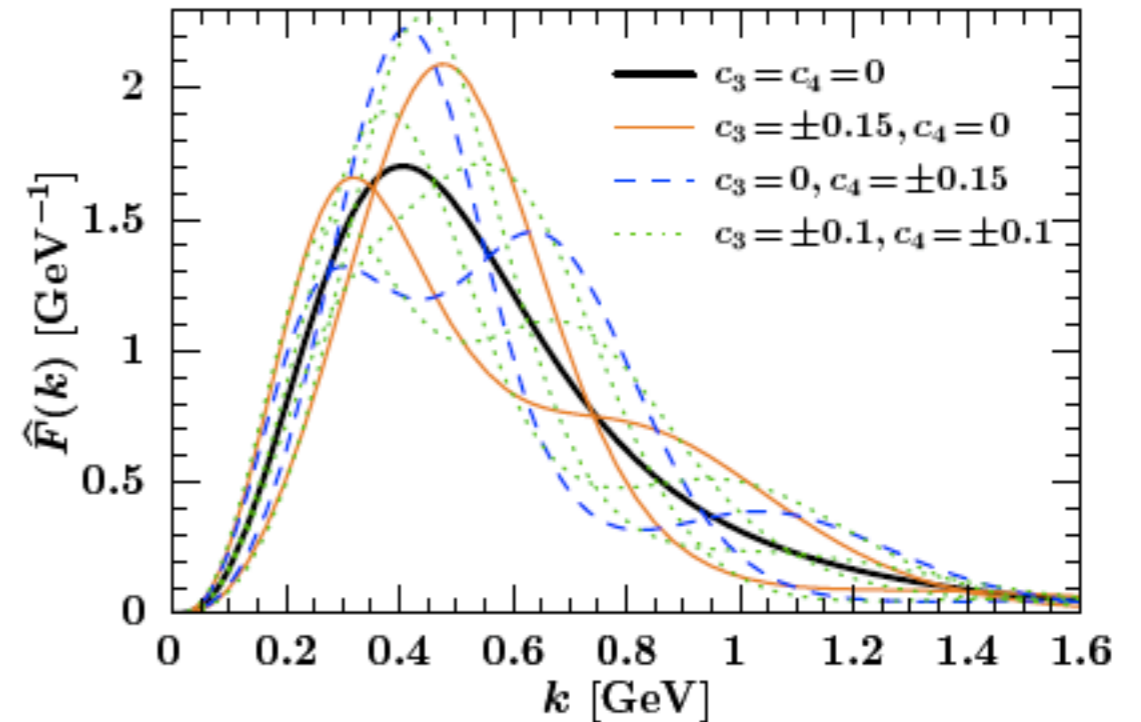
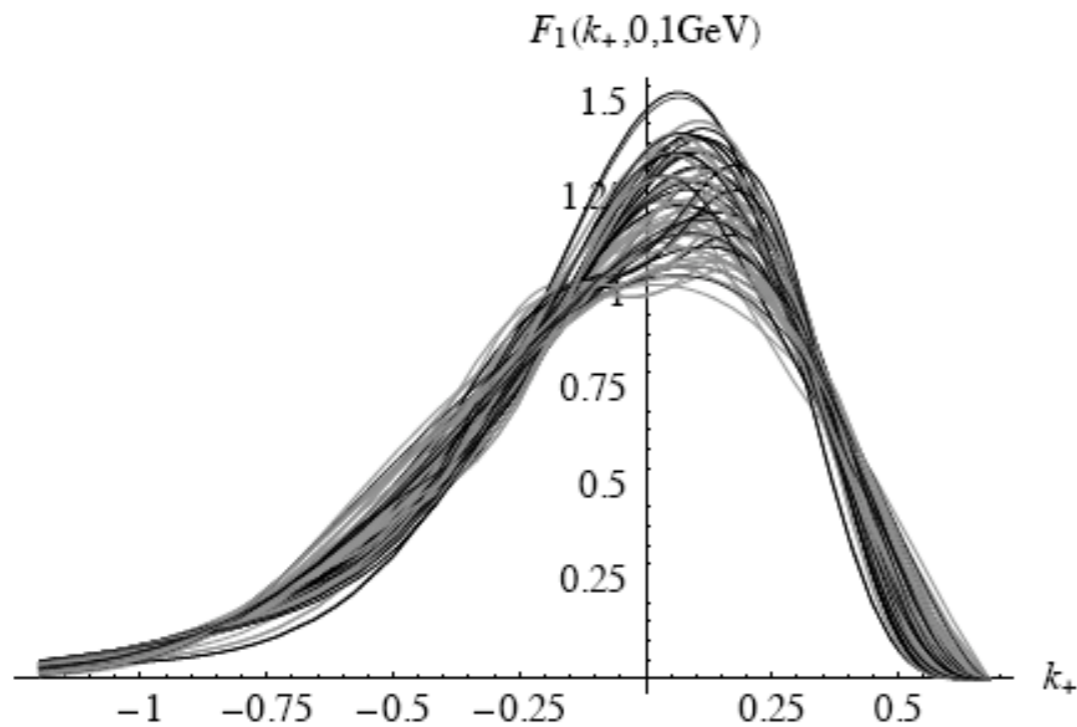
Fermi motion can be parameterized within the OPE like PDFs in DIS. At leading order in m_b only a single universal function of one parameter enters (SF).

*Unlike resummed pQCD, **the OPE does not predict the SF**, only its first few moments. One then **needs an ansatz for its functional form**.*

$$\int dk_+ k_+^n F_i(k_+, q^2) = \text{local OPE prediction} \Leftrightarrow \text{moments fits}$$

*Two very different implementations:
PG, Giordano, Ossola, Uraltsev (GGOU)
Bosch, Lampe, Neubert, Paz (BLNP)*

Functional forms



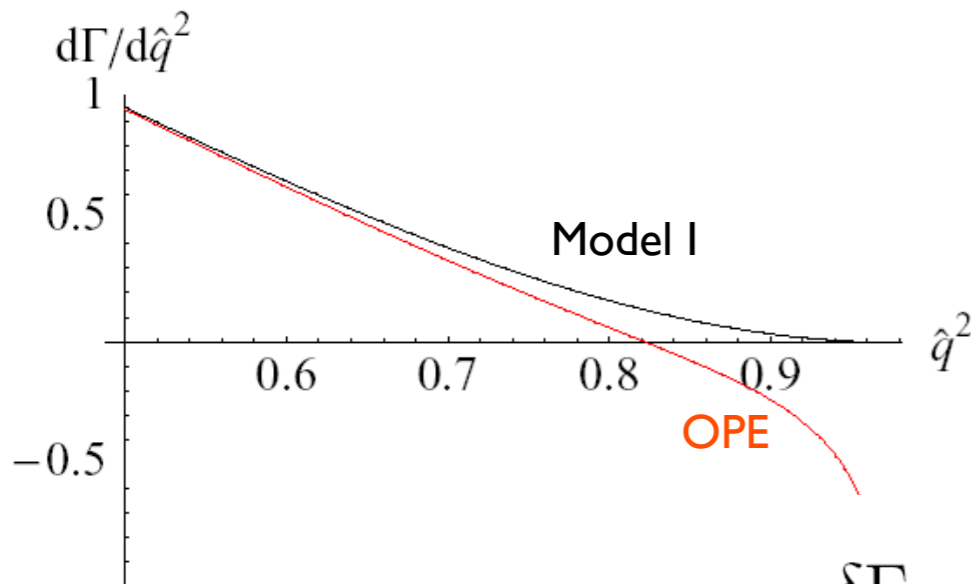
About 100 forms considered in GGOU, large variety, double max discarded. Small uncertainty (1-2%) on V_{ub}

Recent more systematic method by Ligeti et al. arXiv:0807.1926
Plot shows 9 SFs that satisfy all the first three moments

The high q^2 tail

At high q^2 higher dimensional operators are not suppressed leading to pathological features. Origin in the non-analytic square root

$$\frac{d\Gamma}{dq_0 dq^2} \propto \sqrt{q_0^2 - q^2} \quad \Rightarrow \quad \frac{d\Gamma}{dq^2} \sim - \sum_{n=1}^{\infty} \frac{(-1)^n b_n(\hat{q}^2)}{(1 - \hat{q}^2)^{n-2}} \left(\frac{\bar{\Lambda}}{m_b}\right)^n$$



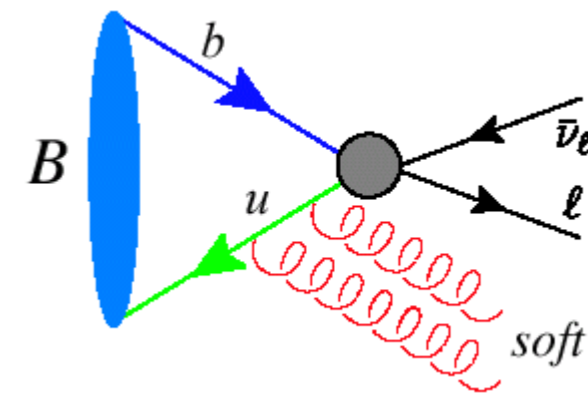
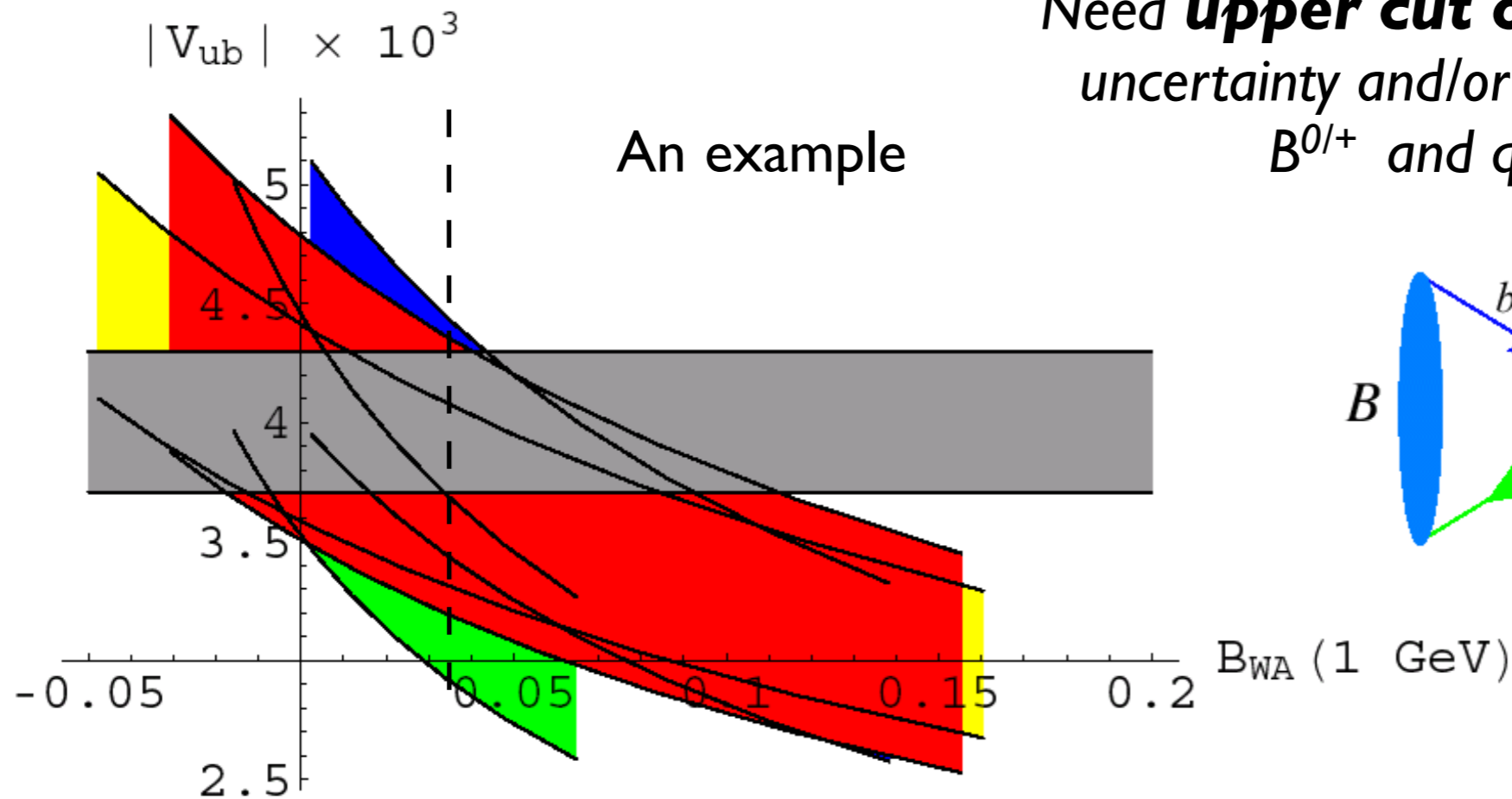
In the integrated rate the $1/m_b^3$ singularity is removed by the WA operator: needs modelling for q^2 spectrum

$$\delta\Gamma \sim \left[C_{\text{WA}} B_{\text{WA}}(\mu_{\text{WA}}) - \left(8 \ln \frac{m_b^2}{\mu_{\text{WA}}^2} - \frac{77}{6} \right) \frac{\rho_D^3}{m_b^3} + \mathcal{O}(\alpha_s) \right]$$

WA matrix element B_{WA} parameterizes global properties of the tail, affects V_{ub} determinations depending on cuts, tends to decrease V_{ub}

Constraining Weak Annihilations

WA may pollute all present estimates, and tend to **decrease** the extracted V_{ub} . Need **upper cut on q^2** to remove this uncertainty and/or constrain WA from $B^{0/+}$ and q^2 spectrum



Comparing the existing approaches at common m_b (HFAG ichep08, CKM08)

$|V_{ub}|$ from DGE

Gardi & Andersen
see Gardi talk

Main features of the spectra
are reproduced $\implies |V_{ub}|$ stable,
small errors and good χ^2

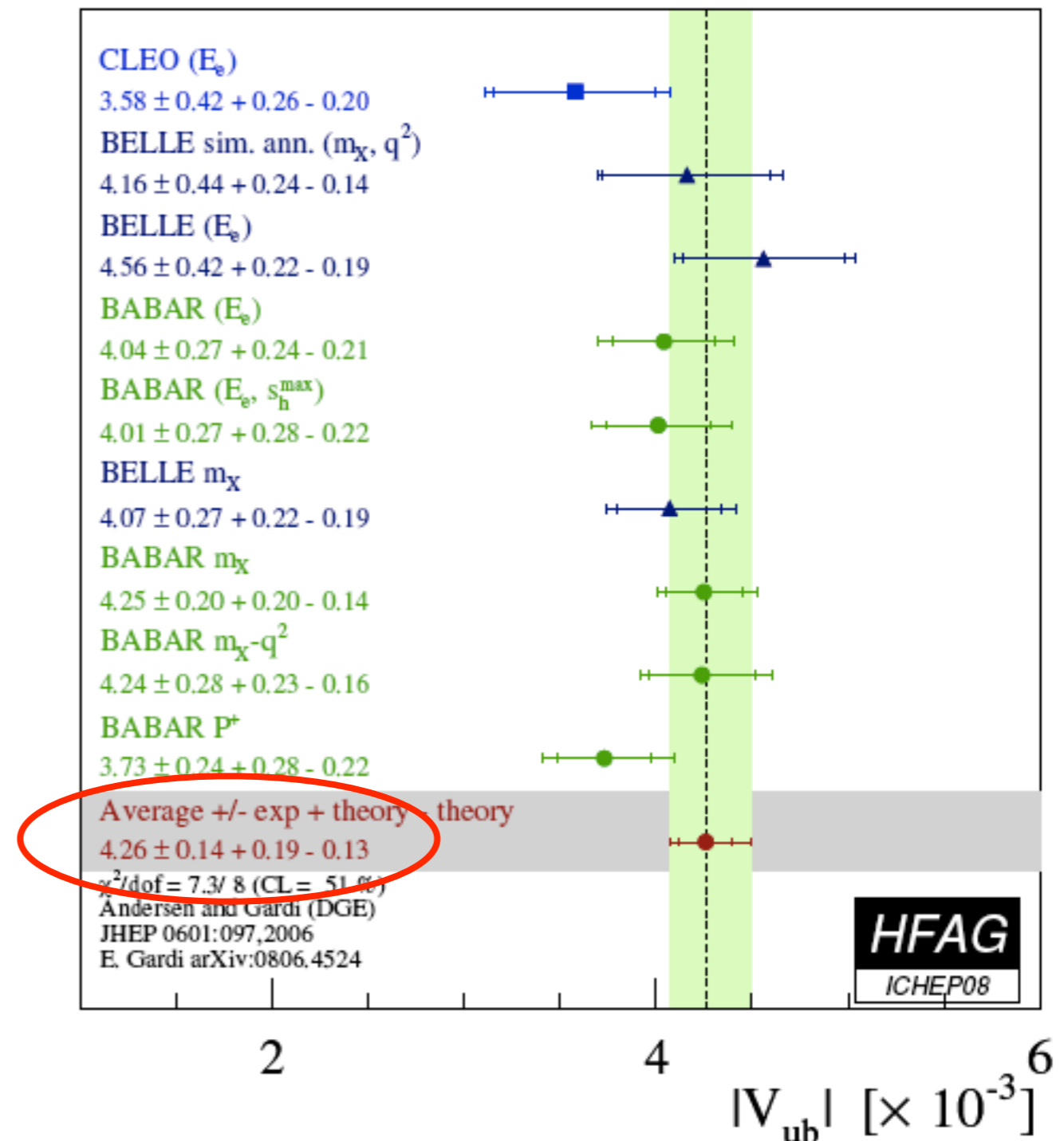
NNLL and $O(\alpha_s^2\beta_0)$ implemented

Power corrections in the SF region
are included here only in theor. err.
No subleading SF.

WA error equal for all cuts.
Matches to local OPE.

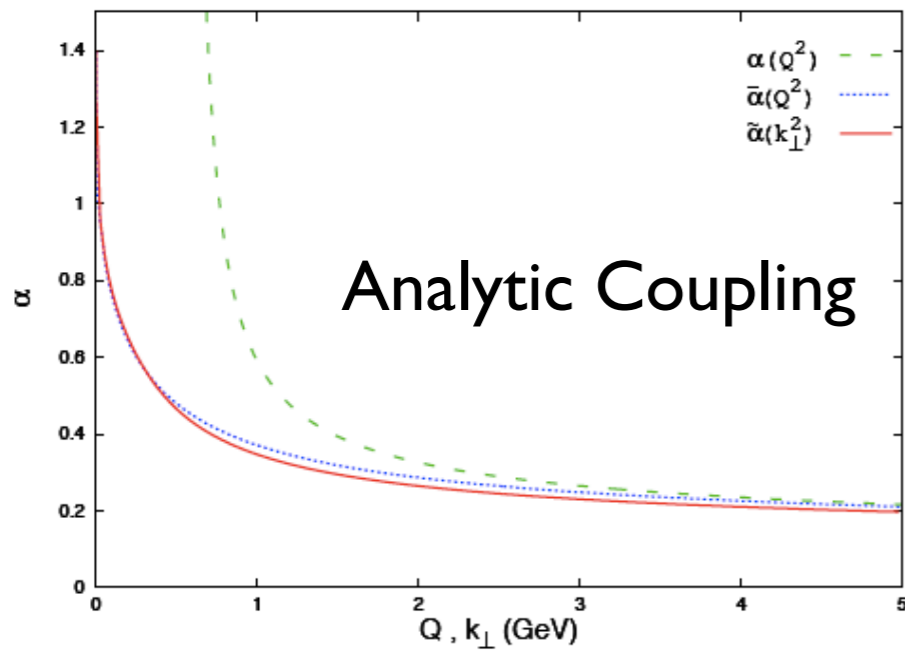
Only input other than α_s
 $m_b(m_b)=4.24(4)$ from global fit

5-6% total error, mostly m_b



$|V_{ub}|$ from ADFR

Aglietti, Di Lodovico, Ferrera, Ricciardi



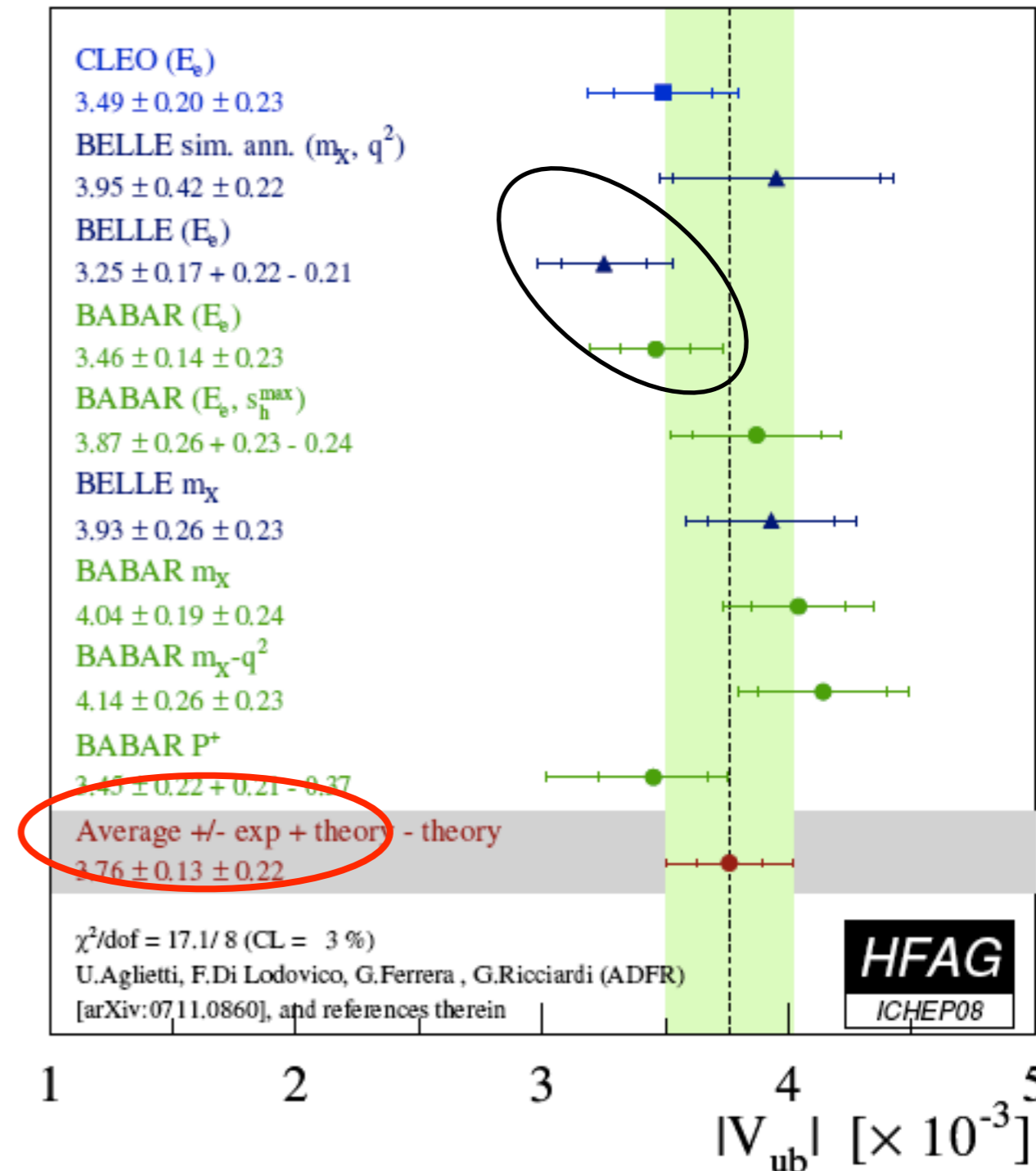
Worse consistency here.

NNLO resummation, NLO constants

Consider E_l cuts higher than 2.3 GeV because their E_l apparently does not reproduce data (see later)

employs M_B in on-shell calculation of spectra: no renormalon cancellation, no convergence to OPE. no model error

~7% total error, mostly m_c



$|V_{ub}|$ in BLNP

Bosch, Lange, Neubert, Paz

$$\tilde{W}_1^{(0)}(P_+, y) = U_y(\mu_h, \mu_i) H(y, \mu_h) \int_0^{P_+} d\hat{\omega} m_b J(y, m_b(P_+ - \hat{\omega}), \mu_i) \hat{S}(\hat{\omega}, \mu_i)$$

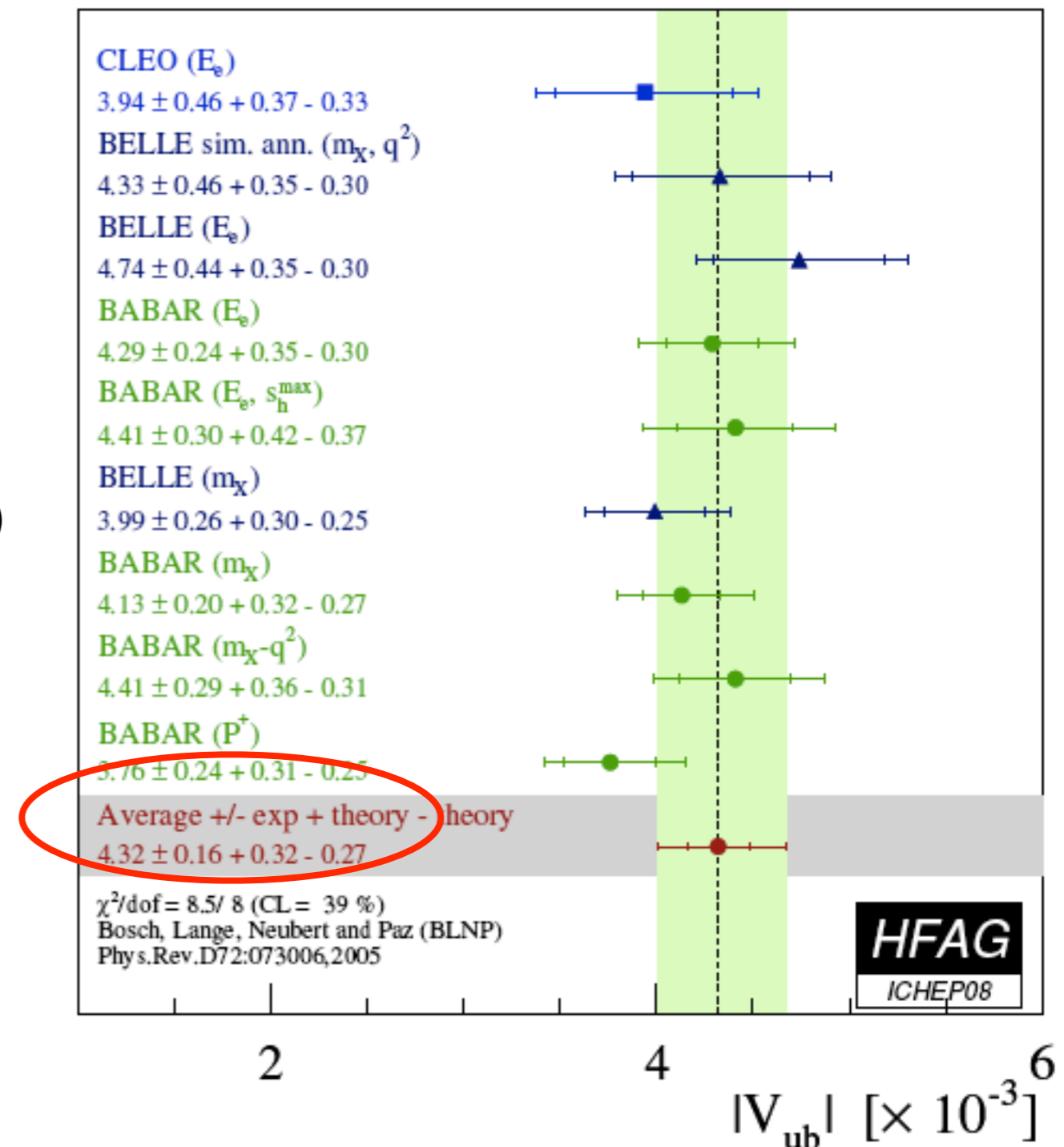
$$d\Gamma = HJ \otimes \hat{S} + \frac{1}{m_b} H'_i J'_i \otimes \hat{S}'_i + \dots$$

Good consistency. Uses elegant multiscale OPE that resums soft-collinear logs, but many largely unconstrained subleading SFs

NNLL resummation, only $O(\alpha_s, \Lambda^2/m_b^2)$ matching to OPE, 3 ffs for leading SF, extensive modelling of SSF.

m_b and μ_π^2 in SF scheme obtained from global fit in the kin scheme

~7-8% total error, main error HQE parameters



SF in GGOU

Leading SF resums leading twist effects, $m_b \rightarrow \infty$
universal, q^2 indep



Finite m_b distribution functions include all $1/m_b$ effects, *non-universal*
no need for subleading SFs

$$F(k_+) \longrightarrow F_i(k_+, q^2, \mu)$$

Structure function
($i = 1, 2, 3$)
q² dependence
cutoff dependence
(gluons with $E_g < \mu$)

$$\frac{d^3\Gamma}{dq^2 dq_0 dE_\ell} = \frac{G_F^2 |V_{ub}|^2}{8\pi^3} \left\{ q^2 W_1 - \left[2E_\ell^2 - 2q_0 E_\ell + \frac{q^2}{2} \right] W_2 + q^2 (2E_\ell - q_0) W_3 \right\}$$

$$W_i(q_0, q^2) = m_b^{n_i}(\mu) \int dk_+ F_i(k_+, q^2, \mu) W_i^{pert} \left[q_0 - \frac{k_+}{2} \left(1 - \frac{q^2}{m_b M_B} \right), q^2, \mu \right]$$

This factorization formula perturbatively defines the distribution functions
see also Benson, Bigi, Uraltsev for bsy

$$\int dk_+ k_+^n F_i(k_+, q^2) = \text{local OPE} \quad \text{Importance of subleading effects}$$

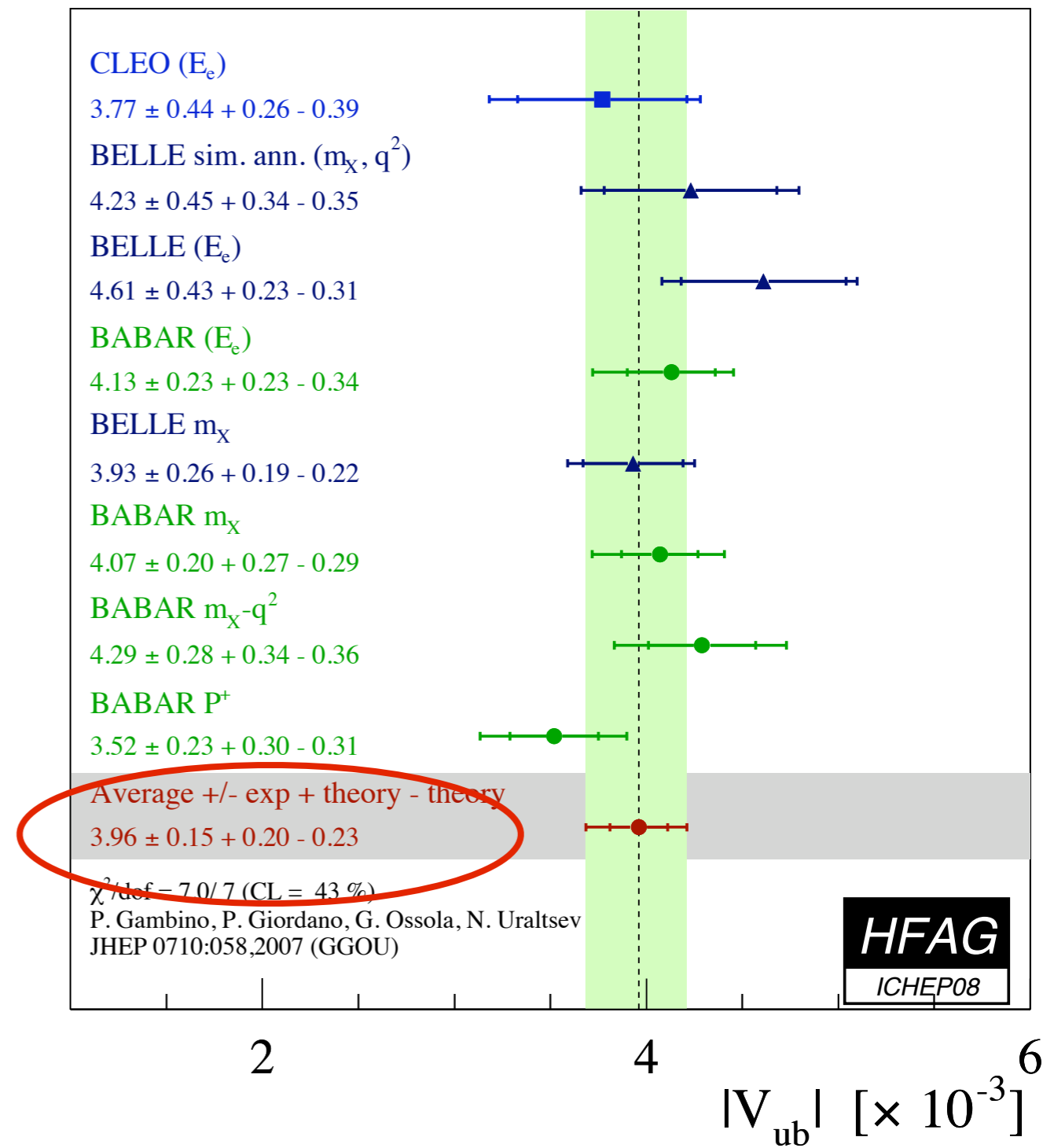
$|V_{ub}|$ in the kinetic scheme -GGOU

PG, Giordano, Ossola, Uraltsev

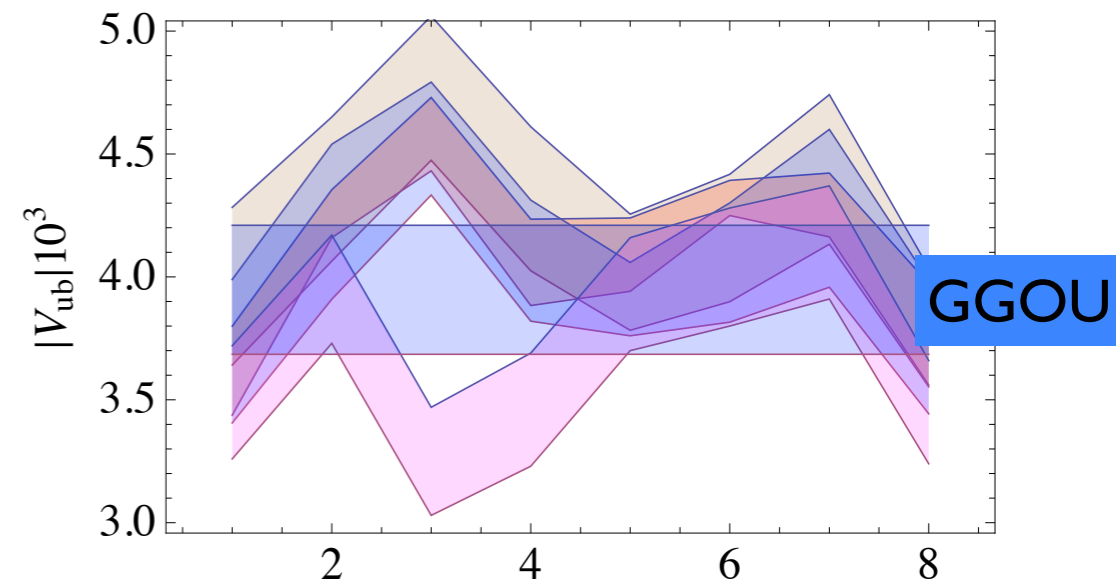
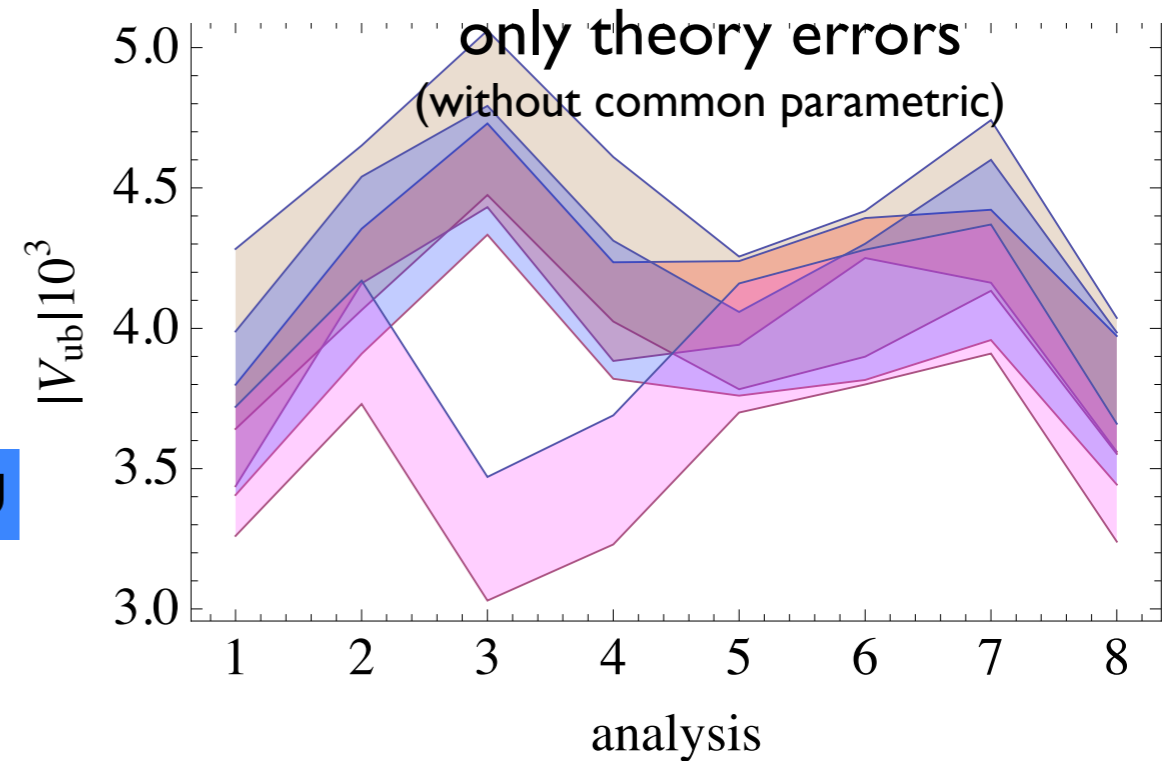
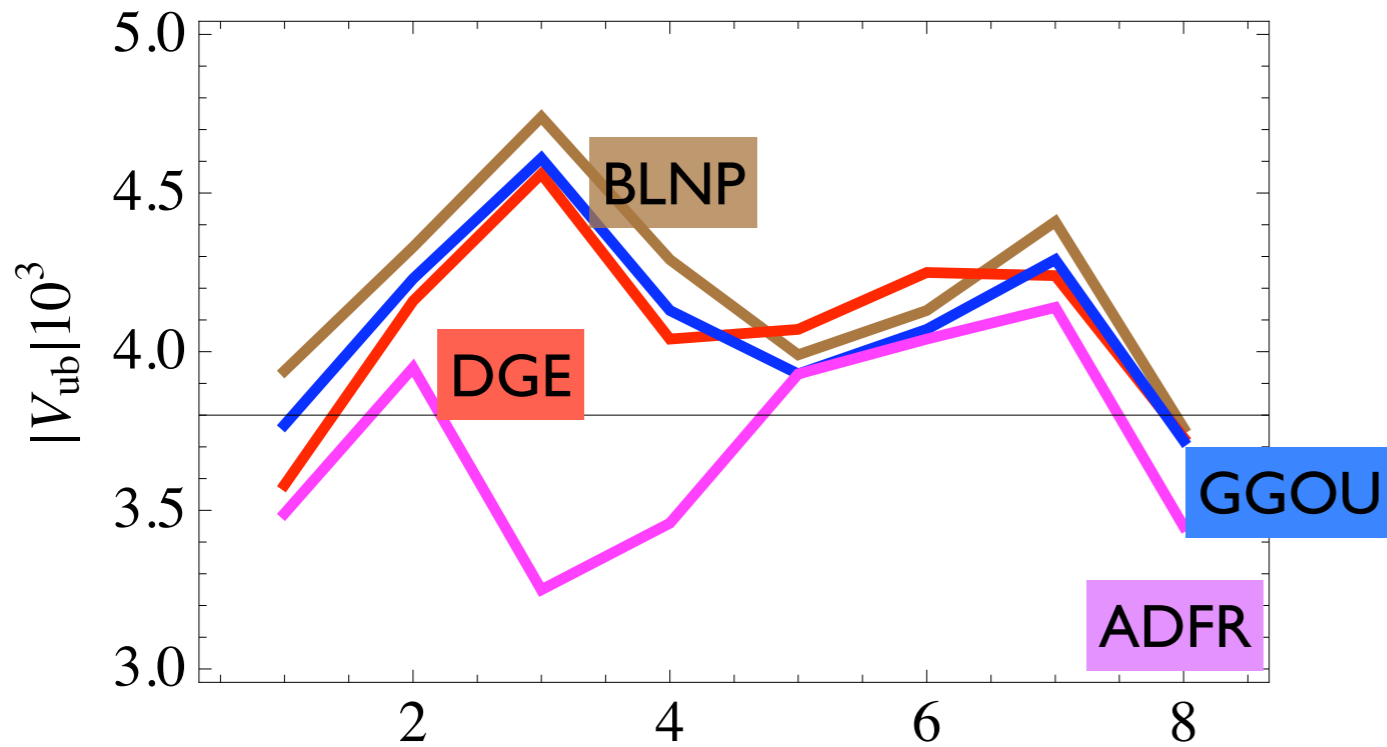
Good consistency & small th error.
 OPE in a scheme with Wilsonian IR cutoff
 $\sim 1 \text{ GeV}$, all subleading $1/m_b$ and $O(\alpha_s^2 \beta_0)$
 terms consistently included,
 careful treatment of high q^2 tail.

Inputs from global fit to the moments

+6.3-7.0% total error



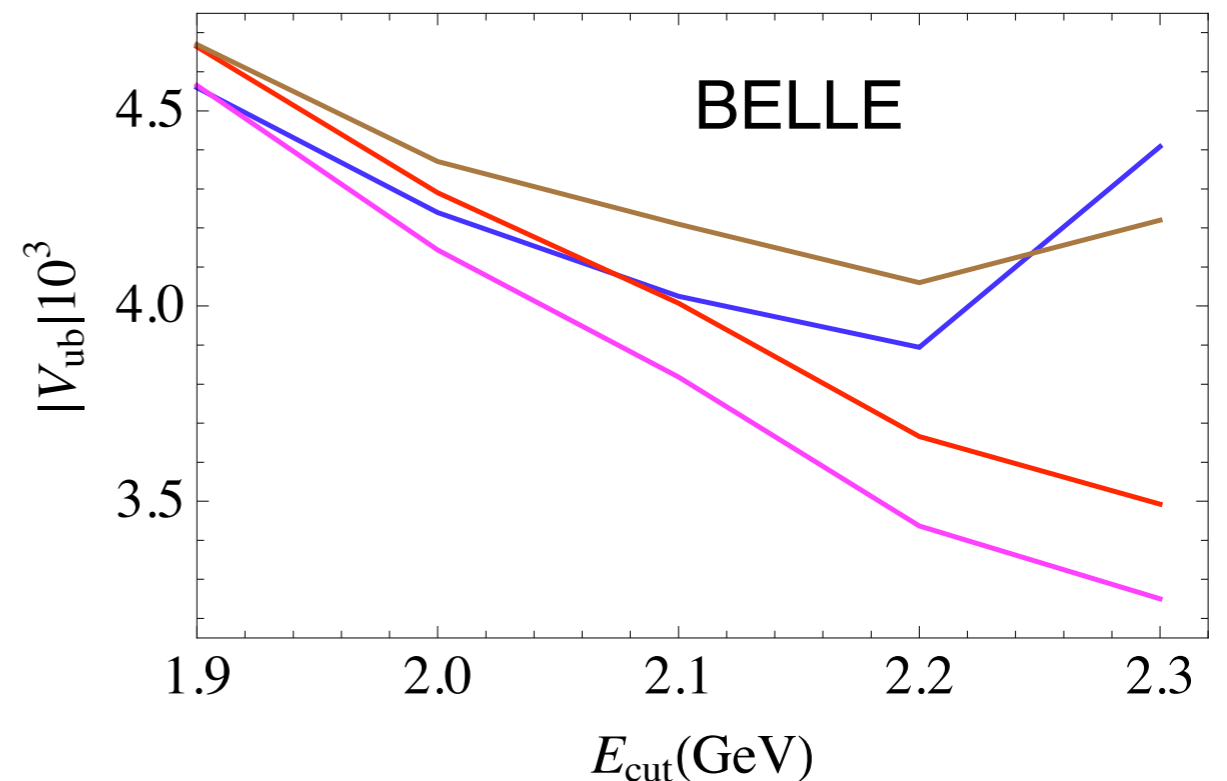
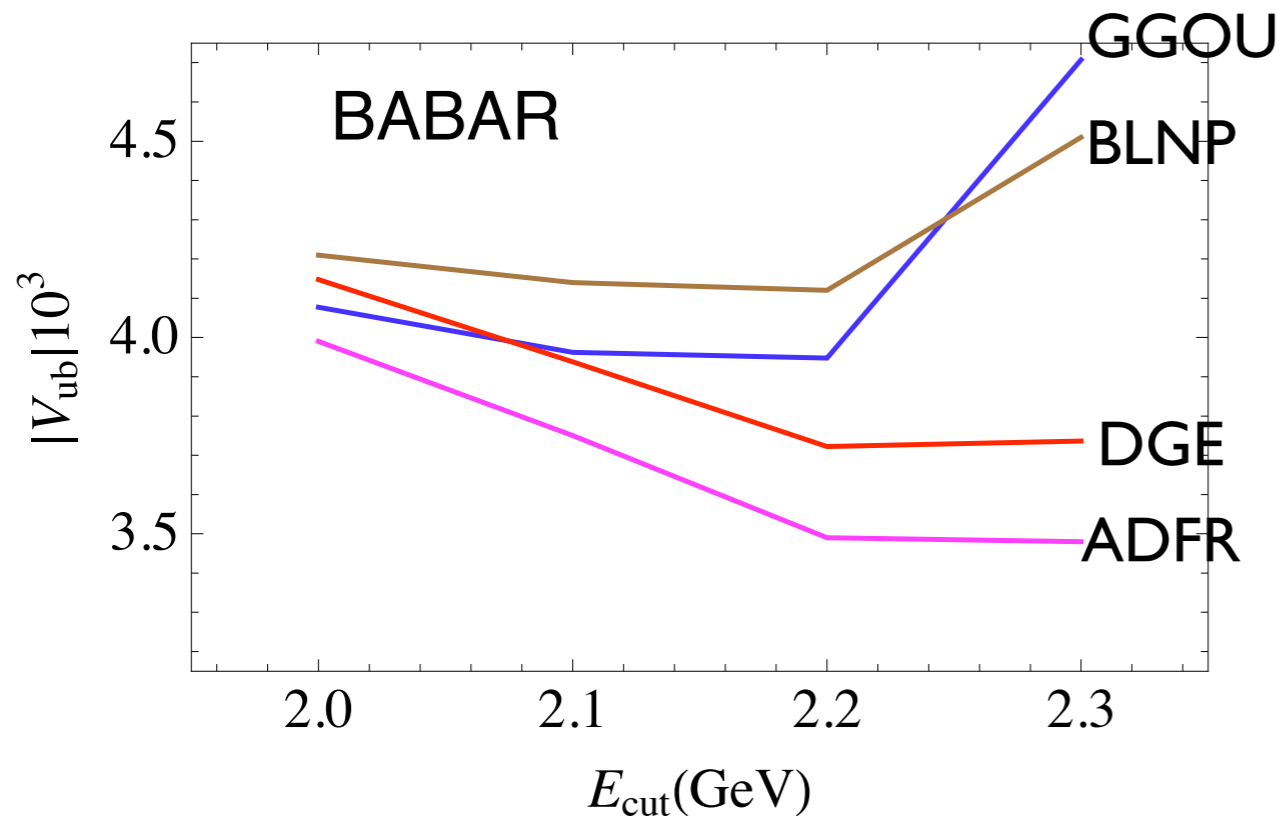
A global comparison



- * common inputs (except ADFR)
- * Overall good agreement with one exception
SPREAD WITHIN THE ERRORS!
- * Systematic offset of central values:
normalization? to be investigated
- * Very different methods, common systematics?
WA, inputs, pert corrections

Why do central values differ up to 9-10%?

The lepton spectrum

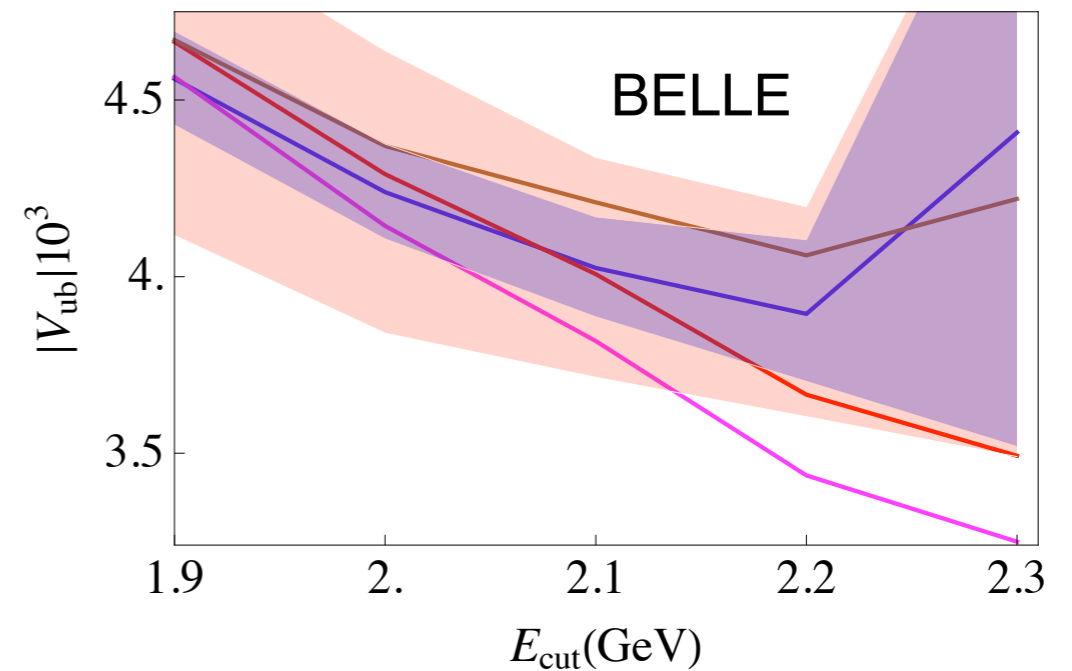
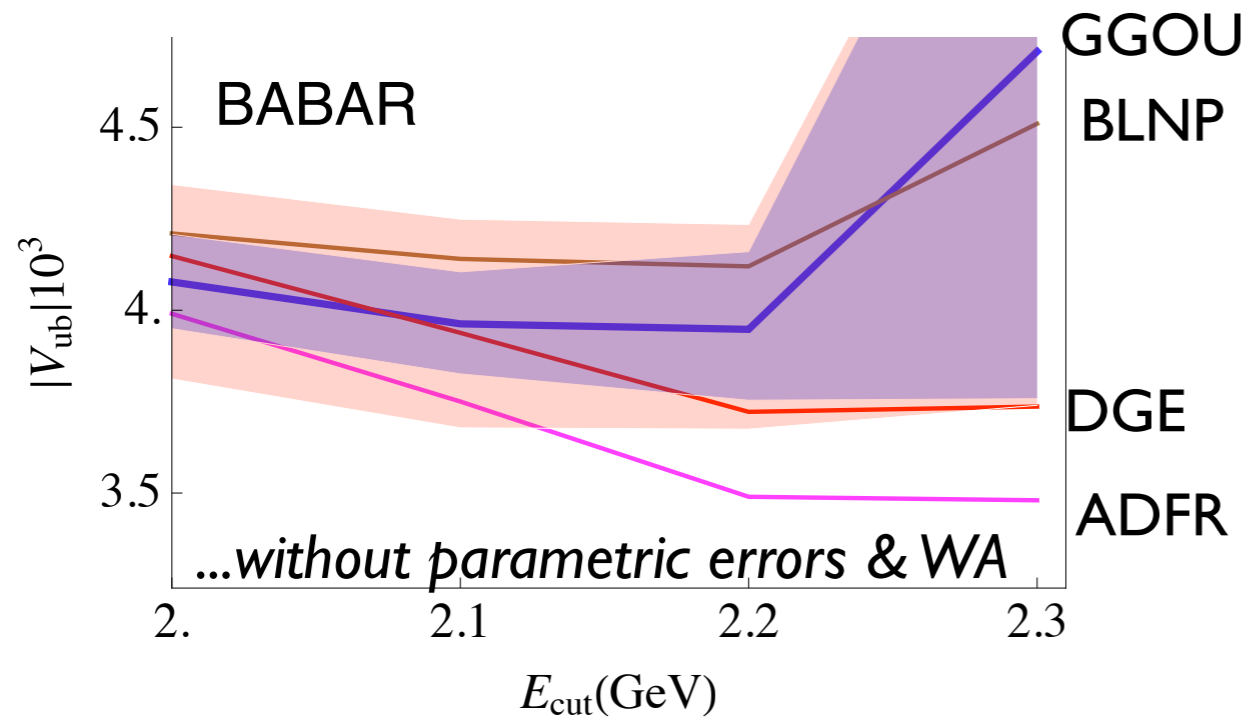


Babar E_l determination

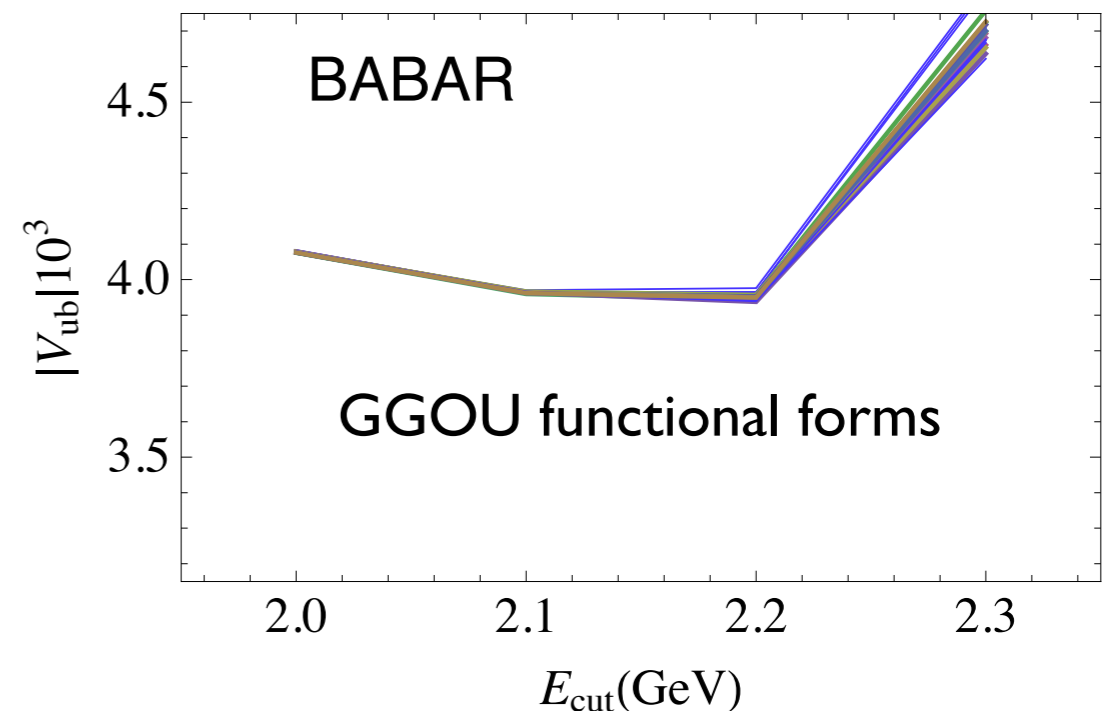
Belle E_l determination

Common inputs, $m_b^{\text{kin}}=4.60\text{GeV}$ or $m_b(m_b)=4.24\text{GeV}$.
Exp analyses depend strongly on generator (slight inconsistency here...)

The lepton spectrum

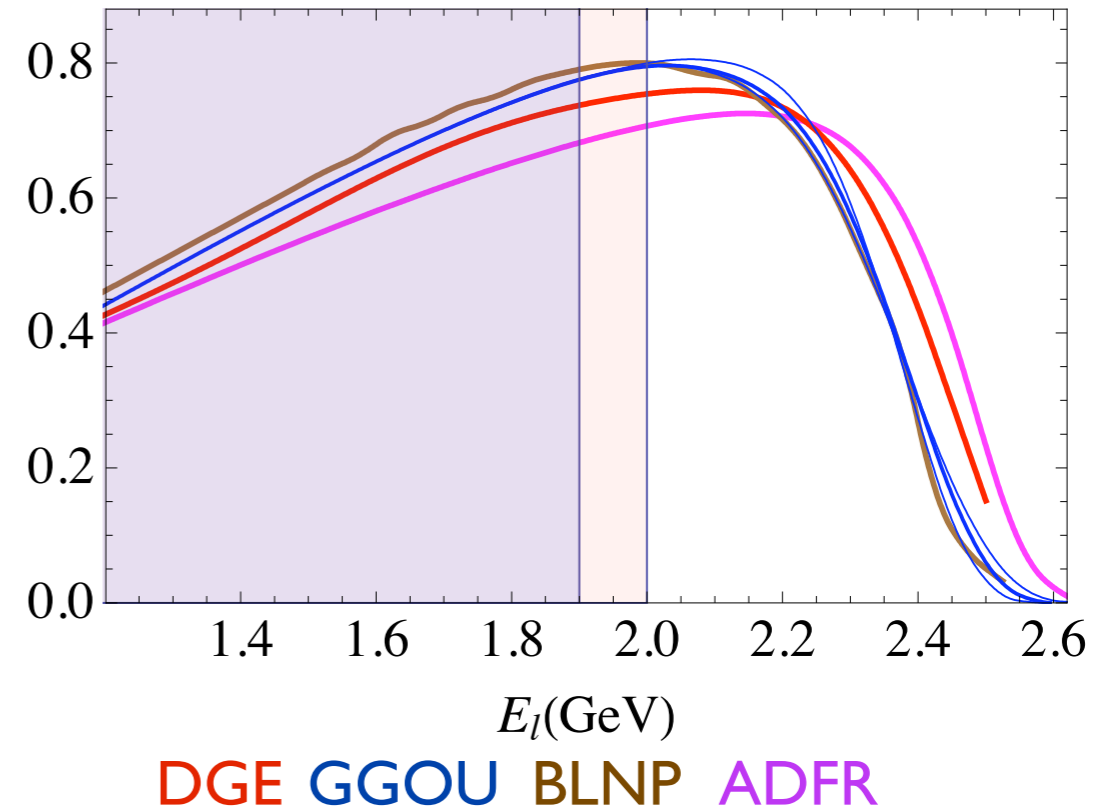
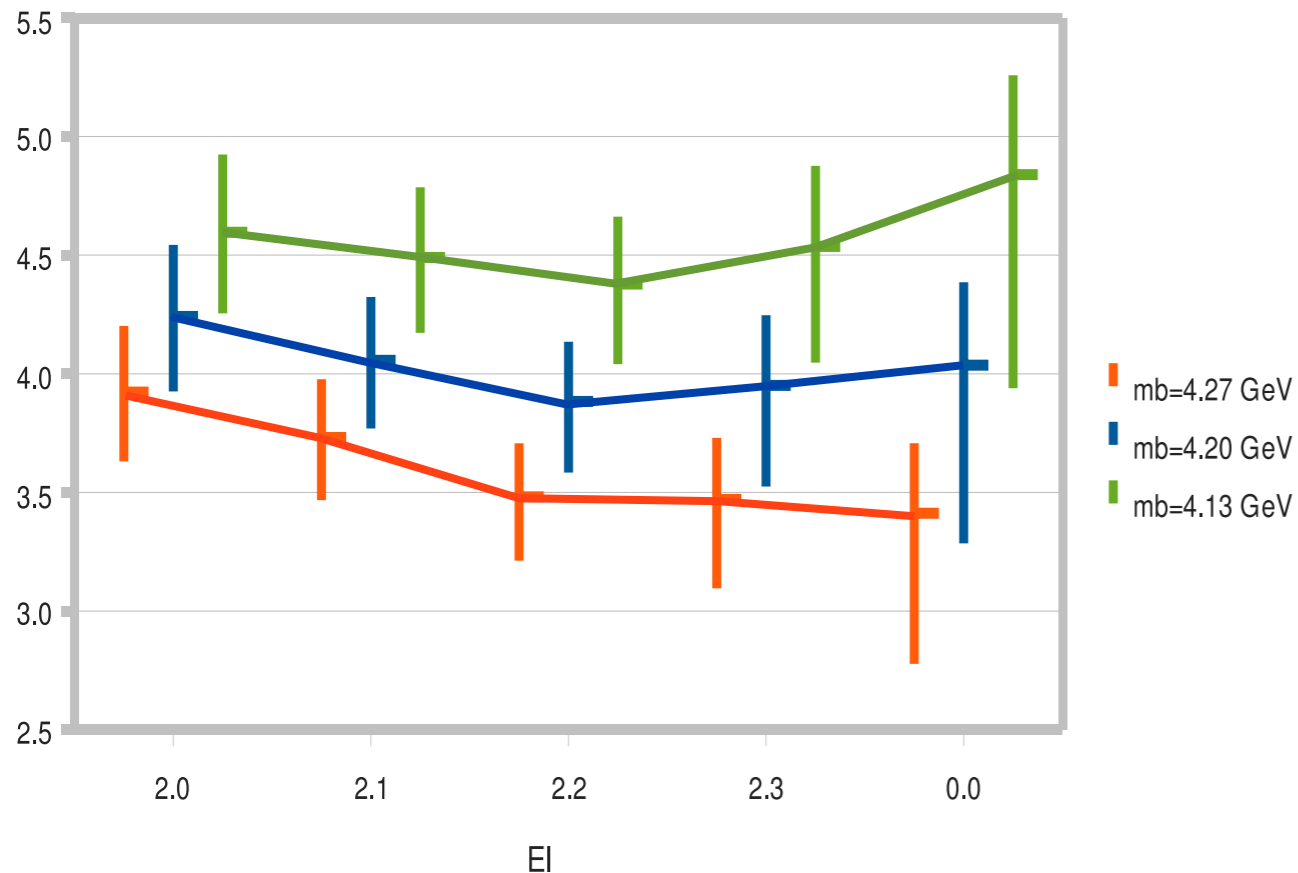


The spectrum does provide information:
 OPE based methods close to each other up to 2.2GeV, resummed methods show larger slope, seem to behave in same way



The lepton spectrum

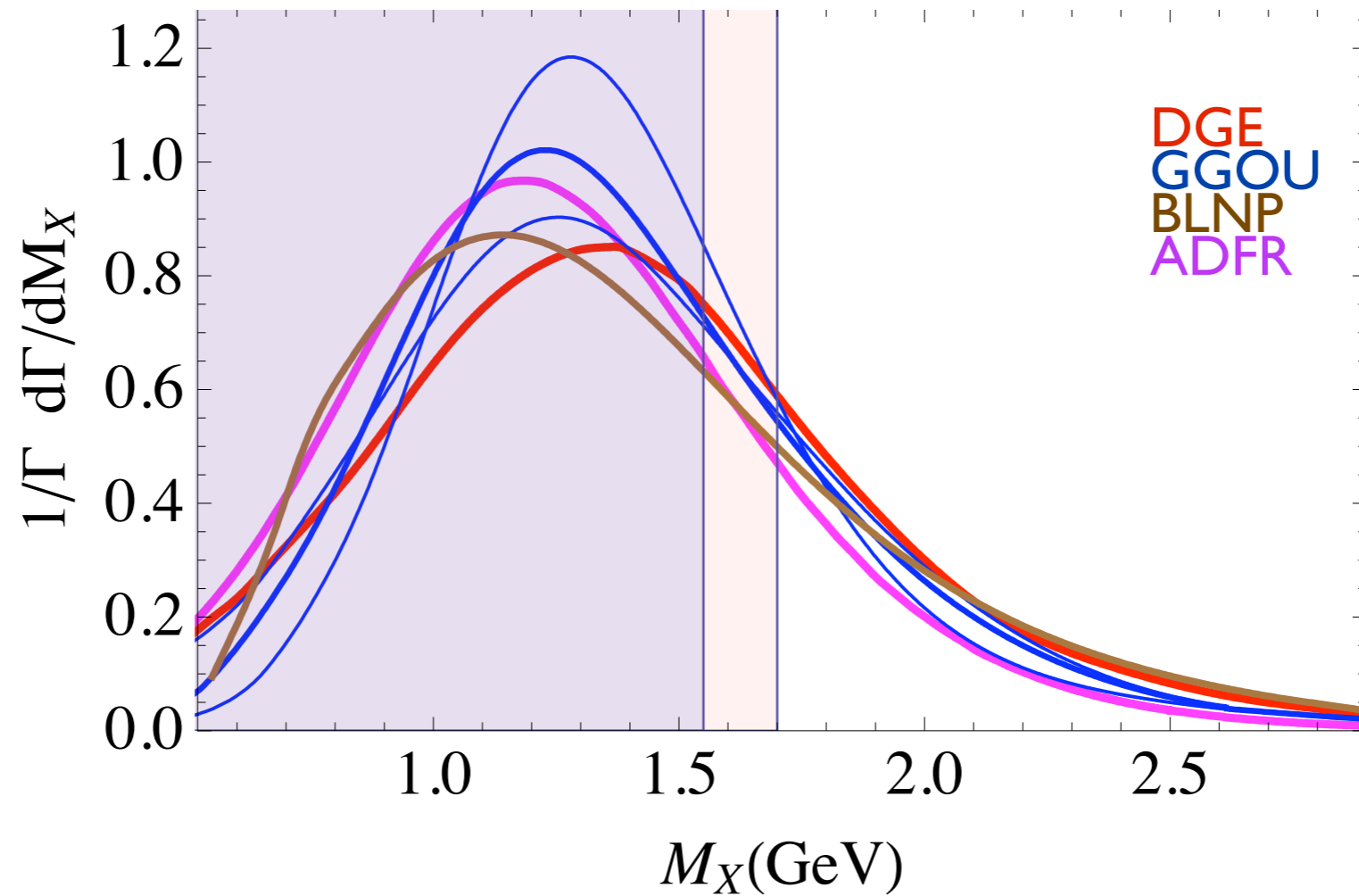
DGE slope vs m_b



*In DGE the slope depends on m_b ,
while ADFR have it fixed*

The leptonic spectrum is not sensitive to the SF except quite close to the endpoint. At 1.5 GeV all methods should agree (it's pQCD after all)

The M_X spectrum



- *Not all observables are equally clean. eg high q^2 tail is sensitive to WA*
- Need spectra: only way to test frameworks (see E_l spectrum).
- *More inclusive measurements, less dependence on m_b*
- Theory errors are partly parametric: m_b dependence is crucial

	Average $ V_{ub} \times 10^3$
DGE	$4.26(14)_{\text{ex}}^{+19}_{-13}$
BLNP	$4.31(16)_{\text{ex}}^{+32}_{-27}$
GGOU	$3.96(15)_{\text{ex}}^{+20}_{-23}$

2.1, 1.9, 1.3 σ from $B \rightarrow \pi l \nu$
(MILC-FNAL)

3.1, 2.4, 1.5 σ from UFit
(because of $\sin 2\beta$)

NEW preliminary Belle Multivariate analysis only $E_l > 1 \text{ GeV}$

$$|V_{ub}| = (4.45 \pm 0.26^{+0.13}_{-0.22}) \times 10^{-3}_{\text{GGOU}}$$

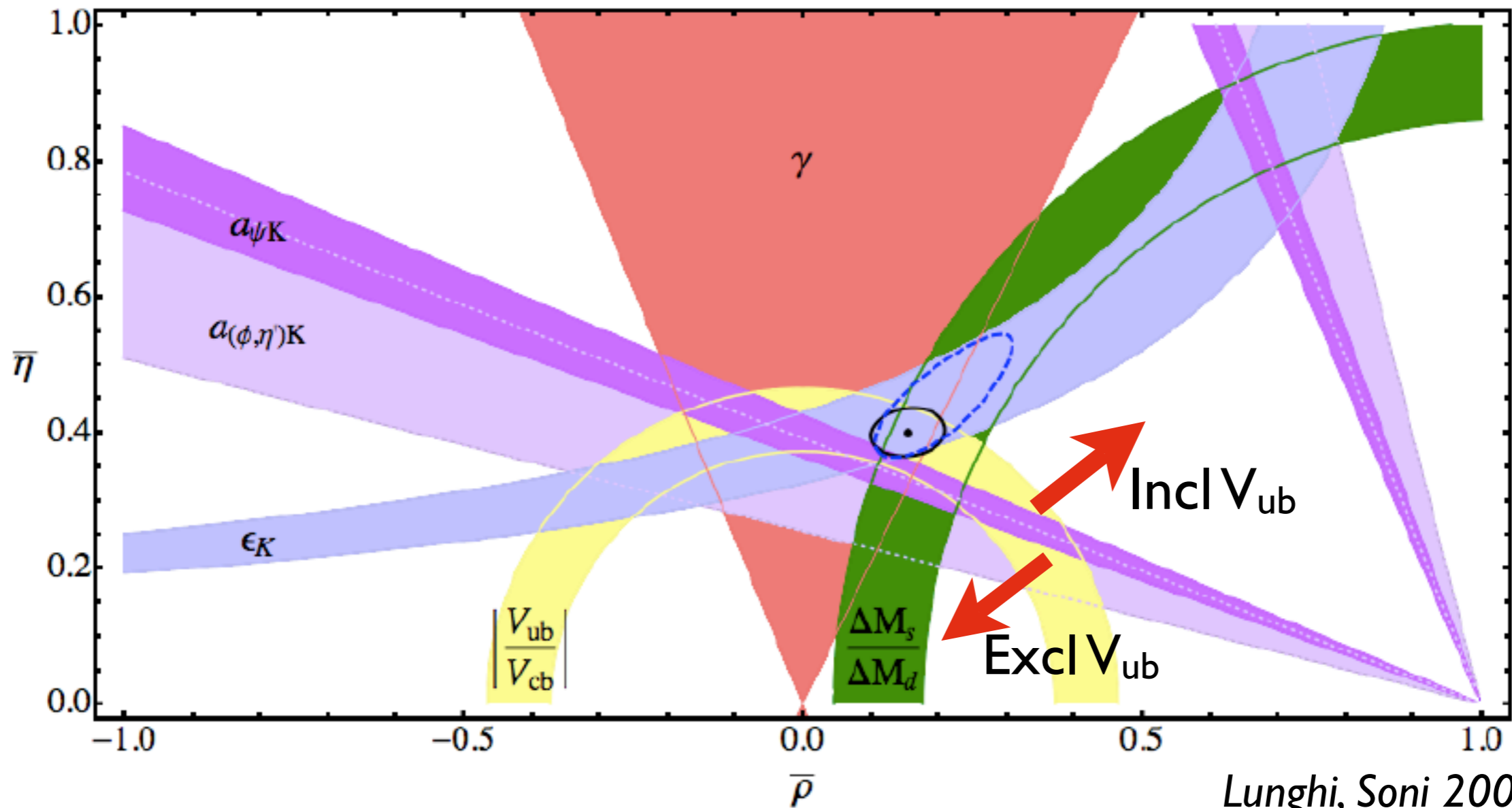
2.1 σ from excl, 2.5 σ from UFit

This includes about 90% of the rate really inclusive measurement, no need for SF. Only crucial input m_b needs to be confirmed!

NEW PHYSICS?

eg LR models Chen, Nam





Recent lattice results for B_K and previously neglected contributions lead to 15% smaller ϵ_K , in $\sim 1.8-2\sigma$ conflict with $\exp \sin 2\beta$.

Perhaps $\sin 2\beta$ is simply too low...

...or incl V_{ub} and latest B_K both wrong

Conclusions

- ➔ MAIN ISSUE for V_{ub} : bottom mass. Can we include additional constraints on m_b in the fit? which ones?
- ➔ *Not all observables are equivalent*, some are cleaner. For ex high q^2 tail is sensitive to WVA: it decreases V_{ub} . Can we drop it? how much do the exp analyses depend on the high q^2 tail?
- ➔ Need spectra and/or analysis with varying cuts: only way to test current frameworks (see trial exercise on E_l spectrum). M_X cuts?
- ➔ More inclusive measurements are welcome: they decrease the dependence of $|V_{ub}|$ on both SF and m_b
- ➔ Frameworks fairly compatible within non-par th errors. Convergence to OPE (normalization) to be checked. Use m_b uncertainty as exp error?
- ➔ The primary goal is the precise determination of $|V_{ub}|$. All frameworks are interesting, but they are **not all equivalent**. After $|V_{ub}|$ is measured we can go back and study models of QCD dynamics.